

Split - 4 October '10

*LHC Physics  
in the  
SM and beyond*

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The LHC physics run is going on, .... finally!

3.5 TeV per beam. Goal  $1\text{fb}^{-1}$  by the end of 2011

Top physics priorities at the LHC (ATLAS&CMS):

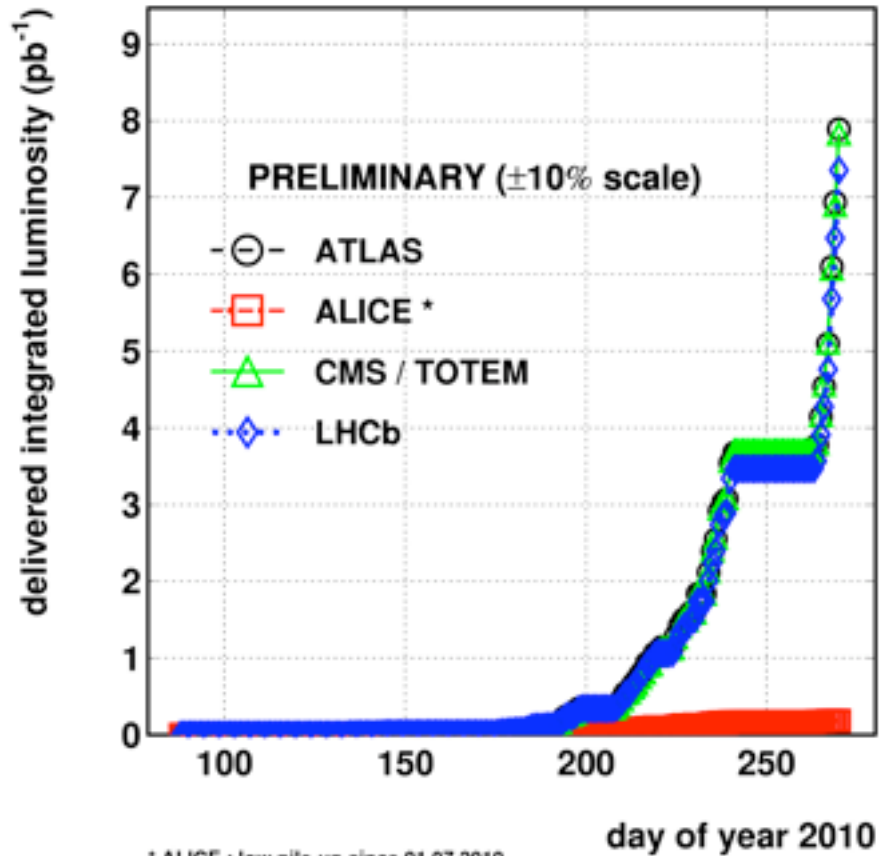
- Clarify the EW symmetry breaking sector
- Search for new physics at the TeV scale
- Identify the particle(s) that make the Dark Matter in the Universe

Also:

- LHCb: precision B physics (CKM matrix and CP violation)
- ALICE: Heavy ion collisions & QCD phase diagram

⊕ At this point, fresh input from experiment is badly needed!!

### LHC 2010 RUN (3.5 TeV/beam)



\* ALICE : low pile-up since 01.07.2010

Total integrated  
luminosity so far  
~ 8 pb<sup>-1</sup>



# Saturday 2<sup>nd</sup> October

02-Oct-2010 05:26:55					Fill #: 1387					Energy: 3500 GeV					I(B1): 1.41e+13					I(B2): 1.43e+13				
Experiment Status		ATLAS PHYSICS			ALICE PHYSICS			CMS PHYSICS			LHCb PHYSICS													
Instantaneous Lumi (ub.s) <sup>-1</sup>		38.149			0.135			39.007			37.199													
BRAN Luminosity (ub.s) <sup>-1</sup>		33.126			0.133			39.134			36.521													
Fill Luminosity (nb) <sup>-1</sup>		36.3			--			31.2			25.1													
BKGD 1		0.045			0.015			8.343			0.291													
BKGD 2		214.000			0.369			0.002			0.002													
BKGD 3		19.000			0.005			0.000			0.217													
LHCb VELO Position		IN			Gap: 0.0 mm			STABLE BEAMS			TOTEM:			STANDBY										

peak luminosity  $\sim 4 \cdot 10^{31} \text{ cm}^{-2}\text{s}^{-1}$   
152 bunches  
 $1.4 \cdot 10^{13}$  protons in total in the beam



## Particle physics at a glance

The SM is a low energy effective theory  
(nobody can believe it is the ultimate theory)

It happens to be renormalizable, hence highly predictive.  
And is well supported by the data.

However, we expect corrections from higher energies

not only from the GUT or Planck scales  
but also from the TeV scale (LHC!)

In fact even just as a low energy effective theory  
the SM it is not satisfactory

⊕ QCD + the gauge part of the EW theory are fine,  
but the Higgs sector is so far only a conjecture

# The Higgs problem is central in particle physics today

A review: G.A. ArXiv:1003.3180

The main problems of the SM show up in the Higgs sector

$$V_{Higgs} = V_0 - \mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 + [\bar{\psi}_{Li} Y_{ij} \psi_{Rj} \phi + h.c.]$$

Vacuum energy  
 $V_{0\text{exp}} \sim (2 \cdot 10^{-3} \text{ eV})^4$

Possible instability  
depending on  $m_H$

Origin of quadratic  
divergences.  
Hierarchy problem

The flavour problem:  
large unexplained ratios  
of  $Y_{ij}$  Yukawa constants



The Standard EW theory:  $\mathcal{L} = \mathcal{L}_{\text{symm}} + \mathcal{L}_{\text{Higgs}}$

$$\mathcal{L}_{\text{symm}} = -\frac{1}{4}[\partial_\mu W_\nu^A - \partial_\nu W_\mu^A - ig\epsilon_{ABC}W_\mu^AW_\nu^B]^2 +$$

$$-\frac{1}{4}[\partial_\mu B_\nu - \partial_\nu B_\mu]^2 +$$

$$+\bar{\psi}\gamma^\mu[i\partial_\mu + gW_\mu^At^A + g'B_\mu\frac{Y}{2}]\psi$$

$$\mathcal{L}_{\text{Higgs}} = |[\partial_\mu - igW_\mu^At^A - ig'B_\mu\frac{Y}{2}]\phi|^2 +$$

$$+ V[\phi^\dagger\phi] + \bar{\psi}\Gamma\psi\phi + \text{h.c}$$

with  $V[\phi^\dagger\phi] = \mu^2(\phi^\dagger\phi)^2 + \lambda(\phi^\dagger\phi)^4$

$\mathcal{L}_{\text{symm}}$ : well tested (LEP, SLC, Tevatron...),  $\mathcal{L}_{\text{Higgs}}$ : ~ untested

**All** we know from experiment about the SM Higgs:

No Higgs seen at LEP2  $\rightarrow m_H > 114.4$  GeV (95%cl) 

Rad. corr's  $\rightarrow m_H < 186$  GeV (95%cl, incl. direct search bound)

$v = \langle\phi\rangle = \sim 174$  GeV ;  $m_W = m_Z \cos\theta_W$   doublet Higgs

That some sort of spontaneous symmetry breaking mechanism is at work has already been established (couplings symmetric, spectrum totally non symmetric)

The question is on the nature of the Higgs mechanism/particle(s)

- One doublet, more doublets, additional singlets?
- SM Higgs or SUSY Higgses
- Fundamental or composite (of fermions, of WW....)
- Pseudo-Goldstone boson of an enlarged symmetry
- A manifestation of extra dimensions (fifth comp. of a gauge boson, an effect of orbifolding or of boundary conditions....)
- ⊕ • Some combination of the above



## Can we do without the Higgs?

Suppose we take the gauge symmetric part of the SM and put masses by hand.

Gauge invariance is broken explicitly. The theory is no more renormalizable. One loses understanding of the observed accurate validity of gauge predictions for couplings.

Still, what is the fatal problem at the LHC scale?

The most immediate disease that needs a solution is the occurrence of unitarity violations in some amplitudes

To avoid this either there is one or more Higgs particles or some new states (e.g. new vector bosons)

Thus something must happen at the few TeV scale!!

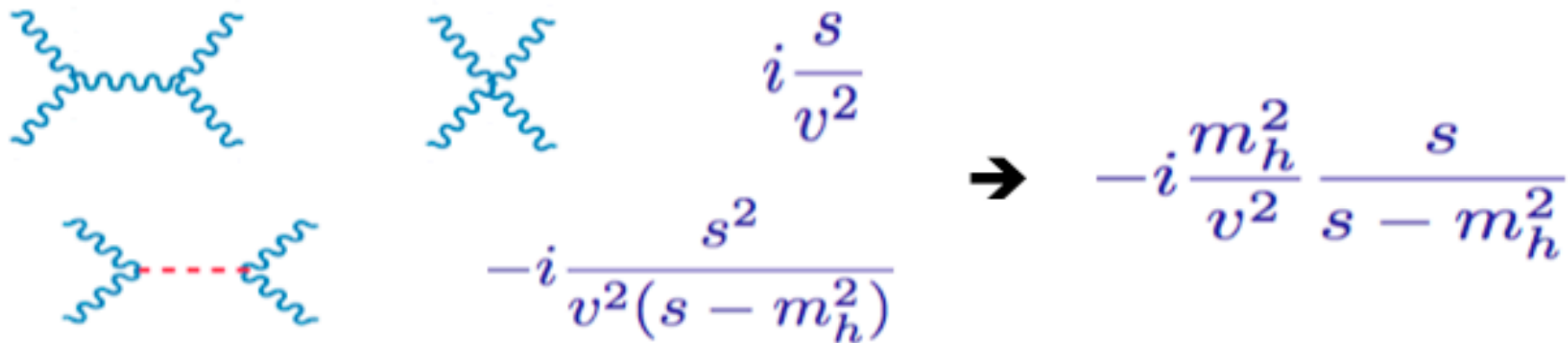


With no Higgs unitarity violations for  $E_{\text{CM}} \sim 1\text{-}3 \text{ TeV}$

**Unitarity** implies that scattering amplitudes cannot grow indefinitely with the centre-of-mass energy  $s$

In the SM, the Higgs particle is essential in ensuring that the scattering amplitudes with longitudinal weak bosons ( $W_L, Z_L$ ) satisfy (tree-level) unitarity constraints  
 [Veltman, 1977; Lee-Quigg-Thacker, 1977; ...] Zwirner

An example:  $\mathcal{A}(W_L^+ W_L^- \rightarrow Z_L Z_L) \quad (s \gg m_W^2)$



If no Higgs then something must happen!

