

Theory Perspectives

Goran Senjanovic
ICTP

talks by Altarelli, Djouadi, Fayet, Nesti,...

Message:

- **NO** theory that **predicts** new physics at LHC energies
- Connected with **NO** problem in the SM
(except neutrino mass -but that can be higher energy)

Thank you !

Well, I mean...

SM: in spite of great success **No Higgs yet**



LHC is a Higgs machine

One could wait to see if it exists...

- instead, one imagines it is there
- then ask: *why* $m_H \ll M_{GUT} (M_{Pl})$?

Funny question,
since:

- GUT never seen
- M_{Pl} not a physical scale

$$M_{Pl} \equiv 1/\sqrt{G_N} \simeq 10^{19} GeV$$

maybe $G_N = \frac{g^2}{M_F^2} \quad (g = M_F/M_{Pl})$

$g \ll 1 \implies M_F \ll M_{Pl}$ Glashow, '85

Realization: Large Extra Dimensions (LED)

$$M_{Pl}^2 = M_F^{2+n} R^n \rightsquigarrow g = (M_F R)^{-n/2}$$

$g \ll 1 \longleftarrow R \gg M_F^{-1} \sim (TeV)^{-1}$

Arkani-Hamed, Dimopoulos, Dvali (ADD), '98

Large number of species

of states : $N = (M_F R)^n = \left(\frac{M_{Pl}}{M_F} \right)^2$

the tower: $\frac{1}{R}, \frac{2}{R}, \frac{3}{R}, \dots, \frac{N}{R} \simeq M_F$

$$M_F \simeq TeV \quad \Rightarrow \quad N \simeq 10^{32}$$

Dvali, '07 - '10

all that matters : large # of species

$$M_F \simeq \frac{M_{Pl}}{\sqrt{N}}$$

$M_F \rightarrow$ the scale of strong gravity

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Dvali, '07 - '10

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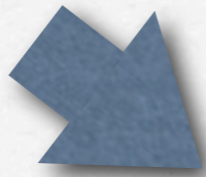
$$M_F \simeq \frac{M_{Pl}}{\sqrt{N}}$$

$M_F \rightarrow$ the scale of strong gravity

Black holes @ LHC ?

$$M_F \simeq TeV$$

ADD, '98



no hierarchy problem

Physics of black holes not
well understood:

how to distinguish from particles ?

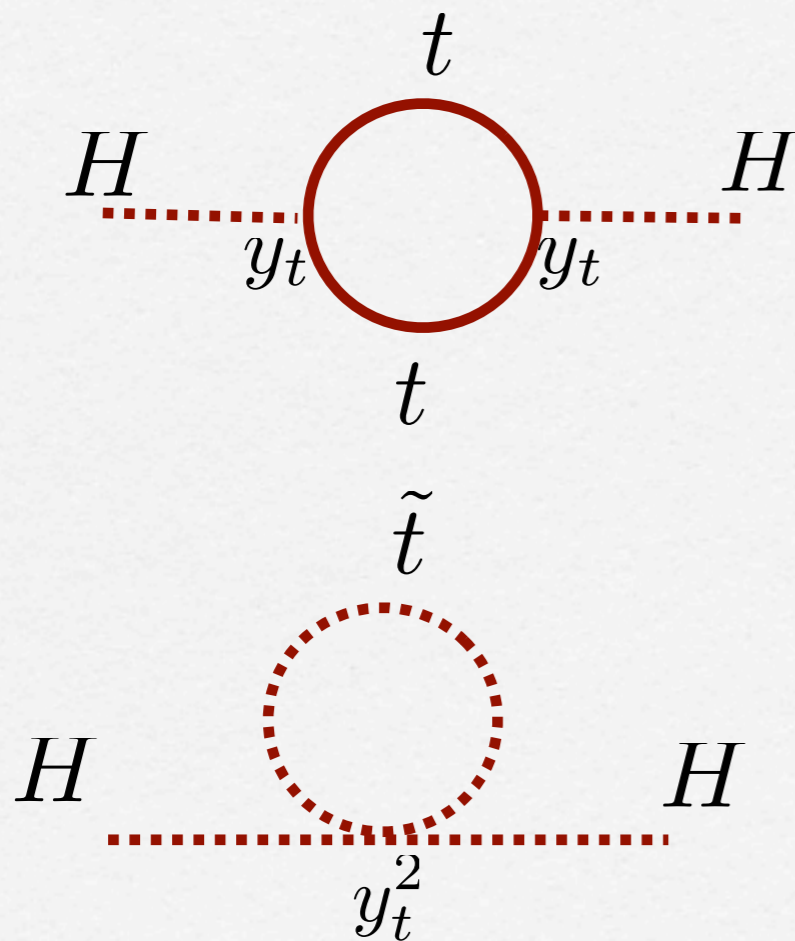
Dvali, Gomez, Mukhanov '10

but the question was:

Why is Higgs light?

Hierarchy problem: SUSY@LHC ?

talk by Fayet



+

$$\delta m_H^2 = \frac{y_t^2}{16\pi^2} (\Lambda^2 + m_t^2)$$

M_{GUT}, M_{Pl}

$$\delta m_H^2 = -\frac{y_t^2}{16\pi^2} (\Lambda^2 + m_{\tilde{t}}^2)$$

$$\delta m_H^2 = -\frac{y_t^2}{16\pi^2} (m_{\tilde{t}}^2 - m_t^2)$$

Maiani; Fayet; Witten '79

large cutoff: $\Lambda \geq 10^{16} \text{ GeV}$

low energy supersymmetry
tailor fit for GUT

$$m_{\tilde{t}} \leq \text{TeV}$$

$$\delta m_H^2 = -\frac{y_t^2}{16\pi^2} (m_{\tilde{t}}^2 - m_t^2)$$

small

$\delta m_H^2 < 0 \implies$ Higgs mechanism
for $y_t \simeq 1$

Alvarez-Gaume, Polchinski, Wise, '82
Inoue, Kakuto, Komatsu, Takeshita, '82

Low energy supersymmetry



unification of gauge couplings

* $\sin^2 \theta_W \simeq 0.23$

needed
 $m_t \simeq 200 \text{ GeV}$

Ibáñez, Ross, '80

Dimopoulos, Rabi, Wilczek, '80

Einhorn, Jones '81

Marciano, G.S., '81 *

$$\rho \simeq 1 + \frac{\alpha}{2\pi} \frac{m_t^2 - m_b^2}{M_W^2}$$

In order to boost $(\sin^2 \theta_W)_{exp} \simeq 0.21$

The light Higgs mass in the MSSM:

$$m_{h^0}^2 \leq M_Z^2 + \frac{3\alpha}{4\pi} \frac{m_t^4}{M_W^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + 3 \right]$$

$$m_{\tilde{t}} \simeq \text{TeV} \quad \Rightarrow \quad m_{h^0} \leq 135 \text{ GeV}$$

talk by Altarelli

Electro-weak baryogenesis:

needs a light stop, lighter than a top

neutrino mass:

R-parity violating couplings

dark matter:

or stable neutralino

long lived gravitino



new physics of
neutrino mass

R-parity and neutrino mass

$$lle^c + ql d^c + u^c d^c d^c$$

product small

neutrino mass

proton decay

all vanish (symmetry)

add ν^c

$lh\nu^c$

symmetry?

neutrino massless

Neutralino DM?

not so convincing in MSSM:

symmetry no guarantee for vanishing
couplings

gravitino necessarily long

lived, DM for $m_{3/2} \lesssim 10 - 100 \text{ GeV}$

Is this convincing ?

SUSY @ LHC ?

In Grand unification
there is fine tuning:

$$m_{h^0}^2(\text{tree}) = -m_0^2 + M_{GUT}^2 \simeq M_W^2$$

only loops are OK

If fine-tuning is a sin, SUSY helps you
to sin only once - and at tree level

□ There are solutions, but more like exercises in
Group Theory

□ Comment by Carl Sagan:

A. Ozpinaci '10



“...I can’t see a thing on the surface of Venus. Why not? Because it’s covered with a dense layer of clouds. Well, what are clouds made of? Water, of course. Therefore Venus must have an awful lot of water on it, therefore the surface must be wet. If the surface is wet there’s probably a swamp, if there’s a swamp there’s ferns, if there’s ferns... maybe there’s even dinosaurs.

**Observation: you couldn’t see a thing.
Conclusion: Dinosaurs “**

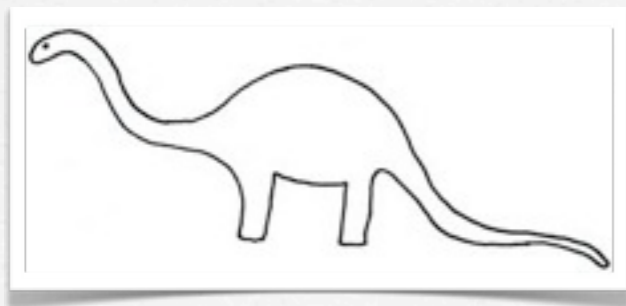
Is our motivation somewhat metaphysical?

We have seen nothing :

- No Higgs
- No new physics

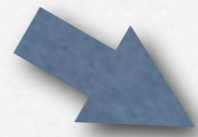
...and yet we speculate on

- for *every* particle a *superpartner*
- *the end of field theory* at TeV: black holes and such at LHC



Well, there **is** new physics :

Neutrino oscillations



Neutrino masses and mixings

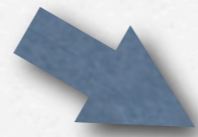
touches into the heart of an old
and fundamental question:

are neutrinos massive
and 'real' particles?

