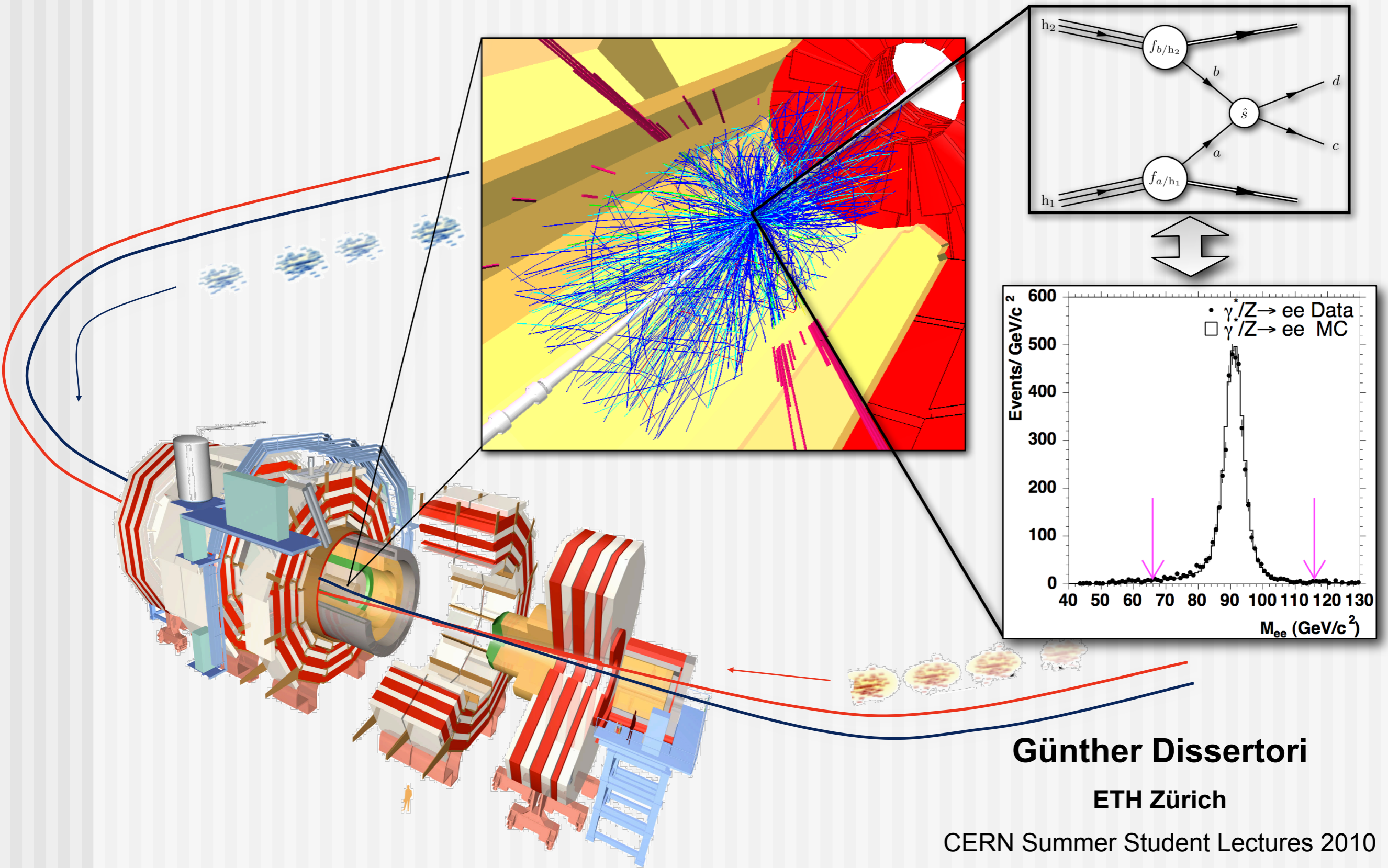


From Raw Data to Physics Results

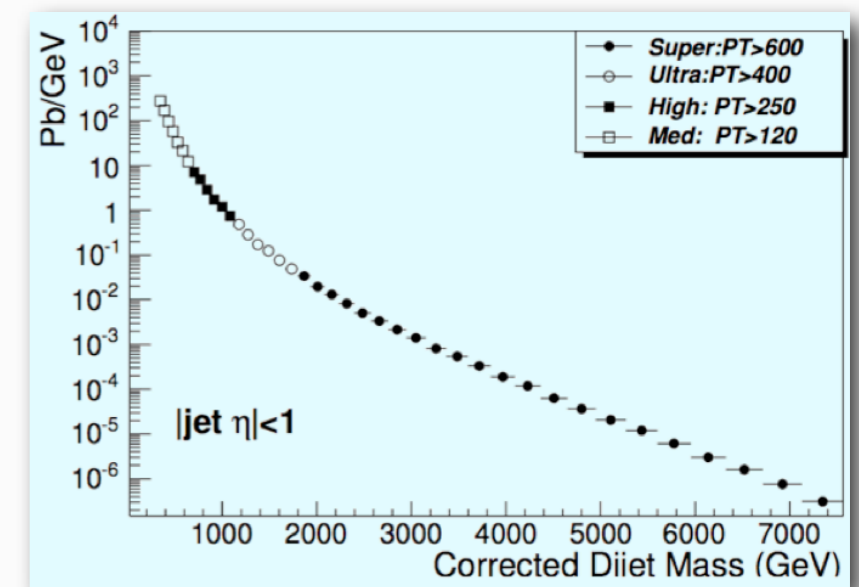
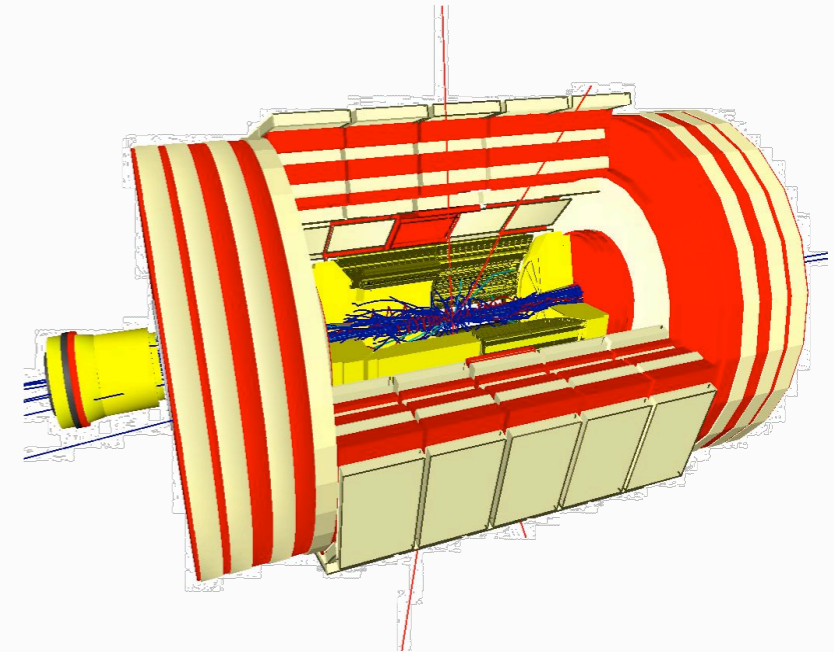


Günther Dissertori

ETH Zürich

CERN Summer Student Lectures 2010

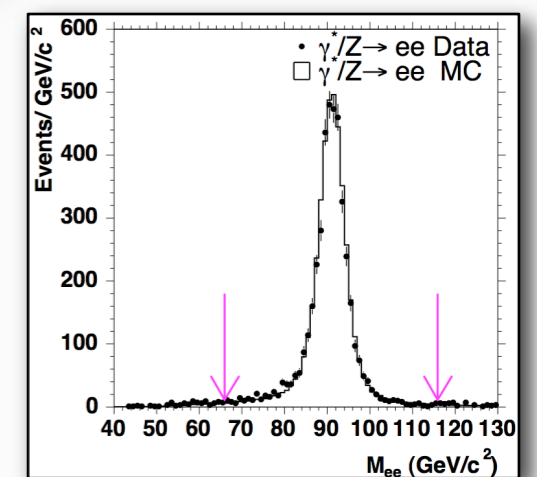
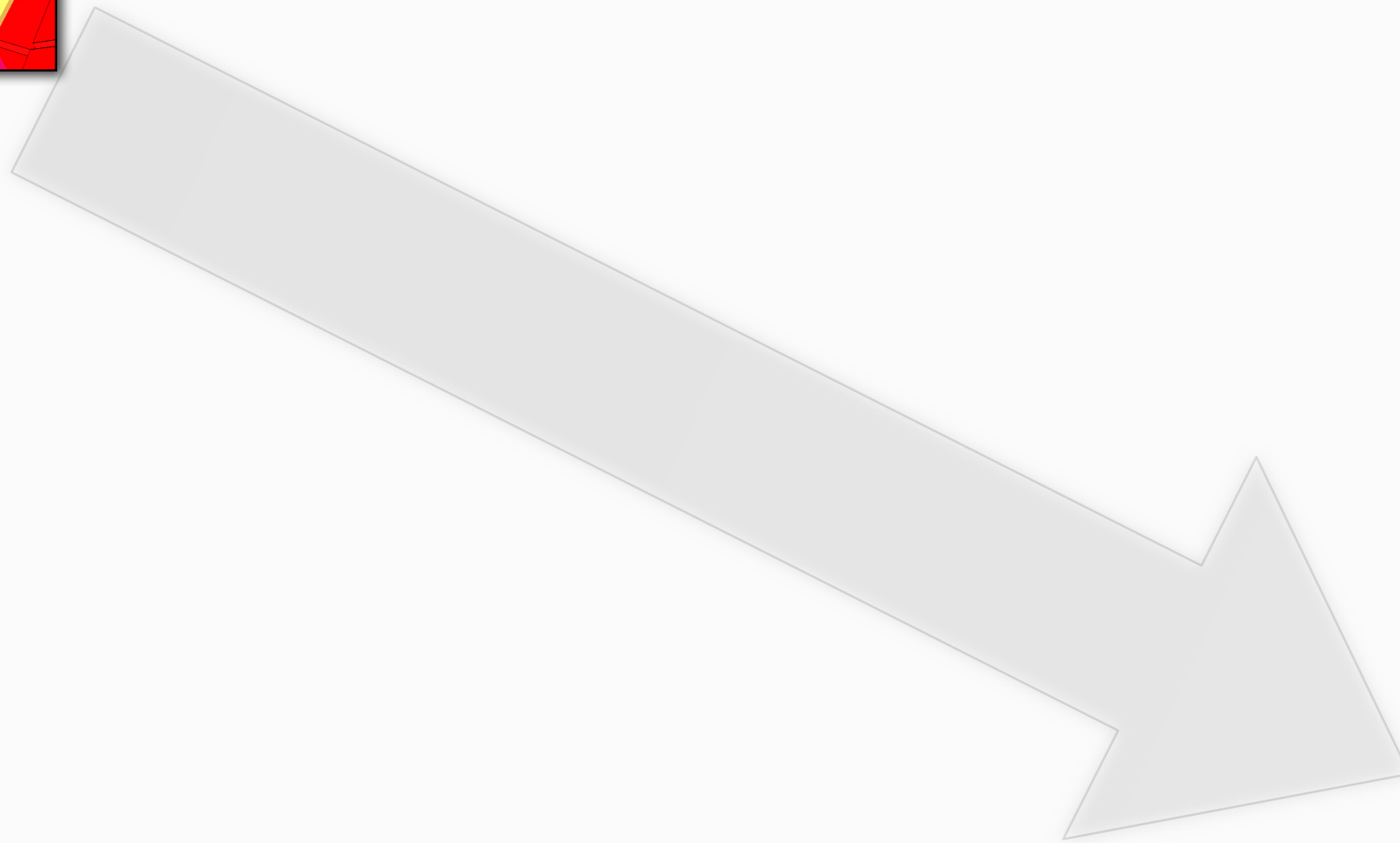
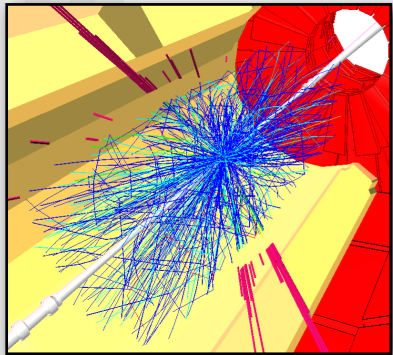
- “Executive Summary”
 - The whole story in a nutshell
- Some more details
 - Introduction
 - A simple example : Z decays
 - A more complicated example : Jets
- How is it done in practice?
 - Track and Calorimeter energy reconstruction
 - High-level algorithms : Jets
 - The computing part
- Summary



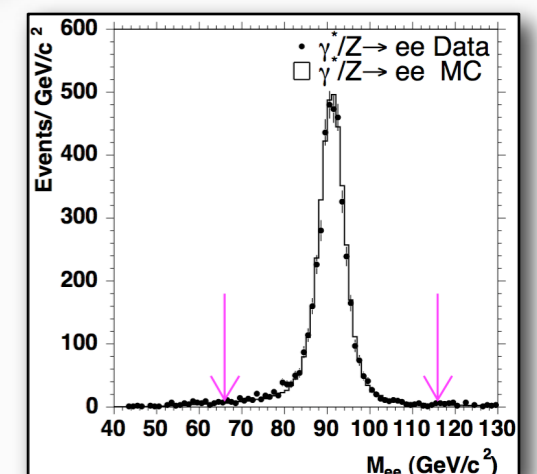
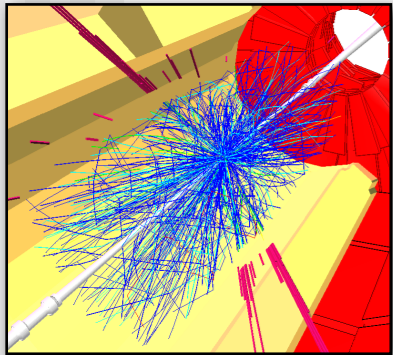
⌘ Disclaimer : Several slides based on past CSS lectures by B. Jacobsen
 ⌘ thanks also to J. Weng, T. Punz, A. Valassi

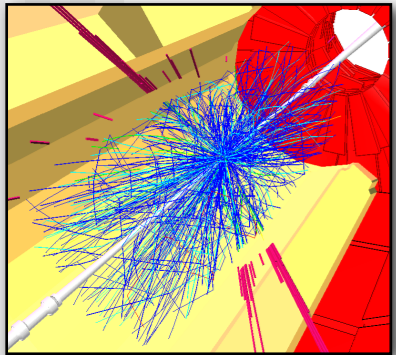
The whole lecture in a nutshell

“But you should not leave immediately after this....”

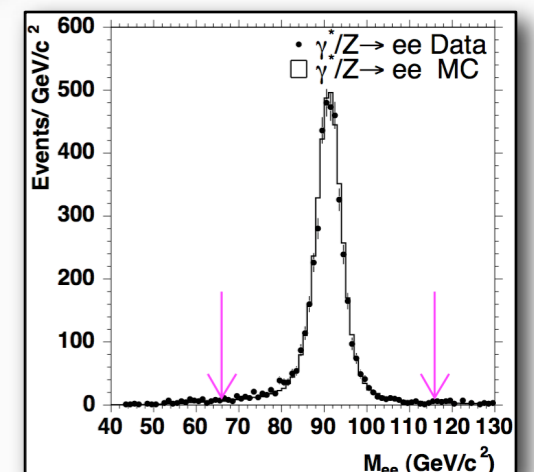
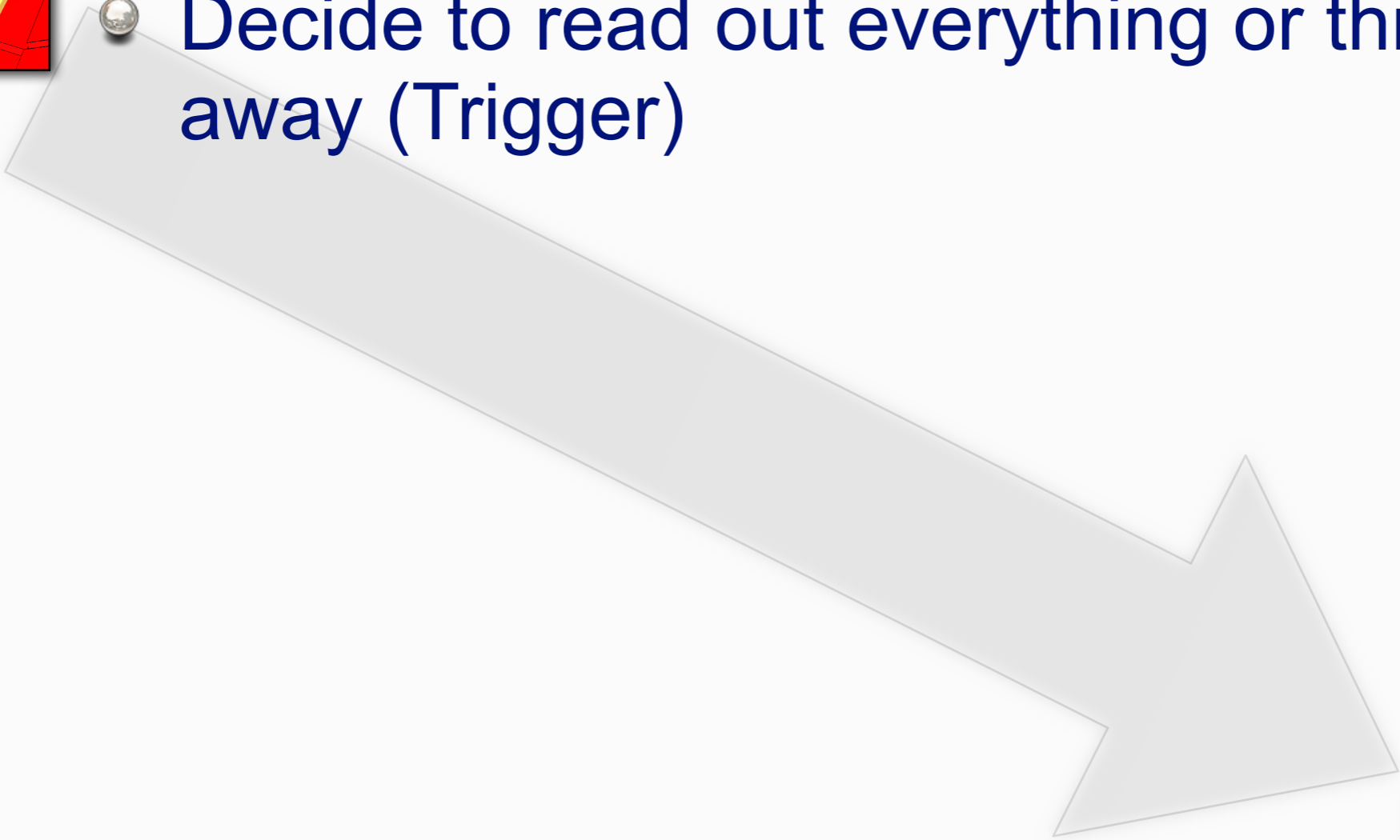


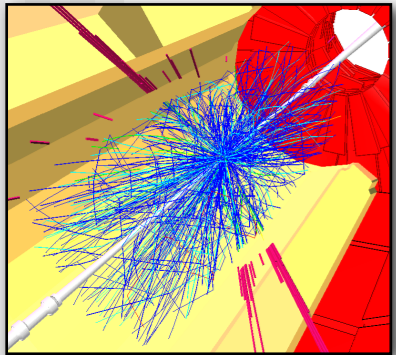
- Have to collect data from many channels on many sub-detectors (millions)



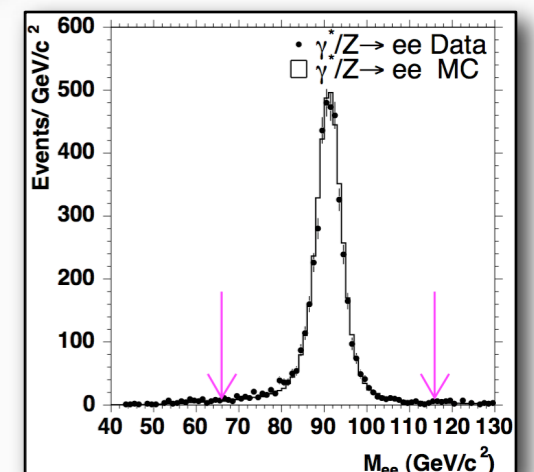
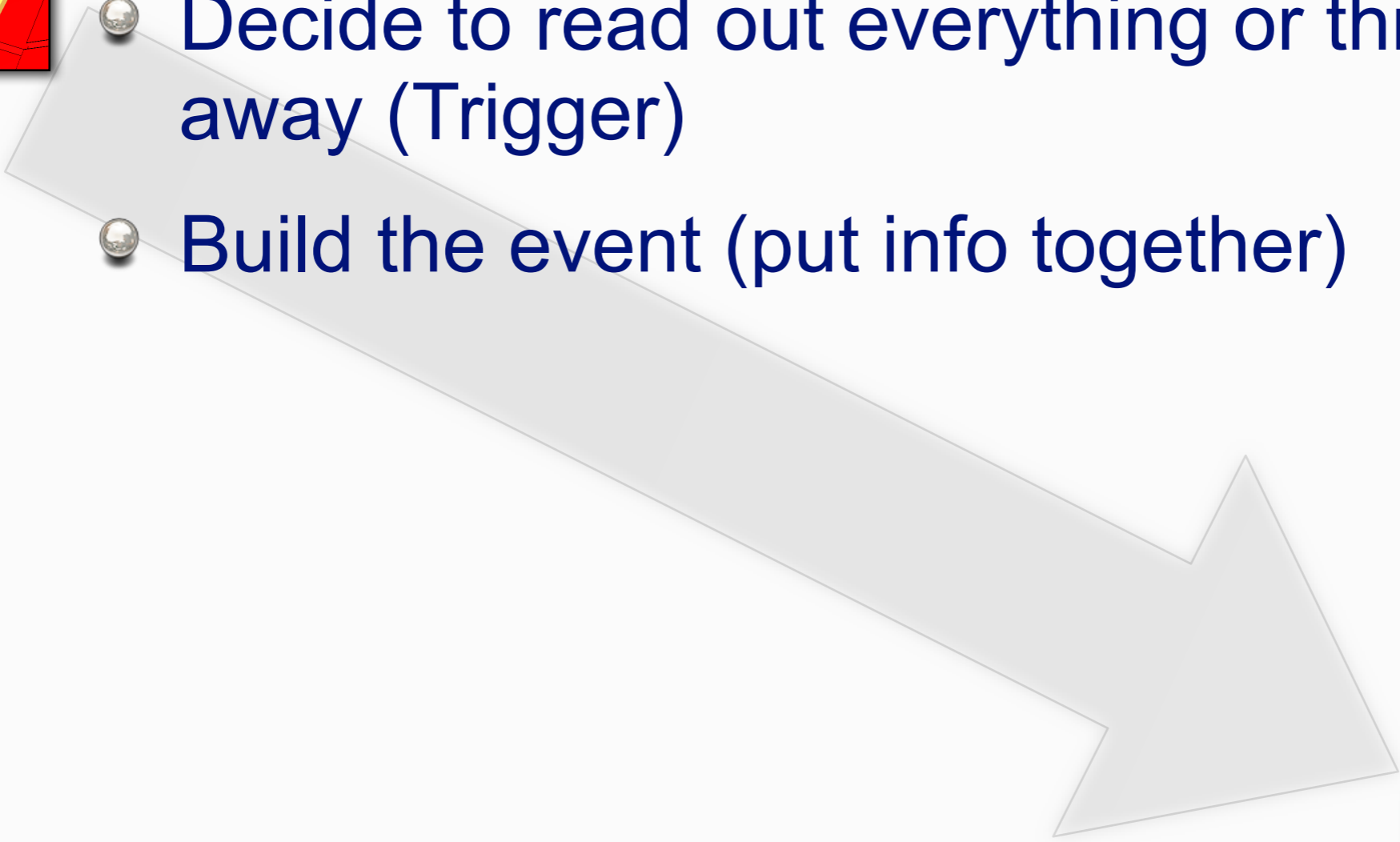


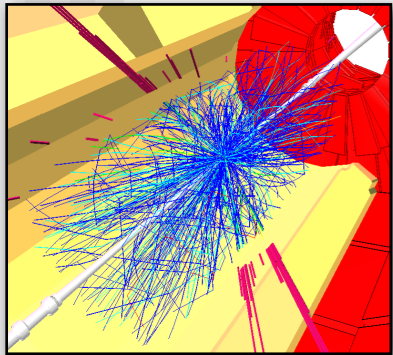
- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)



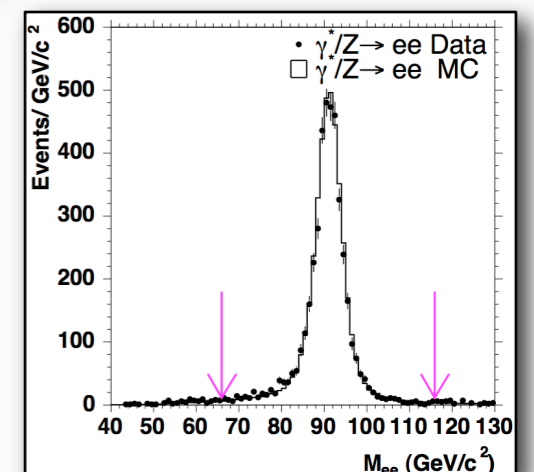
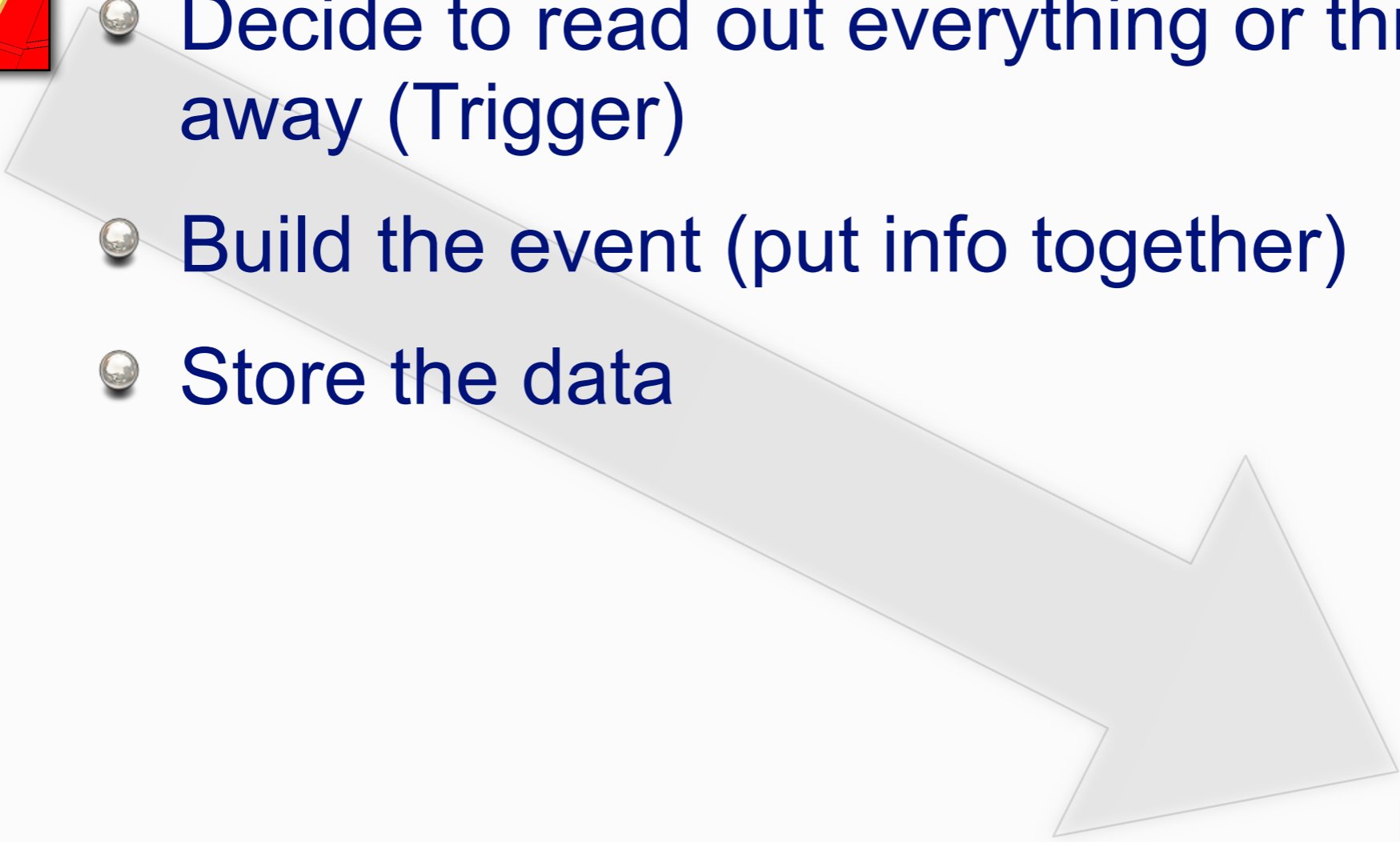


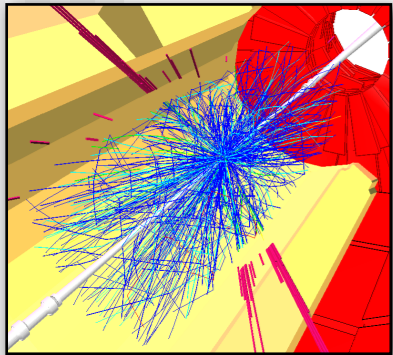
- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)



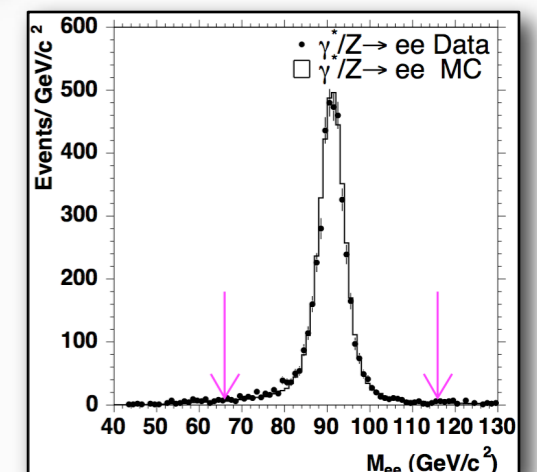


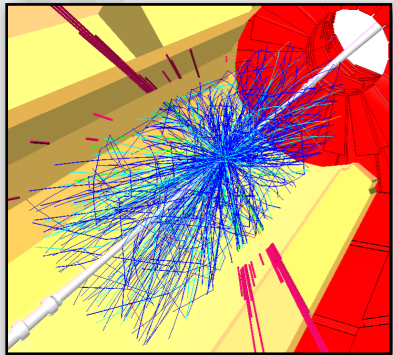
- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)
- Store the data



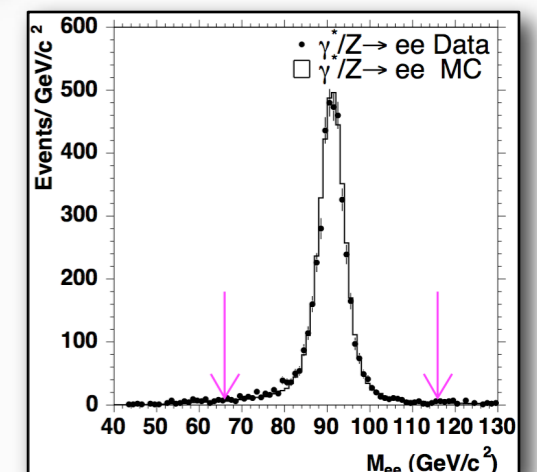


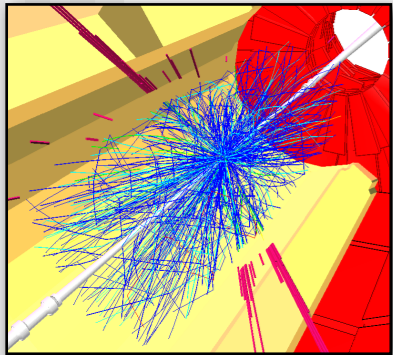
- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)
- Store the data
- Analyze them
 - reconstruction, user analysis algorithms, data volume reduction



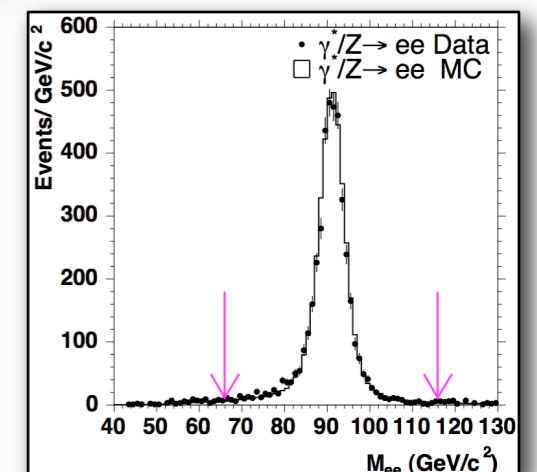


- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)
- Store the data
- Analyze them
 - reconstruction, user analysis algorithms, data volume reduction
- do the same with a simulation
 - correct data for detector effects





- Have to collect data from many channels on many sub-detectors (millions)
- Decide to read out everything or throw event away (Trigger)
- Build the event (put info together)
- Store the data
- Analyze them
 - reconstruction, user analysis algorithms, data volume reduction
- do the same with a simulation
 - correct data for detector effects
- Compare data and theory



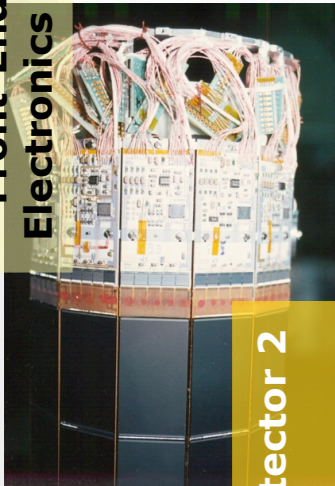


DAQ chain (see lectures by N. Neufeld)

**Detector
Front-End**

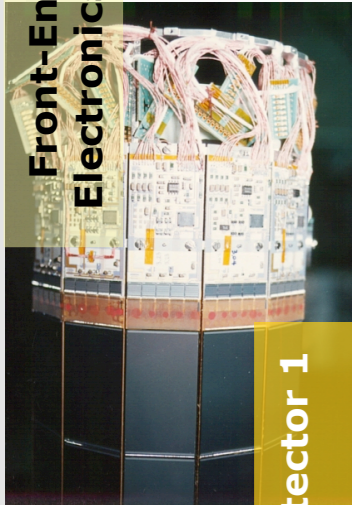
▪
▪
▪
▪

**Front-End
Electronics 2**



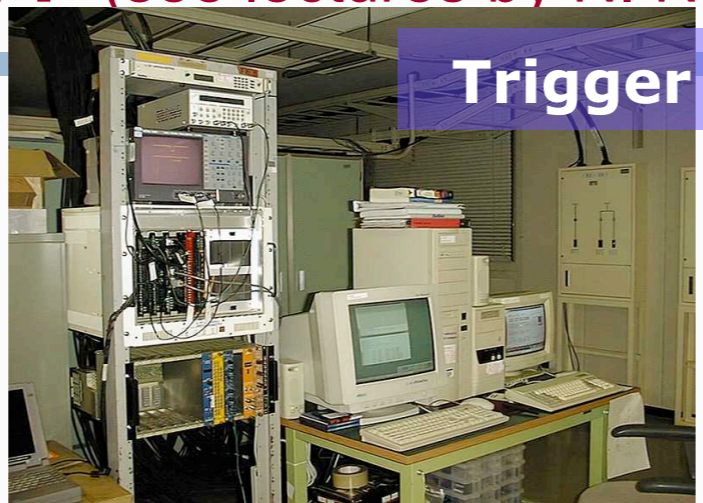
Detector 2

**Front-End
Electronics 1**

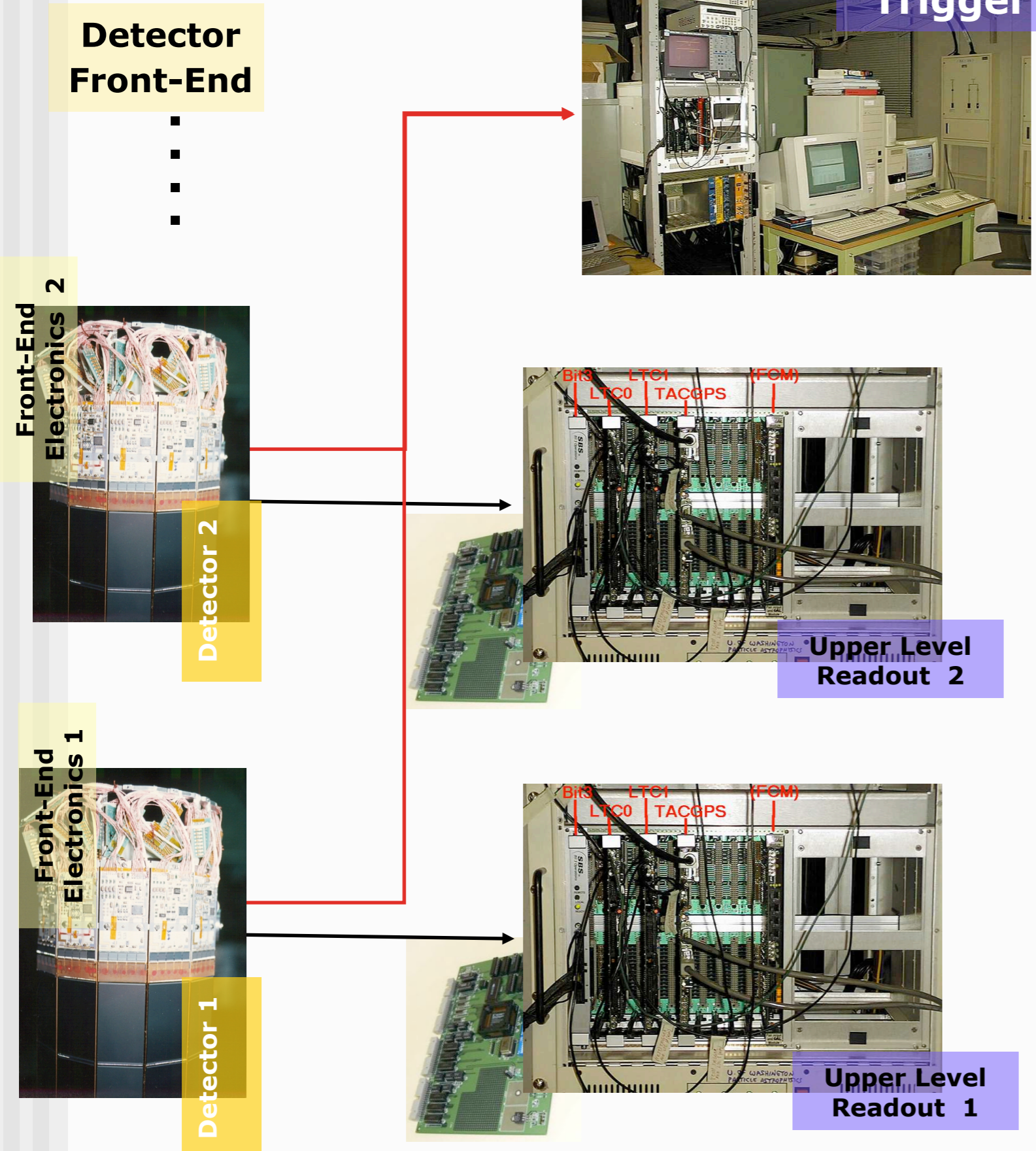


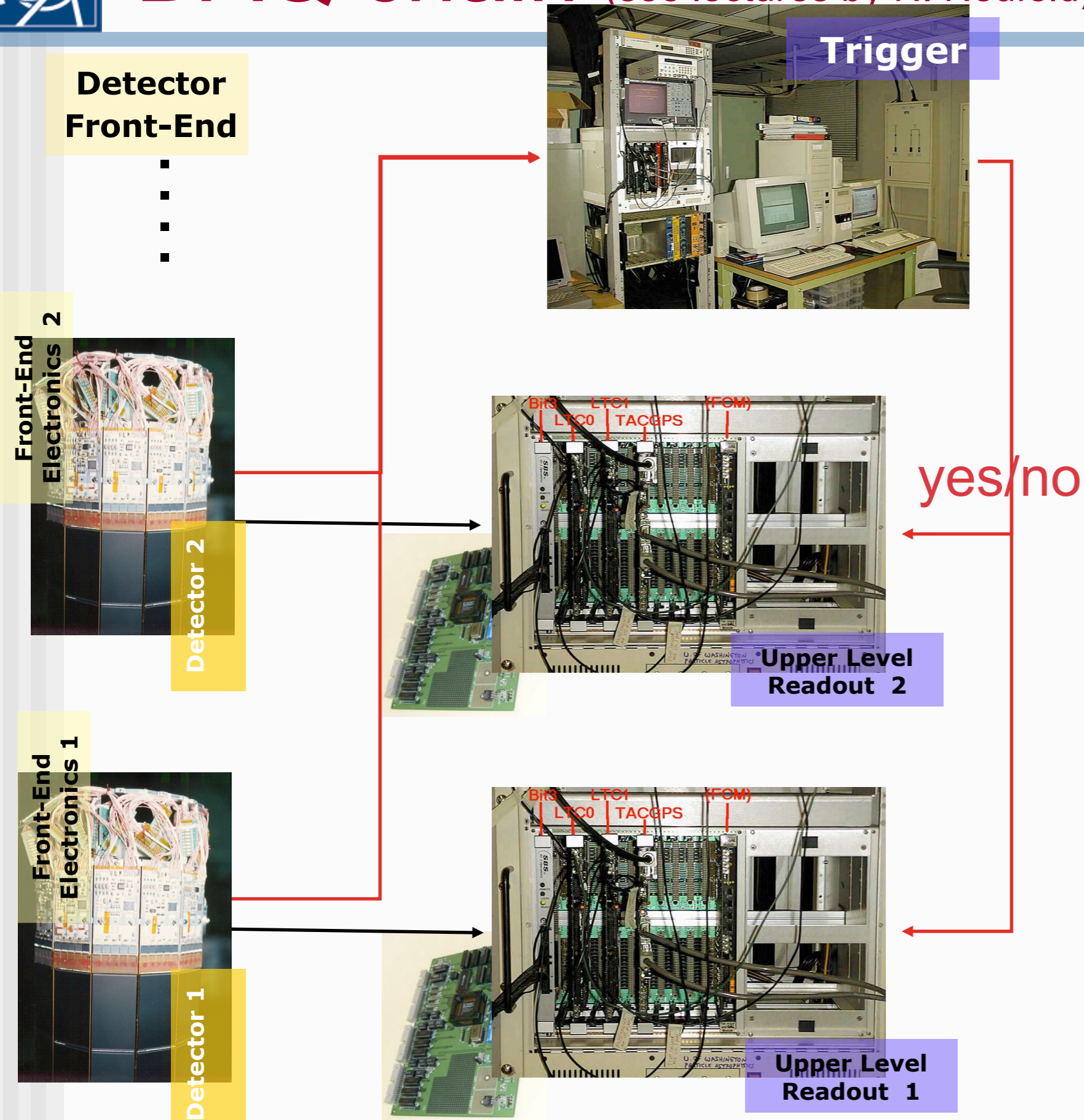
Detector 1

DAQ chain (see lectures by N. Neufeld)

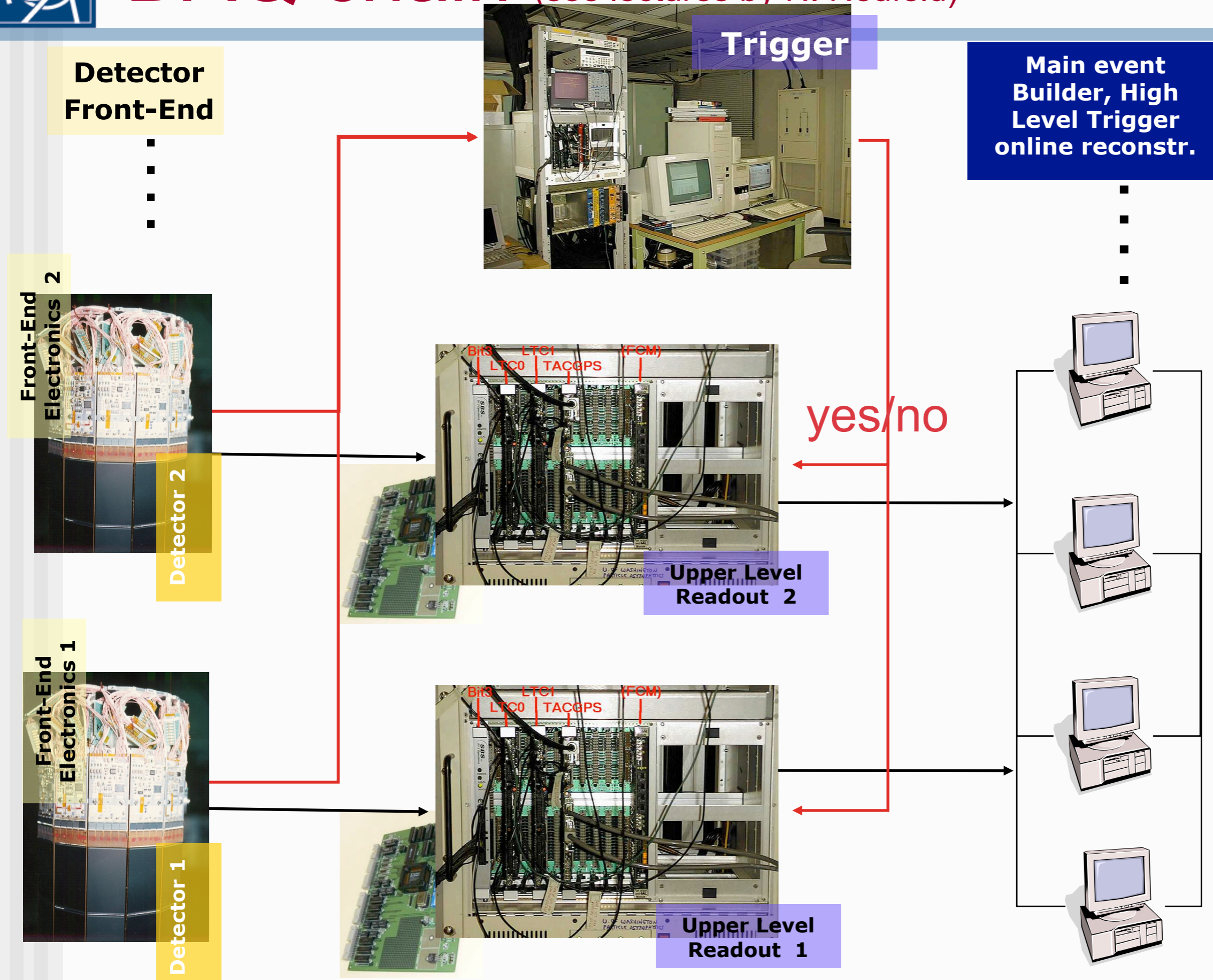


DAQ chain (see lectures by N. Neufeld)

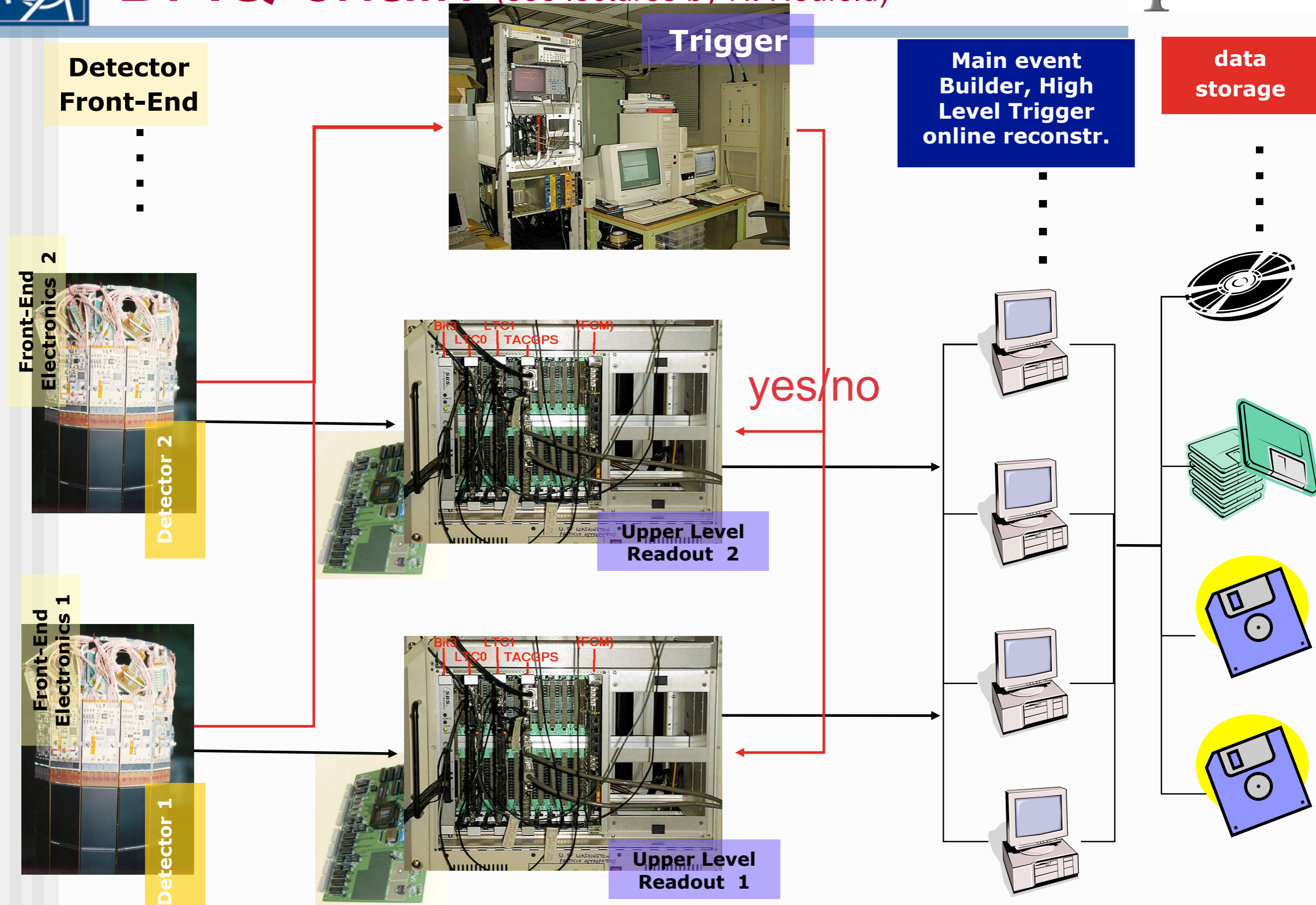




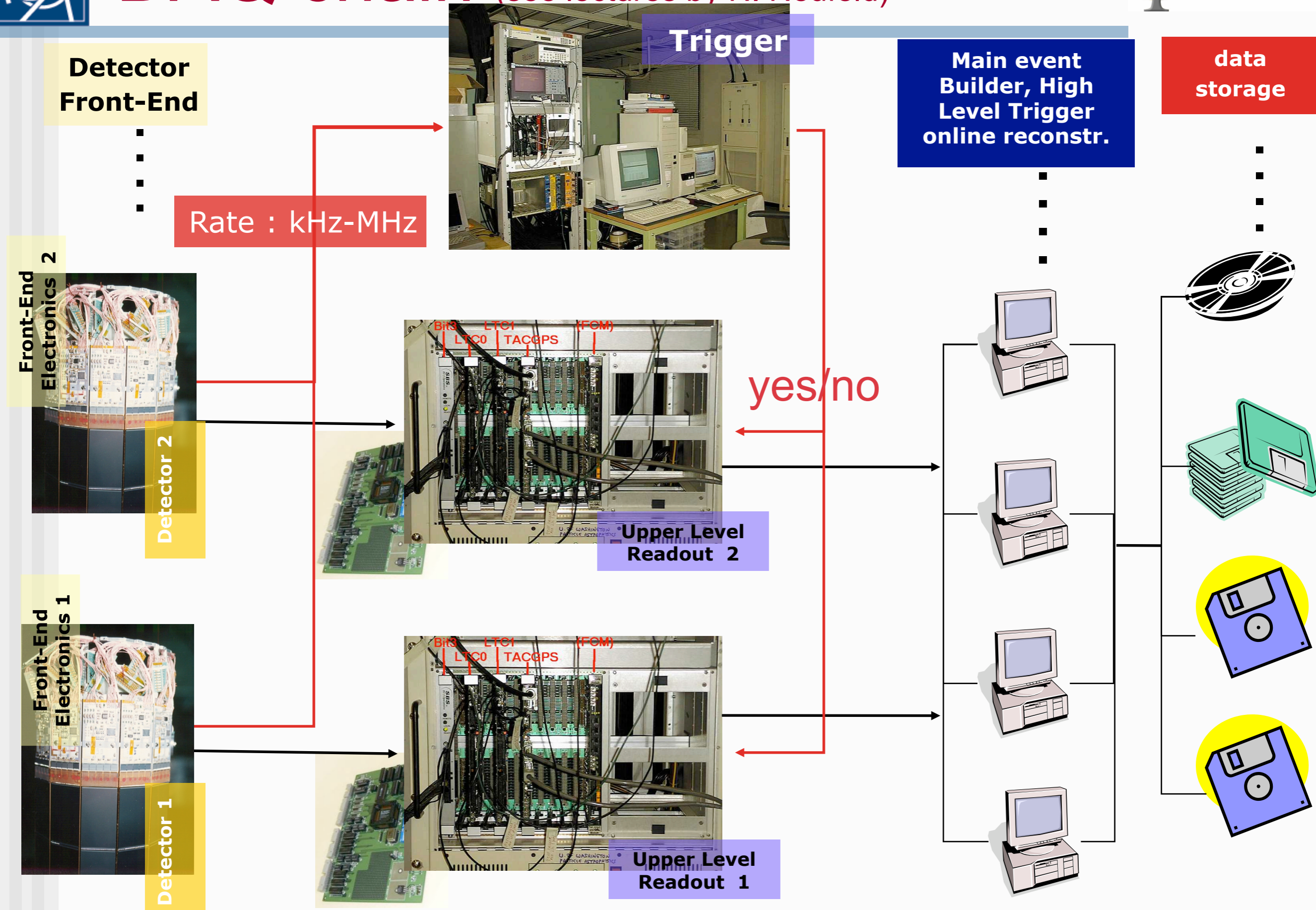
DAQ chain (see lectures by N. Neufeld)



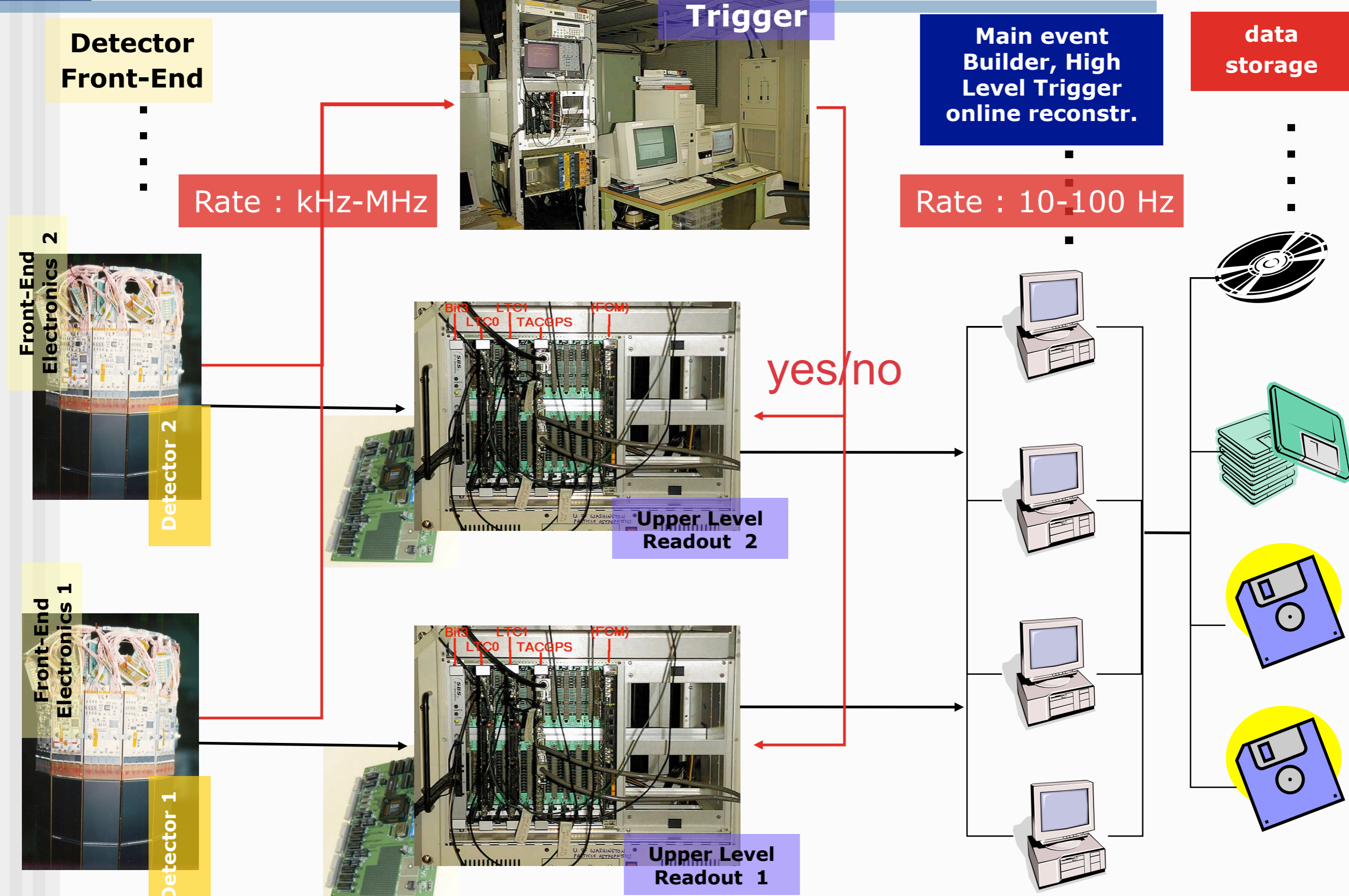
DAQ chain (see lectures by N. Neufeld)



DAQ chain (see lectures by N. Neufeld)



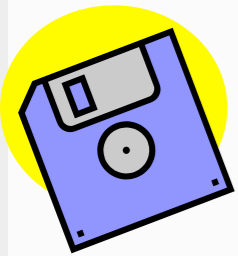
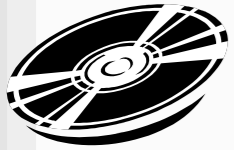
DAQ chain (see lectures by N. Neufeld)





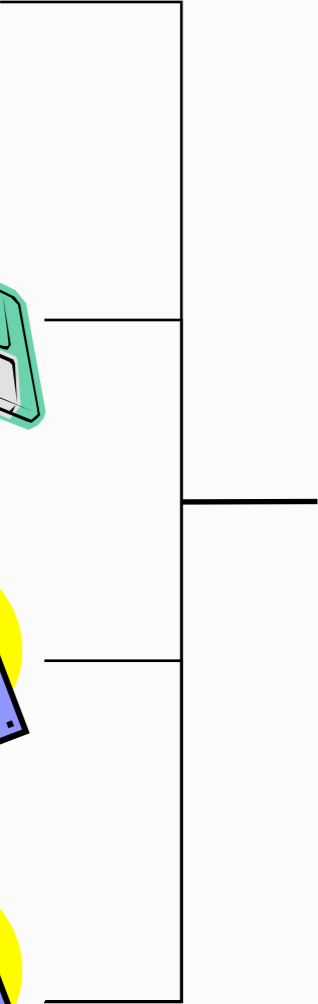
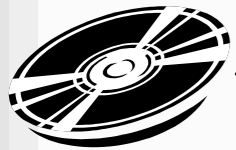
Offline Analysis Chain