## Can we do without the Higgs?

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## A crucial question for the LHC

### What saves unitarity?

- the Higgs
- some new vector boson
  - $W', Z'$
  - KK recurrences
  - resonances from a strong sector
  - .....



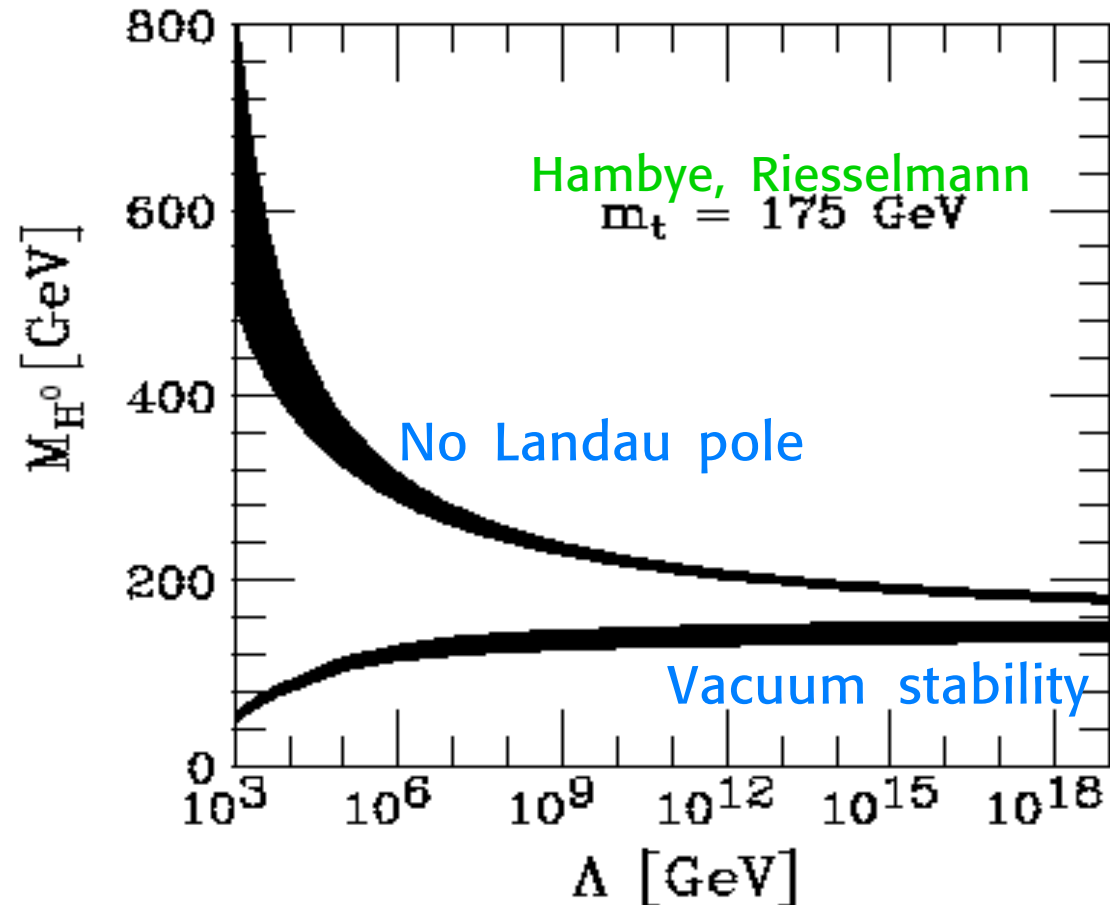
# Theoretical bounds on the SM Higgs mass

$\Lambda$ : scale of new physics beyond the SM

Upper limit: No Landau pole up to  $\Lambda$

Lower limit: Vacuum (meta)stability

The LHC was designed to cover the whole range



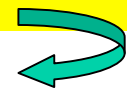
If the SM would be valid up to  $M_{\text{GUT}}$ ,  $M_{\text{Pl}}$  then  $m_H$  would be limited in a small range



Lower now because of  $m_t$



$128 \text{ GeV} < m_H < 180 \text{ GeV}$



# Status of the SM Higgs fit

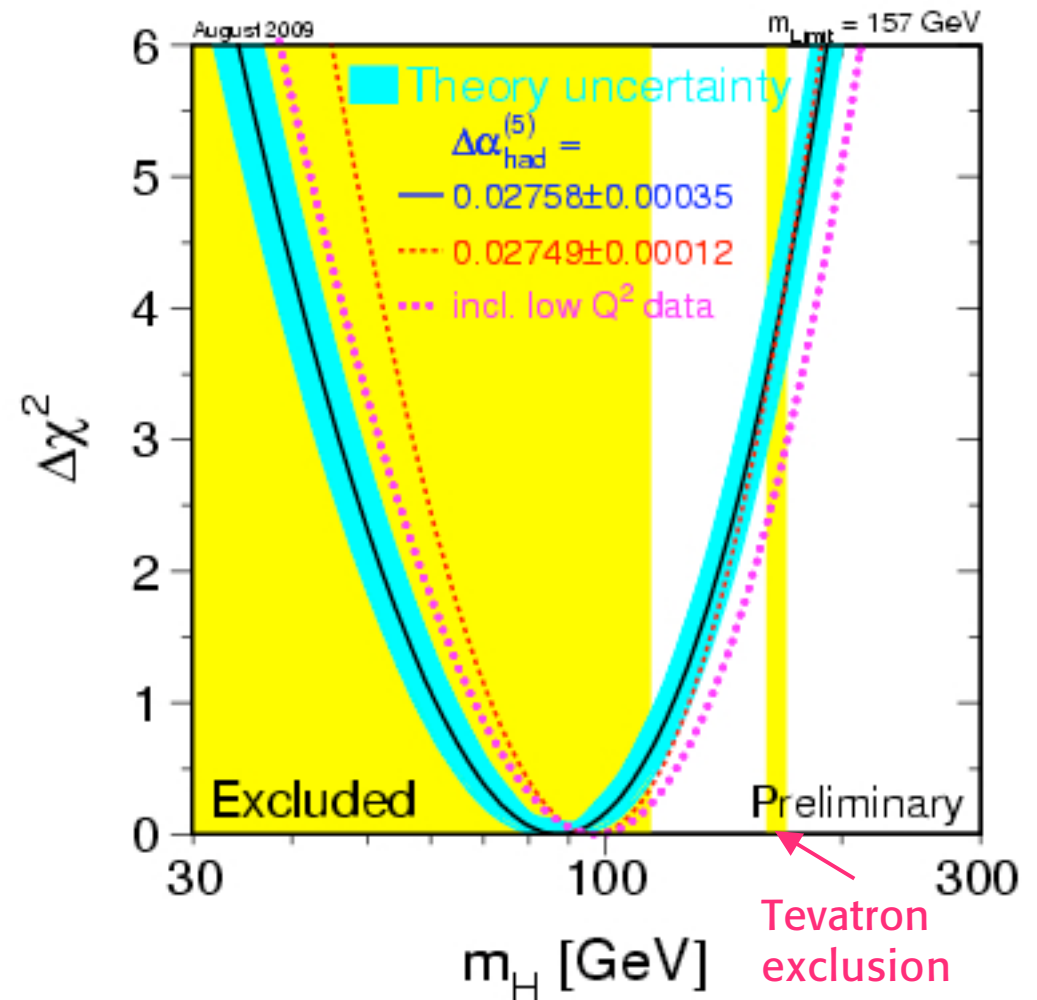
Summer '09

Rad Corr.s -> Sensitive to  $\log m_H$   
 $\log_{10} m_H(\text{GeV}) = 1.94 \pm 0.15$   
 $m_H = 87^{+35}_{-26} \text{ GeV}$

This is a great triumph for the SM: ~right in the narrow allowed range  $\log_{10} m_H \sim 2 - 3$

Direct search:  $m_H > 114.4 \text{ GeV}$

Radiative corr's indicate a light H



At 95 % cl

$m_H < 157 \text{ GeV}$  (rad corr.'s)

$m_H < 186 \text{ GeV}$  (incl. direct search bound)



# Is it possible that the Higgs is not found at the LHC?

Here "Higgs" means the "the EW symmetry breaking mechanism"

Looks pretty unlikely!!

The LHC discovery range is large enough:  $m_H < \sim 1 \text{ TeV}$   
the Higgs should be really heavy!

Rad. corr's indicate a light Higgs (whatever its nature)

A heavy Higgs would make perturbation theory to collapse nearby (violations of unitarity for  $m_H > \sim \text{TeV}$ )

Such nearby collapse of pert. th. is very difficult to reconcile with EW precision tests **plus** simulating a light Higgs

The SM good agreement with the data favours forms of new physics that keep at least some Higgs light

# The Standard Model works very well

So, why not find the Higgs and declare particle physics solved?

First, you have to find it!

→ LHC

Because of both:

## Conceptual problems

- Quantum gravity
- The hierarchy problem
- The flavour puzzle

.....

and experimental clues:

- Neutrino masses
- Coupling unification
- Dark matter
- Baryogenesis
- Vacuum energy

.....

Some of these problems point at new physics at the weak scale: eg Hierarchy  
Dark matter (perhaps)





## Dark Matter

WMAP, SDSS,  
2dFGRS.....

Most of the Universe is not made up of atoms:  $\Omega_{\text{tot}} \sim 1$ ,  $\Omega_{\text{b}} \sim 0.045$ ,  $\Omega_{\text{m}} \sim 0.27$

Most is Dark Matter and Dark Energy

Most Dark Matter is Cold (non relativistic at freeze out)  
Significant Hot Dark matter is disfavoured

Neutrinos are not much cosmo-relevant:  $\Omega_{\nu} < 0.015$

SUSY has excellent DM candidates: eg Neutralinos ( $\rightarrow$  LHC)  
Also Axions are still viable (introduced to solve strong CPV)  
(in a mass window around  $m \sim 10^{-4}$  eV and  $f_a \sim 10^{11}$  GeV  
but these values are simply a-posteriori)

Identification of Dark Matter is a task of enormous importance for particle physics and cosmology



LHC?



