# Progress with the Minimum Bias modeling in Herwig++

Minimum Bias and Underlying Event Working Group

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Karlsruhe Institute of Technology

CERN, 7 September 2010

#### Outline

#### This talk:

- ▶ Introduction Underlying event in Herwig++
- ▶ New data! ATLAS @ 900 GeV and @ 7 TeV
- Colour structure
- ▶ Outlook

## Underlying event in Herwig++

#### UA5 model (deprecated, only for reference)

▶ Included from Herwig++ 2.0.

[Herwig++, hep-ph/0609306]

- ► Little predictive power.
- ▶ Was default in fHerwig. Superseded by JIMMY

[JM Butterworth, JR Forshaw, MH Seymour, ZP C72 637 (1996)]

### Underlying event in Herwig++

[Herwig++, 0711.3137]

#### Semihard UE

- ▶ Default from Herwig++ 2.1.
  - Multiple hard interactions,  $p_t \geqslant p_t^{min}$  [Bähr, Gieseke, Seymour, JHEP 0807:076]
- ► Similar to JIMMY
- ► Good description of harder Run I UE data (Jet20).

### Underlying event in Herwig++

#### Semihard+Soft UE

Default from Herwig++ 2.3.

- [Herwig++, 0812.0529]
- Extension to soft interactions,  $p_t \leqslant p_t^{min}$  [Bähr, Gieseke, Seymour, JHEP 0807:076]
- ▶ Theoretical work with simplest possible extension.

[Bähr, Butterworth, Seymour, JHEP 0901:065]

"Hot Spot" model.

[Bähr, Butterworth, Gieseke, Seymour, 0905.4671]

Starting point: hard inclusive jet cross section.

$$\sigma^{\rm inc}(s; p_t^{\rm min}) = \sum_{i,j} \int_{p_t^{\rm min^2}} \mathrm{d}p_t^2 f_{i/h_1}(x_1, \mu^2) \otimes \frac{\mathrm{d}\hat{\sigma}_{i,j}}{\mathrm{d}p_t^2} \otimes f_{j/h_2}(x_2, \mu^2) \,,$$

 $\sigma^{\rm inc} > \sigma_{\rm tot}$  eventually (for moderately small  $p_t^{\rm min}$ ).

Interpretation:  $\sigma^{\text{inc}}$  counts *all* partonic scatters that happen during a single *pp* collision  $\Rightarrow$  more than a single interaction.

$$\sigma^{\rm inc} = \bar{n}\sigma_{\rm inel}$$
.

Use eikonal approximation (= independent scatters). Leads to Poisson distribution of number m of additional scatters,

$$P_m(\vec{b},s) = \frac{\bar{n}(\vec{b},s)^m}{m!} e^{-\bar{n}(\vec{b},s)} .$$

Then we get  $\sigma_{\text{inel}}$ :

$$\sigma_{\rm inel} = \int \mathrm{d}^2 \vec{b} \, \sum_{n=1}^{\infty} P_m(\vec{b}, s) = \int \mathrm{d}^2 \vec{b} \left( 1 - \mathrm{e}^{-\vec{n}(\vec{b}, s)} \right) \; .$$

Cf.  $\sigma_{\text{inel}}$  from scattering theory in eikonal approx. with scattering amplitude  $a(\vec{b},s)=\frac{1}{2i}(e^{-\chi(\vec{b},s)}-1)$ 

$$\sigma_{inel} = \int d^2\vec{b} \left(1 - e^{-2\chi(\vec{b},s)}\right) \qquad \Rightarrow \quad \chi(\vec{b},s) = \tfrac{1}{2} \vec{n}(\vec{b},s) \; . \label{eq:sigma-inel}$$

 $\chi(\vec{b},s)$  is called *eikonal* function.

From parton model assumptions we get:  $\bar{n}(\overrightarrow{b},s) = A(\overrightarrow{b})\sigma^{inc}(s;p_t^{min})$ 

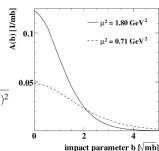
$$A(b) = \int d^{2}\vec{b}'G_{A}(|\vec{b}'|)G_{B}(|\vec{b} - \vec{b}'|)$$

 $G(\vec{b})$  from electromagnetic FF:

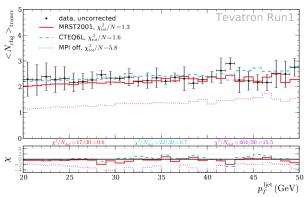
$$G_p(\vec{b}) = G_{\vec{p}}(\vec{b}) = \int \frac{\mathrm{d}^2 \vec{k}}{(2\pi)^2} \frac{\mathrm{e}^{i\vec{k}\cdot\vec{b}}}{(1+\vec{k}^2/\mu^2)^2}$$

But  $\mu^2$  not fixed to the electromagnetic 0.71 GeV<sup>2</sup>. Free for colour charges.