Majorana, '37

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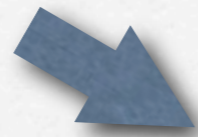
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Majorana, '37

last paper before
his disappearance

The Majorana Program:

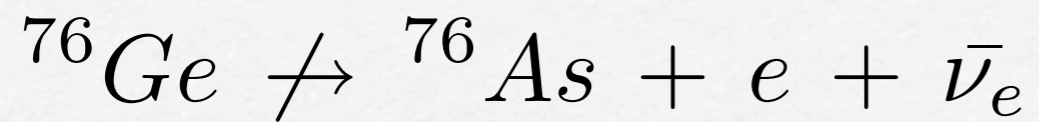
$$\nu_M = \nu_L + \nu_L^* \quad \Leftrightarrow \quad m_\nu^M (\nu_L \nu_L + h.c.)$$

$\Rightarrow \Delta L = 2$ violation of Lepton number

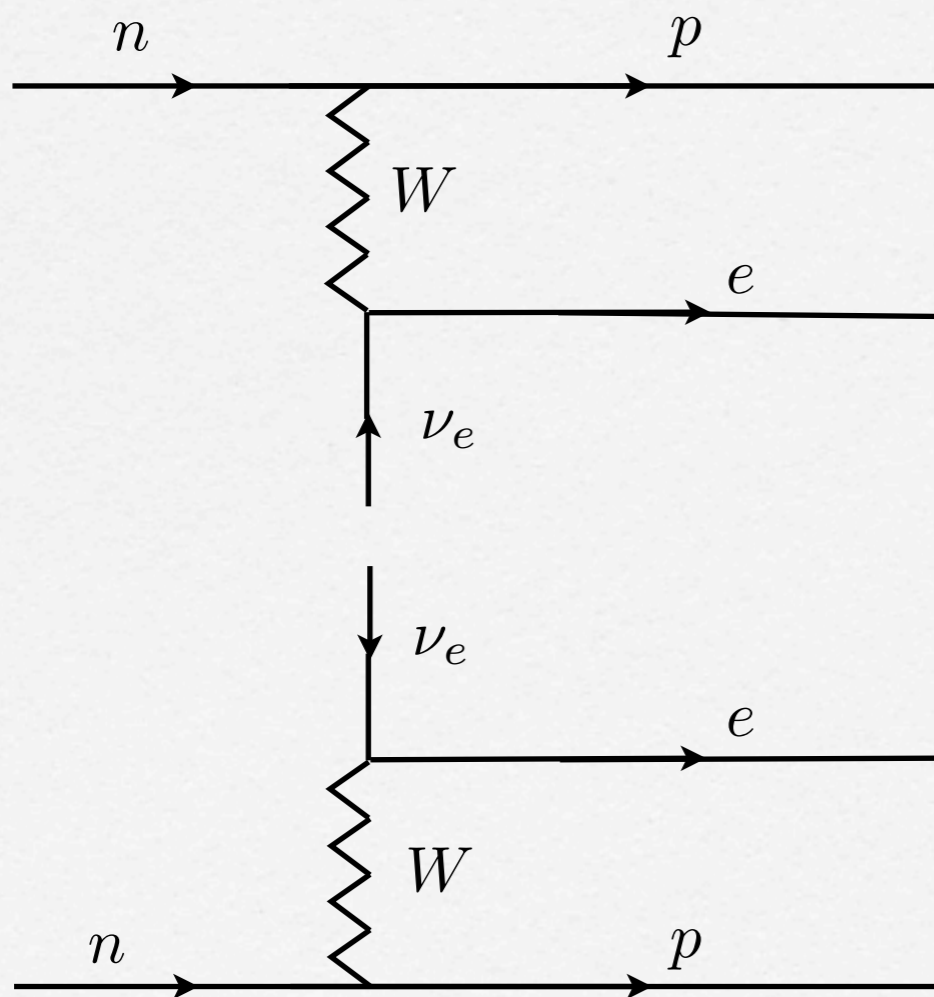
 neutrinoless double-beta decay $(0\nu 2\beta)$

Majorana, '37
Racah, '37 - Furry, '38

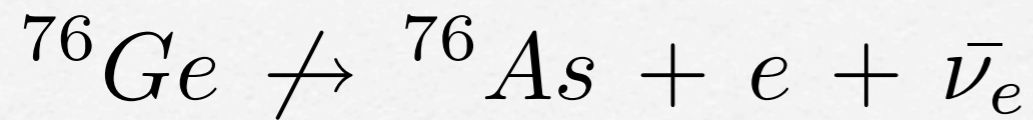
Double-beta decay



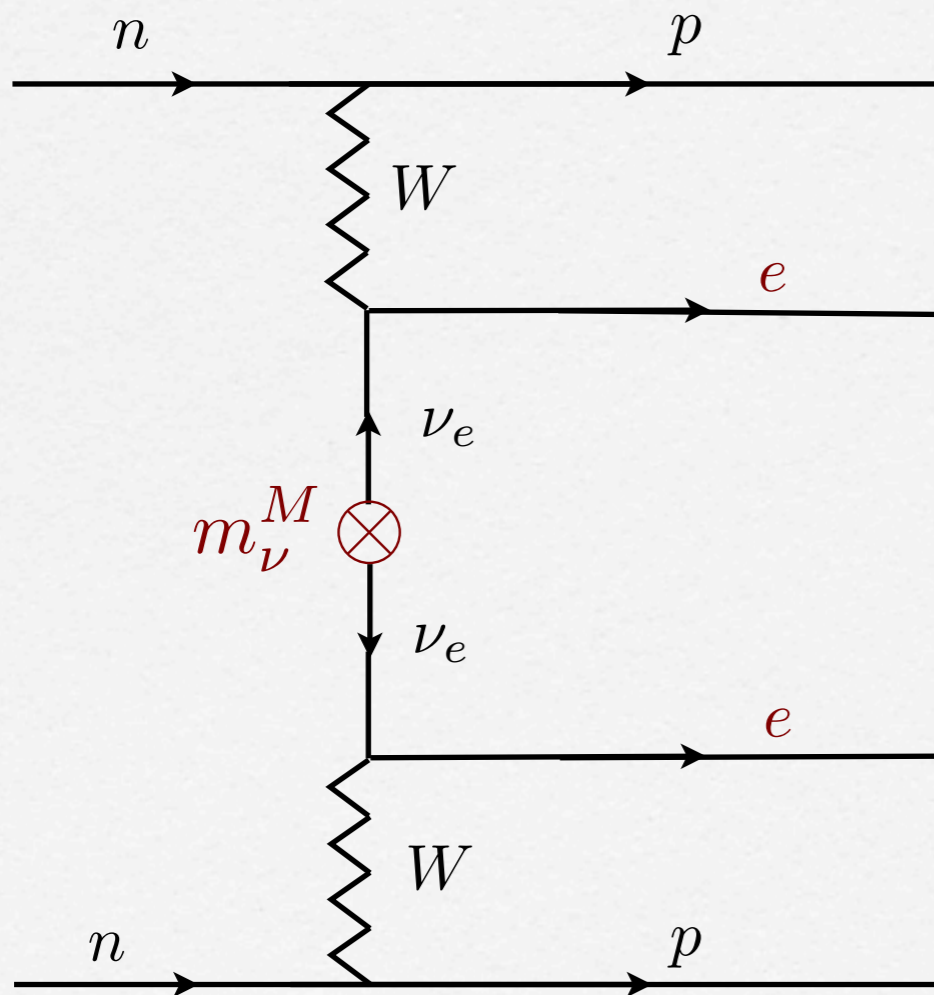
Göppert-Mayer, '35



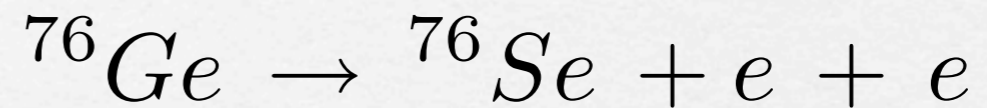
Double-beta decay



Göppert-Mayer, '35



electrons created
out of "nothing"



$$t_{1/2} \geq 10^{24} \text{ yr} \Rightarrow m_\nu^M \lesssim 1 \text{ eV}$$

Why is this relevant for LHC ?

Majorana mass forbidden in the SM

 window to new physics

Conventional wisdom:

add ν_R  see-saw

$$\begin{matrix} \nu_L \\ \nu_R \end{matrix} \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix} \Rightarrow$$

$M_R \gg m_D$

Why is this relevant for LHC ?

Majorana mass forbidden in the SM

➔ *Window to new physics*

Conventional wisdom:

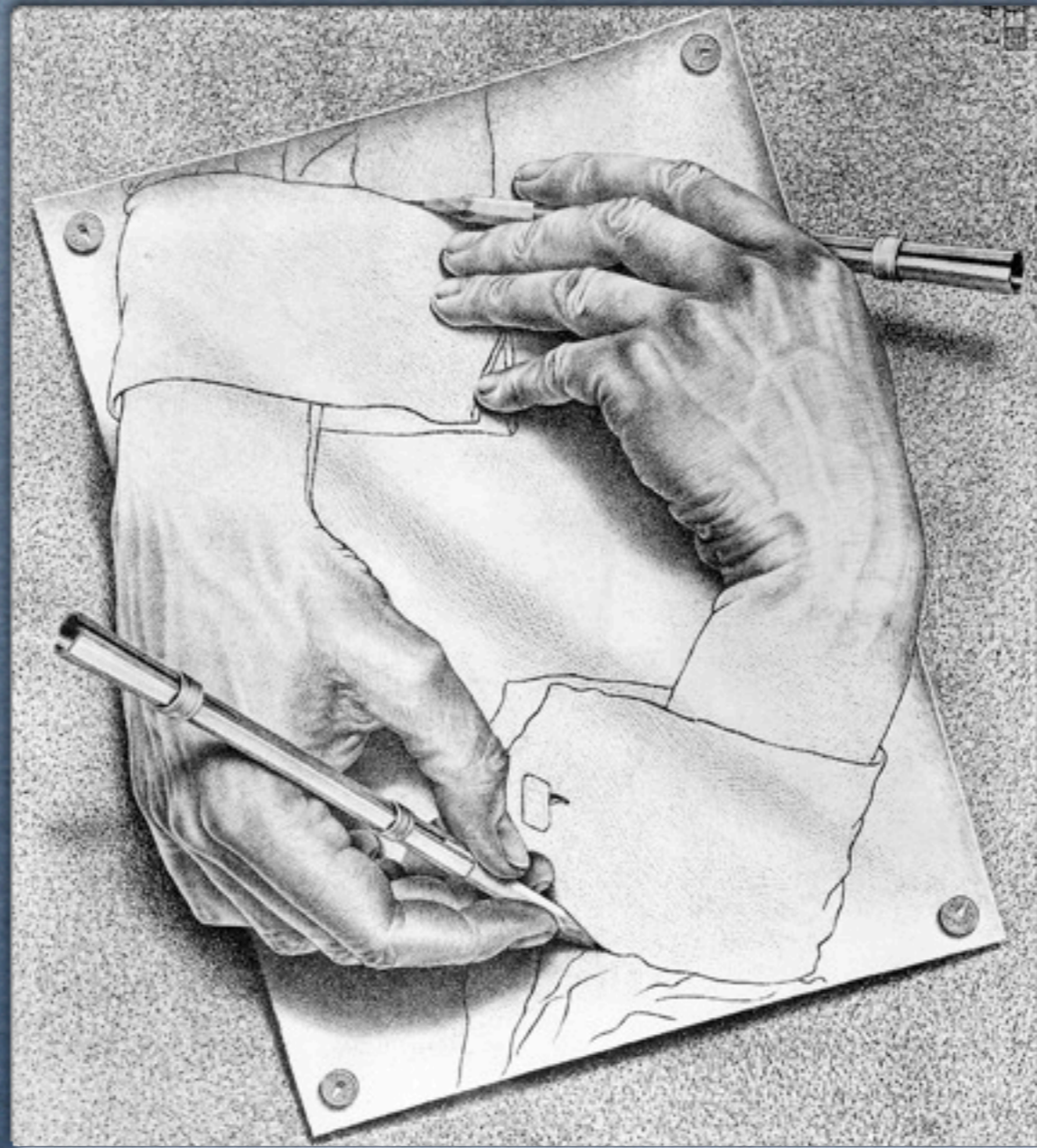
add ν_R ➔ see-saw

$$\begin{matrix} \nu_L \\ \nu_R \end{matrix} \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix} \Rightarrow m_\nu \simeq -\frac{m_D^2}{M_R}$$

$M_R \gg m_D$

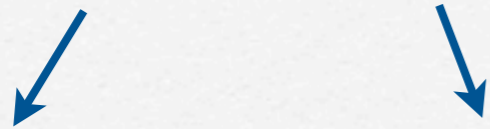


Majorana mass



Why ν_R ?