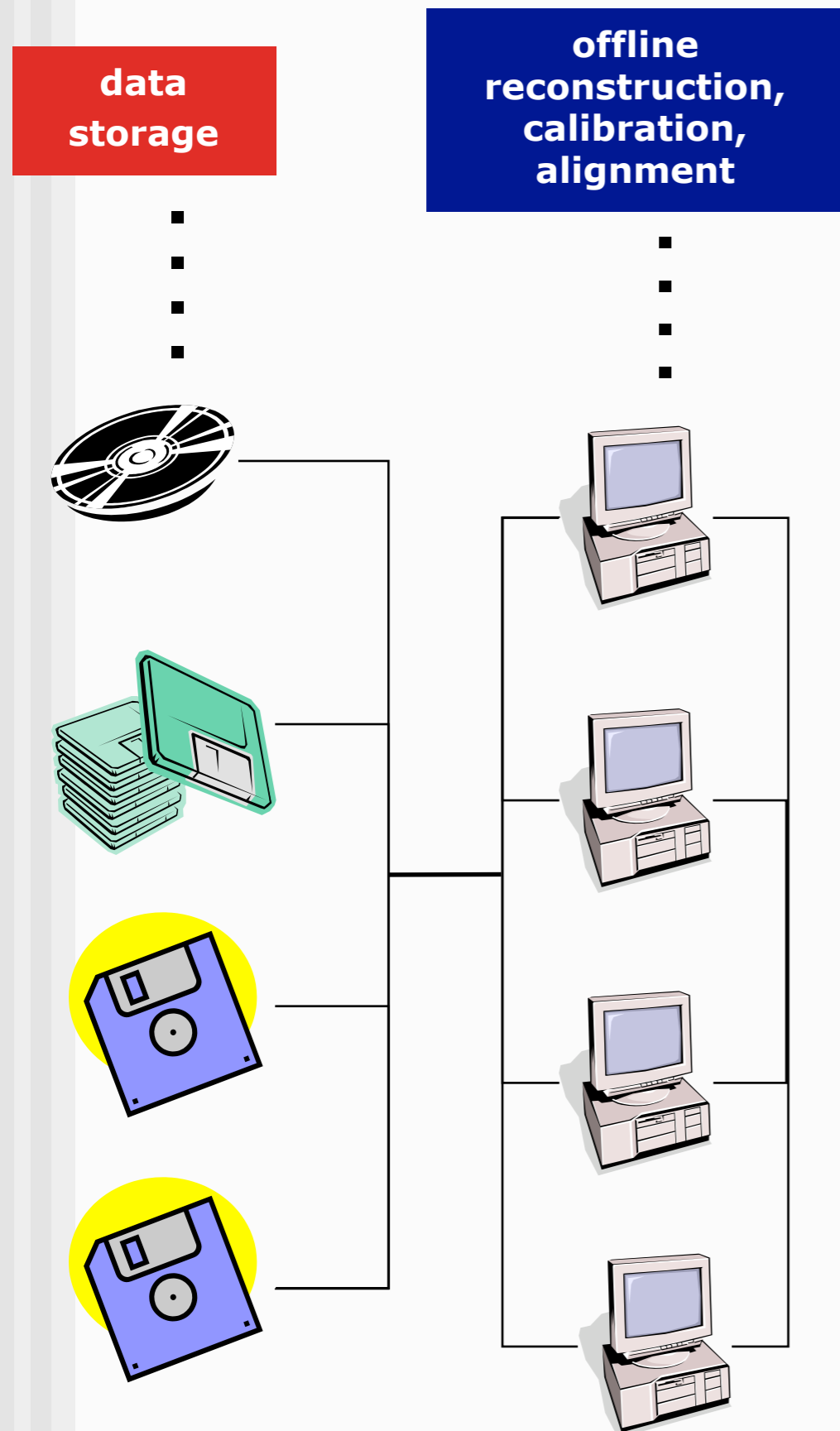
**data
storage**

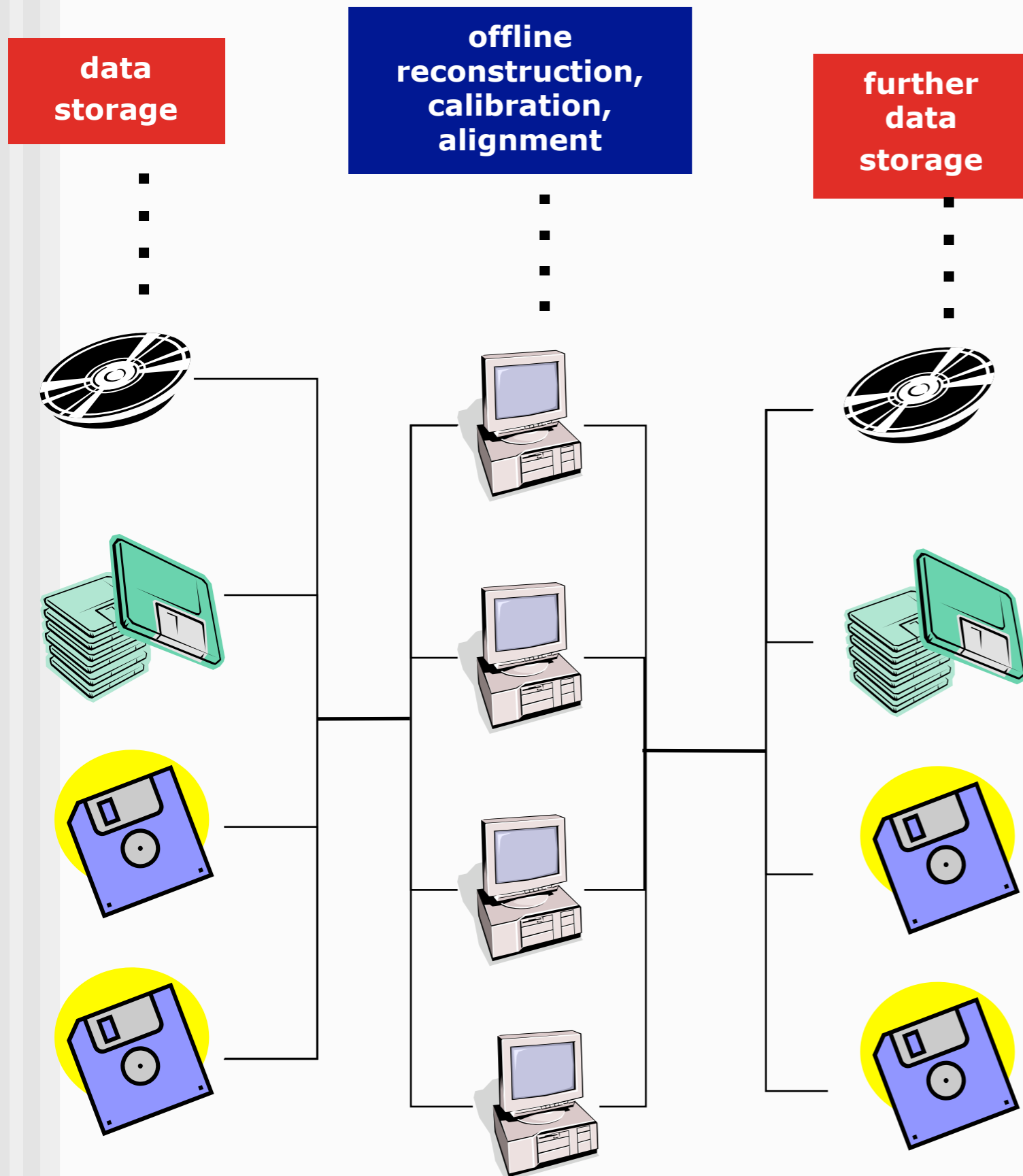


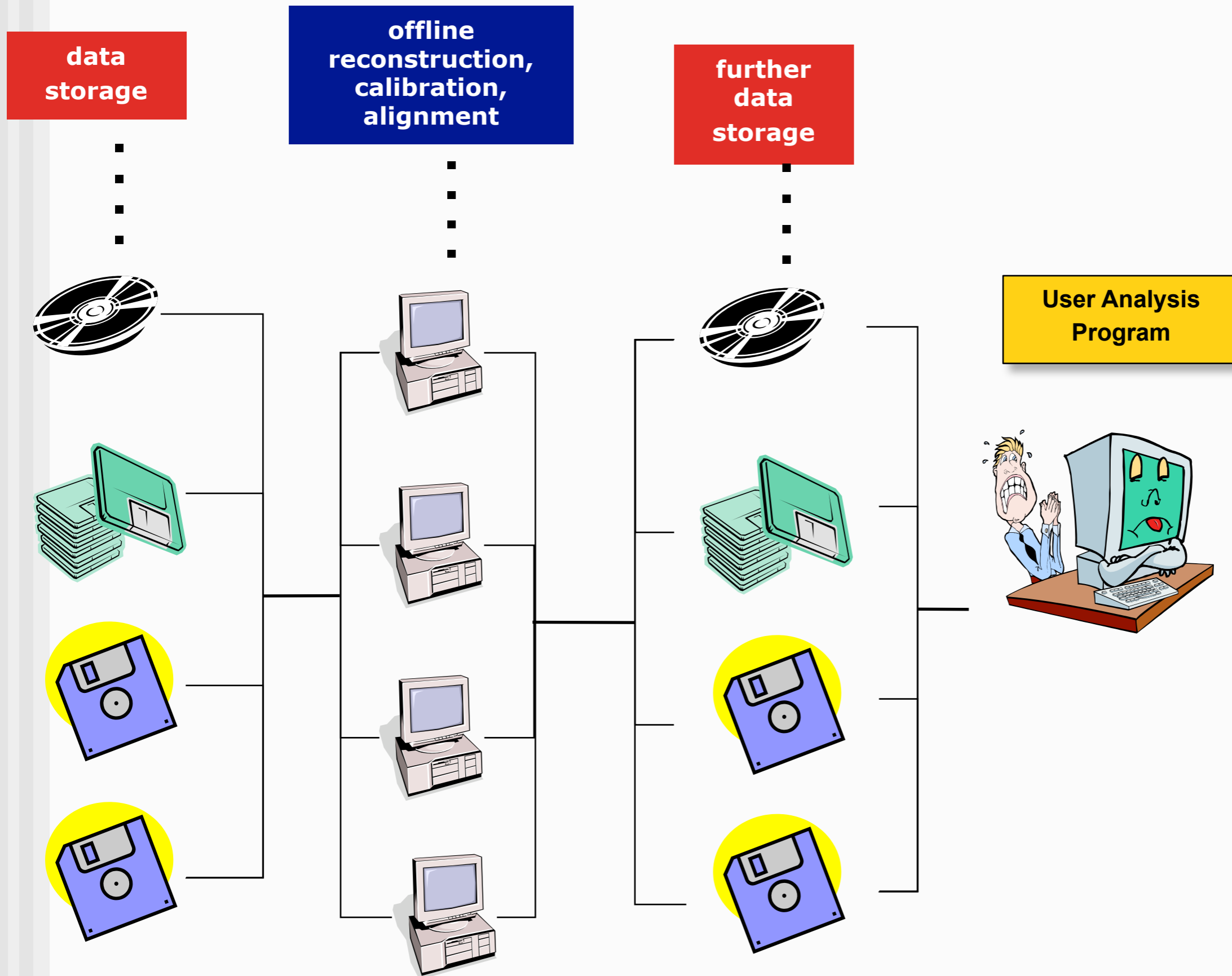
**data
storage**

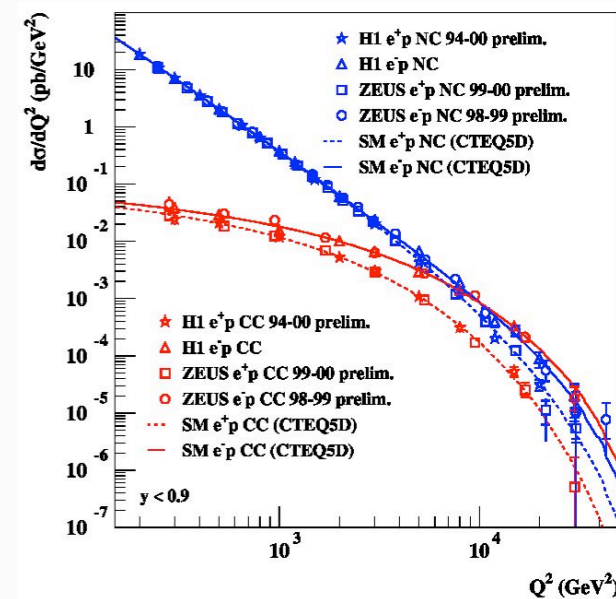
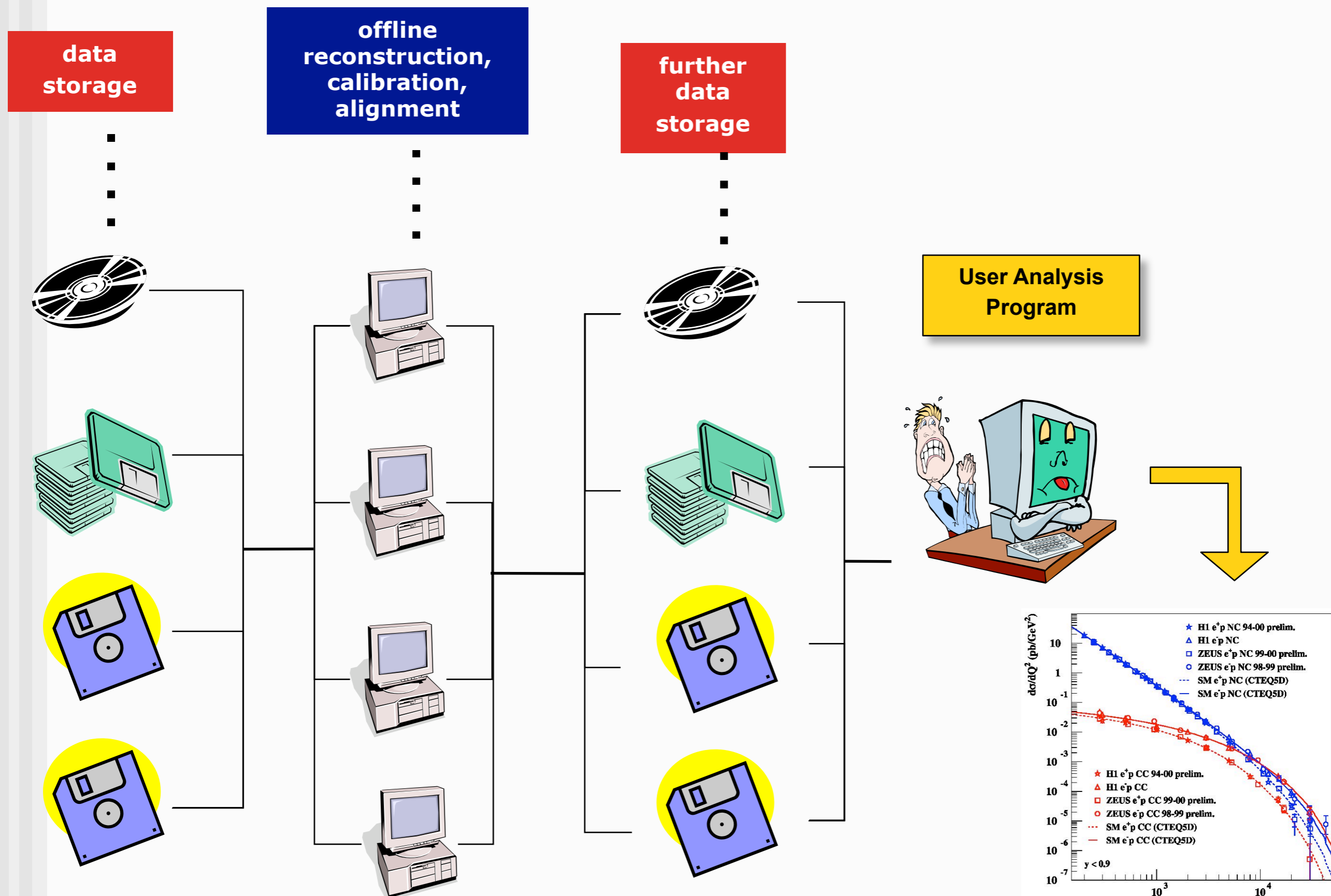
▪
▪
▪
▪

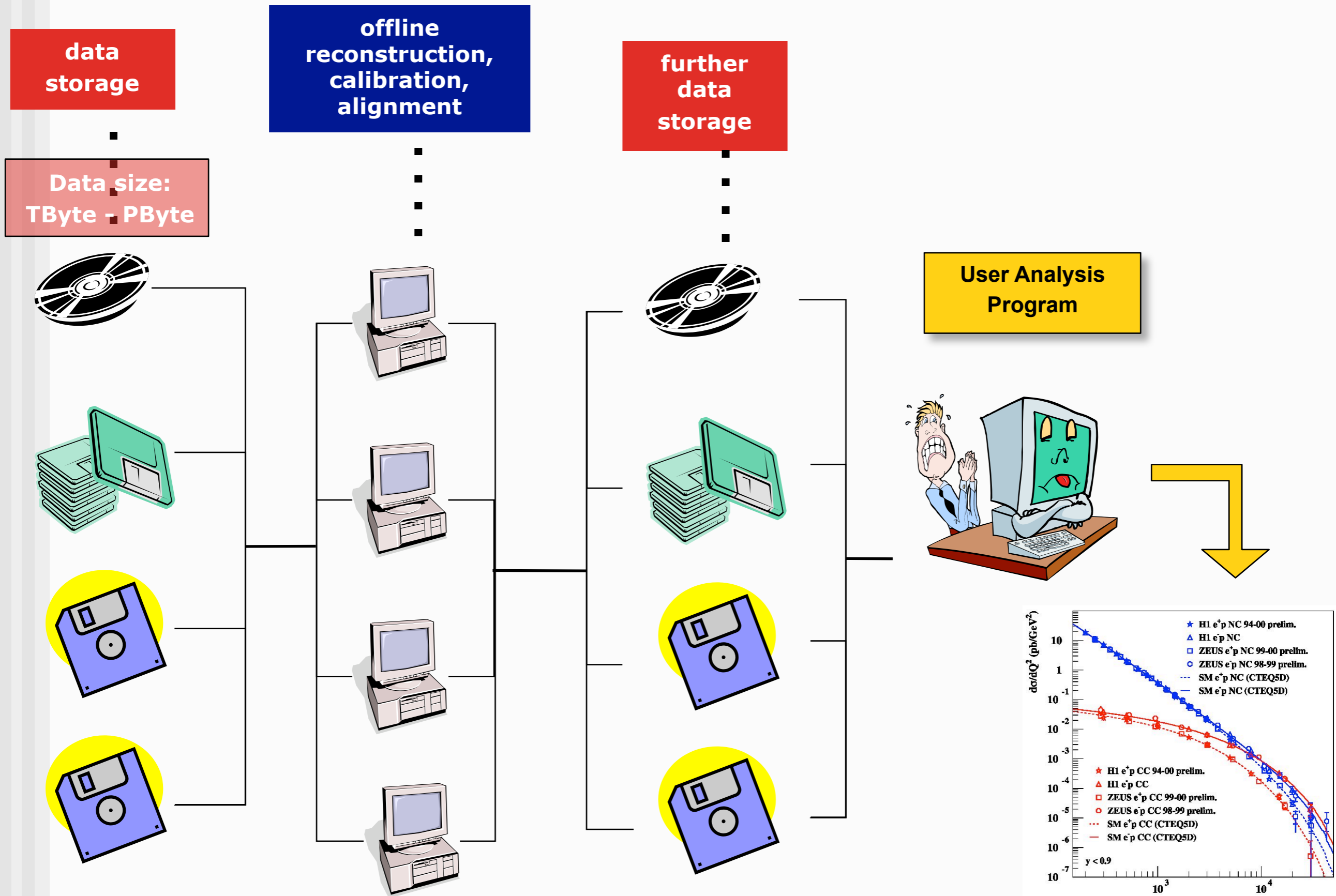


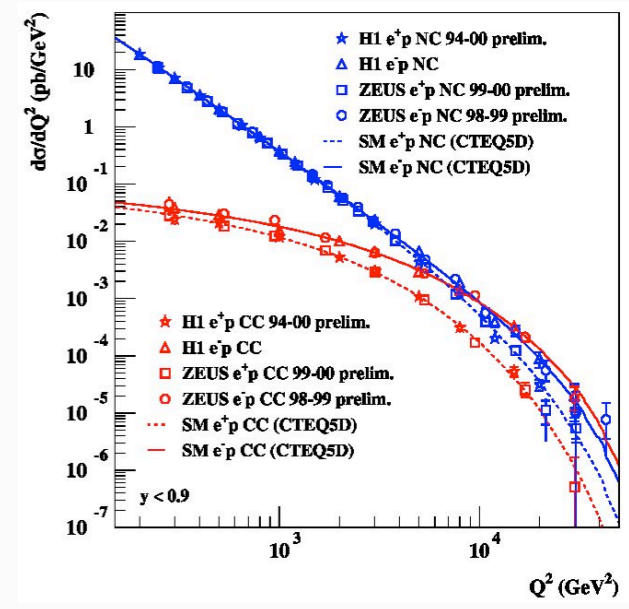
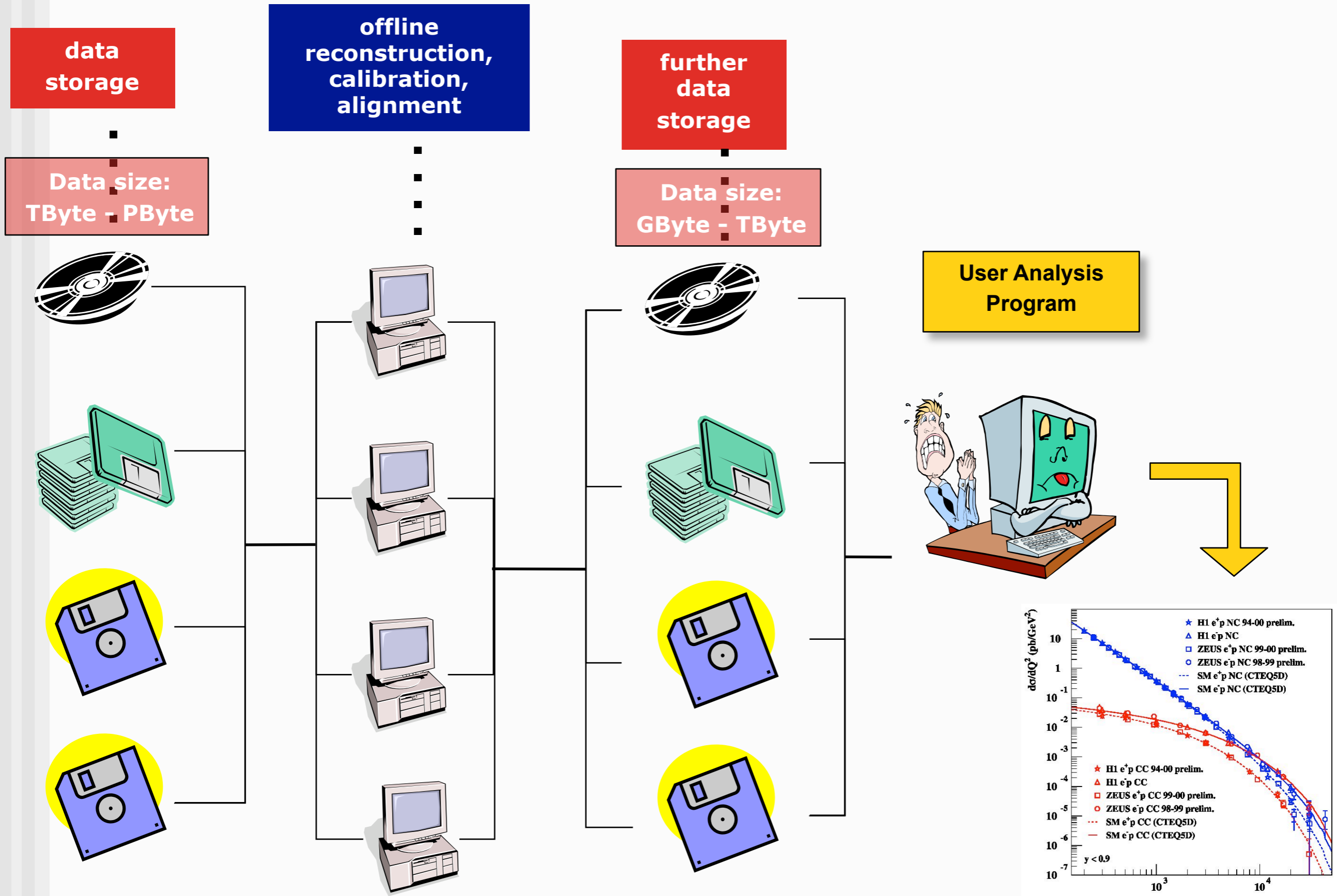


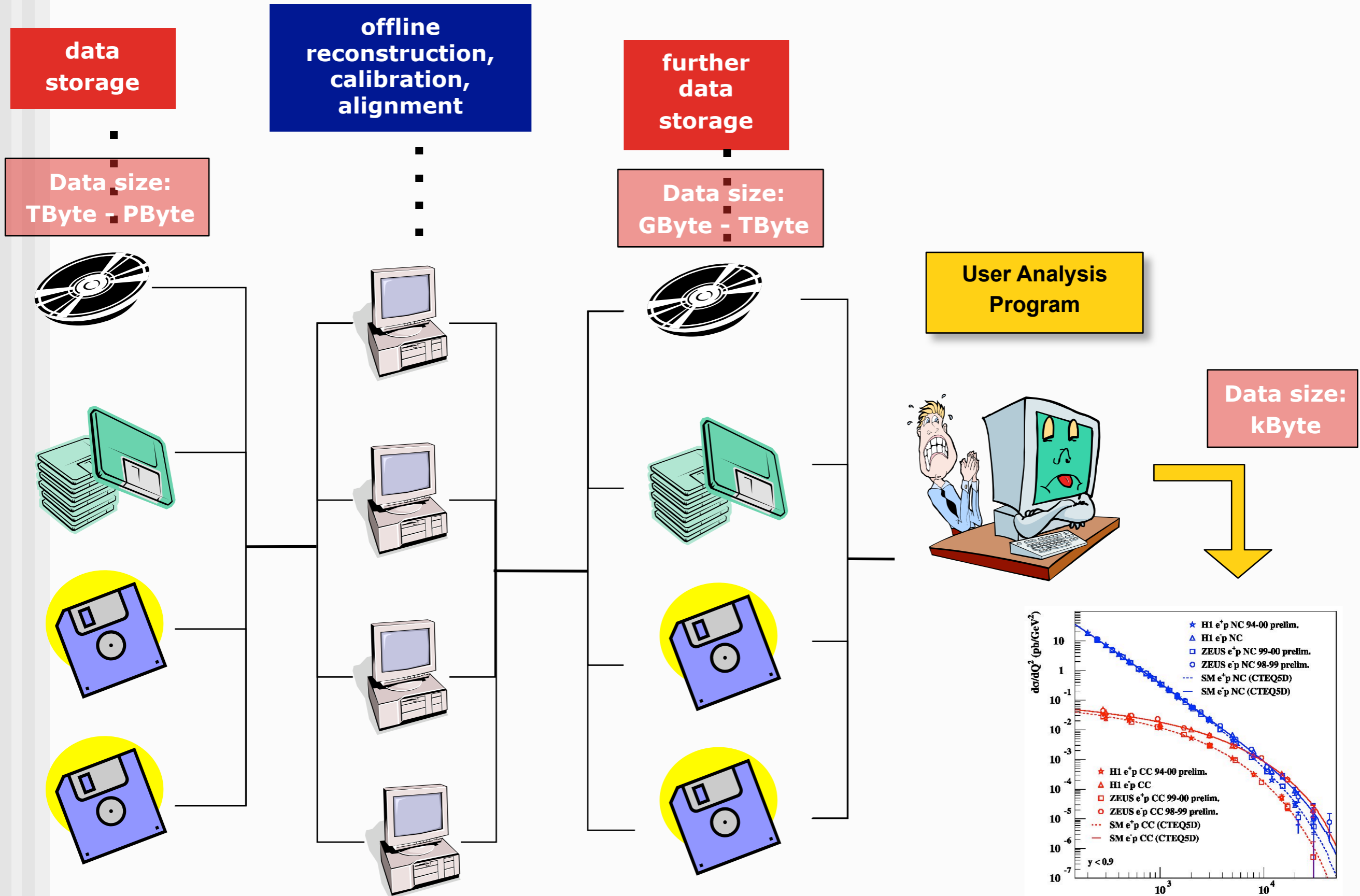








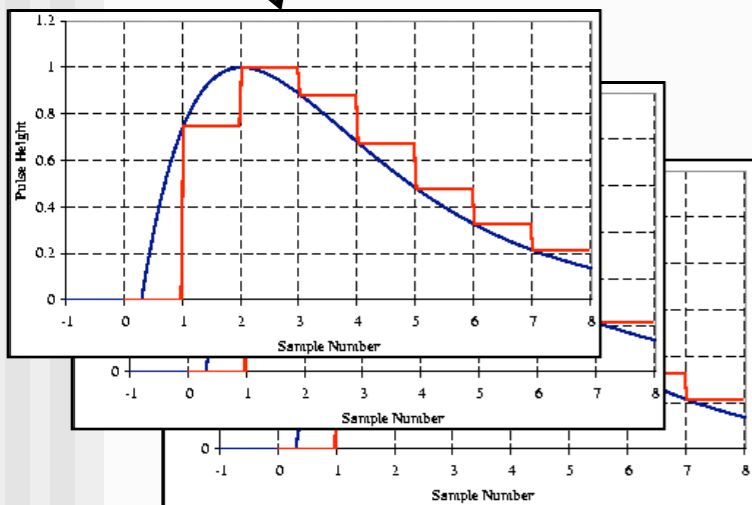
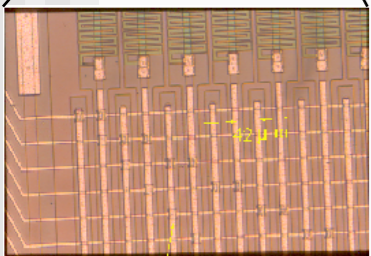
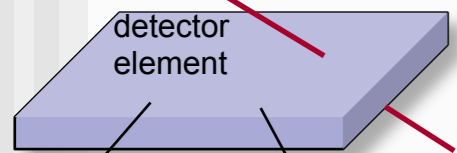






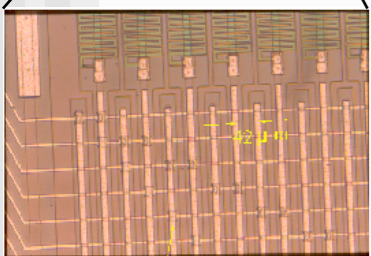
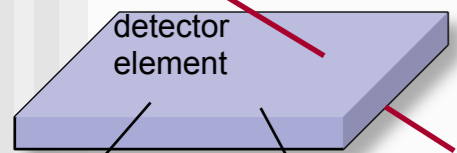
Data reduction/abstraction

particle

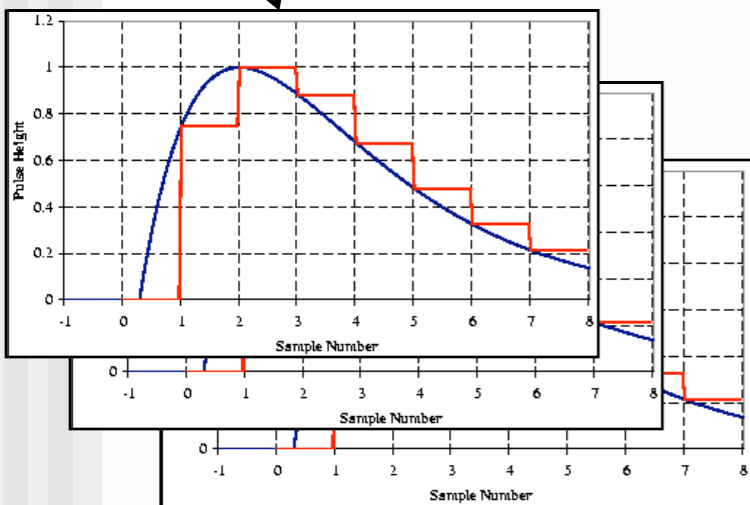


Analog signals

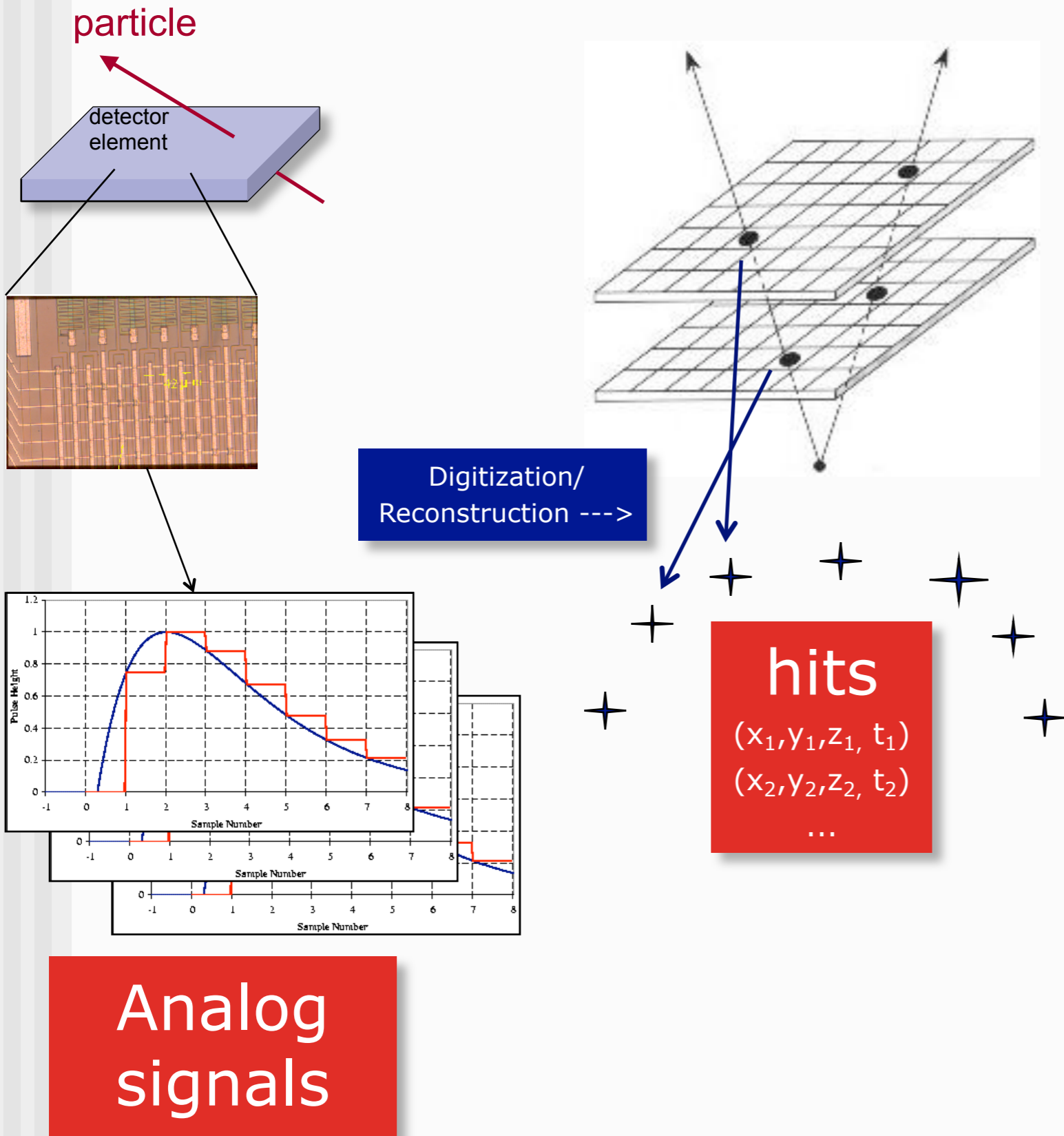
particle

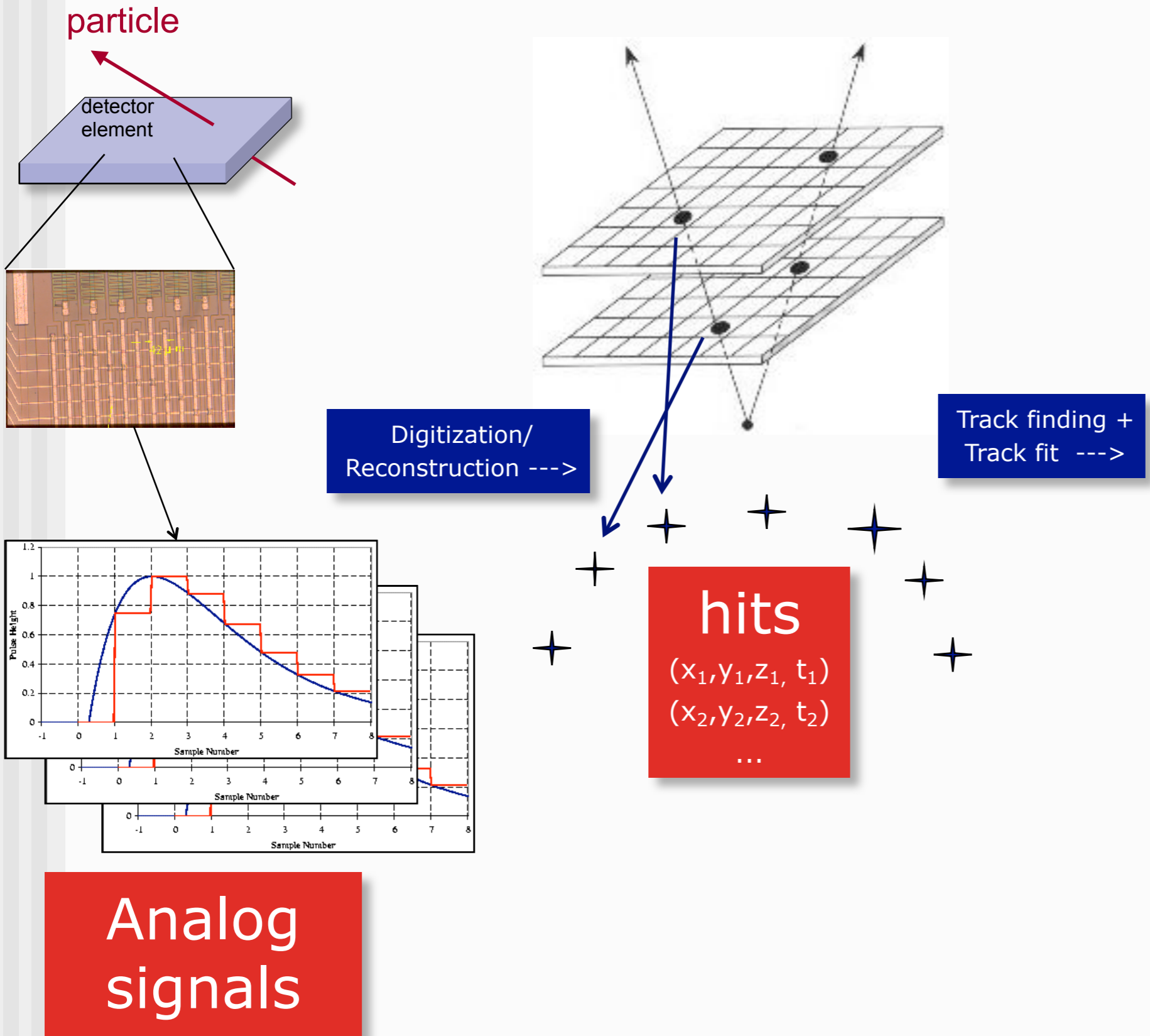


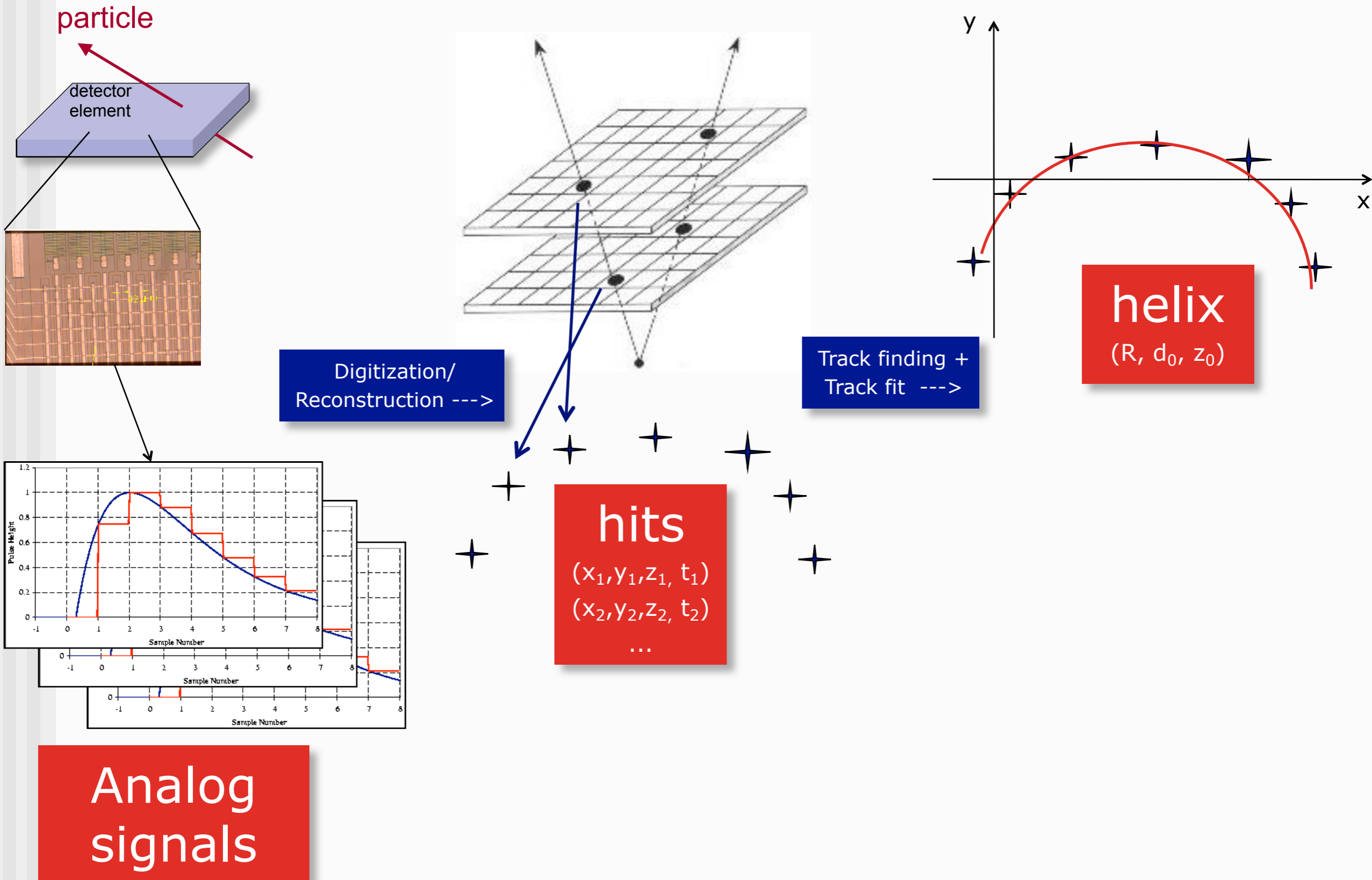
Digitization/
Reconstruction --->

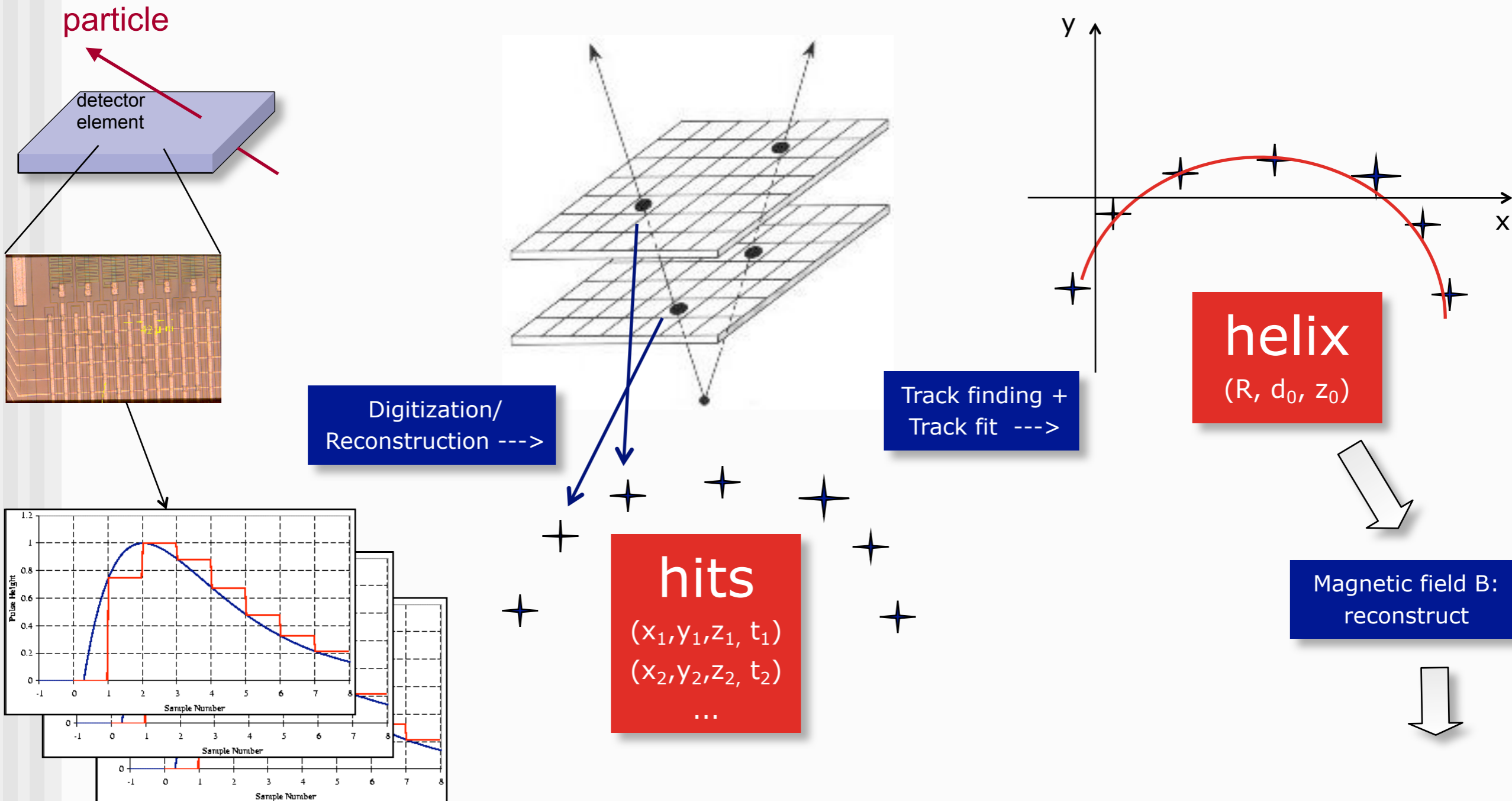


Analog
signals

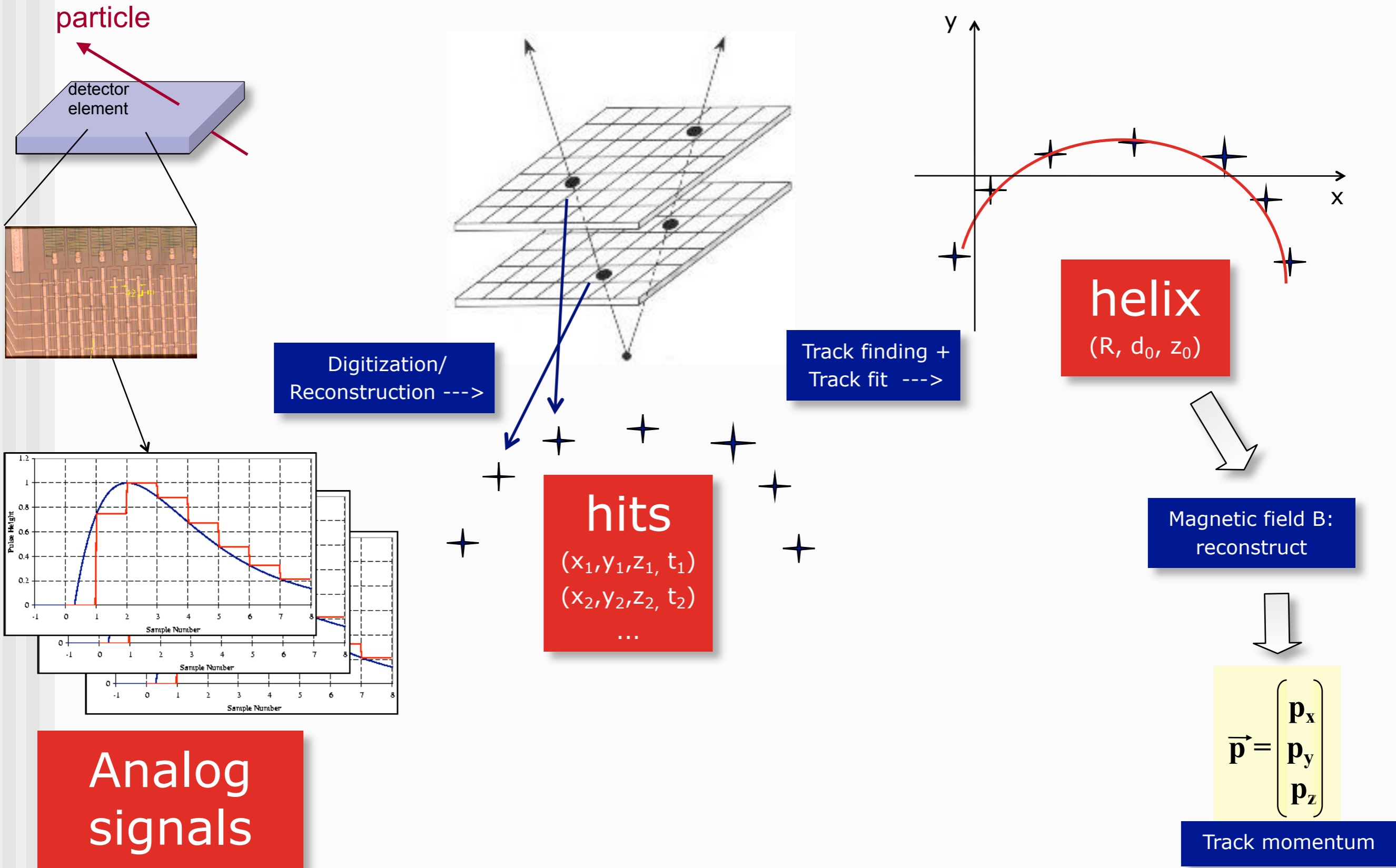


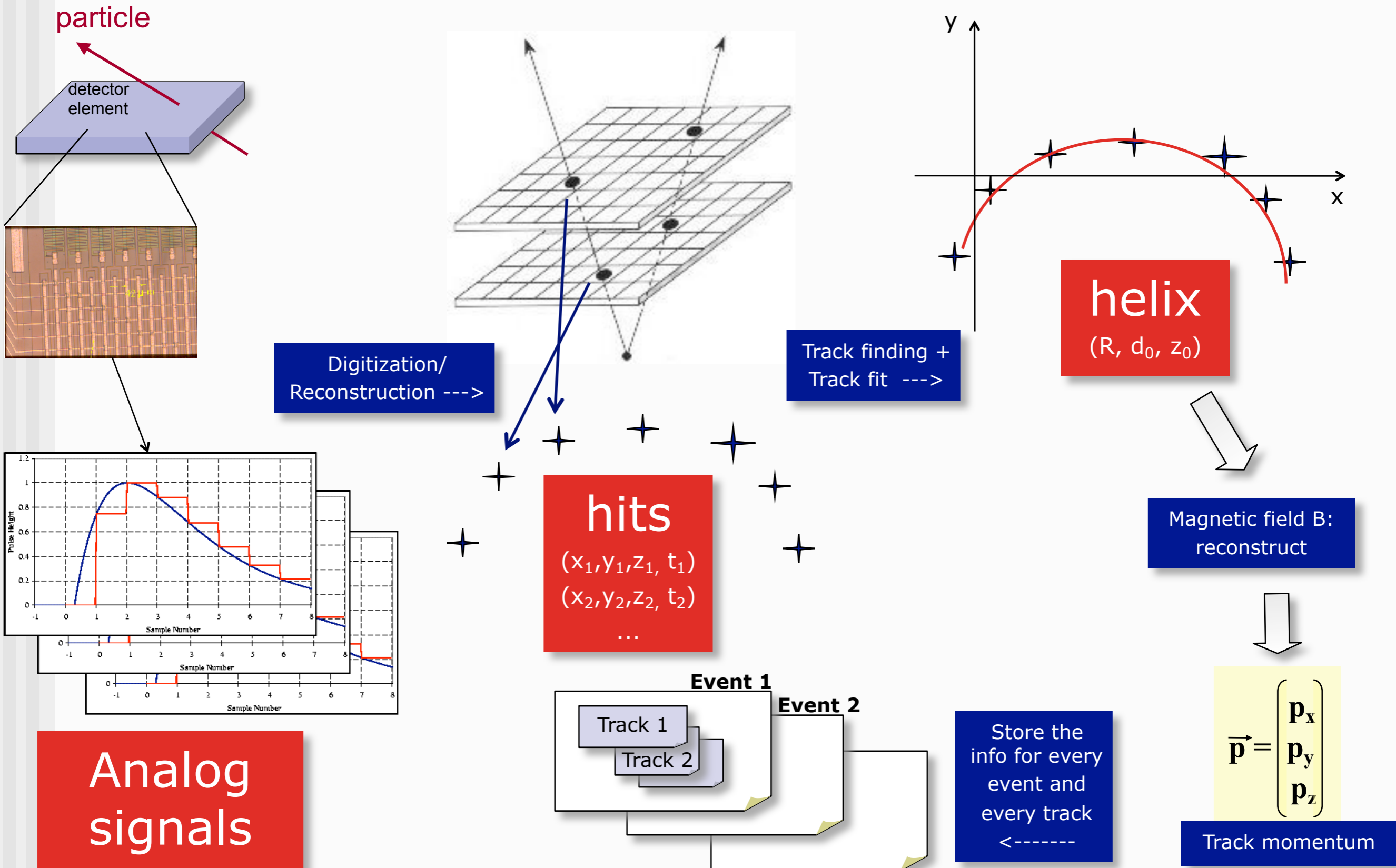






Analog signals





- Data are stored sequentially in files...

File A

- Data are stored sequentially in files...

Event 1

Nch (charged tracks) :
2

Pcha
(Momentum of each track):
{ {"-7.65698", "42.9725", "14.3404"},
 {" 7.54101", "-42.1729", "-14.0108"} }

px

py

pz

Qcha
(Charge of each track):
{ -1, 1 }

File A

- Data are stored sequentially in files...

Event 1

Nch (charged tracks) :
2

Pcha
(Momentum of each track):
{ {"-7.65698", "42.9725", "14.3404"},
 {" 7.54101", "-42.1729", "-14.0108"} }

px py pz

Qcha
(Charge of each track):
{ -1, 1 }

Event 2

Nch (charged tracks) :
3

Pcha
(Momentum of each track):
{ {"-12.9305", "12.2713", "40.5615"},
 {" 12.2469", "-11.606", "-38.7182"},
 {"0.143435", "-0.143435", "-0.497444"} }

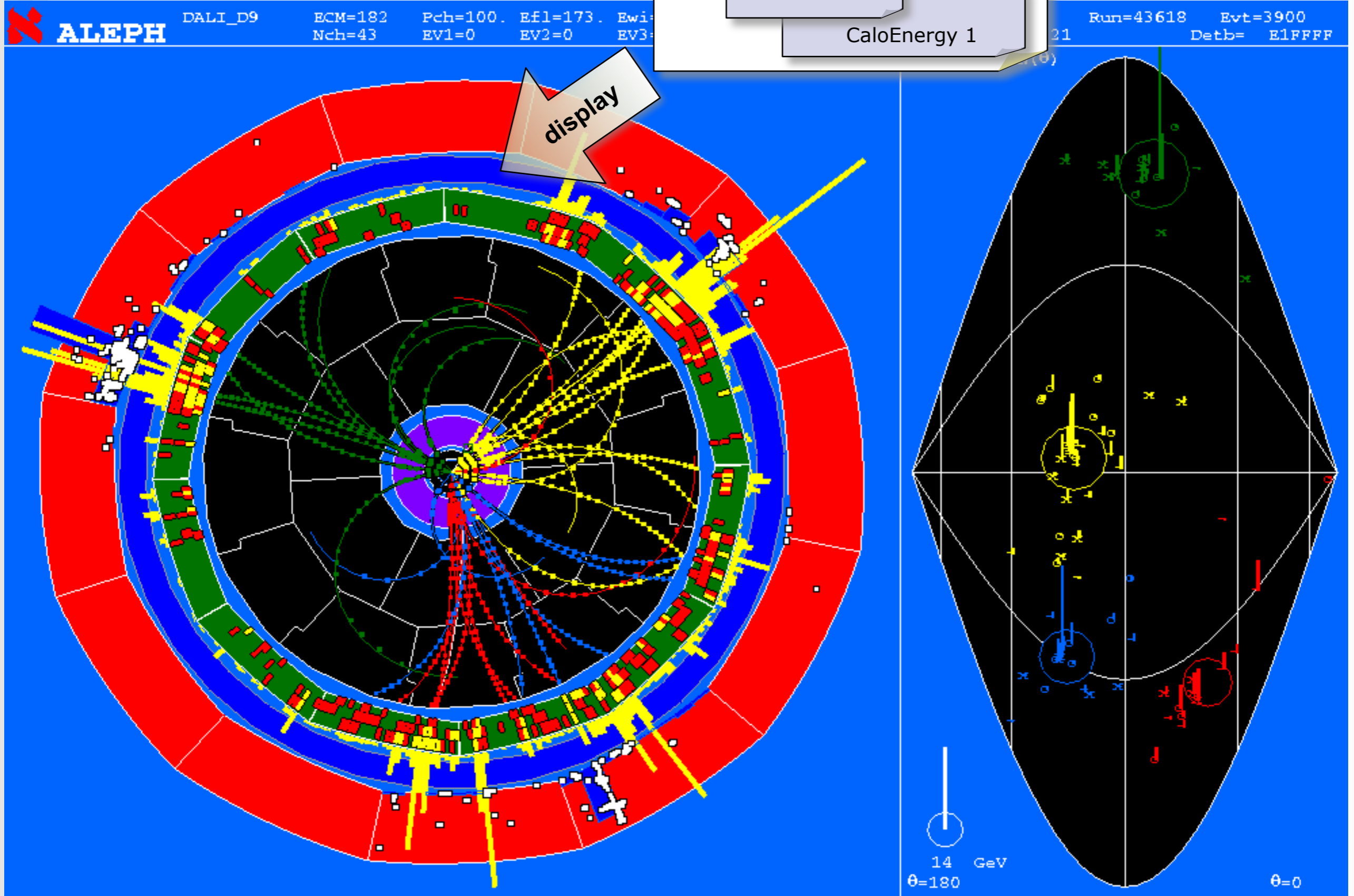
px py pz

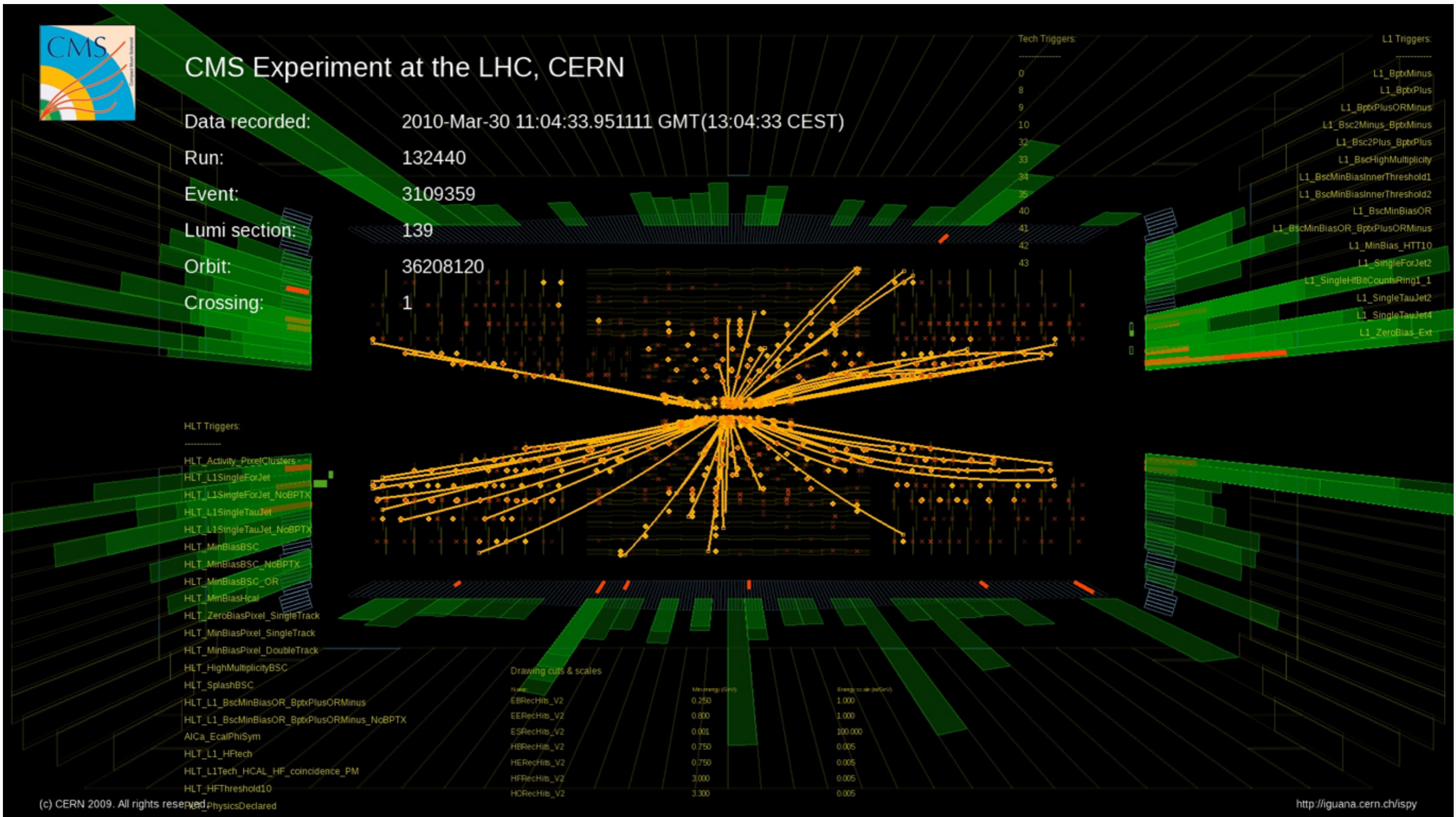
Qcha
(Charge of each track):
{ -1, 1, -1 }

