Higher Order Corrections to the Drell-Yan Cross Section in the Mellin Space

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- The Drell-Yan mechanism
- The Mellin transform
- Calculation of the Drell-Yan structure functions in Mellin space up to NNLO
 - Coefficient functions
 - Parton distribution functions
- Partial results and checks
- Conclusions and outlook

Massive lepton pair production in hadron-hadron collision, $M_{l_1l_2}^2 >> 1 \,\mathrm{GeV}^2$ [Drell, Yan 1970]



Neutral Current

- $pp \to \gamma^* \to ll X$ $M_{l\bar{l}} \ll M_Z$ $M_Z \sim 91.2 \text{ GeV}$
- $pp \to Z^0 \to l\bar{l}X \qquad M_{l\bar{l}} \sim M_Z$
- **Charged Current**
 - $pp \to W^{\pm} \to l\nu X$ $M_{l\bar{l}} \sim M_W$ $M_W \sim 80.4 \,\mathrm{GeV}$

- $M_{l_1 l_2} = Q^2 = (p_2 + p_3)^2$
- CM energy of hadrons $s = (P_1 + P_2)^2$
- $p_i = x_i P_i \qquad x_i \in (0,1)$
- CM energy of partons $\hat{s} = (p_1 + p_2)^2 = sx_1x_2$

The Drell-Yan process at the LHC



- Large total cross sections for W and Z production
- Clear signal
- Background for new phyiscs measurements
- Test of of the Standard Model
- Constraints on parton distribution functions

$$\frac{\mathrm{d}\,\sigma_{DY}^{V}(Q^{2})}{\mathrm{d}\,Q^{2}} = \sum_{a,b} \int_{0}^{1} \mathrm{d}\,x_{1} \int_{0}^{1} \mathrm{d}\,x_{2} \underbrace{f_{a}(x_{1},Q^{2})f_{b}(x_{2},Q^{2})}_{\mathsf{PDFs}} \underbrace{\frac{\mathrm{d}\,\hat{\sigma}_{ab}^{V}(Q^{2})}{\mathrm{d}\,Q^{2}}}_{\mathsf{Hard Cross section}}$$

$$V = \gamma^*, Z^0, W^{\pm}$$
 $a, b = q, \bar{q}, g$ $x_{1,2} \in (0, 1)$

- Higher order NLO QCD corrections [Altarelli, Ellis, Martinelli 1979]
 NNLO [Hamberg, Matsuura, van Neerven 1990; Harlander, Kilgore 2002]
- NNLO fully differential cross section [Ferrera, Grazzini 2010] ⇒ new constraints on PDFs
- Electroweak corrections up to NLO

[Anastasiou, Dixon, Melnikov, Petriello 2004; Catani,

[Alekhin, Melnikov, Petriello 2006]

[Hollik, Wackeroth 1996; Vicini et.al. 2009]

The Drell-Yan Cross Section

$$\frac{\mathrm{d}\,\sigma_{DY}^{V}(Q^{2})}{\mathrm{d}\,Q^{2}} = \sum_{a,b} \int_{0}^{1} \mathrm{d}\,x_{1} \int_{0}^{1} \mathrm{d}\,x_{2} f_{a}(x_{1},Q^{2}) f_{b}(x_{2},Q^{2}) \frac{\mathrm{d}\,\hat{\sigma}_{ab}^{V}(z,\alpha_{s})}{\mathrm{d}\,Q^{2}}$$

$$= xN^{V} \underbrace{\sum_{a,b} C_{ab}^{V}(f_{a} \otimes f_{b} \otimes \Delta_{ab})(x)}_{\mathbf{Structure function } W_{DY}(x,Q^{2})} \quad x = \frac{Q^{2}}{s} \quad z = \frac{x}{x_{1}x_{2}}$$

$$\Delta_{ab} = \sum_{k} \left(\frac{\alpha_{s}}{4\pi}\right)^{k} \Delta_{ab}^{(k)} \quad \text{Coefficient functions}$$

Convolution

$$(f_1 \otimes f_2 \otimes \cdots \otimes f_k)(x)$$

= $\int_0^1 \mathrm{d} x_1 \int_0^1 \mathrm{d} x_2 \dots \int_0^1 \mathrm{d} x_k \,\delta(x - x_1 x_2 \dots x_k) f_1(x_1) f_2(x_2) \dots f(x_k)$

The Convolution and The Mellin Transform

• The Mellin transform - integral transform f(x) ("x-space") $\rightarrow \tilde{f}(N)$ ("N-space")

$$\mathbf{M}[f(x)] = \int_0^1 \mathrm{d} x \, x^{N-1} f(x) = \tilde{f}(N)$$

Mellin transform of the convolution is a product of functions in N-space

$$\mathbf{M}[(f_1 \otimes f_2 \otimes \cdots \otimes f_k)(x)](N) = f_1(N)f_2(N) \dots f_k(N)$$

The Convolution and The Mellin Transform

The structure function

 $W(x,Q^2) = (f_a \otimes f_b \otimes \Delta_{ab}^V)(x) \quad \to \quad f_a(N)f_b(N)\Delta_{ab}^V(N) = \tilde{W}(N,Q^2)$

- Two integrations \rightarrow simple product
- Need analytic expressions for all input functions in x-space for which the Mellin transforms exist [Vermaseren 1998]
 [Moch, Vermaseren, Vogt 2004; Blümlein, Kurth 1998 (2000); Blümlein, Ravindran 2005, ...]
- Original x-space recovered performing the inverse Mellin transform