LHC has good chances because it can reach any kind of WIMP:

WIMP: Weakly Interacting Massive Particle  
with  $m \sim 10^1\text{-}10^3$  GeV

For WIMP's in thermal equilibrium after inflation the density is:

$$\Omega_\chi h^2 \simeq \text{const.} \cdot \frac{T_0^3}{M_{\text{Pl}}^3 \langle \sigma_{Av} \rangle} \simeq \frac{0.1 \text{ pb} \cdot c}{\langle \sigma_{Av} \rangle}$$

can work for typical weak cross-sections!!!

This "coincidence" is a good indication in favour of a WIMP explanation of Dark Matter



A crucial question for the LHC

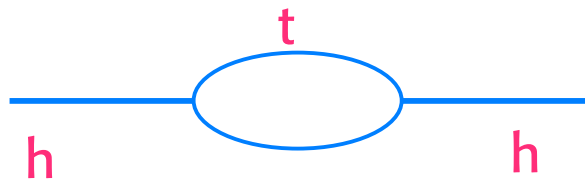
Is Dark Matter a WIMP?

LHC will tell yes or no to WIMPS



# The “little hierarchy” problem

e.g. the top loop (the most pressing):



$$m_h^2 = m_{\text{bare}}^2 + \delta m_h^2$$

$$\delta m_h^2|_{\text{top}} = -\frac{3G_F}{2\sqrt{2}\pi^2} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$$

This hierarchy problem demands new physics near the weak scale

$\Lambda$ : scale of new physics beyond the SM

- $\Lambda \gg m_Z$ : the SM is so good at LEP
- $\Lambda \sim \text{few times } G_F^{-1/2} \sim o(1\text{TeV})$  for a natural explanation of  $m_h$  or  $m_W$

Barbieri, Strumia

◀ **The LEP Paradox:**  $m_h$  light, new physics must be close but its effects were not visible at LEP2

⊕ **The B-factory Paradox:** and not visible in flavour physics

$$\Lambda \sim o(1\text{TeV})$$

## A crucial question for the LHC

What damps the top loop  $\Lambda^2$  dependence?

- the s-top (SUSY)
- some new fermion
  - t' (Little Higgs)
  - KK recurrences of the top (Extra dim.)
  - .....



# Solutions to the hierarchy problem

- Supersymmetry: boson-fermion symm.  
exact (**unrealistic**): cancellation of  $\Lambda^2$  in  $\delta m_h^2$   
approximate (**possible**):  $\Lambda \sim m_{\text{SUSY}} - m_{\text{ord}}$   $\longrightarrow$  top loop  
 $\Lambda \sim m_{\text{stop}}$   
The most widely accepted

- The Higgs is a  $\bar{\psi}\psi$  condensate. No fund. scalars. But needs new very strong binding force:  $\Lambda_{\text{new}} \sim 10^3 \Lambda_{\text{QCD}}$  (technicolor).  
Strongly disfavoured by LEP. Coming back in new forms

- Models where extra symmetries allow  $m_h$  only at 2 loops and non pert. regime starts at  $\Lambda \sim 10 \text{ TeV}$

"Little Higgs" models. Some extra trick needed to solve problems with EW precision tests

- Extra spacetime dim's that "bring"  $M_{\text{Pl}}$  down to  $o(1\text{TeV})$

Exciting. Many facets. Rich potentiality. No baseline model emerged so far

-  Ignore the problem: invoke the anthropic principle

# The anthropic route

The scale of the cosmological constant is a big mystery.

$$\Omega_\Lambda \sim 0.75 \quad \longrightarrow \quad \rho_\Lambda \sim (2 \cdot 10^{-3} \text{ eV})^4 \sim (0.1 \text{ mm})^{-4}$$

In Quantum Field Theory:  $\rho_\Lambda \sim (\Lambda_{\text{cutoff}})^4$  Similar to  $m_\nu$ !?

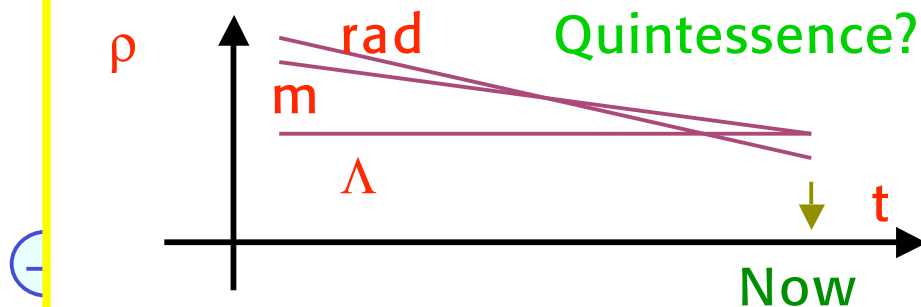
If  $\Lambda_{\text{cutoff}} \sim M_{\text{Pl}}$   $\longrightarrow$   $\rho_\Lambda \sim 10^{123} \rho_{\text{obs}}$

Exact SUSY would solve the problem:  $\rho_\Lambda = 0$

But SUSY is broken:  $\rho_\Lambda \sim (\Lambda_{\text{SUSY}})^4 \sim 10^{59} \rho_{\text{obs}}$

It is interesting that the correct order is  $(\rho_\Lambda)^{1/4} \sim (\Lambda_{\text{EW}})^2 / M_{\text{Pl}}$

Other problem:  
"Why now"?



"Quintessence"  
 $\Lambda$  as a vev of a field  $\phi$ ?

Coupled to gauge singlet matter, eg  $\nu_R$ , to solve magnitude and why now?

# Is naturalness relevant?

Speculative physics reasons to doubt:

- The empirical value of the cosmological constant  $\Lambda$  poses a tremendous, unsolved naturalness problem yet the value of  $\Lambda$  is close to the Weinberg upper bound for galaxy formation
- Possibly our Universe is just one of infinitely many continuously created from the vacuum by quantum fluctuations
- Different physics in different Universes according to the multitude of string theory solutions ( $\sim 10^{500}$ )

Perhaps we live in a very unlikely Universe but one that allows our existence





I find applying the anthropic principle to the SM hierarchy problem not appropriate

After all we can find plenty of models that reduce the fine tuning from  $10^{14}$  to  $10^2$ :  
so why make our Universe so terribly unlikely?

The case of the cosmological constant is a lot different: the context is not as fully specified as the for the SM (quantum gravity, string cosmology, branes in extra dims., wormholes thru different Universes....)



## SUSY: boson fermion symmetry

The hierarchy problem:  $\delta m_{h|top}^2 = -\frac{3G_F}{2\sqrt{2}\pi} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$

In broken SUSY  $\Lambda^2$  is replaced by  $(m_{stop}^2 - m_t^2) \log \Lambda$

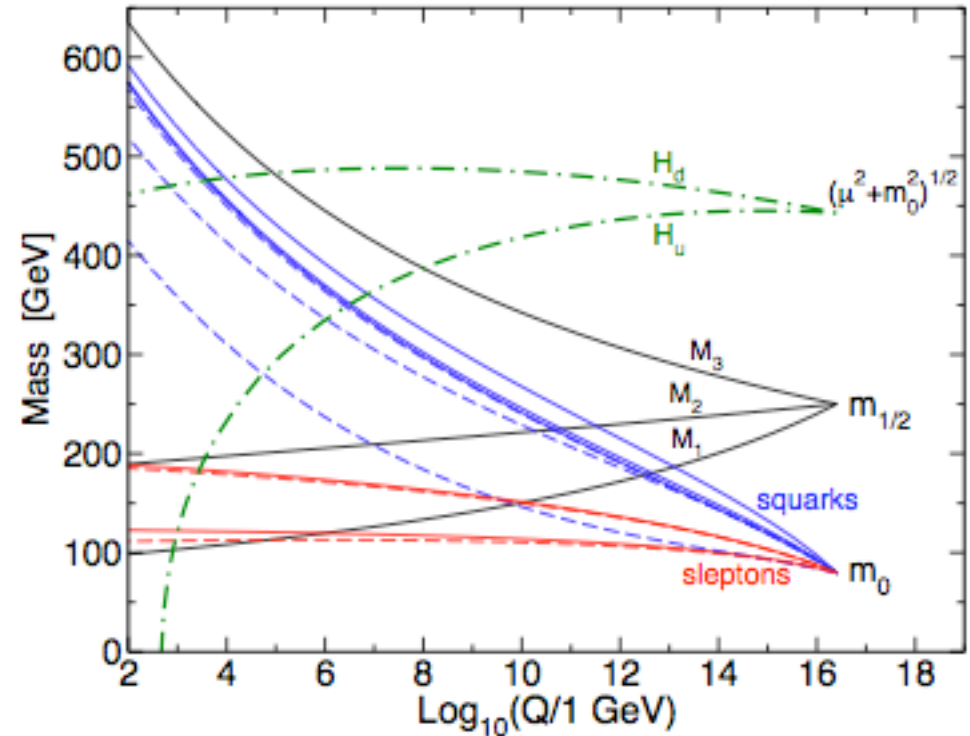
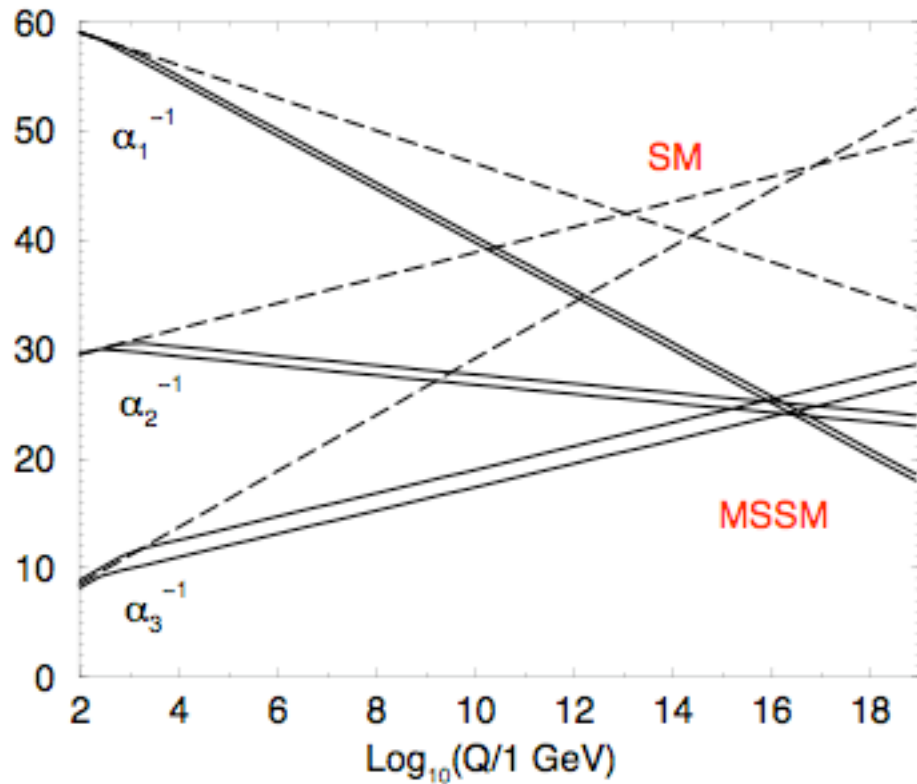
$m_H > 114.4$  GeV,  $m_{\chi_+} > 100$  GeV, EW precision tests, success of CKM, absence of FCNC, all together, impose sizable Fine Tuning (FT) particularly on **minimal** realizations (MSSM, CMSSM...).

