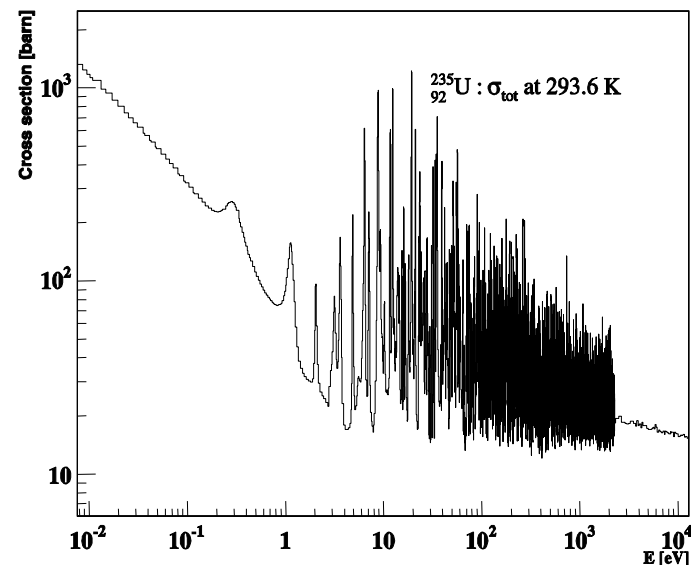
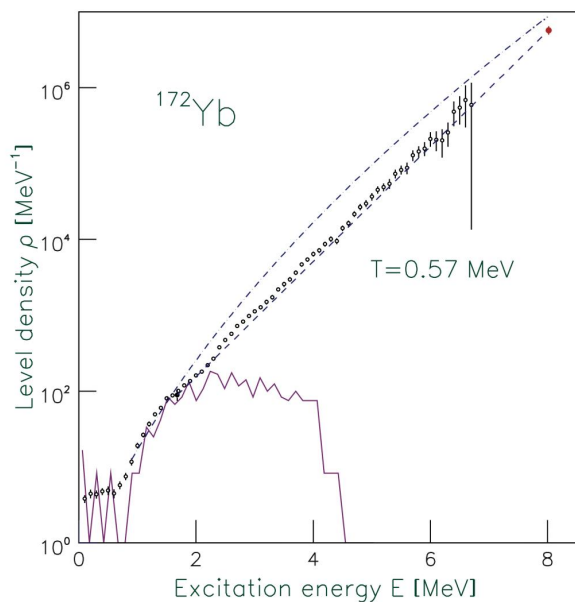


Level Densities, Decay Probabilities and Cross sections in the Actinide Region



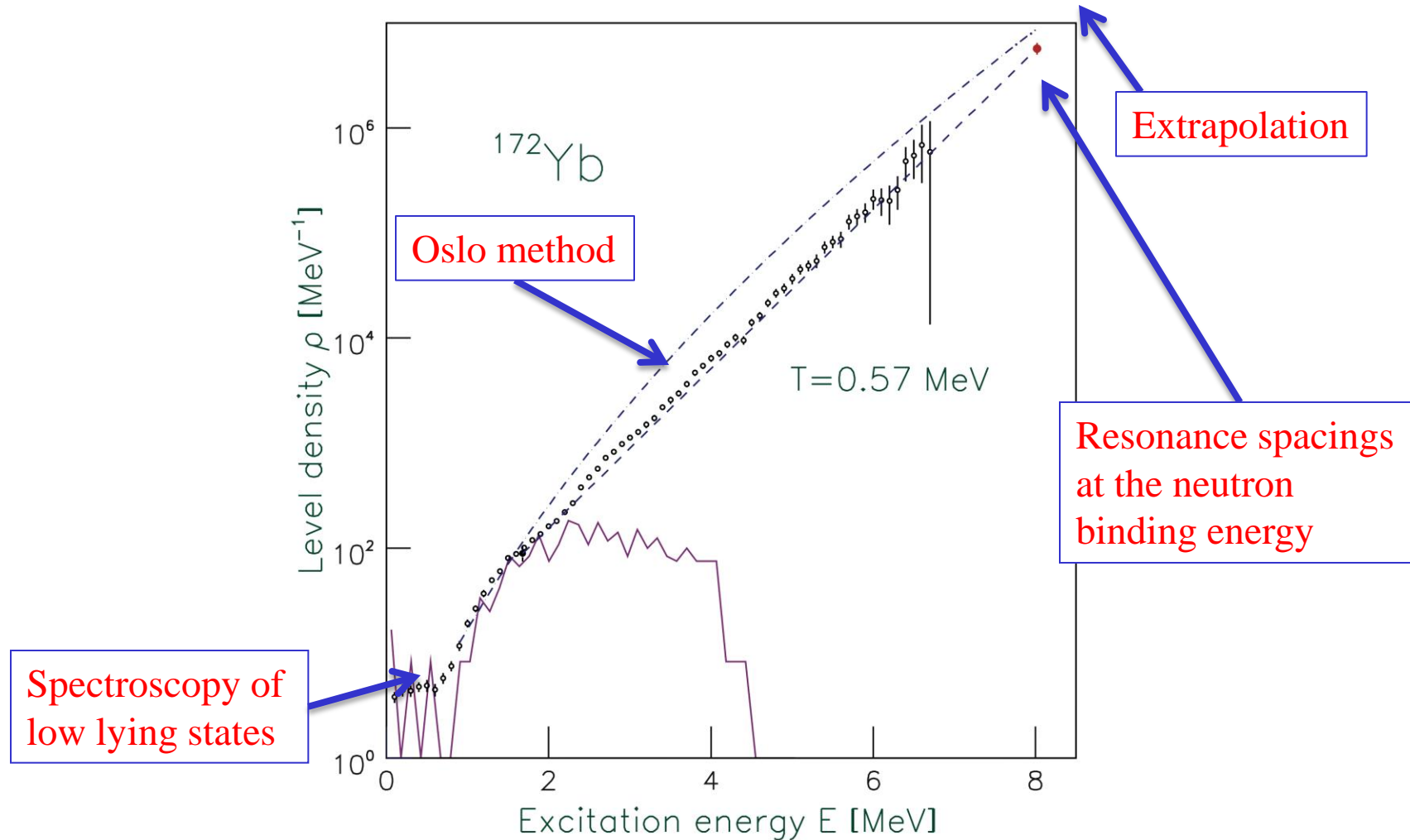
J.N. Wilson

Institut de Physique Nucléaire, Orsay

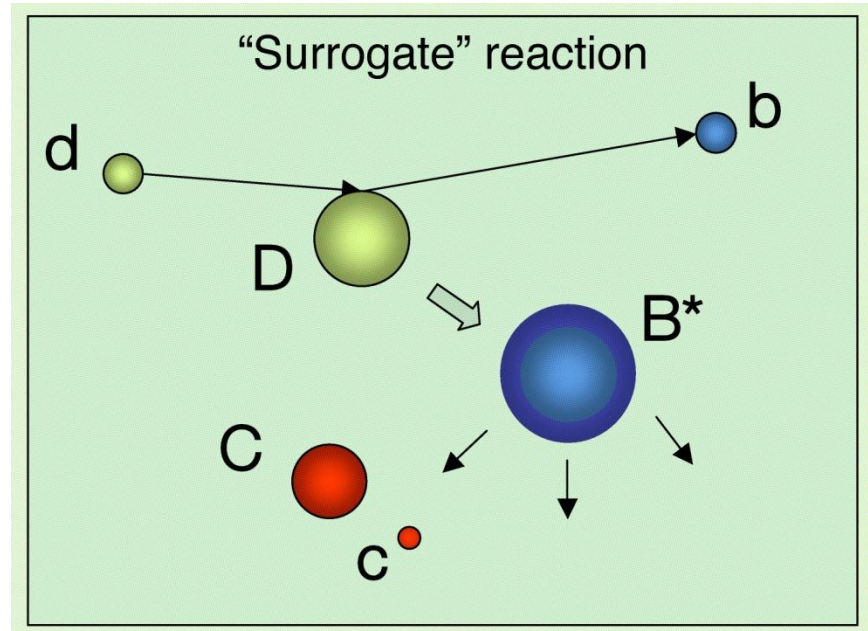
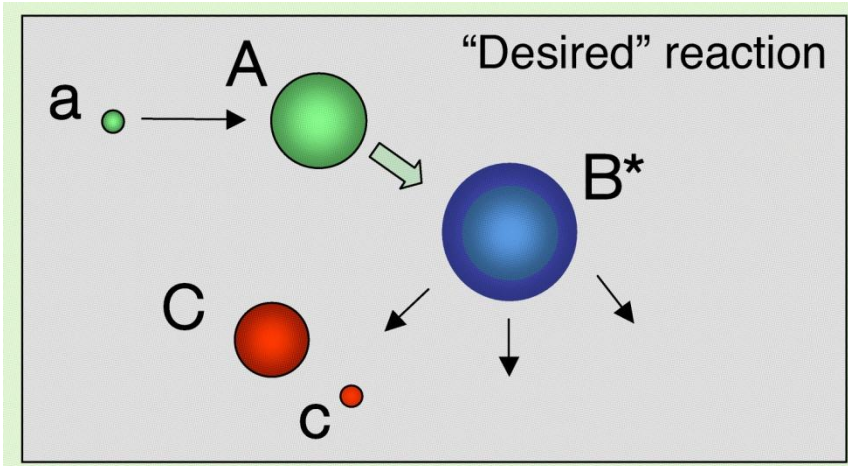
NLD Measurements in the Actinides: Goals & Motivations

1. Fundamental physics - NLD and GSF fine structure, statistical properties of the hot compound nucleus
2. Indirect cross section measurements via surrogate reactions
3. Level density knowledge helps cross section calculations where direct measurements are difficult or impossible

1. Fundamental Physics



2. Indirect Measurements

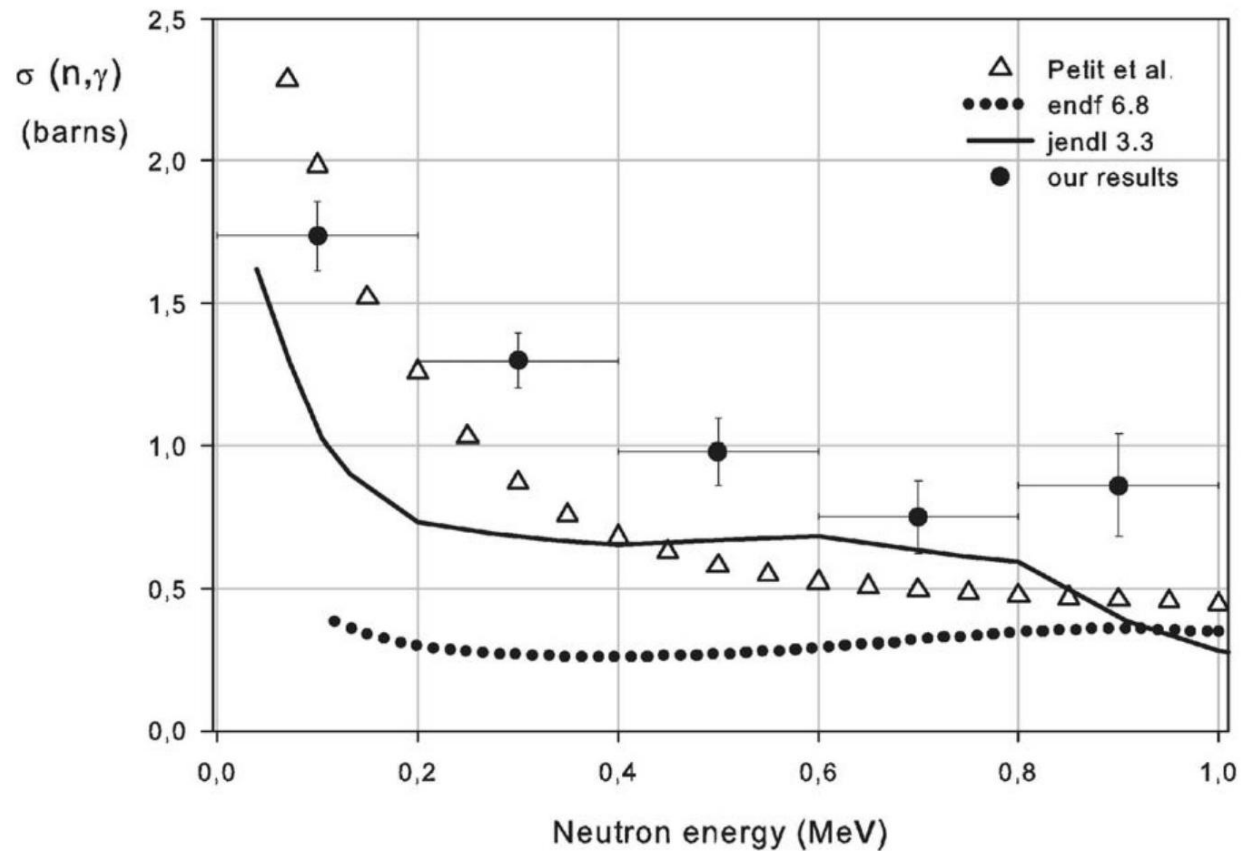


$$\sigma_{n,f}(E_n) = \underbrace{\sigma_{CN}(E_n)}_{\text{Optical model calculation}} \cdot \underbrace{P_f(E_n)}_{\text{Measured}}$$