The Inverse Mellin Transform



$$W(\tau) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \mathrm{d} N \, x^{-N} \tilde{W}(N)$$



- All poles of $\tilde{W}(N)$ lie to the left from the contour c
- In general inverse MT are performed numerically

The Inverse MT - numerical evaluation

Rewriting the integral

$$\frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \mathrm{d}\, N\, x^{-N} \tilde{f}(N) \quad = \quad \frac{1}{\pi} \int_0^\infty \mathrm{d}\, \rho \, \mathrm{Im}[e^{i\varphi} x^{-c-\rho e^{i\varphi}} f(N)]$$

 \rightarrow integral over a real variable ρ $N=c+\rho e^{i\phi} \text{ becomes complex variable}$

- Integration only up to ρ_{\max} due to the factor $\exp[\rho \ln(1/x) \cos \varphi]$
- φ , c and number of N's can be tuned to get better accuracy
- Need analytic continuations of all N-space functions to the complex plane



Coefficient functions in N-space

- Starting from x-space expressions
 [Ravindran, unpublished]

 Mellin transforms obtained using harmpol package
 [Remiddi, Vermaseren 2000]
- N-space results (NNLO) have been published

[Blümlein, Ravindran 2005]

- $\rightarrow N$ -space results mostly in terms of harmonic sums $S_k(N)$ with "most right" pole at N =1.
- Analytic continuations of harmonic sums in terms of polygamma functions
- Finite set of more complicated harmonic sums needs special treatment in order to get the analytic continuations, e.g. ancont
 [Blümlein, 2000]

Evolution of Parton Distribution Functions

• Scale dependence of $f_a(x, Q^2)$ described by evolution equations (DGLAP)

$$\frac{\partial f_a(x,Q^2)}{\partial \ln Q^2} = \sum_b \underbrace{P_{ab}(x,Q^2)}_{\text{Splitting functions}} \otimes f_b(x,Q^2)$$

The solution

$$f_a(x,Q^2) = E\Big[P_{ab}(x), \alpha_s(Q^2), \underbrace{f_b(x,Q_0^2)}_{Q_0^2} \sim \text{few GeV}\Big]$$

Parton Distribution Functions in *N***-space**

- Evolution in N-space
 - Choose initial distribution $f(x, Q_0^2)$, typical parametrization

$$\begin{split} xf(x,\mu_0^2) &= ax^b(1-x)^c(1+dx^f+gx), \qquad a,\dots,g \in \mathbb{R} \\ f(N,Q_0^2) &= a\Big[\beta(N+b-1,c+1)+\beta(N+b+f-1,c+1)+\\ &\qquad \beta(N+b,c+1)\Big] \end{split}$$

• Splitting functions in *N*-space

Evolve PDFs in N-space

[Moch, Vermaseren, Vogt 2004] [QCD-PEGASUS, Vogt 2004]

$$f_a(N,Q^2) = E\Big[P_{ab}(N), \alpha_s(Q^2), f_b(N,Q_0^2)\Big]$$

• Evolution in x-space (use LHAPDF) \rightarrow find an interpolation function for $f(x, Q^2)$ and transform it to N-space.

The Calculation - Indirect Check

- Choose a toy initial parton distribution function, perform evolution in $N\text{-space with } \text{QCD-PEGASUS} \to \tilde{\sigma}_{\text{toy}}(N,Q^2)$
- Choose the same initial distribution in x-space and do a standard convolution $\to \sigma_{toy}(x,Q^2)$ [CUBA, Hahn 2005]
- Compare $\mathbf{M}[\tilde{\sigma}_{toy}(N,Q^2)] \Leftrightarrow \sigma_{toy}(x,Q^2).$
- Replace toy pdfs in x-space with real ones (LHAPDF) $\sigma_{toy}(x, Q^2) \rightarrow \sigma_{LHAPDF}(x, Q^2)$ and compare to publicly available codes
 - MCFM (NLO)

ZWPROD

- [Hamberg, Matsuura, van Neerven, 2002]
- DYNNLO
 [Catani, Cieri, Ferrera, de Florian, Grazzini 2009; Catani, Grazzini 2007]
- FEWZ

[Melnikov, Petriello]

[Campbell, Ellis]

The Calculation - Direct Check

• Choose a real initial parton distribution function \rightarrow Mellin transform, perform evolution in N-space with QCD-PEGASUS and compare $IM[\sigma(N, Q^2)]$ to cross section calculated by publicly available codes

- Needs to extend slightly QCD-PEGASUS
- Not possible for any distributions in LHAPDF some initial parametrizations are not in form that Mellin transform on them can be performed

• Alternatively, access directly evolved PDFs from LHAPDF grids \rightarrow interpolation \rightarrow Mellin transform

QG contribution to the **DY** structure function



QG contribution to the DY structure function



Summary and Outlook

- Higher order corrections for Drell-Yan process are necessary for precise predictions
- Calculation in Mellin space is fast and accurate

- NNLO cross section
- Interface for PDFs accessing directly LHAPDF grids
- Implementation of DIS
- Release the code

Backup

Analytic continuation of a Harmonic Sum

$$S_k(N) = \frac{(-1)^{k+1}}{(k-1)!} \psi^{k-1}(N+1) + \zeta(k)$$

The Euler Beta function

$$\beta(a,b) = \frac{\Gamma(p)\Gamma(q)}{\Gamma(p+q)}$$