Yet SUSY is a completely specified, consistent, computable model, perturbative up to  $M_{Pl}$  quantitatively in agreement with coupling unification (GUT's) **(unique among NP models)** and has a good DM candidate: the neutralino (actually more than one).



Remains the reference model for NP

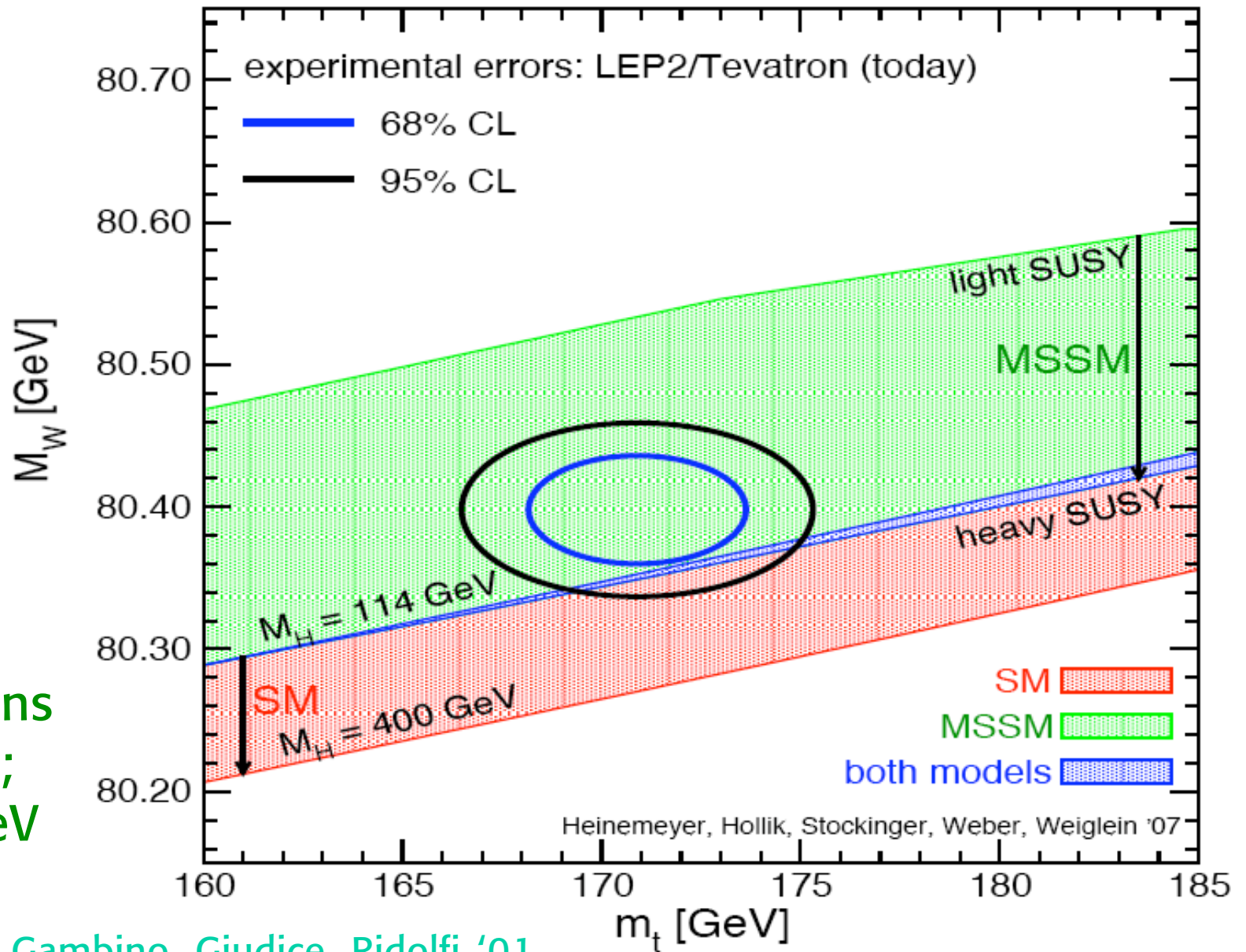
SUSY is unique in providing a weakly interacting theory up to the GUT/Planck scale



Other BSM models (little Higgs, composite Higgs, Higgsless....) all become strongly interacting at a multi-TeV scale



# SUSY effects could modify the SM fit



“light SUSY” =  
 = light s-leptons  
 and charginos;  
 s-quarks  $\sim 1 \text{ TeV}$

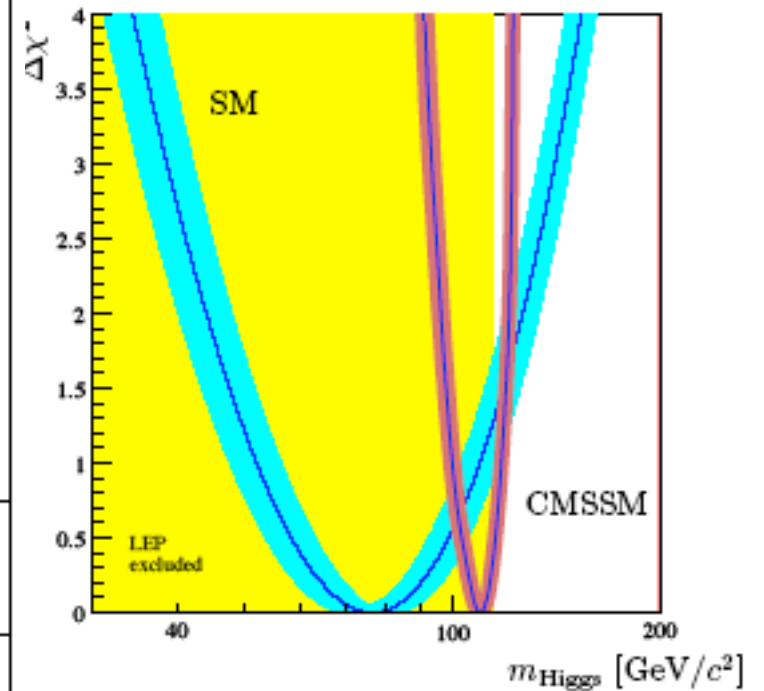
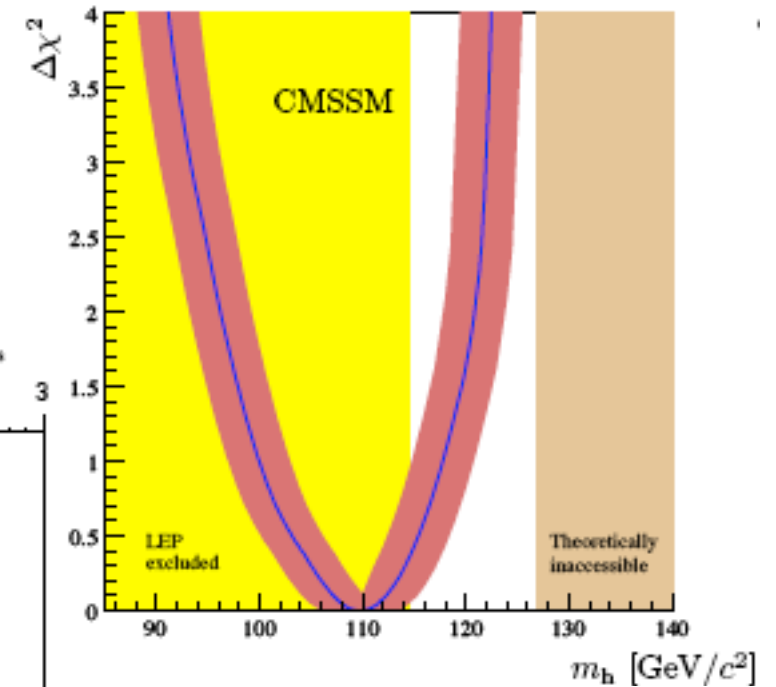
⊕ G.A. Caravaglios, Gambino, Giudice, Ridolfi '01

A recent study indicates that  $m_h$  goes up in CMSSM when  $b \rightarrow s\gamma$ ,  $a_\mu$ ,  $\Omega_{DM}$  are added