Left - Right symmetry



$$SU(2)_L \times U(1) \times SU(2)_R$$

$$\nu_L \Rightarrow \nu_R$$

Patil, Salam '74
Mohapatra, G.S. '75

talk by Nesti

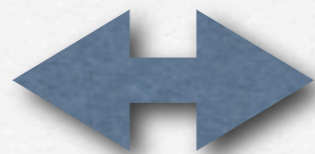
$$m_{\nu_R} \propto M_{W_R} \Rightarrow m_\nu \propto m_D^2 / M_{W_R}$$

$$\rightarrow 0$$

$$M_{W_R} \rightarrow \infty$$

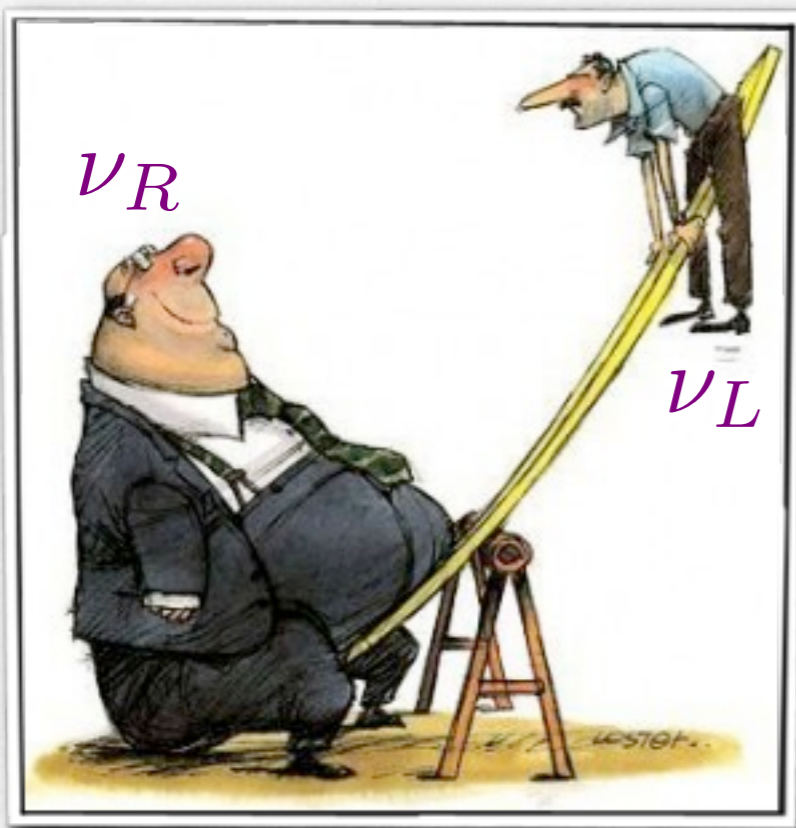
SM limit

neutrino
mass



Parity
violation

Minkowski '77
Mohapatra, G.S. '79



$$m_\nu \leq 1\text{eV} \longleftrightarrow M_{W_R} \geq 10^3 \text{ GeV}$$

see-saw mechanism

Minimal model:
theoretical limit

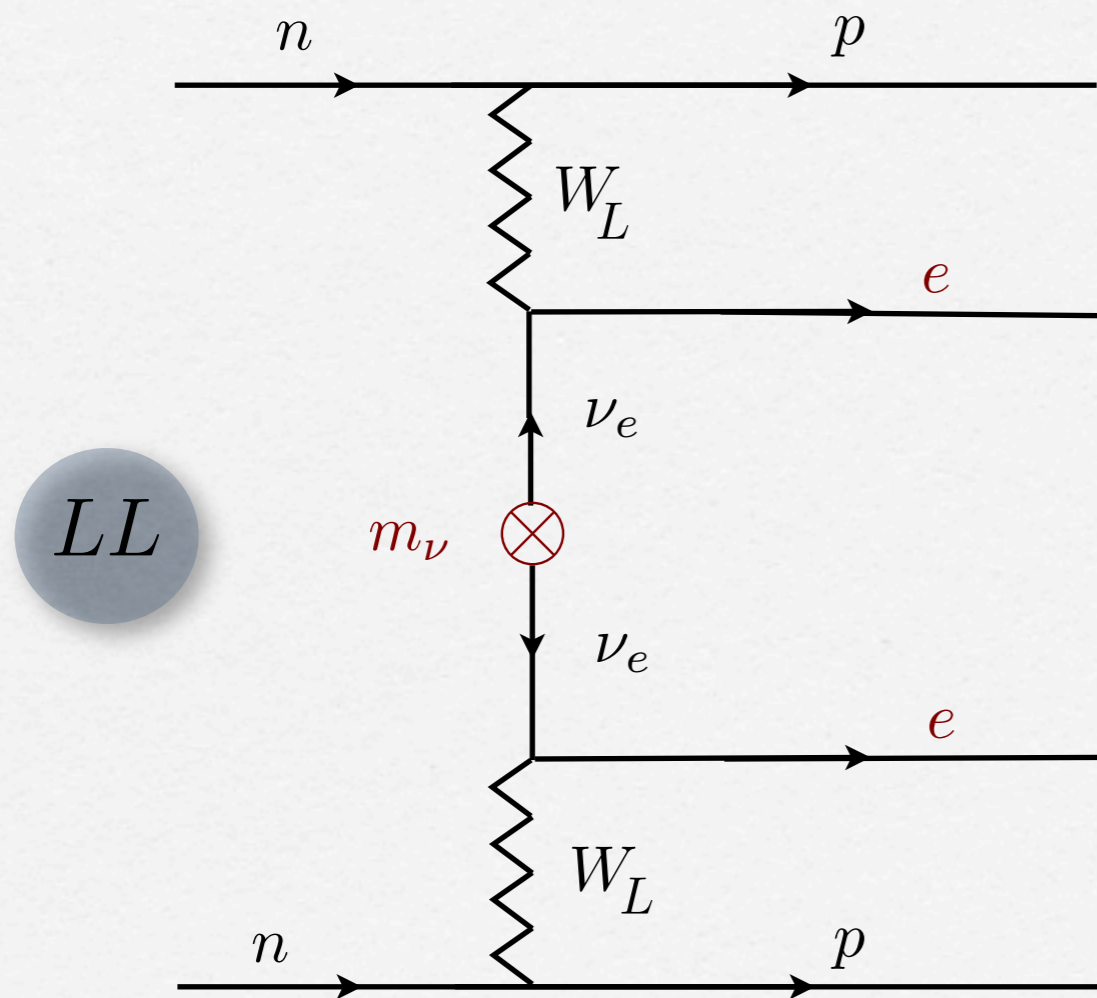