File A

Event 1

Hits Track 1
Hits Track 2
CaloEnergy 1



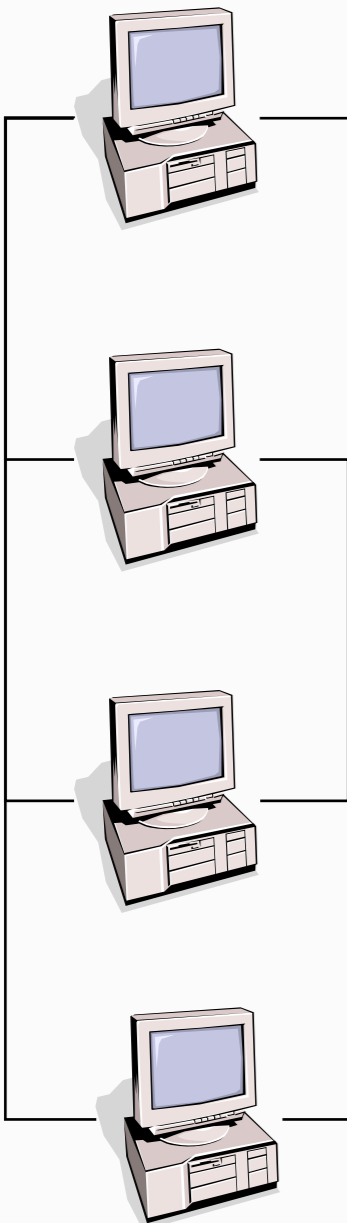




Simulation

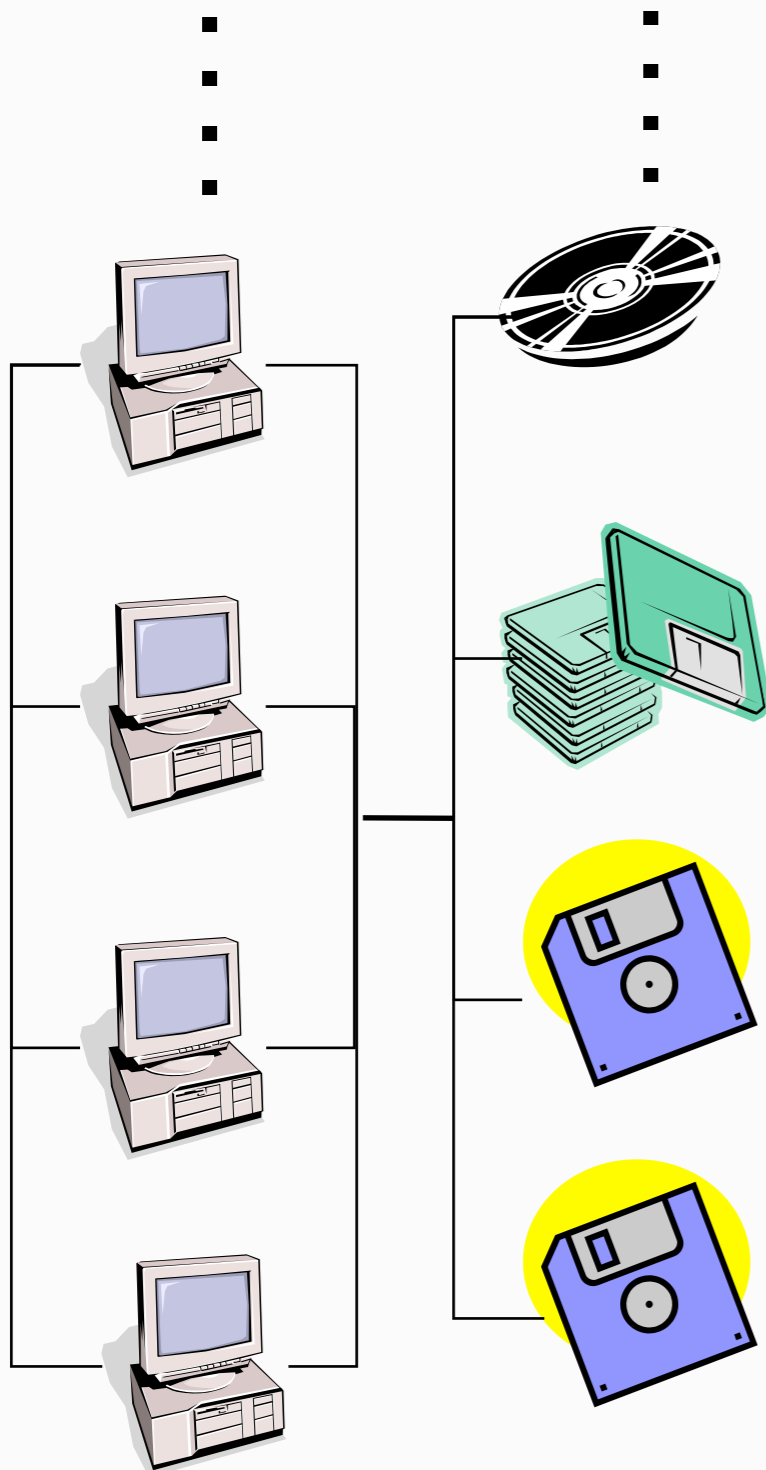
process and detector simulation

■
■
■
■



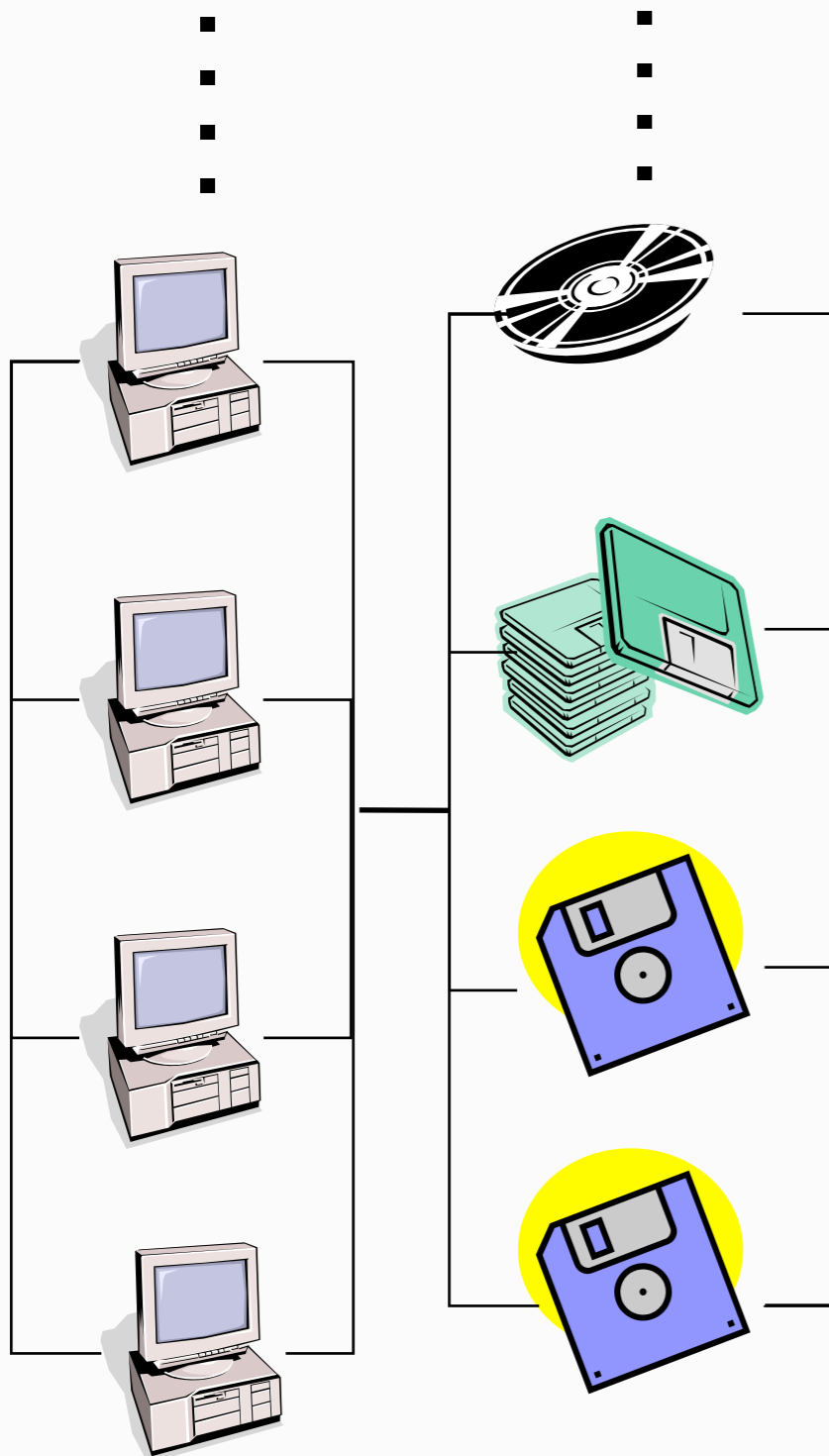
**process and
detector
simulation**

**data
storage**



process and
detector
simulation

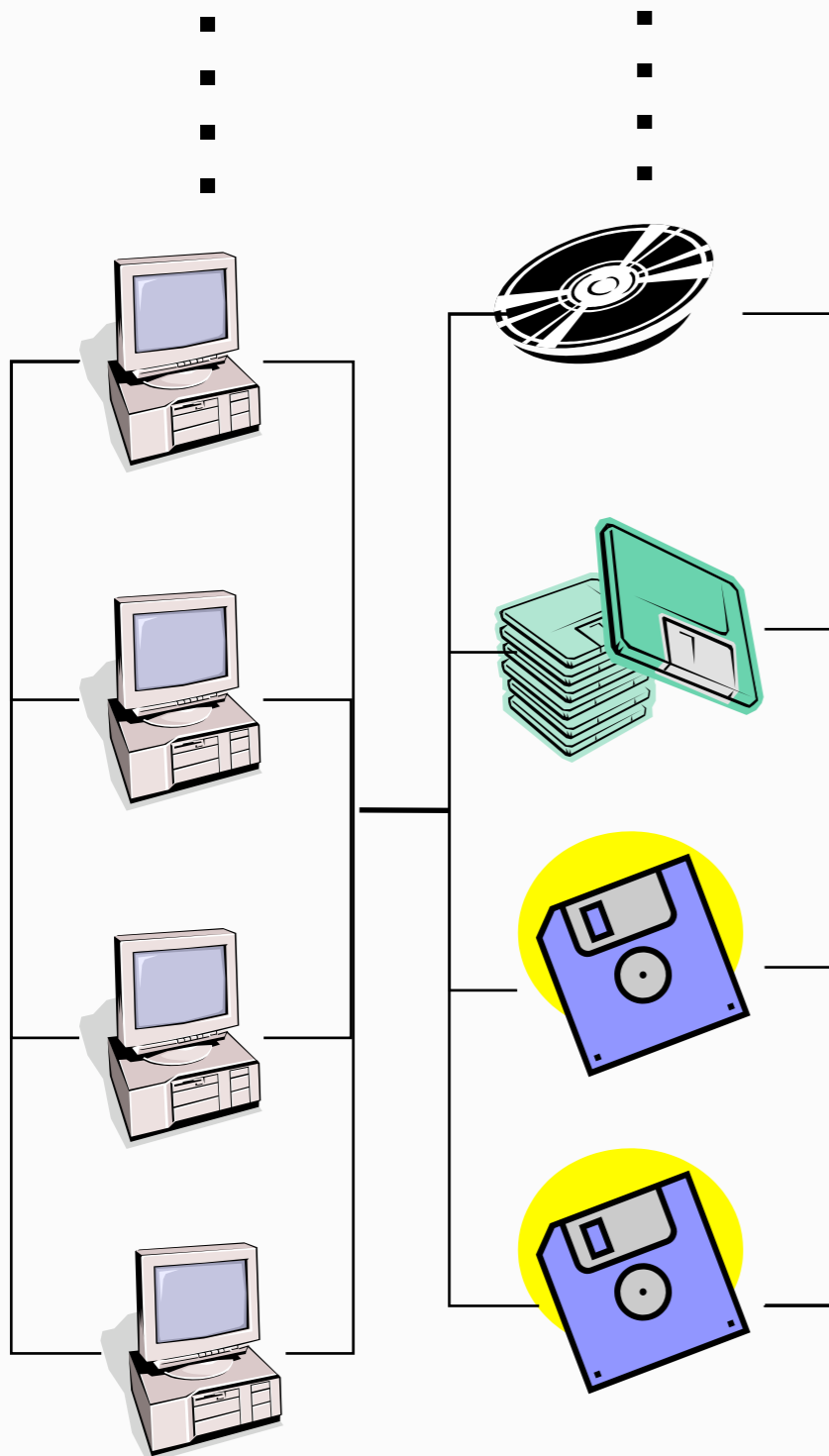
data
storage



Exactly
the same
steps as
for the
data

process and detector simulation

data storage

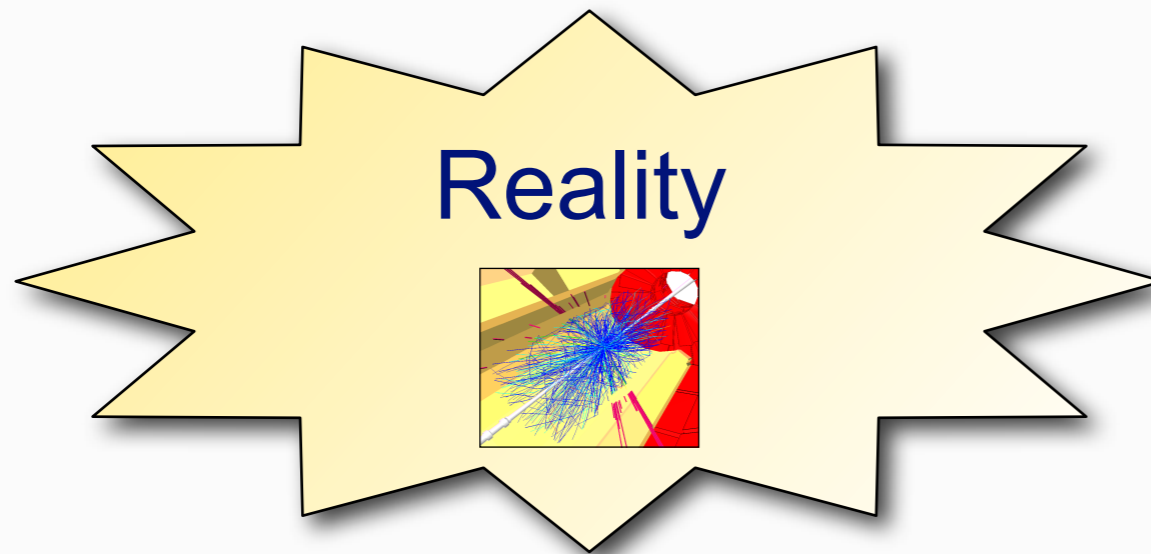


Exactly the same steps as for the data

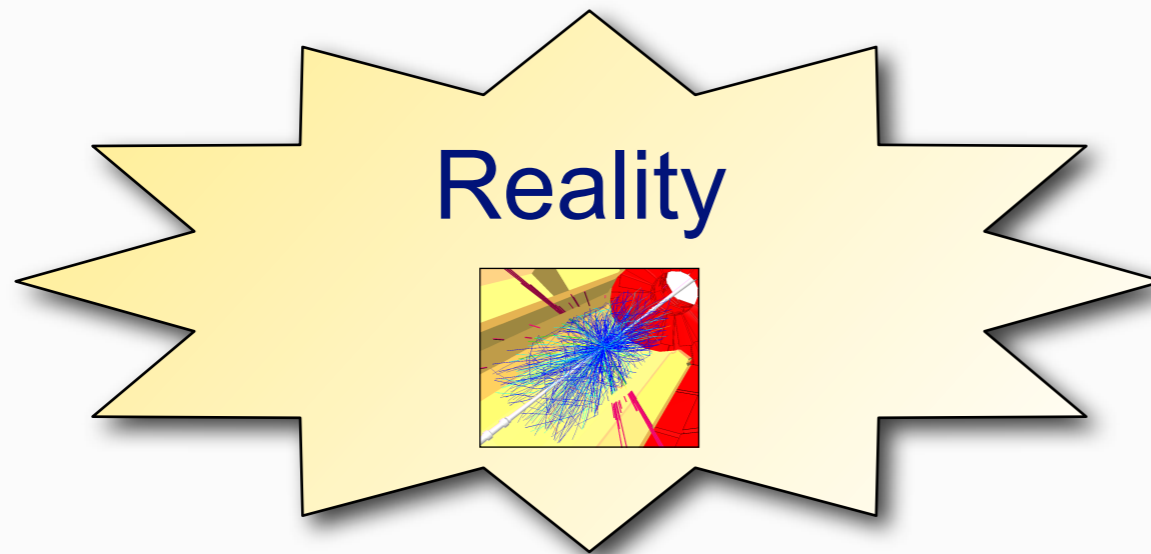
Simulation of many (millions) of events

- simulate physics process
e.g. $e^+e^- \rightarrow \text{hadrons}$
or $p p \rightarrow \text{jets}$
- plus the detector response to the produced particles
- understand detector response and analysis parameters (lost particles, resolution, efficiencies, backgrounds)
- and compare to real data
- Note : simulations present from beginning to end of experiment, needed to make design choices

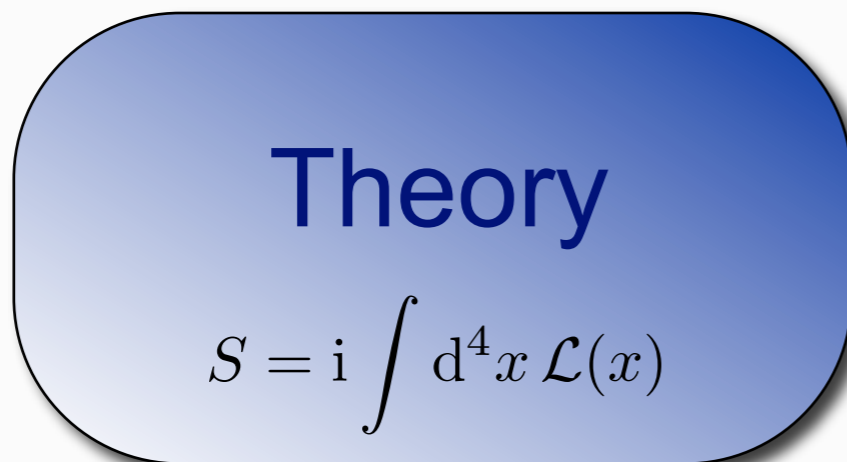
**And now let's go a
little bit more into the
details ...**



We use experiments
to inquire about what
“reality” (nature) does

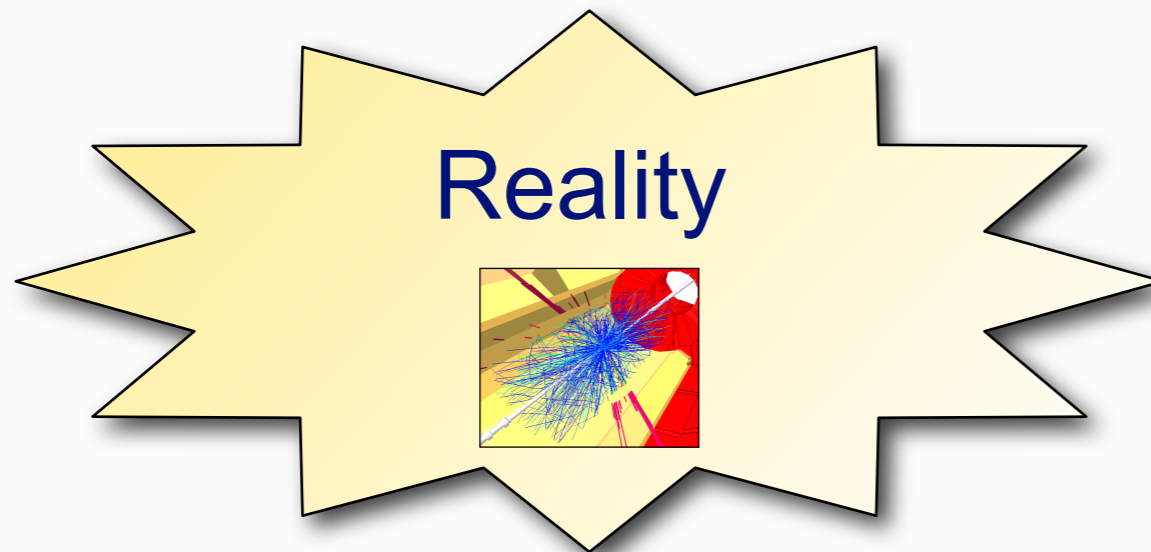


We use experiments to inquire about what “reality” (nature) does

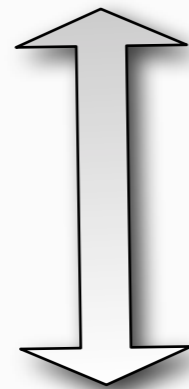


The goal is to understand in the most general; that’s usually also the simplest.

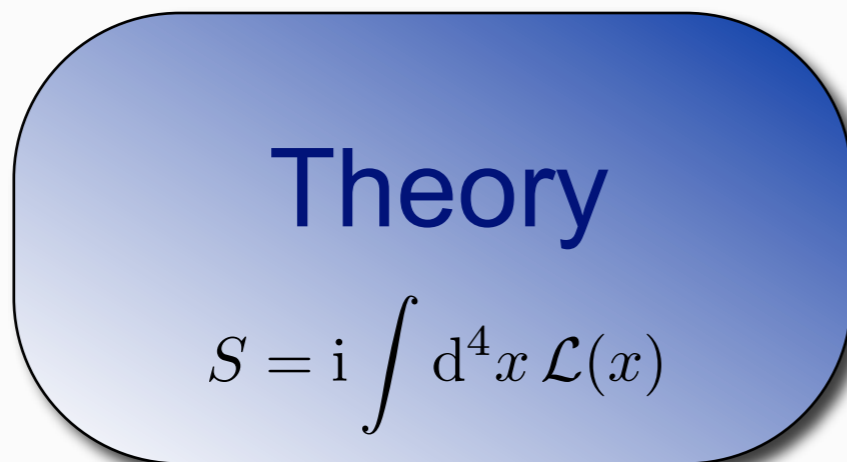
- A. Eddington



We use experiments to inquire about what “reality” (nature) does



We intend to fill this gap



The goal is to understand in the most general; that's usually also the simplest.

- A. Eddington

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} \mathbf{W}_{\mu\nu} \cdot \mathbf{W}^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ kinetic} \\ \text{energies and} \\ \text{self-interactions} \end{array} \right. \\
 & + \bar{L} \gamma^\mu \left(i \partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu \right) L \\
 & + \bar{R} \gamma^\mu \left(i \partial_\mu - g' \frac{Y}{2} B_\mu \right) R & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{kinetic energies} \\ \text{and their} \\ \text{interactions with} \\ W^\pm, Z, \gamma \end{array} \right. \\
 & + \left| \left(i \partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu \right) \phi \right|^2 & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ and} \\ \text{Higgs masses} \\ \text{and couplings} \end{array} \right. \\
 & - V(\phi) \\
 & - (G_1 \bar{L} \phi R + G_2 \bar{L} \phi_c R + h.c.) & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{masses and} \\ \text{coupling to Higgs} \end{array} \right.
 \end{aligned}$$

L ... left-handed fermion (l or q) doublet
 R ... right-handed fermion singlet

\mathcal{L} from QCD:

$$\mathcal{L} = \underbrace{\bar{q} (i \gamma^\mu \partial_\mu - m) q}_{E_{\text{kin}}(q)} - g \underbrace{(\bar{q} \gamma^\mu T_a q) G_\mu^a}_{\text{Interaction } q, g} - \frac{1}{4} \underbrace{G_{\mu\nu}^a G_a^{\mu\nu}}_{\substack{E_{\text{kin}}(g) \\ \text{includes} \\ \text{self-interaction} \\ \text{between gluons}}}$$

eg.
the Standard Model

$$\begin{aligned}
 \mathcal{L} = & -\frac{1}{4} \mathbf{W}_{\mu\nu} \cdot \mathbf{W}^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ kinetic} \\ \text{energies and} \\ \text{self-interactions} \end{array} \right. \\
 & + \bar{L} \gamma^\mu \left(i \partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu \right) L \\
 & + \bar{R} \gamma^\mu \left(i \partial_\mu - g' \frac{Y}{2} B_\mu \right) R & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{kinetic energies} \\ \text{and their} \\ \text{interactions with} \\ W^\pm, Z, \gamma \end{array} \right. \\
 & + \left| \left(i \partial_\mu - g \frac{1}{2} \boldsymbol{\tau} \cdot \mathbf{W}_\mu - g' \frac{Y}{2} B_\mu \right) \phi \right|^2 & \left\{ \begin{array}{l} W^\pm, Z, \gamma \text{ and} \\ \text{Higgs masses} \\ \text{and couplings} \end{array} \right. \\
 & - V(\phi) \\
 & - (G_1 \bar{L} \phi R + G_2 \bar{L} \phi_c R + h.c.) & \left\{ \begin{array}{l} \text{lepton and quark} \\ \text{masses and} \\ \text{coupling to Higgs} \end{array} \right.
 \end{aligned}$$

L ... left-handed fermion (l or q) doublet
 R ... right-handed fermion singlet

\mathcal{L} from QCD:

$$\mathcal{L} = \bar{q} \underbrace{\left(i \gamma^\mu \partial_\mu - m \right)}_{E_{\text{kin}}(q)} q - g \underbrace{\left(\bar{q} \gamma^\mu T_a q \right) G_\mu^a}_{\text{Interaction } q, g} - \frac{1}{4} \underbrace{G_{\mu\nu}^a G_a^{\mu\nu}}_{E_{\text{kin}}(g)}$$

$E_{\text{kin}}(g)$ includes self-interaction between gluons

eg. the Standard Model

has parameters

coupling constants

masses

predicts:
cross sections,
branching ratios,
lifetimes, ...

```
0x01e84c10: 0x01e8 0x8848 0x01e8 0x83d8 0x6c73 0x6f72 0x7400 0x0000
0x01e84c20: 0x0000 0x0019 0x0000 0x0000 0x01e8 0x4d08 0x01e8 0x5b7c
0x01e84c30: 0x01e8 0x87e8 0x01e8 0x8458 0x7061 0x636b 0x6167 0x6500
0x01e84c40: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c50: 0x01e8 0x8788 0x01e8 0x8498 0x7072 0x6f63 0x0000 0x0000
0x01e84c60: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c70: 0x01e8 0x8824 0x01e8 0x84d8 0x7265 0x6765 0x7870 0x0000
0x01e84c80: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c90: 0x01e8 0x8838 0x01e8 0x8518 0x7265 0x6773 0x7562 0x0000
0x01e84ca0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cb0: 0x01e8 0x8818 0x01e8 0x8558 0x7265 0x6e61 0x6d65 0x0000
0x01e84cc0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cd0: 0x01e8 0x8798 0x01e8 0x8598 0x7265 0x7475 0x726e 0x0000
0x01e84ce0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cf0: 0x01e8 0x87ec 0x01e8 0x85d8 0x7363 0x616e 0x0000 0x0000
0x01e84d00: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d10: 0x01e8 0x87e8 0x01e8 0x8618 0x7365 0x7400 0x0000 0x0000
0x01e84d20: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d30: 0x01e8 0x87a8 0x01e8 0x8658 0x7370 0x6c69 0x7400 0x0000
0x01e84d40: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d50: 0x01e8 0x8854 0x01e8 0x8698 0x7374 0x7269 0x6e67 0x0000
0x01e84d60: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d70: 0x01e8 0x875c 0x01e8 0x86d8 0x7375 0x6273 0x7400 0x0000
0x01e84d80: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d90: 0x01e8 0x87c0 0x01e8 0x8718 0x7377 0x6974 0x6368 0x0000
```

eg.
**1/30th of an event in
the BaBar detector**

 get about 100 evts/sec


```

0x01e84c10: 0x01e8 0x8848 0x01e8 0x83d8 0x6c73 0x6f72 0x7400 0x0000
0x01e84c20: 0x0000 0x0019 0x0000 0x0000 0x01e8 0x4d08 0x01e8 0x5b7c
0x01e84c30: 0x01e8 0x87e8 0x01e8 0x8458 0x7061 0x636b 0x6167 0x6500
0x01e84c40: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c50: 0x01e8 0x8788 0x01e8 0x8498 0x7072 0x6f63 0x0000 0x0000
0x01e84c60: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c70: 0x01e8 0x8824 0x01e8 0x84d8 0x7265 0x6765 0x7870 0x0000
0x01e84c80: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84c90: 0x01e8 0x8838 0x01e8 0x8518 0x7265 0x6773 0x7562 0x0000
0x01e84ca0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cb0: 0x01e8 0x8818 0x01e8 0x8558 0x7265 0x6e61 0x6d65 0x0000
0x01e84cc0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cd0: 0x01e8 0x8798 0x01e8 0x8598 0x7265 0x7475 0x726e 0x0000
0x01e84ce0: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84cf0: 0x01e8 0x87ec 0x01e8 0x85d8 0x7363 0x616e 0x0000 0x0000
0x01e84d00: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d10: 0x01e8 0x87e8 0x01e8 0x8618 0x7365 0x7400 0x0000 0x0000
0x01e84d20: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d30: 0x01e8 0x87a8 0x01e8 0x8658 0x7370 0x6c69 0x7400 0x0000
0x01e84d40: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d50: 0x01e8 0x8854 0x01e8 0x8698 0x7374 0x7269 0x6e67 0x0000
0x01e84d60: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d70: 0x01e8 0x875c 0x01e8 0x86d8 0x7375 0x6273 0x7400 0x0000
0x01e84d80: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d90: 0x01e8 0x87c0 0x01e8 0x8718 0x7377 0x6974 0x6368 0x0000

```

eg.
**1/30th of an event in
the BaBar detector**

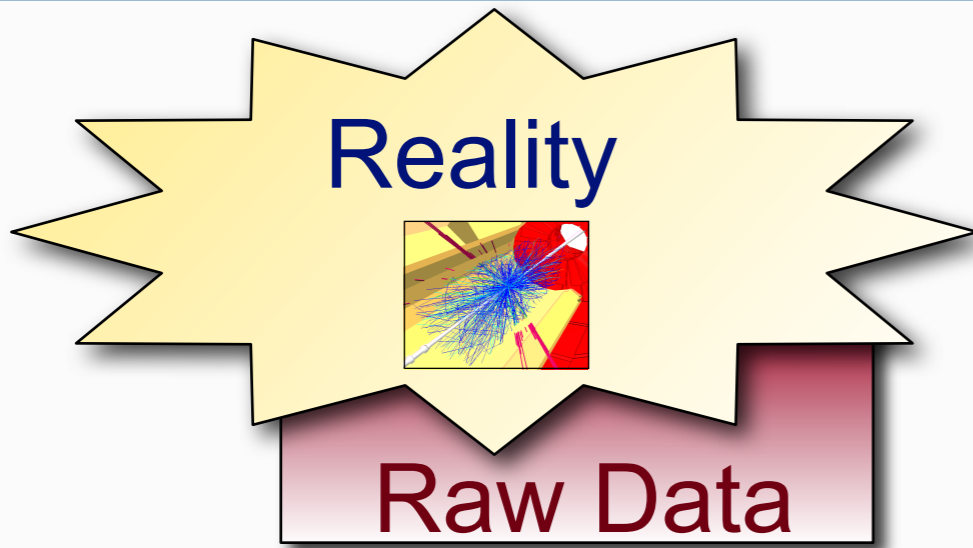
get about 100 evts/sec

“Address” :

which detector element
took the reading

“Value(s)” :

what the electronics
wrote out

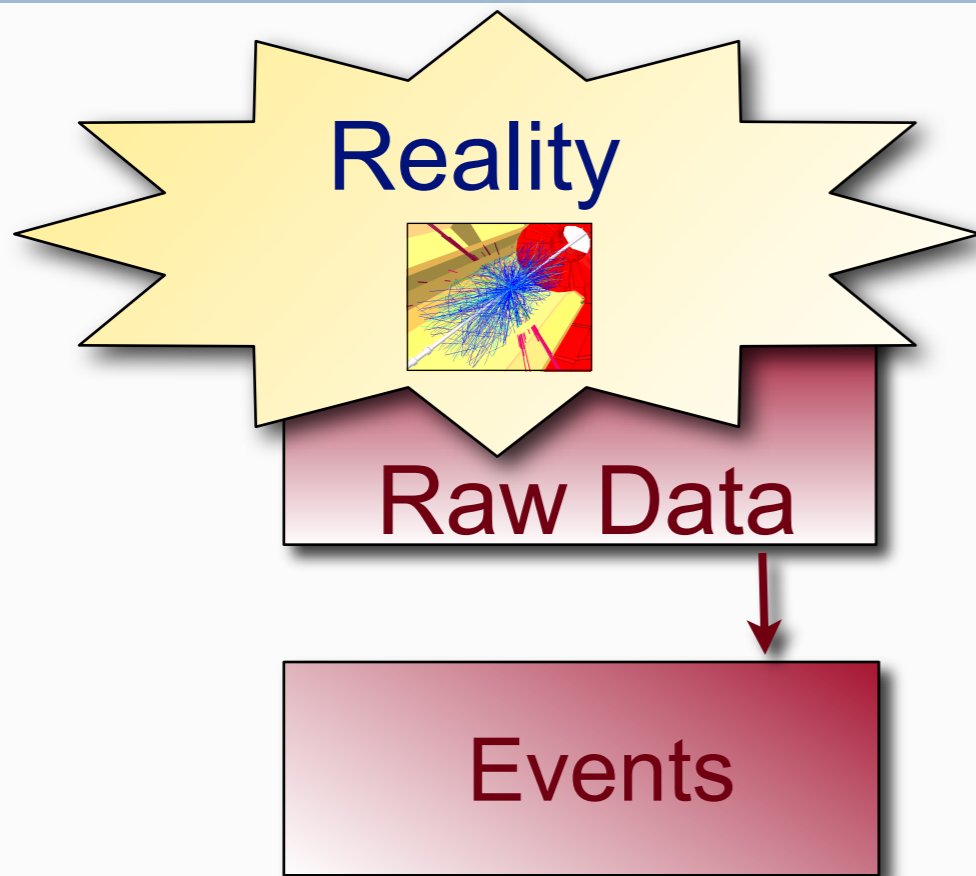


The imperfect measurement of a (set of) interactions in the detector

Theory

$$S = i \int d^4x \mathcal{L}(x)$$

A small number of general equations, with some parameters (poorly or not known at all)



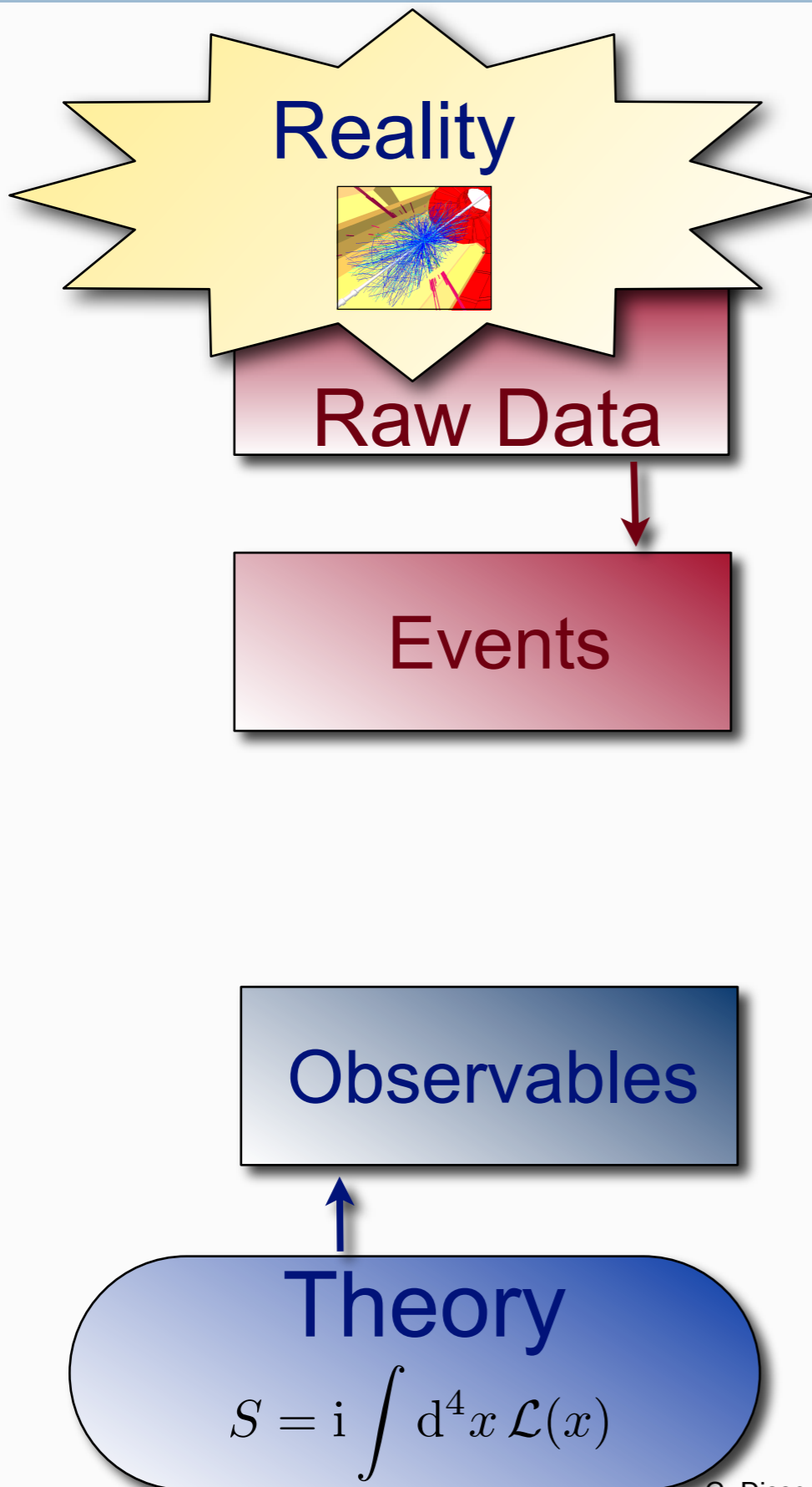
The imperfect measurement of a (set of) interactions in the detector

A unique happening:
 eg. Run 23458, event 1345
 which contains a $Z \rightarrow \mu^+ \mu^-$ decay

Theory

$$S = i \int d^4x \mathcal{L}(x)$$

A small number of general equations, with some parameters (poorly or not known at all)

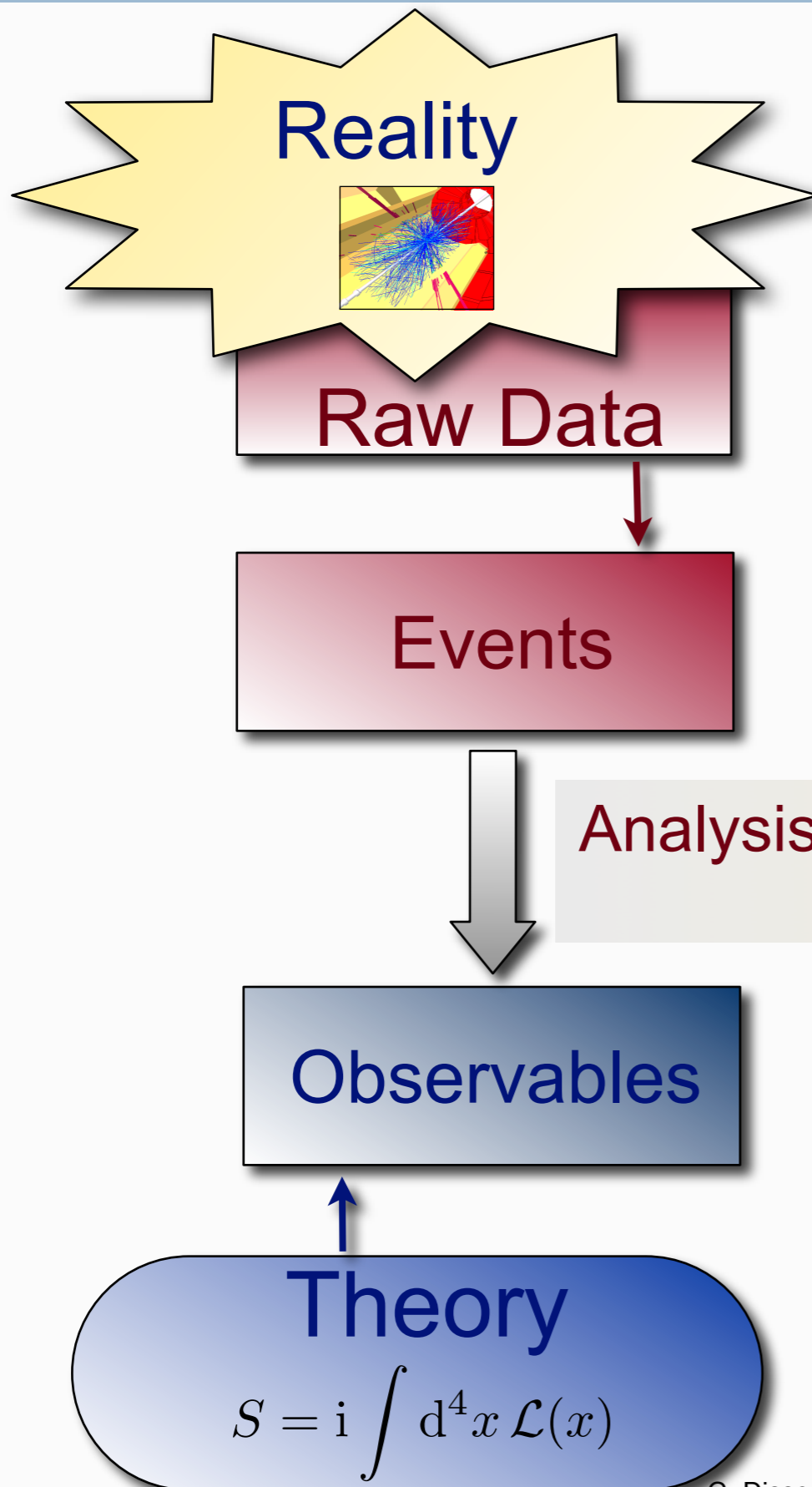


The imperfect measurement of a (set of) interactions in the detector

A unique happening:
eg. Run 23458, event 1345
which contains a $Z \rightarrow \mu^+ \mu^-$ decay

cross sections (probabilities for interactions),
branching ratios (BR), ratios of BRs, specific
lifetimes, ...

A small number of general equations, with
some parameters (poorly or not known at all)



The imperfect measurement of a (set of) interactions in the detector

A unique happening:
eg. Run 23458, event 1345
which contains a $Z \rightarrow \mu^+ \mu^-$ decay

Analysis : We “confront theory with experiment” by comparing the measured quantity (observable) with the prediction.

cross sections (probabilities for interactions),
branching ratios (BR), ratios of BRs, specific
lifetimes, ...

A small number of general equations, with
some parameters (poorly or not known at all)

A “simple” example

- Measurement of e^+e^- annihilation into hadrons and muons:

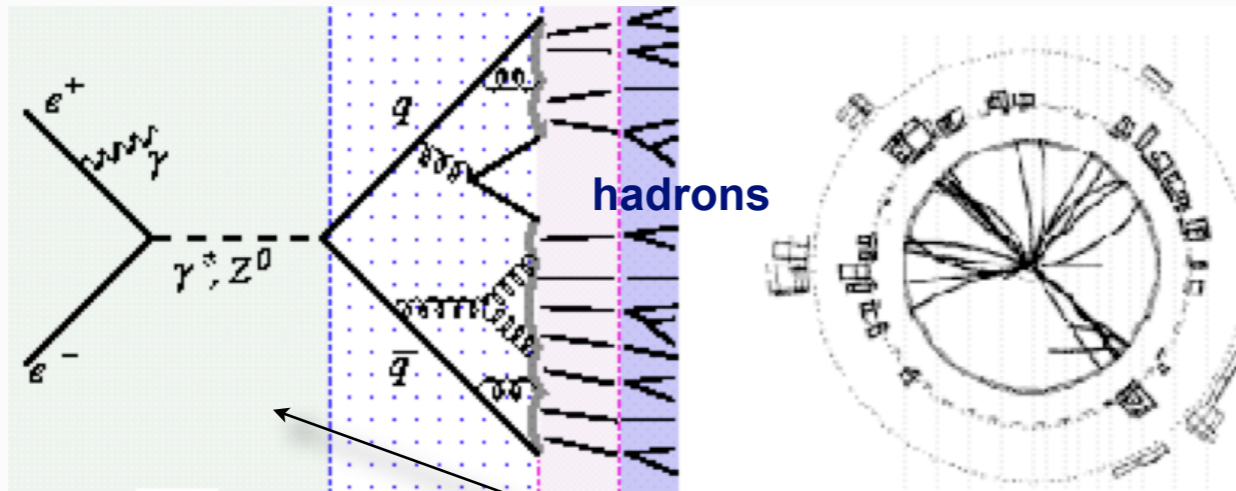
sum over all quark flavours, which can be produced at a certain e^+e^- centre-of-mass energy E_{CM} , eg. d, u, s, c, b, t

$$R := \frac{\sigma(e^+e^- \rightarrow q_f \bar{q}_f)}{\sigma(e^+e^- \rightarrow \mu^+ \mu^-)} = N_c \sum_f z_f^2$$

Number of colours

electric charges of quarks, in units of electron charge

Measurement of e^+e^- annihilation into hadrons and muons:



Hadronic final state

- many charged tracks ($> \sim 10$)
- sum of energy deposits in calorimeters not too far from centre-of-mass energy

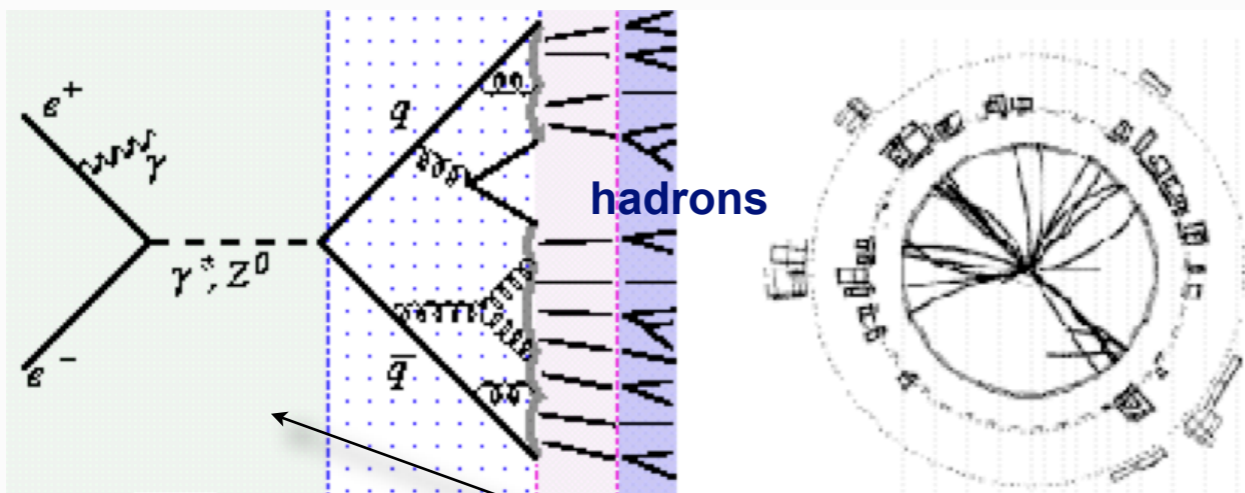
sum over all quark flavours, which can be produced at a certain e^+e^- centre-of-mass energy E_{CM} , eg. d, u, s, c, b, t

$$R := \frac{\sigma(e^+e^- \rightarrow q_f \bar{q}_f)}{\sigma(e^+e^- \rightarrow \mu^+ \mu^-)} = N_c \sum_f z_f^2$$

Number of colours

electric charges of quarks, in units of electron charge

Measurement of e^+e^- annihilation into hadrons and muons:



Hadronic final state

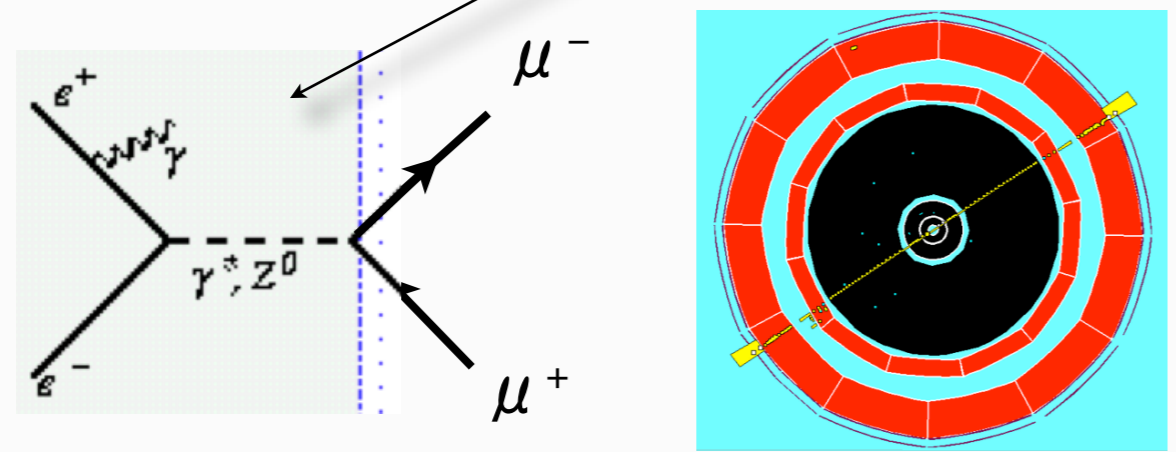
- many charged tracks ($> \sim 10$)
- sum of energy deposits in calorimeters not too far from centre-of-mass energy

sum over all quark flavours, which can be produced at a certain e^+e^- centre-of-mass energy E_{CM} , eg. d, u, s, c, b, t

$$R := \frac{\sigma(e^+e^- \rightarrow q_f \bar{q}_f)}{\sigma(e^+e^- \rightarrow \mu^+ \mu^-)} = N_c \sum_f z_f^2$$

Number of colours

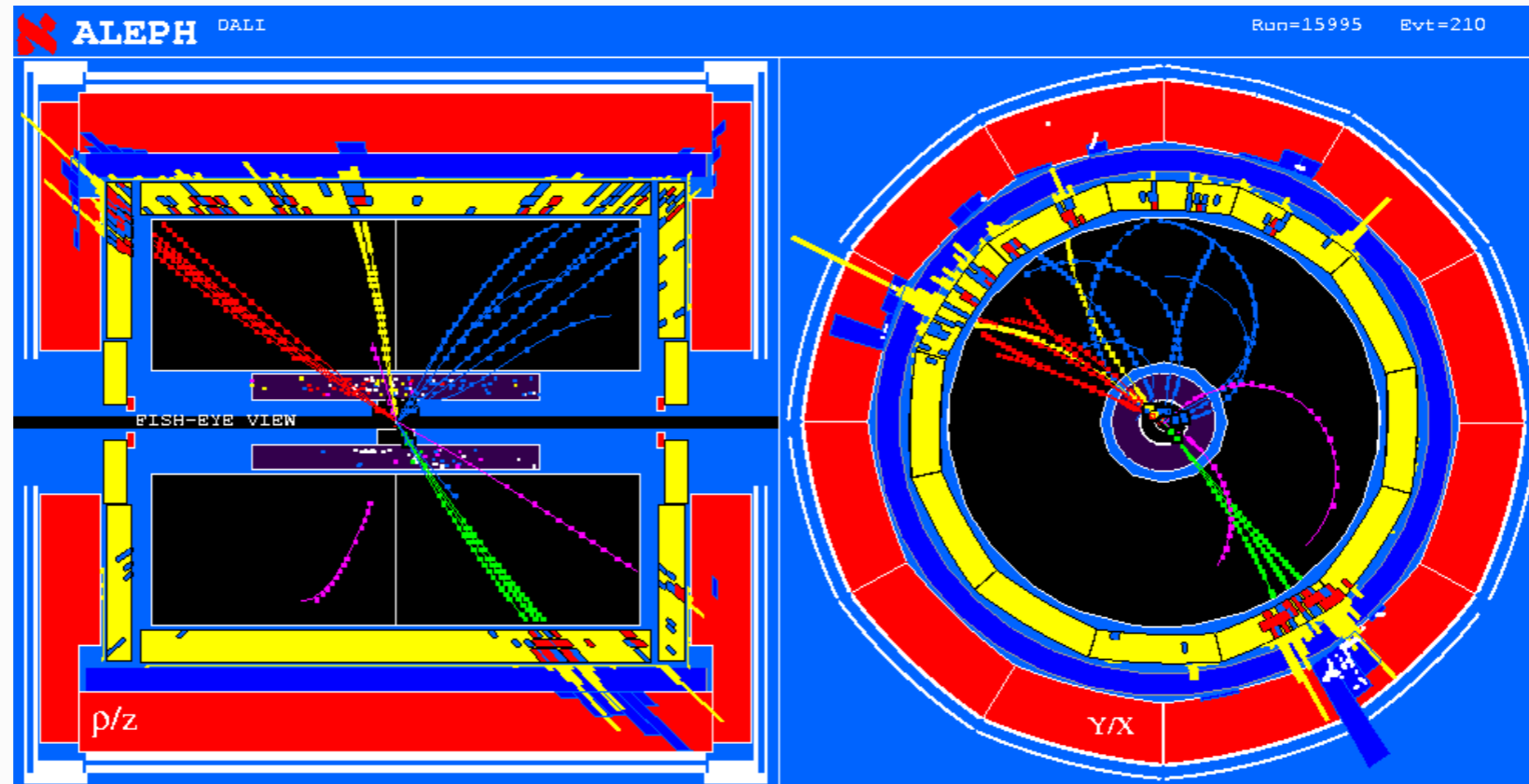
electric charges of quarks, in units of electron charge



Muonic final state

- two charged tracks, approx. back-to-back, with expected momentum ($\sim 1/2 E_{CM}$)
- right number of muon hits in outer layers (muons very penetrating, traverse whole detector)
- expected energy in calorimeter (electrons deposit all their energy, muons leave little)