O. Buchmuller et al '07, '08 [0808.4128]

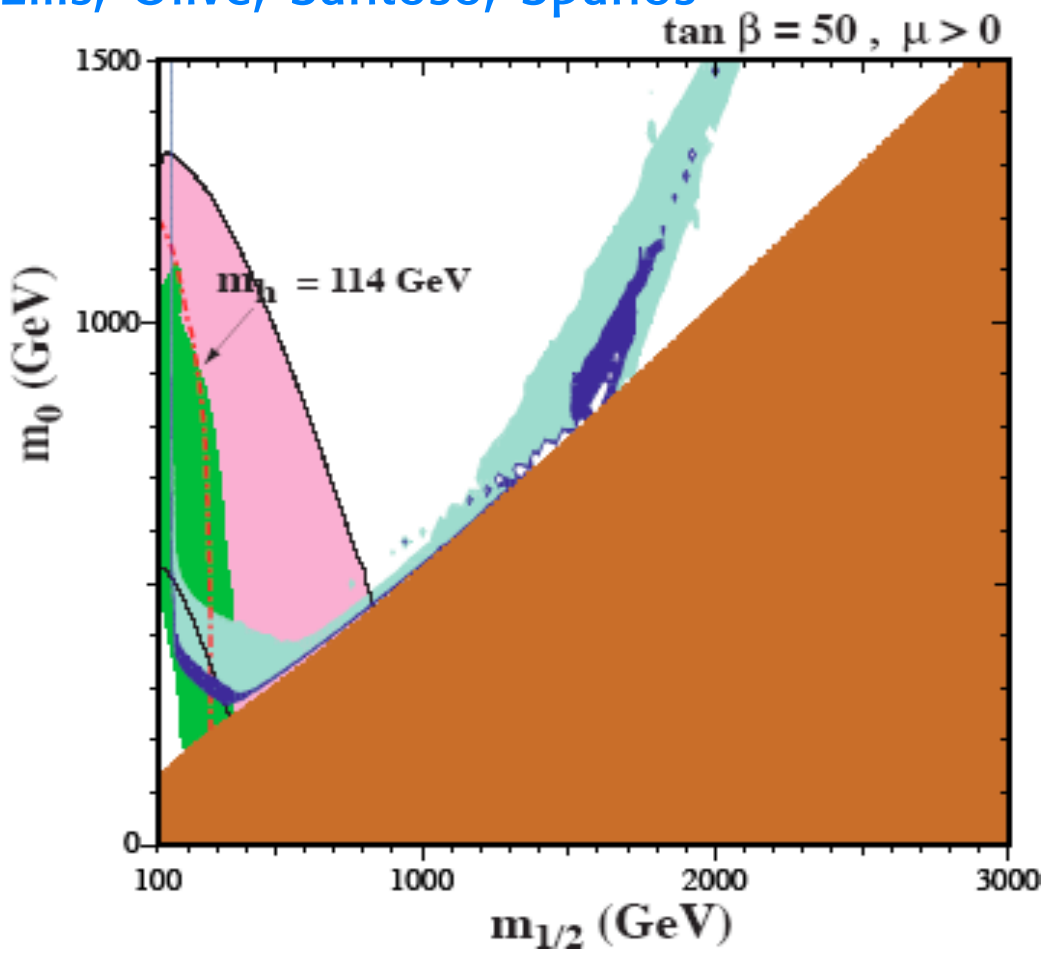
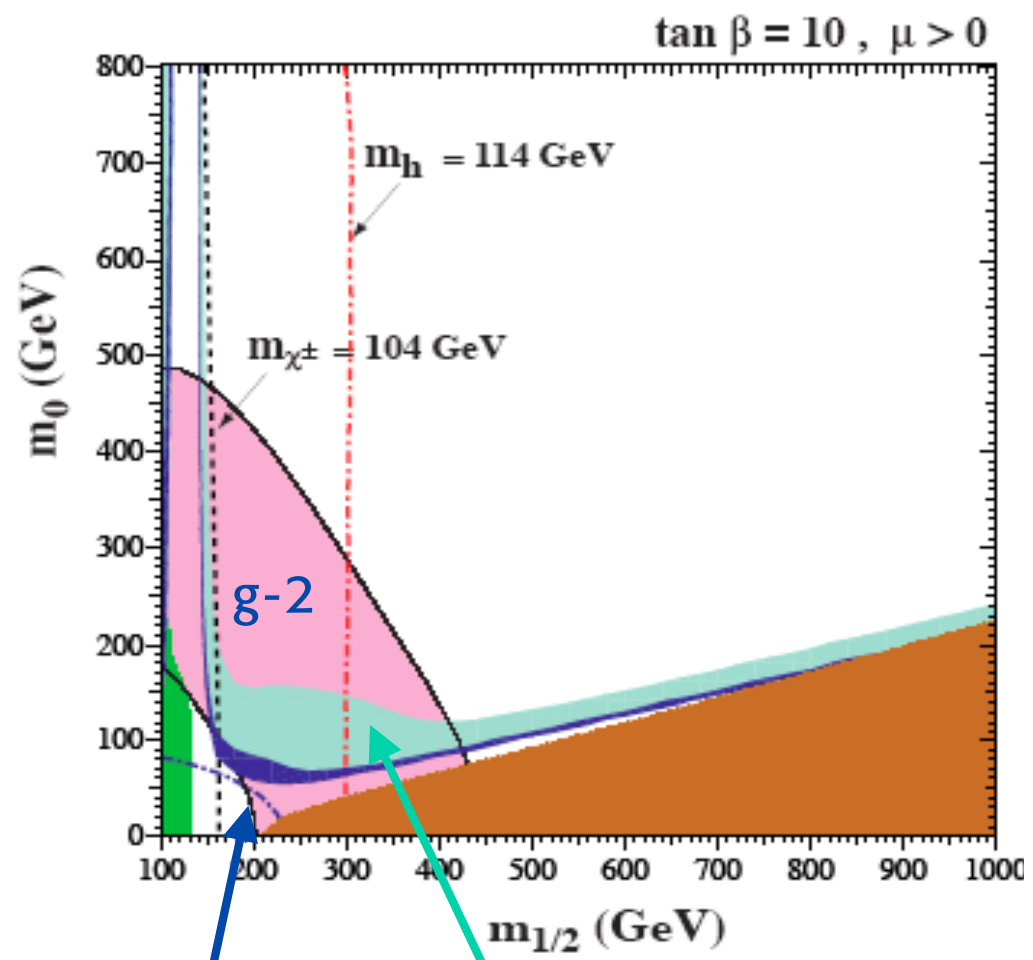
also:  
J. Ellis et al '07

CMSSM			$10^{(meas - O^{fit})/\sigma_{meas}}$			
Variable	Measurement	Fit	0	1	2	3
$\Delta\alpha_{had}^{(S)}(m_Z)$	$0.02758 \pm 0.00035$	0.02774	[Bar]			
$m_Z$ [GeV]	$91.1875 \pm 0.0021$	91.1873	[Bar]			
$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	2.4952	[Bar]			
$\sigma_{had}^0$ [nb]	$41.540 \pm 0.037$	41.486	[Bar]			
$R_1$	$20.767 \pm 0.025$	20.744	[Bar]			
$A_{fb}^{0,l}$	$0.01714 \pm 0.00095$	0.01641	[Bar]			
$A_1(P_\tau)$	$0.1465 \pm 0.0032$	0.1479	[Bar]			
$R_b$	$0.21629 \pm 0.00066$	0.21613	[Bar]			
$R_c$	$0.1721 \pm 0.0030$	0.1722	[Bar]			
$A_{fb}^{0,b}$	$0.0992 \pm 0.0016$	0.1037	[Bar]			
$A_{fb}^{0,c}$	$0.0707 \pm 0.0035$	0.0741	[Bar]			
$A_b$	$0.923 \pm 0.020$	0.935	[Bar]			
$A_c$	$0.670 \pm 0.027$	0.668	[Bar]			
$A_1(SLD)$	$0.1513 \pm 0.0021$	0.1479	[Bar]			
$\sin^2\theta_{eff}^{lep}(Q_b)$	$0.2324 \pm 0.0012$	0.2314	[Bar]			
$m_W$ [GeV]	$80.398 \pm 0.025$	80.382	[Bar]			
$m_t$ [GeV]	$170.9 \pm 1.8$	170.8	[Bar]			
$R(b \rightarrow s\gamma)$	$1.13 \pm 0.12$	1.12	[Bar]			
$B_s \rightarrow \mu\mu$ [ $\times 10^{-8}$ ]	$< 8.00$	0.33	N/A (upper limit)			
$\Delta a_\mu$ [ $\times 10^{-9}$ ]	$2.95 \pm 0.87$	2.95	[Bar]			
$\Omega h^2$	$0.113 \pm 0.009$	0.113	[Bar]			



SUSY Dark Matter: best candidate the neutralino  
 [in SUSY the gravitino is a non-WIMP alternative]