#### Good description of Run I Underlying event data ( $\chi^2 = 1.3$ ).



Only  $p_T^{ljet} > 20 \,\text{GeV}$ .

So far only hard MPI. Now extend to soft interactions with

$$\chi_{\text{tot}} = \chi_{OCD} + \chi_{\text{soft}}$$
.

Similar structures of eikonal functions:

$$\chi_{\text{soft}} = \frac{1}{2} A_{\text{soft}}(\vec{b}) \sigma_{\text{soft}}^{\text{inc}}$$

Simplest possible choice:  $A_{\rm soft}(\vec{b};\mu)=A_{\rm hard}(\vec{b};\mu)=A(\vec{b};\mu).$  Then

$$\chi_{\text{tot}} = \frac{A(\vec{b}; \mu)}{2} \left( \sigma_{\text{hard}}^{\text{inc}} + \sigma_{\text{soft}}^{\text{inc}} \right) .$$

One new parameter  $\sigma_{\text{soft}}^{\text{inc}}$ .

Taking the Tevatron data together with the wide range of possible values of  $\sigma_{tot}$  considered at LHC, we see that this model is to simple.

Extension: Relax the constraint of identical overlap functions:

$$A_{soft}(b) = A(b, \mu_{soft})$$

Fix the two parameters  $\mu_{\text{soft}}$  and  $\sigma_{\text{soft}}^{\text{inc}}$  in

$$\chi_{\rm tot}(\vec{b},s) = \frac{1}{2} \left( A(\vec{b};\mu) \sigma^{\rm inc} {\rm hard}(s;p_t^{\rm min}) + A(\vec{b};\mu_{\rm soft}) \sigma^{\rm inc}_{\rm soft} \right)$$

from two constraints. Require simultaneous description of  $\sigma_{\rm tot}$  and  $b_{\rm el}$  (measured/well predicted),

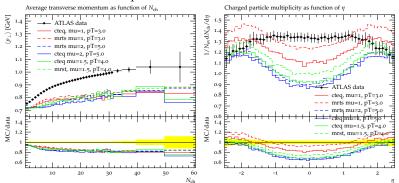
$$\begin{split} &\sigma_{tot}(s) \stackrel{!}{=} 2 \int d^2\vec{b} \left(1 - \mathrm{e}^{-\chi_{tot}(\vec{b},s)}\right) \,, \\ &b_{el}(s) \stackrel{!}{=} \int d^2\vec{b} \frac{b^2}{\sigma_{tot}} \left(1 - \mathrm{e}^{-\chi_{tot}(\vec{b},s)}\right) \,. \end{split}$$

#### Sum up:

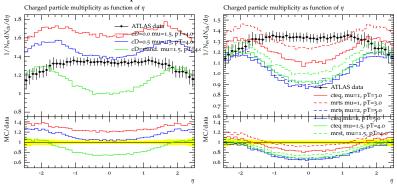
 $\Rightarrow$ at the end of the day we have two main parameters:  $\mu^2$ ,  $p_t^{min}$ .

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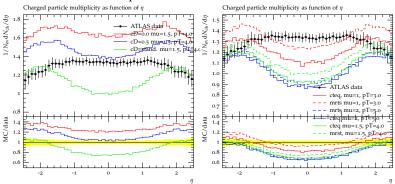


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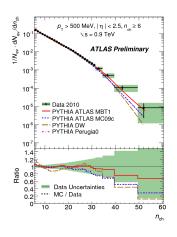


Colour structure of soft events.  $p_{disruvt} = \text{probability of disruption (default} = 1, \text{ completely disconected)}.$ 

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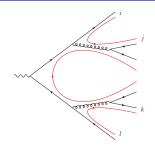


- Colour structure of soft events.  $p_{disrupt}$  = probability of disruption (default = 1, completely disconected).
- ▶ Problem: diffraction  $\Rightarrow$  Diffractive suppressed with cut:  $N_{ch} \ge 6$



- ► We used a diffractive suppressed sample with cut:  $N_{ch} \ge 6$
- Attention: The ATLAS graphs for N<sub>ch</sub> ≥ 6 are public, but the data points are not. We read the data points from the plots using:
  - ► EasyNData Peter Uwer [arXiv:0710.2896]
  - DataThief B. Tummers, http://datathief.org/
  - g3data J. Frantz, http://www.frantz.fi/software/g3data.php
  - some other tricks ...
  - question to the collaborations: can we do something about this?