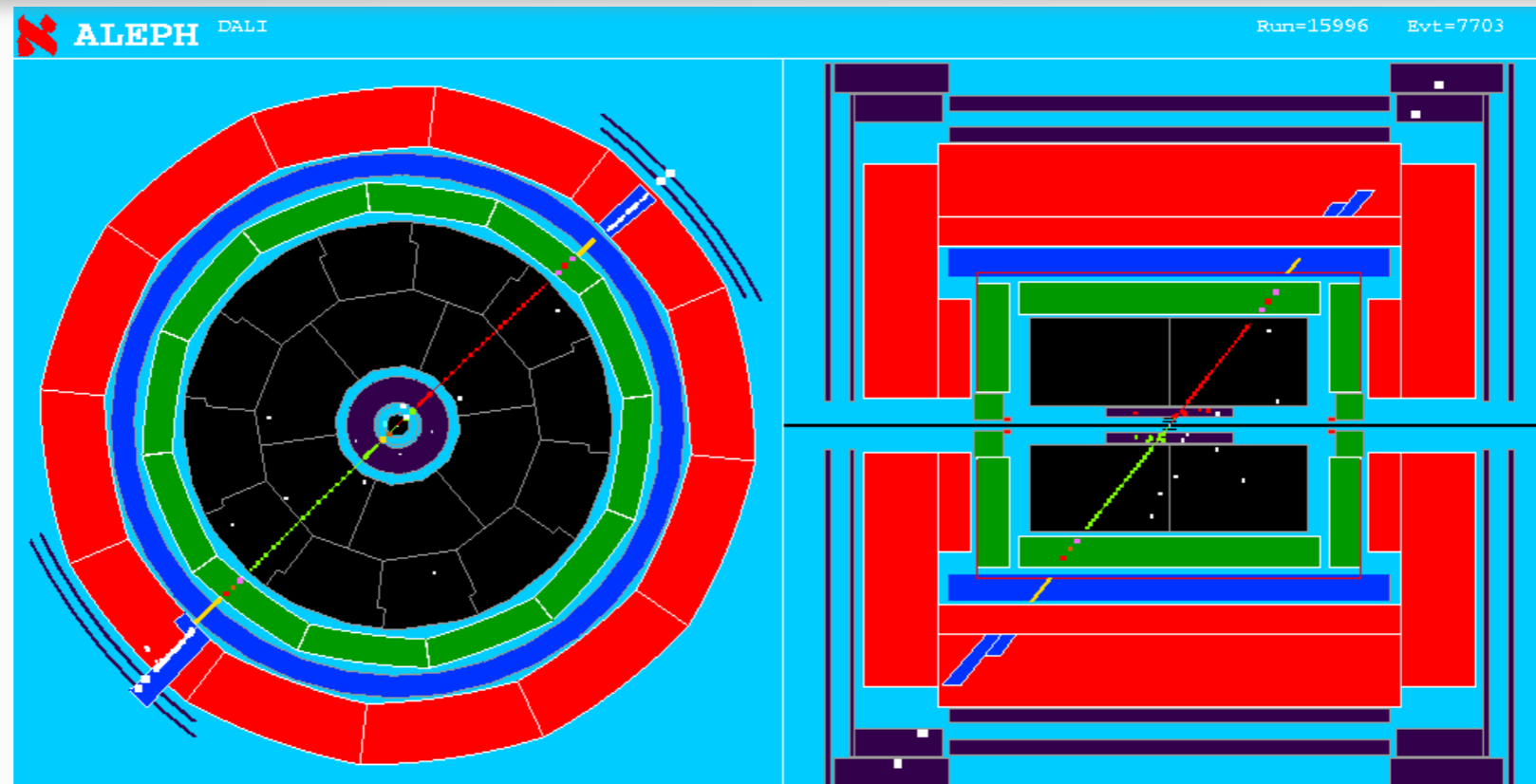
of



Hadron
final
states

R =

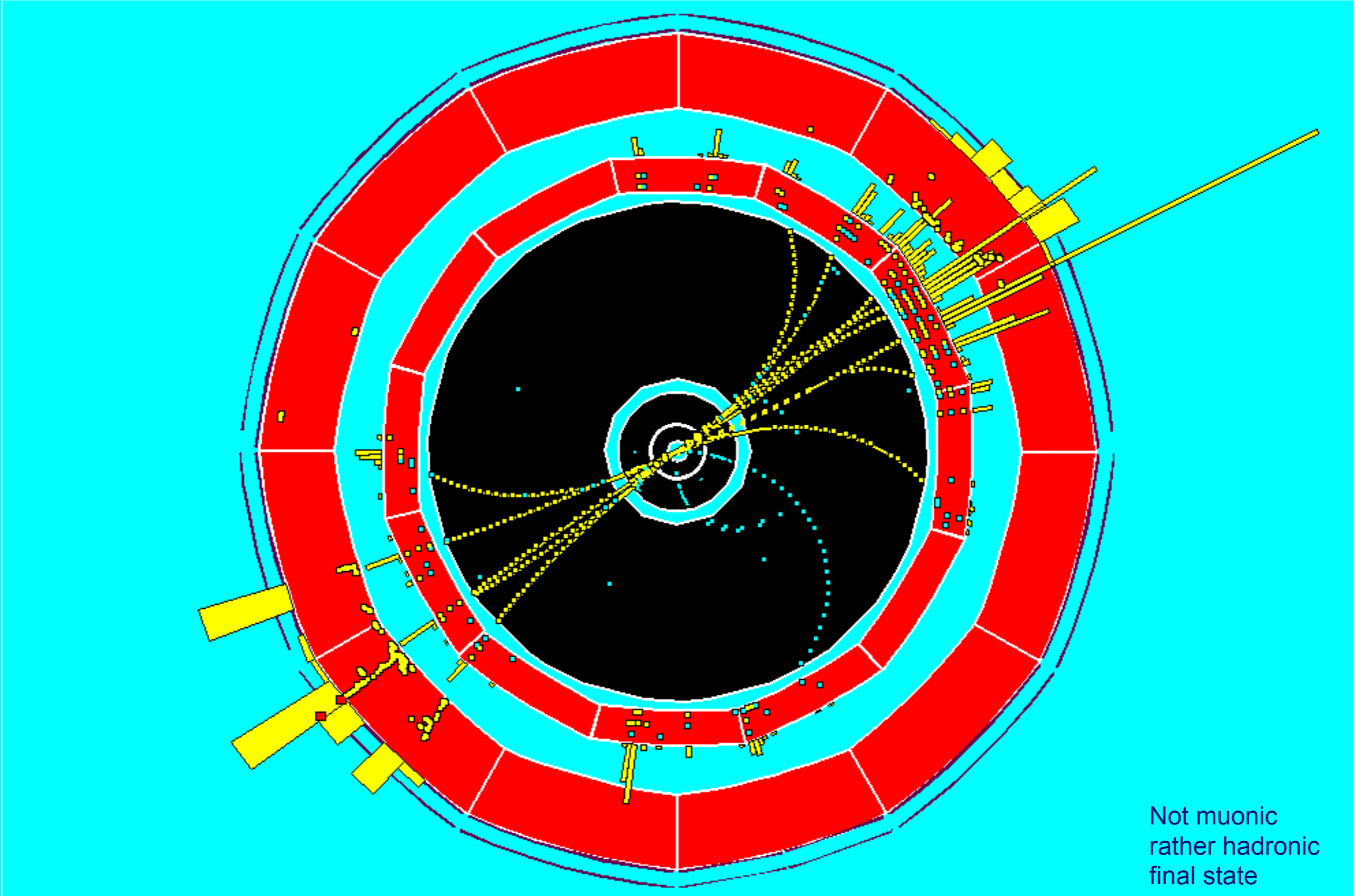
of



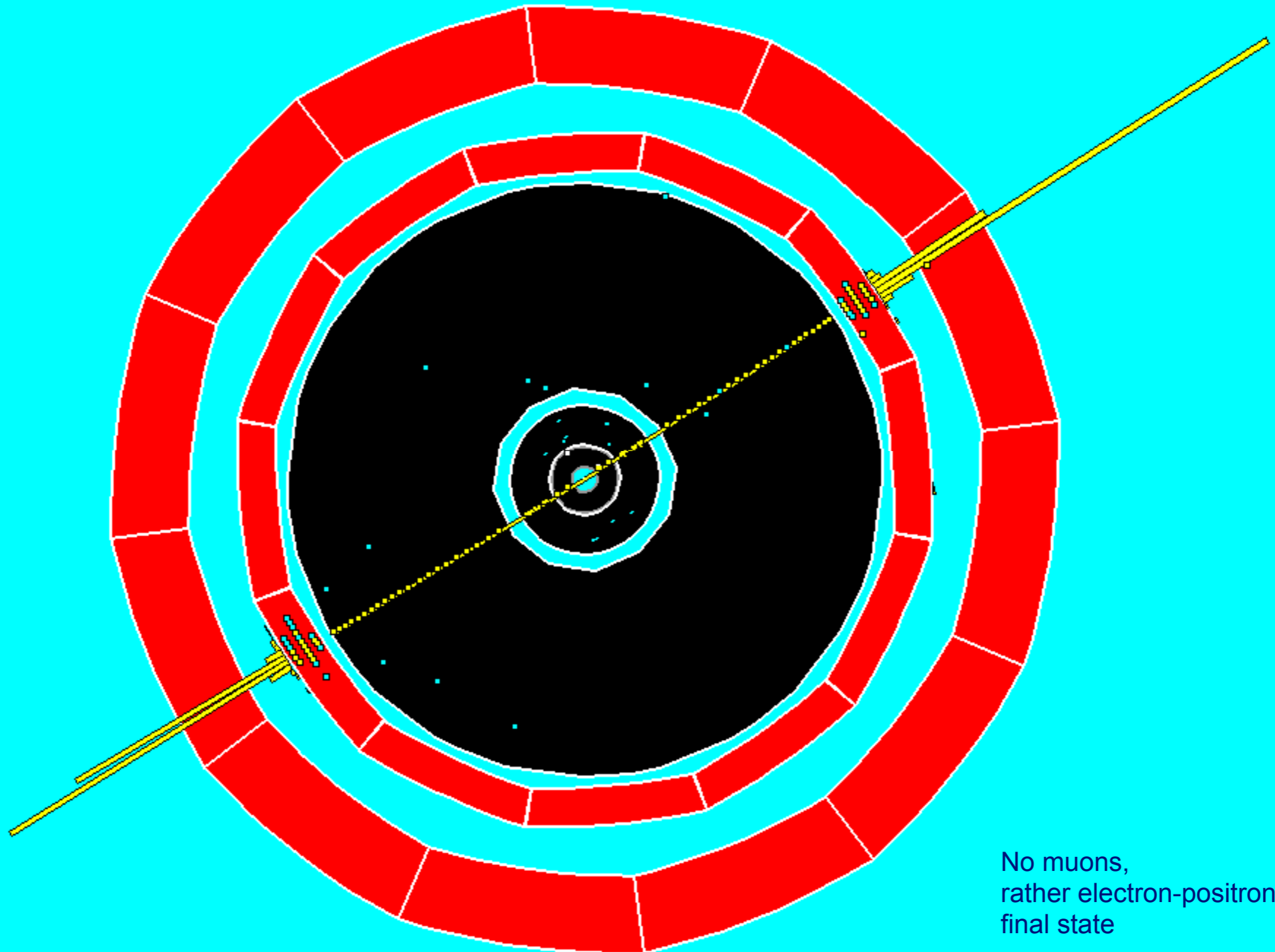
Muon
final
states

 **ALEPH** DALI

Run=15768 Evt=5906



Not muonic
rather hadronic
final state

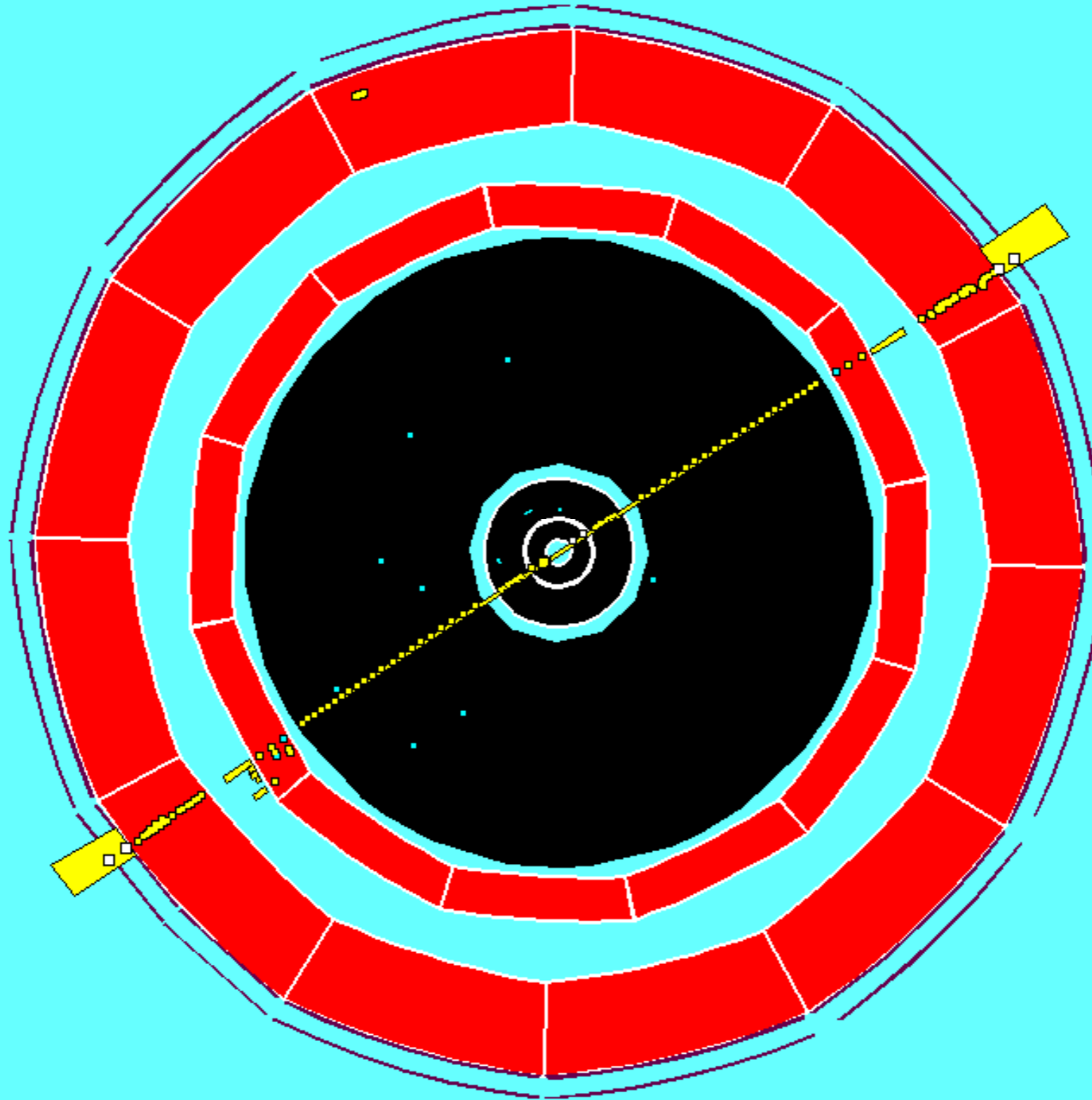


No muons,
rather electron-positron
final state

$$Z \rightarrow \mu^+ \mu^-$$

ALEPH DALI

Run=15995 Evt=835

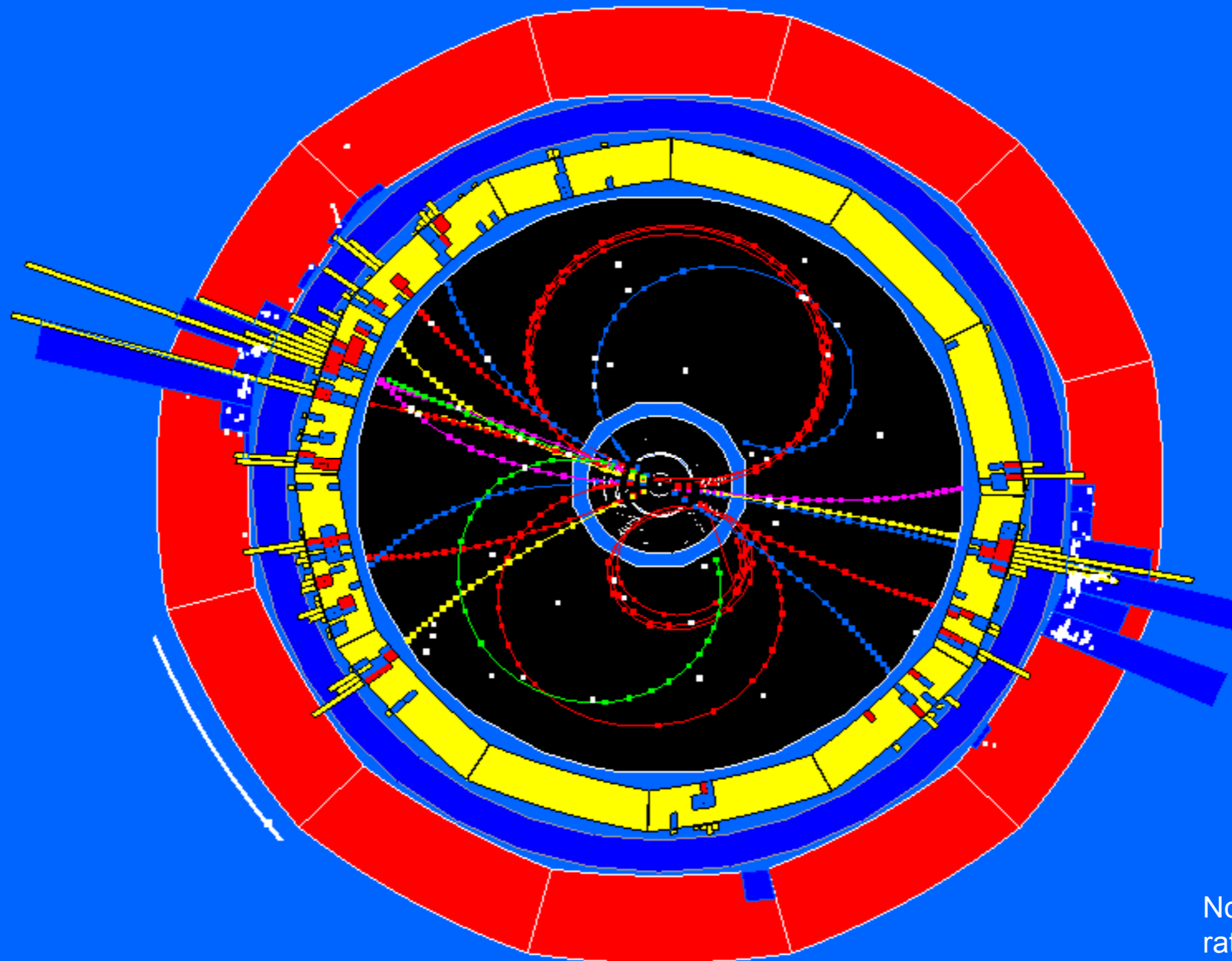


ALEPH DALI

$e^+ e^- \rightarrow q \bar{q} \rightarrow \text{hadrons}$

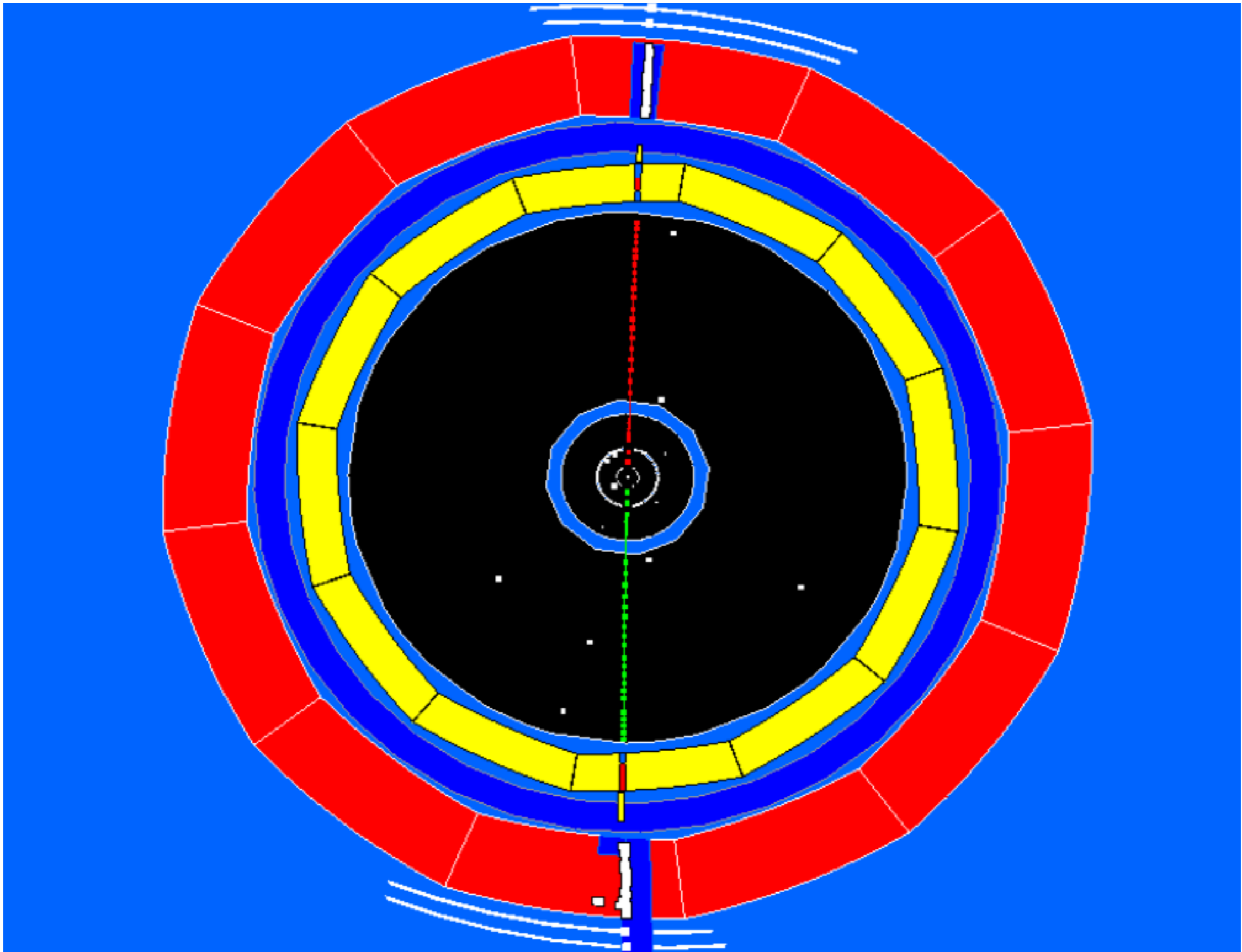
Run=15995

Evt=55



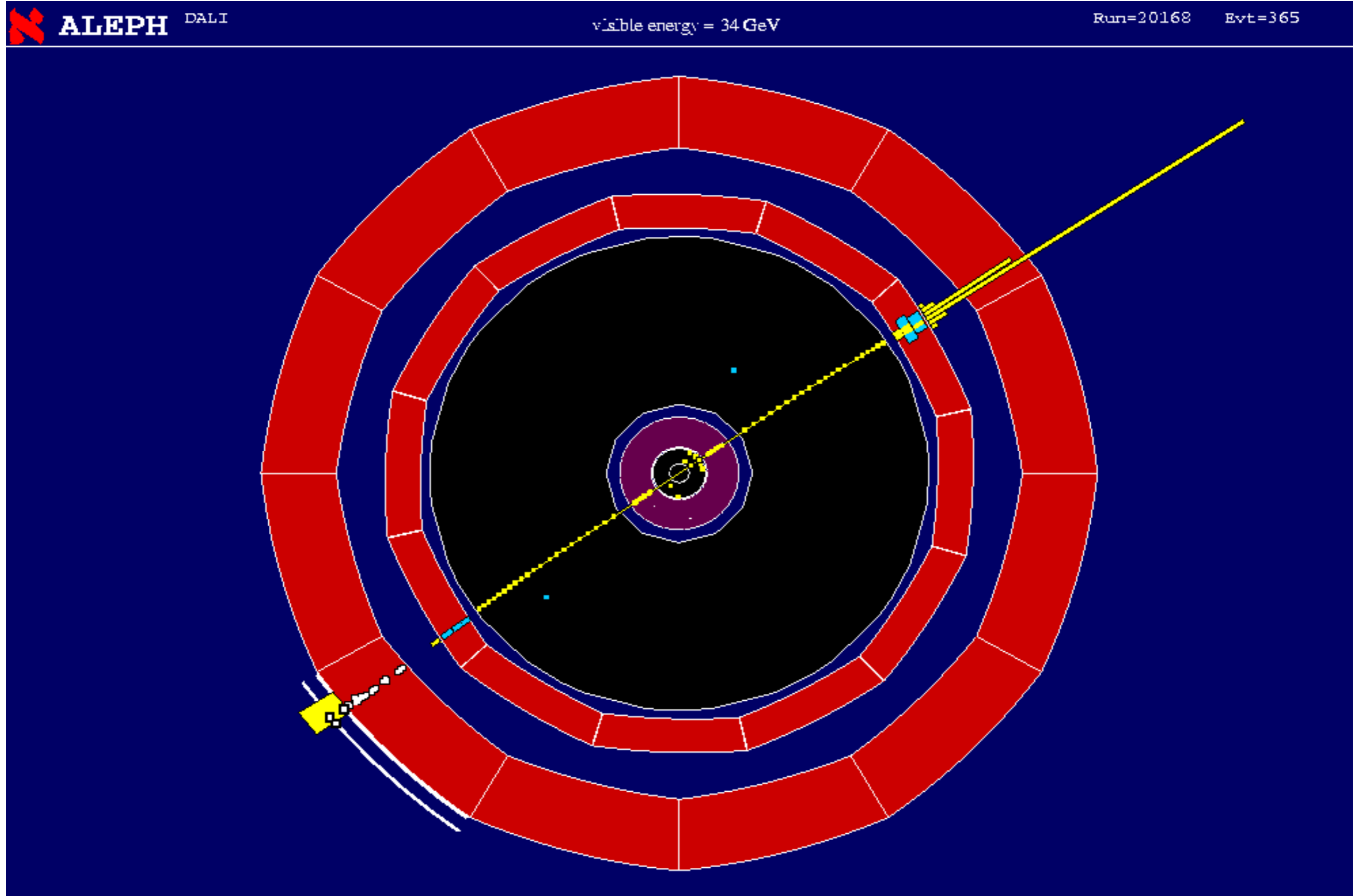
Not muonic
rather hadronic
final state

$$Z \rightarrow \mu^+ \mu^-$$

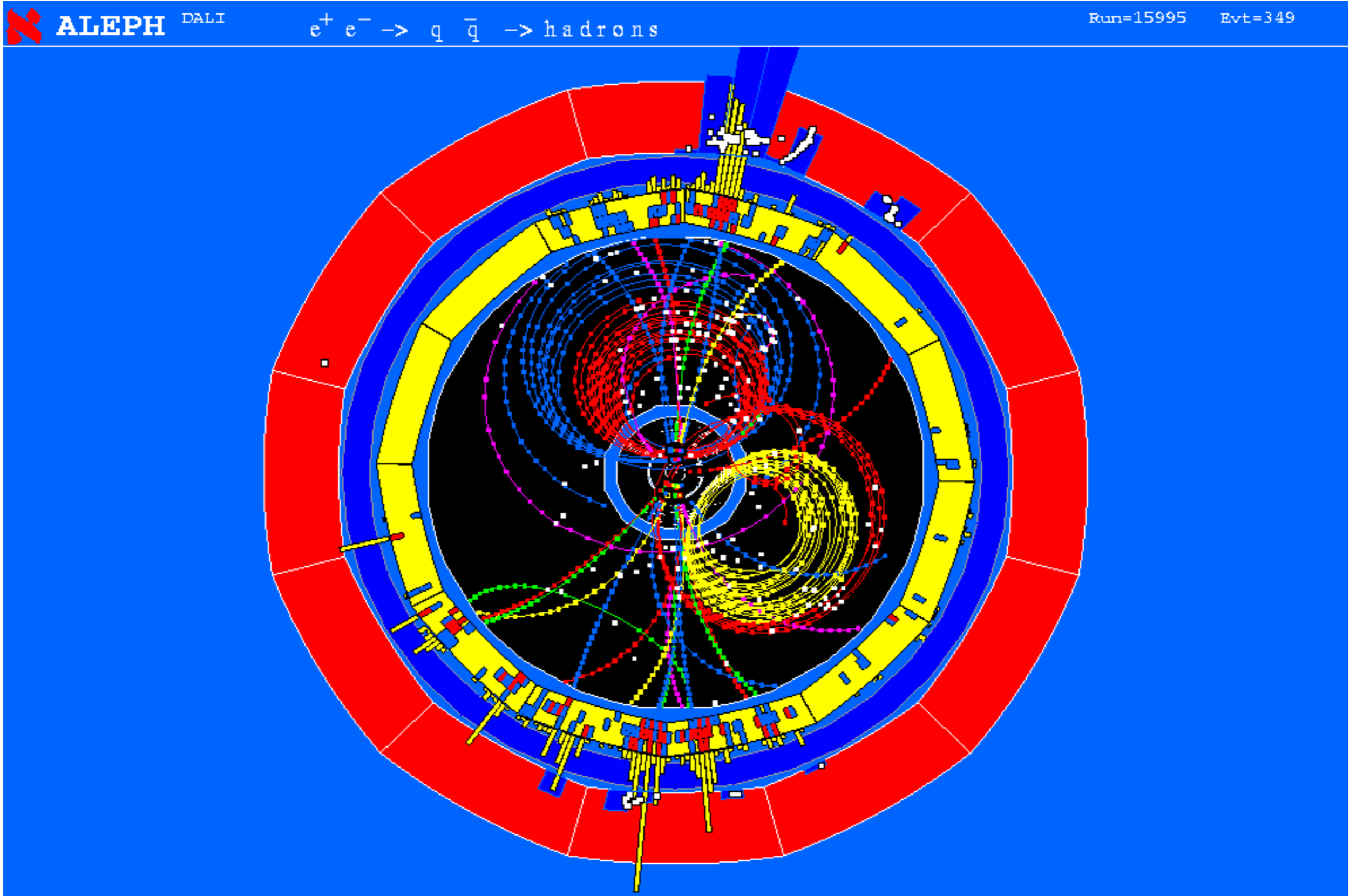


rather Z decay to $\tau^+ \tau^-$,
one tau decayed to electron + 2 neutrinos
the other tau decayed to muon + 2 neutrinos

not $Z \rightarrow \mu^+ \mu^-$



Not muonic, rather hadronic final state



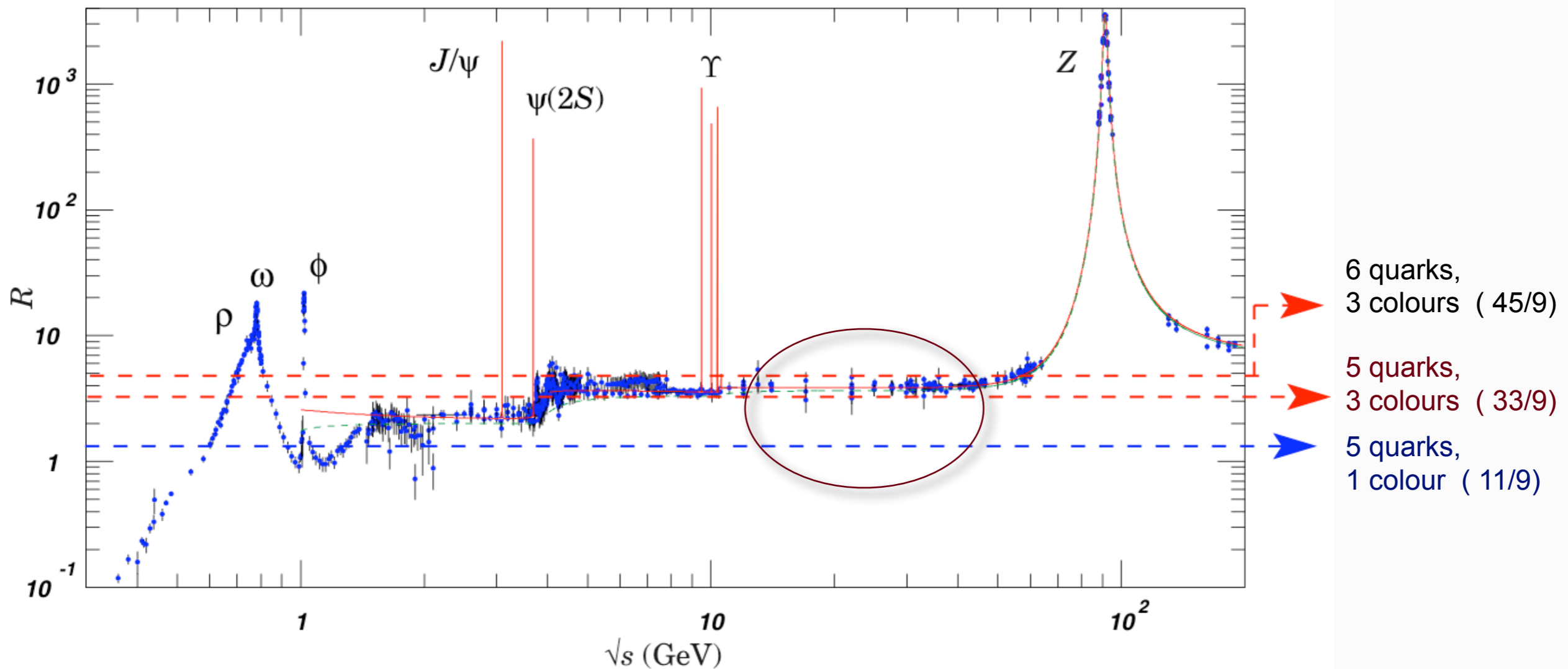
And so on.... 26

For E_{CM} below the Z peak and above the Υ resonance we expect:

$$R = N_c \sum_f z_f^2 = N_c \cdot \left[\left(\frac{2}{3}\right)^2_{\text{u}} + \left(-\frac{1}{3}\right)^2_{\text{d}} + \left(-\frac{1}{3}\right)^2_{\text{s}} + \left(\frac{2}{3}\right)^2_{\text{c}} + \left(-\frac{1}{3}\right)^2_{\text{b}} \right] = N_c \cdot \frac{11}{9}$$

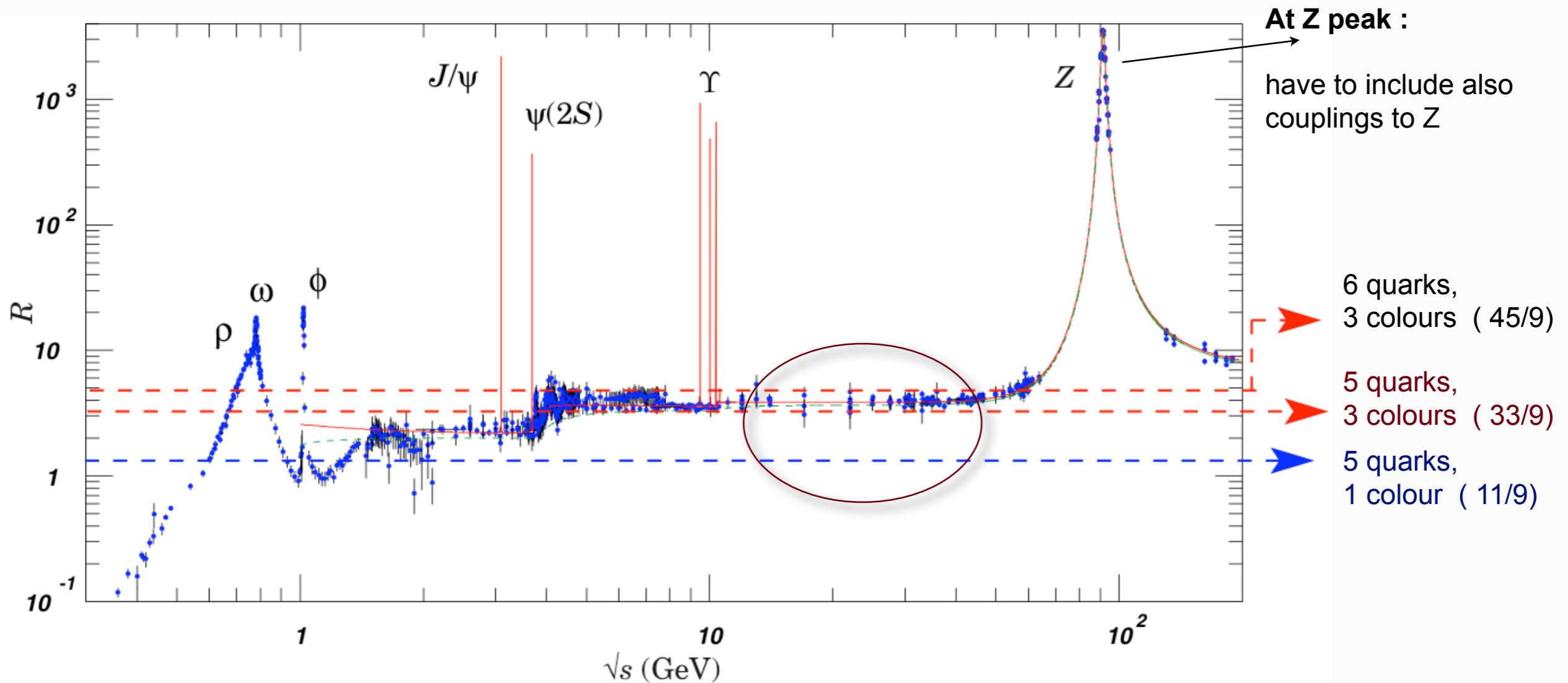
For E_{CM} below the Z peak and above the Υ resonance we expect:

$$R = N_c \sum_f z_f^2 = N_c \cdot \left[\underbrace{\left(\frac{2}{3}\right)^2}_u + \underbrace{\left(-\frac{1}{3}\right)^2}_d + \underbrace{\left(-\frac{1}{3}\right)^2}_s + \underbrace{\left(\frac{2}{3}\right)^2}_c + \underbrace{\left(-\frac{1}{3}\right)^2}_b \right] = N_c \cdot \frac{11}{9}$$



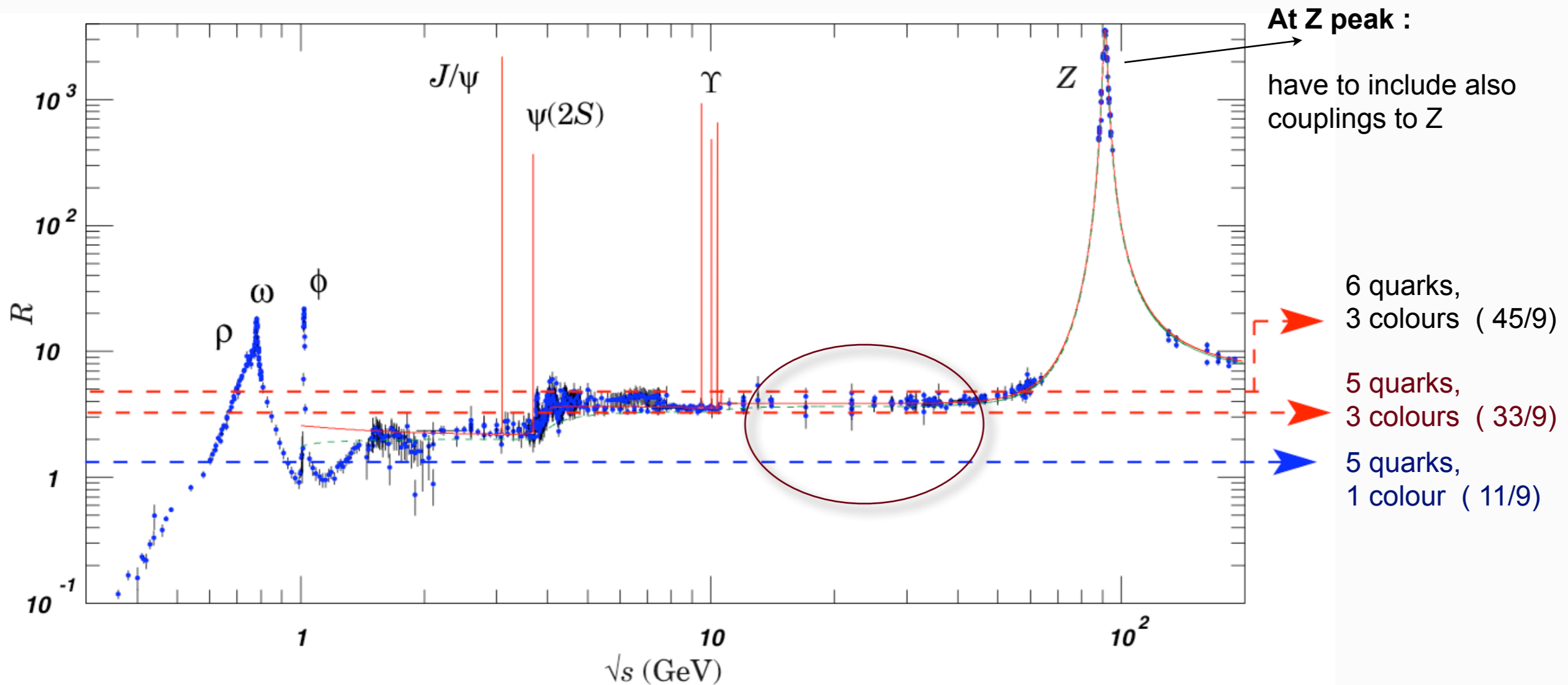
For E_{CM} below the Z peak and above the Υ resonance we expect:

$$R = N_c \sum_f z_f^2 = N_c \cdot \left[\underbrace{\left(\frac{2}{3}\right)^2}_u + \underbrace{\left(-\frac{1}{3}\right)^2}_d + \underbrace{\left(-\frac{1}{3}\right)^2}_s + \underbrace{\left(\frac{2}{3}\right)^2}_c + \underbrace{\left(-\frac{1}{3}\right)^2}_b \right] = N_c \cdot \frac{11}{9}$$



For E_{CM} below the Z peak and above the Υ resonance we expect:

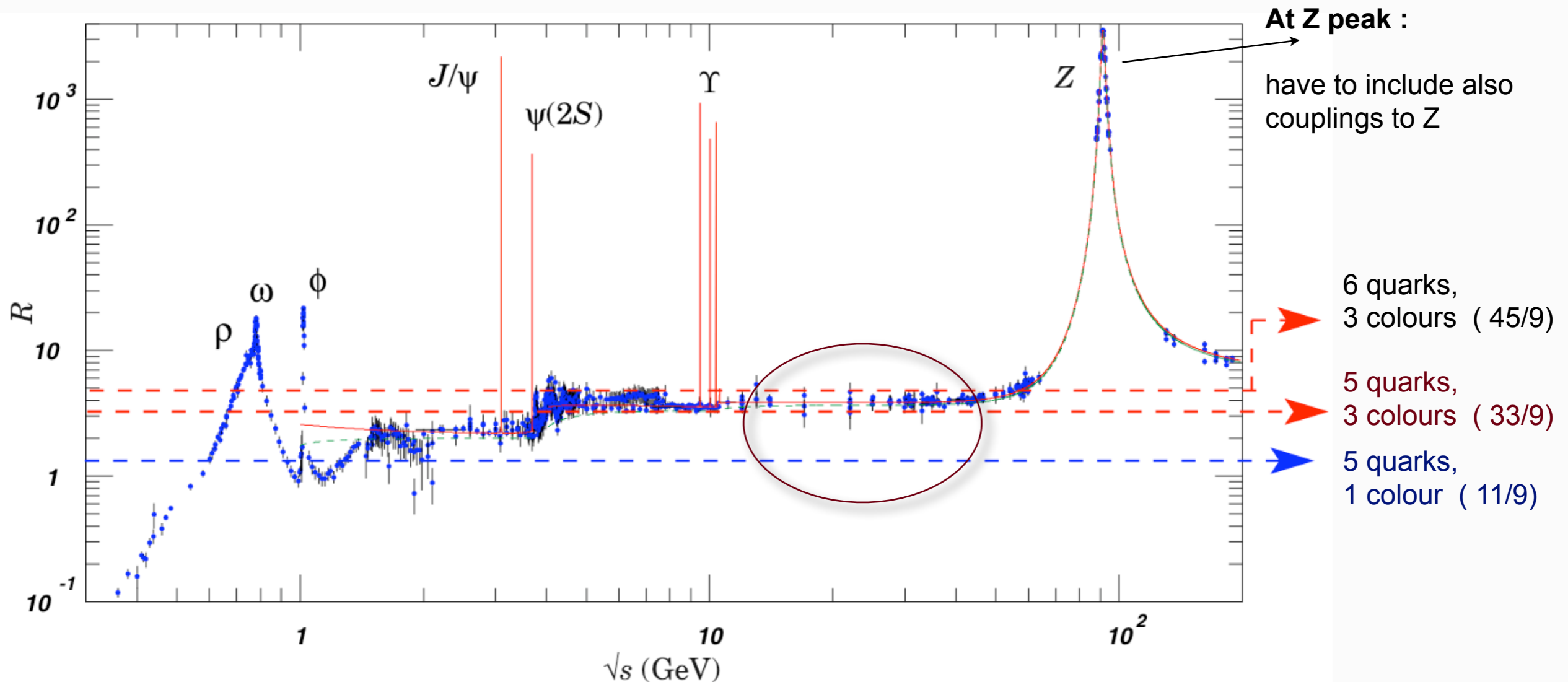
$$R = N_c \sum_f z_f^2 = N_c \cdot \left[\underbrace{\left(\frac{2}{3}\right)^2}_u + \underbrace{\left(-\frac{1}{3}\right)^2}_d + \underbrace{\left(-\frac{1}{3}\right)^2}_s + \underbrace{\left(\frac{2}{3}\right)^2}_c + \underbrace{\left(-\frac{1}{3}\right)^2}_b \right] = N_c \cdot \frac{11}{9}$$



Confirmation of : Number of colours = 3 !

For E_{CM} below the Z peak and above the Υ resonance we expect:

$$R = N_c \sum_f z_f^2 = N_c \cdot \left[\left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 + \left(\frac{2}{3}\right)^2 + \left(-\frac{1}{3}\right)^2 \right] = N_c \cdot \frac{11}{9}$$



Confirmation of : Number of colours = 3 !

Note : small remaining difference : because of QCD correction (gluon radiation) = $1 + \alpha_s / \pi$

- Just having a “counting result” is not all, there’s lot more to do!

- Just having a “counting result” is not all, there’s lot more to do!

- Statistical error**

- We saw 2 muon events, could easily have been 1 or 3
- Those fluctuations go like the square-root of the number of events

$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{\mu\mu}}{N_{total}} \pm \frac{\sqrt{N_{\mu\mu}}}{N_{total}}$$

- To reduce this uncertainty, you need to record lots (millions) of events in the detector, and process them

- Just having a “counting result” is not all, there’s lot more to do!

Statistical error

- We saw 2 muon events, could easily have been 1 or 3
- Those fluctuations go like the square-root of the number of events

$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{\mu\mu}}{N_{total}} \pm \frac{\sqrt{N_{\mu\mu}}}{N_{total}}$$

- To reduce this uncertainty, you need to record lots (millions) of events in the detector, and process them

Systematic error

- What if you only see 50% of the $\mu^+ \mu^-$ events?

$$N_{\mu\mu_{seen}} = \overset{\text{“efficiency”}}{\varepsilon} N_{\mu\mu}$$

- because of event selection (cut), detector imperfections, poor understanding, etc.

- Just having a “counting result” is not all, there’s lot more to do!

Statistical error

- We saw 2 muon events, could easily have been 1 or 3
- Those fluctuations go like the square-root of the number of events

$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{\mu\mu}}{N_{total}} \pm \frac{\sqrt{N_{\mu\mu}}}{N_{total}}$$

- To reduce this uncertainty, you need to record lots (millions) of events in the detector, and process them

Systematic error

- What if you only see 50% of the $\mu^+ \mu^-$ events?

“efficiency”

$$N_{\mu\mu_{seen}} = \epsilon N_{\mu\mu}$$

- because of event selection (cut), detector imperfections, poor understanding, etc.

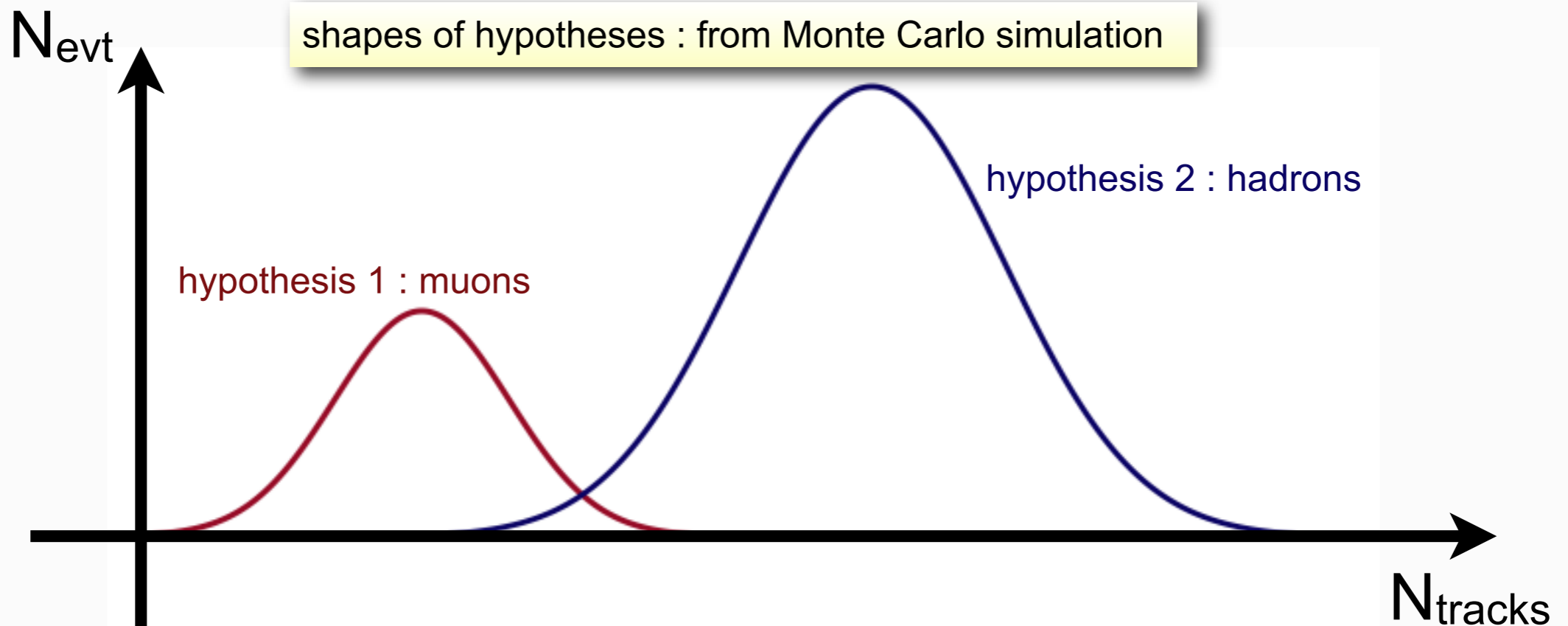
$$BR(Z^0 \rightarrow \mu^+ \mu^-) = \frac{N_{seen}/\epsilon}{N_{total}}$$

$$\epsilon = 0.50 \pm 0.05$$

- from statistical error of detector simulation
- imperfect modeling of geometry in simulation
- model of muon interactions in simulation, etc

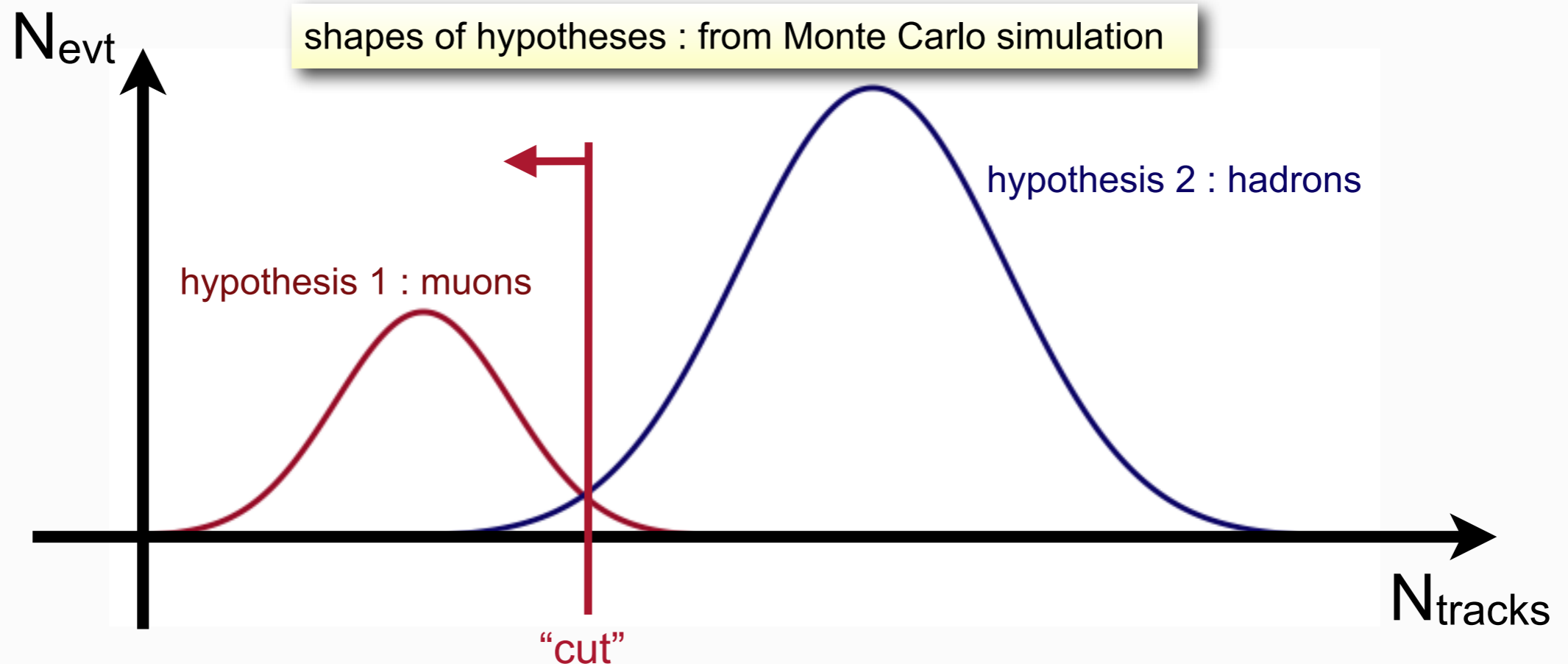
- Event per event have to decide how to categorize it
 - eg. do we call it a muon event, or a hadronic event?
 - how do we estimate the efficiency?
 - Define an **event selection**, eg. “cut-based”
 - see statistics lectures, *hypothesis testing* etc...

- Event per event have to decide how to categorize it
 - eg. do we call it a muon event, or a hadronic event?
 - how do we estimate the efficiency?
 - Define an **event selection**, eg. “cut-based”
 - see statistics lectures, *hypothesis testing* etc...



Event per event have to decide how to categorize it

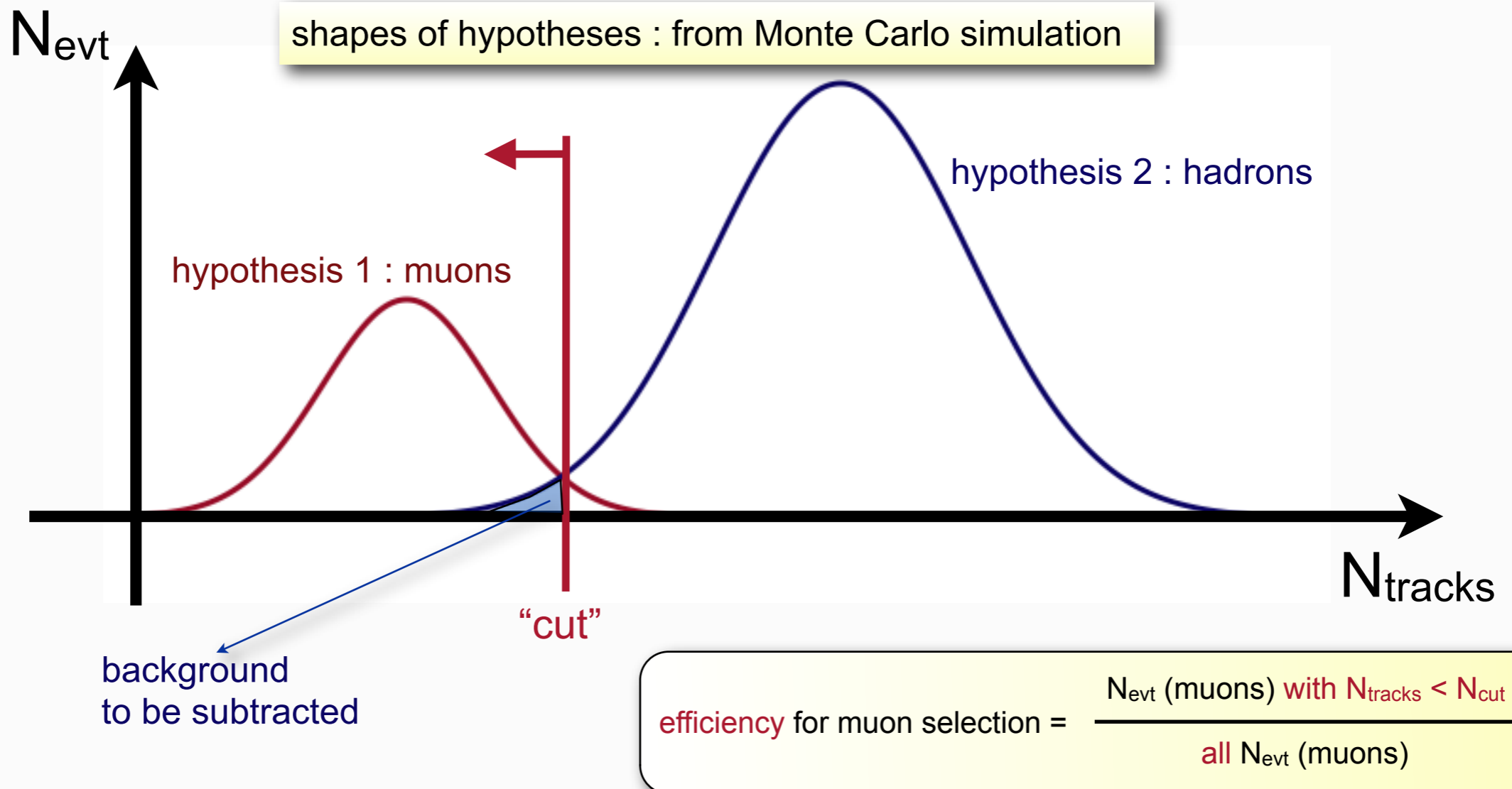
- eg. do we call it a muon event, or a hadronic event?
- how do we estimate the efficiency?
- Define an **event selection**, eg. “cut-based”
- see statistics lectures, *hypothesis testing* etc...



$$\text{efficiency for muon selection} = \frac{N_{evt} \text{ (muons) with } N_{tracks} < N_{cut}}{\text{all } N_{evt} \text{ (muons)}}$$

Event per event have to decide how to categorize it

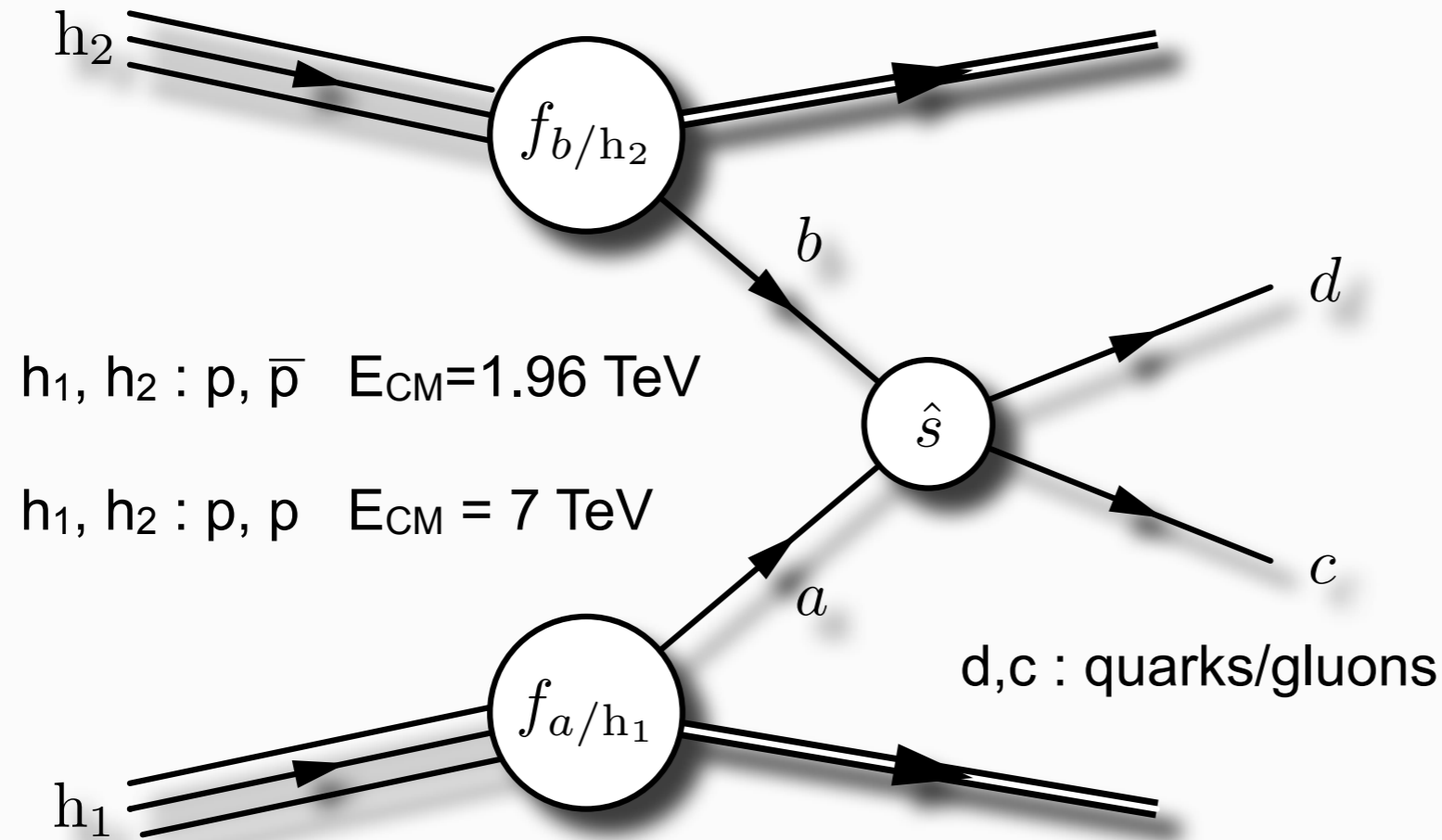
- eg. do we call it a muon event, or a hadronic event?
- how do we estimate the efficiency?
- Define an **event selection**, eg. “cut-based”
- see statistics lectures, *hypothesis testing* etc...



A “more complicated” example

“The greater the obstacle, the more glory in overcoming it.” (Moliere)

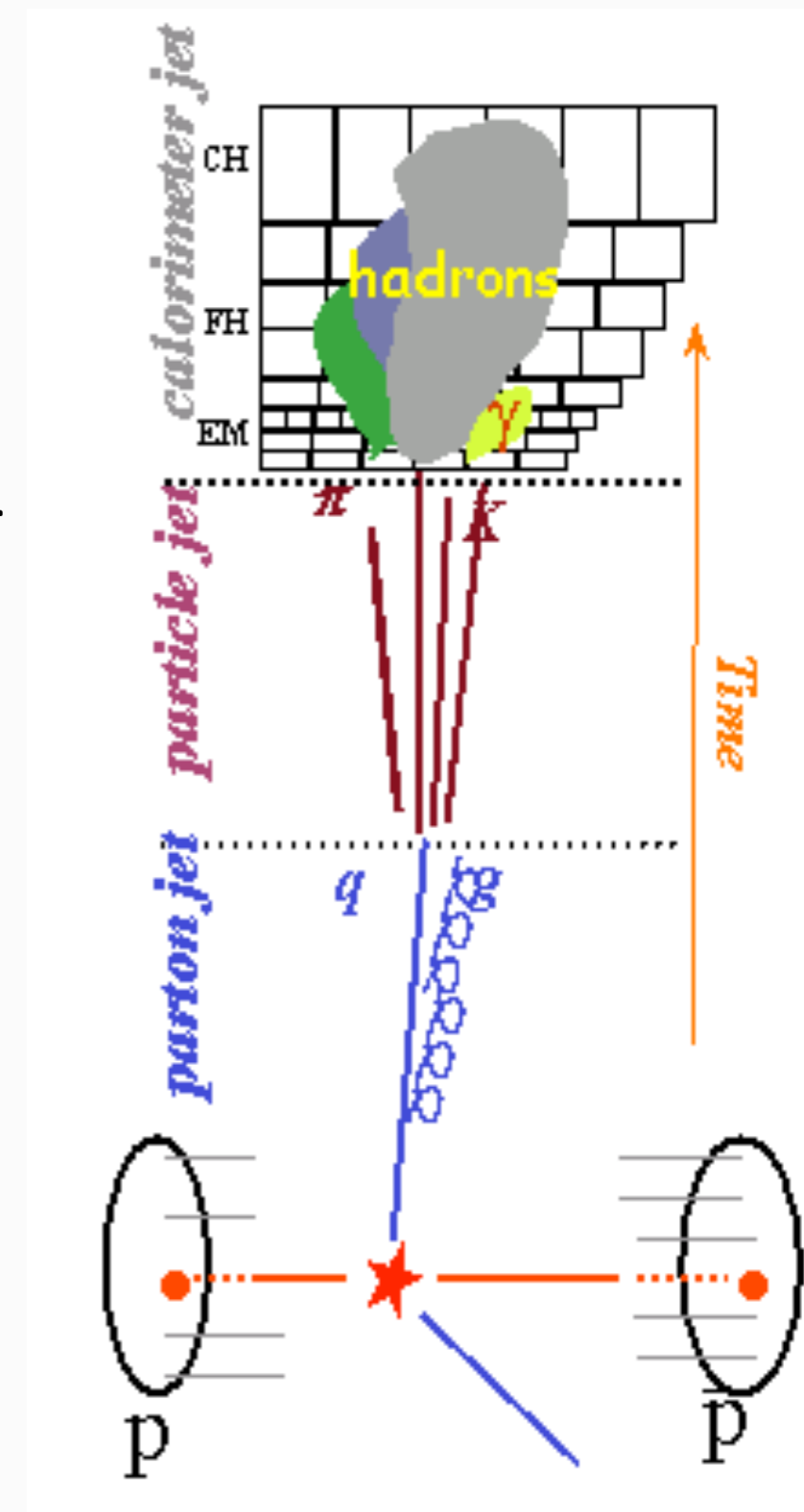
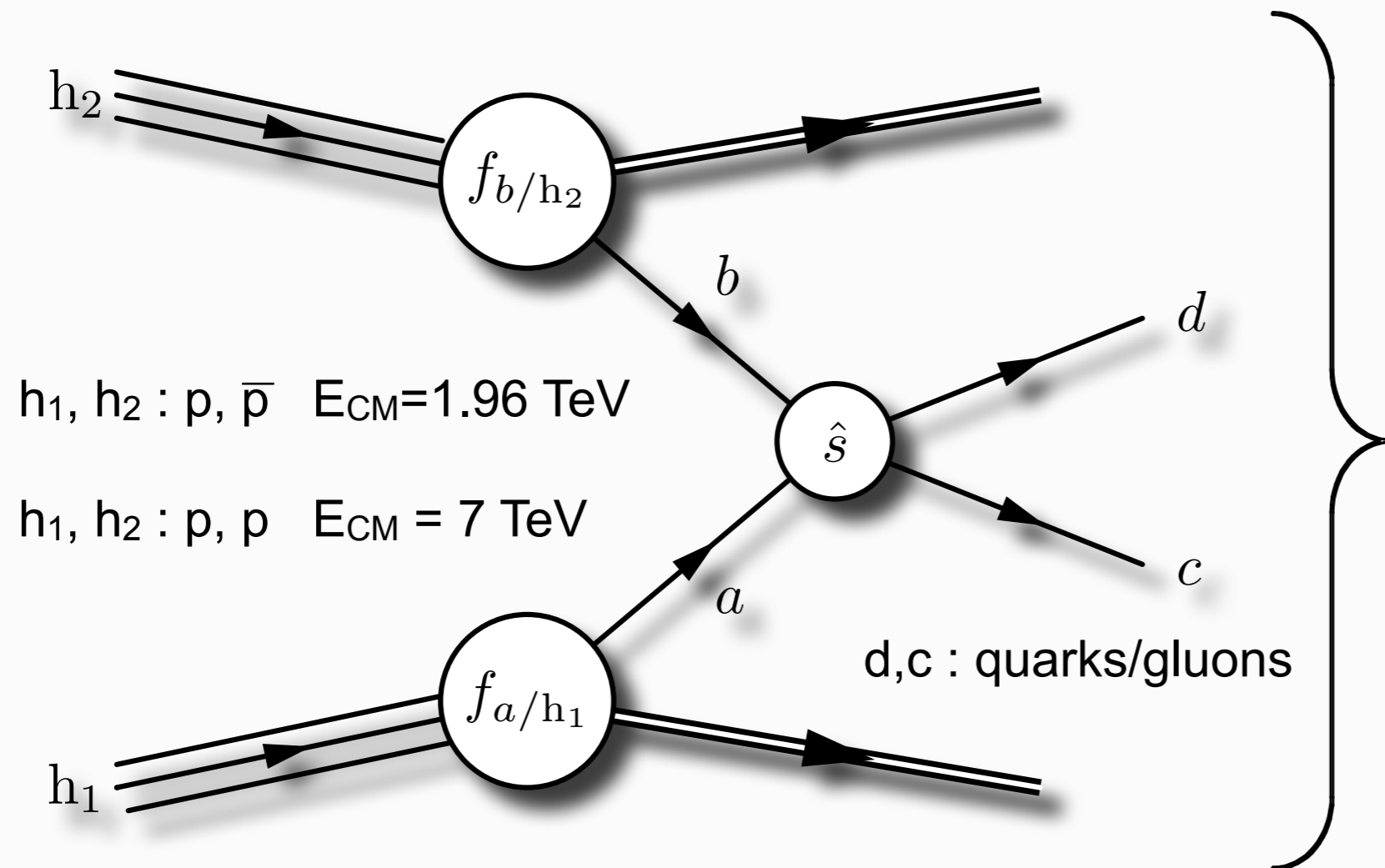
at the Tevatron, or now at the LHC



see also lectures by B. Heinemann and by S. Ellis

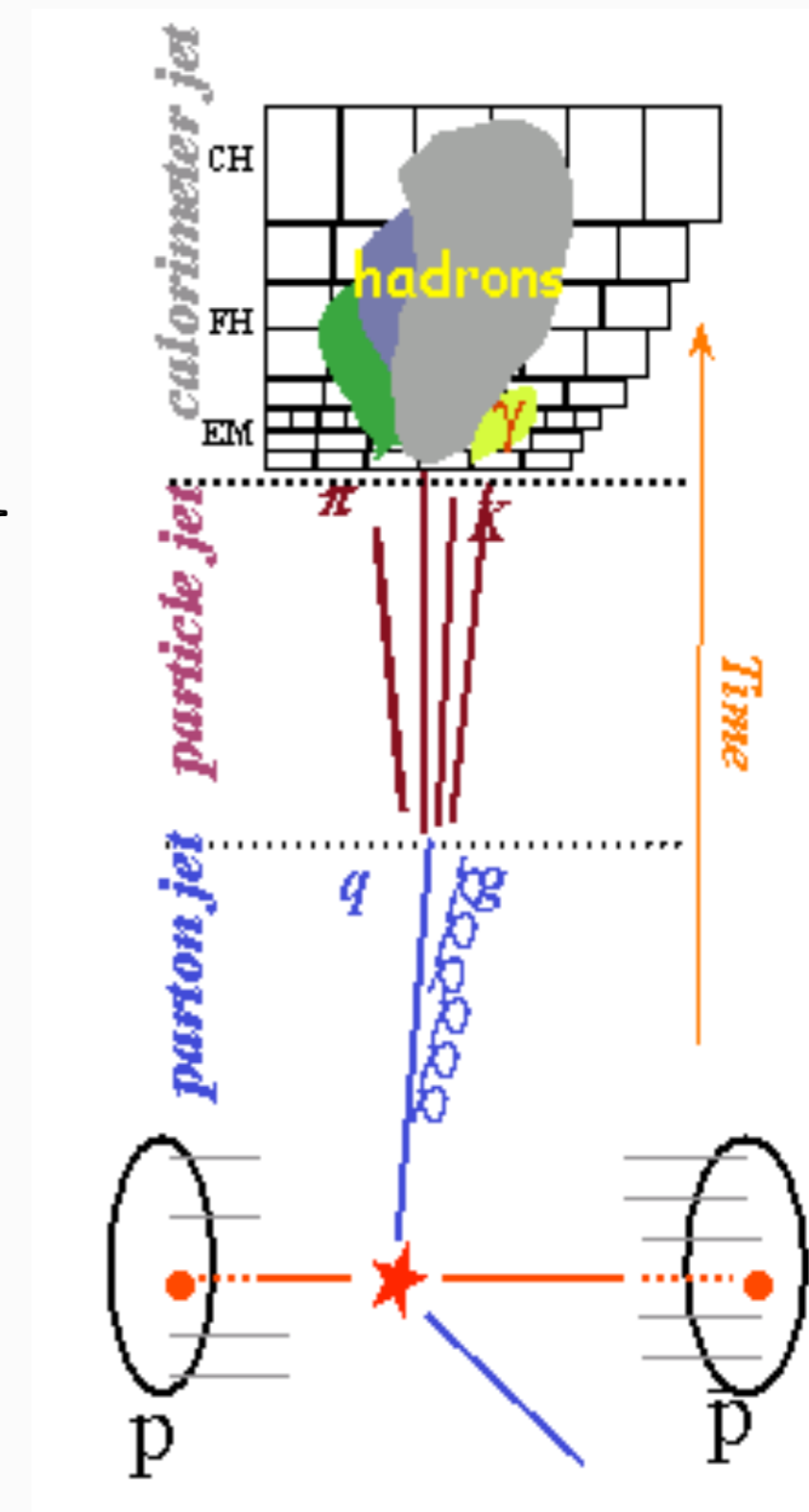
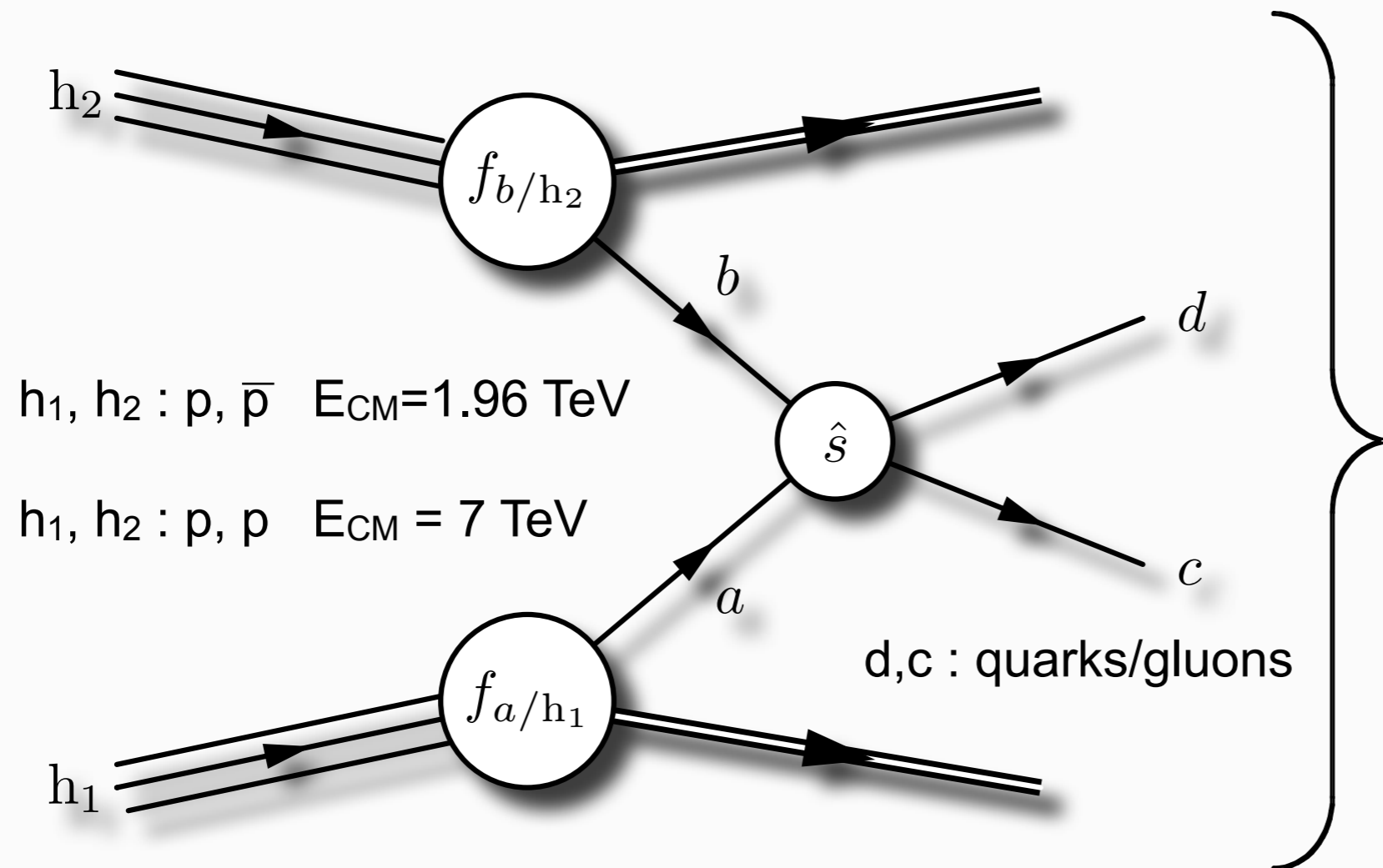
see also lectures by B. Heinemann and by S. Ellis

at the Tevatron, or now at the LHC



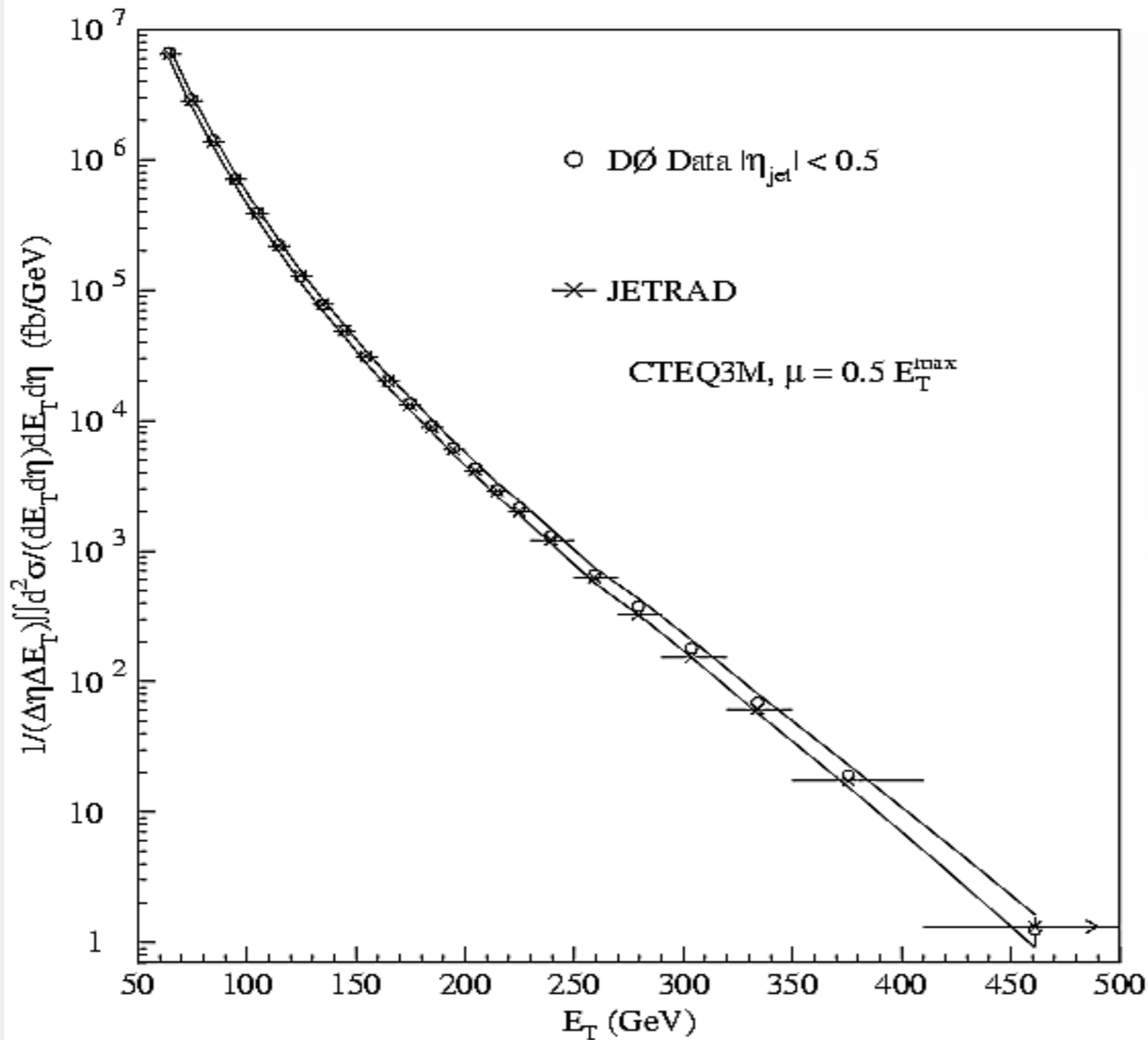
see also lectures by B. Heinemann and by S. Ellis

at the Tevatron, or now at the LHC



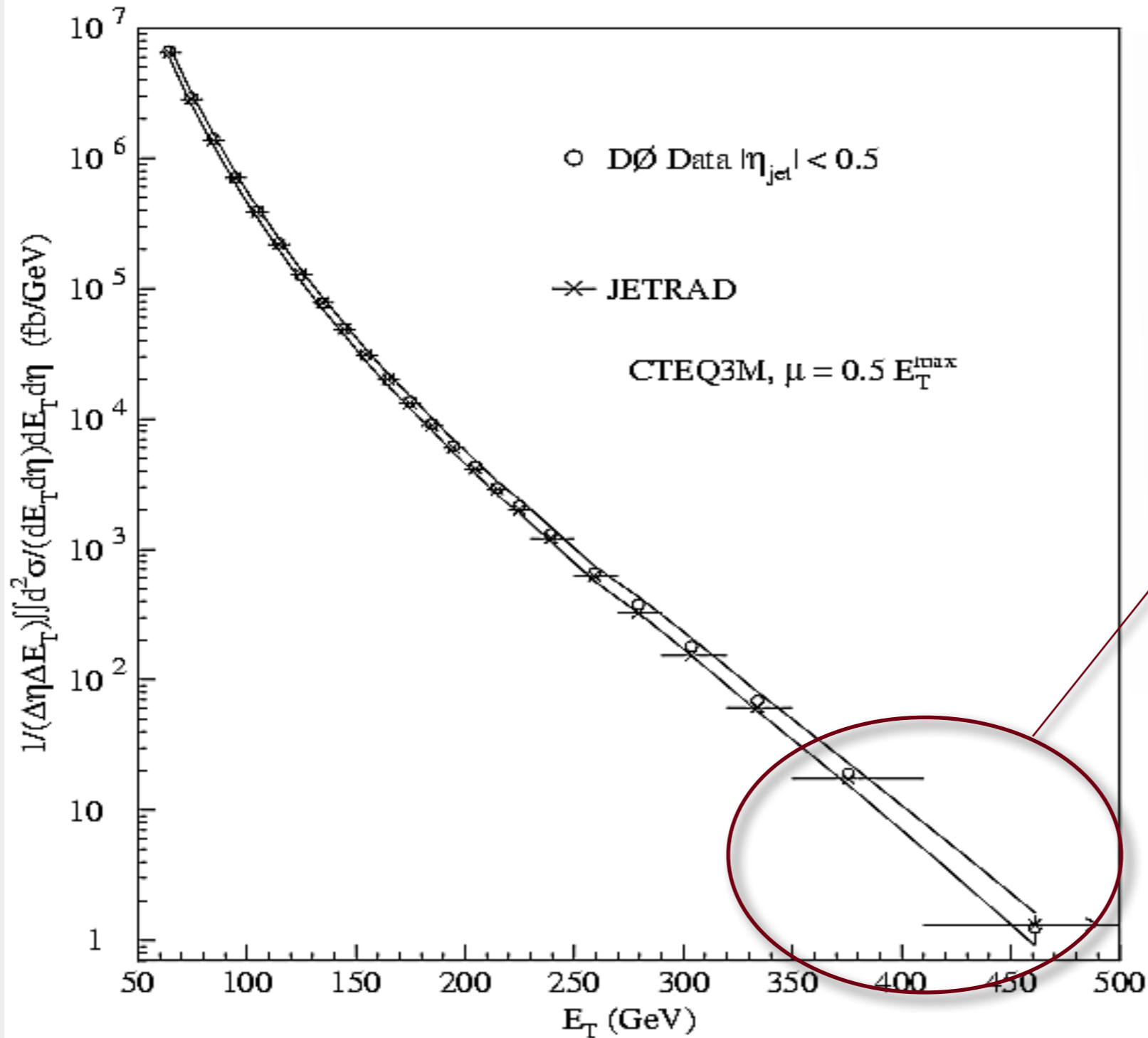
Goal

- measure probability that **quarks/gluons** are produced with a certain energy, at a certain angle
- Problem** : do not observe quarks and gluons directly, only hadrons, which appear collimated into **jets**
- Reconstruct tracks and/or energy clusters in the calorimeter