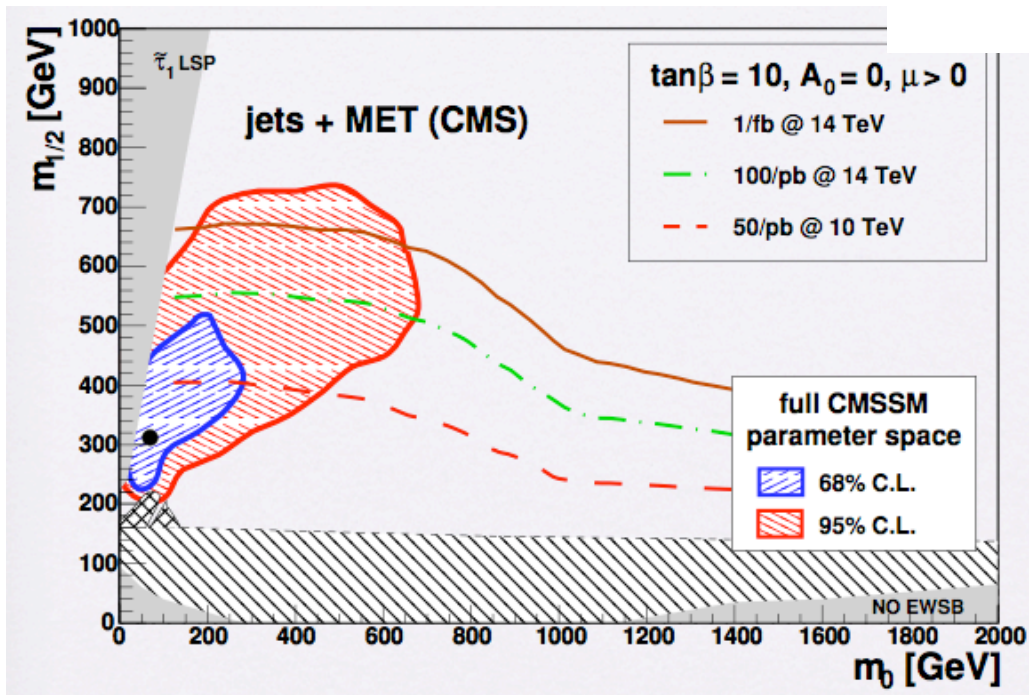
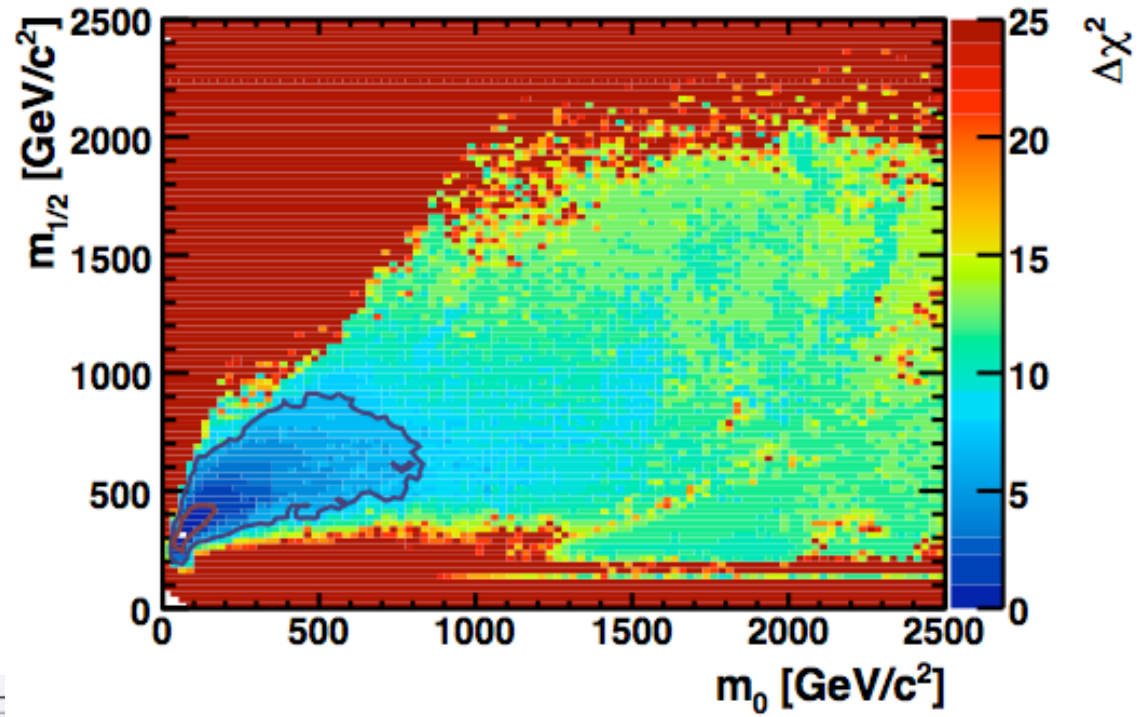
Ellis, Olive, Santoso, Spanos



This is for the CMSSM  
 With less constraints, more space



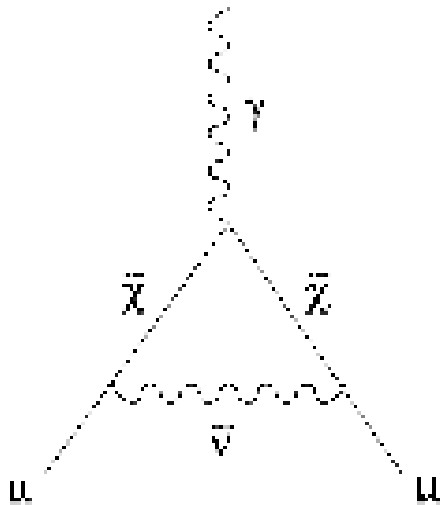
relatively light SUSY  
is indicated



accessible at the LHC

## Muon g-2 and SUSY

Could be new physics  
eg light SUSY



Observed Difference with Experiment:

$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{SM}} = (27.5 \pm 8.4) \times 10^{-10}$$

➔ 3.3 "standard deviations"

$$\delta a_{\mu} = 13 \cdot 10^{-10} \left( \frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \text{tg} \beta$$

$a_{\mu}$  is a plausible  
location for a  
new physics signal!!





# Little Higgs Models

Georgi (moose)/Arkani-Hamed et al/Low, Skiba, Smith/Kaplan, Schmaltz/Chang,Wacker/Gregoire et al

$$G \supset [SU(2) \otimes U(1)]^2 \supset SU(2) \otimes U(1)$$

↑
↑
↑

global
gauged
SM

H is (pseudo)-Goldstone boson of G: takes mass only at 2-loops (needs breaking of 2 subgroups or 2 couplings)

recall:  $\delta m_{h|top}^2 = -\frac{3G_F}{2\sqrt{2}\pi} m_t^2 \Lambda^2 \sim -(0.2\Lambda)^2$        $G_F \sim g^2 \rightarrow g^4$

cutoff  $\Lambda$

$\sim 10$  TeV

$\Lambda^2$  divergences canceled by:

- |                        |  |   |                |
|------------------------|--|---|----------------|
| $\delta m_{H top}^2$   | new coloured fermion $\chi$ with $Q=2/3$ | } | $\sim 1$ TeV   |
| $\delta m_{H gauge}^2$ | $W', Z', \gamma'$                        |   |                |
| $\delta m_{H Higgs}^2$ | new scalars                              |   |                |
| $\oplus$               | 2 Higgs doublets                         |   | $\sim 0.2$ TeV |

With some tension, Little Higgs models technically can work.

T parity interchanges the two  $SU(2) \times U(1)$  groups Cheng, Low

Standard gauge bosons are T even, heavy ones are T odd

Lightest T-odd particle stable  $\rightarrow$  Dark Matter

Technically sophisticated. But the main drawback is:

Little Higgs provides just a postponement:

UV completion beyond  $\sim 10$  TeV? GUT's?

Still important as it offers well specified signals and signatures for searching at the LHC:

a light Higgs, a new top-like fermion  $\chi$  to damp the top loop, new  $W', Z'$  for the  $W, Z$  loops,.....



# Extra Dimensions (ED)

String Theory ---> ED at  $M_{Pl}$

Perhaps ED have a direct impact on physics below  $M_{Pl}$

Exciting possibilities (a large domain of contemporary BSM)

- GUT's in ED ( $M_{GUT}$ )
- ED as (part of the) solution of the hierarchy problem ( $M_{EW}$ )
- EW symmetry breaking from ED ( $M_{EW}$ )

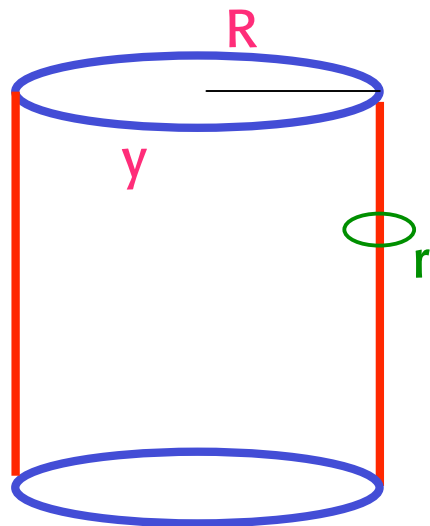


## Early formulation

Solve the hierarchy problem by bringing gravity down from  $M_{\text{Pl}}$  to  $o(1\text{TeV})$

Arkani-Hamed, Dimopoulos/ Dvali+Antoniadis

- **Large** compactified extra dimensions:  $1/R \ll \sim 1\text{TeV}$
- SM fields are on a brane
- Gravity propagates in the whole bulk



y: extra dimension  
R: compact'n radius

y=0 "our" brane (possibly with thickness r)

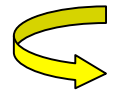
$G_N \sim 1/M_{\text{Pl}}^2$ :  
Newton const.  
 $M_{\text{Pl}}$  large as  
 $G_N$  weak

The idea is that gravity appears weak as a lot of lines of force escape in extra dimensions



# Generic feature of extra dim. models:

compact dim.  $\longrightarrow$  Kaluza-Klein (KK) modes



$$p=n/R$$

$$m^2=n^2/R^2$$

(quantization in a box)

Many possibilities:

emerges as the most promising formulation



- SM fields on a brane or in bulk
- cfr: • Gravity always on bulk

- Factorized metric:

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu + h_{ij}(y) dy^i dy^j$$

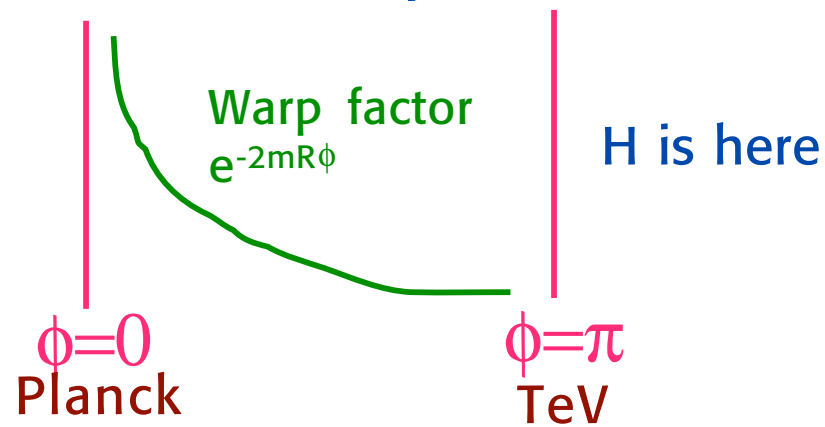
- Warped metric: Randall-Sundrum (R-S)

$$ds^2 = e^{-2mR|\varphi|} \eta_{\mu\nu} dx^\mu dx^\nu - R^2 \varphi^2$$



$$m_{\text{weak}} = M_{\text{pl}} \exp(-mR\pi) \longrightarrow Rm \sim 12$$

## A more promising scheme (warped ED)



All SM particles in bulk except the H

$$m \sim M_{Pl} \text{ for all } mR: m^2 \sim M_{Pl}^2 (1 - e^{-2mR\phi})$$

All 4-dim masses  $m_4$  are scaled down with respect to 5-dim masses  $m_5 \sim M_{Pl}$  by the warp factor:  $m_4 = M_{Pl} e^{-mR\pi}$

The hierarchy problem demands that  $mR \sim 12$ : not too large!!  
R not large in this case!