$$M_{W_R} \geq 2500 \text{ GeV}$$

Beall, Bander, Soni '81
Maiezza, Nemevsek, Nesti, G.S. '10

New source for $0\nu 2\beta$

Mohapatra, GS '81

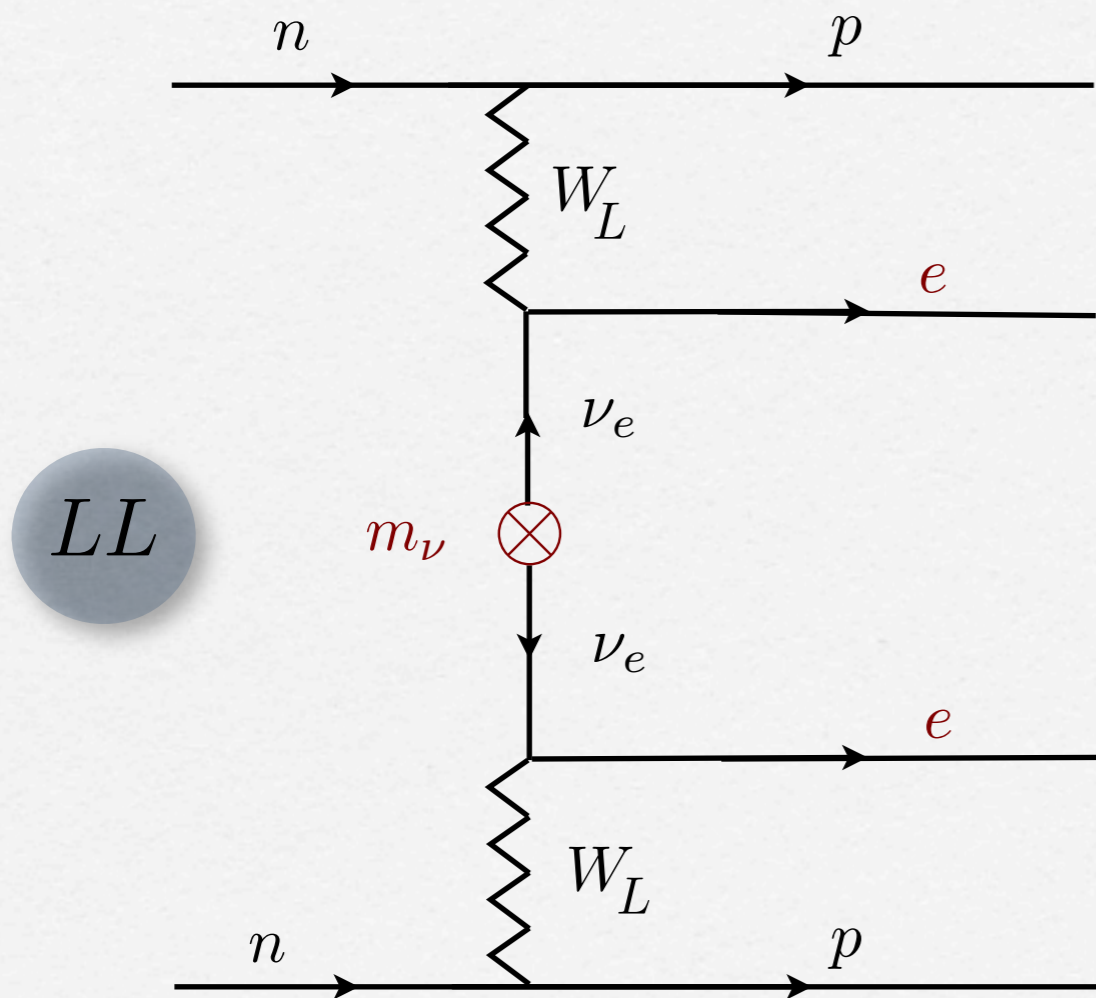


$$LL \propto \frac{1}{M_{W_L}^4} \frac{m_\nu}{p^2}$$

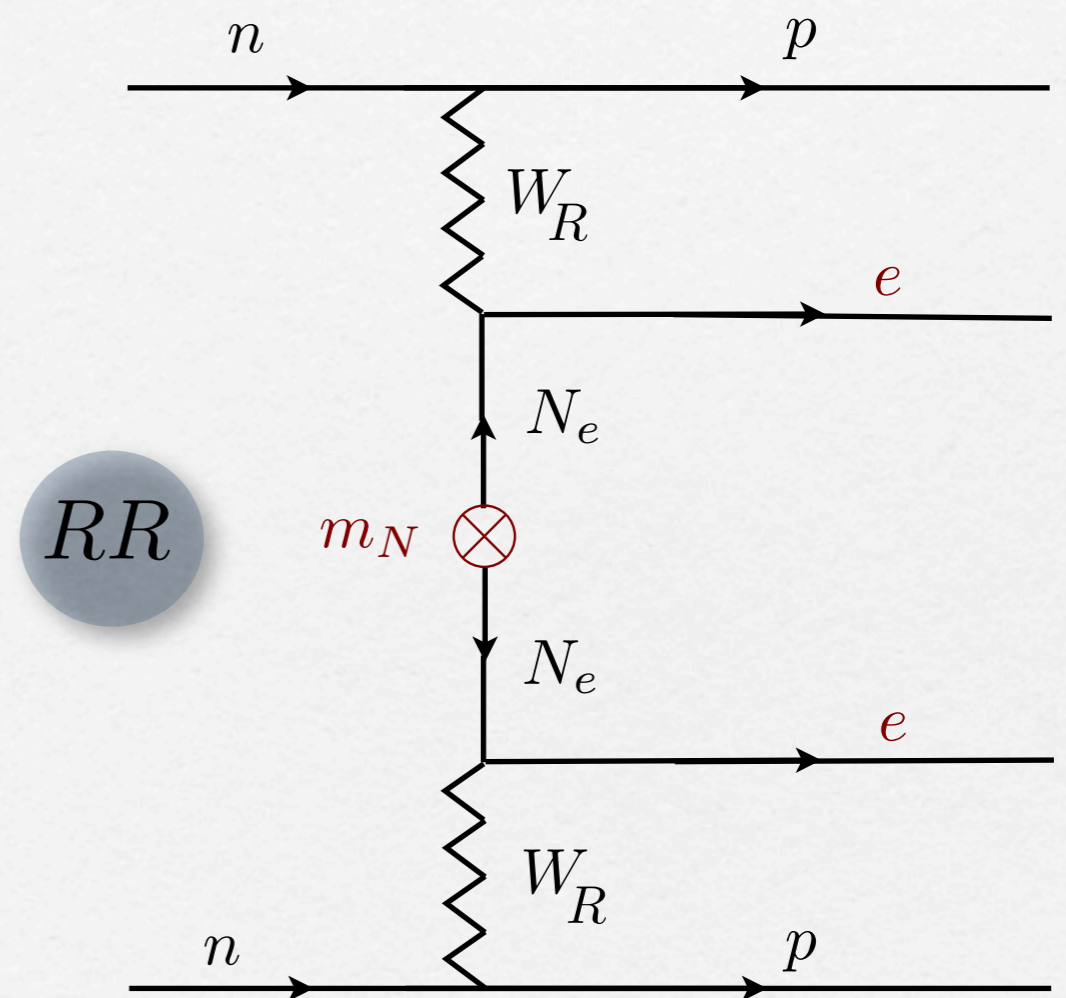
$$p \simeq 100 \text{ MeV}$$

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Mohapatra, GS '81



+



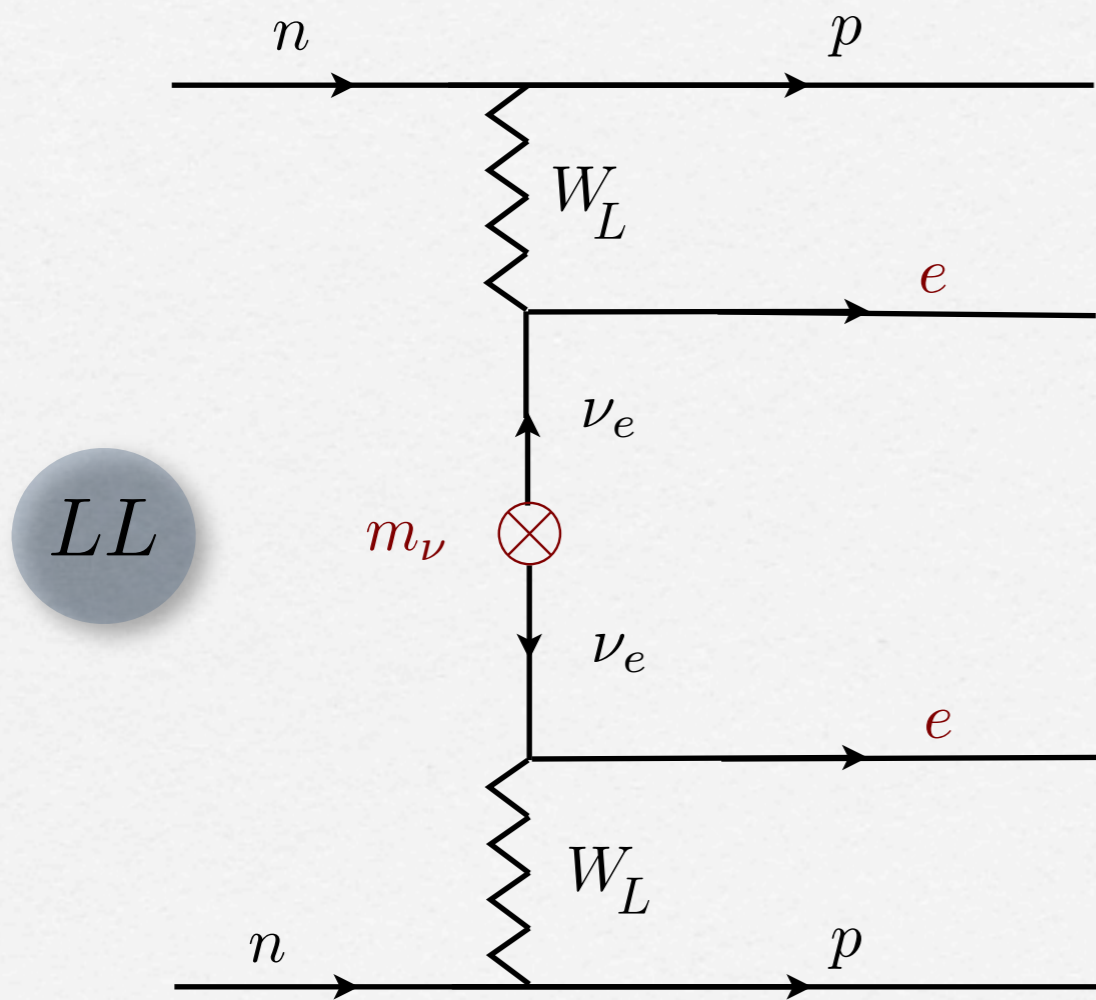
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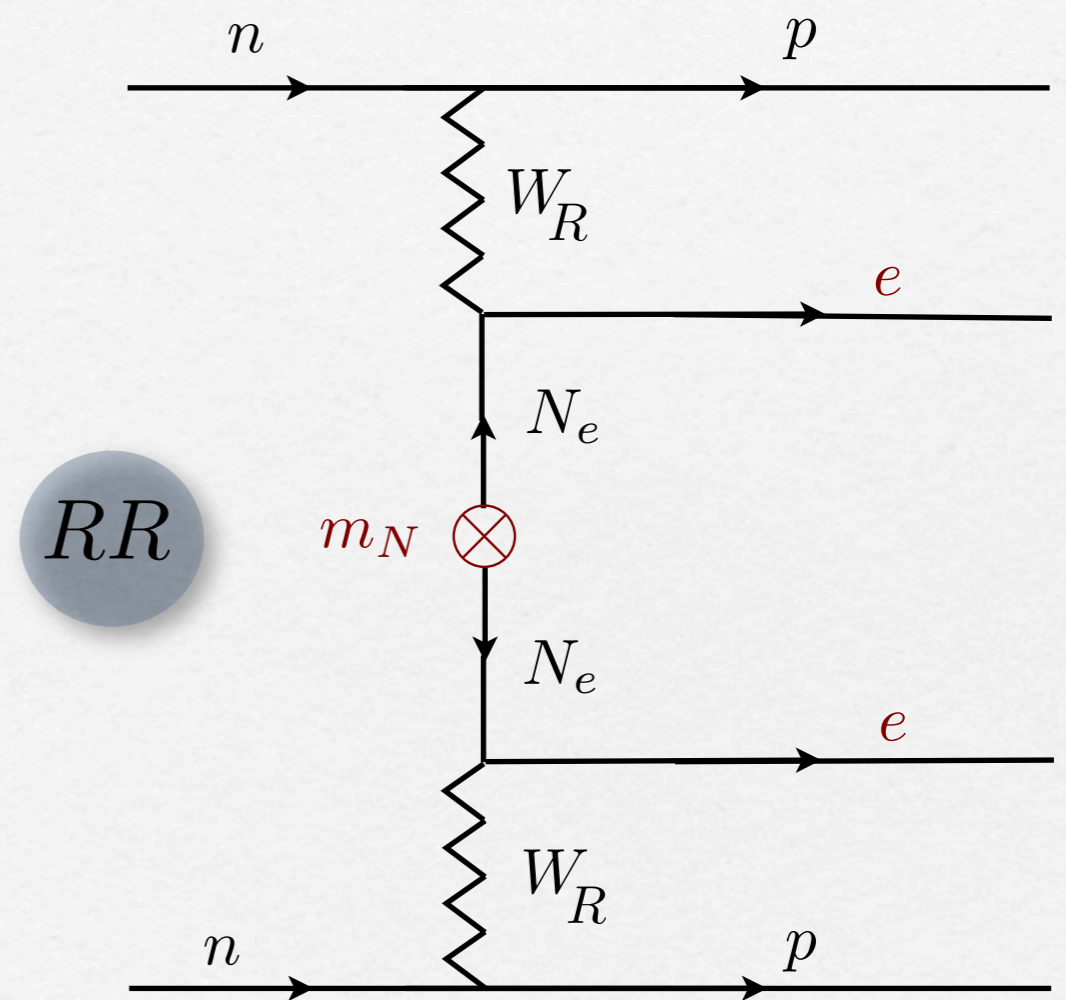
$$RR \propto \frac{1}{M_{W_R}^4} \frac{1}{m_N}$$