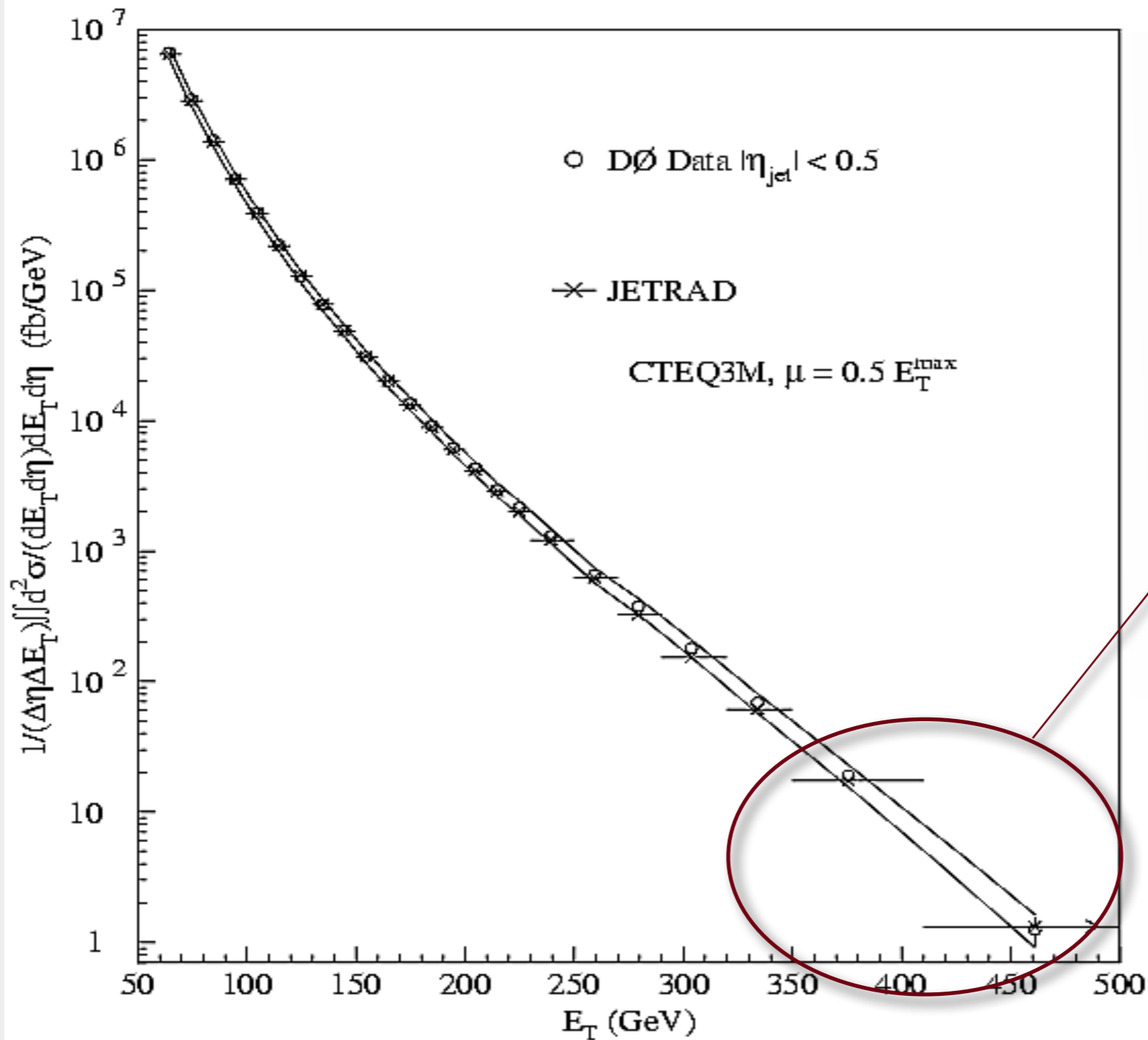
Goal

- measure **cross section** (probability) that **jets** are produced with a certain **transverse energy E_T** , within a certain **rapidity range**
- Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)



Goal

- measure **cross section** (probability) that **jets** are produced with a certain **transverse energy E_T** , within a certain **rapidity range**
- **Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)

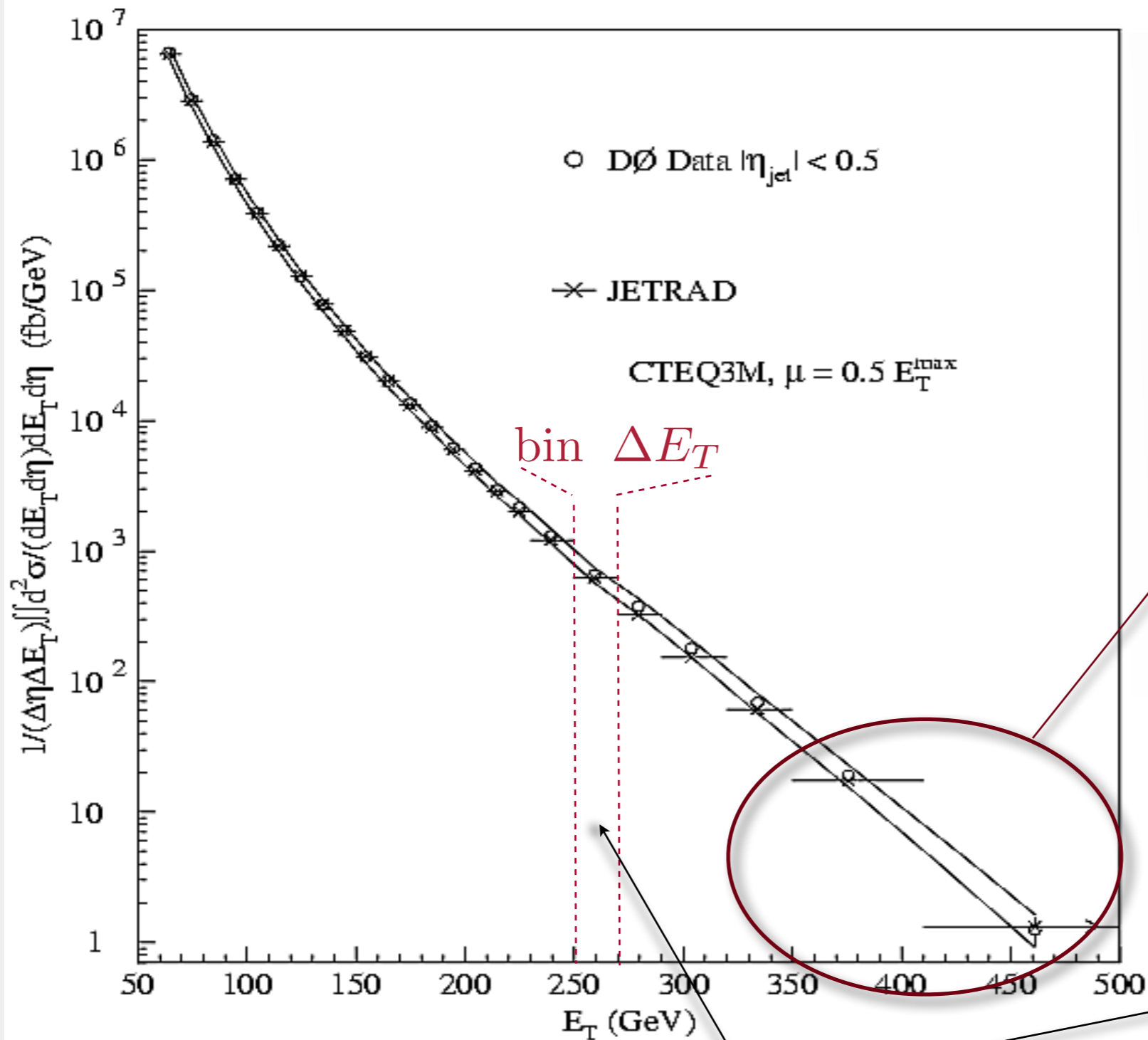


Goal

- measure **cross section** (probability) that **jets** are produced with a certain **transverse energy E_T** , within a certain **rapidity range**
- Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)

can be calculated in pert. QCD

$$\left\langle \frac{d^2\sigma}{dE_T d\eta} \right\rangle = \frac{N}{\Delta E_T \Delta \eta \epsilon \mathcal{L}}$$



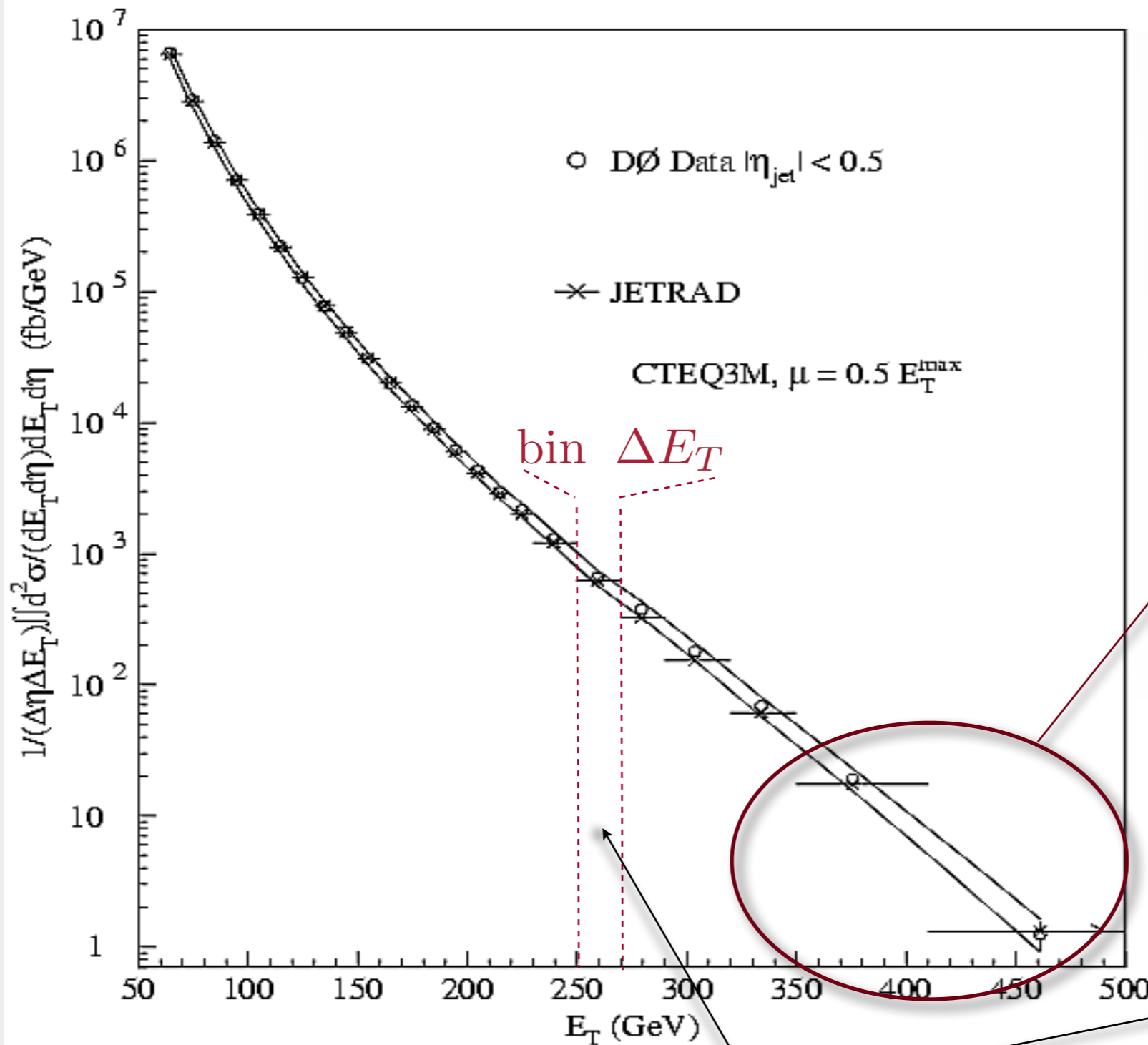
Goal

- measure **cross section** (probability) that **jets** are produced with a certain **transverse energy** E_T , within a certain **rapidity range**
- Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)

can be calculated in pert. QCD

$$\left\langle \frac{d^2\sigma}{dE_T d\eta} \right\rangle = \frac{N}{\Delta E_T \Delta \eta \epsilon \mathcal{L}}$$

- count** number of events, N , in this bin
- for a certain range in **rapidity** (angle) $\Delta \eta$



Goal

- measure **cross section** (probability) that **jets** are produced with a certain **transverse energy E_T** , within a certain **rapidity range**
- Test of perturbative QCD**, over many orders of magnitude!
- Look at **very high energy tail**, **new physics** could show up there in form of excess (eg. sub-structure of quarks?)

can be calculated in pert. QCD

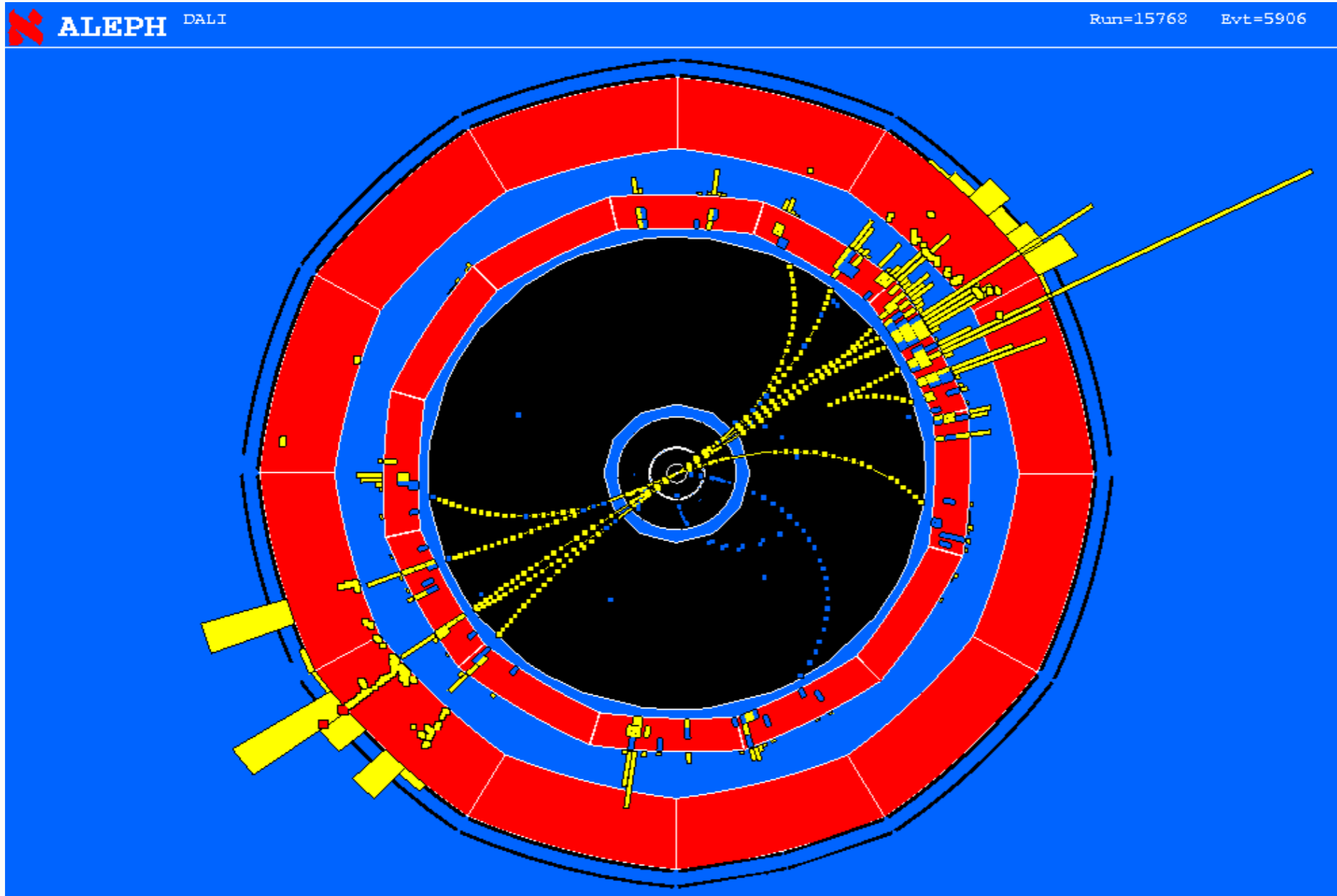
$$\left\langle \frac{d^2\sigma}{dE_T d\eta} \right\rangle = \frac{N}{\Delta E_T \Delta \eta \epsilon \mathcal{L}}$$

count number of events, N , in this bin

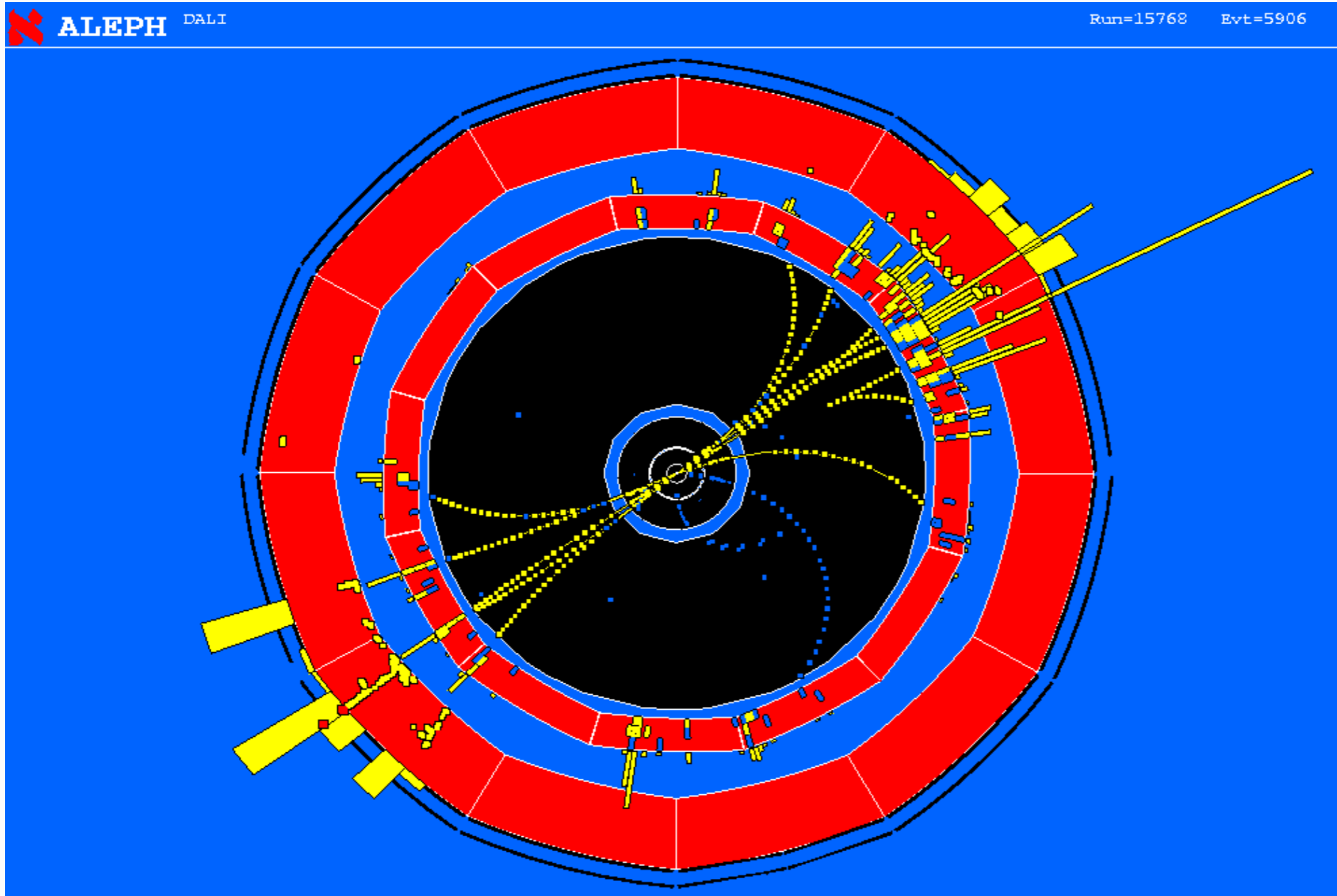
for a certain range in **rapidity** (angle) $\Delta \eta$

efficiency to reconstruct jets

integrated accelerator **luminosity**

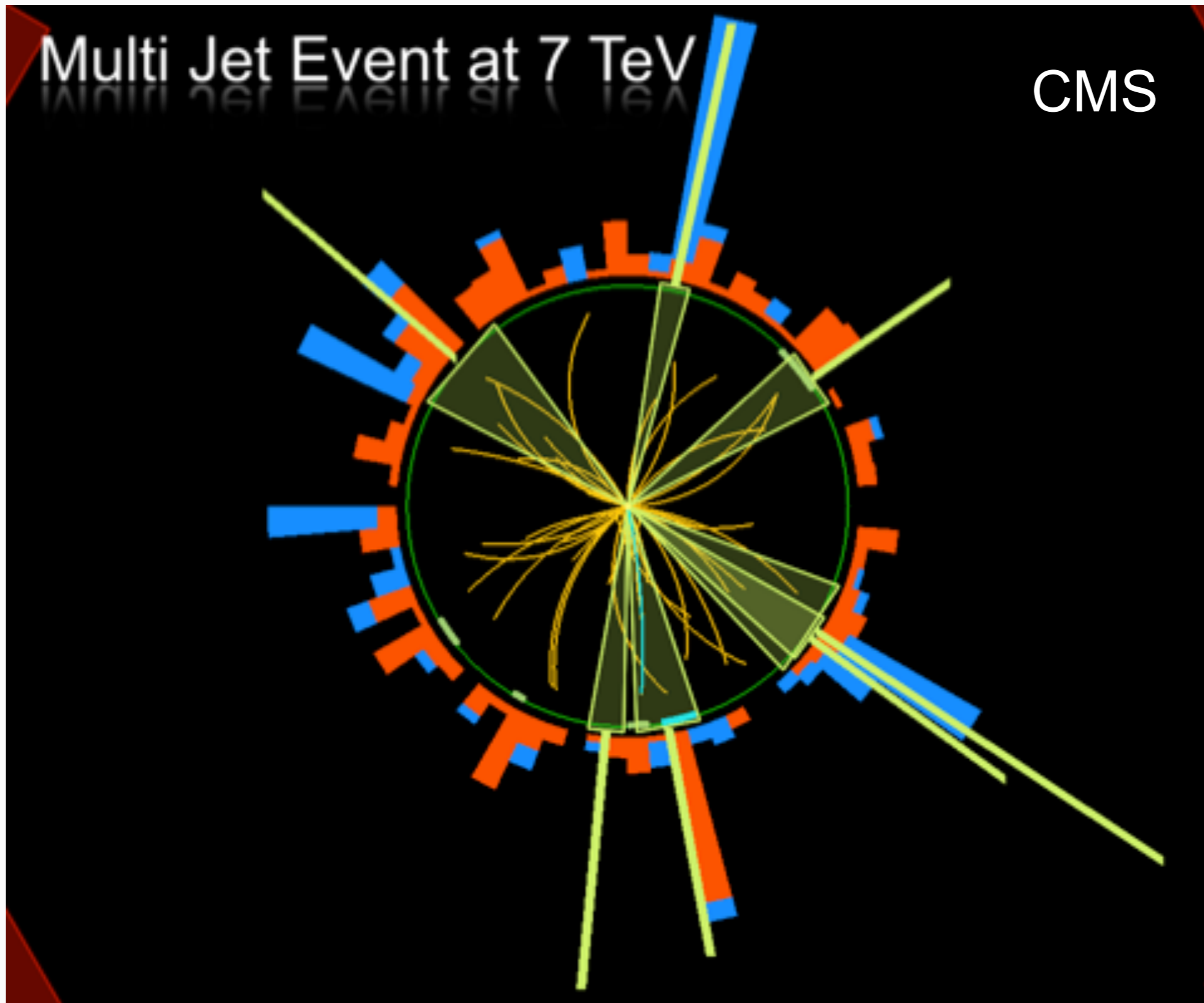


“cluster/spray of particles (tracks, calorimeter deposits) or flow of energy in a restricted angular region”



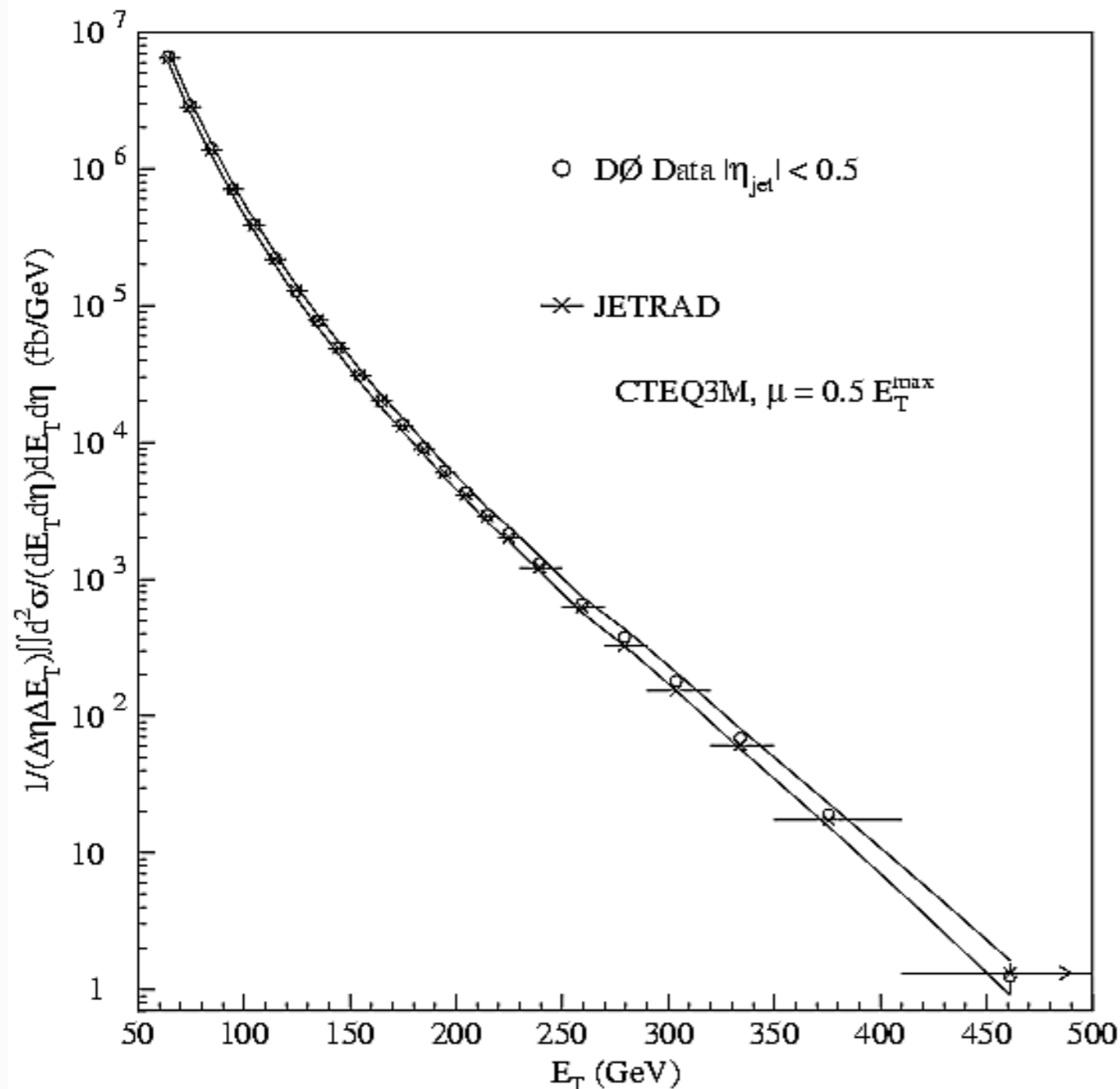
“cluster/spray of particles (tracks, calorimeter deposits) or flow of energy in a restricted angular region”

- clear : need some algorithmic definition.

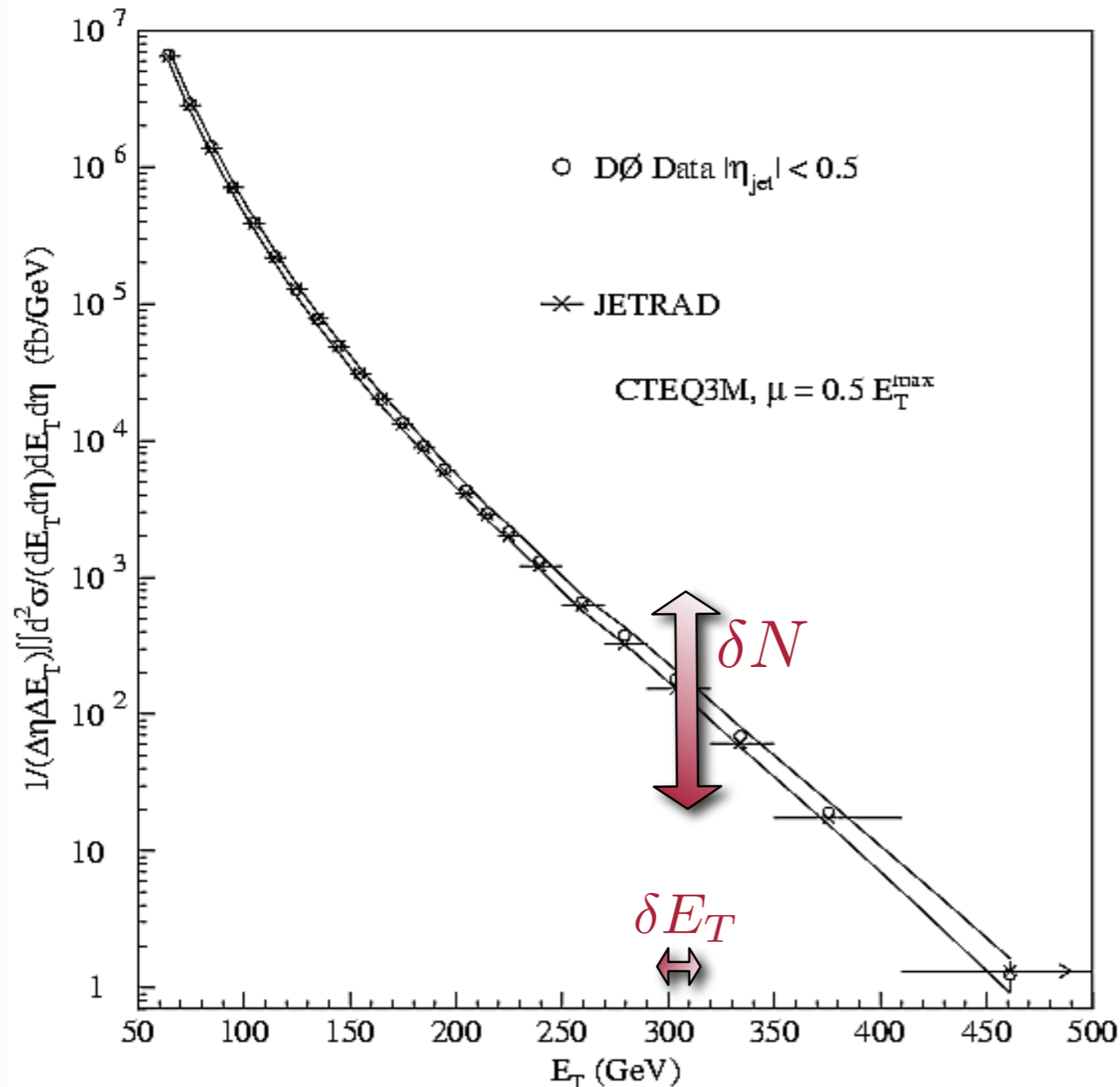


- clear : need some algorithmic definition.

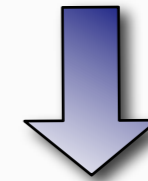
- Question : how well do we know the **energy calibration**?
- Critical because of very steeply falling spectrum!



- Question : how well do we know the **energy calibration**?
- Critical because of very steeply falling spectrum!



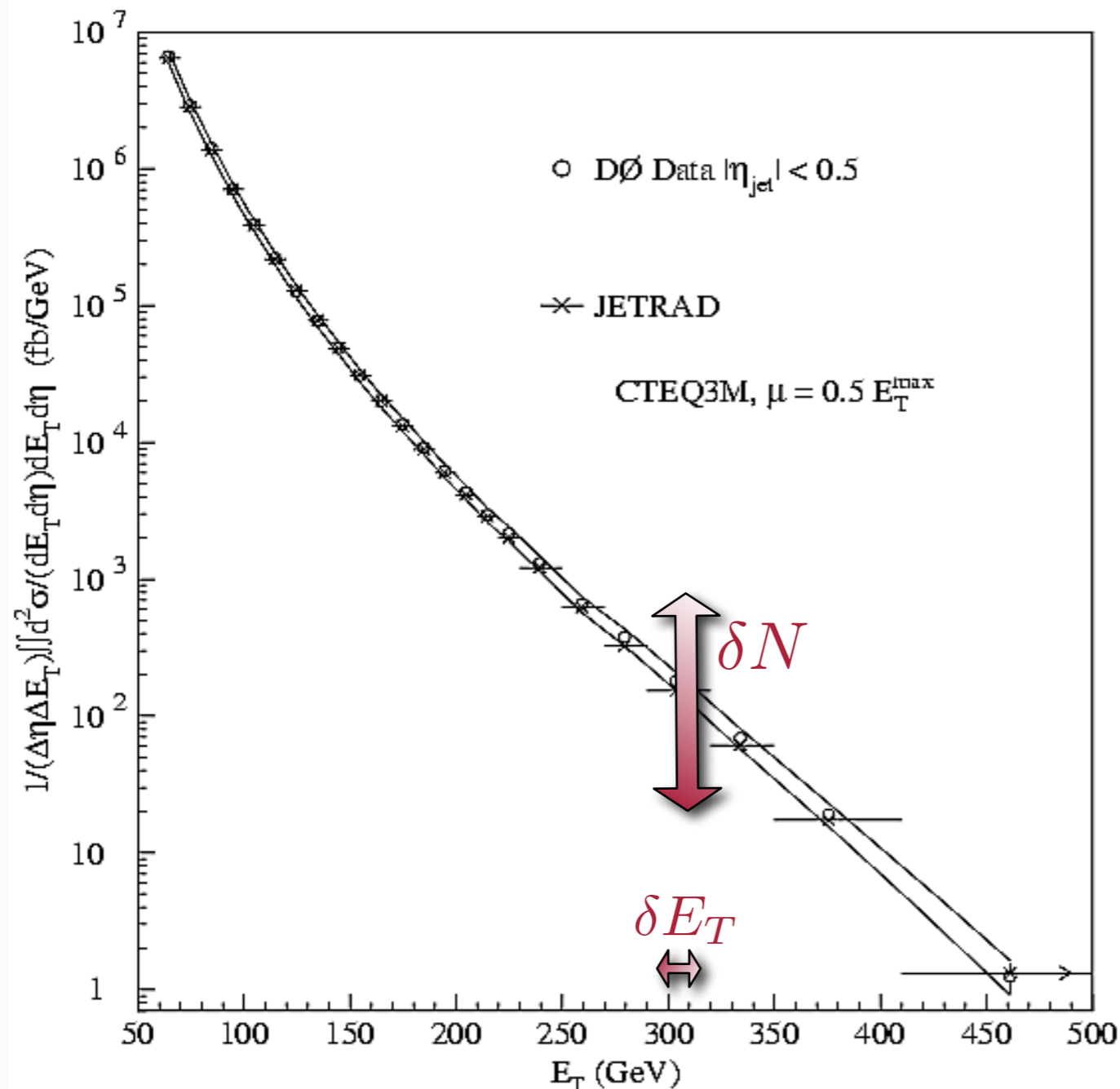
$$\frac{d^2\sigma}{dE_T d\eta} \approx \text{const} \cdot E_T^{-6}$$



relative uncertainties

$$\frac{\delta N}{N} \approx 6 \cdot \frac{\delta E_T}{E_T}$$

- Question : how well do we know the **energy calibration**?
- Critical because of very steeply falling spectrum!



$$\frac{d^2\sigma}{dE_T d\eta} \approx \text{const} \cdot E_T^{-6}$$



relative uncertainties

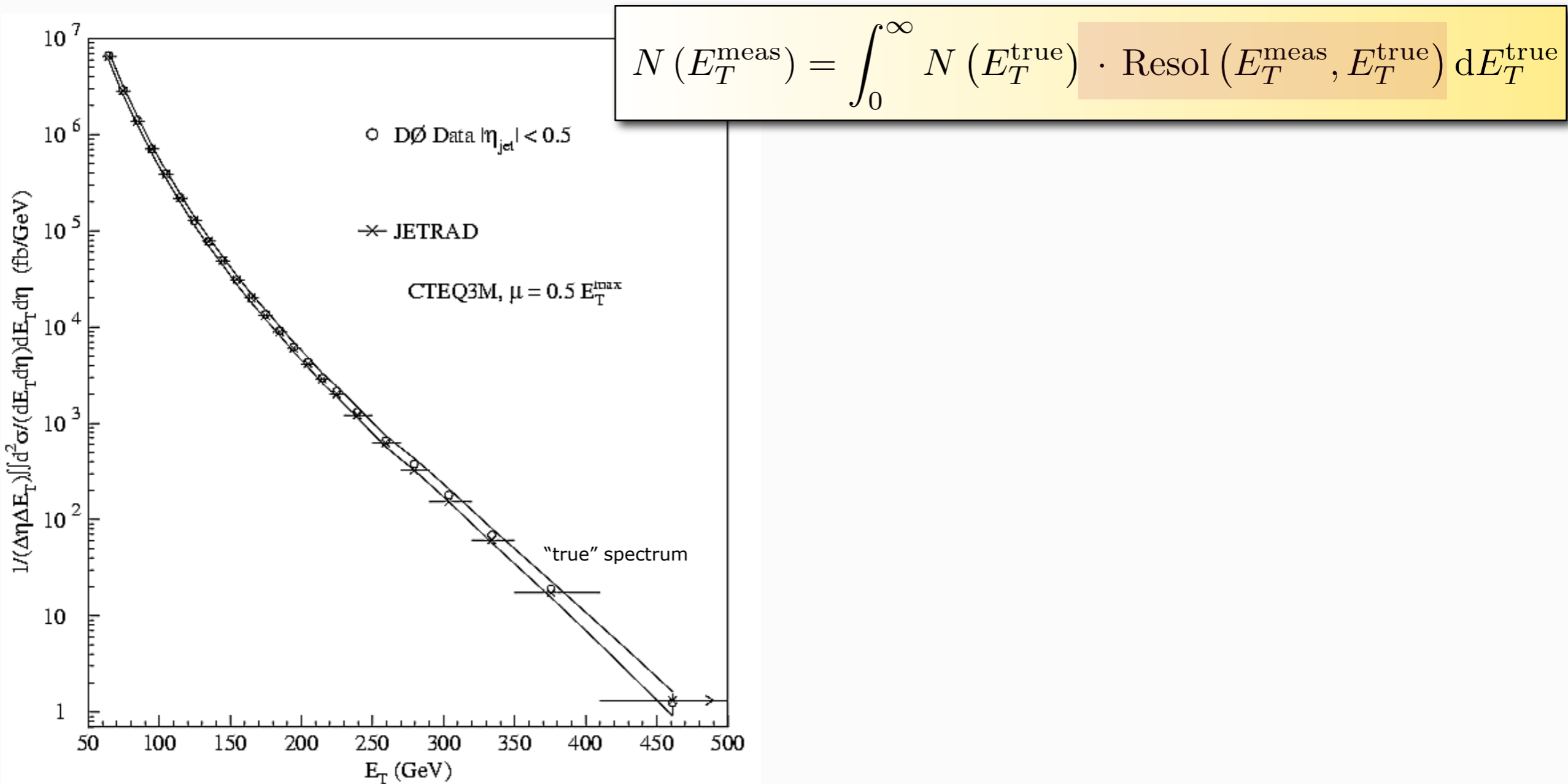
$$\frac{\delta N}{N} \approx 6 \cdot \frac{\delta E_T}{E_T}$$

so beware:

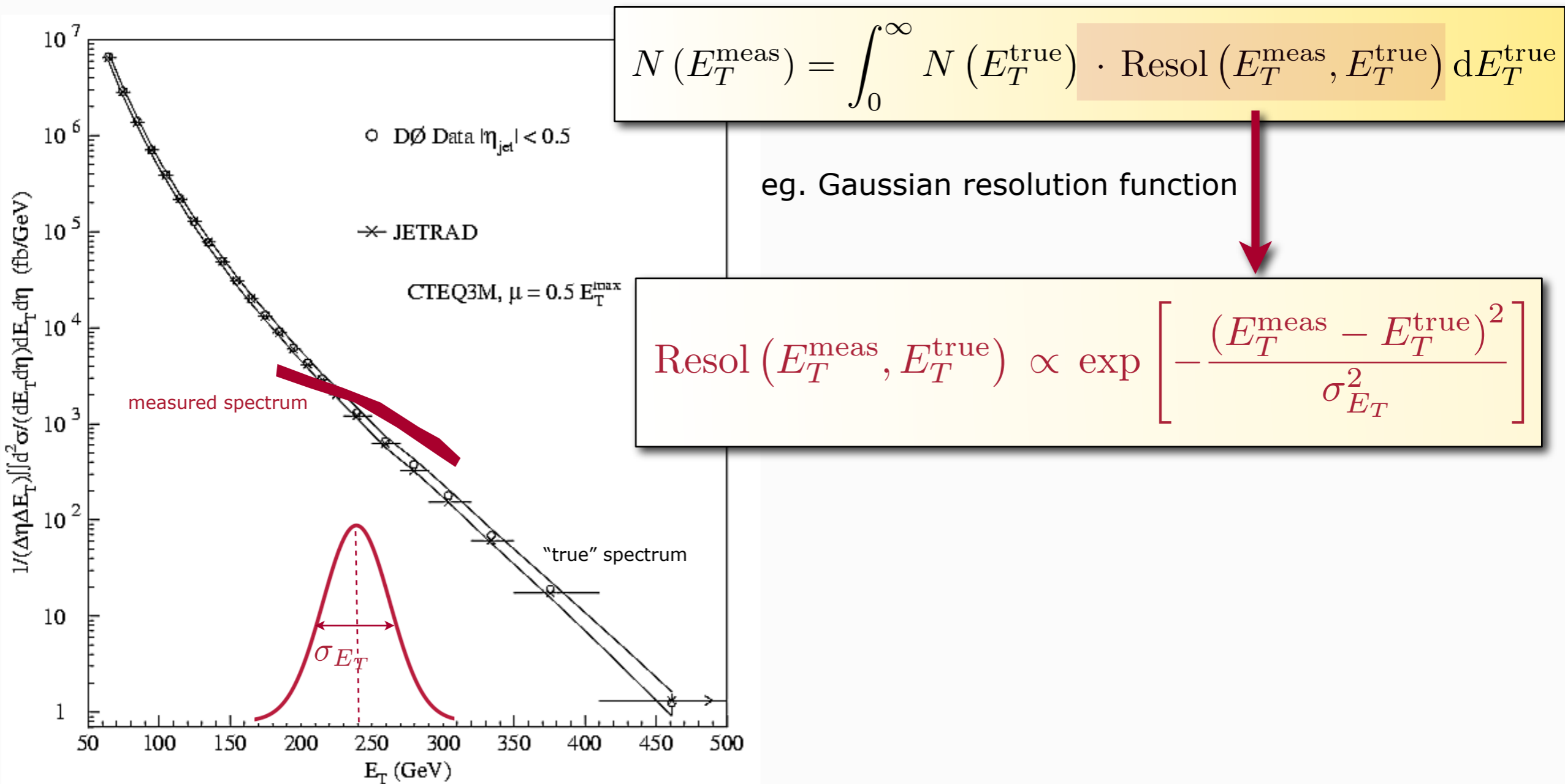
eg. an uncertainty of **5%** on absolute energy scale (calibration)

→ an uncertainty of **30%** (!) on the measured cross section

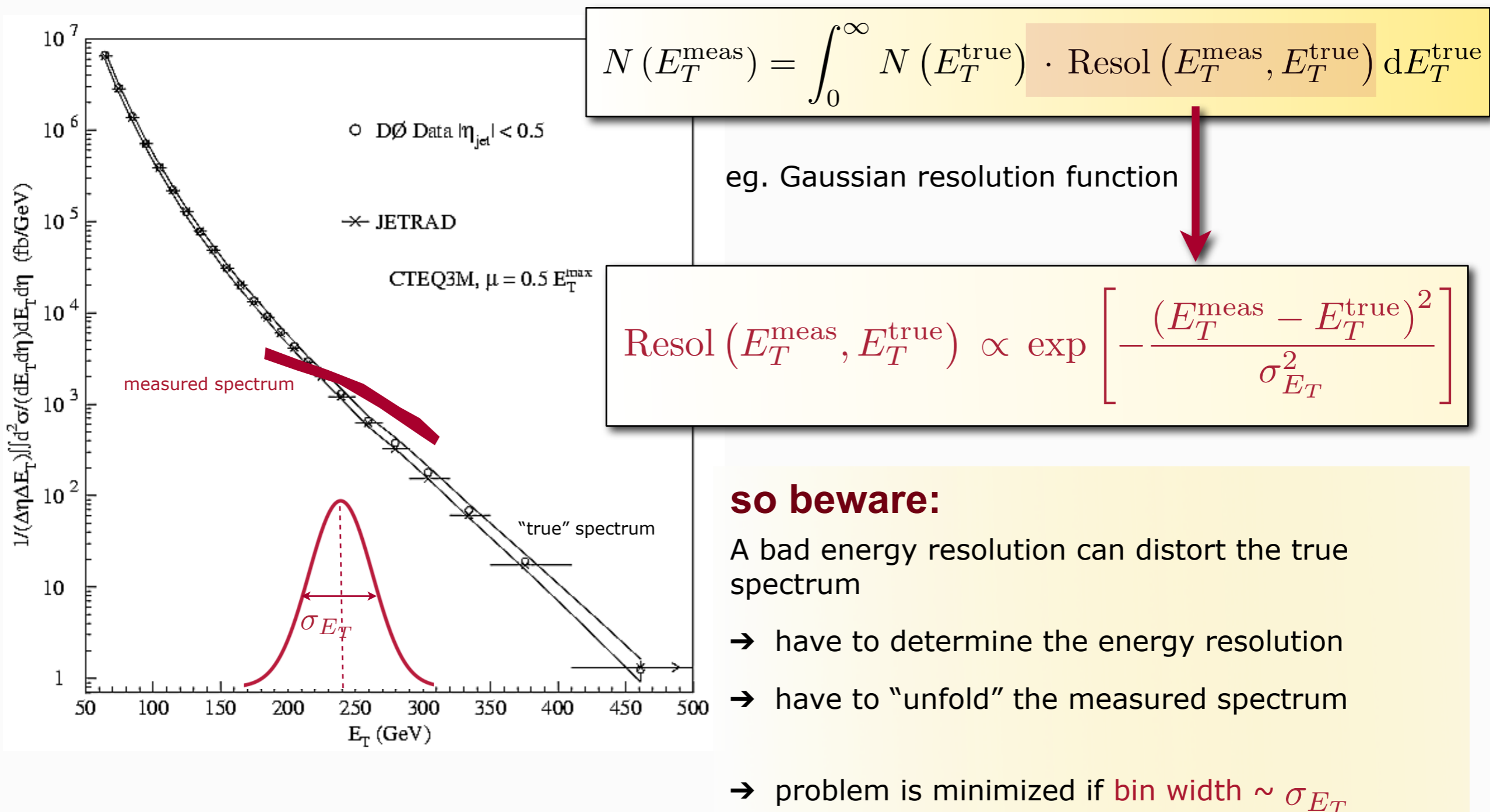
- The energy resolution can distort the spectrum
- Again : Critical because of very steeply falling spectrum!



- The energy resolution can distort the spectrum
- Again : Critical because of very steeply falling spectrum!



- The energy resolution can distort the spectrum
- Again : Critical because of very steeply falling spectrum!



- After data flow from DAQ: data reduction and abstraction

- After data flow from DAQ: data reduction and abstraction
 - **reconstruct** tracks, energy deposits (clusters) in calorimeters

- After data flow from DAQ: data reduction and abstraction
 - **reconstruct** tracks, energy deposits (clusters) in calorimeters
 - calculate “**high-level**” **physics quantities**
 - eg. momentum of charged particles, energy of neutral particles

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation
 - if distributions are measured : take care of absolute calibrations and effects because of detector resolution/smearing
 - correct for these effects

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation
 - if distributions are measured : take care of absolute calibrations and effects because of detector resolution/smearing
 - correct for these effects
 - determine statistical and systematic uncertainties

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation
 - if distributions are measured : take care of absolute calibrations and effects because of detector resolution/smearing
 - correct for these effects
 - determine statistical and systematic uncertainties
 - compare with theory, found a deviation, something new?

- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation
 - if distributions are measured : take care of absolute calibrations and effects because of detector resolution/smearing
 - correct for these effects
 - determine statistical and systematic uncertainties
 - compare with theory, found a deviation, something new?
 - if yes, book the ticket to Stockholm

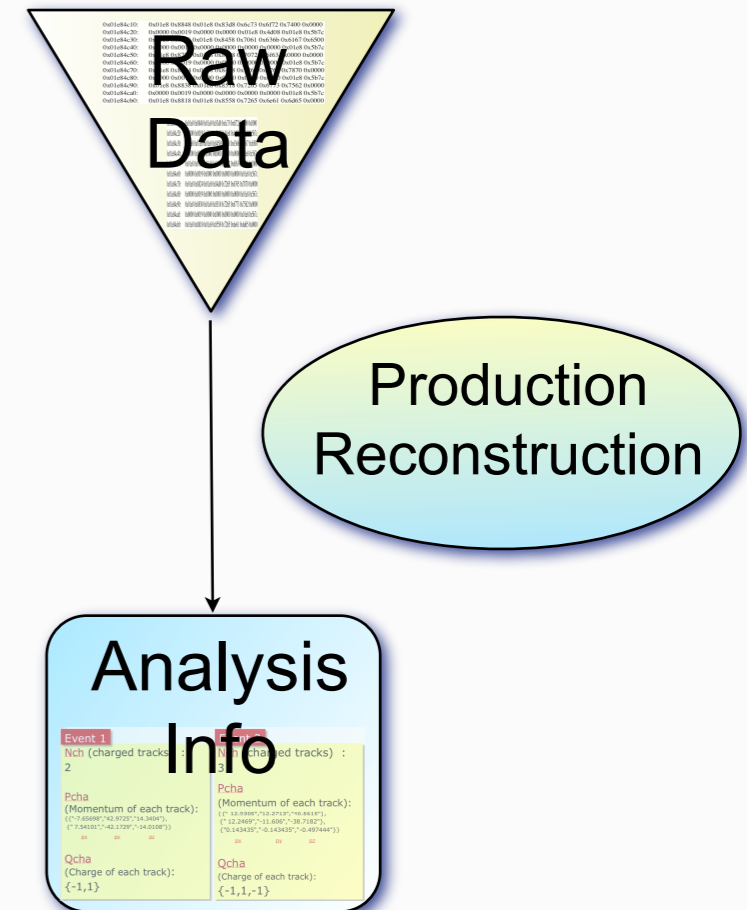
- After data flow from DAQ: data reduction and abstraction
 - reconstruct tracks, energy deposits (clusters) in calorimeters
 - calculate “high-level” physics quantities
 - eg. momentum of charged particles, energy of neutral particles
 - apply even higher-level algorithms, eg. jet finding
 - store all these quantities/objects event per event
- The data analysis
 - define the theoretically computed observable(s) to be measured
 - apply event selection (cuts)
 - estimate efficiencies and backgrounds, eg. from MC simulation
 - if distributions are measured : take care of absolute calibrations and effects because of detector resolution/smearing
 - correct for these effects
 - determine statistical and systematic uncertainties
 - compare with theory, found a deviation, something new?
 - if yes, book the ticket to Stockholm
 - determine parameters, eg. by fitting the prediction to the data

How is all this done in practice?

“The only place you’ll find SUCCESS before WORK is in the dictionary” (May B. Smith)

The reconstruction step is usually done in common

- “Tracks”, “particle ID”, “calorimeter towers” etc are general concepts, not analysis-specific. Common algorithms make it easier to understand how well they work
- “very coordinated” data access

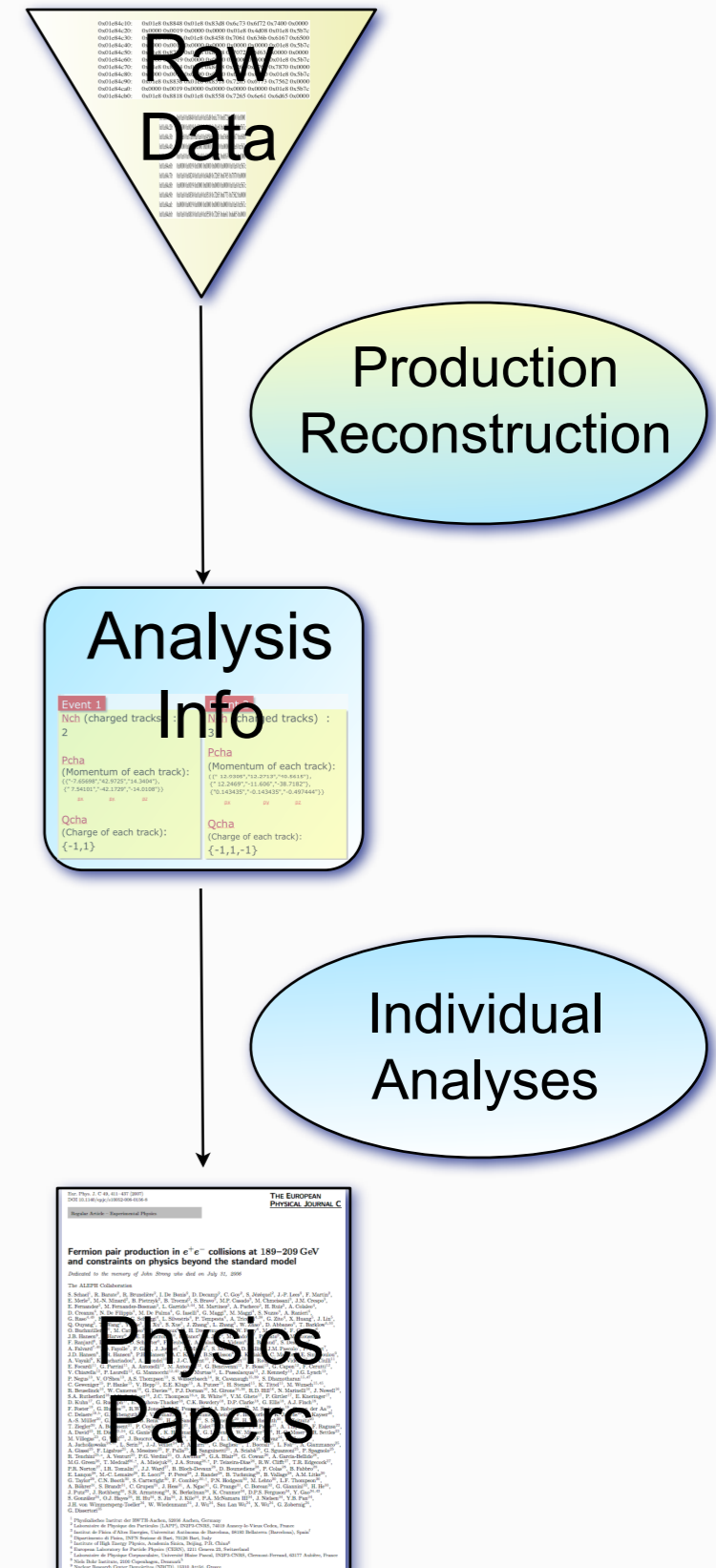


The reconstruction step is usually done in common

- “Tracks”, “particle ID”, “calorimeter towers” etc are general concepts, not analysis-specific. Common algorithms make it easier to understand how well they work
- “very coordinated” data access

Analysis is a very individual thing

- Many different measurements being done at once
- Small groups working on topics they are interested in
- Many different time scales for these efforts
- “chaotic” data access



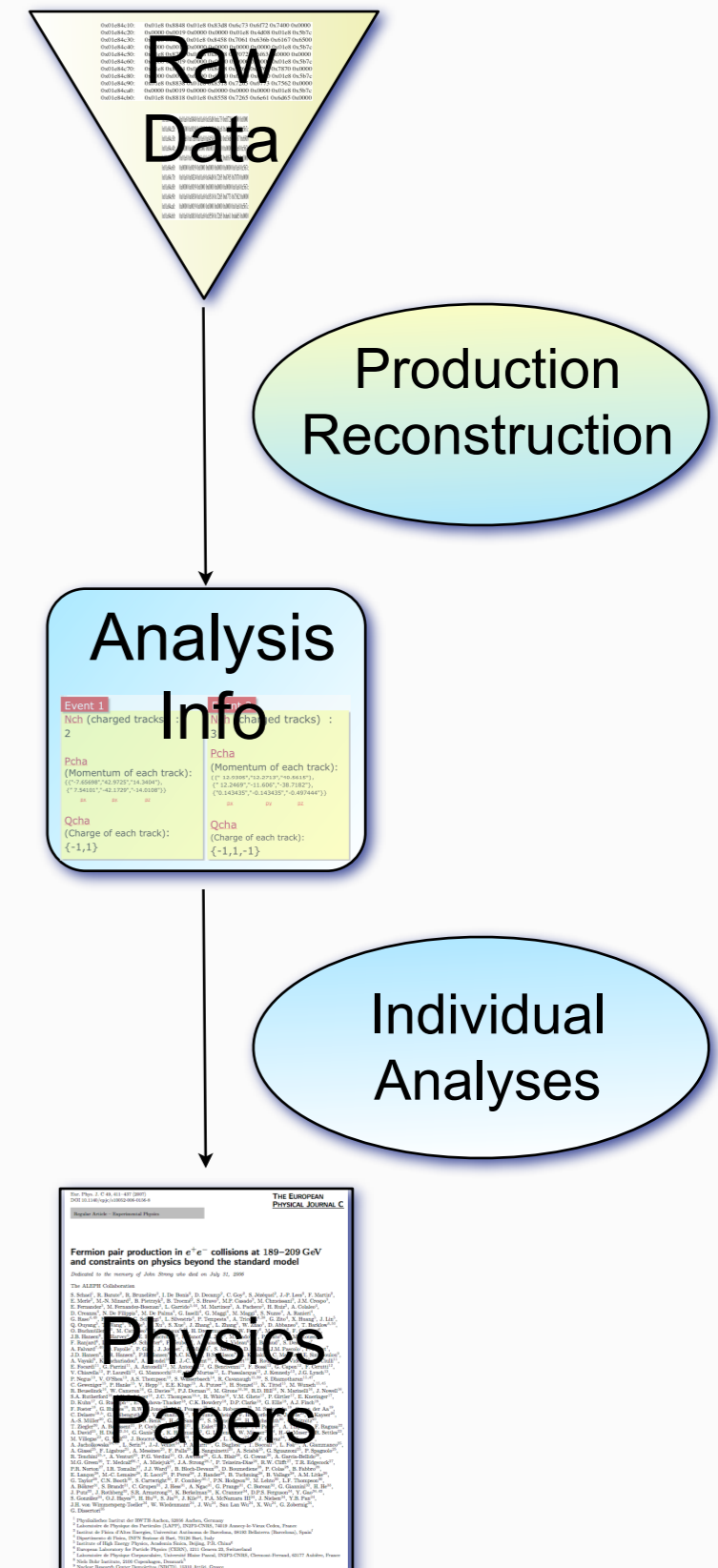
The reconstruction step is usually done in common

- “Tracks”, “particle ID”, “calorimeter towers” etc are general concepts, not analysis-specific. Common algorithms make it easier to understand how well they work
- “very coordinated” data access

Analysis is a very individual thing

- Many different measurements being done at once
- Small groups working on topics they are interested in
- Many different time scales for these efforts
- “chaotic” data access

Collaborations build offline computing systems to handle all this



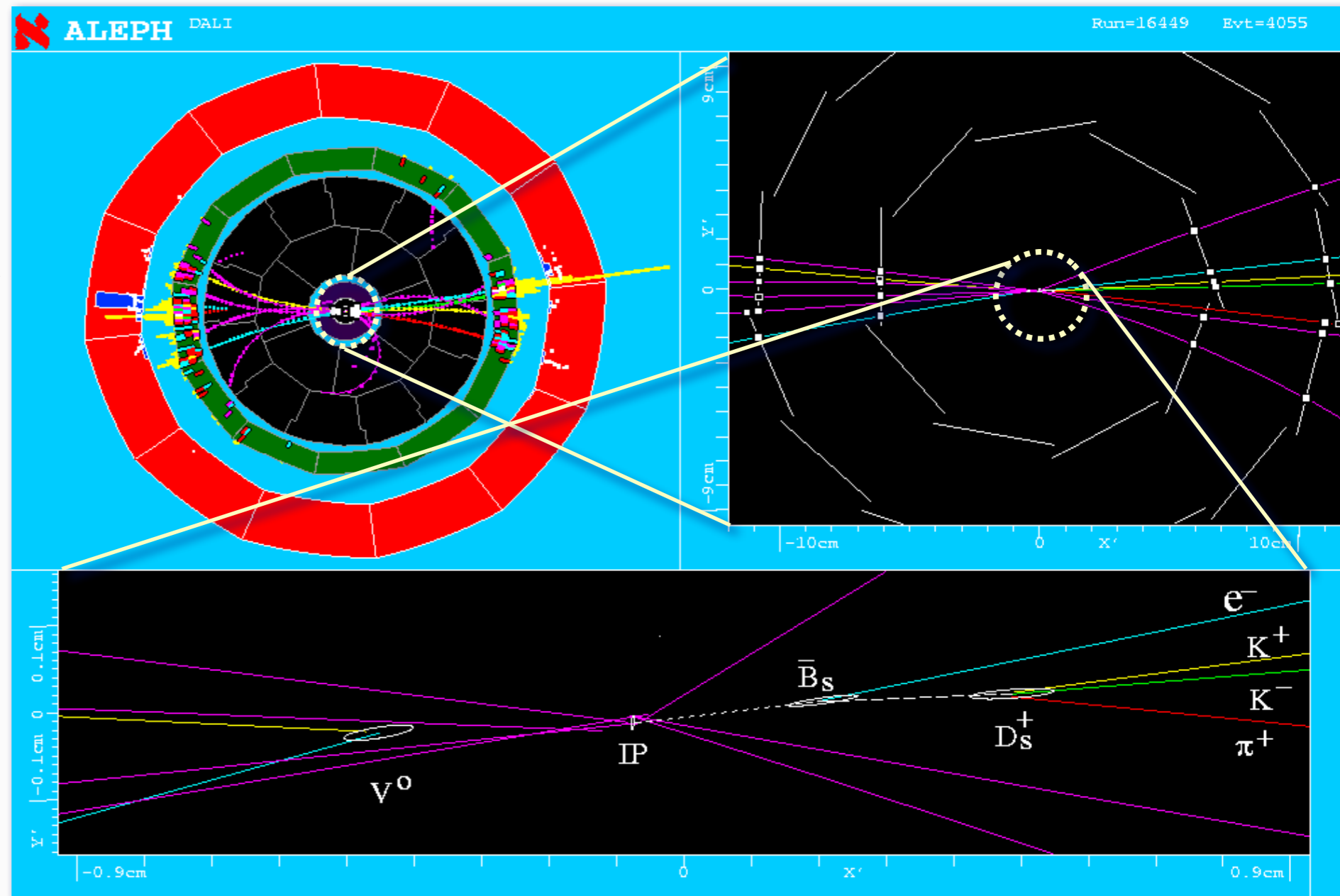
Track finding

- Determine how many charged particles were created in an event
- Measure their momentum
 - direction, magnitude
 - combine these to look for decays of particles with known masses
 - only final stable particles are visible

- Determine how many charged particles were created in an event
- Measure their momentum
 - direction, magnitude
 - combine these to look for decays of particles with known masses
 - only final stable particles are visible

● Measure spatial trajectories

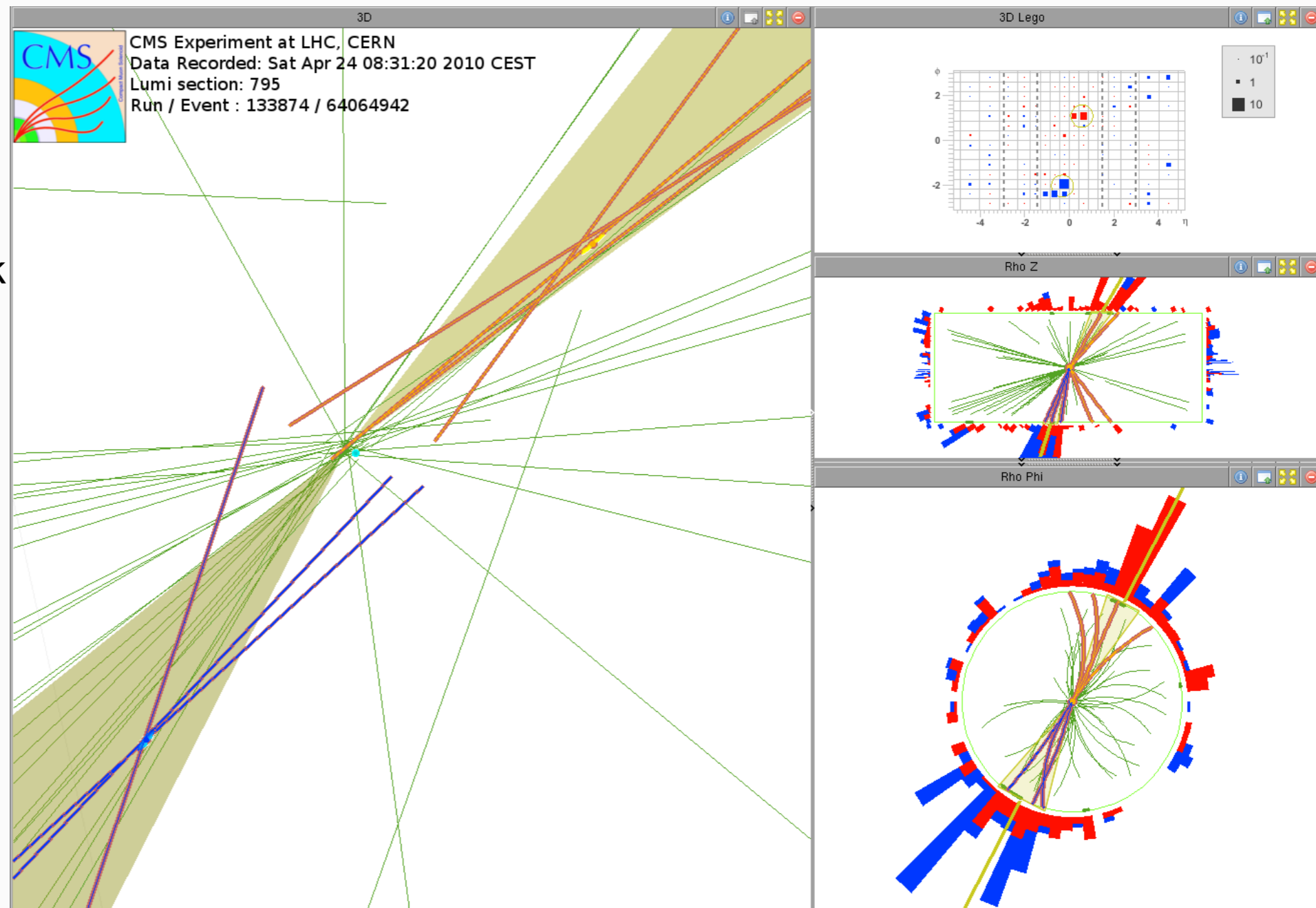
- combine to look for separated vertices, indicating particles with long lifetimes



- Determine how many charged particles were created in an event
- Measure their momentum
 - direction, magnitude
 - combine these to look for decays of particles with known masses
 - only final stable particles are visible

Measure spatial trajectories

- combine to look for separated vertices, indicating particles with long lifetimes



G. Dissertori : From raw data to physics results

- 1D straight line fit as simple case

- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate



- Imperfect** measurements give less precise results

- 1D straight line fit as simple case

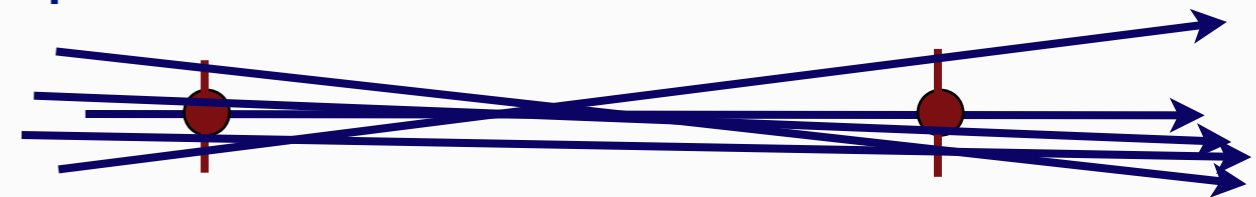
- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate



- Imperfect** measurements give less precise results

- the farther you extrapolate, the less you know



- 1D straight line fit as simple case

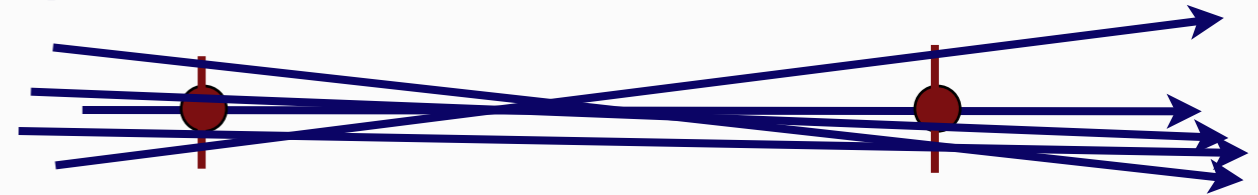
- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate

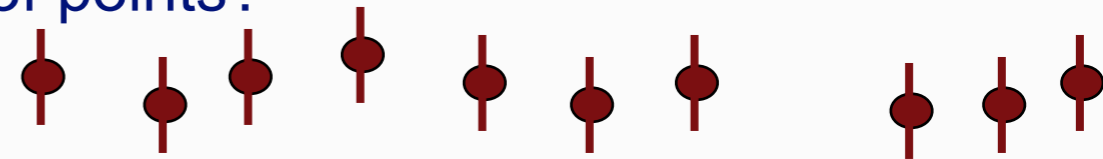


- Imperfect** measurements give less precise results

- the farther you extrapolate, the less you know



- Smaller errors and more points help to constrain the possibilities.
But how to find the best point from a large set of points?