Optical model calculation

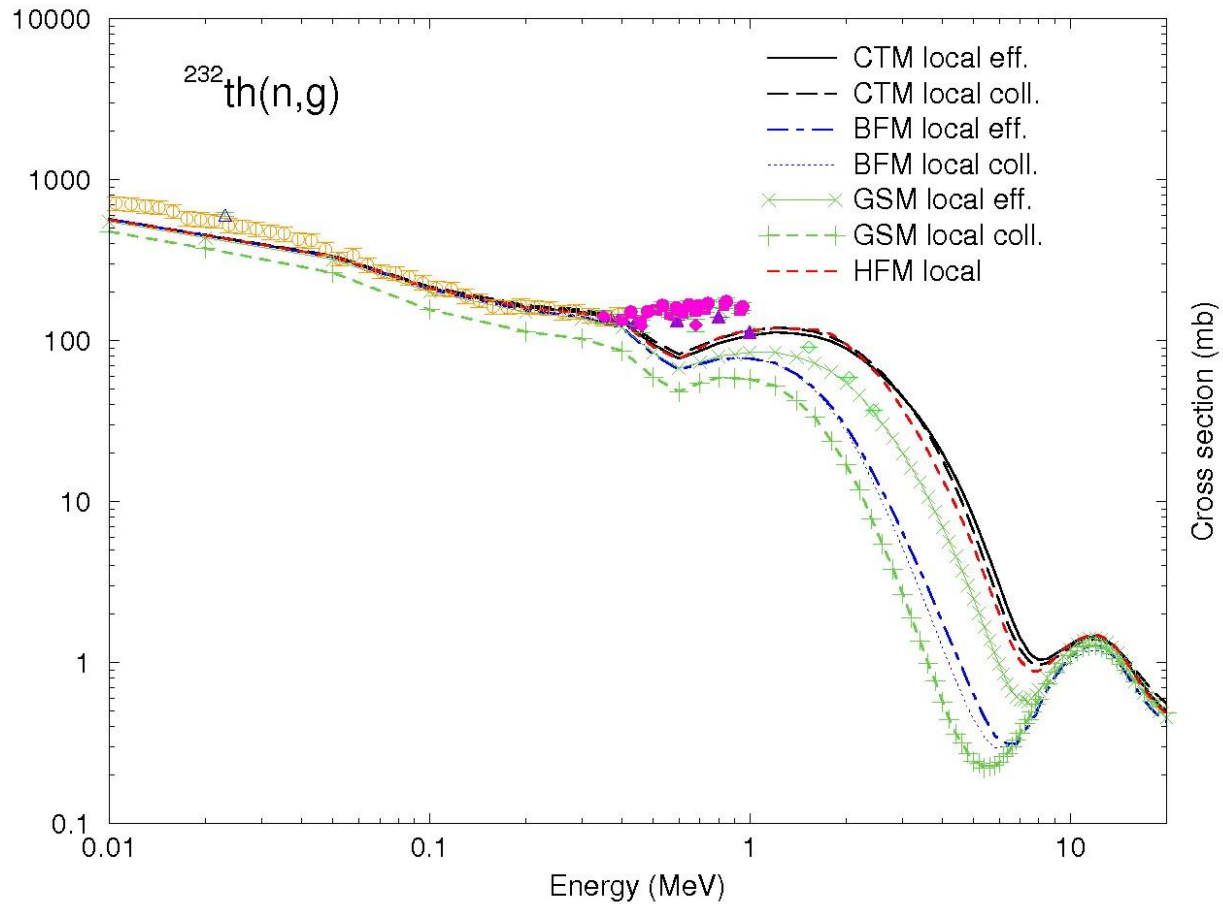
Measured

2. ^{233}Pa Capture Experiment Results



S. Boyer et al. Nucl.Phys. A775, 175 (2006)

3. x-section calculations

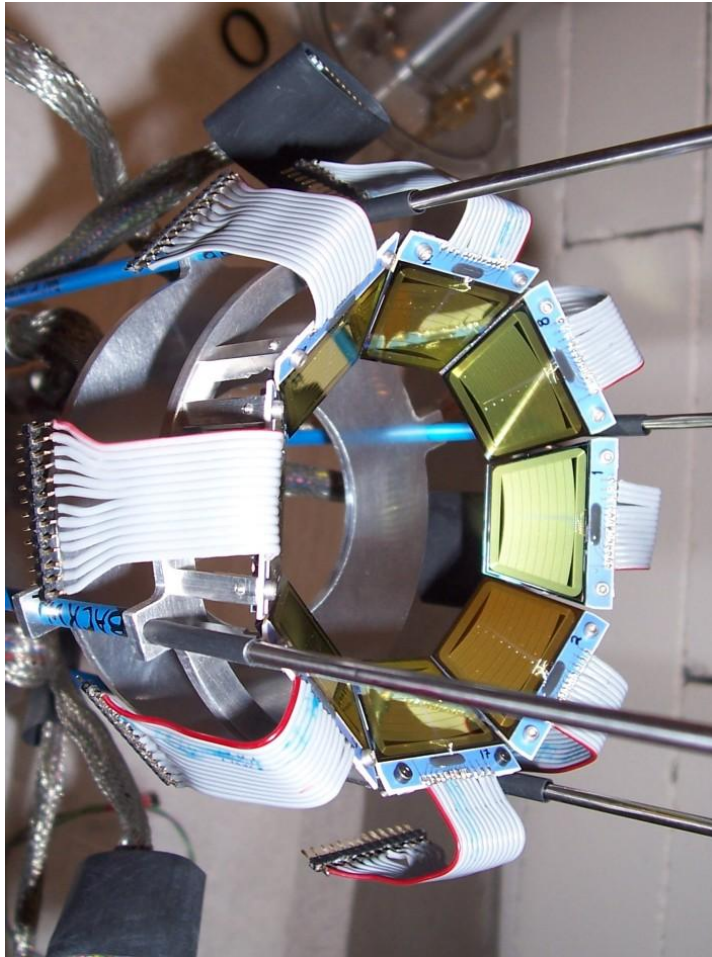


Part I

First Experimental Data

Oslo Cyclotron Experiments

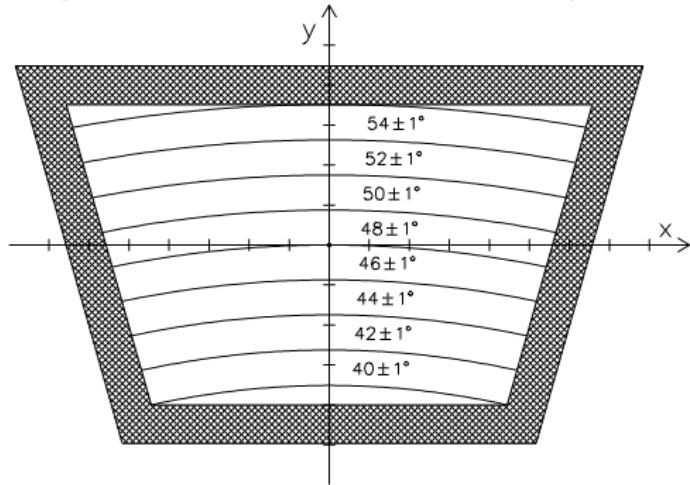
Si Ring



Cactus NaI array



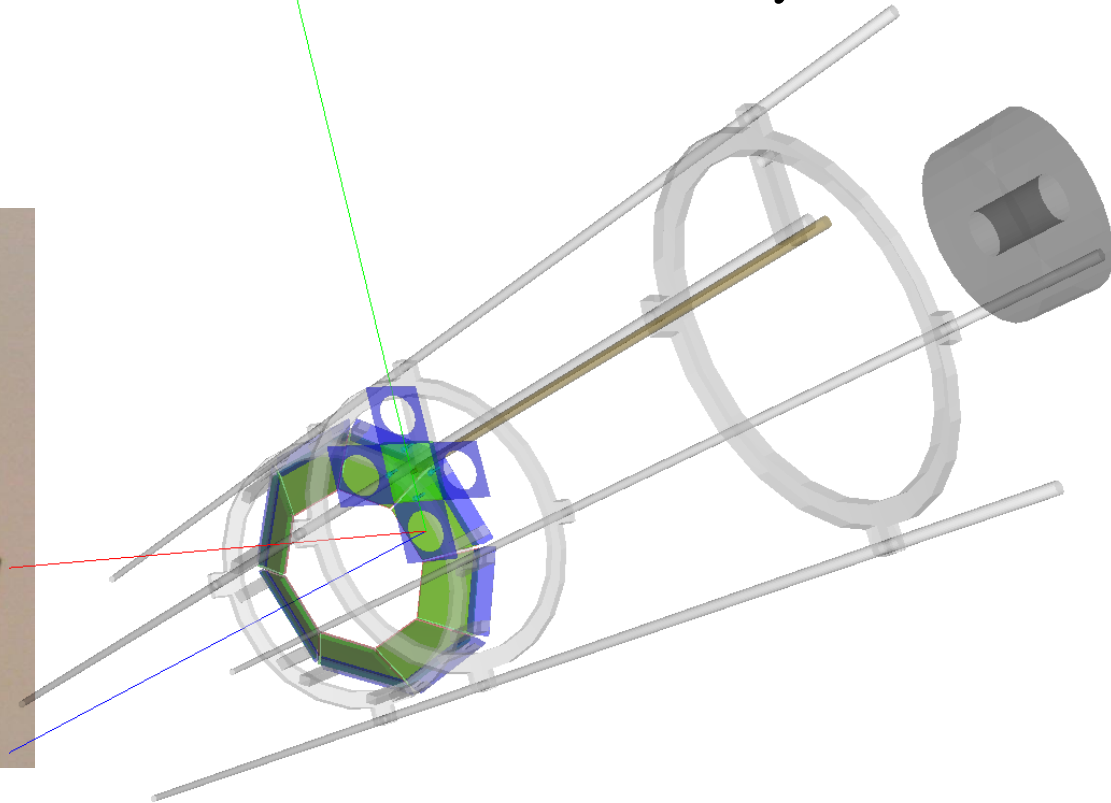
Angle between side rim and vertical = 15.77 degrees



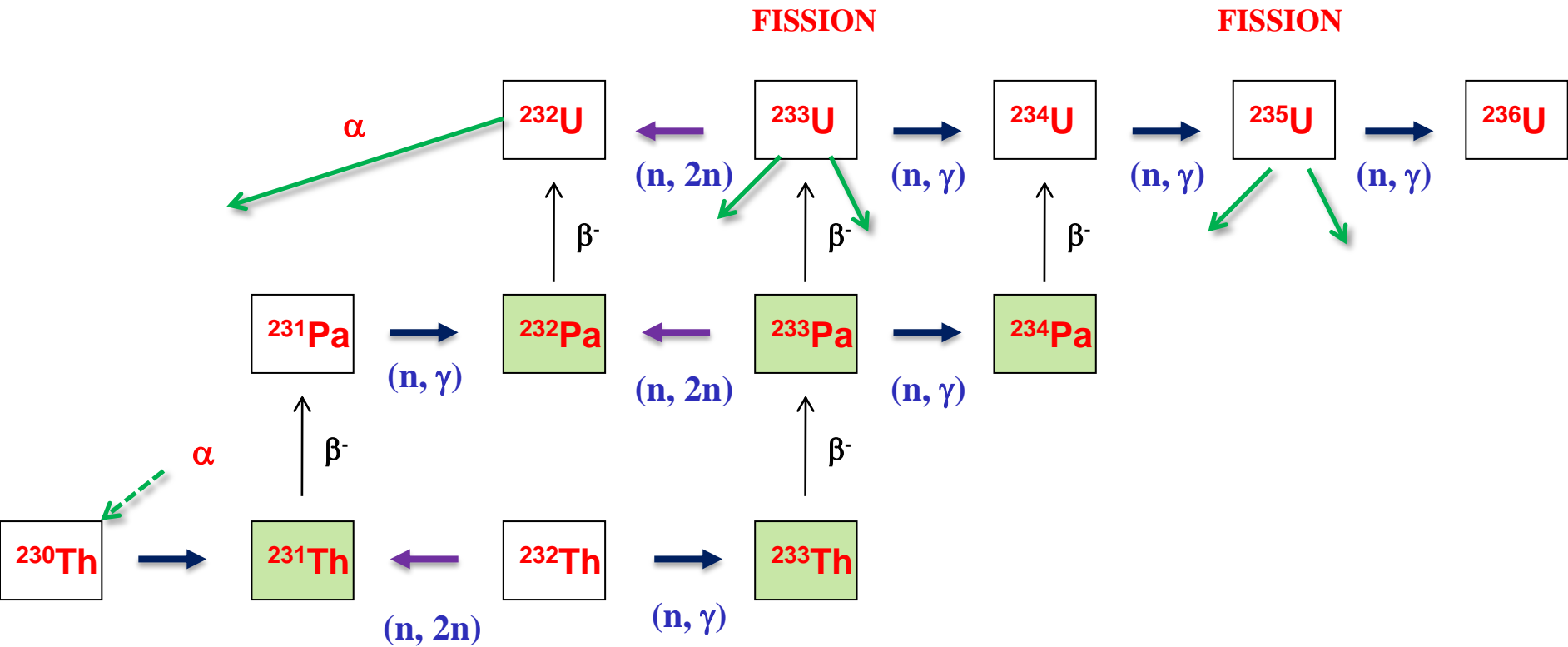
One tick on the axes corresponds to 2 mm in reality

The Silicon Ring

- High efficiency/high angular resolution detection of the outgoing particles ($< 2^\circ$)
- $E/\Delta E$ collimators not necessary



2. The Thorium Cycle



Experiment ^{232}Th : 01/12/09 – 13/12/09

$^{232}\text{Th}(d,x)$ @ 12 MeV

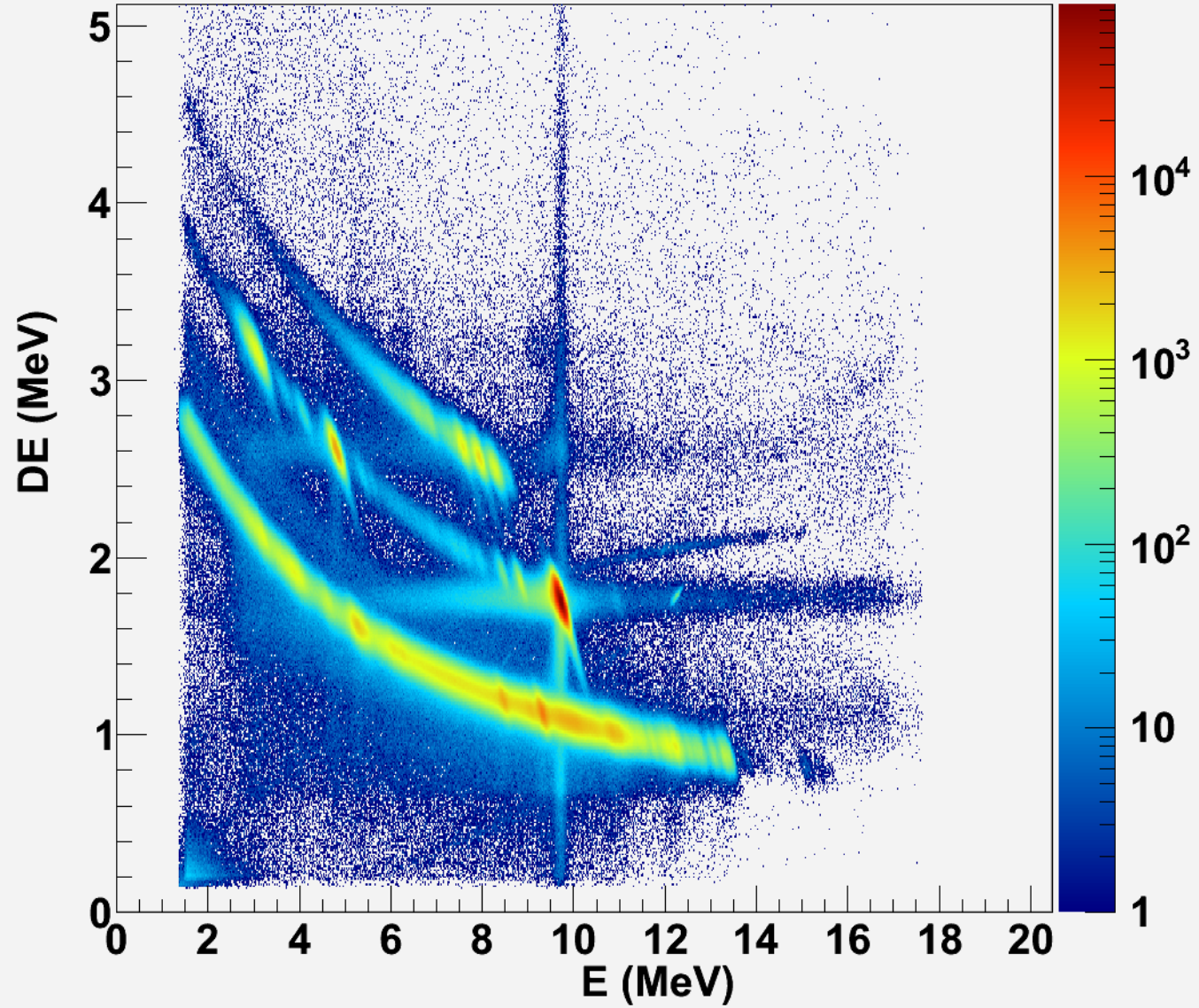
351 M events

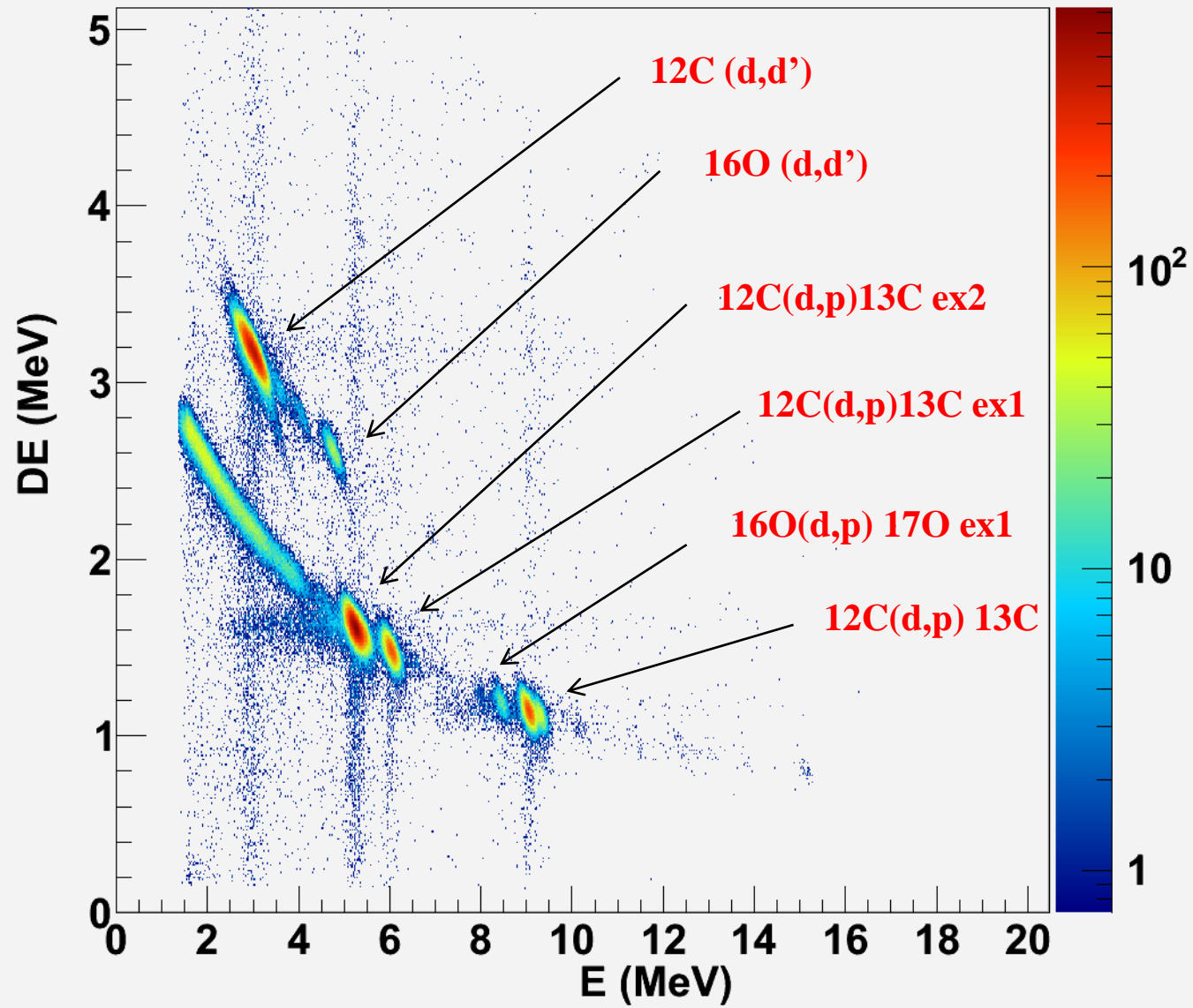
Reaction	Compound Nucleus	Ex-Eg counts
d,p	^{233}Th	23 M
d,d'	^{232}Th	0.24 M
d,t	^{231}Th	1.2 M