New source for $0\nu 2\beta$

Mohapatra, GS '81



+



$$LL \propto \frac{1}{M_{W_L}^4} \frac{m_\nu}{p^2}$$

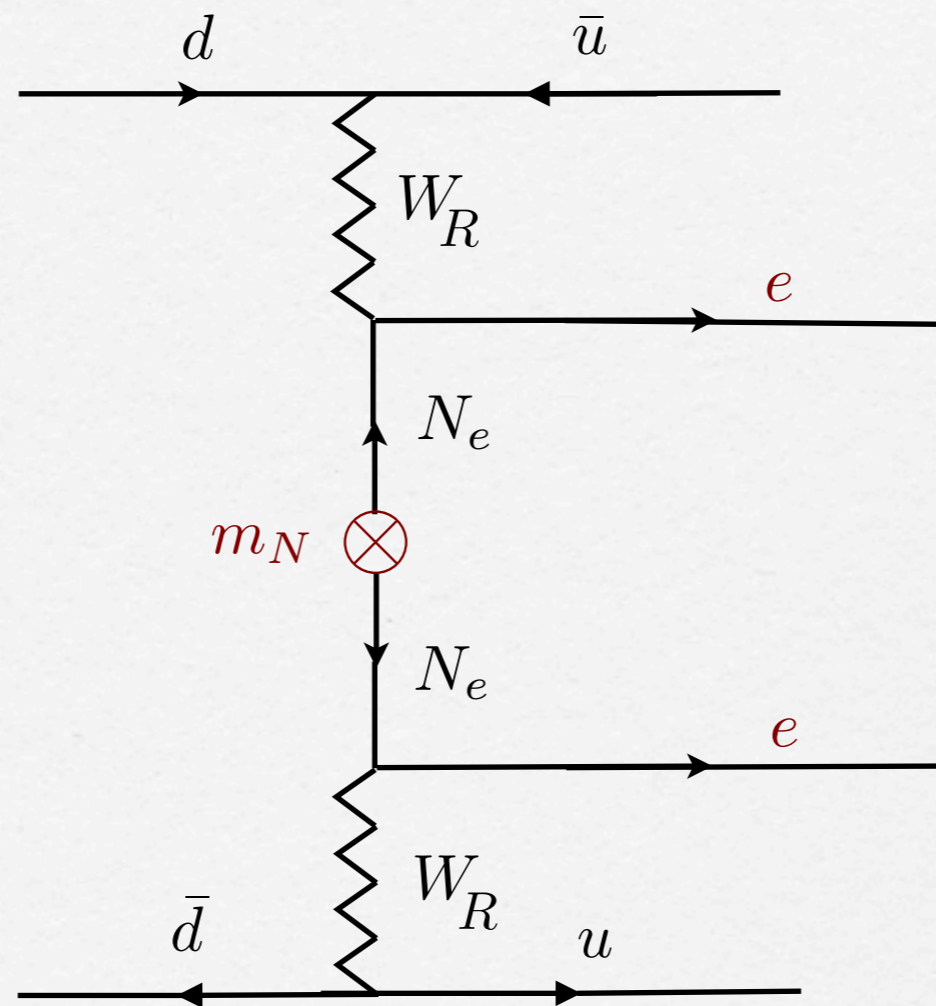
$$p \simeq 100 \text{ MeV}$$

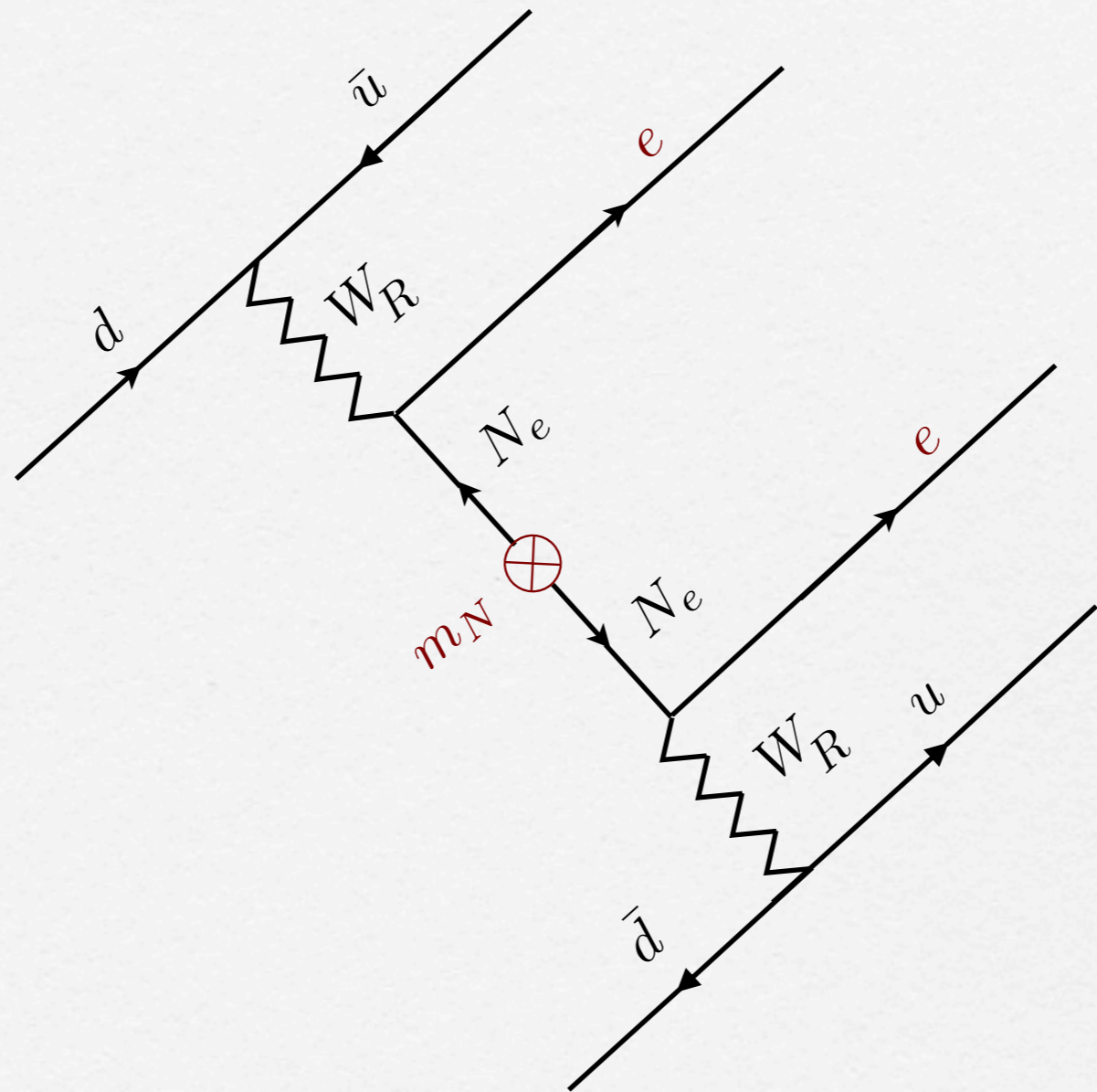
$$\frac{RR}{LL} \simeq O(1)$$

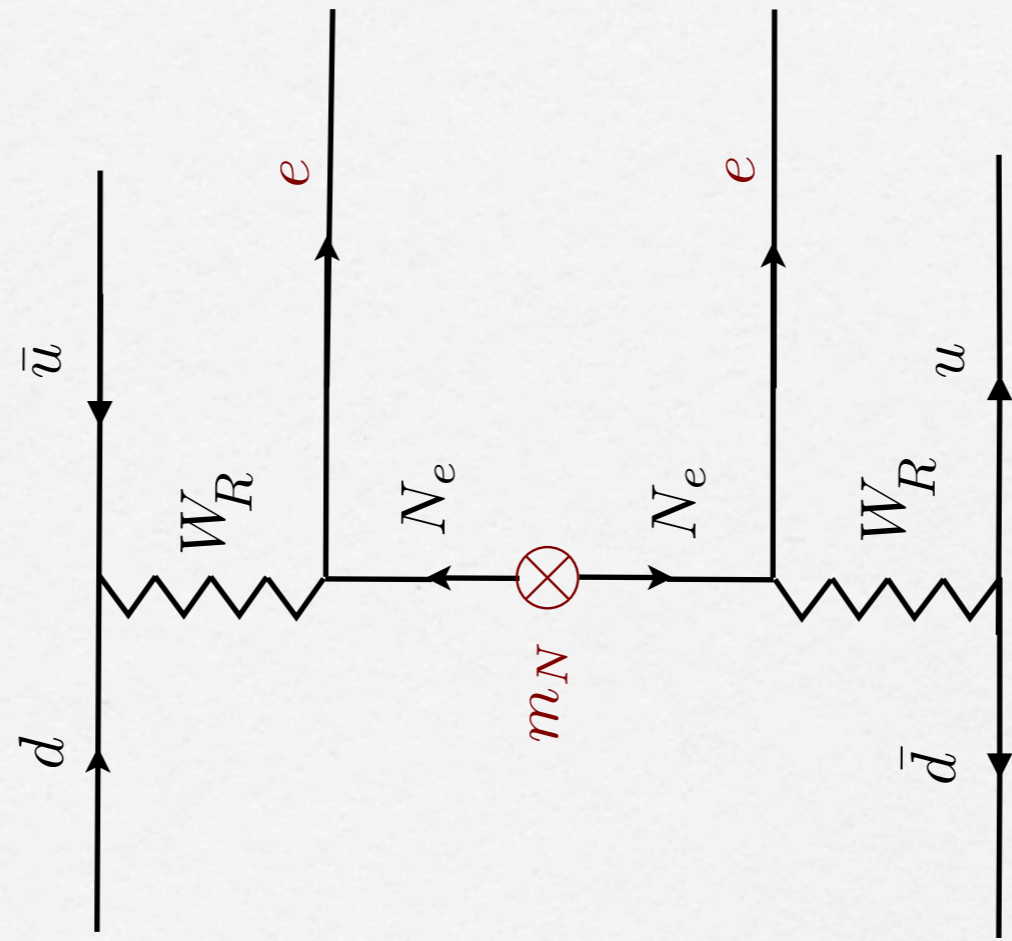
$$M_{W_R} \simeq m_N \simeq 10 M_{W_L}$$