- 1D straight line fit as simple case

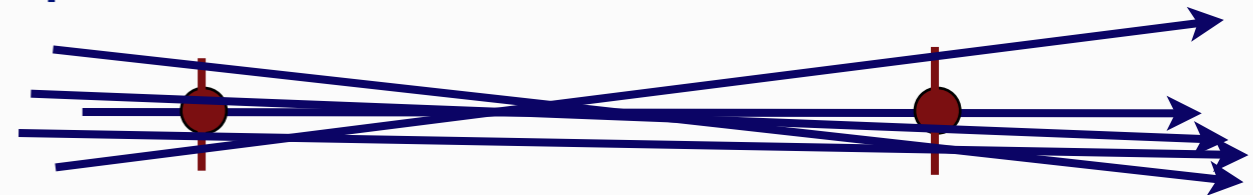
- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate

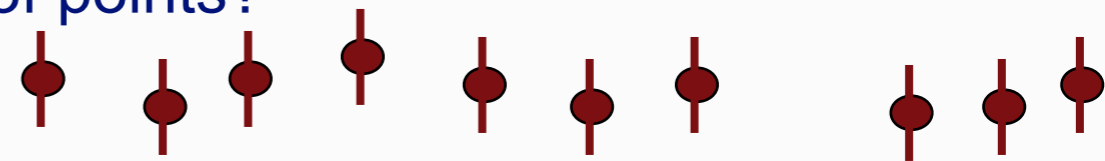


- Imperfect** measurements give less precise results

- the farther you extrapolate, the less you know



- Smaller errors and more points help to constrain the possibilities.
But how to find the best point from a large set of points?



- Quantitatively**

- parameterize a track:
- In case of straight line $y(x) = \theta x + d$ or, eg., **helix** in case of magnetic field present

- 1D straight line fit as simple case

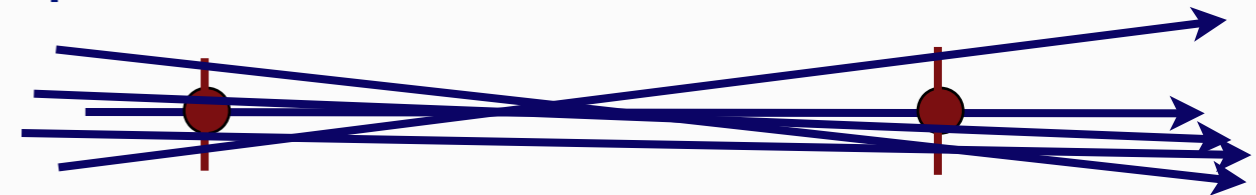
- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate

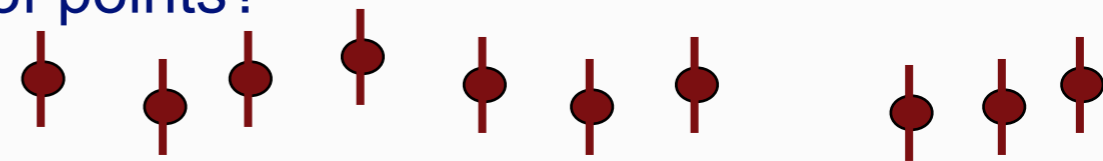


- Imperfect** measurements give less precise results

- the farther you extrapolate, the less you know



- Smaller errors and more points help to constrain the possibilities. But how to find the best point from a large set of points?



- Quantitatively**

- parameterize a track:
- In case of straight line $y(x) = \theta x + d$ or, eg., **helix** in case of magnetic field present

- Find track parameters by **Least-Squares-Minimization**

$$\chi^2 = \sum_{i=1}^{n_{\text{hits}}} \frac{(y_i - y(x_i))^2}{\sigma_i^2}$$

position of i^{th} hit $\rightarrow x_i$
 predicted track position at i^{th} hit $\rightarrow y(x_i)$
 uncertainty of i^{th} measurement $\rightarrow \sigma_i$

see also lecture by G. Cowan

- 1D straight line fit as simple case

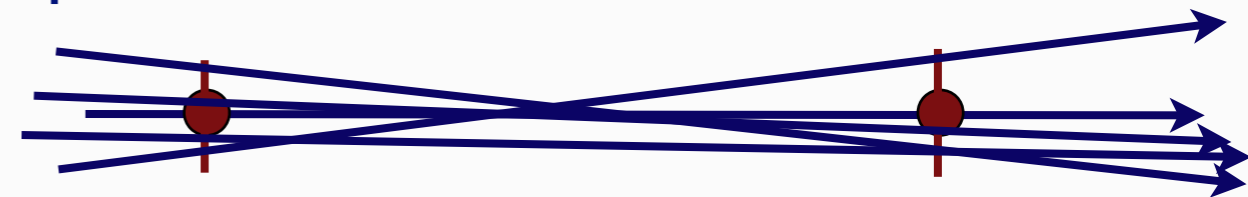
- Two **perfect** measurements

- away from interaction point
- no measurement uncertainty
- just draw a straight line through them and extrapolate

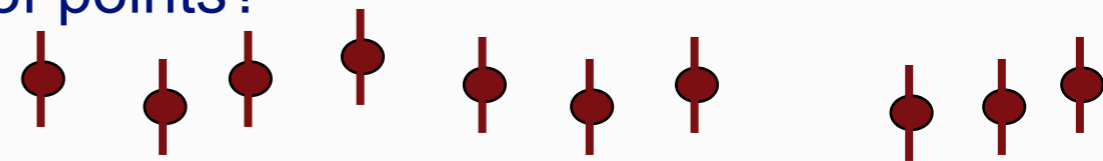


- Imperfect** measurements give less precise results

- the farther you extrapolate, the less you know



- Smaller errors and more points help to constrain the possibilities. But how to find the best point from a large set of points?



- Quantitatively**

- parameterize a track:
- In case of straight line $y(x) = \theta x + d$ or, eg., **helix** in case of magnetic field present

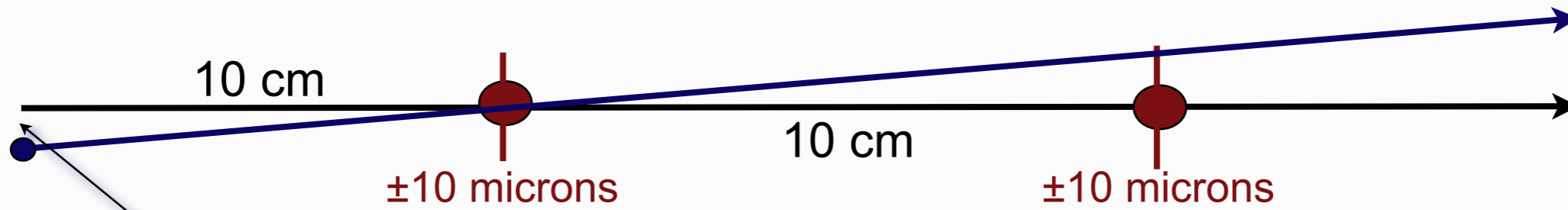
- Find track parameters by **Least-Squares-Minimization**
- Obtain also uncertainties on track parameters

$$\delta\theta \quad \delta d$$

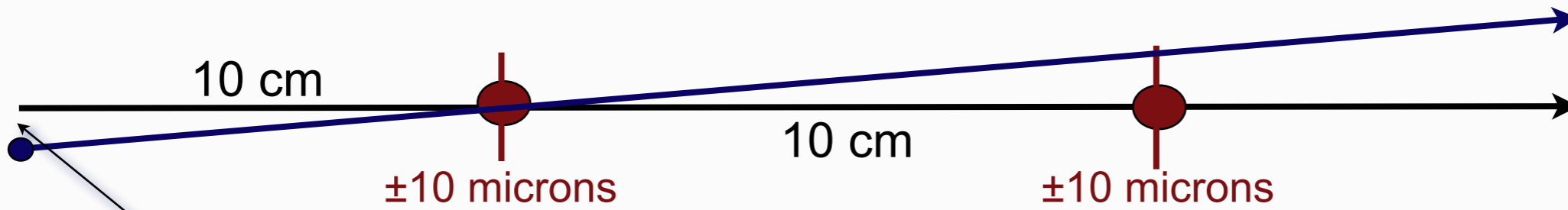
$$\chi^2 = \sum_{i=1}^{n_{\text{hits}}} \frac{(y_i - y(x_i))^2}{\sigma_i^2}$$

position of i^{th} hit
predicted track position at i^{th} hit

uncertainty of i^{th} measurement



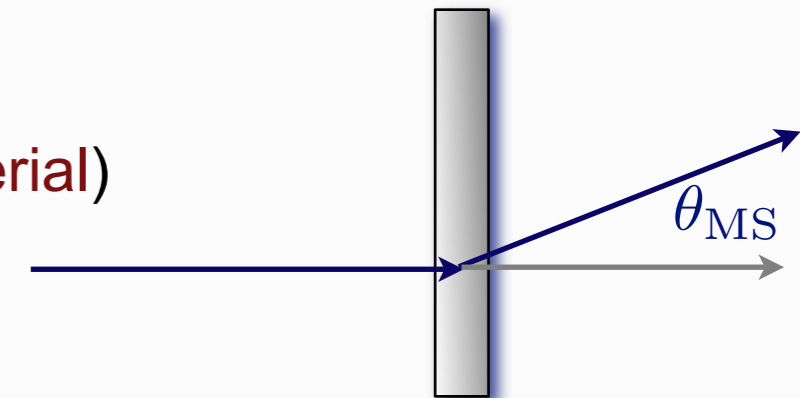
- Error δd on position is about ± 10 microns
- Error $\delta \theta$ on angle is about ± 0.1 milliradians (± 0.002 degrees)
- Satisfyingly small errors
 - allows separation of tracks that come from different particle decays (which can be separated at the order of mm)
- However
 - we “see” particles by interaction with a detector (=material)

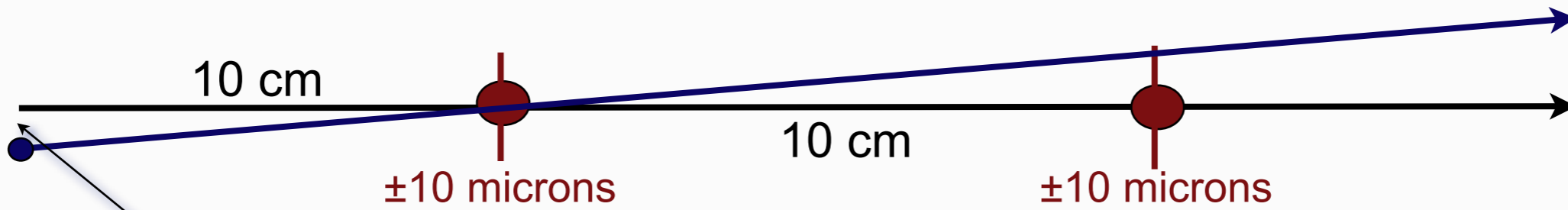


- Error δd on position is about ± 10 microns
- Error $\delta \theta$ on angle is about ± 0.1 milliradians (± 0.002 degrees)
- Satisfyingly small errors
 - allows separation of tracks that come from different particle decays (which can be separated at the order of mm)

However

- we “see” particles by interaction with a detector (=material)
- interaction leads to : energy loss, change in direction

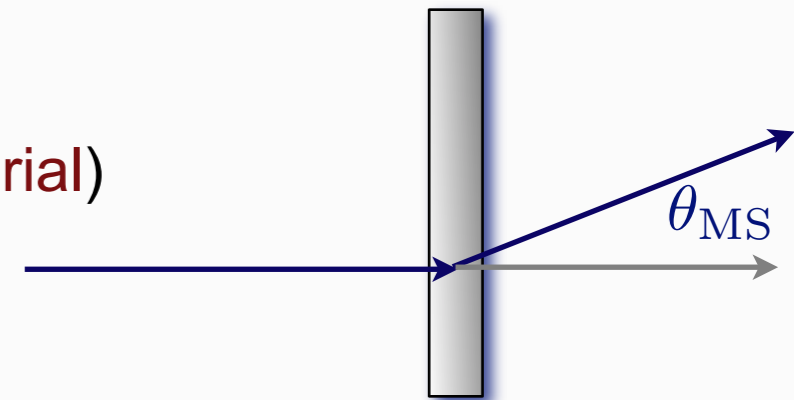




- Error δd on position is about ± 10 microns
- Error $\delta\theta$ on angle is about ± 0.1 milliradians (± 0.002 degrees)
- Satisfyingly small errors
 - allows separation of tracks that come from different particle decays (which can be separated at the order of mm)

However

- we “see” particles by interaction with a detector (=material)
- interaction leads to : energy loss, change in direction
- This is **Multiple Scattering**



- Charged particles passing through matter “scatter” by a random angle

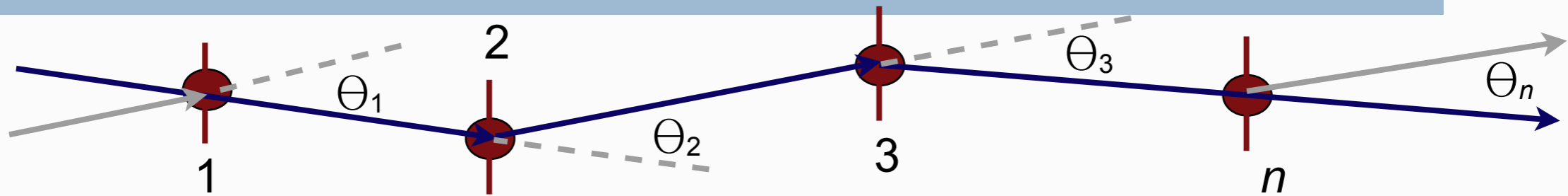
$$\sqrt{\langle \theta_{MS}^2 \rangle} = \frac{15 \text{ MeV}/c}{\beta p} \sqrt{\frac{\text{thickness}}{X_{rad}}}$$

examples:

300 micron Si : RMS = 0.9 mrad / βp
 1 mm Be : RMS = 0.8 mrad / βp

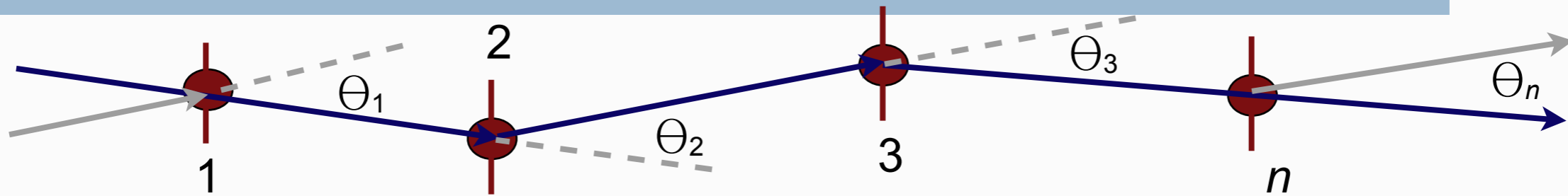
→ leads to additional position errors

Kalman filter



- So? Could extend track parameterization to take this into account
 - n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_1(x - x_2) + \dots$
- and include the multiple scattering information into the Least-Squares (n equations, n unknowns)
 - For large n , computing time grows like $O(n^3)$, quickly un-practicable

Kalman filter



So? Could extend track parameterization to take this into account

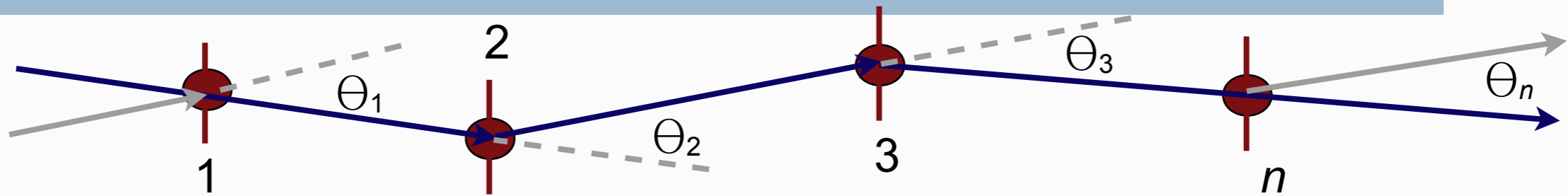
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

Kalman filter



So? Could extend track parameterization to take this into account

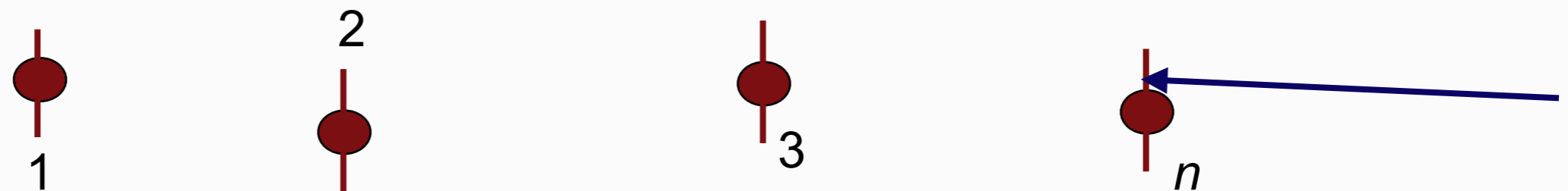
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

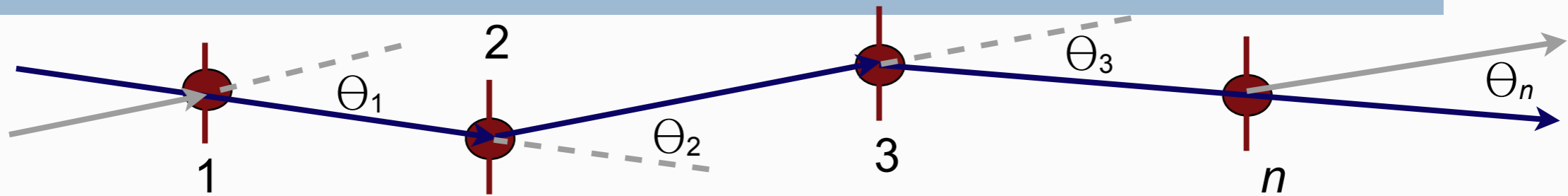
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

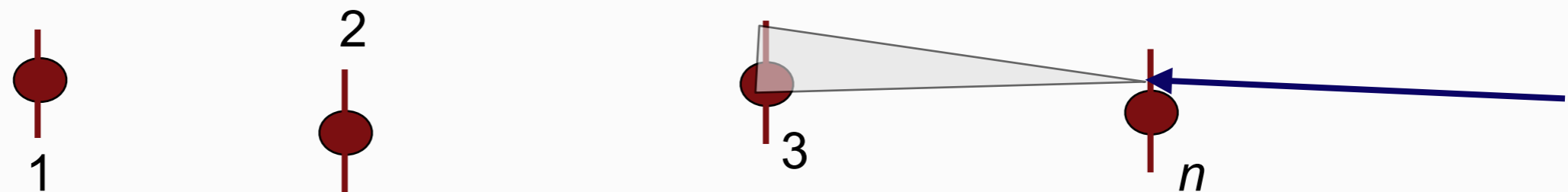
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

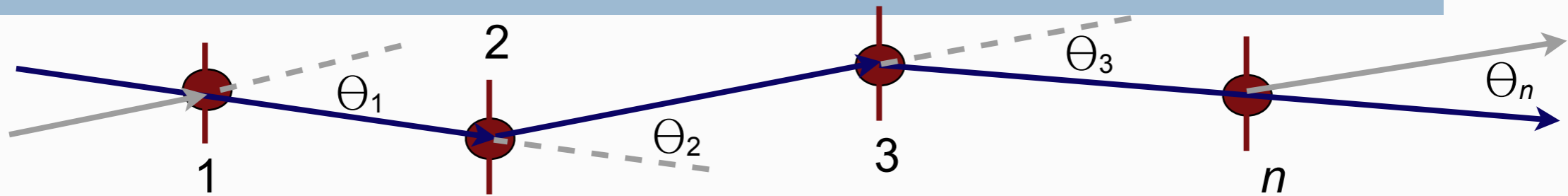
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

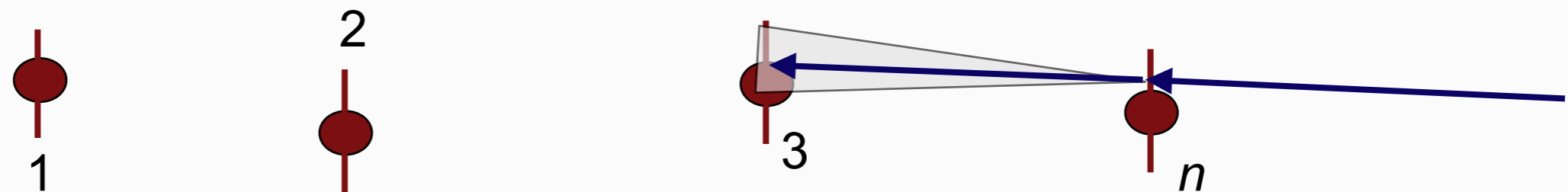
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

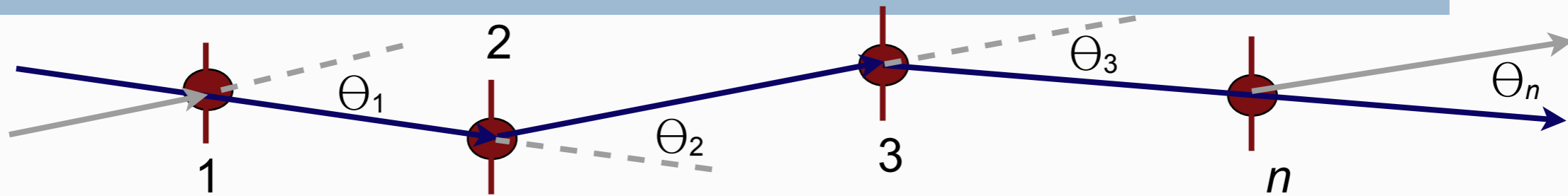
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

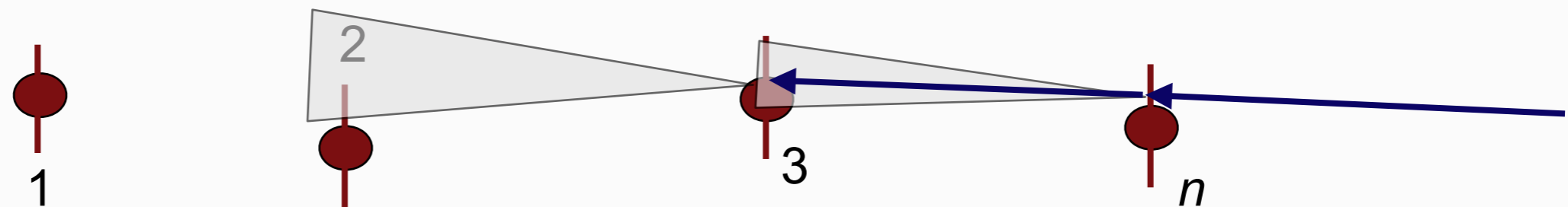
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

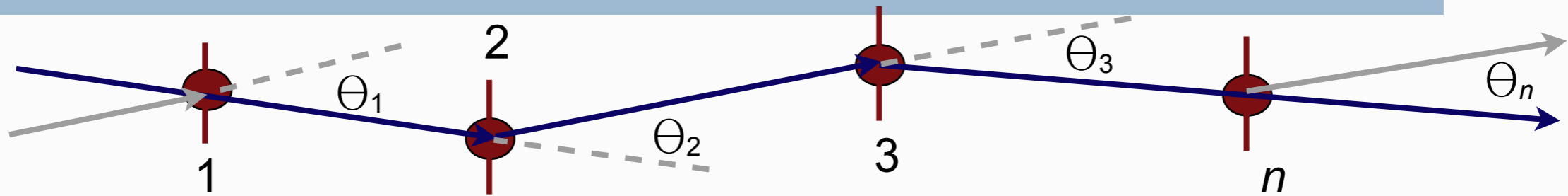
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

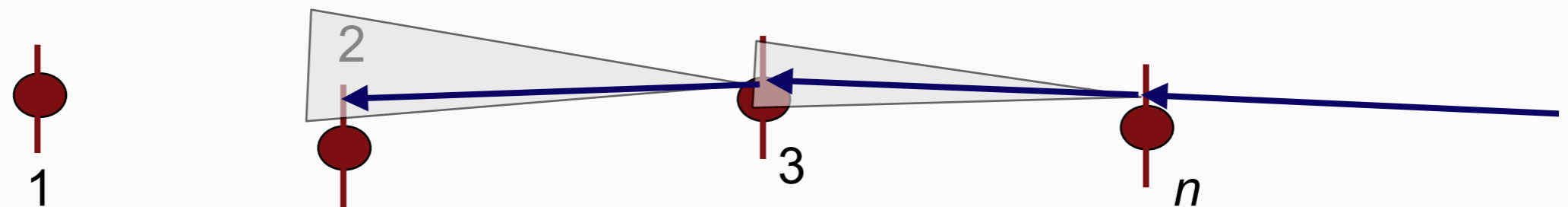
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

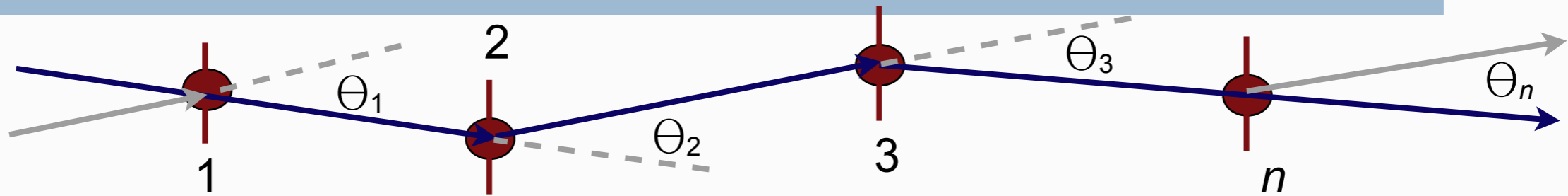
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

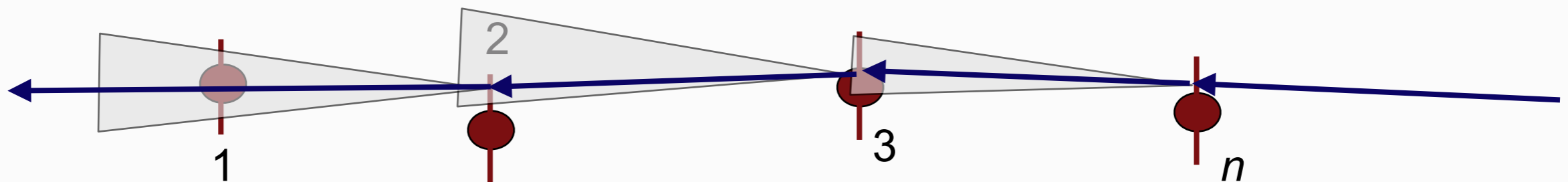
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

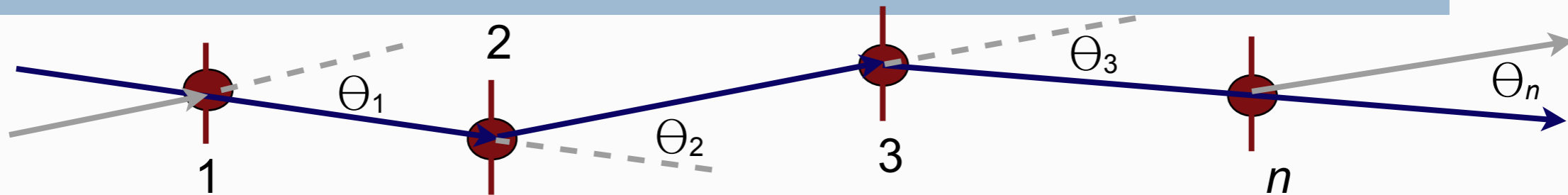
$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

Instead, approximate, work inward N times



Kalman filter



So? Could extend track parameterization to take this into account

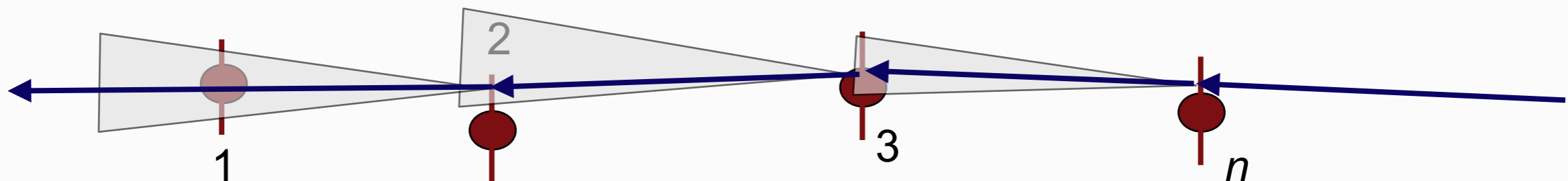
n additional parameters $y(x) = d + \theta x + \Theta(x - x_1)\theta_1(x - x_1) + \Theta(x - x_2)\theta_2(x - x_2) + \dots$

and include the multiple scattering information into the Least-Squares (n equations, n unknowns)

$$\chi^2 = \chi_{\text{old}}^2 + \sum_i \frac{\theta_i^2}{\sigma_{\text{MS}}^2}$$

- For large n , computing time grows like $O(n^3)$, quickly un-practicable
- anyway, not interested in all these angles, only in parameters at the origin

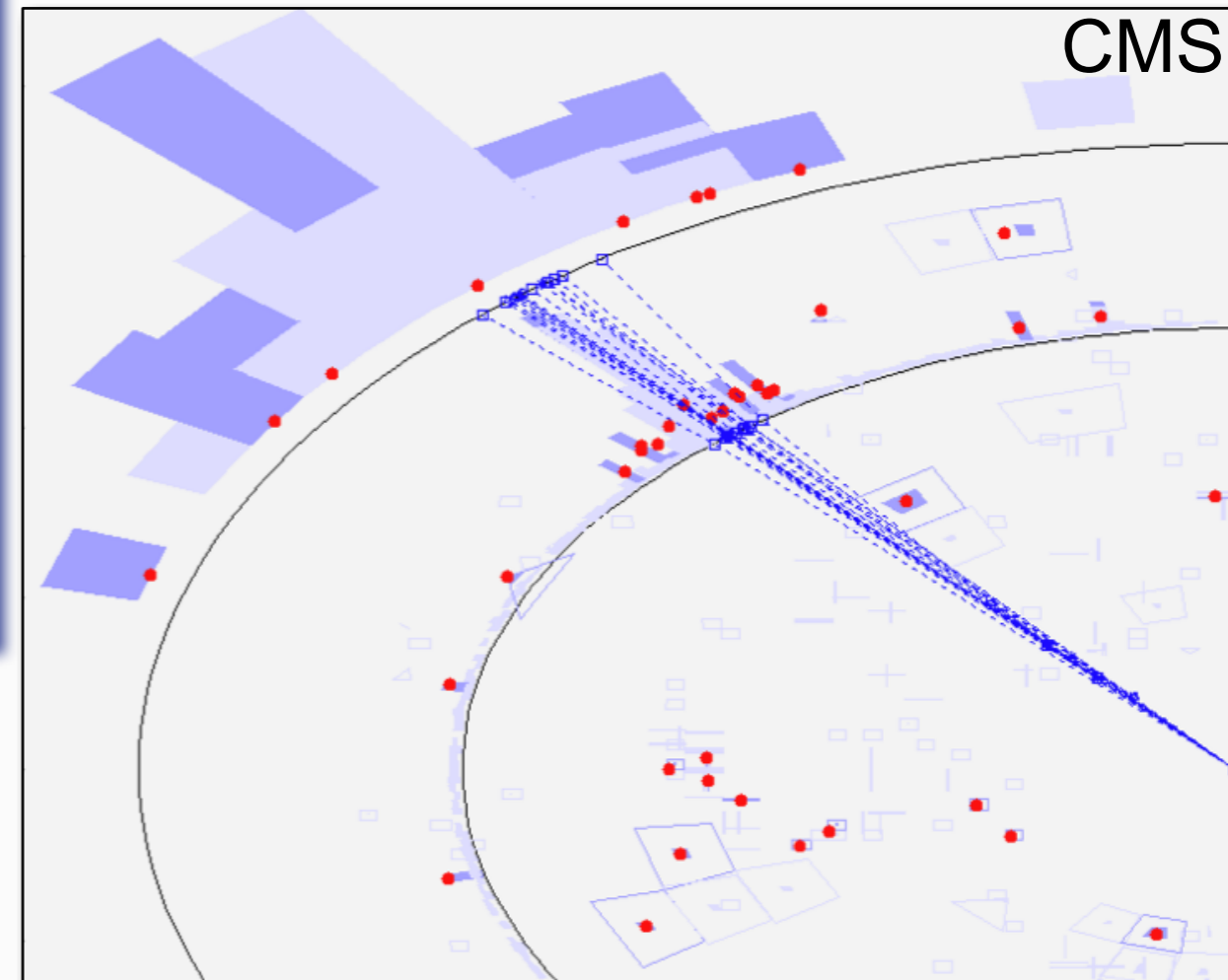
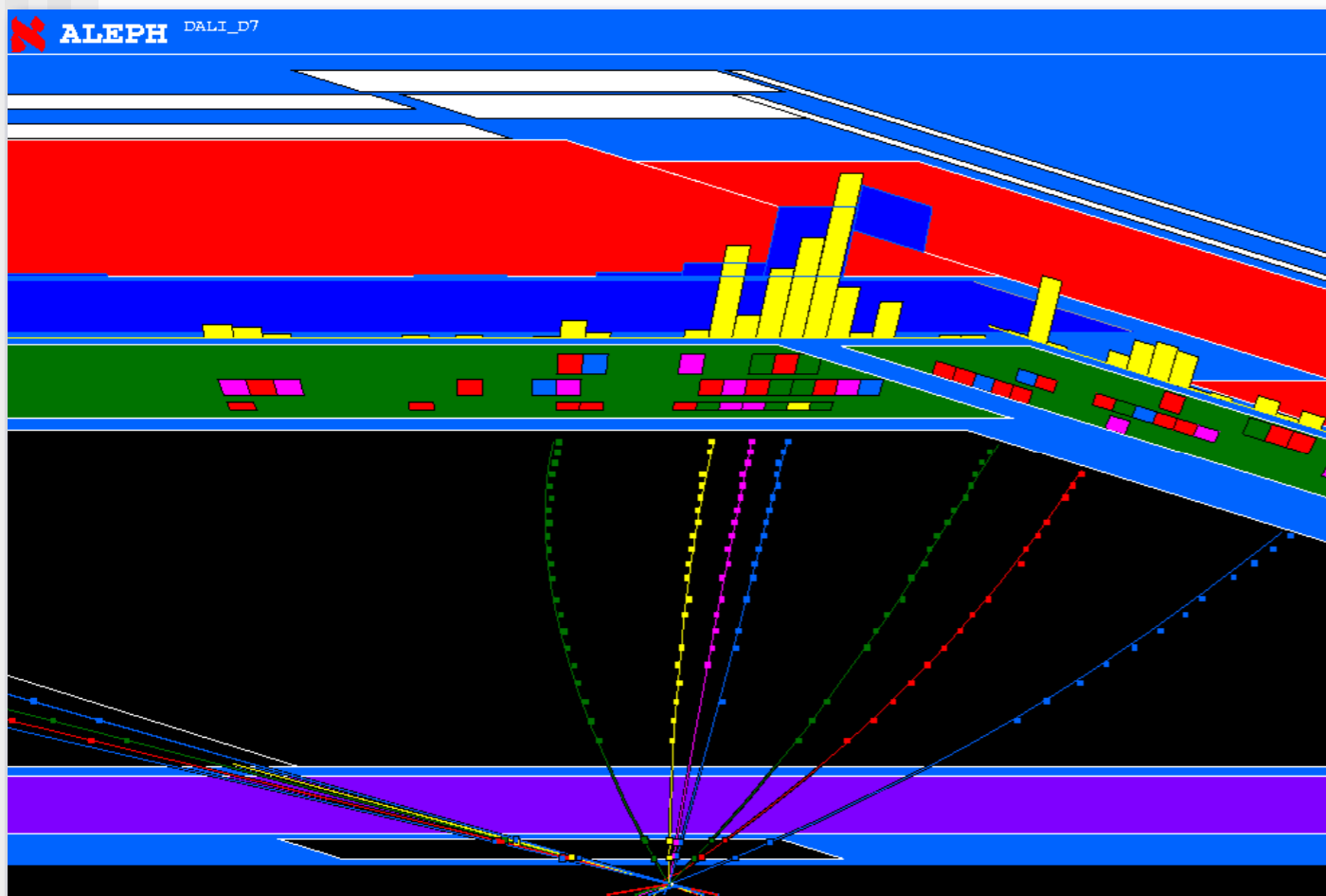
Instead, approximate, work inward N times



- leads to $O(n)$ computations!
- in each step, make extrapolation to next layer, using information from current track parameters, expected scattering error, and measurement at next layer
- Needs a starting estimate (**seed**) and may need some **iterations**, smoothing
- This method is based on theory of the **Kalman Filter**

Calorimeter energy reconstruction

- Reconstruct energy deposited by charged and neutral particles
- Determine position of deposit, direction of incident particles
- Be insensitive to noise and “un-wanted” (un-correlated) energy

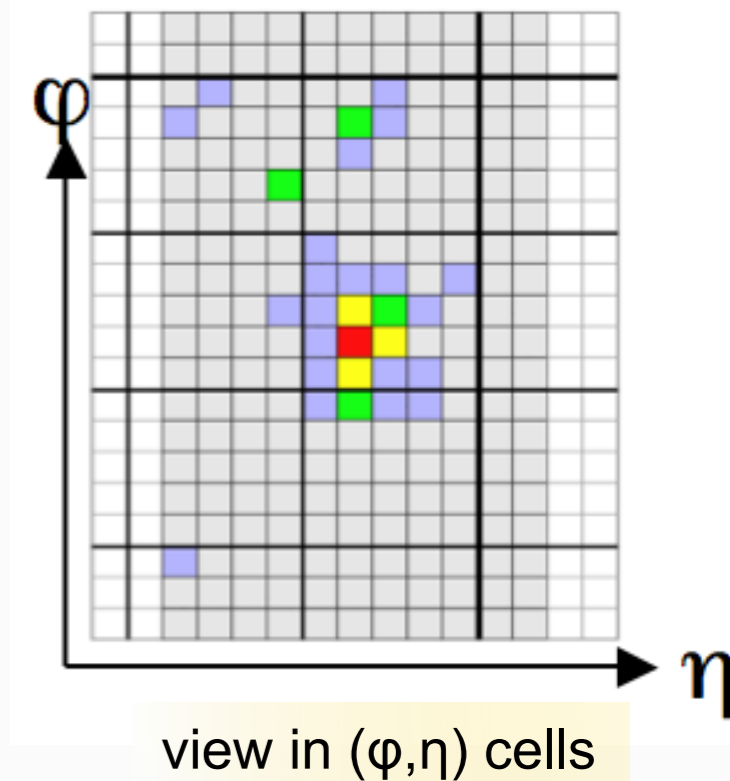
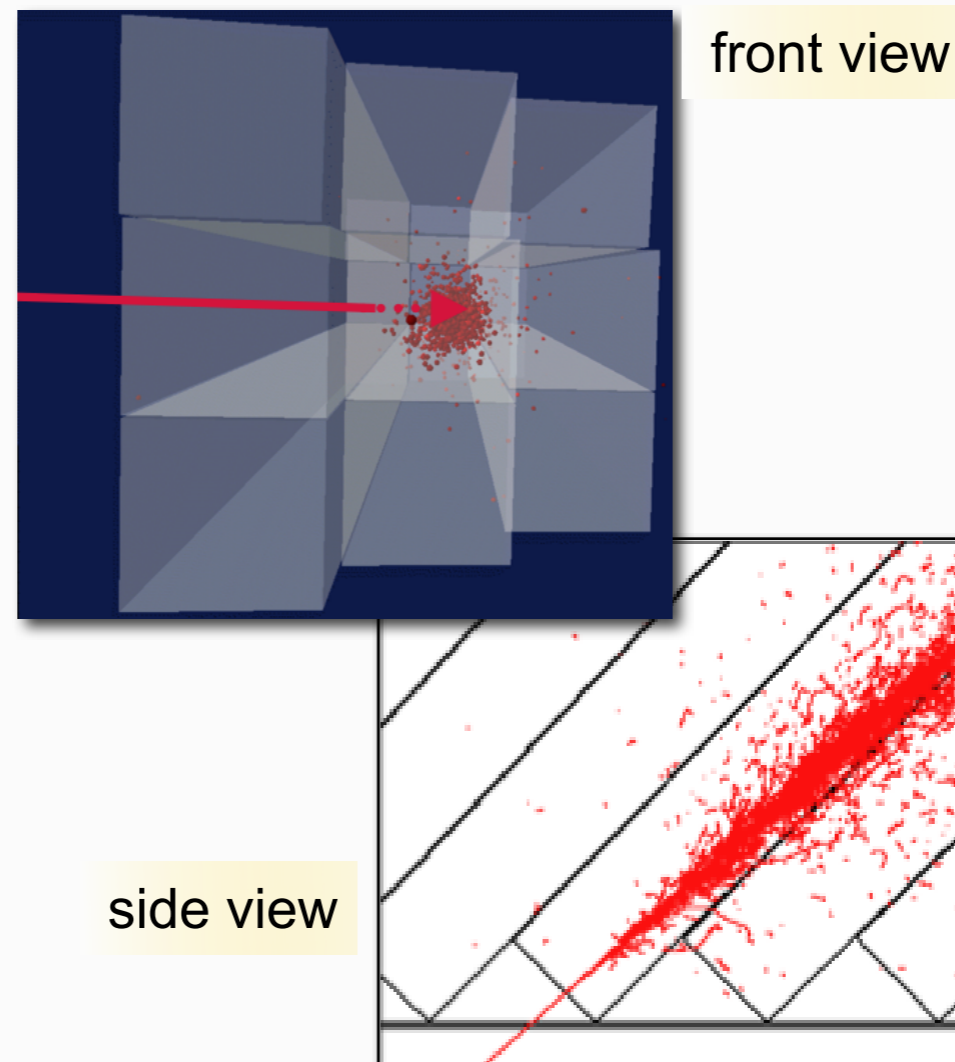
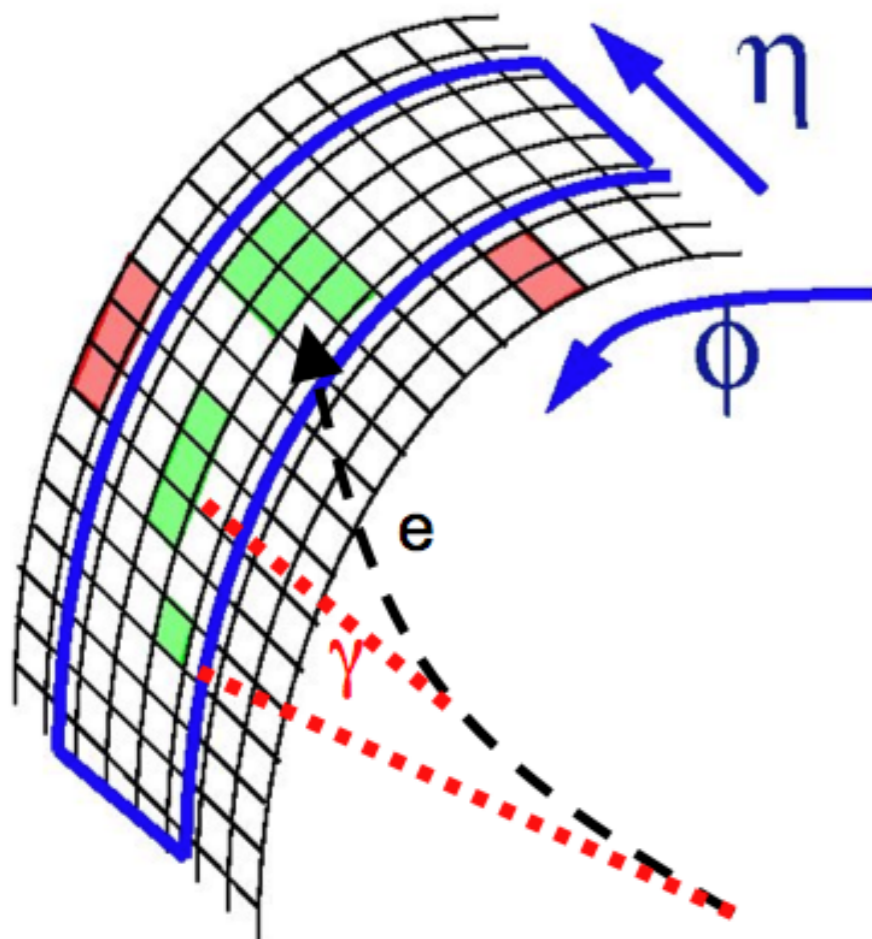


- and obtain the best possible resolution!

- Calorimeters are segmented in **cells**
- Typically a shower extends over several cells
 - Useful to reconstruct precisely the impact point from the “center-of-gravity” of the deposits in the various cells

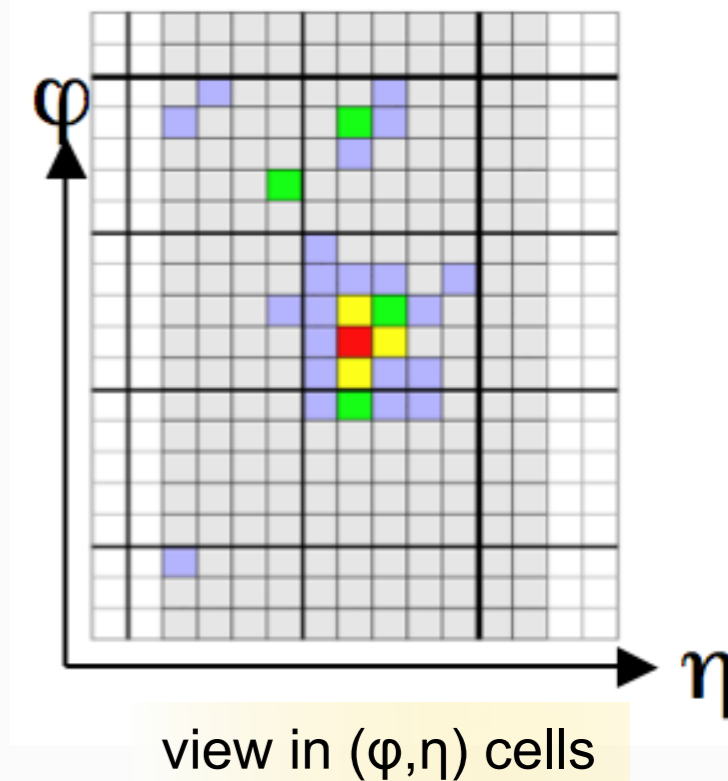
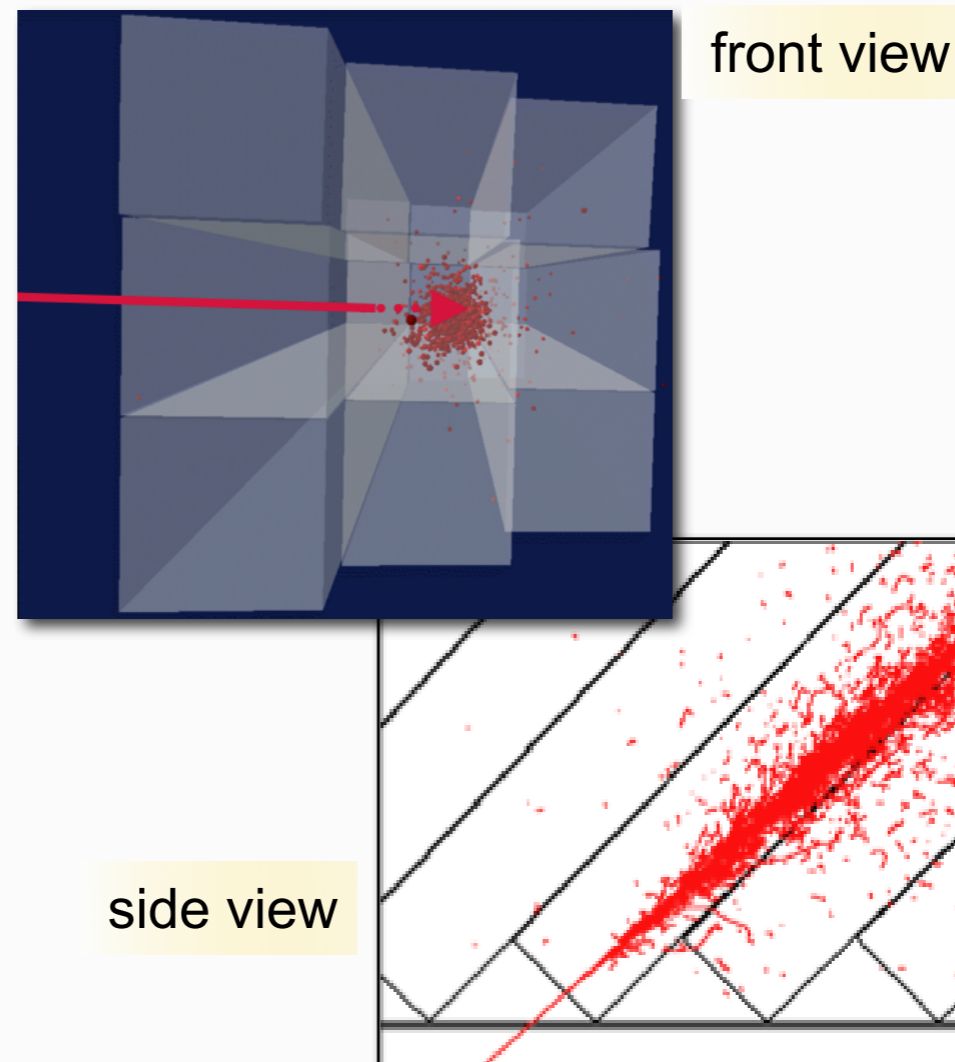
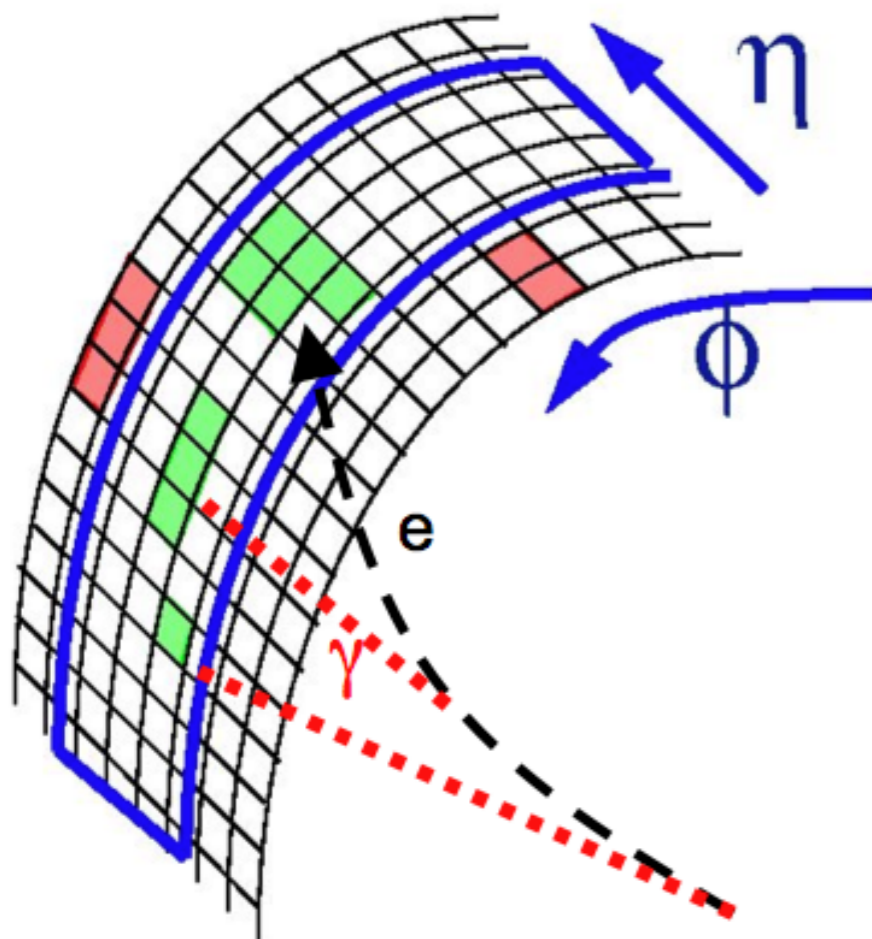
Clusters of energy

- Calorimeters are segmented in **cells**
- Typically a shower extends over several cells
 - Useful to reconstruct precisely the impact point from the “center-of-gravity” of the deposits in the various cells
- Example CMS Crystal Calorimeter:**
 - electron energy in central crystal $\sim 80\%$, in 5×5 matrix around it $\sim 96\%$

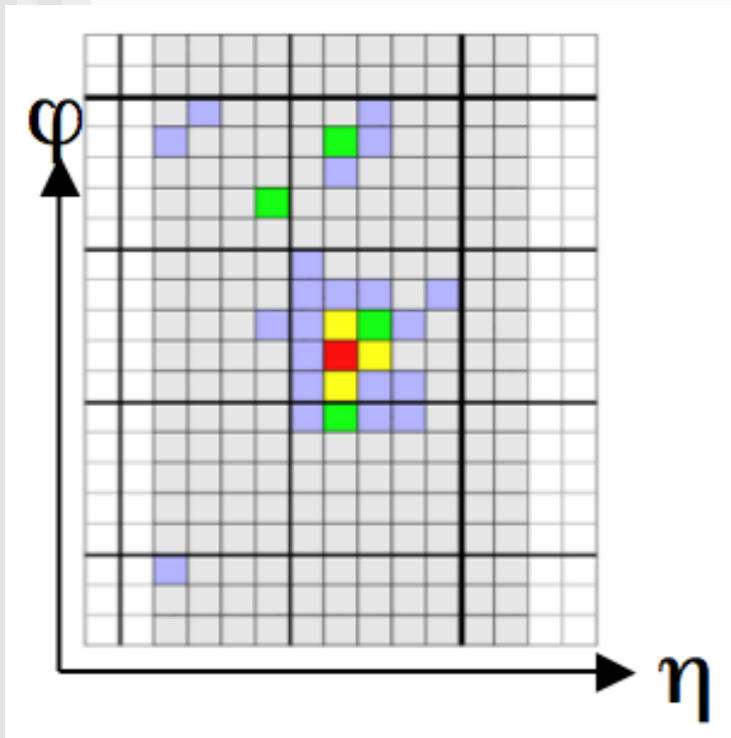


Clusters of energy

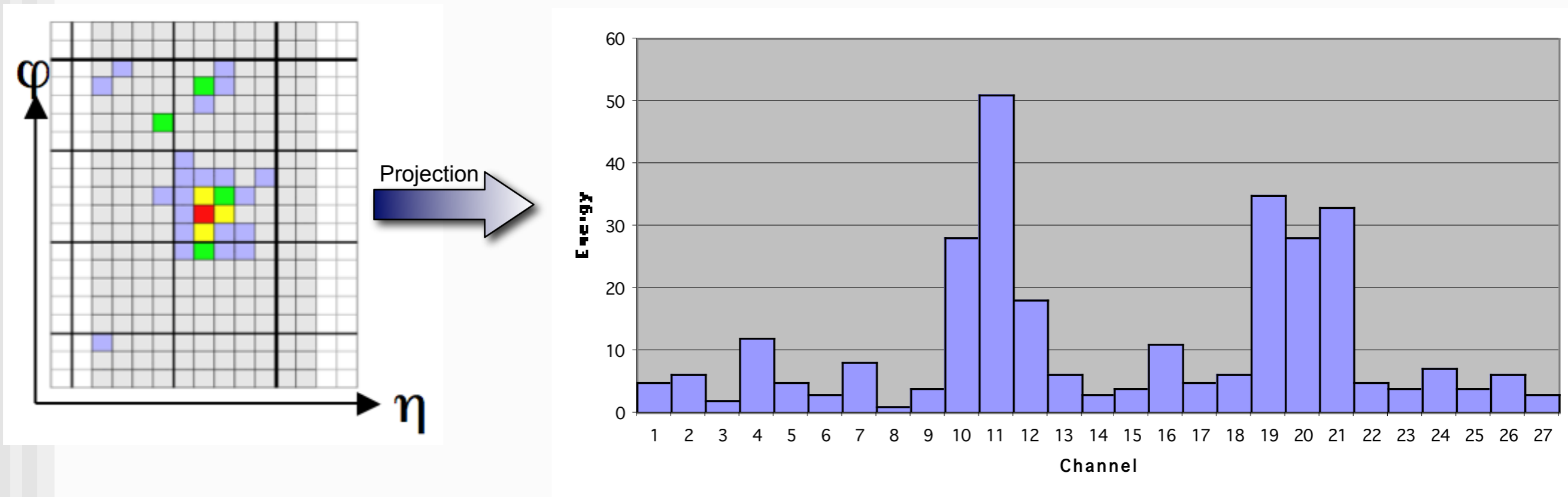
- Calorimeters are segmented in **cells**
- Typically a shower extends over several cells
 - Useful to reconstruct precisely the impact point from the “center-of-gravity” of the deposits in the various cells
- Example CMS Crystal Calorimeter:**
 - electron energy in central crystal $\sim 80\%$, in 5×5 matrix around it $\sim 96\%$
- So **task** is : identify these clusters and reconstruct the energy they contain