$^{232}\text{Th}(^3\text{He},x)$ @ 24 MeV

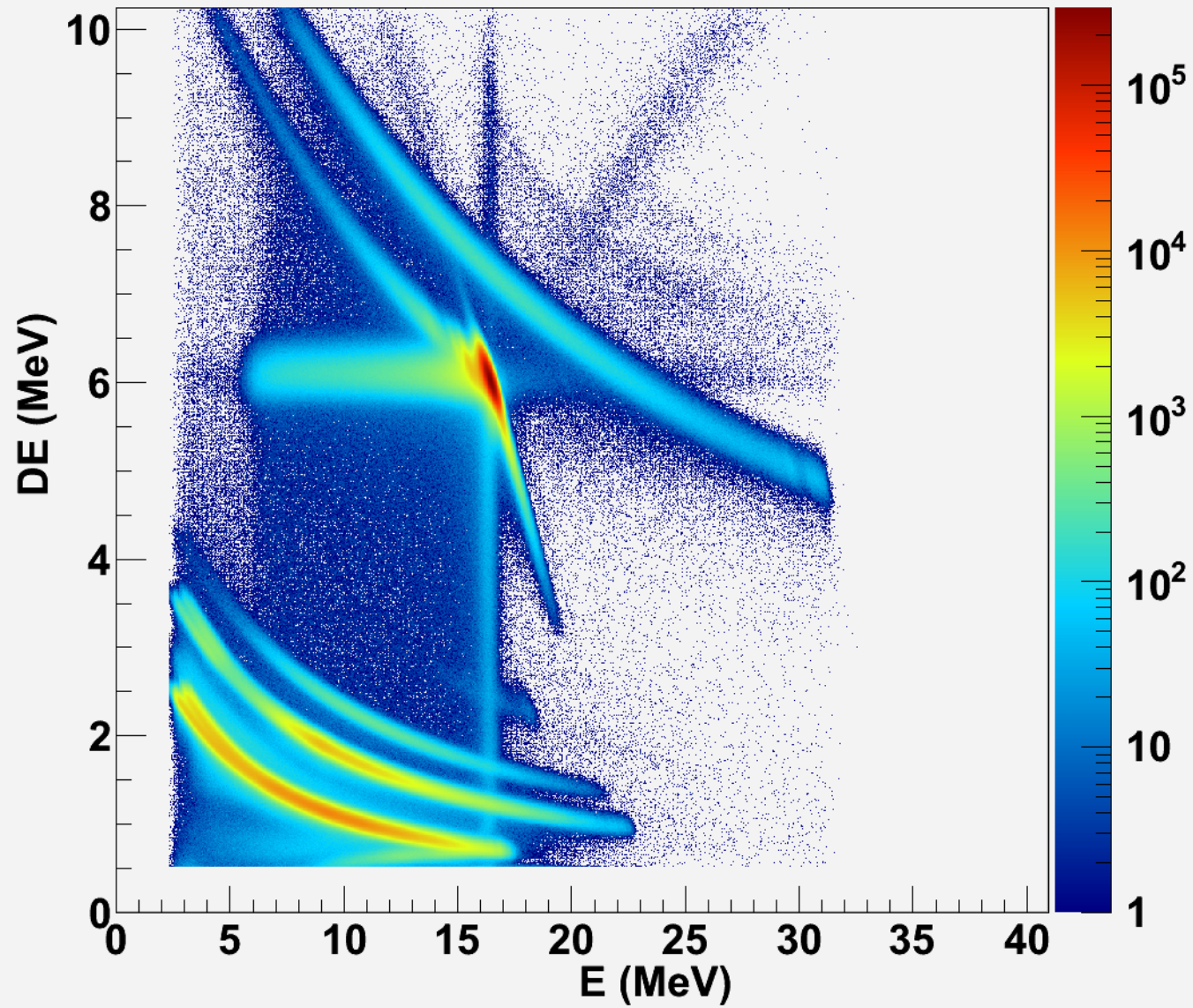
135 M events

Reaction	Compound Nucleus	Ex-Eg counts
$^3\text{He},p$	^{234}Pa	13 M
$^3\text{He},d$	^{233}Pa	6.0 M
$^3\text{He},t$	^{232}Pa	0.57 M
$^3\text{He},a$	^{231}Th	0.79 M





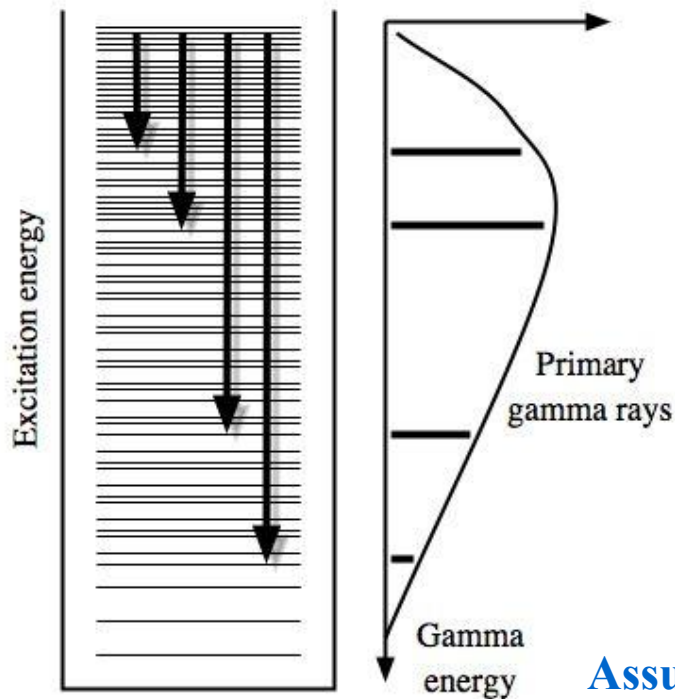
$^{232}\text{Th}(^3\text{He},x) @ 24\text{MeV}$