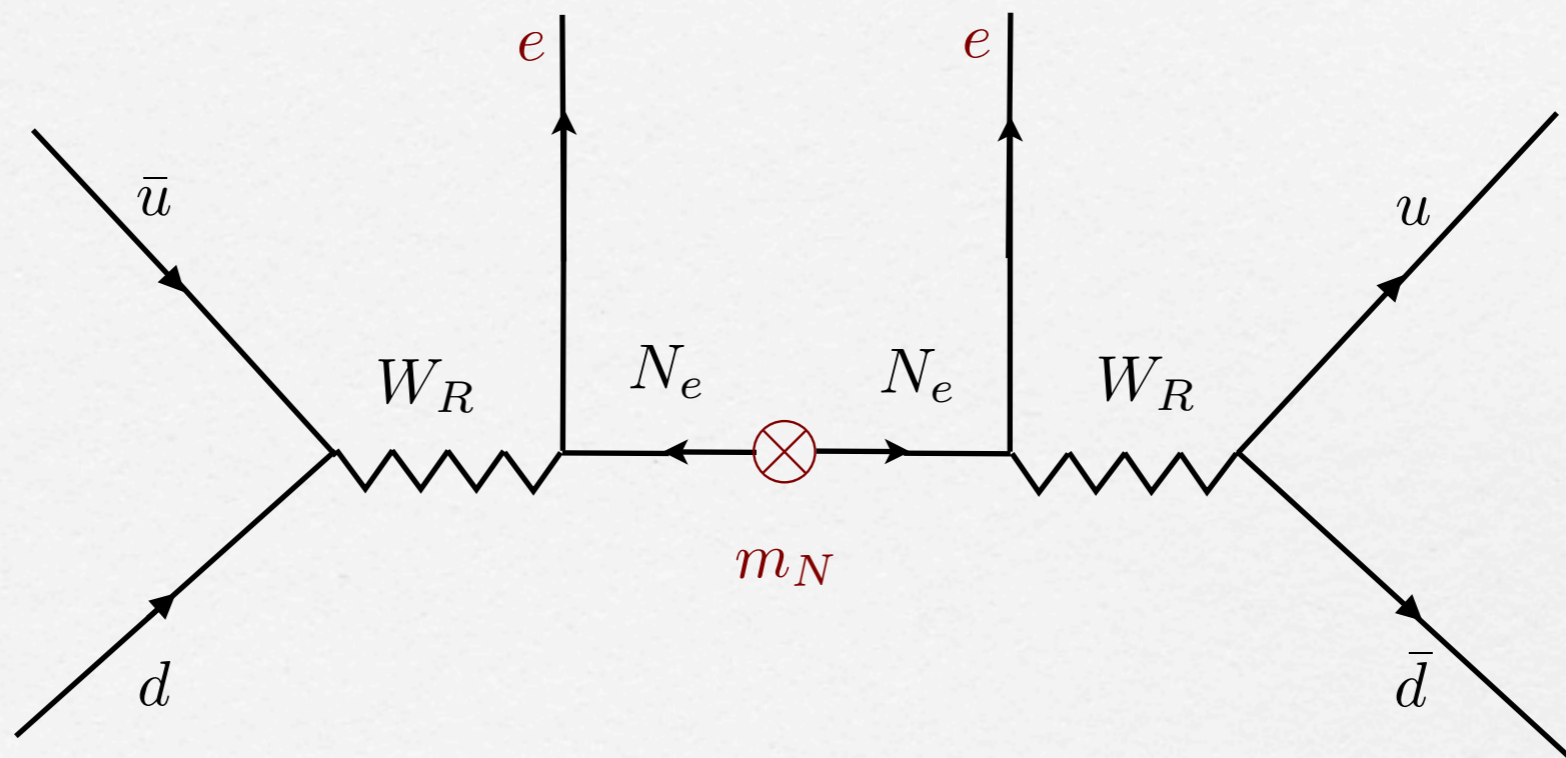
$$m_\nu \simeq 1 \text{ eV}$$

$$RR \propto \frac{1}{M_{W_R}^4} \frac{1}{m_N}$$

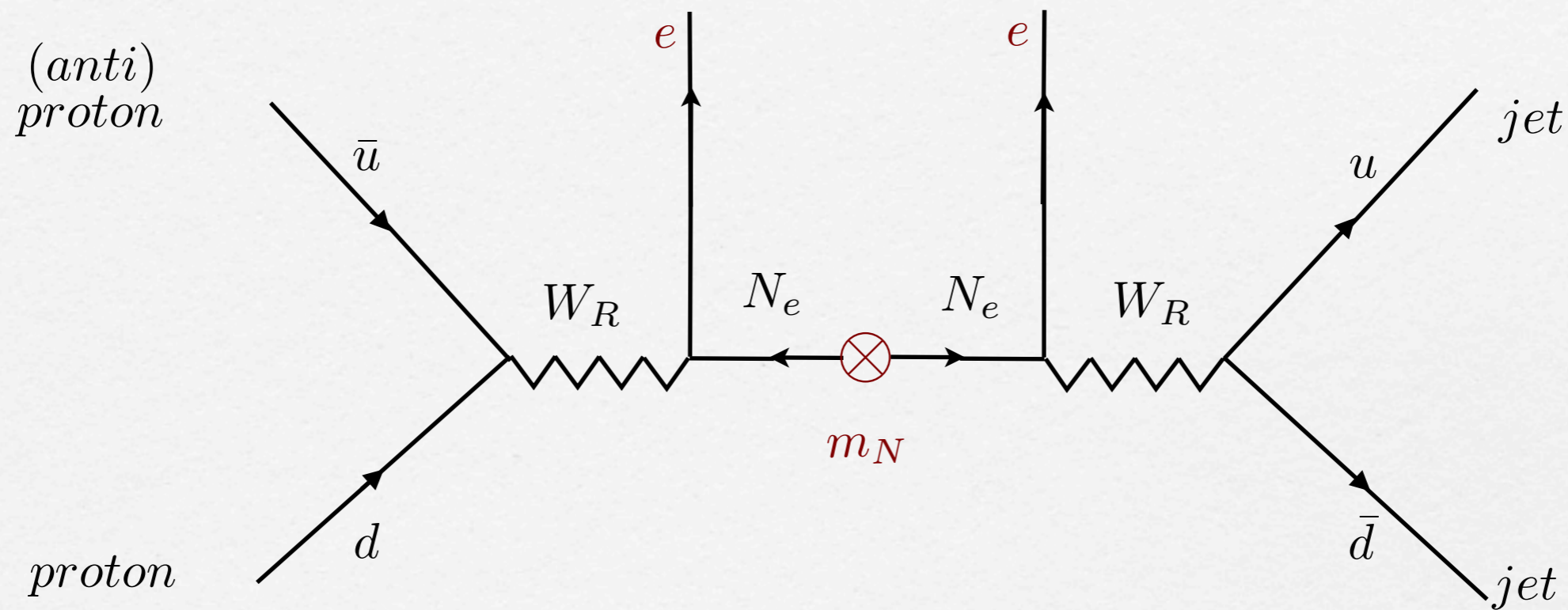






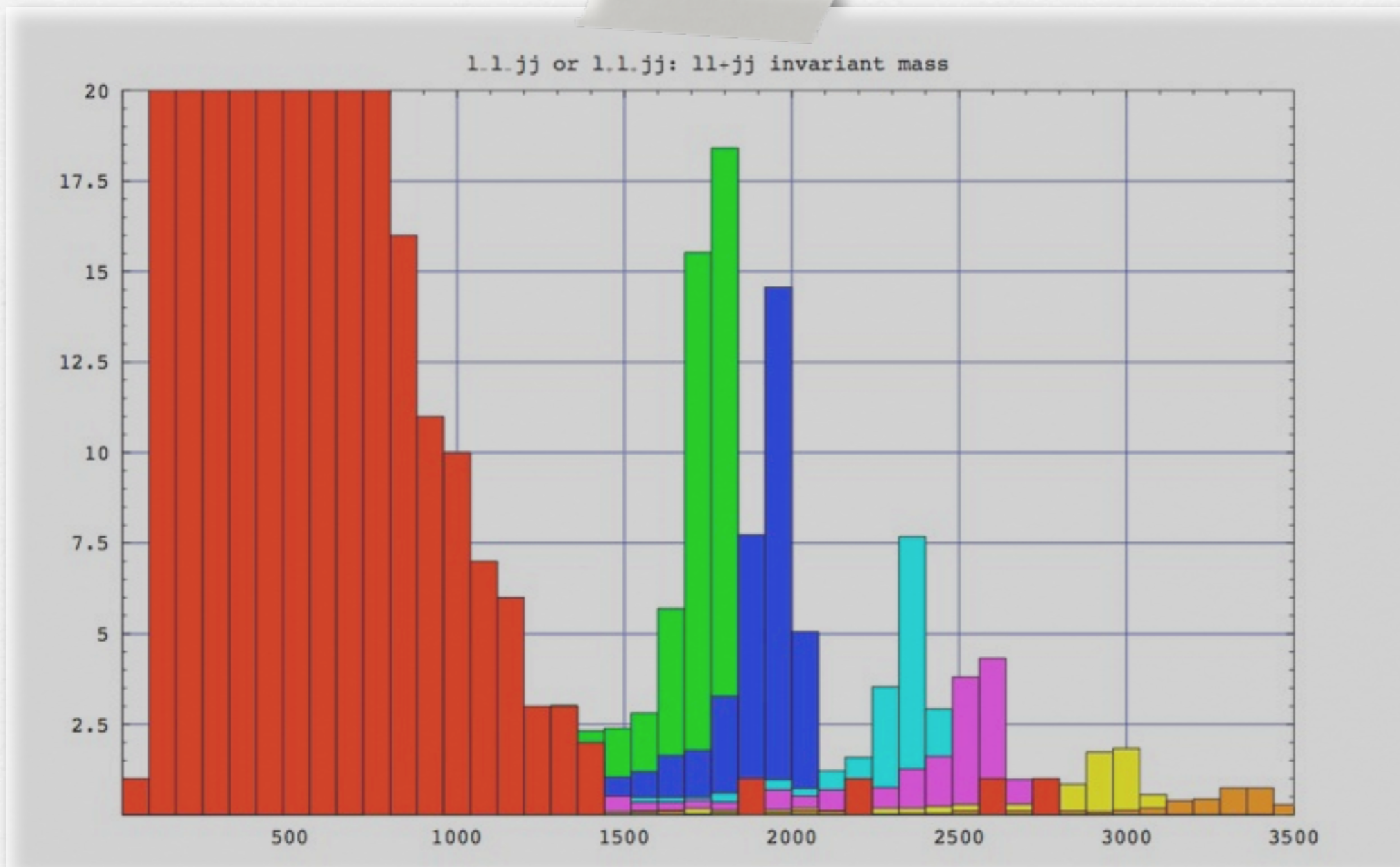


W_R production @ colliders



- Parity restoration
- same-sign dileptons + jets

Keung, G.S. '83



of events as a function of energy (GeV) for $L = 8\text{fb}^{-1}$

M_R (TeV): 1.8; 2, 0; 2.4; 2.6; 3, 0; 3.4

talk by Nesti

Spectacular signatures

BUT

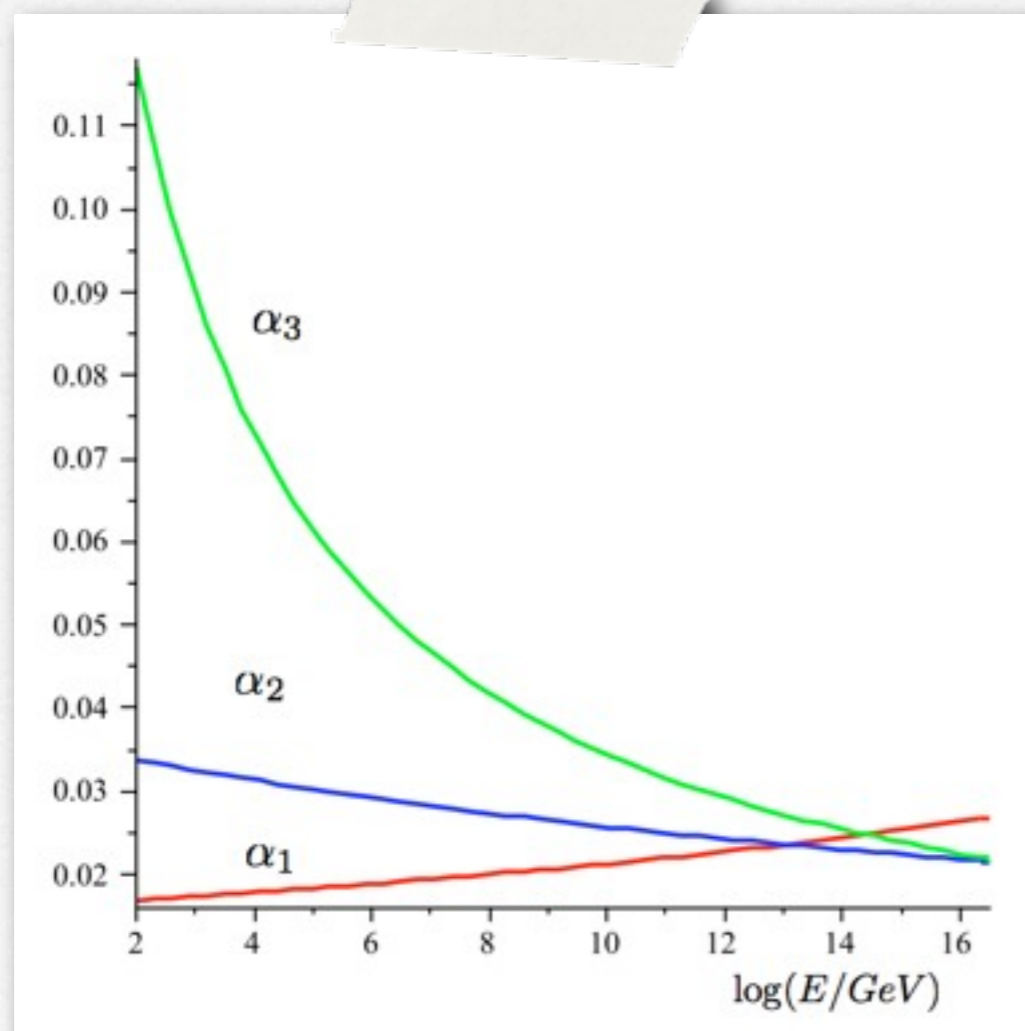
Scale not predicted

Typical of most
Beyond Standard Model physics

Example of a predictive theory

SU(5) grand unified theory: minimal

Georgi, Glashow '74



- massless neutrinos*
- no unification*

talk by Doršner

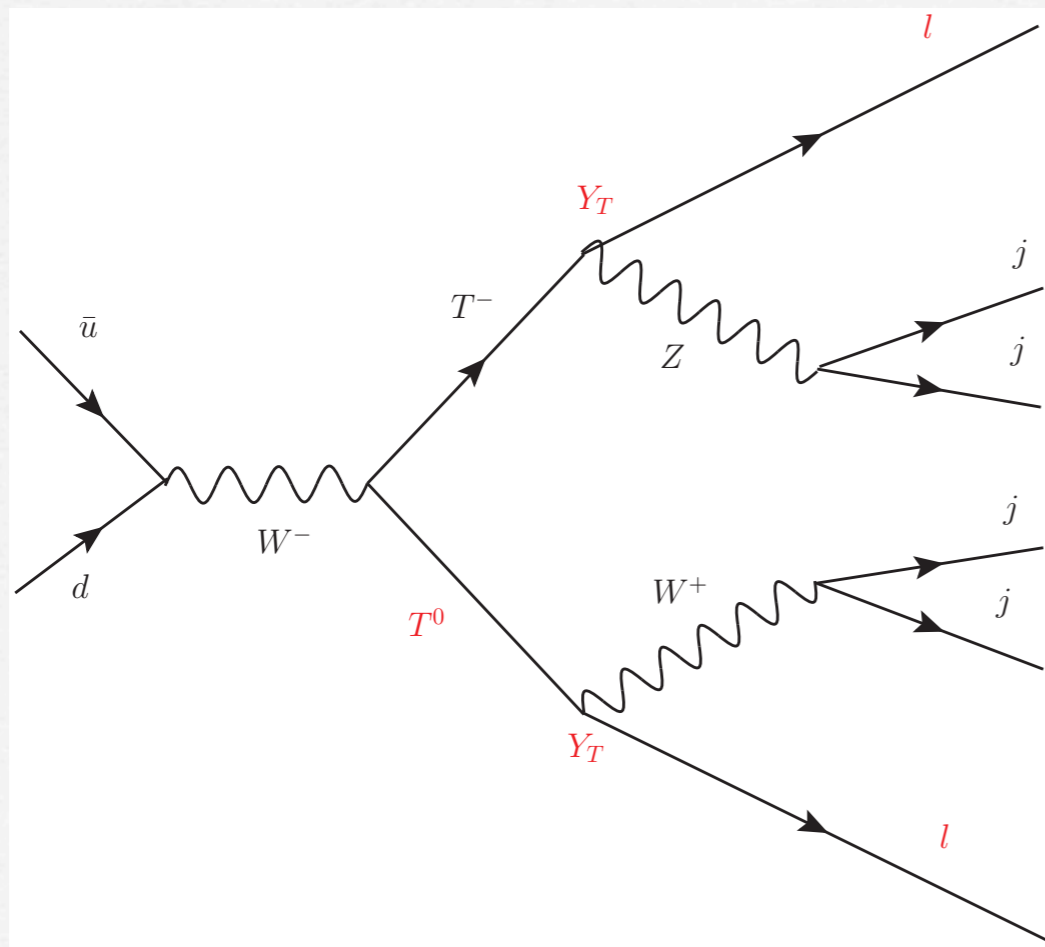
Simple predictive extension: add 24_F Bajc, GS '06

- one ν_R
- one fermion triplet - light



one massless neutrino

(T^+, T^0, T^-) : weak triplet



$$m_T < TeV$$

Bajc, Nemevšek, GS '07

@ LHC:

triplet decays through Yukawas



probe of neutrino masses and mixings

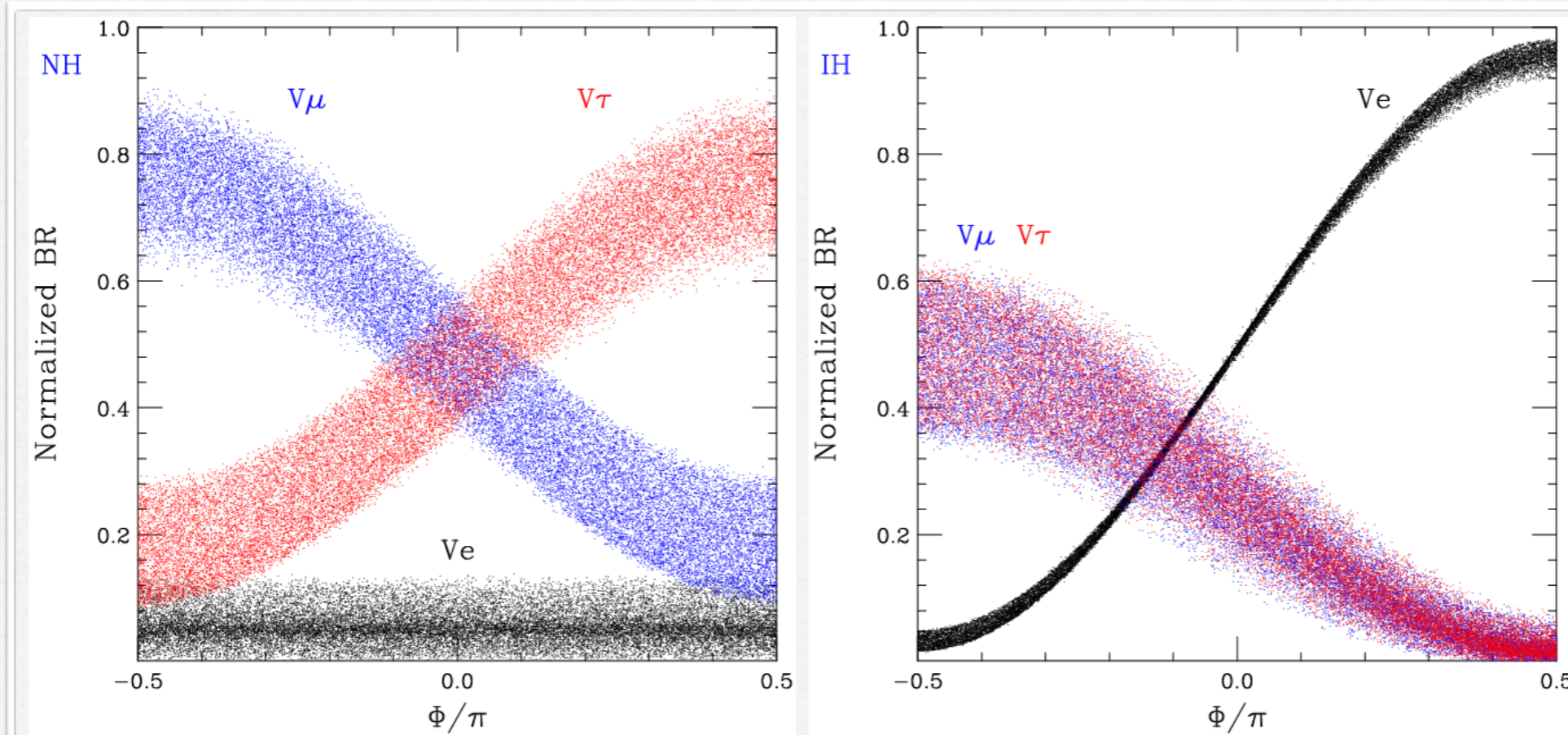


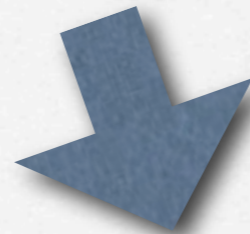
FIG. 10: Normalized branching versus Majorana phase for NH (left) and IH (right). $\text{Im}(z) \geq 2$.

Arhrib, Bajc, Ghosh, Han, Huang,
Puljak, G.S. '09

Message

LHC:

Higgs Hunter Machine



can probe the origin of
neutrino mass

Still, are we seeing dinosaurs ?

lots of important physics within SM - as discussed amply at LHC days

- how many generations?
- limit on Higgs mass?
- Higgs structure: how many doublets?

until a few years ago PDG claimed no new families from high precision: S , T , U

strong impact on the Higgs

high precision physics (& direct search):



upper limit on the Higgs mass

$$m_h \lesssim 180 \text{ GeV}$$

talk by Djouadi

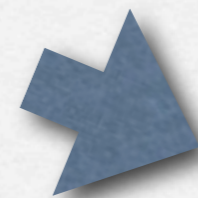
Fourth generation?

talk by Dorigo

$m_4 \lesssim 500 \text{ GeV} \Rightarrow$ perturbativity limit
700 GeV ?