Stabilization of  $mR$  at a compatible value can be assured by a scalar field in the bulk with a suitable potential

⊕ "radion"

Goldberger, Wise

$$ds^2 = e^{-2mR|\phi|} \eta_{\mu\nu} dx^\mu dx^\nu - R^2 \phi^2$$

Randall-Sundrum:

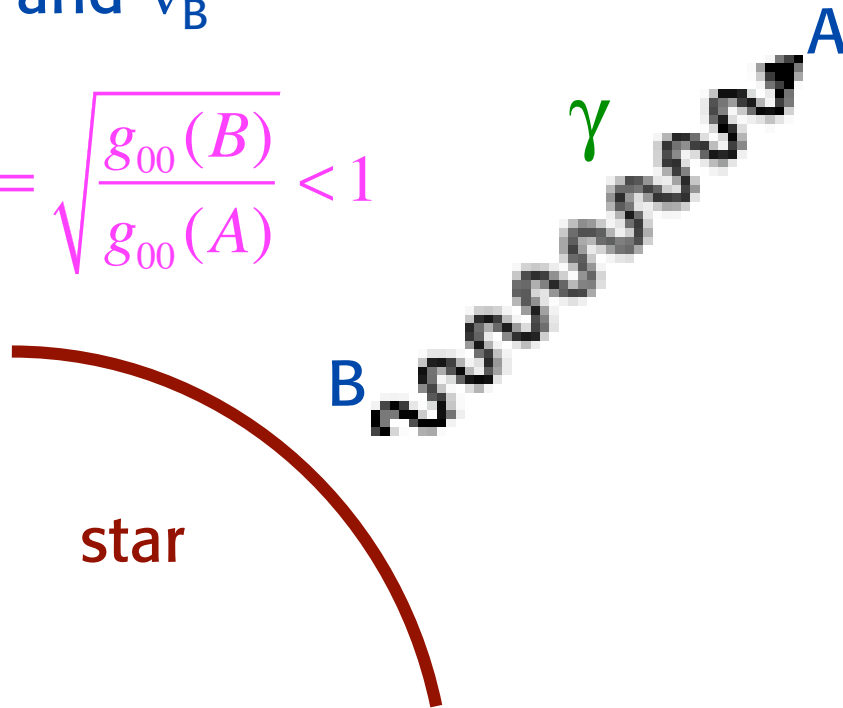
This non-fact.ble metric is solution of Einstein eq.s with 2 branes at  $\phi=0, \pi$  and specified 5-dim cosmological term

2 identical atoms in  
A and B emit light  
with frequencies

$\nu_A$  and  $\nu_B$

seen from A the B frequency is smaller:  
as if the photon kinetic energy lost by  
climbing out of grav. field

$$\frac{\nu_B}{\nu_A} = \sqrt{\frac{g_{00}(B)}{g_{00}(A)}} < 1$$



Similarly in RS  $mc^2$  is smaller  
by the corresponding factor  
 $g_{00}^{1/2} \rightarrow m_4 = M_{pl} e^{-mR\pi}$

Good tutorials:  
R. Sundrum '04  
TASI lectures  
R. Rattazzi '05  
Cargese Lectures  
Csaki et al '05  
Gherghetta'06

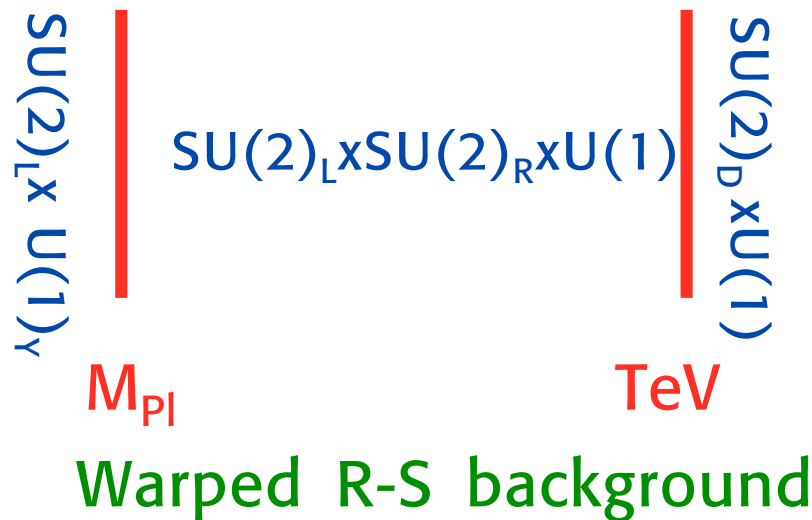


# Applications

- **Gauge Symmetry Breaking (Higgsless theories)**

Csaki et al/Grojean/Papucci/Nomura/Davoudiasl et al/Barbieri, Pomarol, Rattazzi;....

The only models where no Higgs would be found at LHC.  
But signals of new physics would be observed



Symmetries broken by  
Boundary Conditions (BC)  
on the branes

Altogether only  $U(1)_Q$   
unbroken

- Unitarity breaking (no Higgs) delayed by KK recurrences
- Dirac fermions on the bulk (L and R doublets). Only one  
⊕ chirality has a zero mode on the brane



With no Higgs unitarity violations, eg:

$$A(W_L^+ W_L^- \rightarrow Z_L Z_L) = \frac{G_F E^2}{8\sqrt{2}\pi}$$

At  $E \sim 1.2$  TeV unitarity is violated

In Higgsless models unitarity breaking is delayed by the exchange of KK recurrences

Cancellation guaranteed  
by sum rules implied  
by 5-dim symmetry

$Z_k = k_{th}$  KK

$$g_{WWWW}^2 - e^2 - \sum_k g_{WWZ_k}^2 = 0 ;$$

$$4M_W^2 g_{WWWW}^2 - 3 \sum_k g_{WWZ_k}^2 M_{Z_k}^2 = 0 .$$

The small  $W, Z$  mass implies a small KK gap  $\rightarrow W', Z'$  at the LHC

⊕ Higgsless models can also be formulated in 4 dimensions  
(pioneered by Casalbuoni, De Curtis, Dominici, Gatto '85)

No convincing, realistic Higgsless model for EW symmetry breaking emerged so far

Serious problems with EW precision tests

e.g. Chivukula et al '02, Barbieri, Pomarol, Rattazzi '03;

also with  $Z \rightarrow bb$

$m_W$  fixes the KK gap and it is not sufficiently large

Substantial fine tuning required

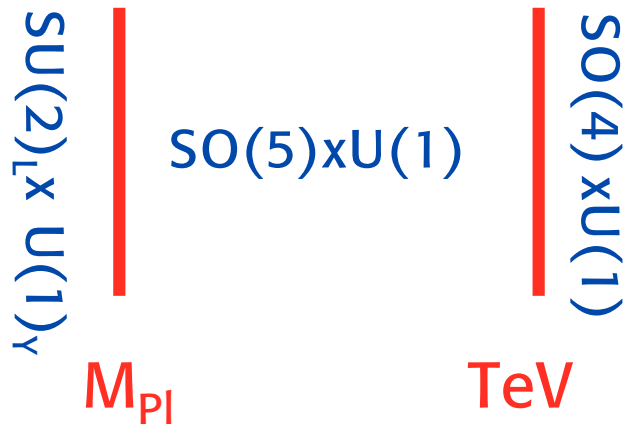
Best try: Cacciapaglia et al '06

However be alerted of possible signals at the LHC: no Higgs but KK recurrences of W, Z and/or additional gauge bosons



- Composite Higgs in a 5-dim holographic theory

Agashe, Contino, Pomarol.....



A new way to look at walking technicolor using AdS/CFT corresp.

All SM fields in the bulk (but the Higgs is localised on the TeV brane)

Warped R-S background

As in Little Higgs models

The Higgs is a PGB and EW symmetry breaking is triggered by bulk effects (in 4-dim the bulk appears as a strong sector).