Part II

Extracting the Level Density

The Brink Axel hypothesis



$$P(E_i, E_\gamma) \propto \rho(E_f) \cdot T(E_\gamma)$$

Level density

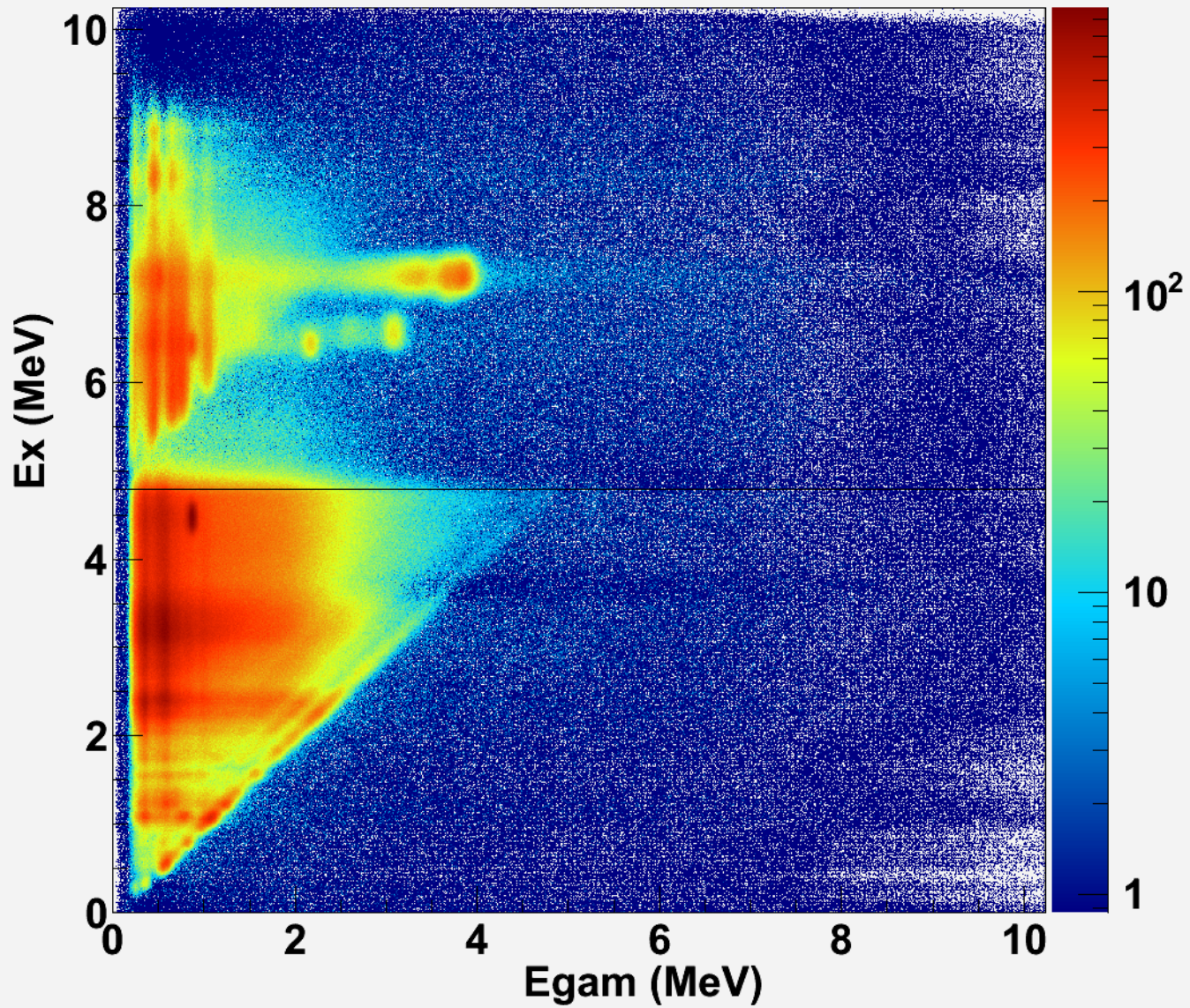
Transmission Coefficient

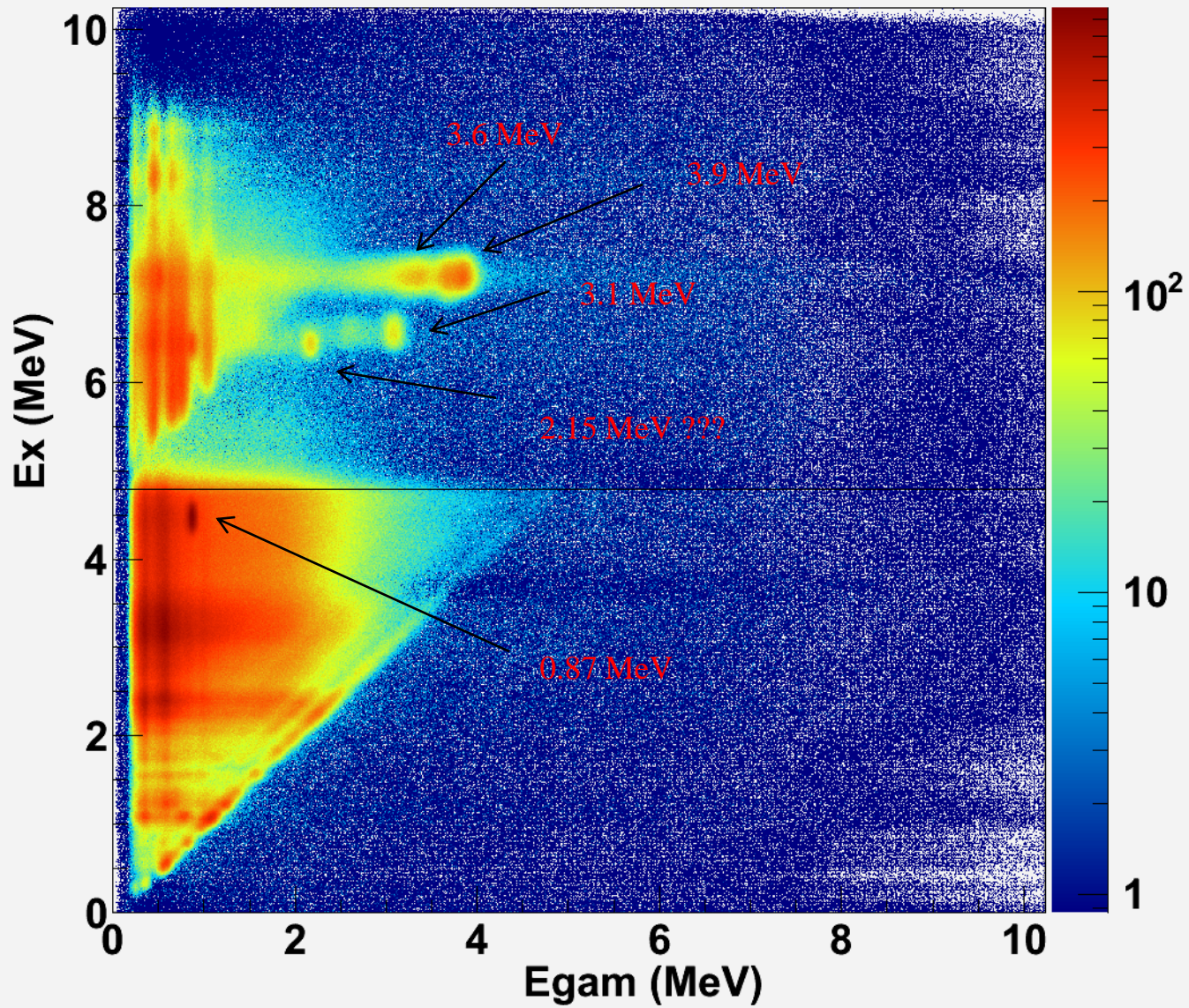
Gamma strength function

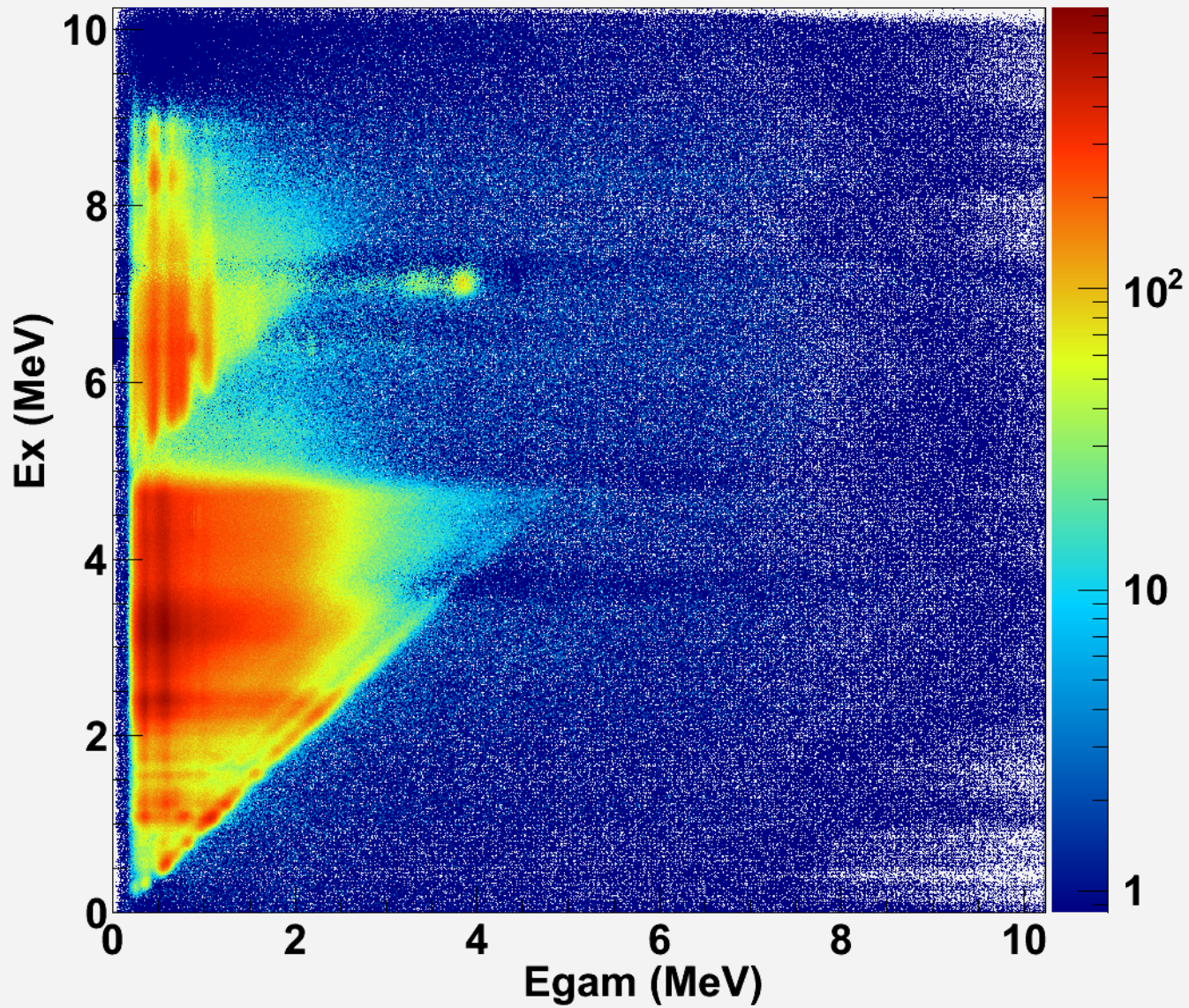
$$T(E_\gamma) = 2\pi \sum_{XL} E_\gamma^{2L+1} f_{XL}(E_\gamma)$$

Assuming dominance of dipole radiation (*E1 and M1*)

$$f(E_\gamma) \simeq \frac{1}{2\pi} \frac{T(E_\gamma)}{E_\gamma^3}$$

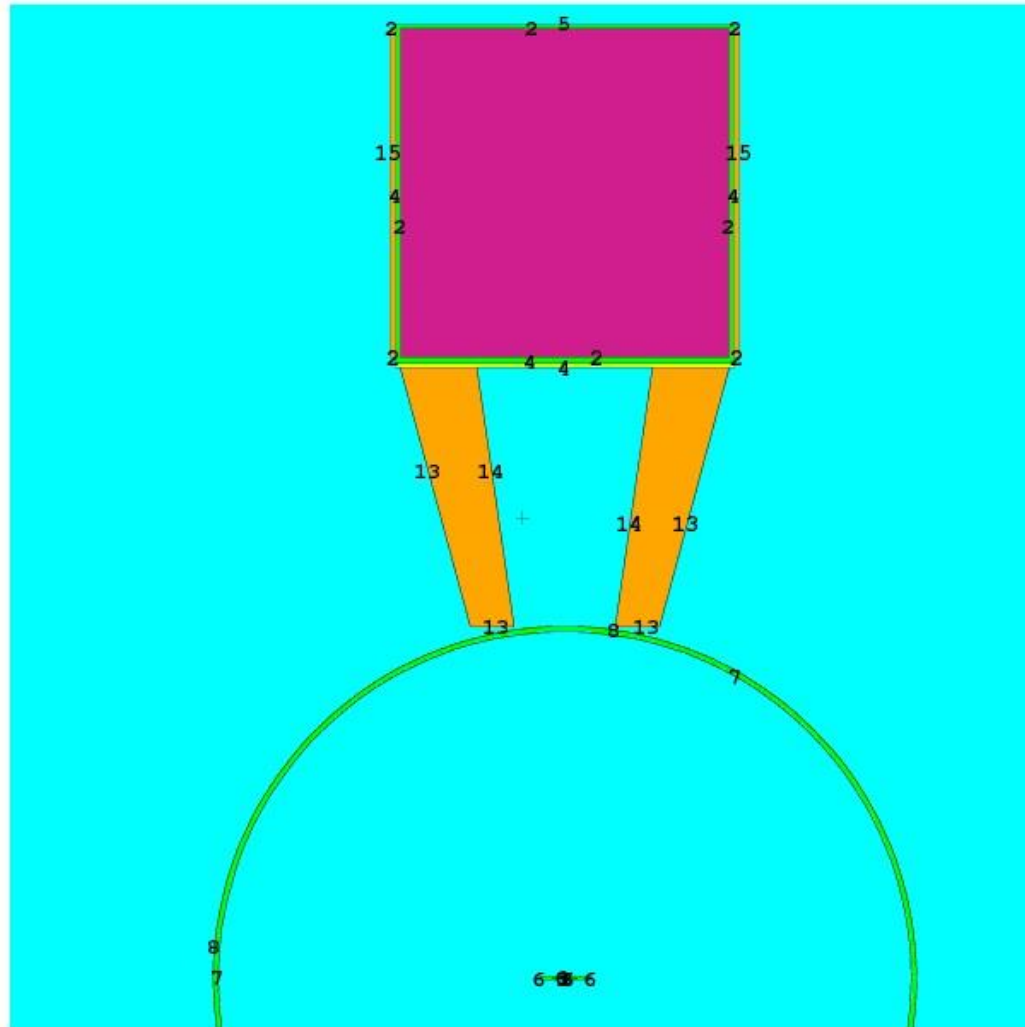


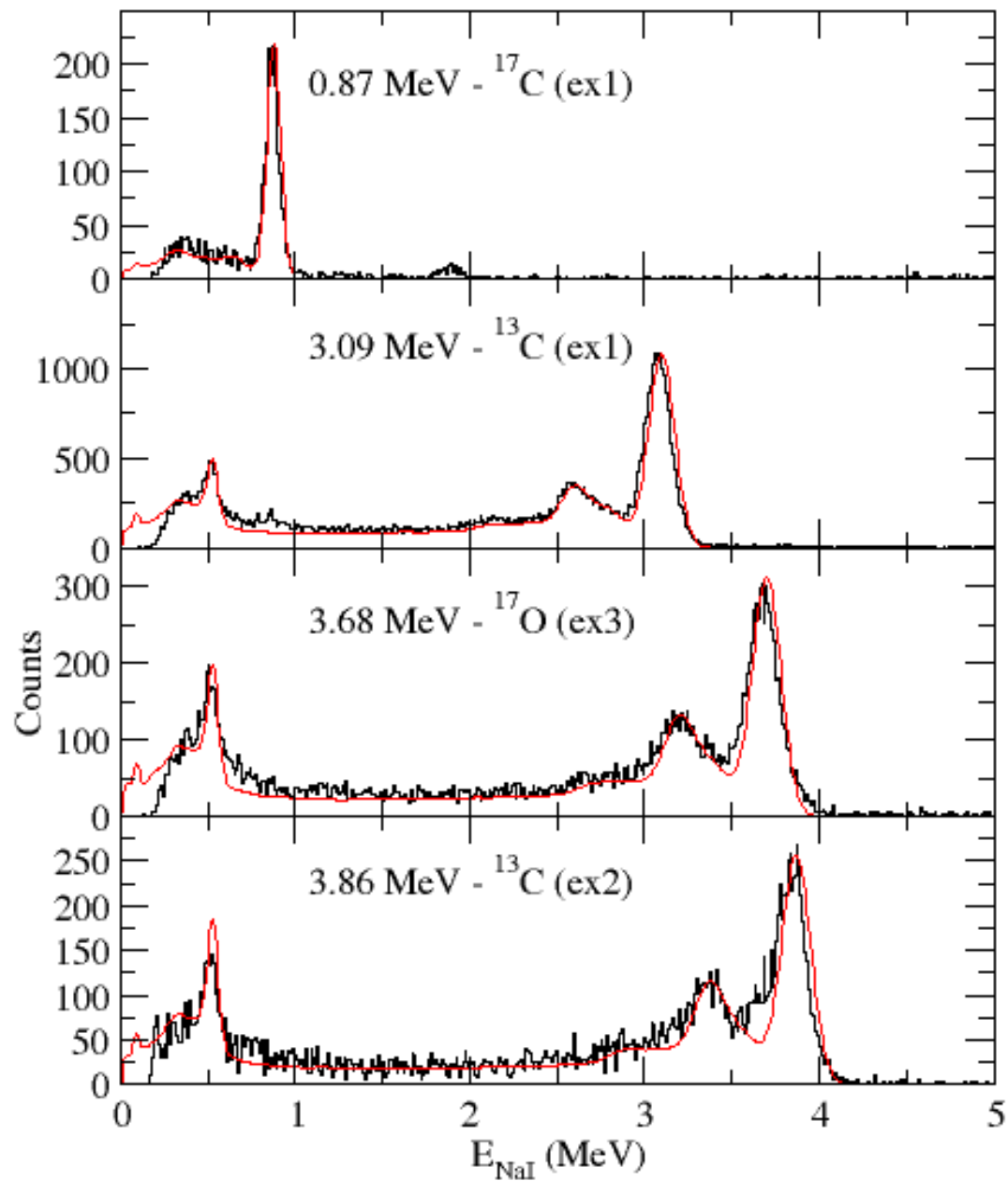


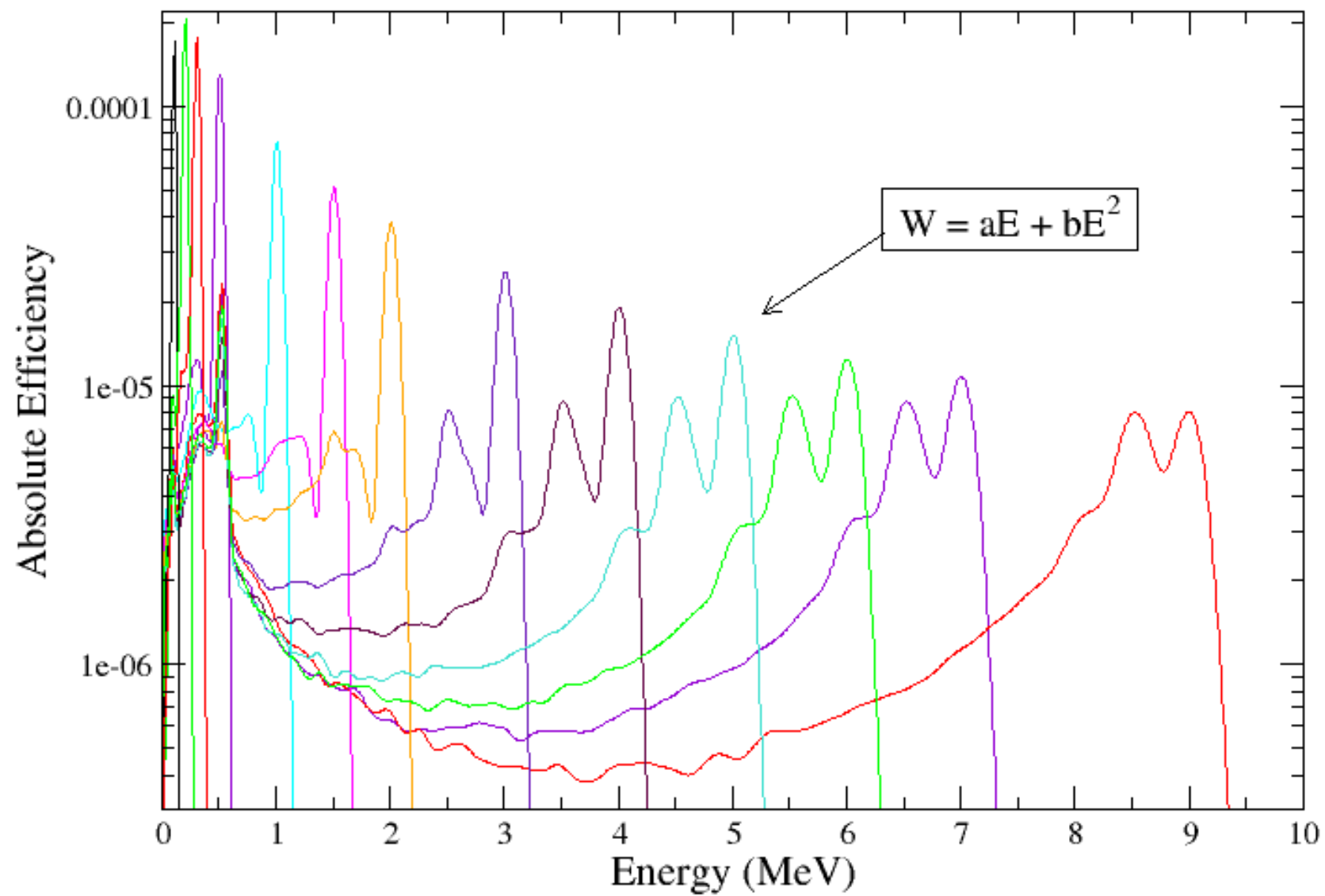


MCNP simulations of Cactus detector response

```
06/18/10 17:22:13
gamma-ray efficiency: NaI -
DETECTOR + THORIUM + COPPER +
Al CAN
probid = 06/18/10 17:20:58
basis: YZ
( 0.000000, 1.000000, 0.000000)
( 0.000000, 0.000000, 1.000000)
origin:
( 0.00, -1.62, 17.80)
extent = ( 19.83, 19.83)
```