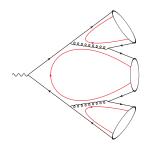
I am happy to provide data points with corresponding Rivet analyses if someone needs it.



Extending the hadronization model in Herwig(++):

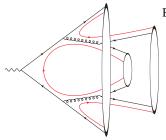
▶ QCD parton showers provide pre-confinement ⇒ colour-anticolour pairs form highly excited hadronic states, the clusters

For details look at Christians Röhr's Diploma thesis



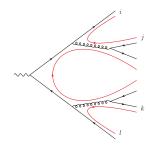
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- ► Physical motivation: exchange of soft gluons during non-perturbative hadronization phase



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- QCD parton showers provide pre-confinement
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- ► CR in the cluster hadronization model: allow reformation of clusters, e.g. (il) + (jk)
- ▶ Physical motivation: exchange of soft gluons during non-perturbative hadronization phase

### Implementation<sup>1</sup>

► Allow CR if the cluster mass decreases,

$$M_{il} + M_{ki} < M_{ii} + M_{kl},$$

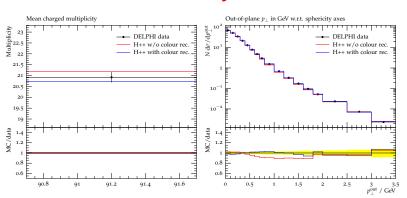
where  $M_{ab}^2 = (p_a + p_b)^2$  is the (squared) cluster mass

Accept alternative clustering with probability  $p_{\text{reco}}$  (model parameter)  $\Rightarrow$  this allows to switch on CR smoothly

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Can we still describe the LEP data similar to Herwig++ w/o colour reconnection?

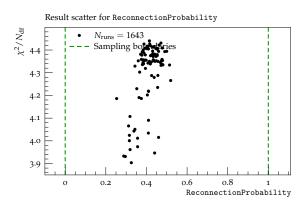
## Preliminary results



We repeated tuning of the hadronization to the LEP data (above 2 examples). Both tunings (with and w/o colour rec.) seems to describe the data at the same level.

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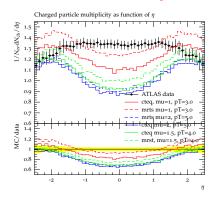
## Preliminary results



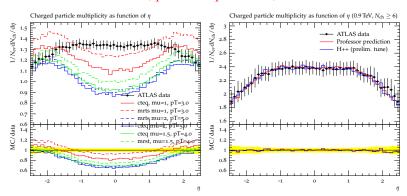
Prefered by LEP data is:  $0.2 \le p_{\text{reco}} \le 0.6$ 

## Preliminary results

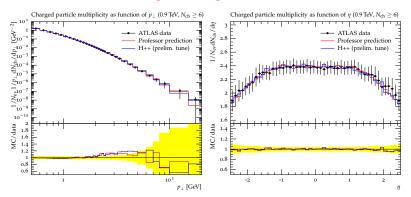
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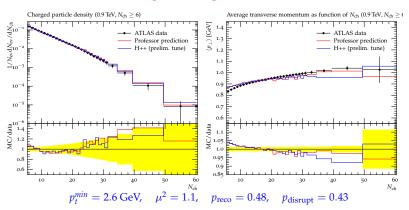


## Preliminary results



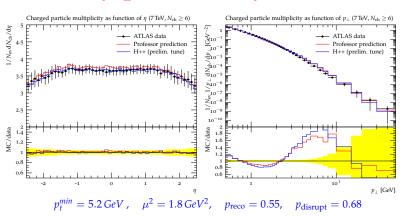
## Preliminary results

(space for improvement)



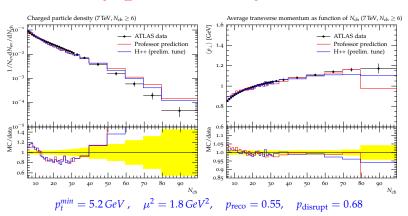
Many thanks to the Professor team for help and hints how to use their program!

## Very preliminary results



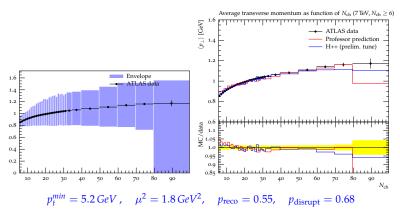
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### Summary

- New implementation of colour reconnection is in validation and seems to work very well!
- New "tuning" gives a good description of the ATLAS 900 GeV data!
- ATLAS 7TeV data will be investigated in more details soon.
- Still space for improvements: better LEP tune, treatment of remnants pdf, more involved overlap function, energy dependent parameters...
- Minimum bias/underlying event/diffraction under constant improvement!
- Stay tuned!