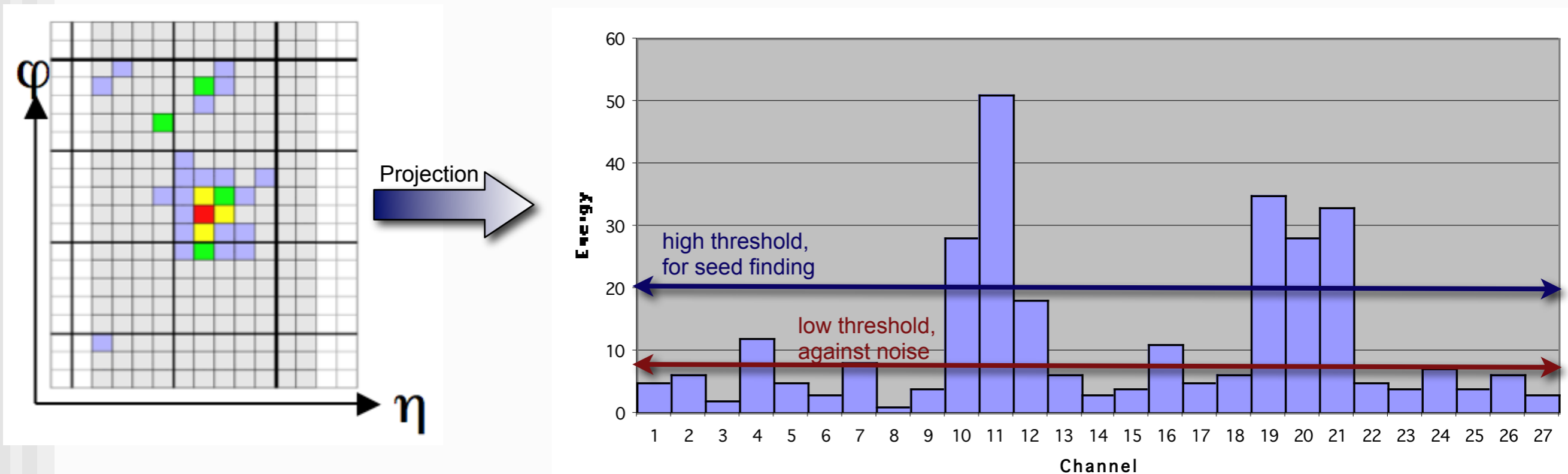
- Clusters of energy in a calorimeter are due to the original particles
 - Clustering algorithm groups individual channel energies
 - Don't want to miss any; don't want to pick up fakes



- Clusters of energy in a calorimeter are due to the original particles
 - Clustering algorithm groups individual channel energies
 - Don't want to miss any; don't want to pick up fakes



- Clusters of energy in a calorimeter are due to the original particles
 - Clustering algorithm groups individual channel energies
 - Don't want to miss any; don't want to pick up fakes

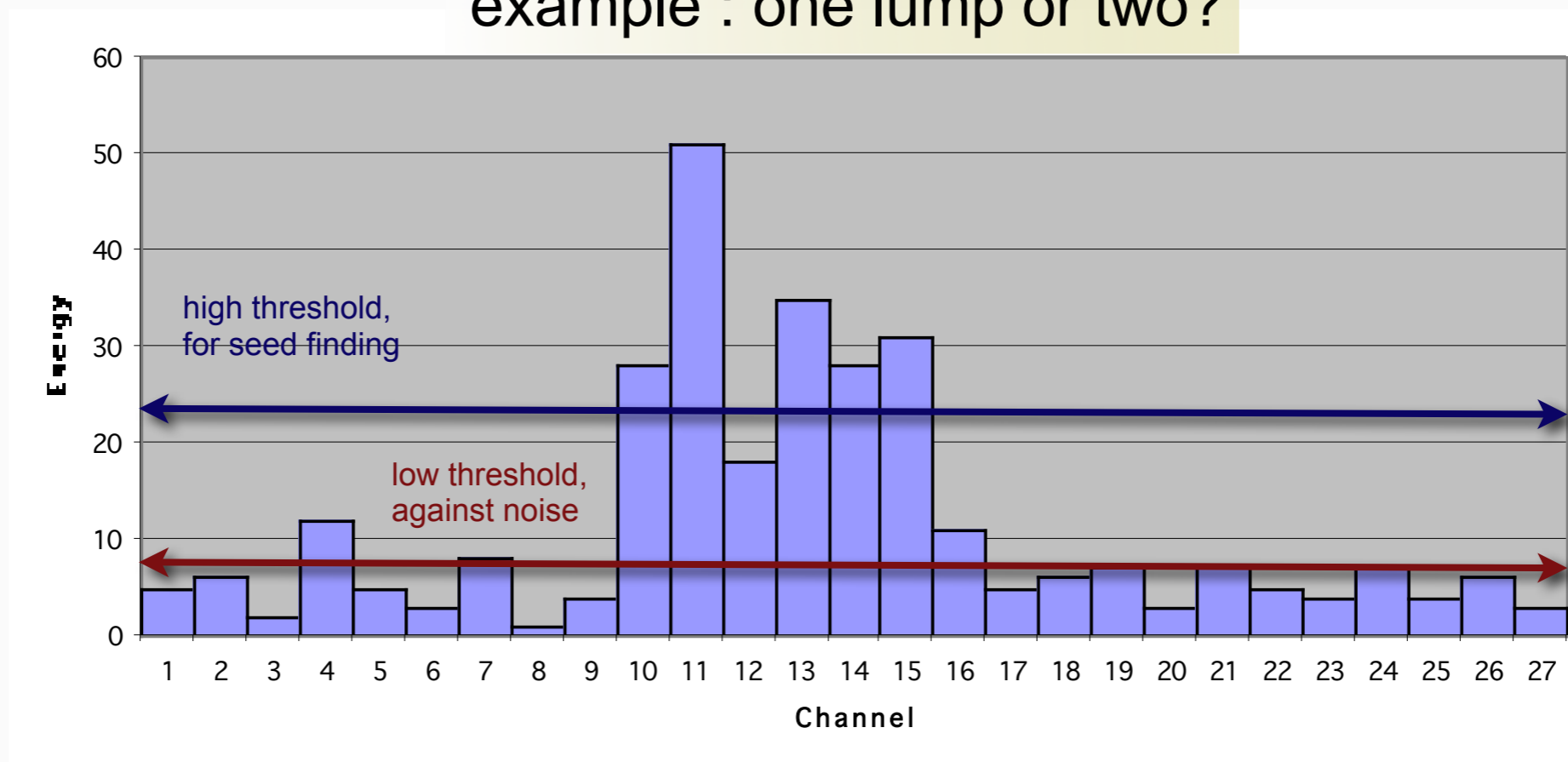


Simple example of an algorithm

- Scan for **seed** crystals = local energy maximum above a defined **seed threshold**
- Starting from the seed position, adjacent crystals are examined, scanning first in ϕ and then in η
- Along each scan line, crystals are **added to the cluster if**
 - The crystal's energy is above the **noise level (lower threshold)**
 - The crystal has not been assigned to another cluster already
 - The previous crystal added (in the same direction) has higher energy

- **Careful tuning of thresholds needed**
 - needs usually learning phase
 - adapt to noise conditions
 - **too low** : pick up too much unwanted energy
 - **too high** : loose too much of “real” energy. Corrections/Calibrations will be larger
- **Sometimes several clustering stages, in order separate or combine nearby clusters**

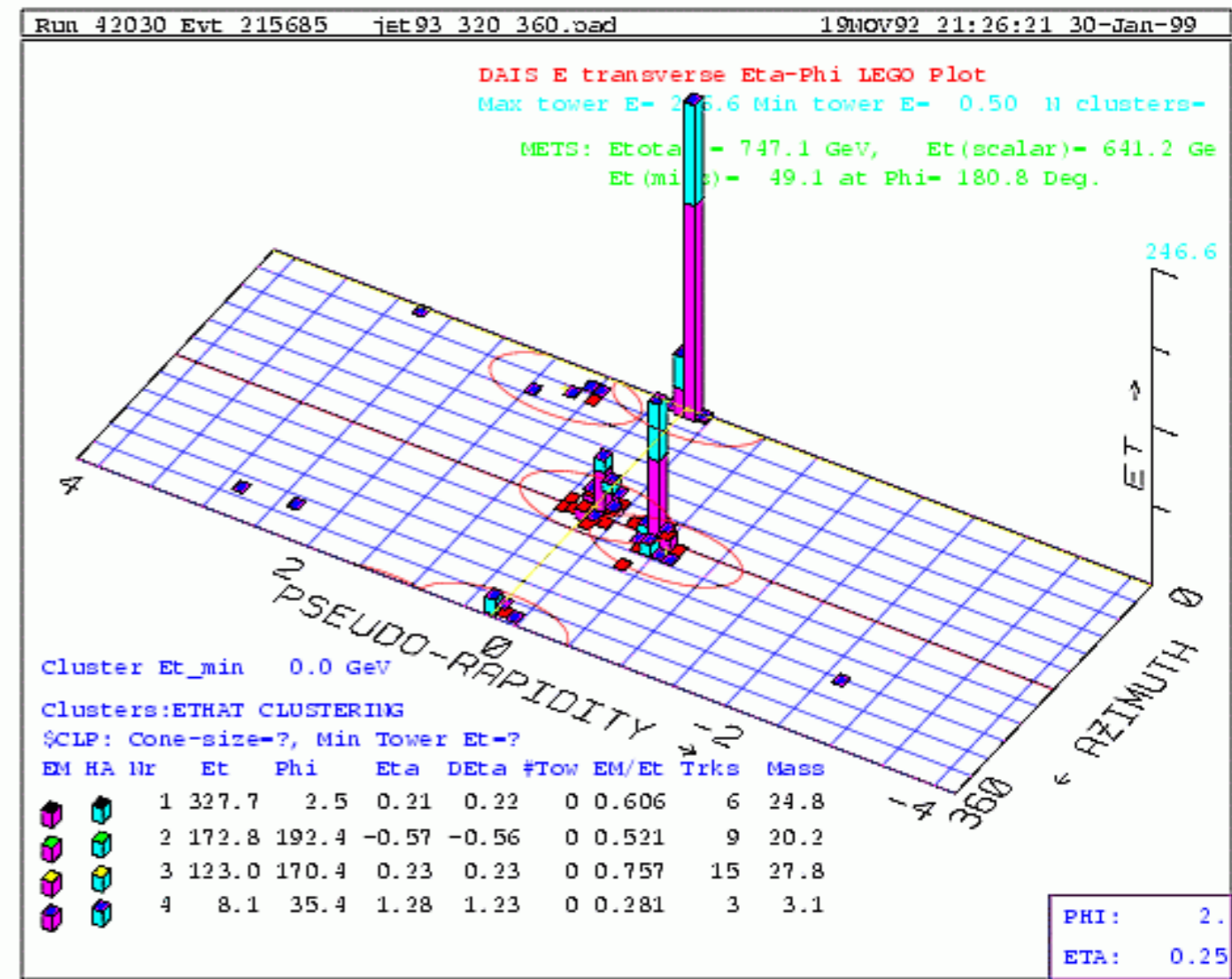
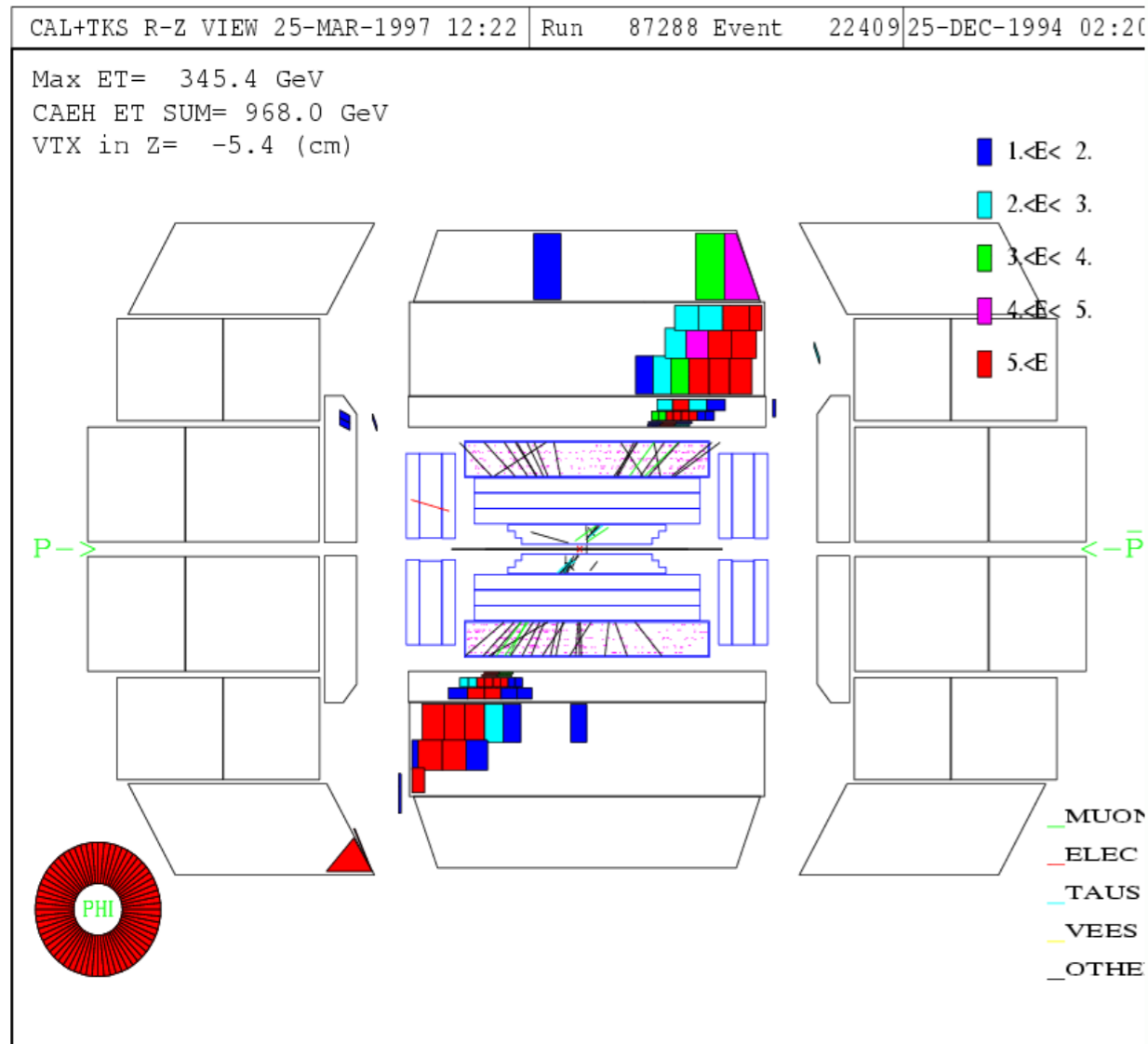
example : one lump or two?



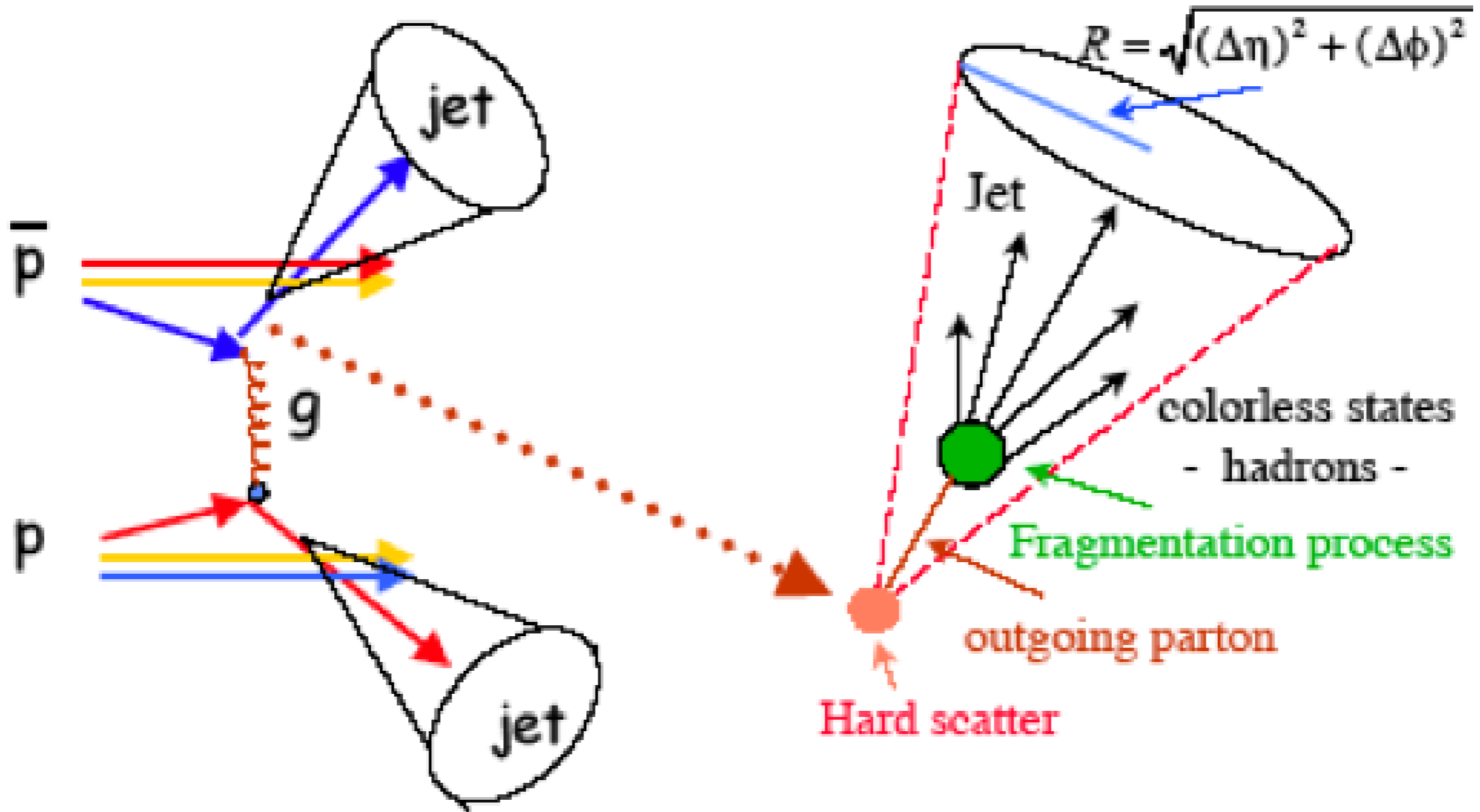
Jet Algorithms

Jets in DØ

CDF

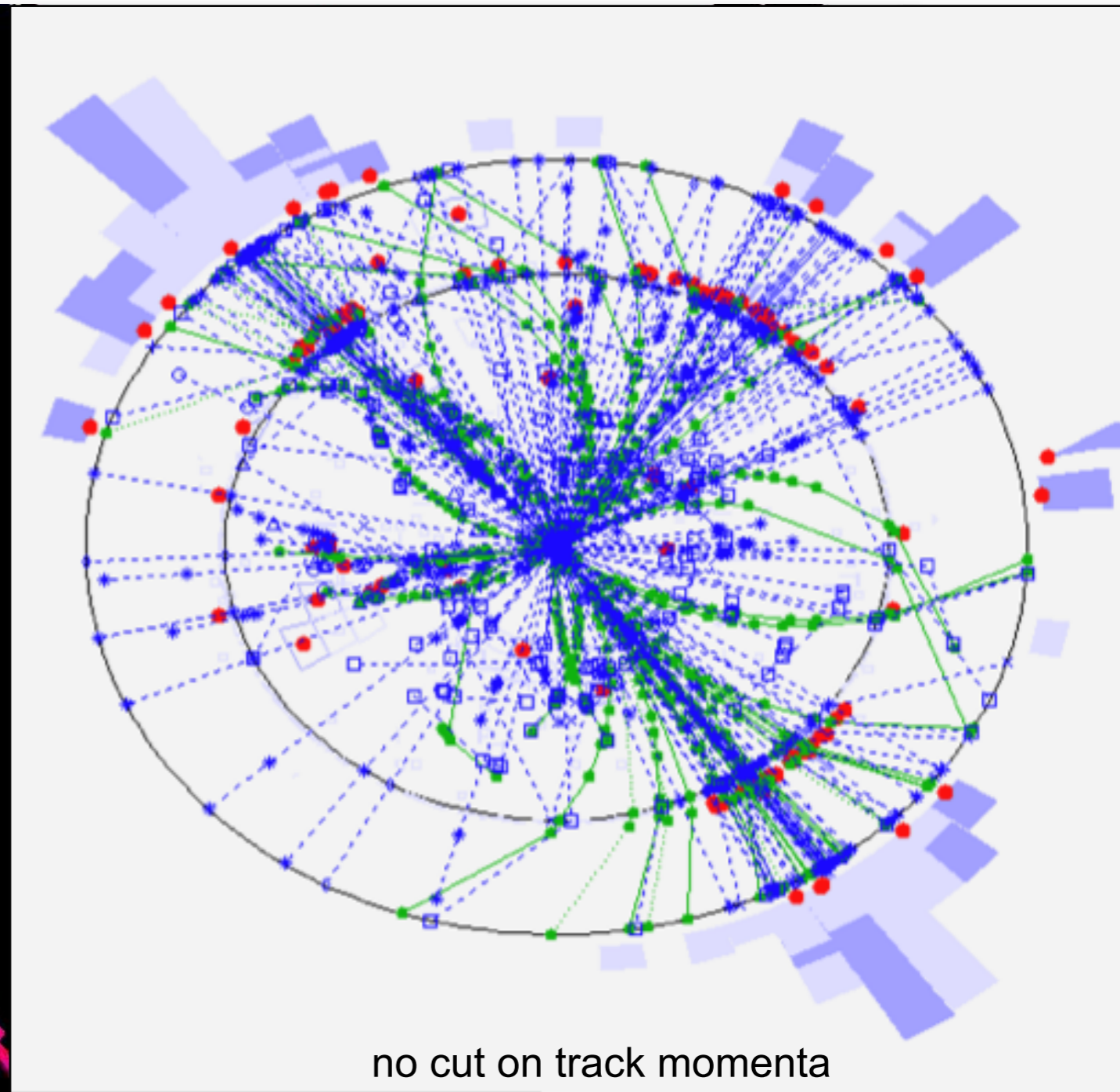
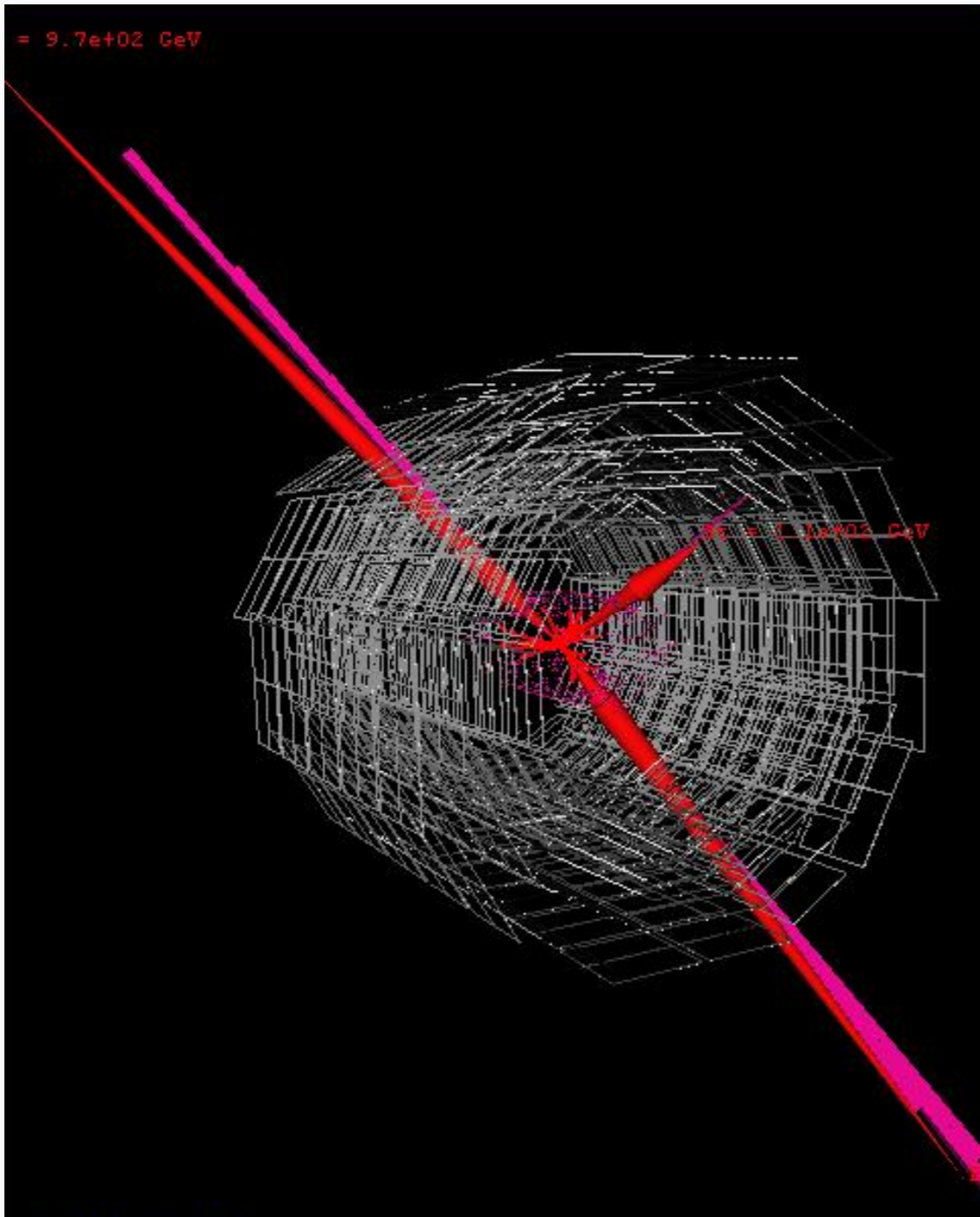


- Introducing a cone prescription seems “natural”...
- But how to make it more quantitative?
 - don't want people “guessing” at whether there are 2,3, ... jets

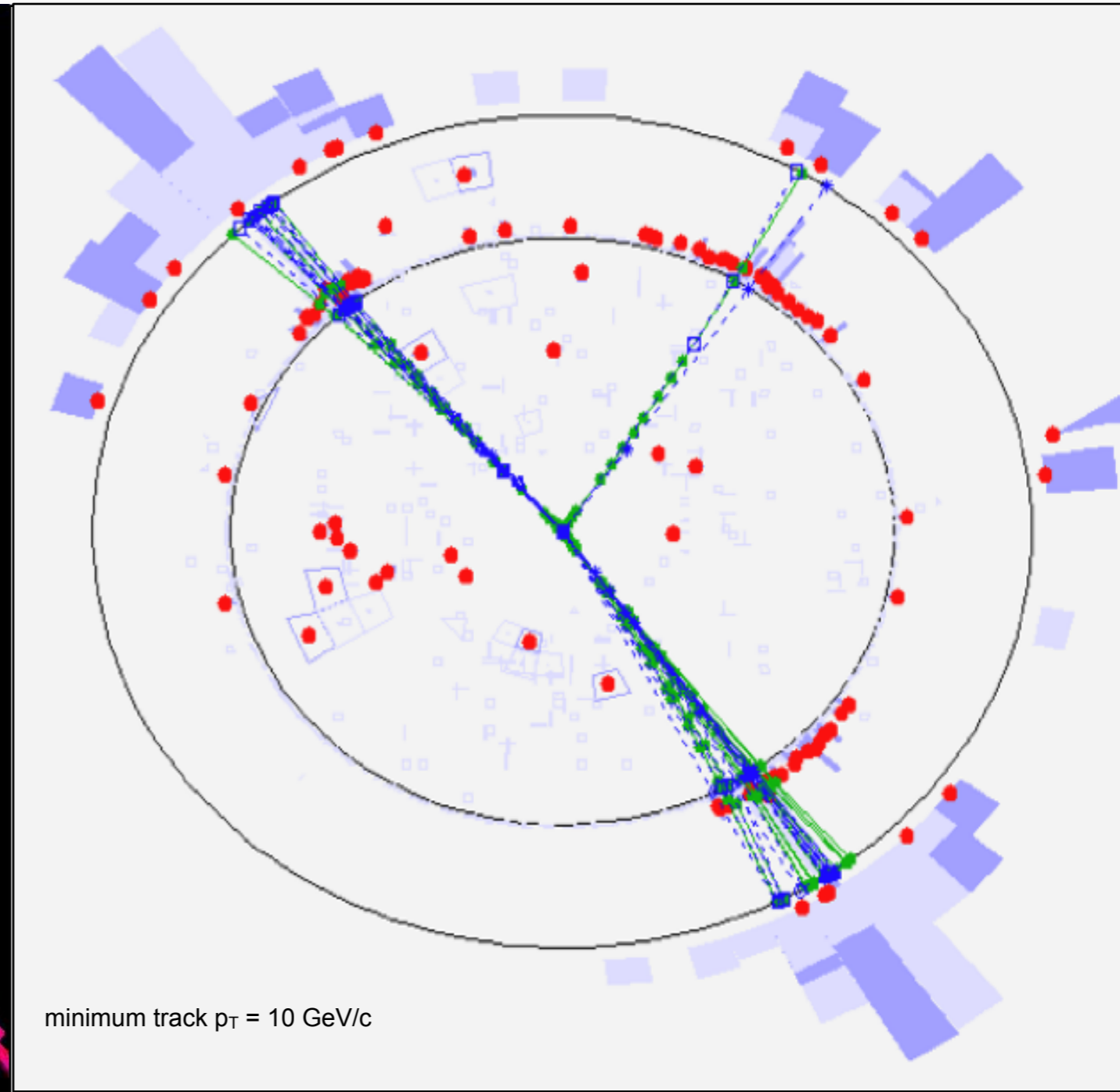
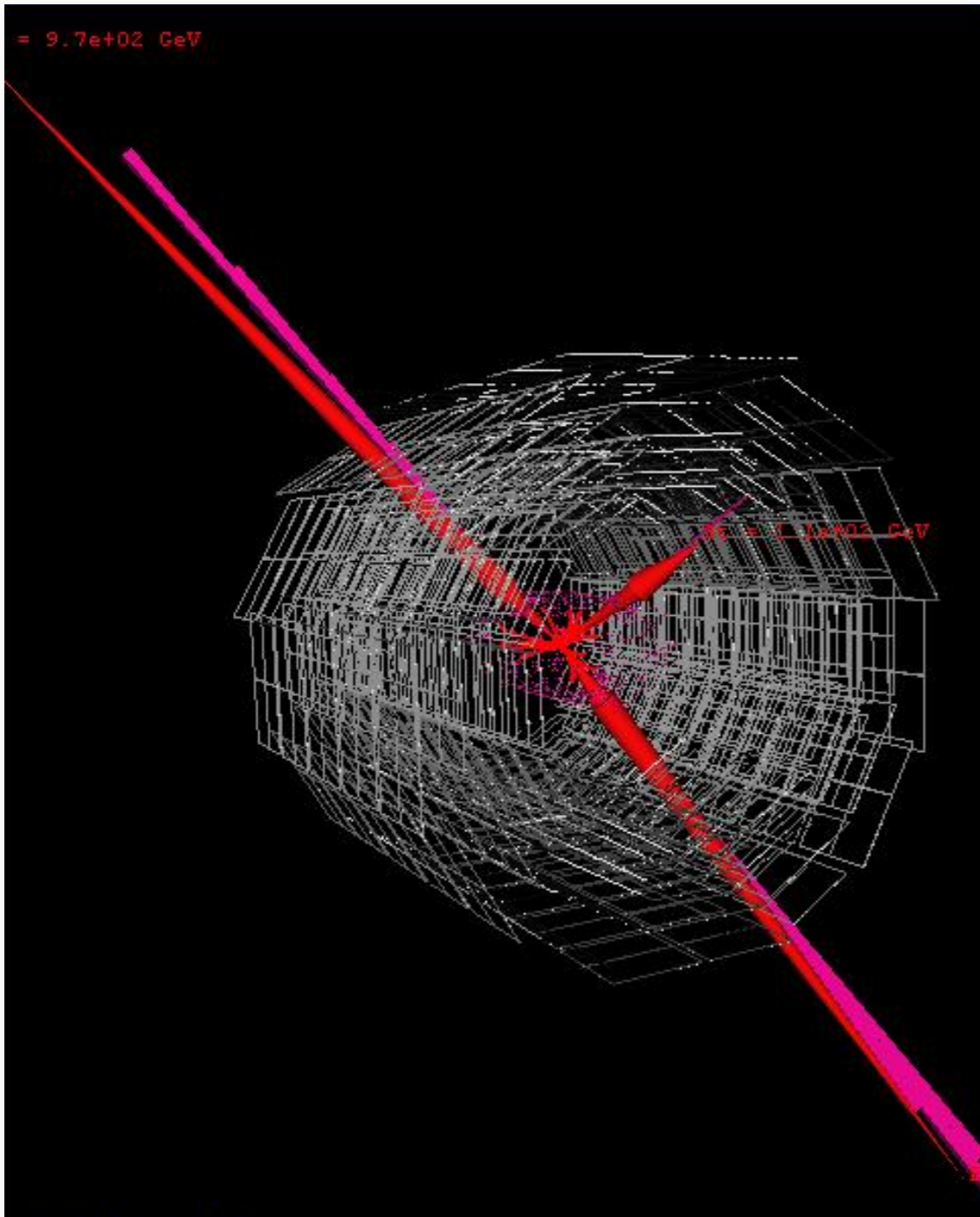


The natural (?) definition of a jet in a hadron collider environment

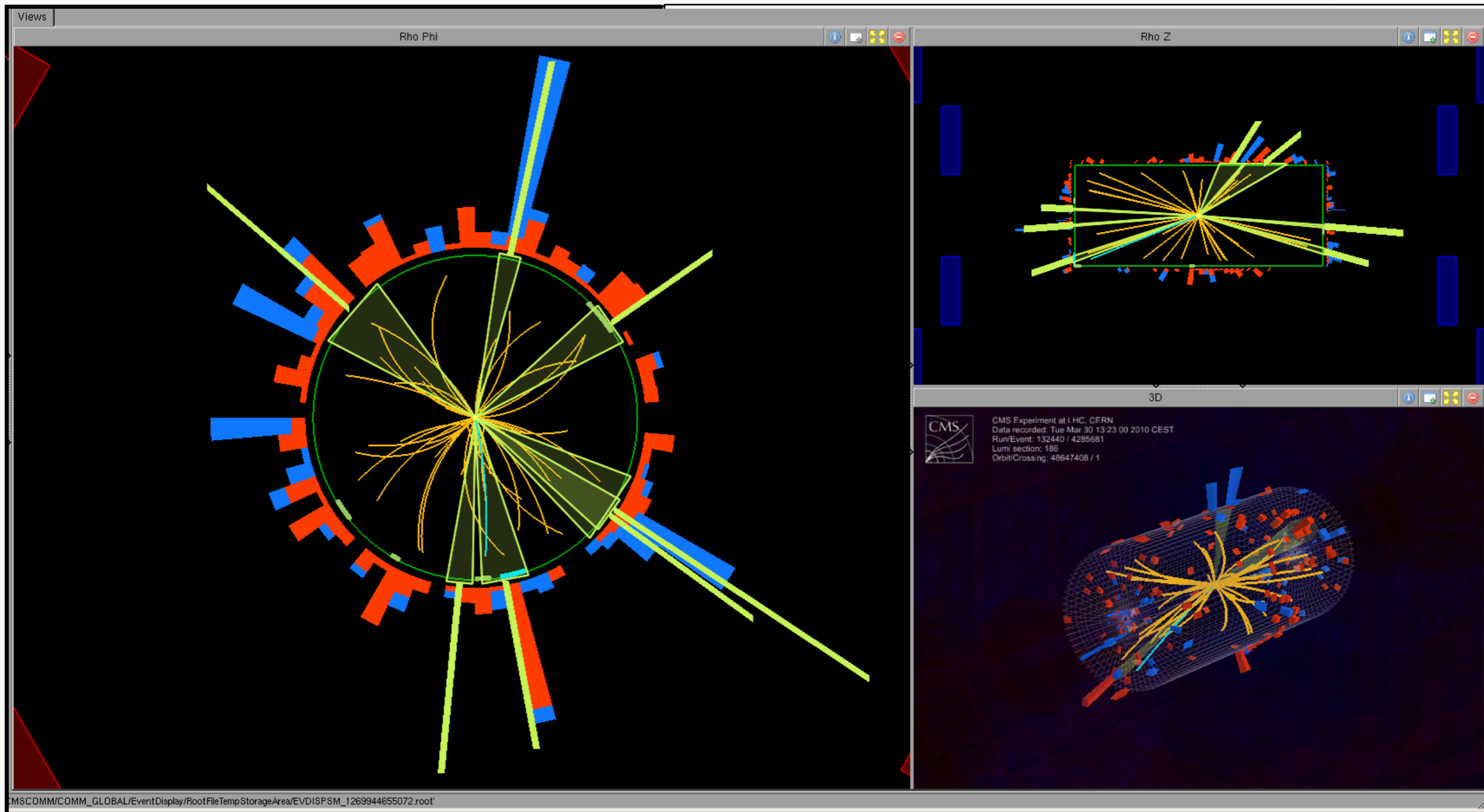
Jets in CMS



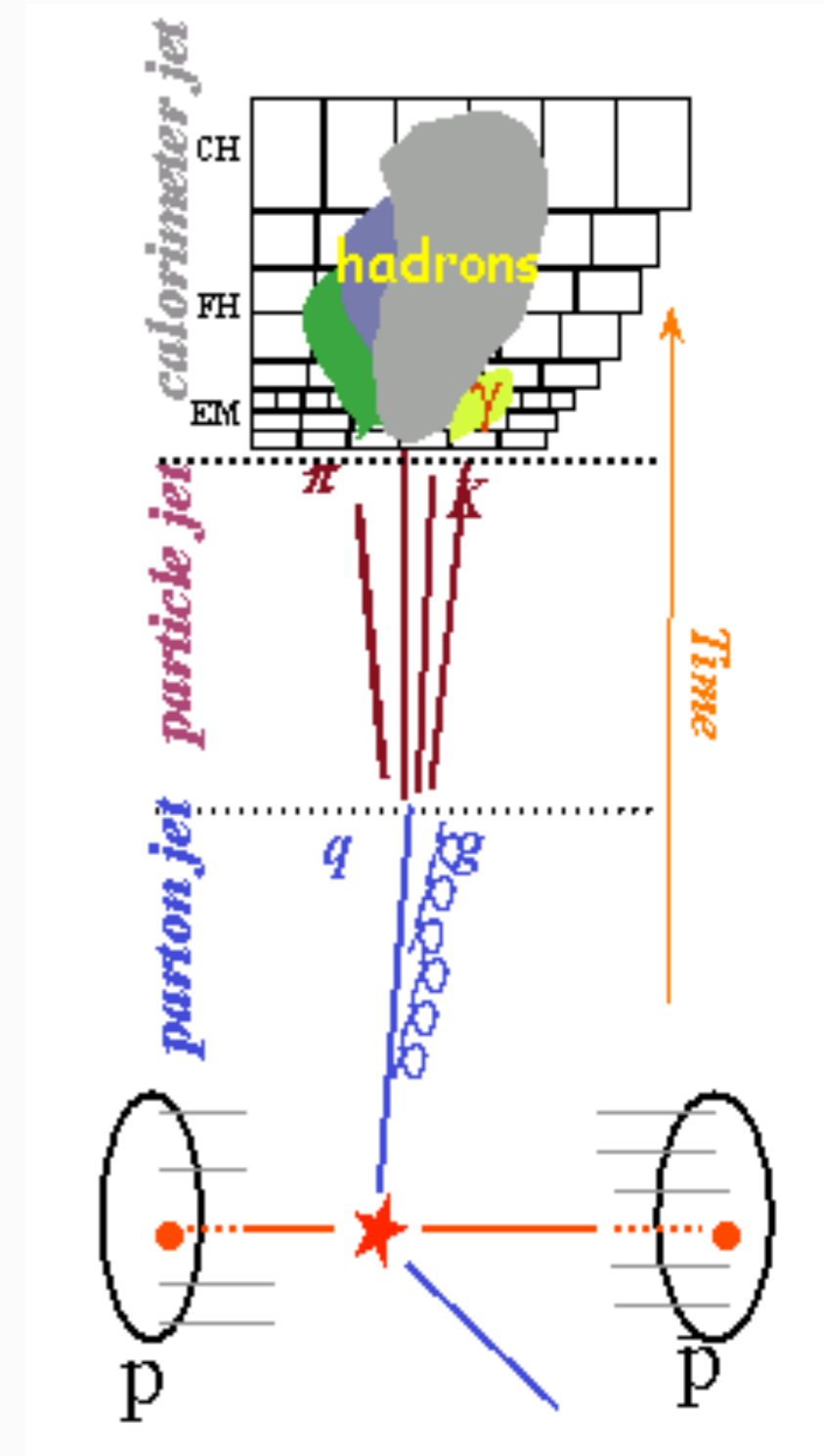
Jets in CMS



Jets in CMS

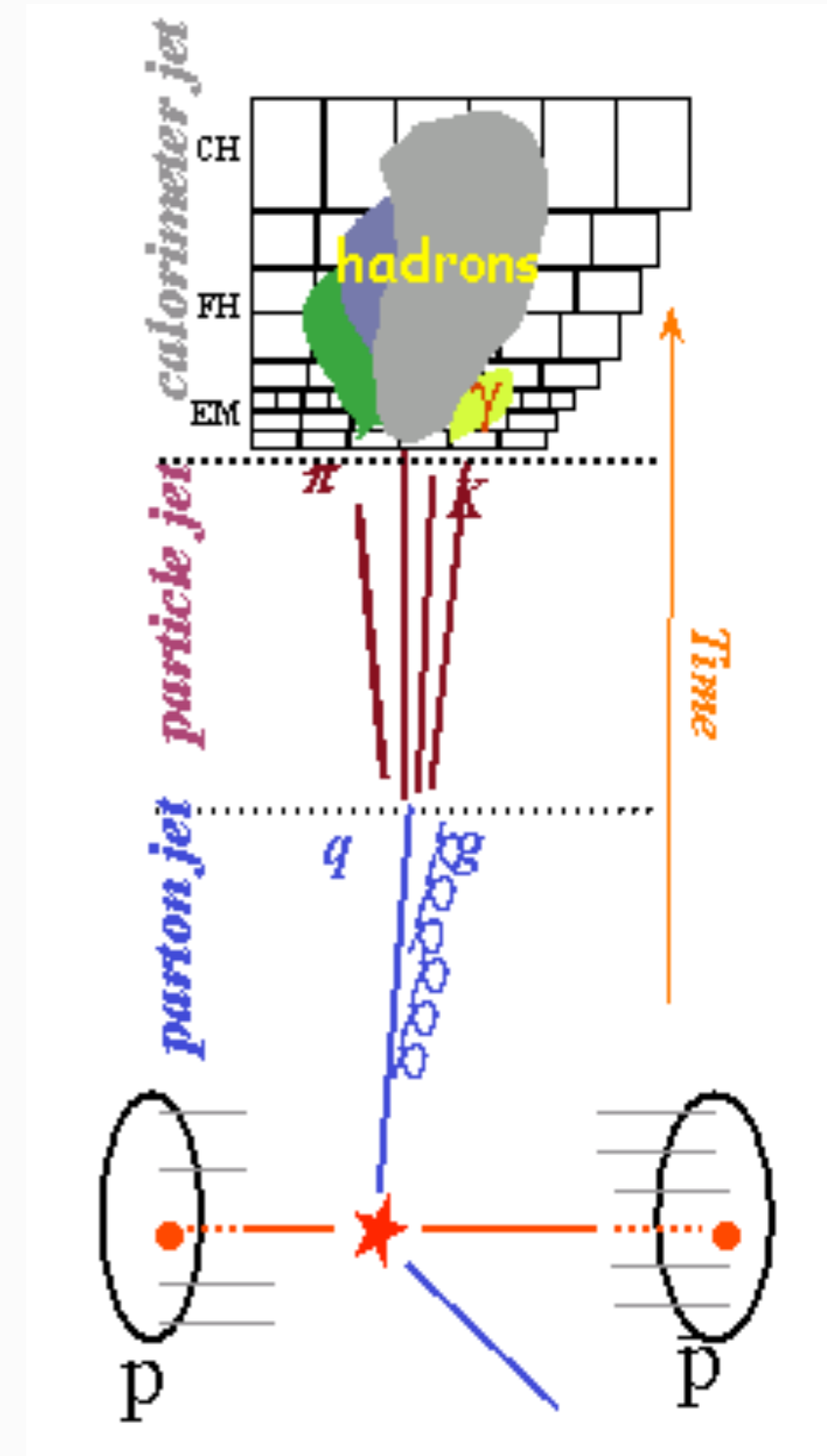


- Applicable at all levels
 - partons, stable particles
 - ◆ for theoretical calculations
 - measured objects (calorimeter objects, tracks, etc)
 - and always **find the same jet**



- **Applicable at all levels**
 - partons, stable particles
 - ◆ for theoretical calculations
 - measured objects (calorimeter objects, tracks, etc)
 - and always **find the same jet**

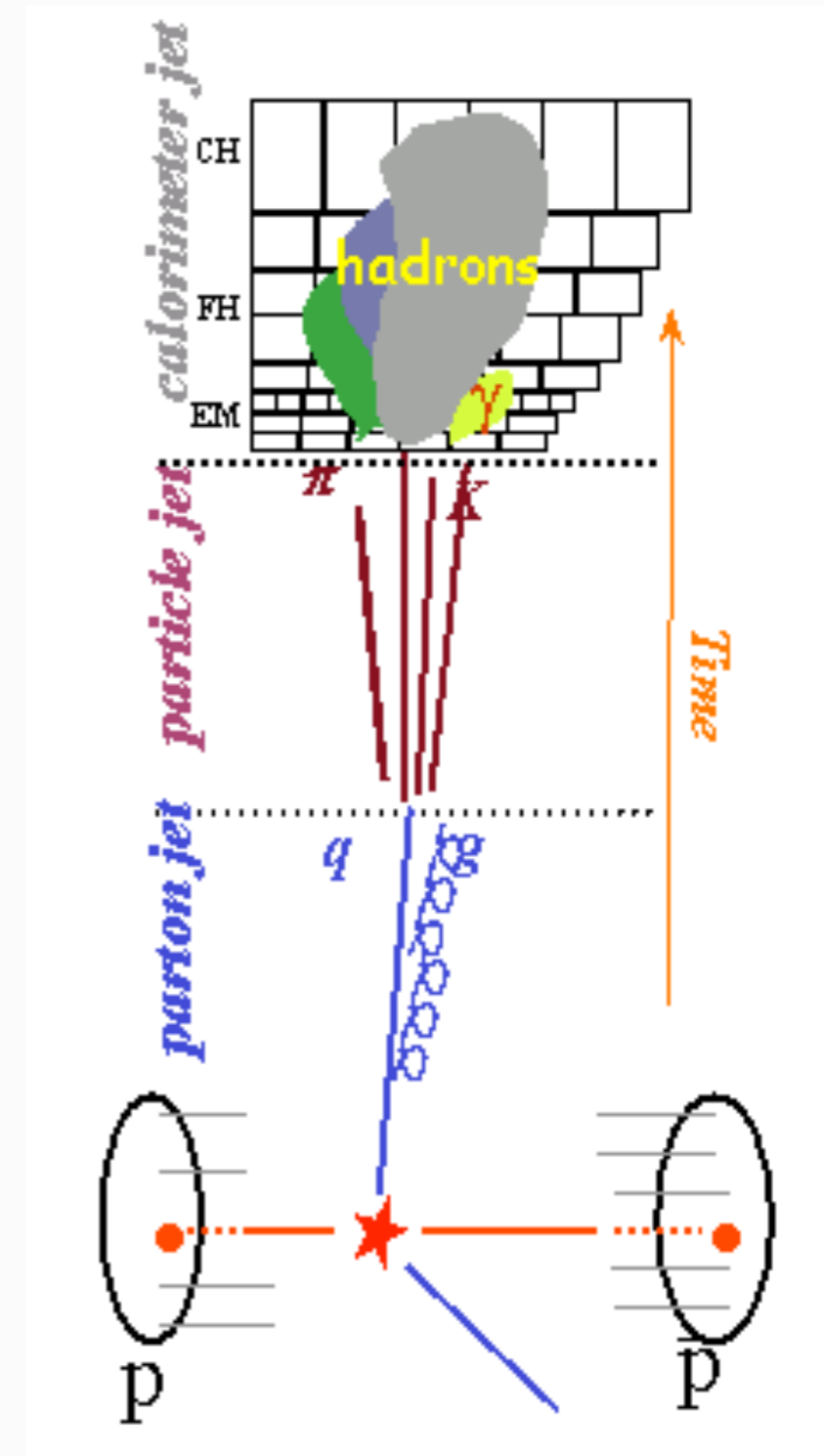
- **Independent of the very details of the detector**
 - example : granularity of the calorimeter, energy response, ...



- Applicable at all levels**
 - partons, stable particles
 - for theoretical calculations**
 - measured objects (calorimeter objects, tracks, etc)
 - and always **find the same jet**

- Independent of the very details of the detector**
 - example : granularity of the calorimeter, energy response, ...

- Easy to implement !**



- Applicable at all levels
 - partons, stable particles
 - ◆ for theoretical calculations
 - measured objects (calorimeter objects, tracks, etc)
 - and always **find the same jet**

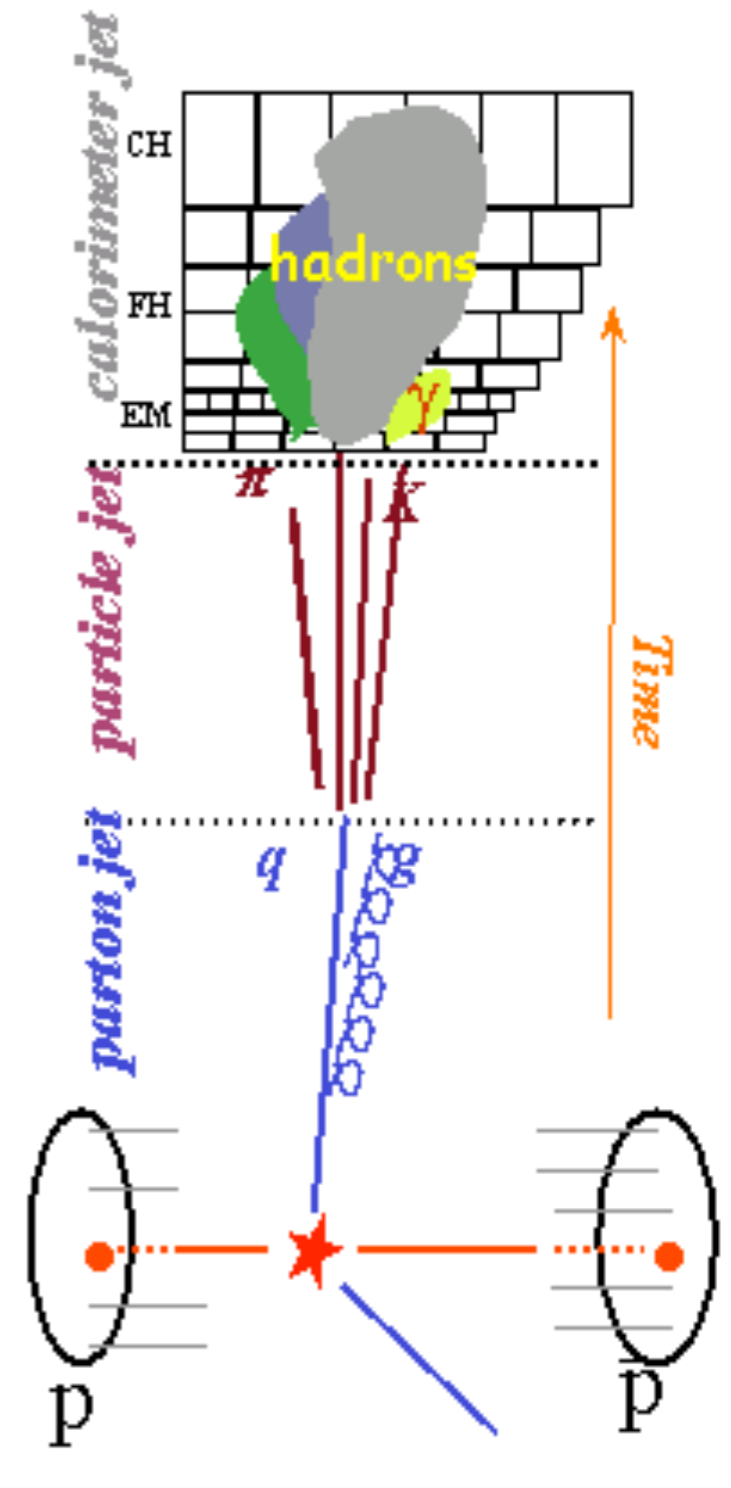
- Independent of the very details of the detector
 - example : granularity of the calorimeter, energy response,...

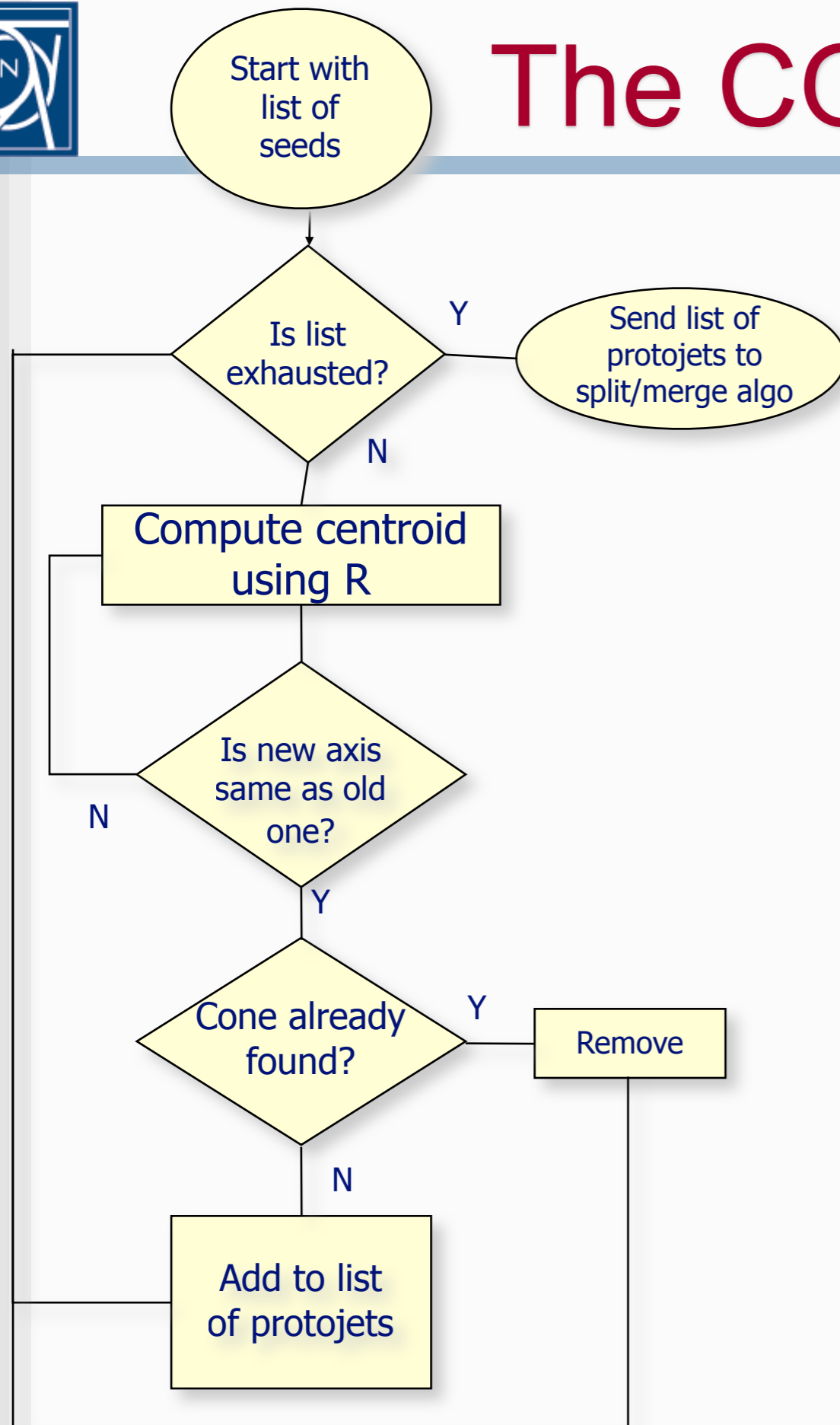
- Easy to implement !

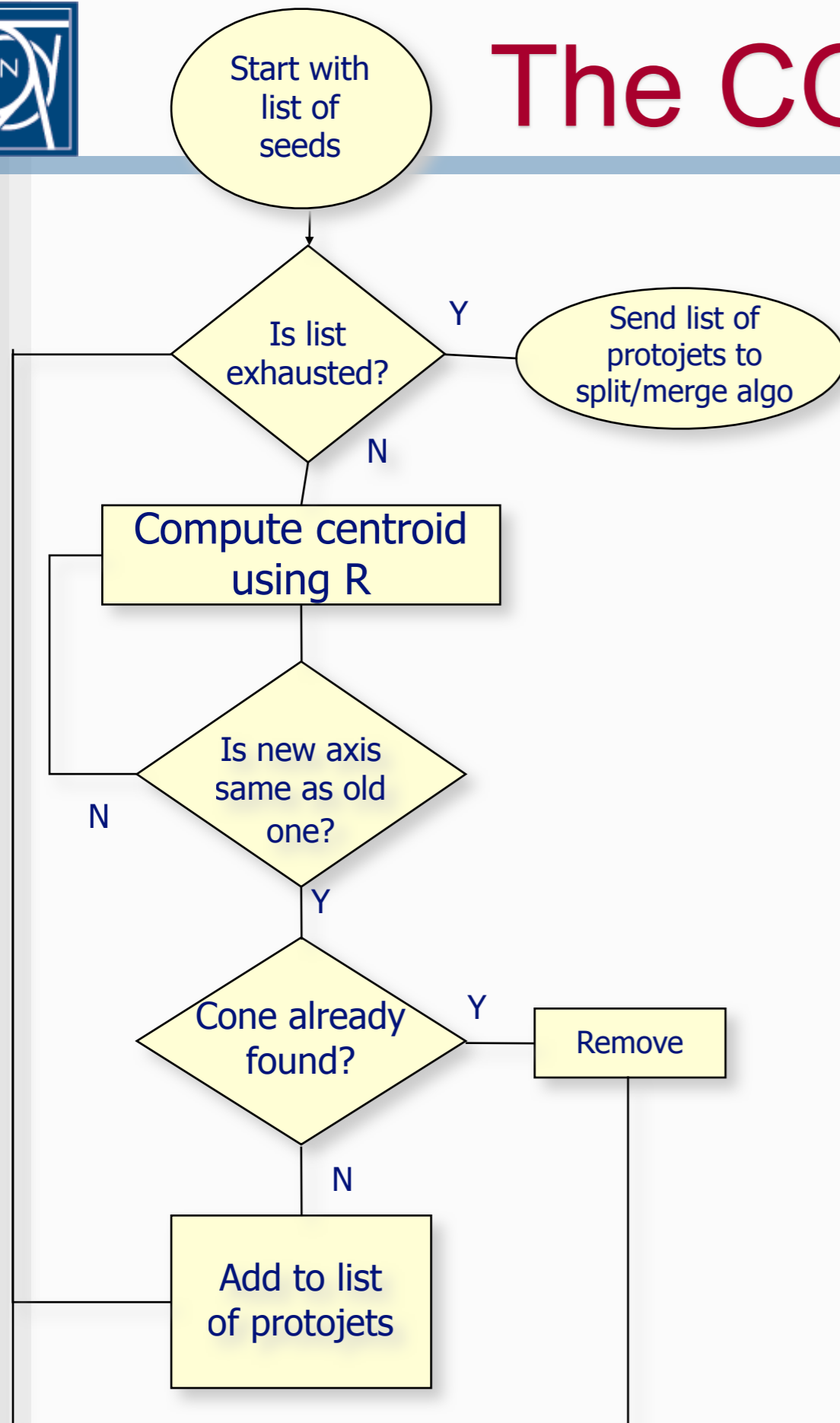
- Close correspondence between

$$P_{\text{parton}} \longleftrightarrow P_{\text{jet}}$$

Energy
Momentum
angle

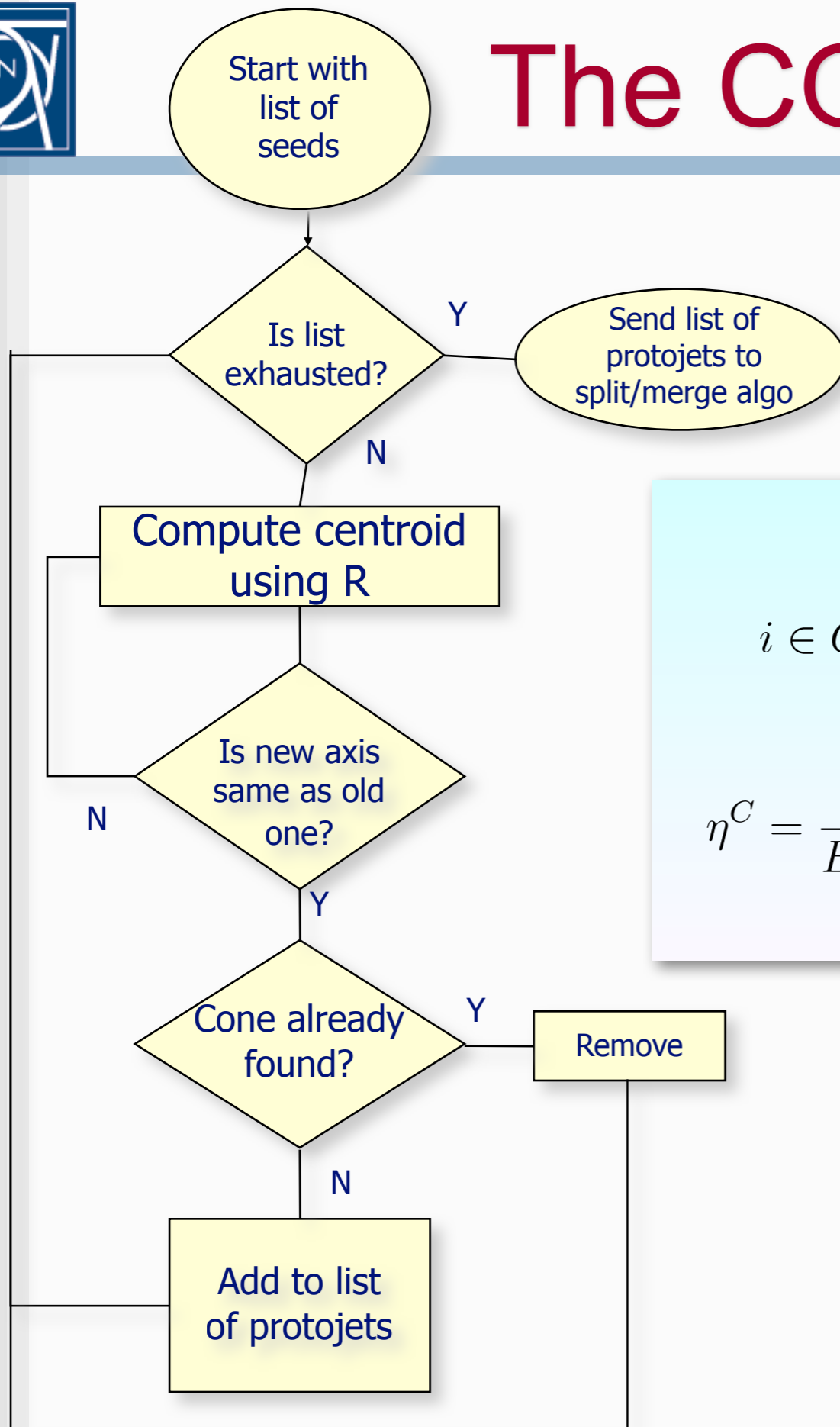






Seeds: for example, energy deposits with transverse energy ($E_T = E \sin\theta$) > 2 GeV in a tower of the calorimeter

Seeds: for example, energy deposits with transverse energy ($E_T = E \sin\theta$) > 2 GeV in a tower of the calorimeter



Centroid (one possible def) :

$$i \in C : \sqrt{(\eta^i - \eta^C)^2 + (\Phi^i - \Phi^C)^2} \leq R \quad \text{cone radius}$$