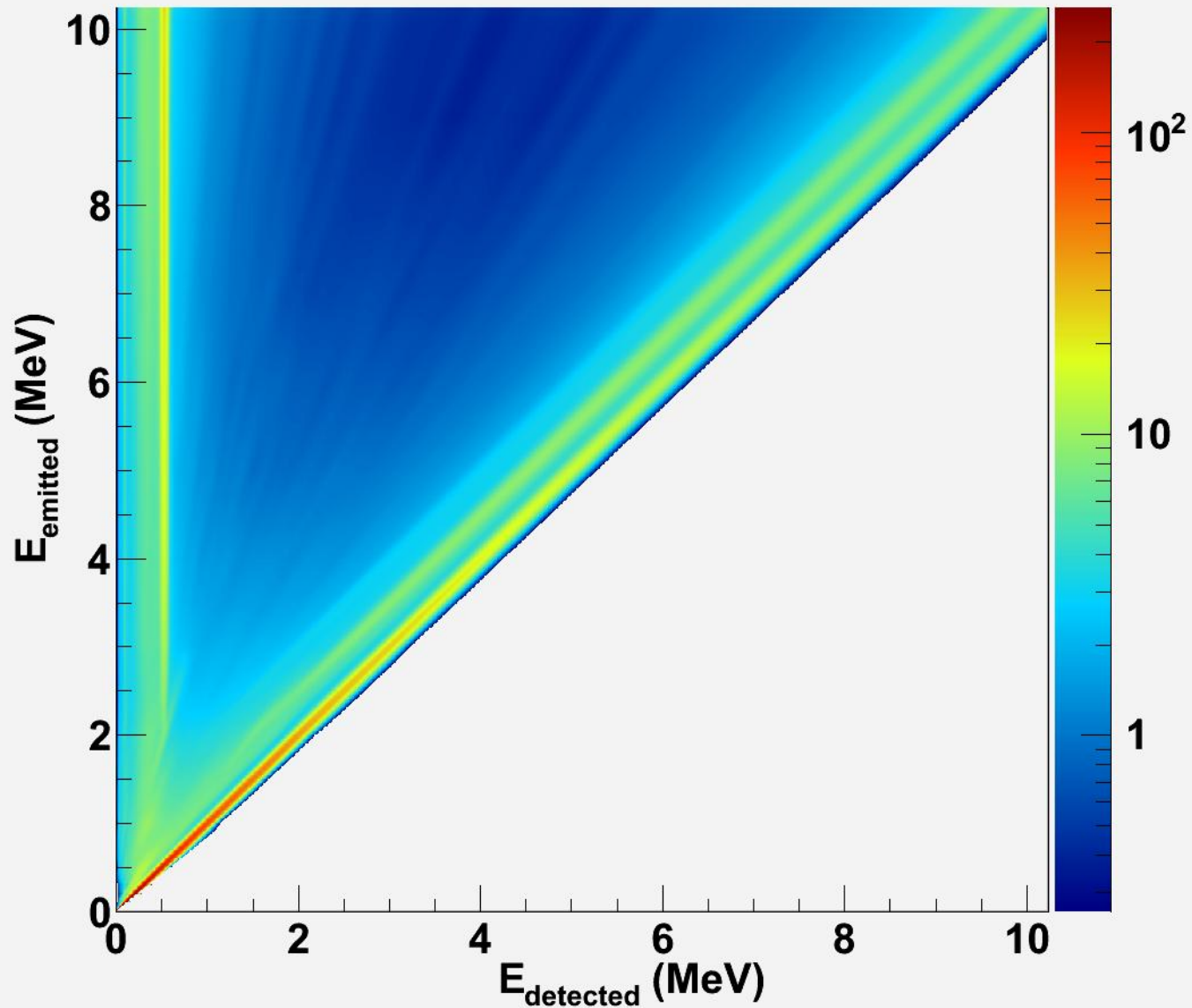


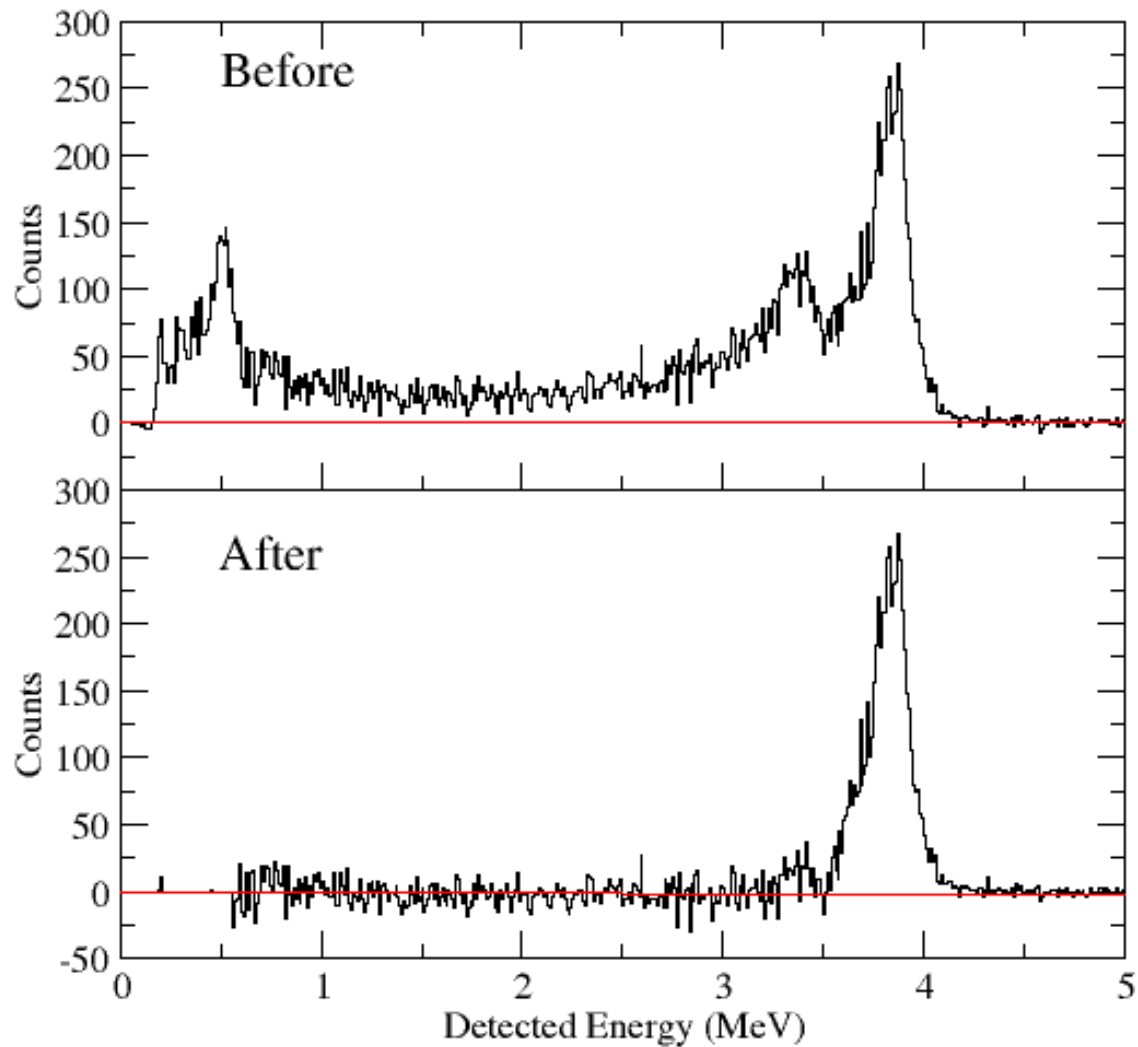




Cactus Response Matrix

E_{det}-E_{emit}

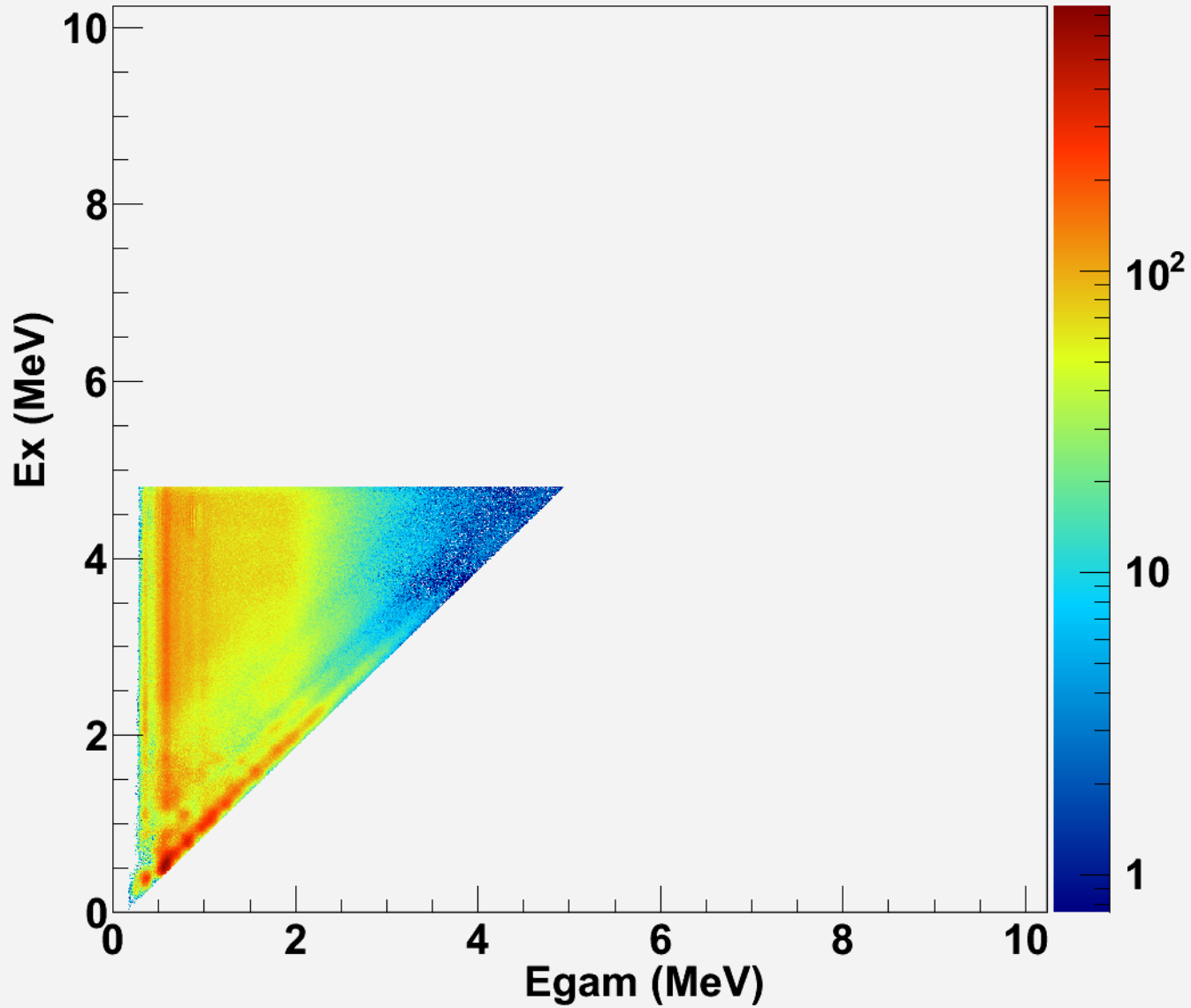




Matrix *response;
Spec *after=before->SubtractComptons(response);

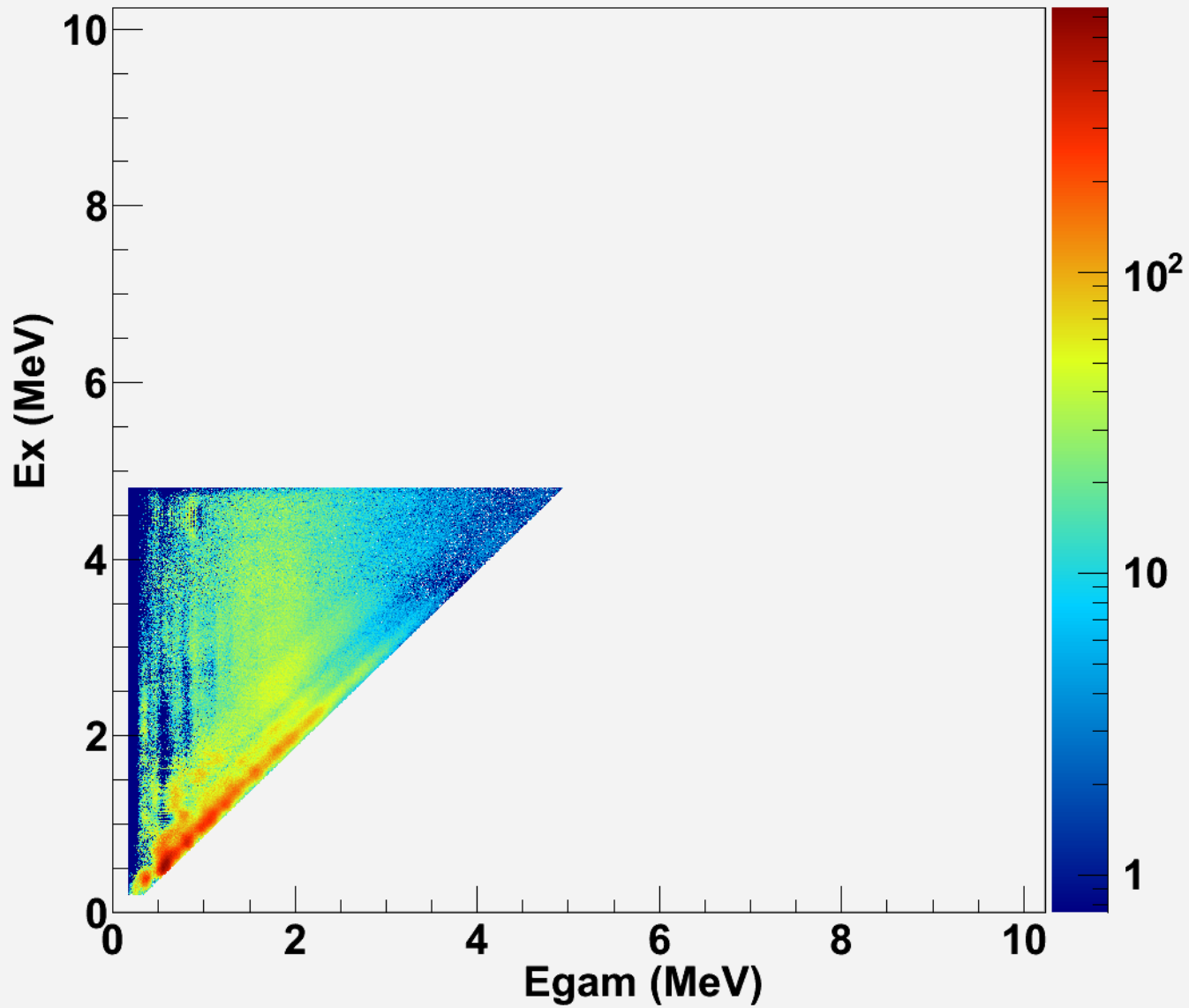
$^{233}\text{Th}^*$ - After Compton Subtraction

Egam-Ex



$^{233}\text{Th}^*$ - Primary gamma rays

Egam-Ex



Part III

The surrogate method

Decay probability measurement

Efficiency of detecting a cascade, ϵ_c of m gammas is:

$$\epsilon_c = 1 - \prod_{j=1..m} (1 - \epsilon_j)$$

If ϵ is small, then:

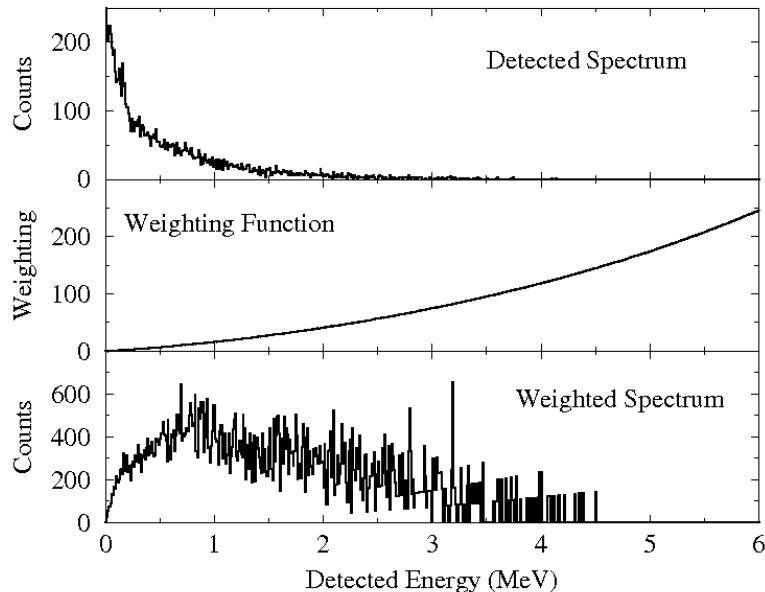
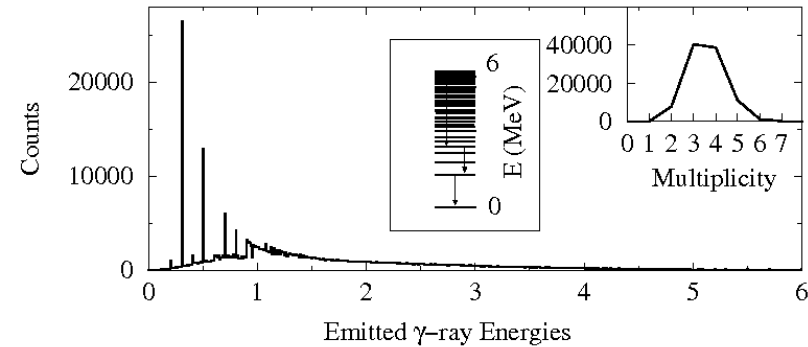
$$\epsilon_c \approx \sum_{j=1..m} \epsilon_j$$