Marciano, Valencia, Willenbrock '89

parameter set	m_{u_4}	m_{d_4}	m_H	ΔS_{tot}	ΔT_{tot}
(a)	310	260	115	0.15	0.19
(b)	320	260	200	0.19	0.20
(c)	330	260	300	0.21	0.22
(d)	400	350	115	0.15	0.19
(e)	400	340	200	0.19	0.20
(f)	400	325	300	0.21	0.25



LHC

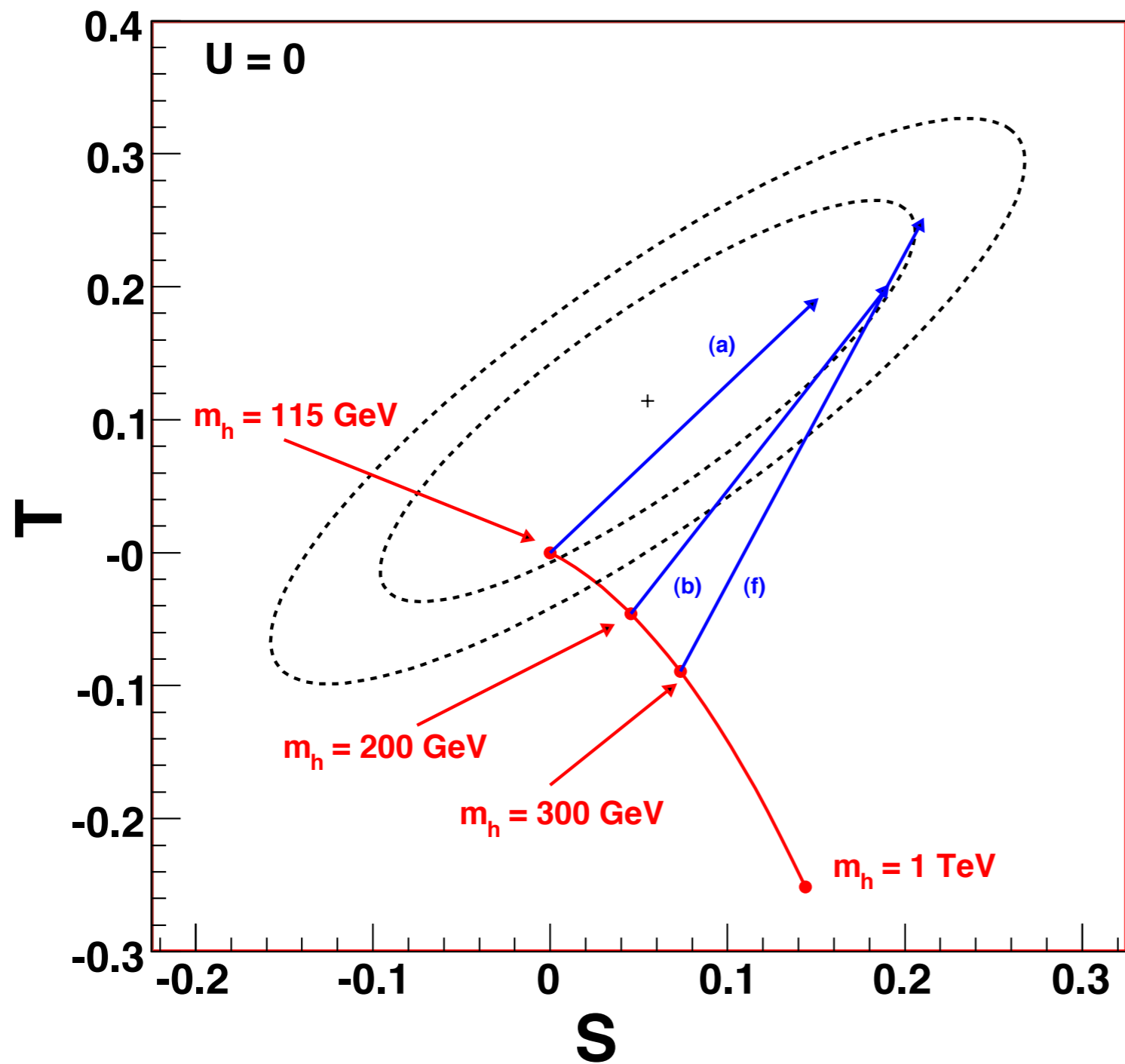
direct CDF limits:

$m_4 \geq 300 \text{ GeV} ? !$

Hung and Sher '07 ?

Flacco et al '10 !

Kribs, et al '07



Kribs, et al '07

heavier Higgs OK



315 (750) GeV
@68(95)% CL

changes:
production,
decay rates

more on 4th

- detailed study of masses and mixings

Eberhardt, Lenz, Rohrwild '10

- CP

Eilam, Melic, Trampetic '09

- hierarchy problem

Hung, Xiong '10

new limits

gluon fusion
strongly affected



heavy quark loops

recent Tevatron

CDF & D0 '10

excludes the region:

$$131 \text{ GeV} \lesssim m_h \lesssim 204 \text{ GeV}$$

@95% CL

MSSM?

Dawson, Jaiswal '10

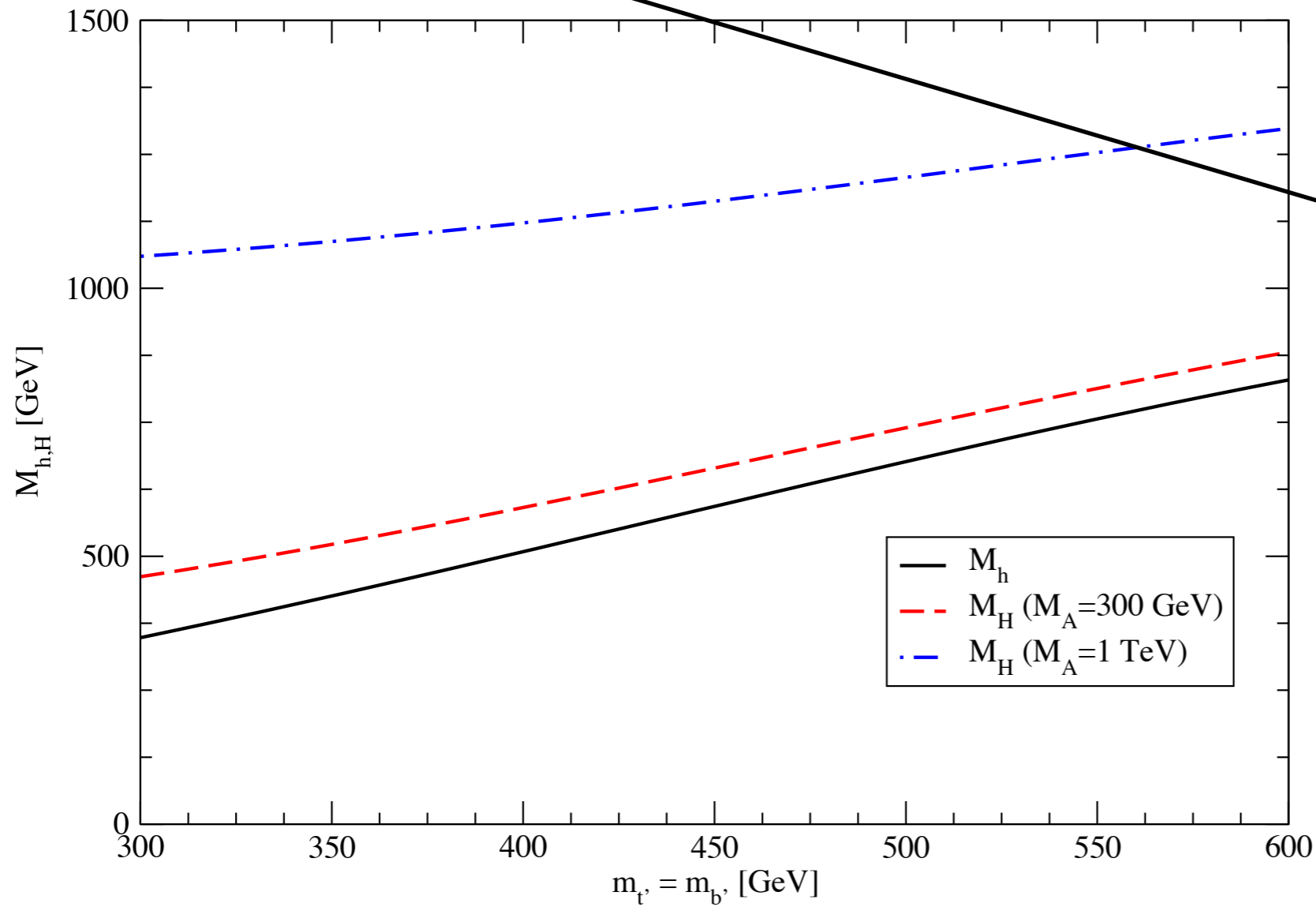
- the mass of $d_4 \Rightarrow \tan\beta \simeq 1$ (perturbativity)
- Higgs mass limit dramatically changed

heavy quark

$$m_{h^0}^2 \leq M_Z^2 + \frac{3\alpha}{4\pi} \frac{m_t^4}{M_W^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + 3 \right]$$

Dawson, Jaishwal '10

$\tan \beta = 1$
 $m_e = 250 \text{ GeV}, m_\nu = 230 \text{ GeV}$



Carpenter,
Rajaraman,
Whiteson,
Thursday

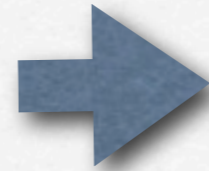
$e_4 \rightarrow 250 \text{ GeV}$

LHC@7 TeV

1 ftp^{-1}

(current around
100 GeV)

Mirror families?



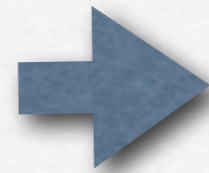
restore parity

Lee, Yang '56

$$\begin{pmatrix} \nu \\ e \end{pmatrix}_L \quad e_R$$

$$\begin{pmatrix} u \\ d \end{pmatrix}_L \quad u_R \quad d_R$$

Mirror families?



restore parity

Lee, Yang '56

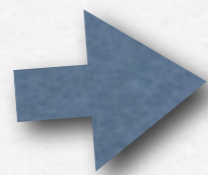
$$\begin{pmatrix} \nu \\ e \end{pmatrix}_L \quad e_R$$

$$\begin{pmatrix} u \\ d \end{pmatrix}_L \quad u_R \quad d_R$$

$$E_L \quad \begin{pmatrix} N \\ E \end{pmatrix}_R$$

$$D_L \quad U_L \quad \begin{pmatrix} U \\ D \end{pmatrix}_R$$

- gauge B, L - required to cancel anomalies
- Kaluza-Klein theories
- $N=2$ supersymmetry
- $SO(10+2N)$ unification



perfectly OK - if two Higgs

He, Polonsky, Su '01
update?

Message

experimentalist's task:

hard to see a black cat in
a dark room,

...especially if it is not there



your turn

physics not only:
why we are here,
why three dimensions,
why are numbers what they are ...