The 5-dim theory is weakly coupled so that the Higgs potential and EW observables can be computed

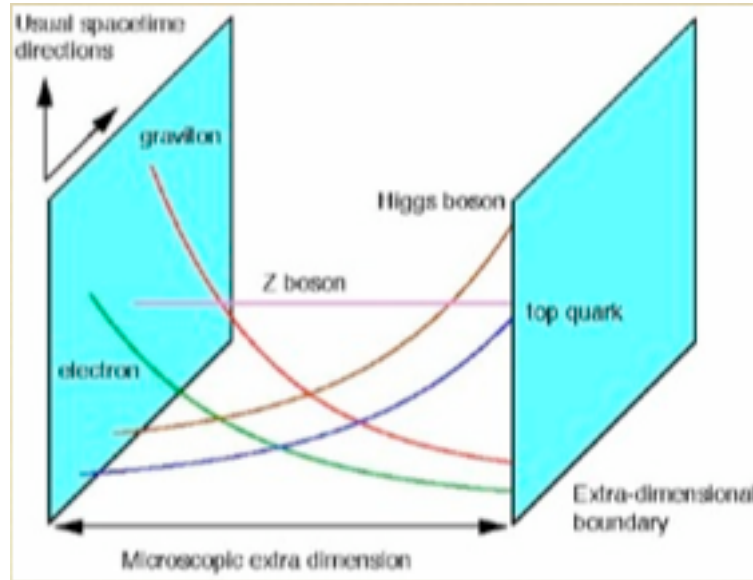
The Higgs is rather light:  $m_H < 185$  GeV



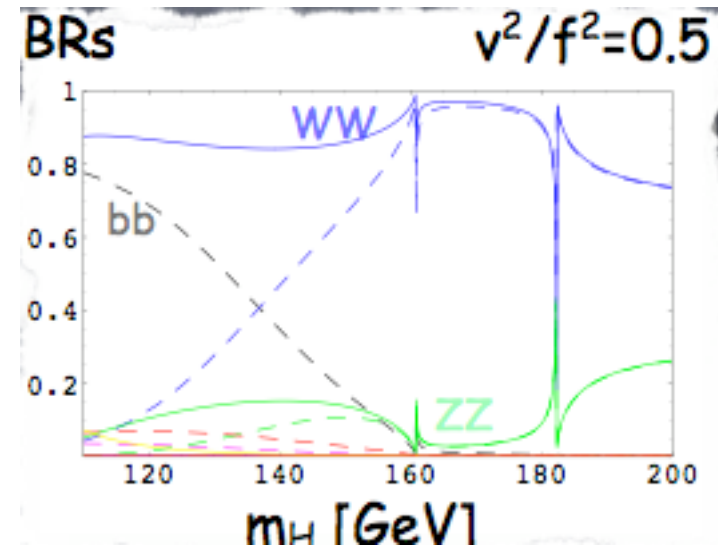
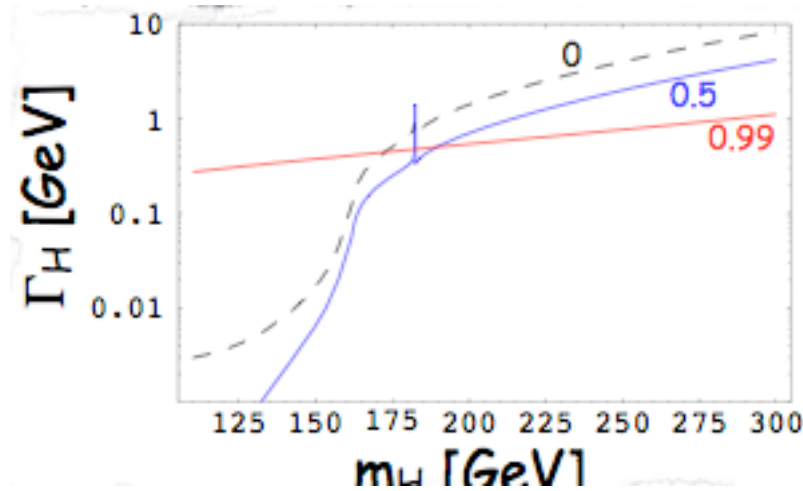
Also in these models a sizable fine-tuning is required

# A qualitative description of flavour

The fermion mass hierarchies explained by exp warp factors with  $o(1)$  exponents



## Higgs couplings modified



- Composite Higgs: a more model indep. approach

Georgi, Kaplan '84

The light Higgs is a bound state of a strongly interacting sector.  
Pseudo-Goldstone boson of an enlarged symmetry.

eg.  $SO(5)/SO(4)$

Agashe/ Contino/Pomarol/Sundrum/ Grojean/Rattazzi....

$v \sim$  EW scale       $f \sim$  SI scale

$$\sim f < m_\rho < \sim 4\pi f$$

$$\xi = (v/f)^2$$

$\xi$  interpolates between SM [ $\xi \sim 0$ ]

and some degree of  
compositeness

[ $\xi \sim o(1)$  limited by precision EW tests]

$m_\rho$



$m_H$

$m_W$

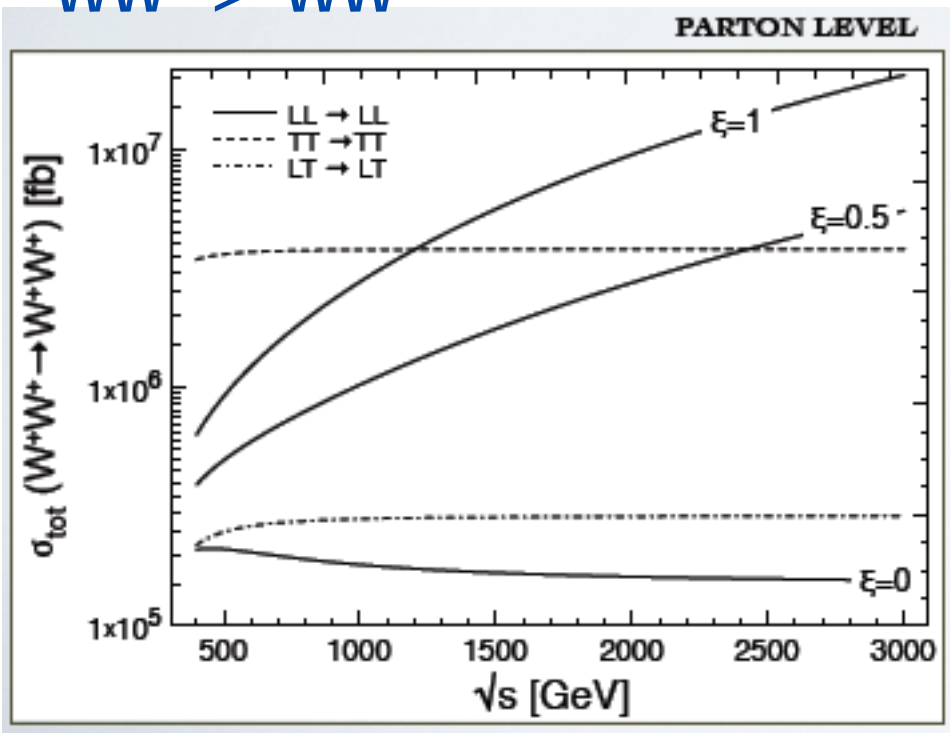


# Detectable $\xi$ effects at the LHC

- Higgs couplings
- WW scattering
- 2-Higgs Production

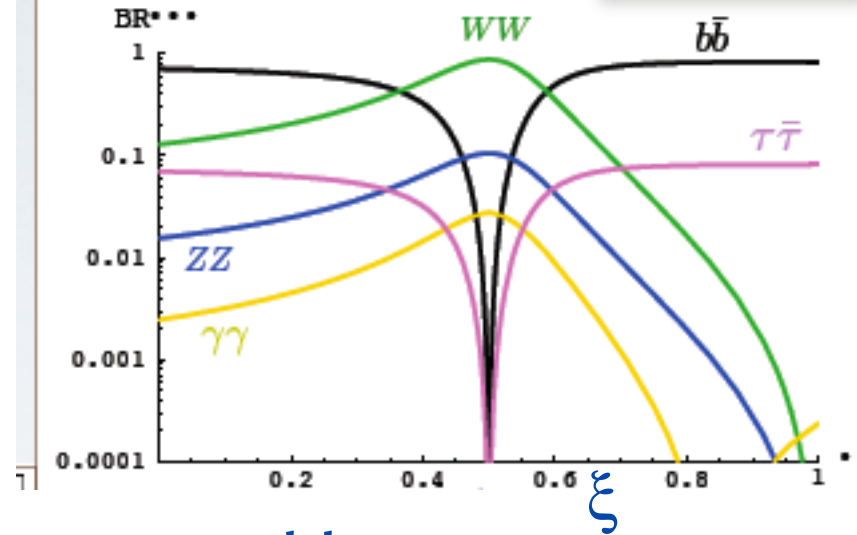
WW  $\rightarrow$  WW

Contino et al

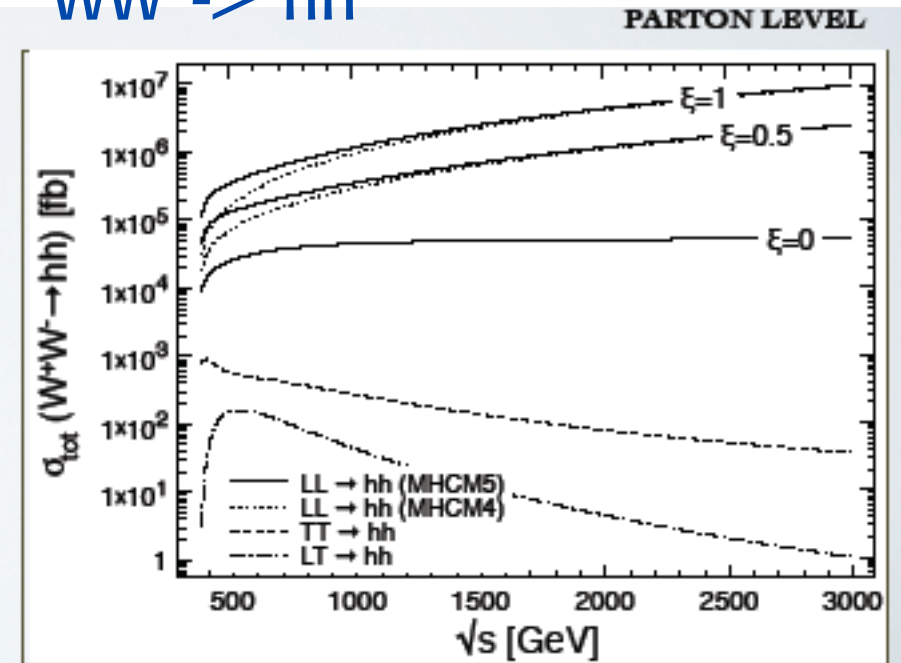


# H Br Ratios

$m_h = 120 \text{ GeV}$



WW  $\rightarrow$  hh



## Lessons from model building

In all the new physics models we mentioned

there is a light Higgs ( $< 200$  GeV)

[except in Higgsless models (if any) but new light new vector bosons exist in this case]

there is at least a % fine tuning

Fine tuning appears to be imposed on us by the data



## In conclusion

Is it possible that the LHC does not find the Higgs particle?

Yes, it is possible, but then something else must be found

Is it possible that the LHC finds the Higgs particle but no other new physics (pure and simple SM)?

Yes, it is technically possible but it is not natural

Is it possible that the LHC finds neither the Higgs nor new physics?



No, it is "approximately impossible"