Suppose detector has this property:

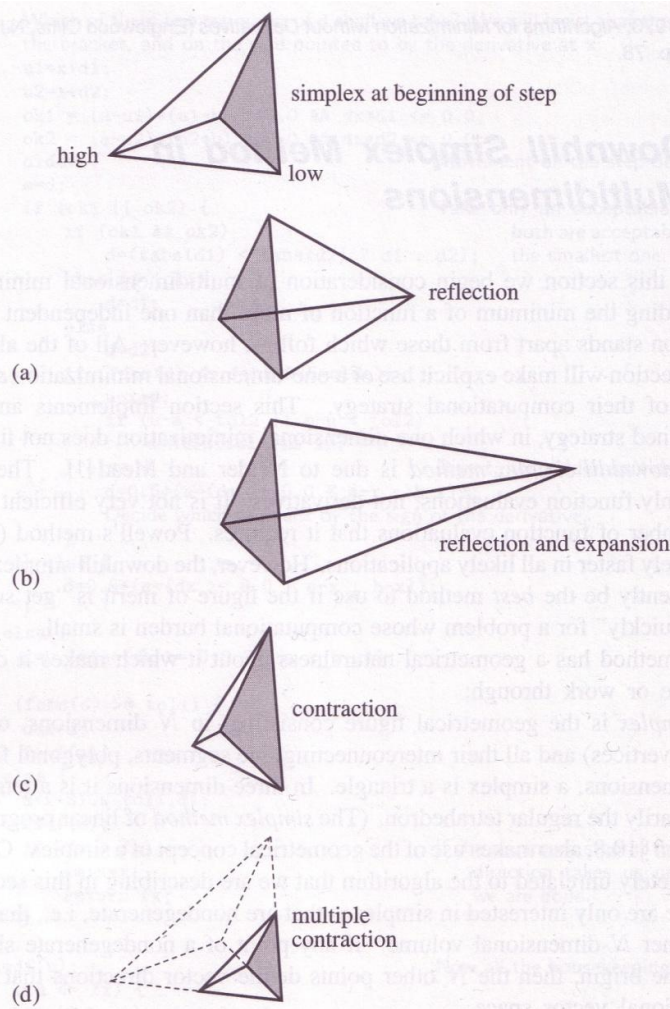
$$\epsilon = k E_\gamma$$

Then efficiency of detecting a cascade becomes proportional to cascade energy and independent of cascade path!!

$$\sum_d W(E_\gamma) R(E_d, E_\gamma) = k E_\gamma$$



Minimisation with the Downhill simplex method



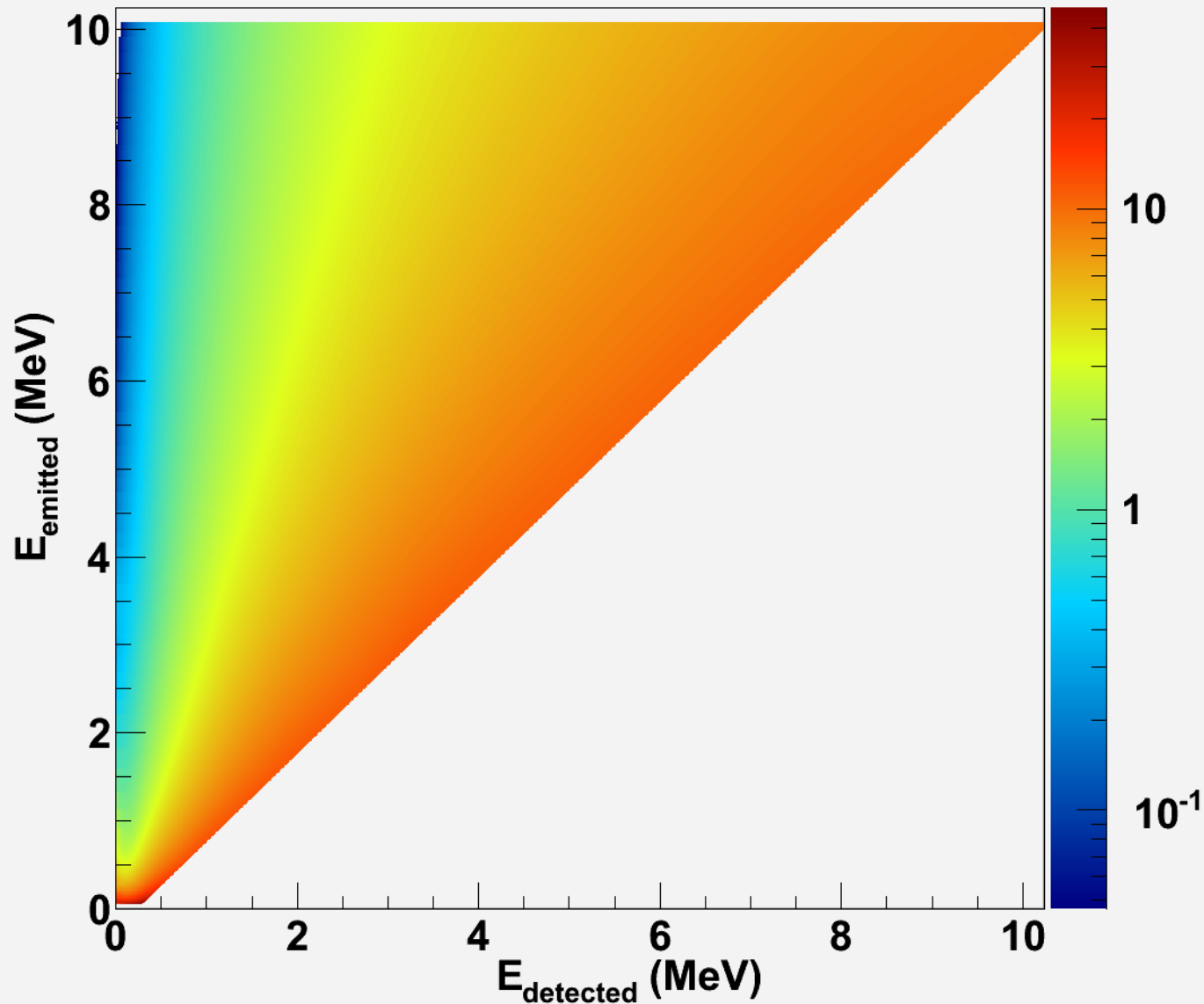
$$\sum_d W(E_\gamma) R(E_d, E_\gamma) - kE_\gamma = 0$$

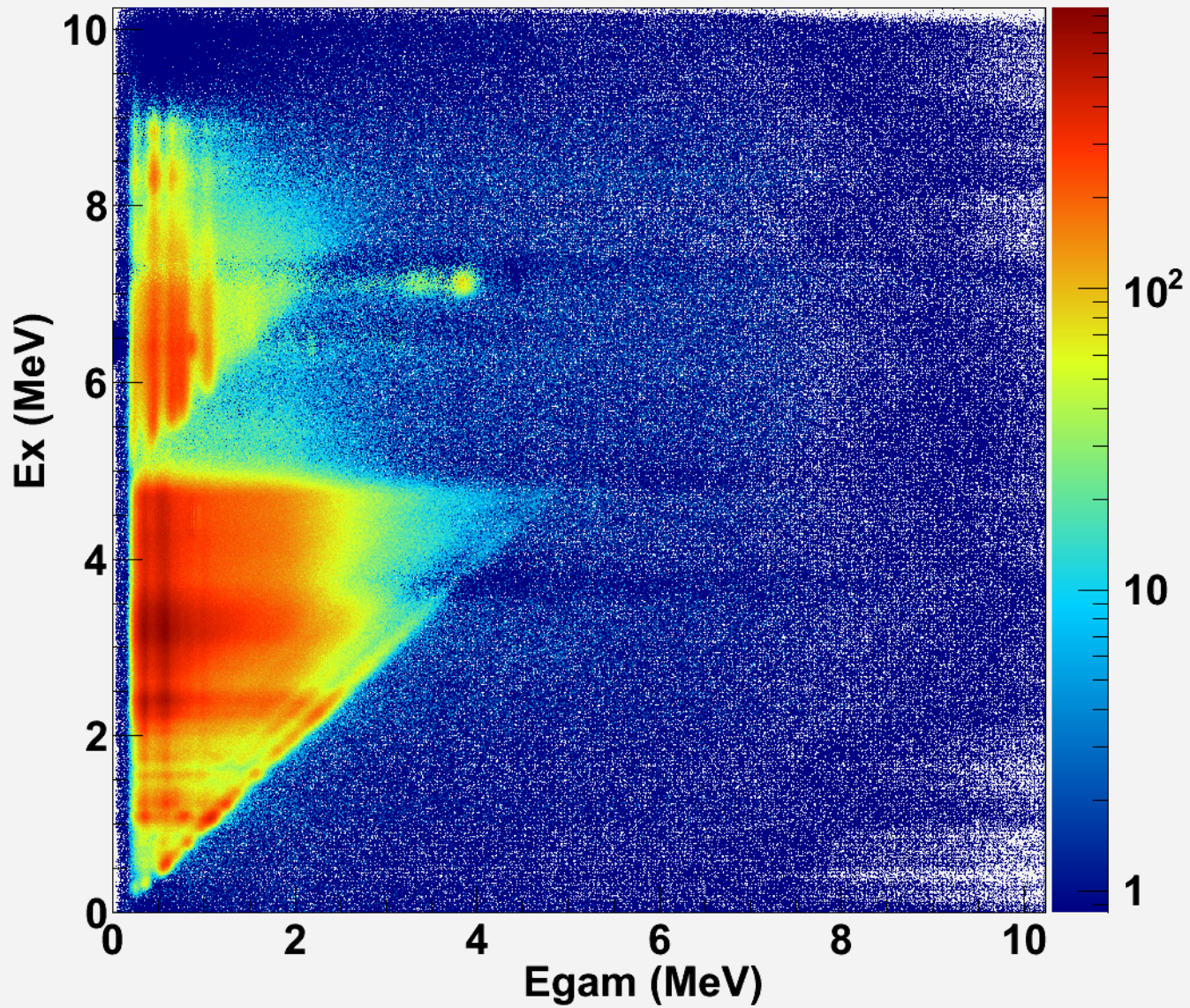
« With four parameters, I could fit an elephant, with five I could make him wiggle his trunk »

$$W(E) = a + bE + cE^2 + dE^3 + eE^4 + fE^5$$

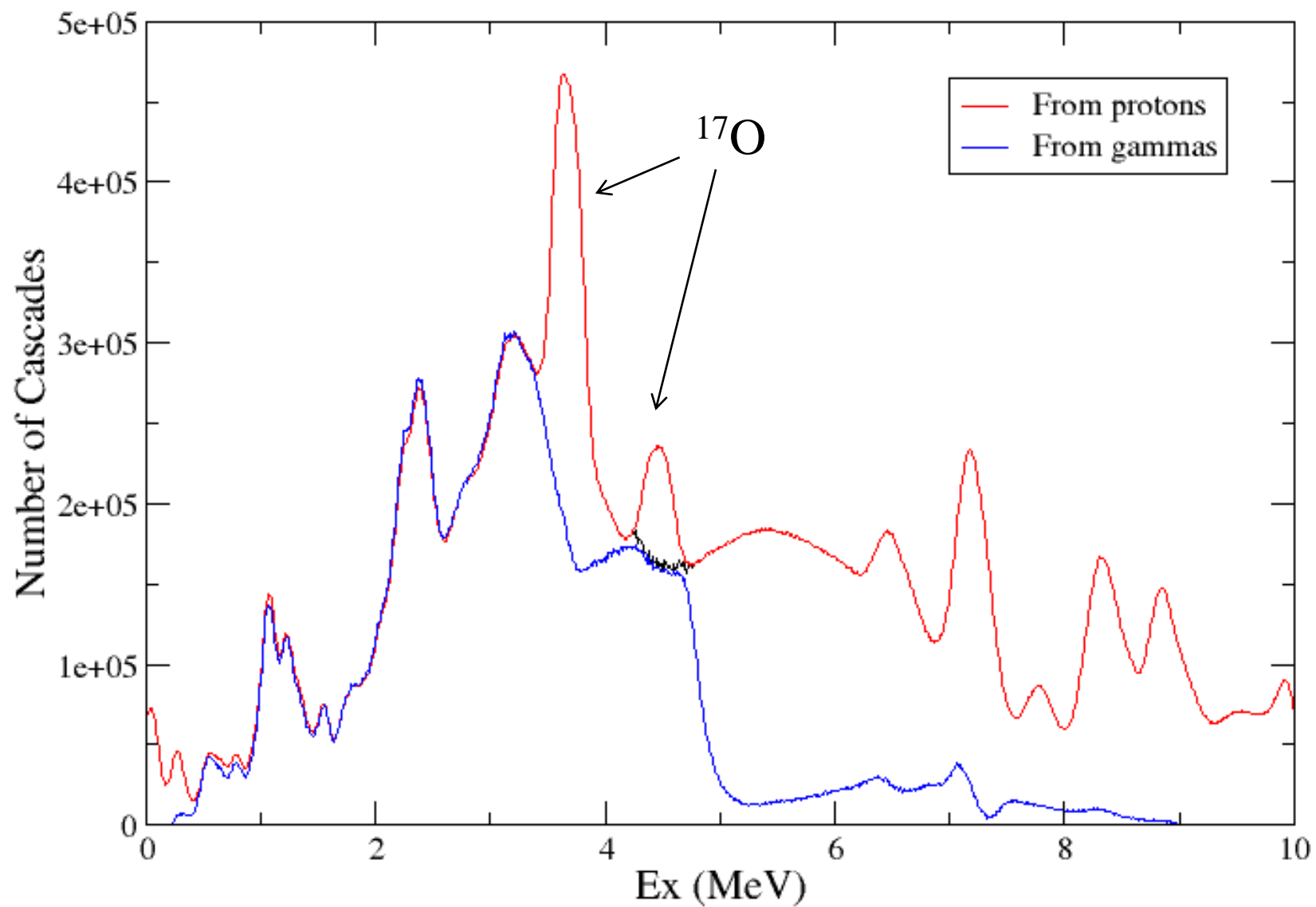
Weighting Function Matrix

Edet-Eemit

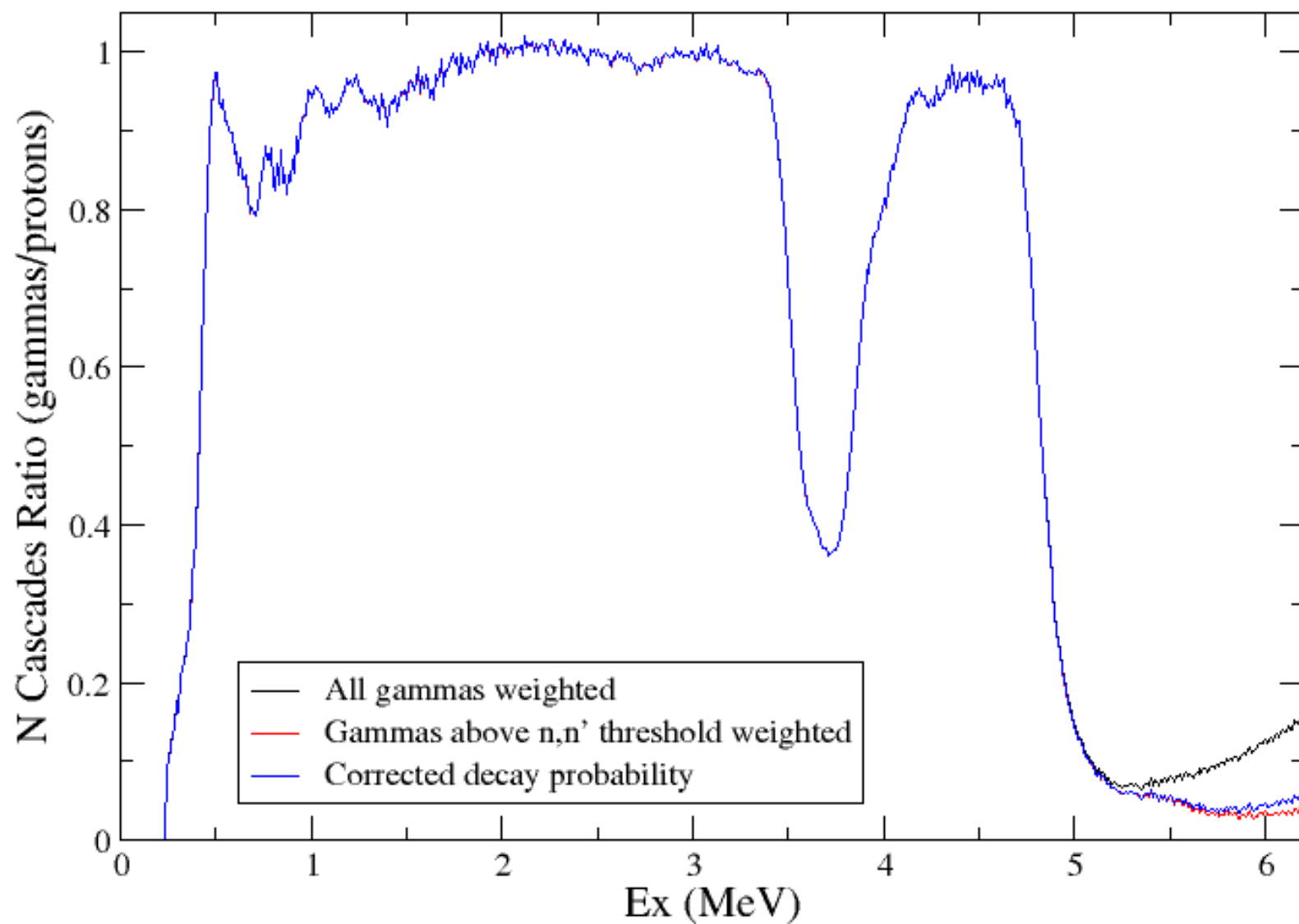


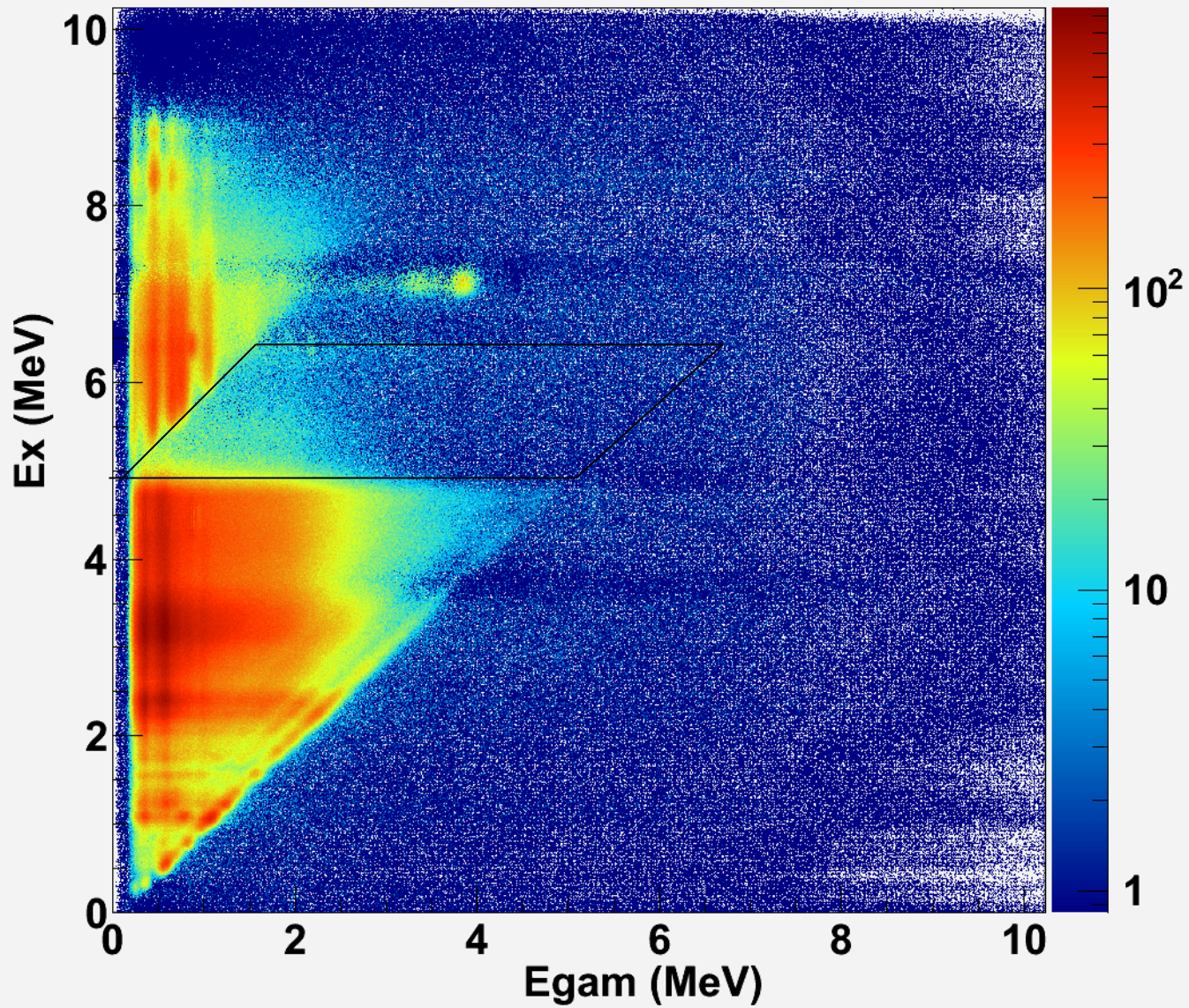


Cascade counting



$^{233}\text{Th}^*$



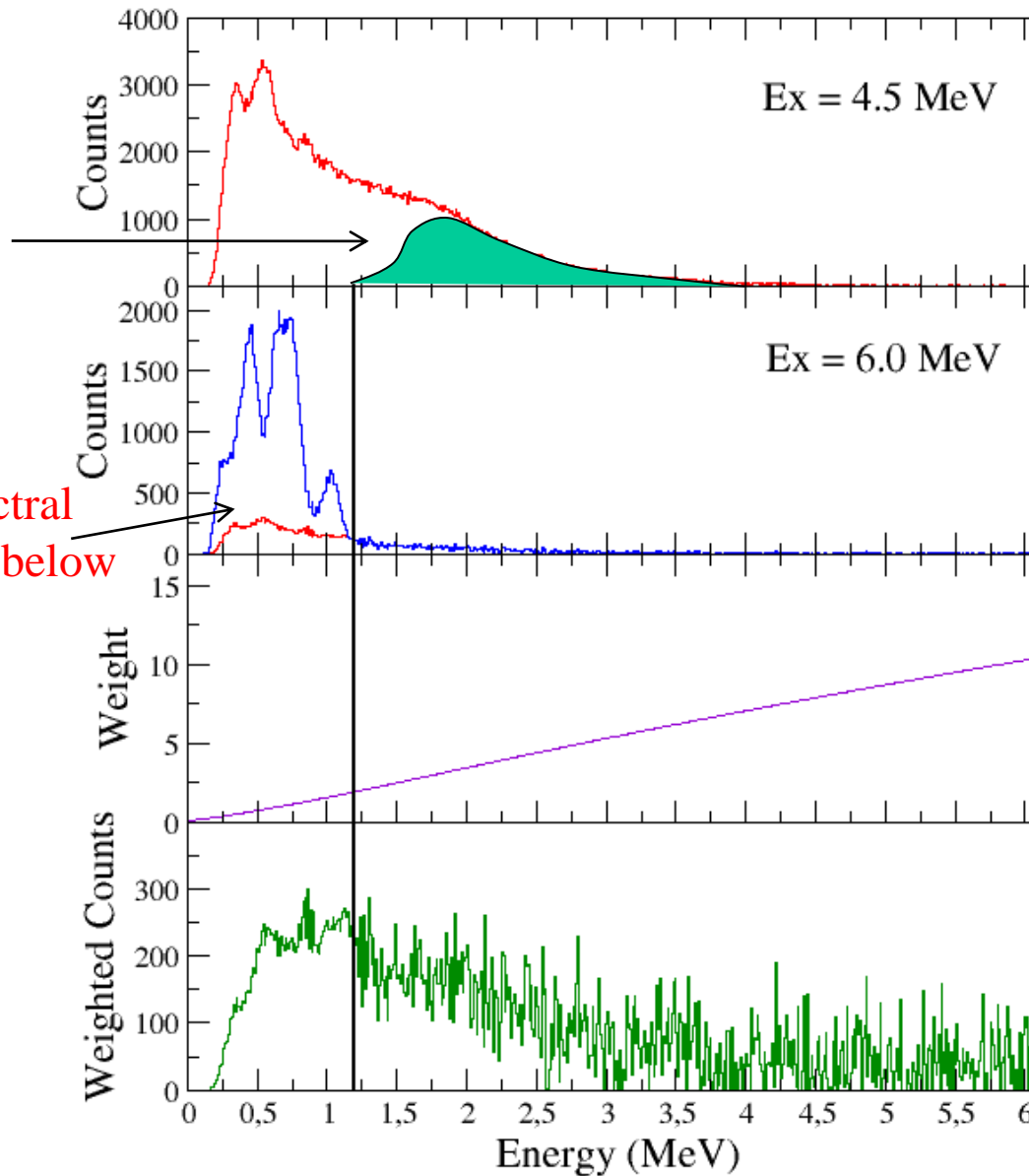


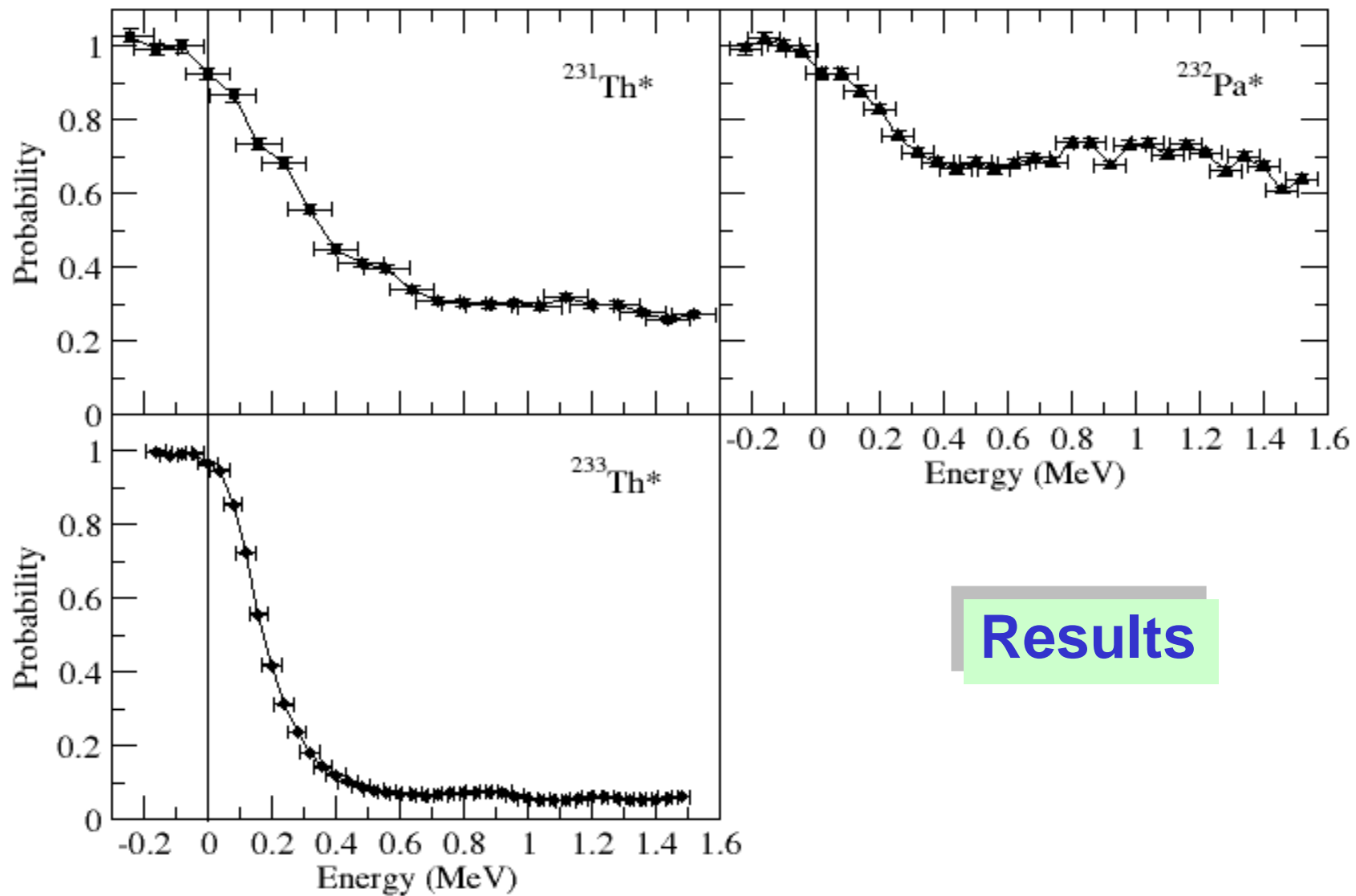
n,n' correction to measured decay probability

(New Technique)

Primary
gamma rays

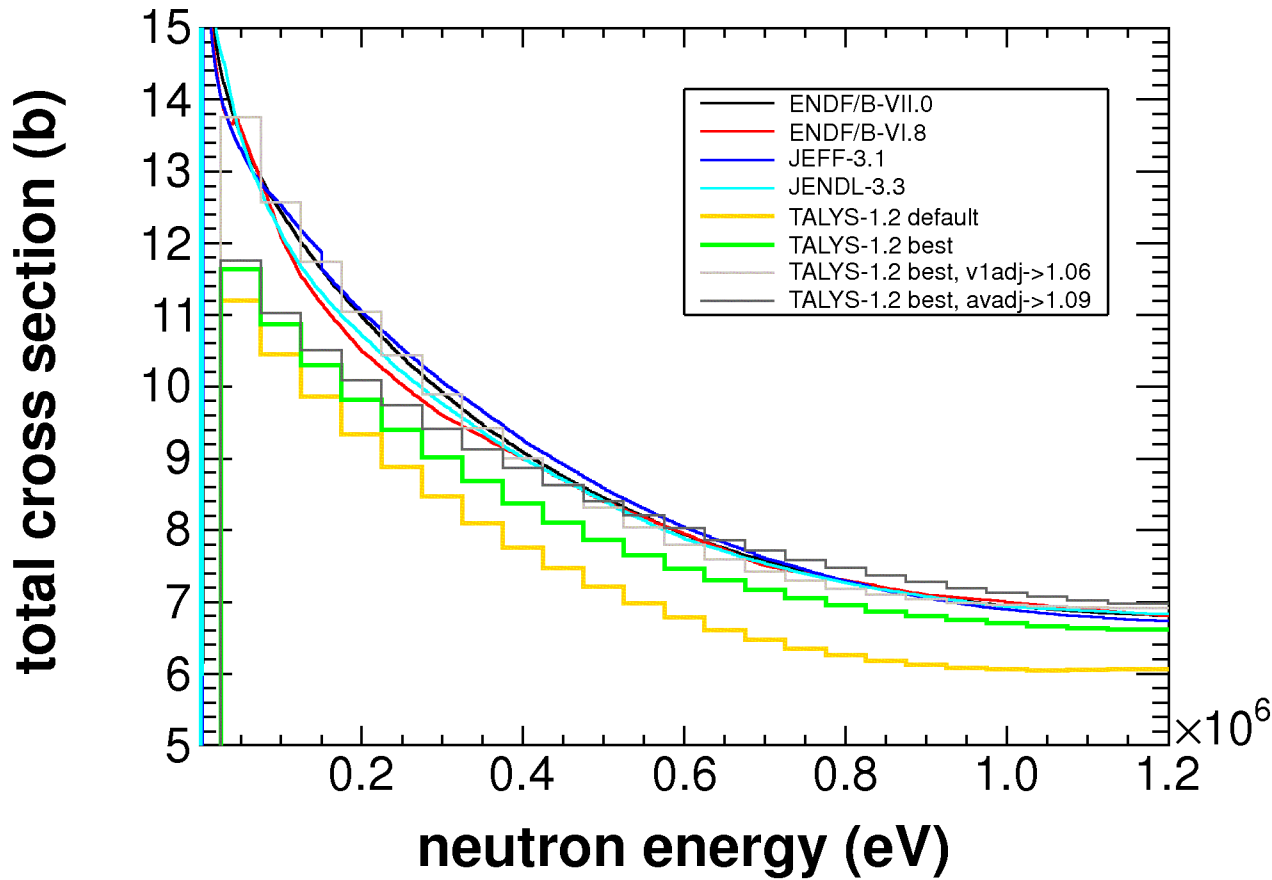
n,gamma spectral
extrapolation below
n,n'threshold





Results

TALYS calculations



Compare results with TALYS calculations: Since $^{232}\text{Th}(n,g)$ has been measured directly at nTOF we can perform a direct test of the optical model and the validity of the surrogate method.

Conclusions

Level density measurements possible in:

^{233}Th - d,p (0- 5 MeV)

^{232}Th - d,d' (0- 3.5 MeV)

^{231}Th - d,t (0 – 3 MeV)

Surrogate reaction & decay probability measurements possible in:

$^{233}\text{Th}^*$ - d,p

~~$^{230}\text{Ac}^*$ - d,a~~

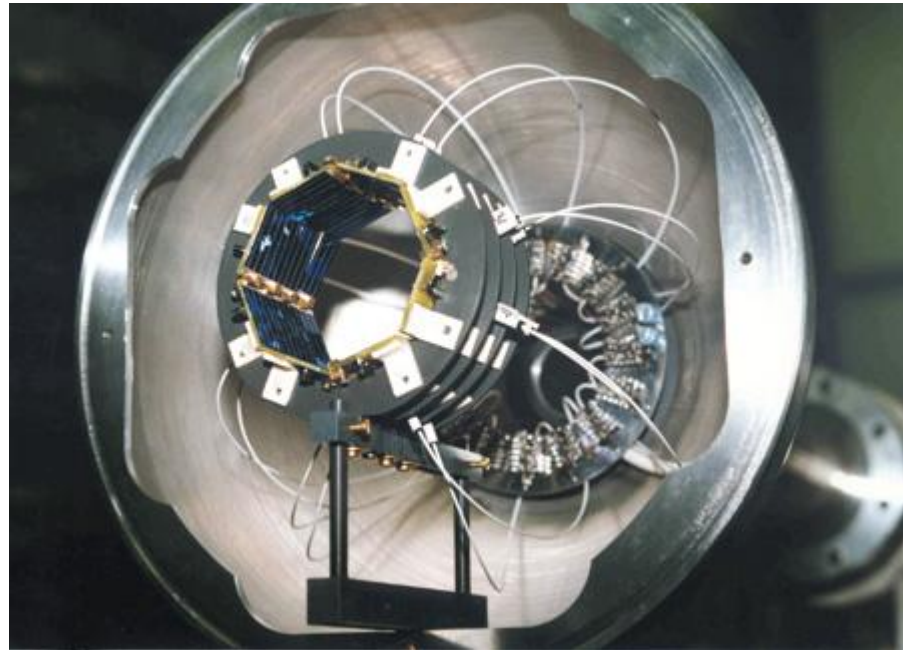
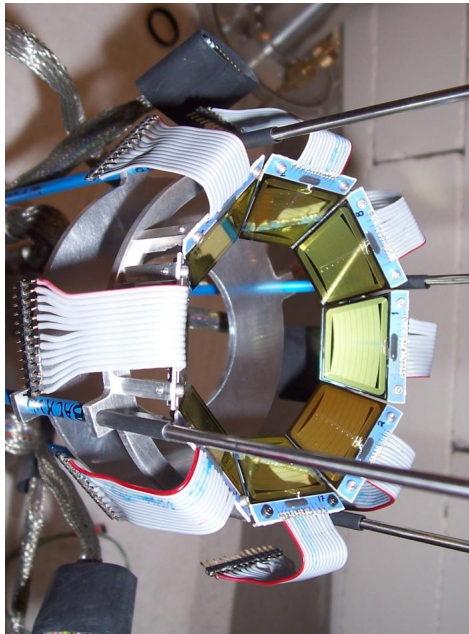
~~$^{234}\text{Pa}^*$ - ^3He ,p~~

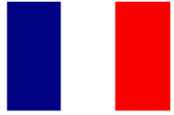
$^{232}\text{Pa}^*$ - ^3He ,t

^{231}Th - ^3He , ^4He

The Future

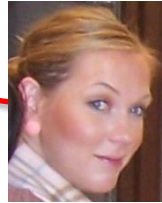
- ^{233}U , ^{231}Pa , ^{234}U targets
- Fission detectors





IPN Orsay

J.N. Wilson
B. Leniau

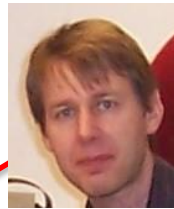


Oslo Cyclotron

S. Siem
S. Rose
A. Bürger
M. Guttormsen
A-C. Larsen
T. Wiborg
H. Toft

CEA Saclay

F. Gunsing
A. Georgen



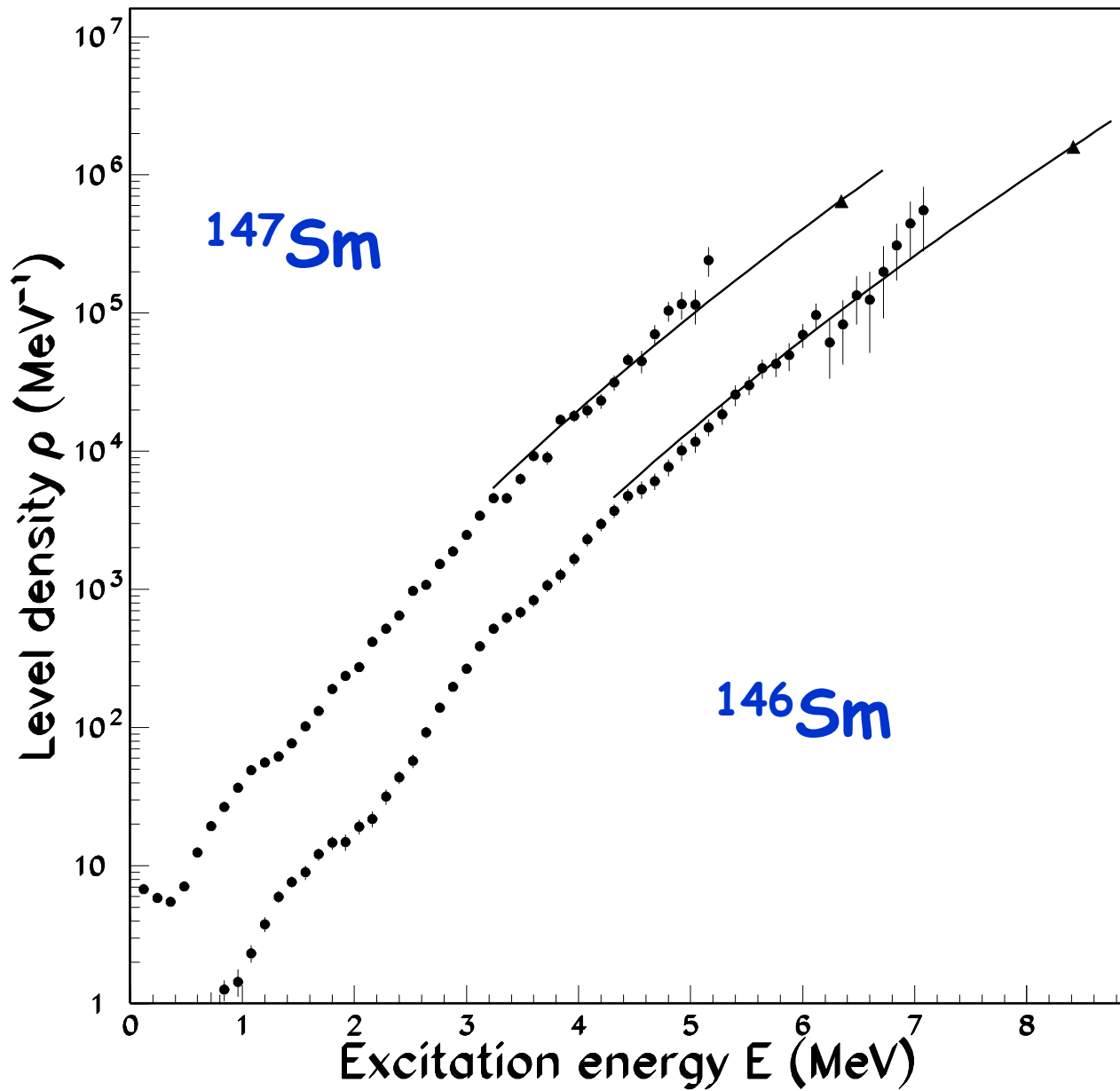
**Lawrence
Livermore Lab**

M. Weideking
L. Bernstein

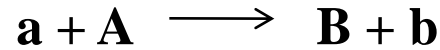
Thankyou

Takk

Merci



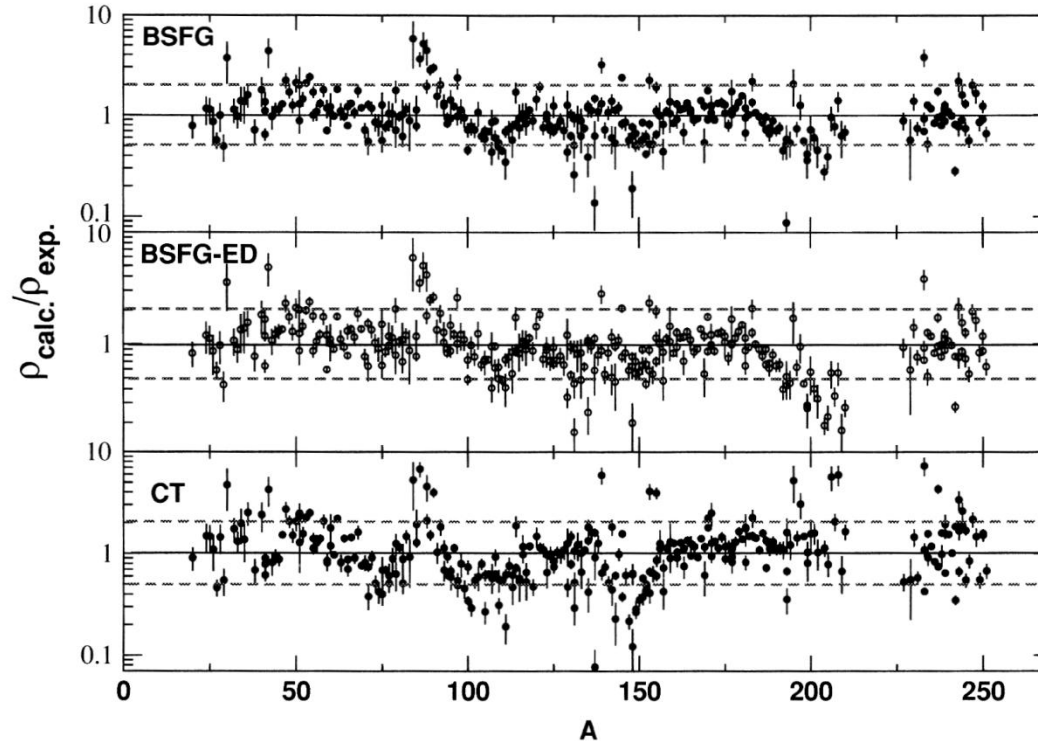
Hauser-Feshbach Formalism



$$\frac{d\sigma}{d\varepsilon_b}(\varepsilon_a, \varepsilon_b) = \sum_{J\pi} \sigma^{\text{CN}}(\varepsilon_a) \frac{\sum_{l\pi} \Gamma_b(U, J, \pi, E, l, \pi) \rho_b(E, l, \pi)}{\Gamma(U, J, \pi)}$$

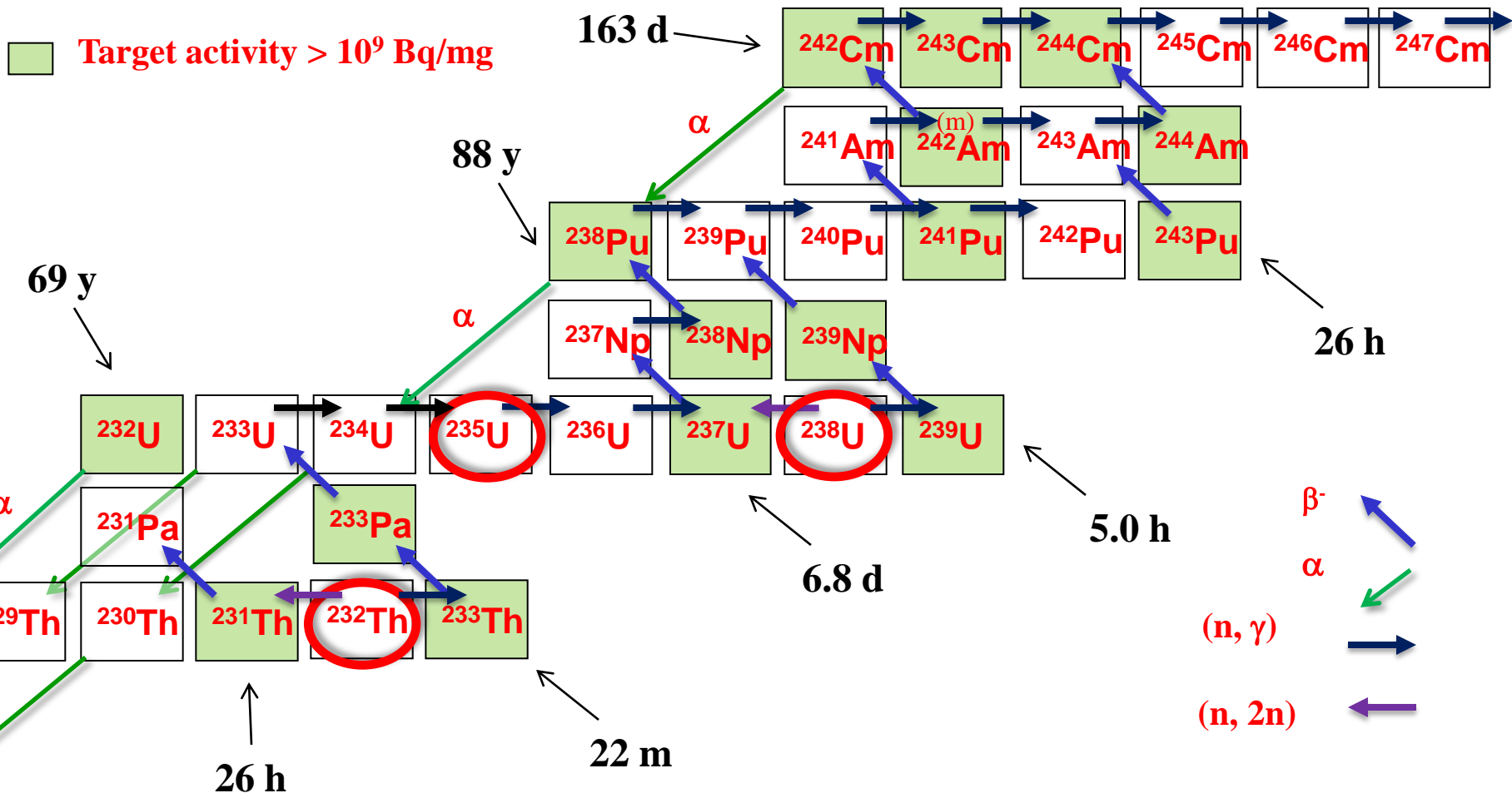
Differential cross section depends on transmission coefficients, Γ and level densities of the residual nuclei, ρ

Level Densities Theory and Experiment



- Level density changes due to collectivity, deformation etc.
- Large effects can occur over a small number of nucleons

Nuclei where direct measurements are difficult



$$\text{Decay probability} = \frac{\sigma_{\gamma}}{\sigma_c}$$

$$\sigma_T = \sigma_{\gamma} + \sigma_{n,n'} + \sigma_f$$

$$\sigma_T = \sigma_c + \sigma_{pot}$$

$$\sigma_{n,n'} = \sigma_{n,n'} \text{ compound} + \sigma_{n,n'} \text{ direct}$$

$$\sigma_{n,n'} = \sigma_{n,n'} \text{ compound} + \sigma_{pot}$$