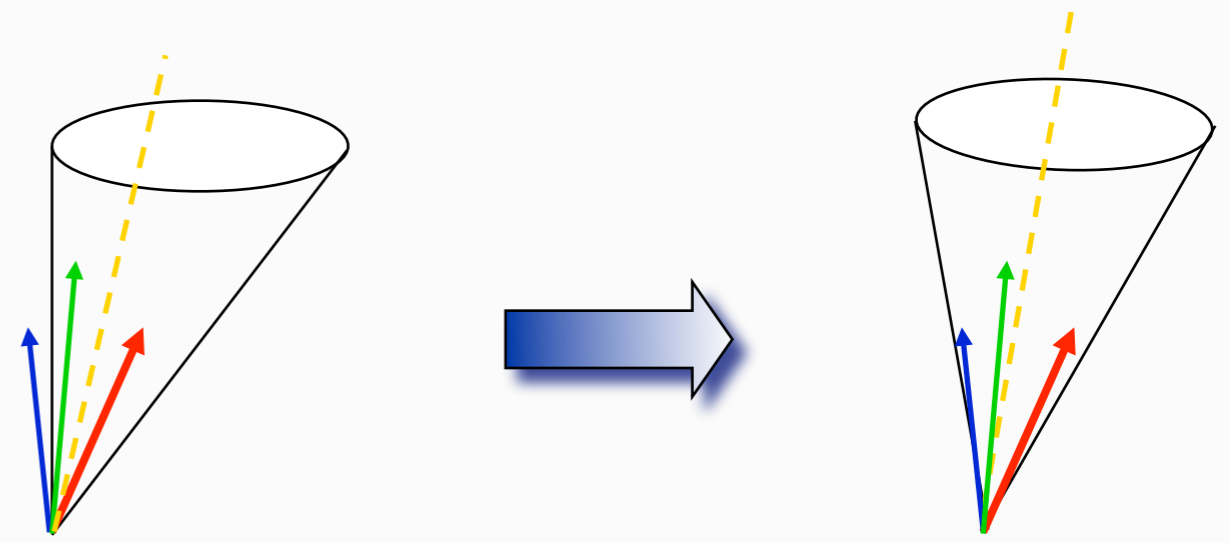
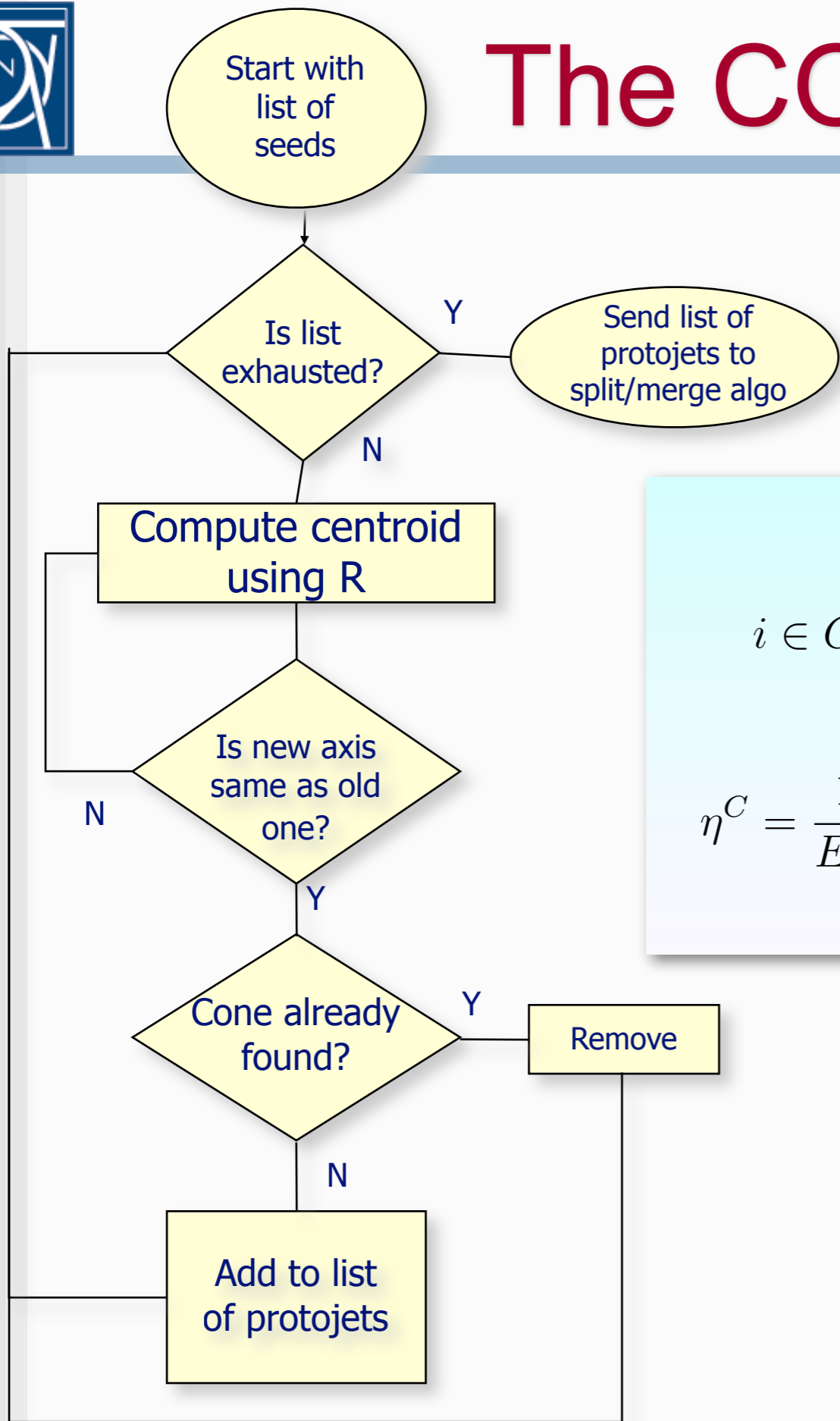
$$\eta^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \eta^i \quad ; \quad \Phi^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \Phi^i \quad ; \quad E_T^C = \sum_{i \in C} E_T^i$$

Seeds: for example, energy deposits with transverse energy ($E_T = E \sin\theta$) > 2 GeV in a tower of the calorimeter

Centroid (one possible def) :

$$i \in C : \sqrt{(\eta^i - \eta^C)^2 + (\Phi^i - \Phi^C)^2} \leq R \quad \text{cone radius}$$

$$\eta^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \eta^i \quad ; \quad \Phi^C = \frac{1}{E_T^C} \sum_{i \in C} E_T^i \Phi^i \quad ; \quad E_T^C = \sum_{i \in C} E_T^i$$



- **Pile Up** : many additional soft proton-proton interactions
 - up to 20 at highest LHC luminosity

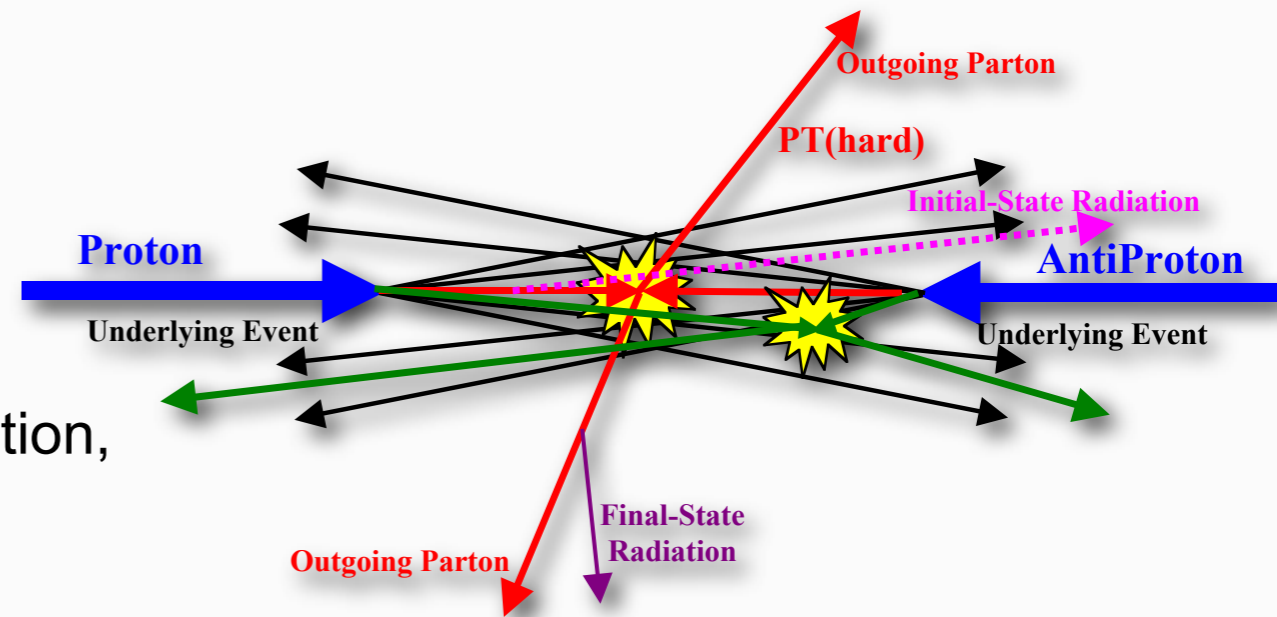
Further difficulties

Pile Up : many additional soft proton-proton interactions

- up to 20 at highest LHC luminosity

Underlying event

- beam-beam remnants, initial state radiation, multiple parton interactions
- gives additional energy in the event



Further difficulties

- **Pile Up** : many additional soft proton-proton interactions

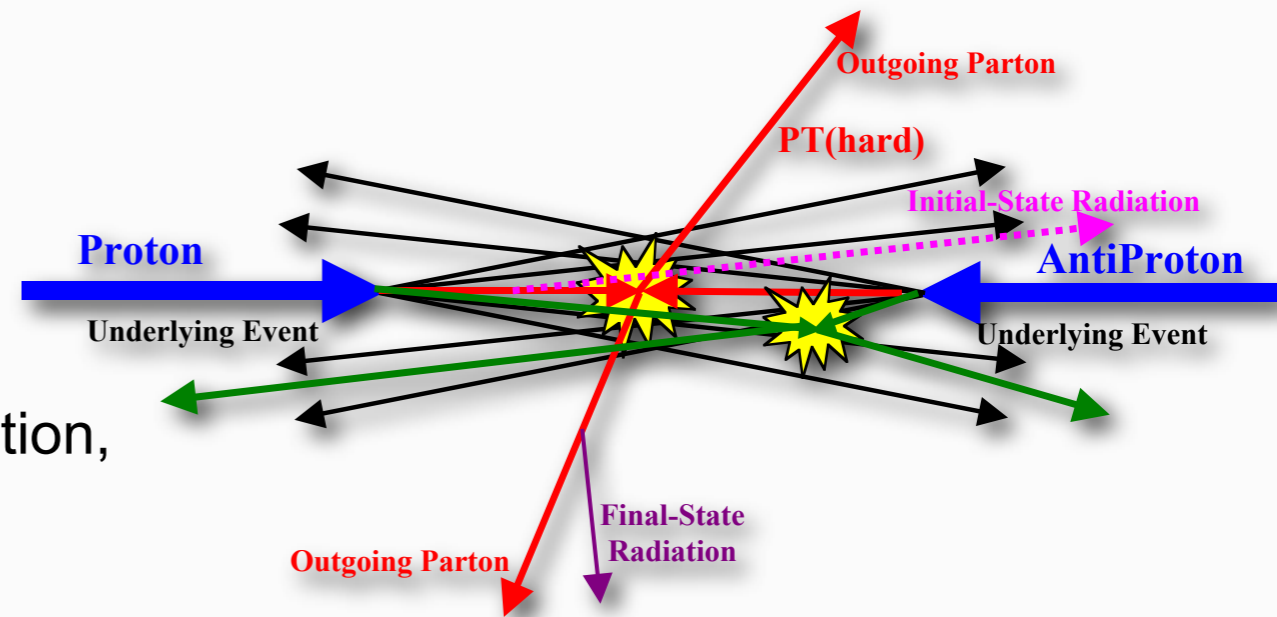
- up to 20 at highest LHC luminosity

- **Underlying event**

- beam-beam remnants, initial state radiation, multiple parton interactions
- gives additional energy in the event

- **All this additional energy has nothing to do with jet energies**

- **have to subtract it**



Further difficulties

- Pile Up** : many additional soft proton-proton interactions

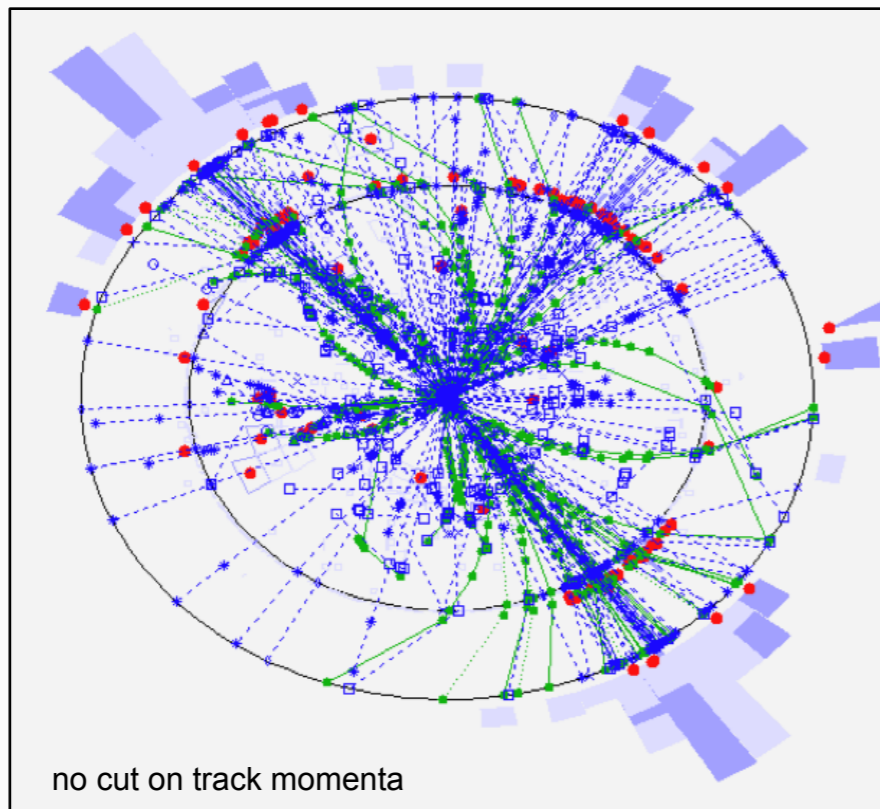
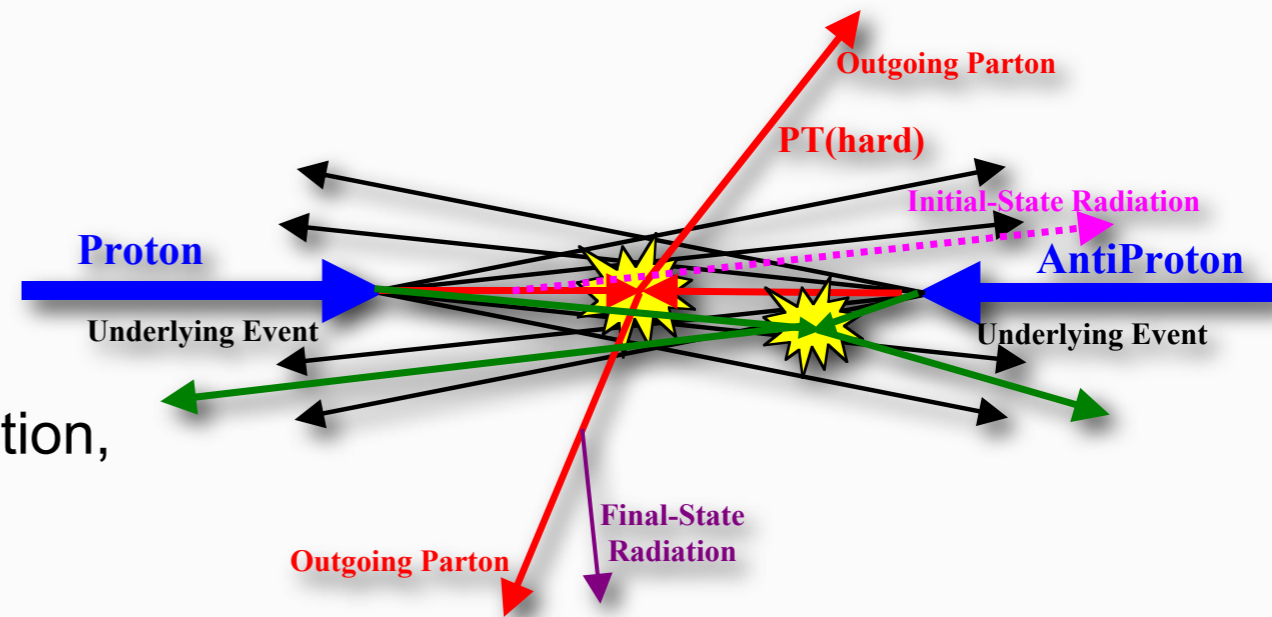
- up to 20 at highest LHC luminosity

- Underlying event**

- beam-beam remnants, initial state radiation, multiple parton interactions
 - gives additional energy in the event

- All this additional energy has nothing to do with jet energies**

- have to subtract it**



Further difficulties

- Pile Up** : many additional soft proton-proton interactions

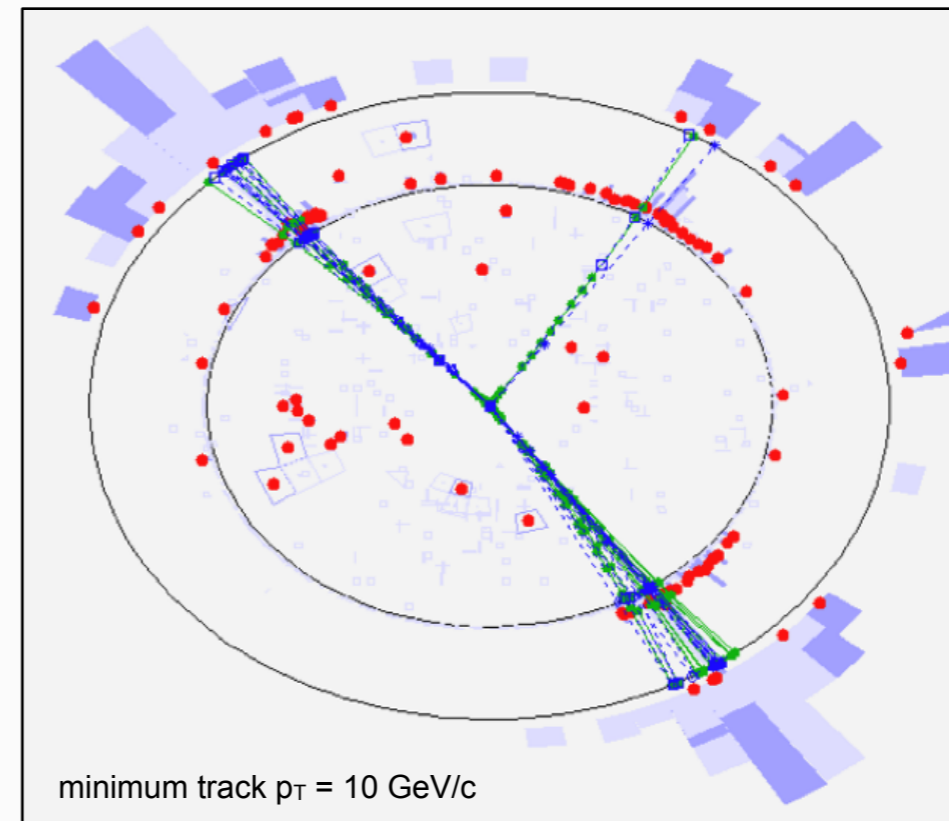
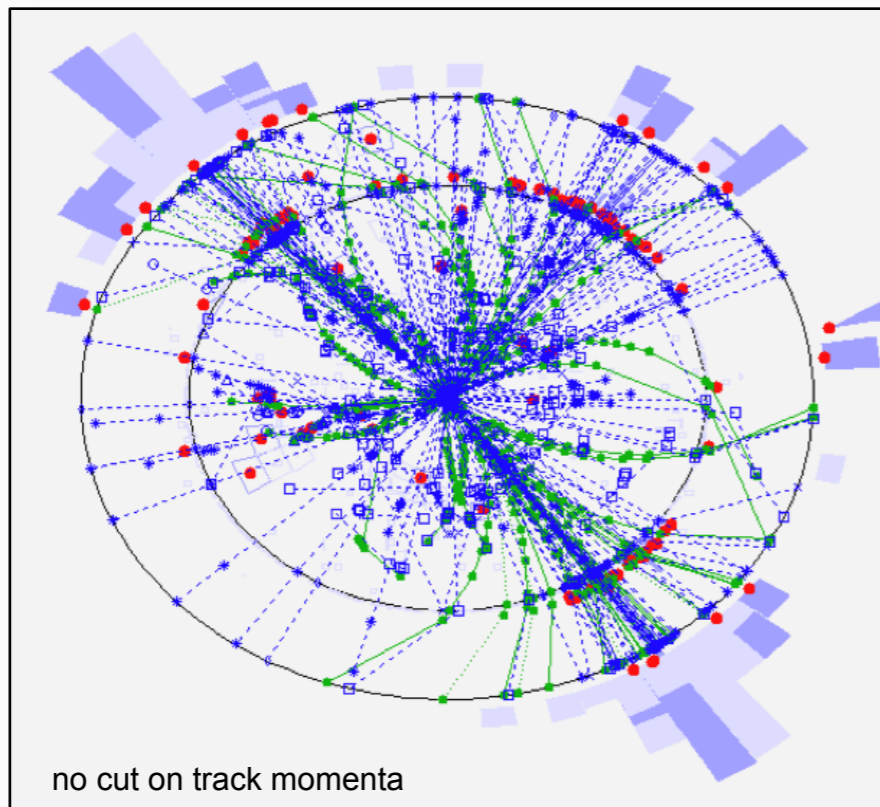
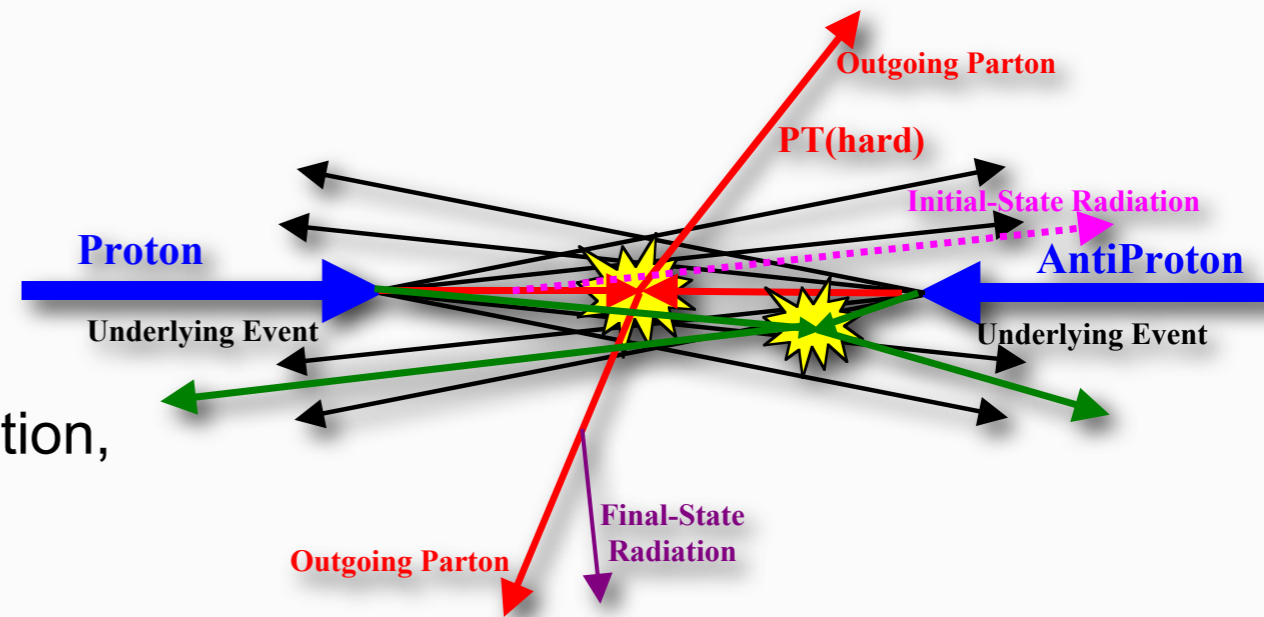
- up to 20 at highest LHC luminosity

- Underlying event**

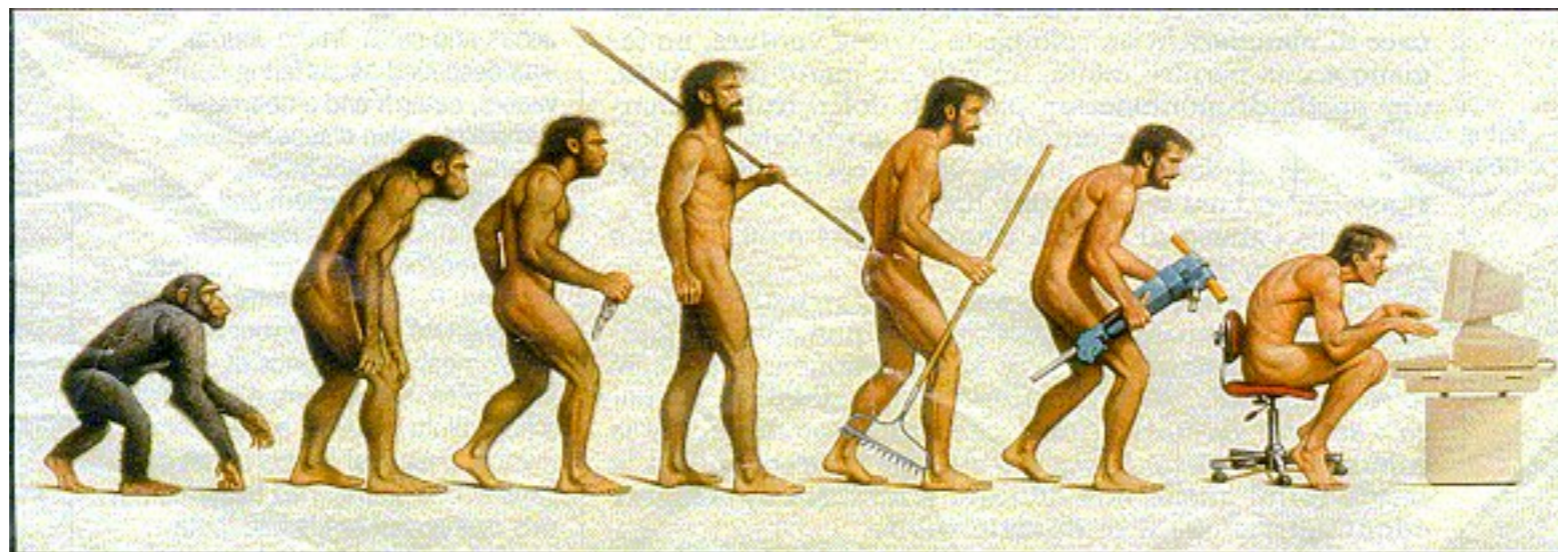
- beam-beam remnants, initial state radiation, multiple parton interactions
 - gives additional energy in the event

- All this additional energy has nothing to do with jet energies**

- have to subtract it**



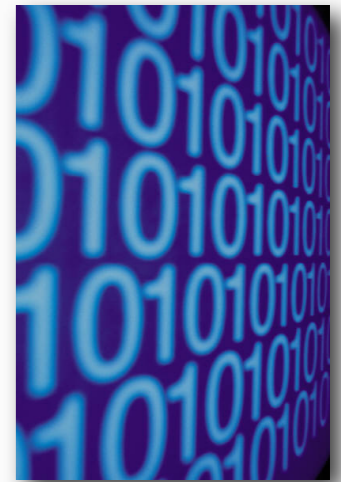
The computing behind all this



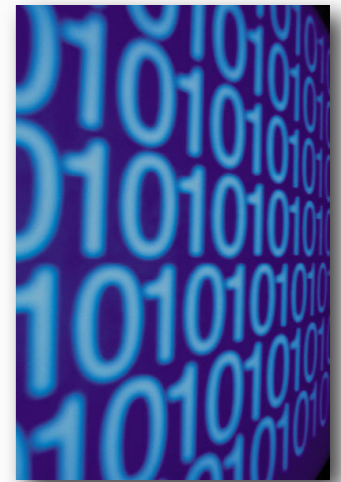
Somewhere, something went terribly wrong

Examples from CMS

- **Rate** of events streaming out from High-Level Trigger farm **~150 Hz**
- each event has a size of the order of **200 kByte**



- **Examples from CMS**
 - **Rate** of events streaming out from High-Level Trigger farm **~150 Hz**
 - each event has a size of the order of **200 kByte**
- at high Lumi : **CMS will record ~100k top-quark events per day**
 - among about 10^7 events in total per day
 - will have roughly 150 “physics” days per year
 - thus about 10^9 evts/year, a few **Pbyte**



Examples from CMS

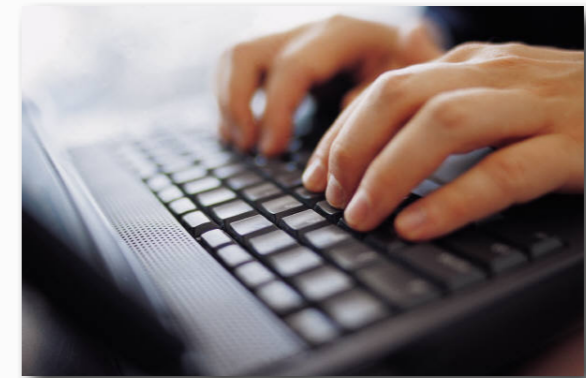
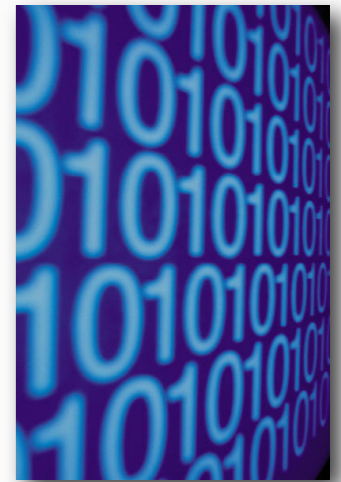
- **Rate** of events streaming out from High-Level Trigger farm ~ 150 Hz
- each event has a size of the order of **200 kByte**

at high Lumi : CMS will record ~ 100 k top-quark events per day

- among about 10^7 events in total per day
- will have roughly 150 “physics” days per year
- thus about 10^9 evts/year, a few **Pbyte**

“prompt” processing

- first reprocessing step within one day (within hours...)
- Reco time per event on std. CPU: < 5 sec (on Ixplus)
- Note : will have to reprocess several times
 - new/better algorithms, updated calibrations, etc.



Examples from CMS

- **Rate** of events streaming out from High-Level Trigger farm ~ 150 Hz
- each event has a size of the order of **200 kByte**

at high Lumi : CMS will record ~ 100 k top-quark events per day

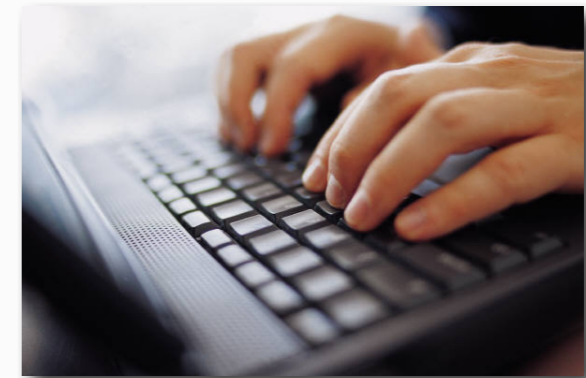
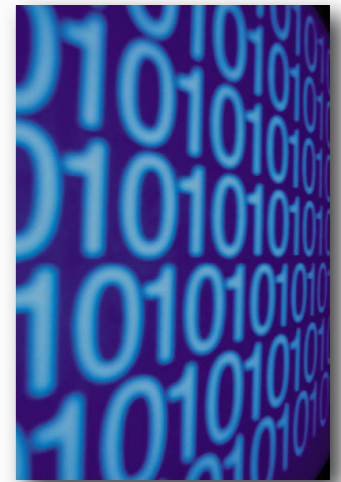
- among about 10^7 events in total per day
- will have roughly 150 “physics” days per year
- thus about 10^9 evts/year, a few **Pbyte**

“prompt” processing

- first reprocessing step within one day (within hours...)
- Reco time per event on std. CPU: < 5 sec (on lxplus)
- Note : will have to reprocess several times
 - new/better algorithms, updated calibrations, etc.

simulating several 100s to 1000s of millions of events

- are mostly done at computing centres outside CERN
- Simulation time per event now ~ 100 secs (eg. for QCD or top evts)



Examples from CMS

- **Rate** of events streaming out from High-Level Trigger farm **~150 Hz**
- each event has a size of the order of **200 kByte**

at high Lumi : CMS will record ~100k top-quark events per day

- among about 10^7 events in total per day
- will have roughly 150 “physics” days per year
- thus about 10^9 evts/year, a few **Pbyte**

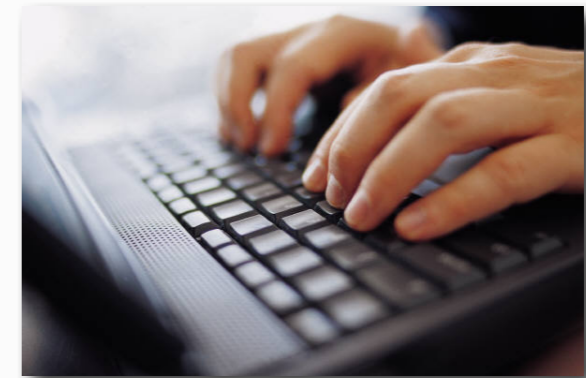
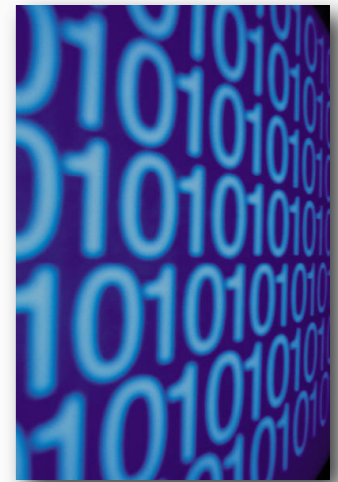
“prompt” processing

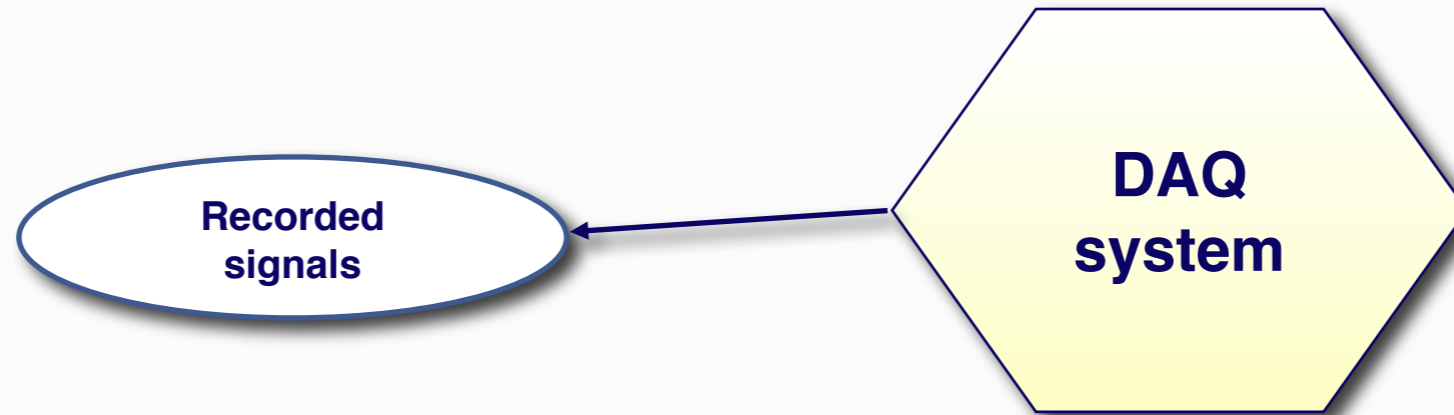
- first reprocessing step within one day (within hours...)
- Reco time per event on std. CPU: < 5 sec (on Ixplus)
- Note : will have to reprocess several times
 - new/better algorithms, updated calibrations, etc.

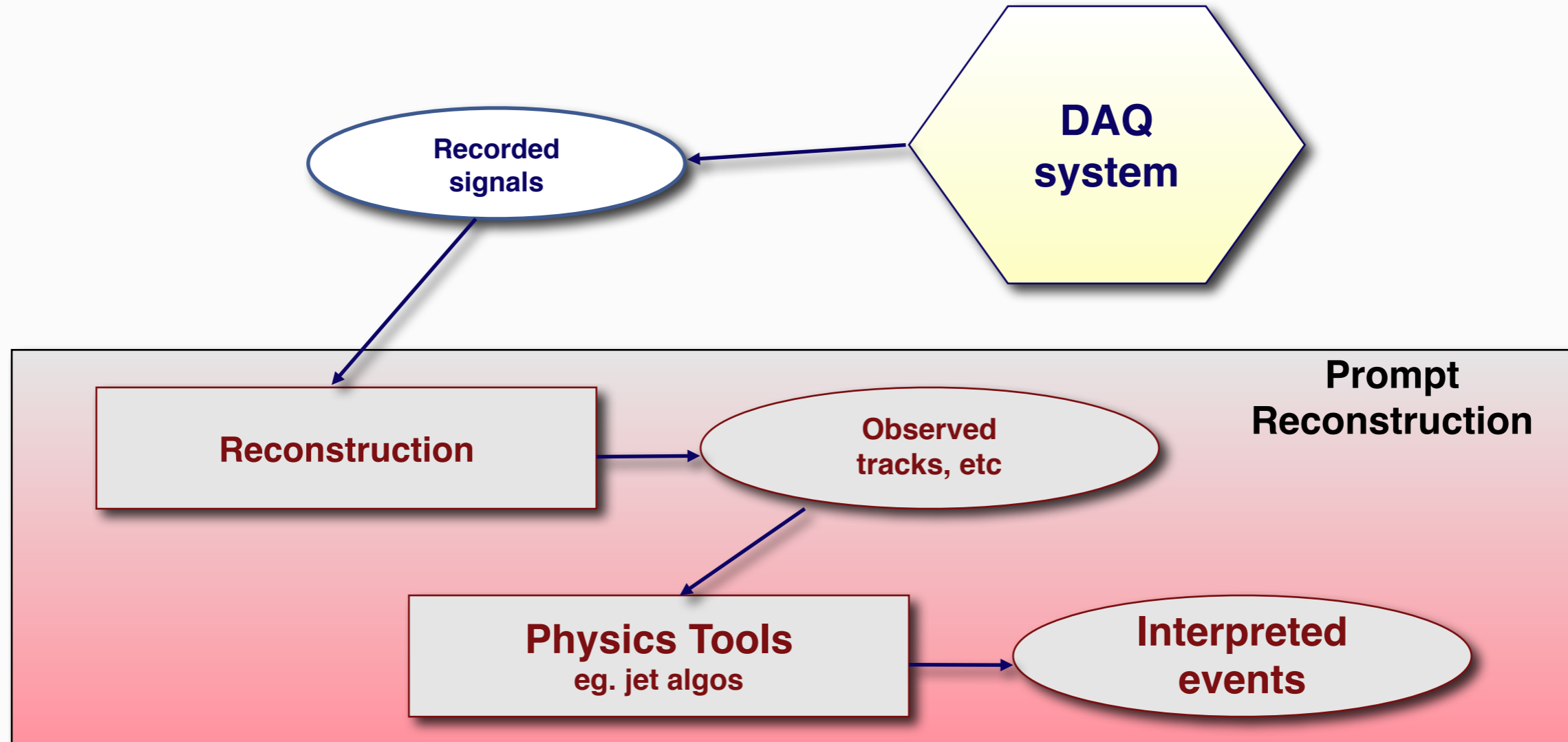
simulating several 100s to 1000s of millions of events

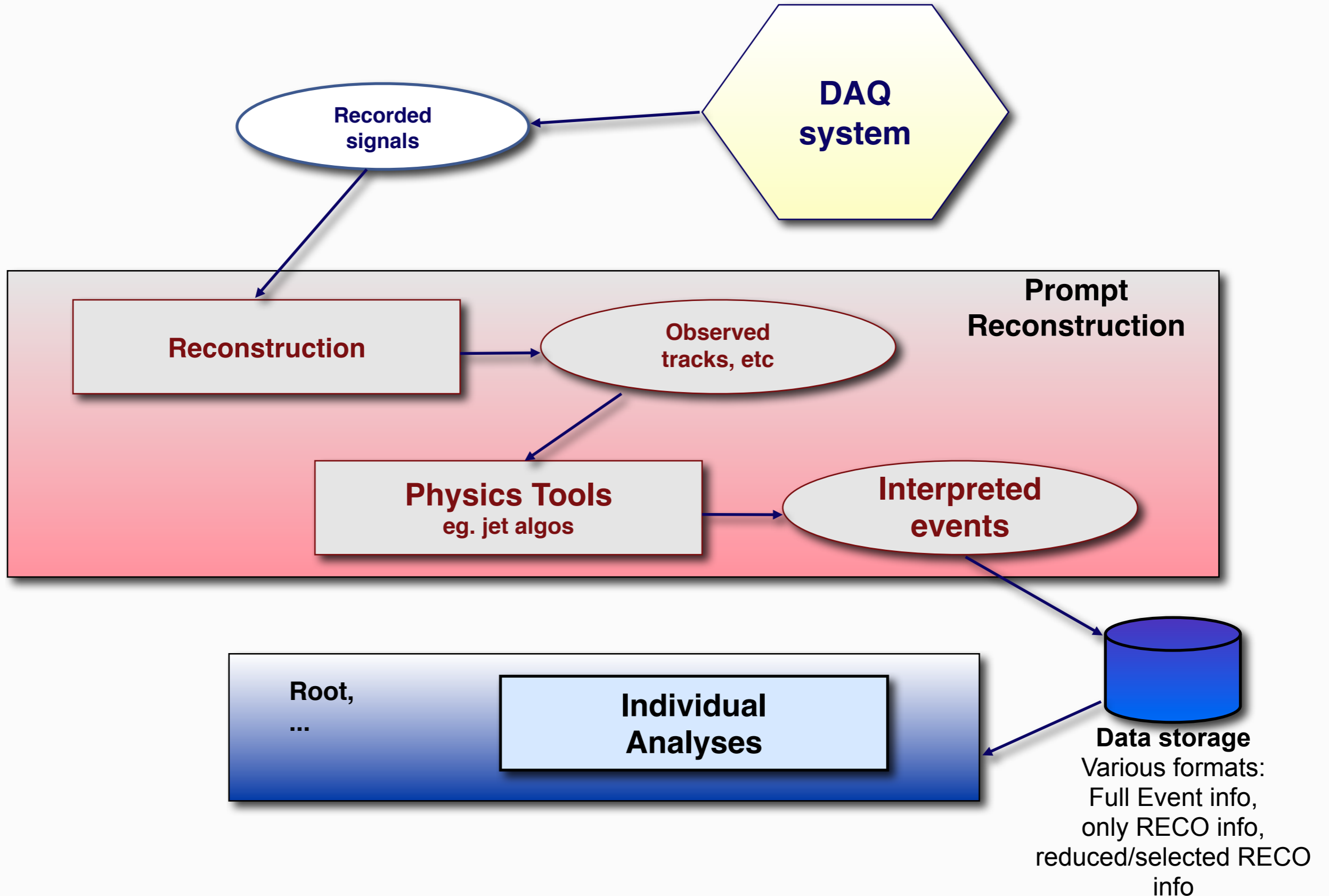
- are mostly done at computing centres outside CERN
- Simulation time per event now ~ 100 secs (eg. for QCD or top evts)

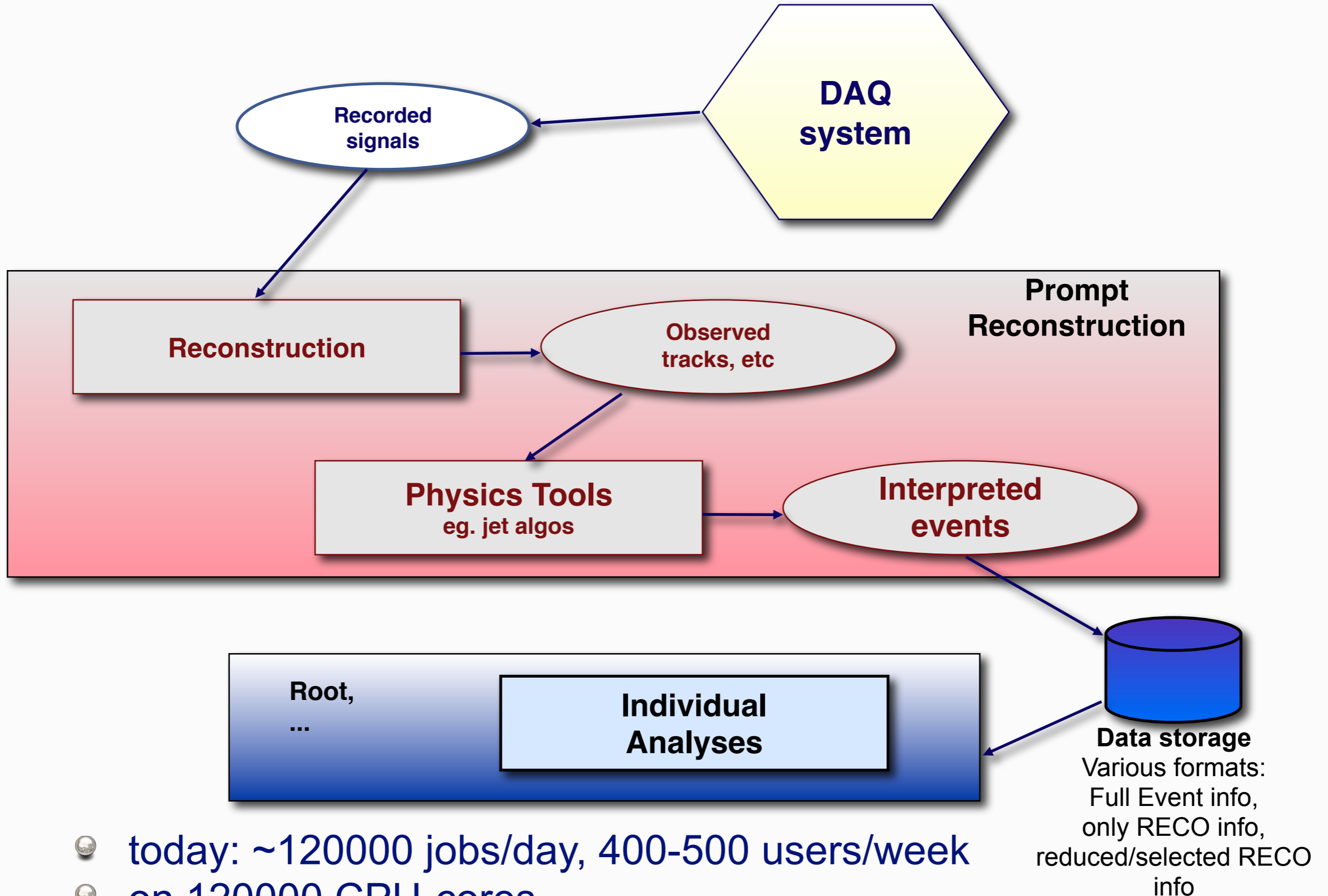
~2 million lines of code (reconstruction and simulation)



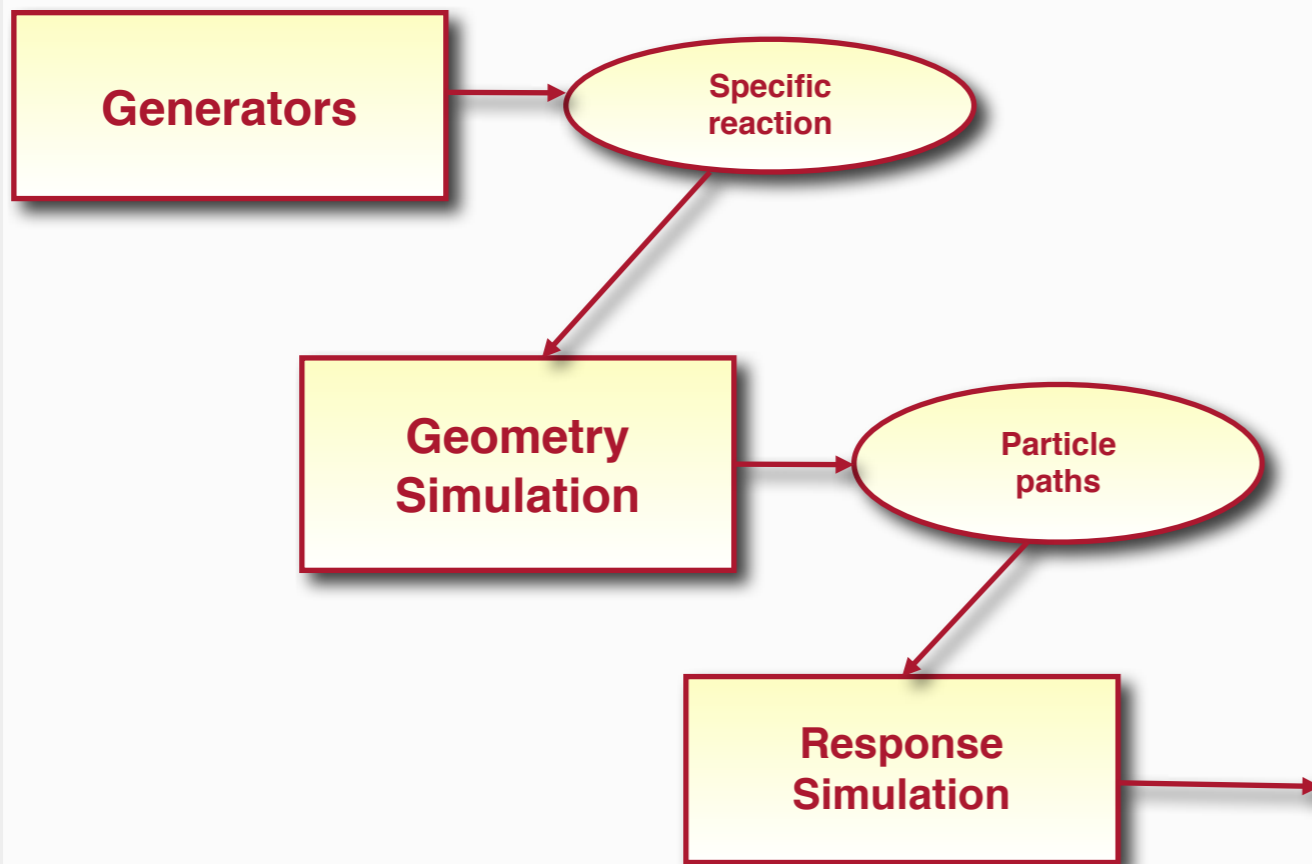




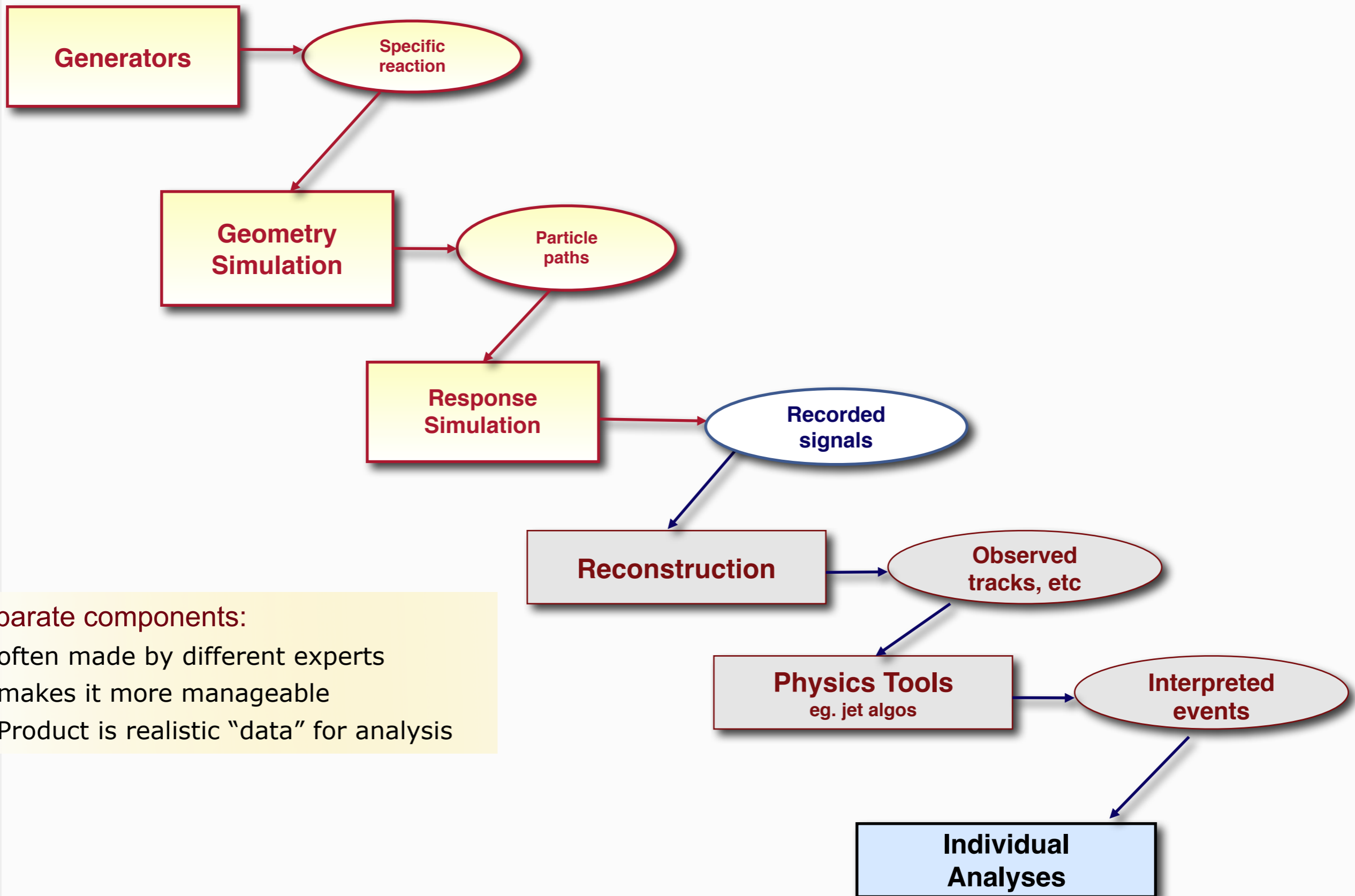




- today: ~120000 jobs/day, 400-500 users/week
- on 120000 CPU-cores



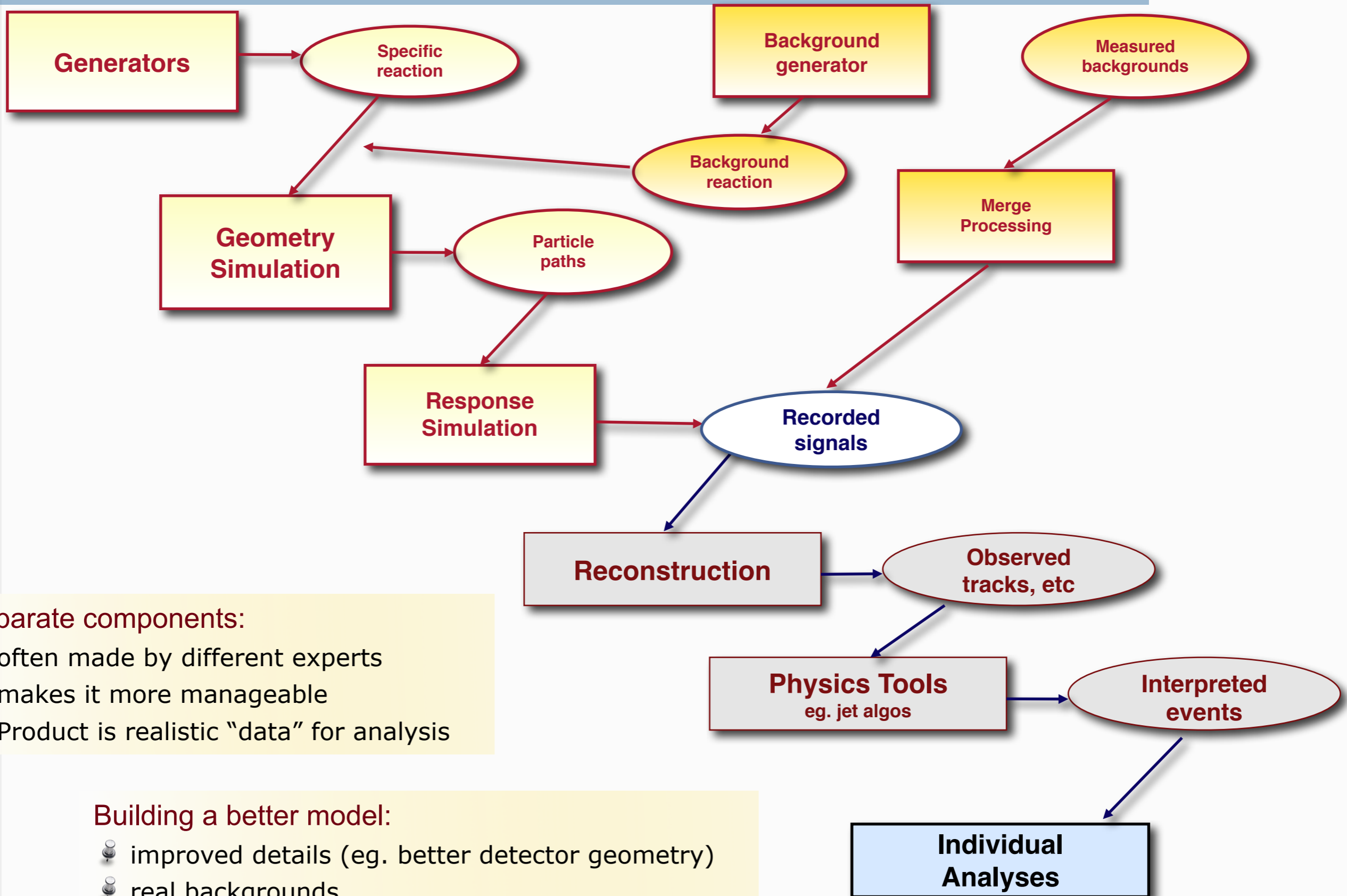
Flow of simulated data



Separate components:

- often made by different experts
- makes it more manageable
- Product is realistic "data" for analysis

Flow of simulated data

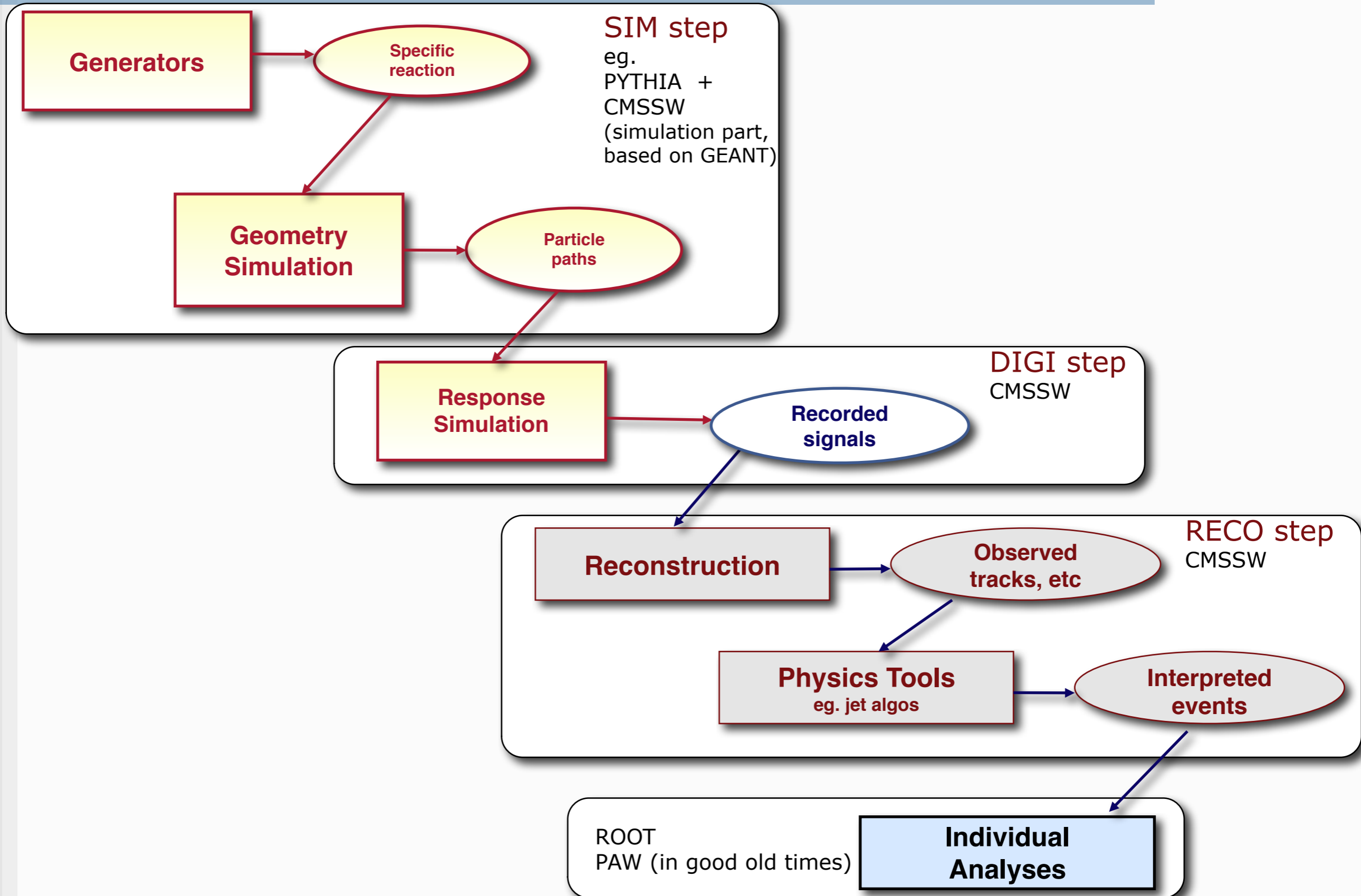


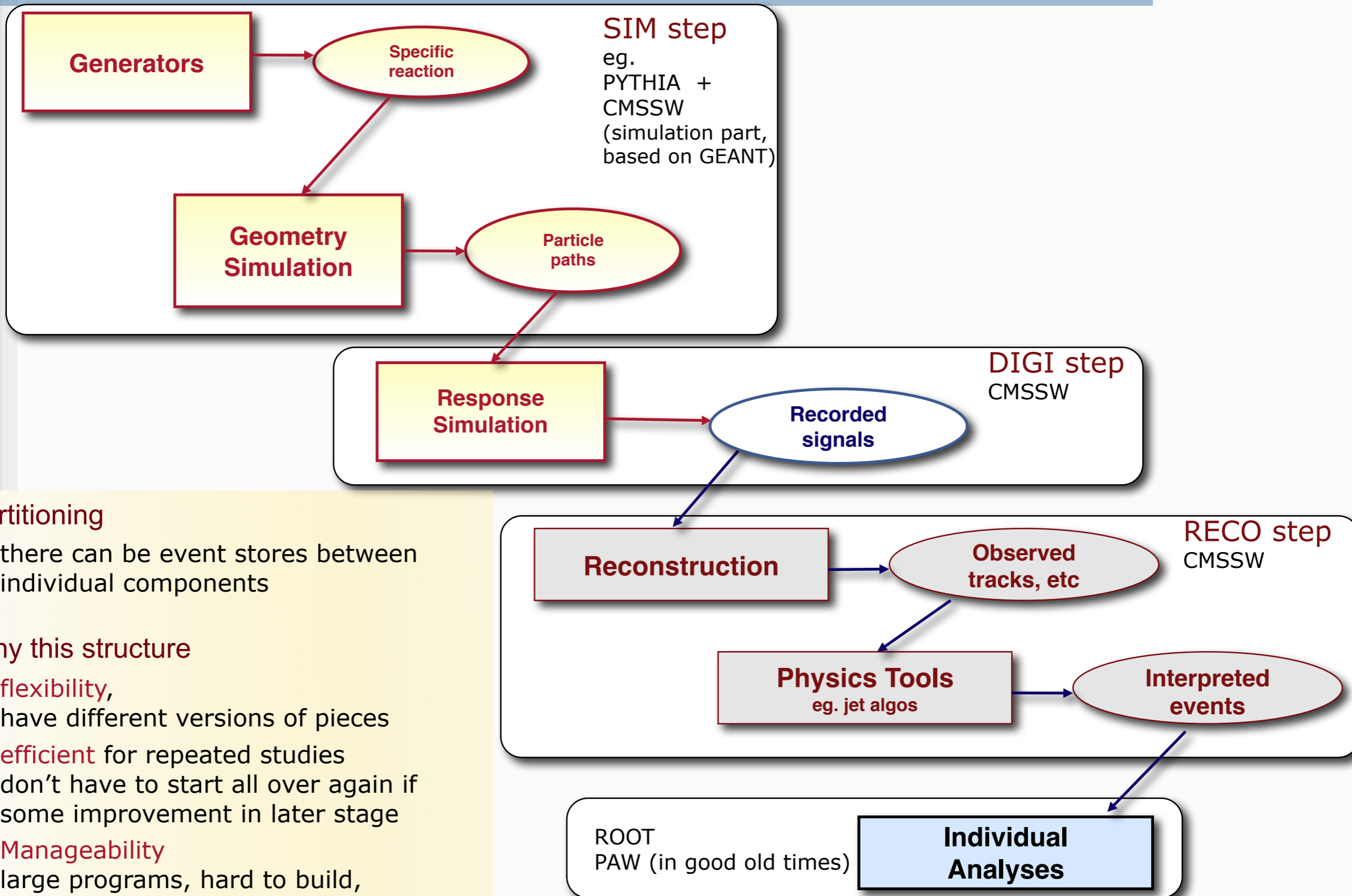
Separate components:

- often made by different experts
- makes it more manageable
- Product is realistic "data" for analysis

Building a better model:

- improved details (eg. better detector geometry)
- real backgrounds



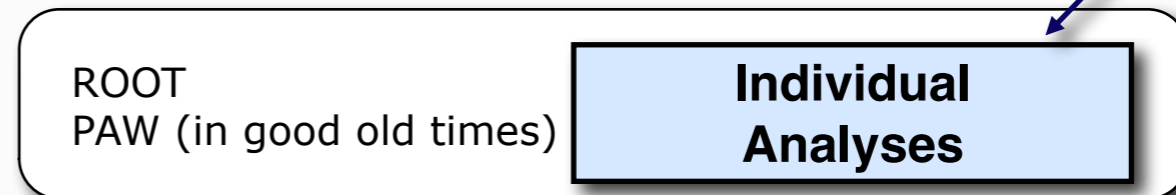
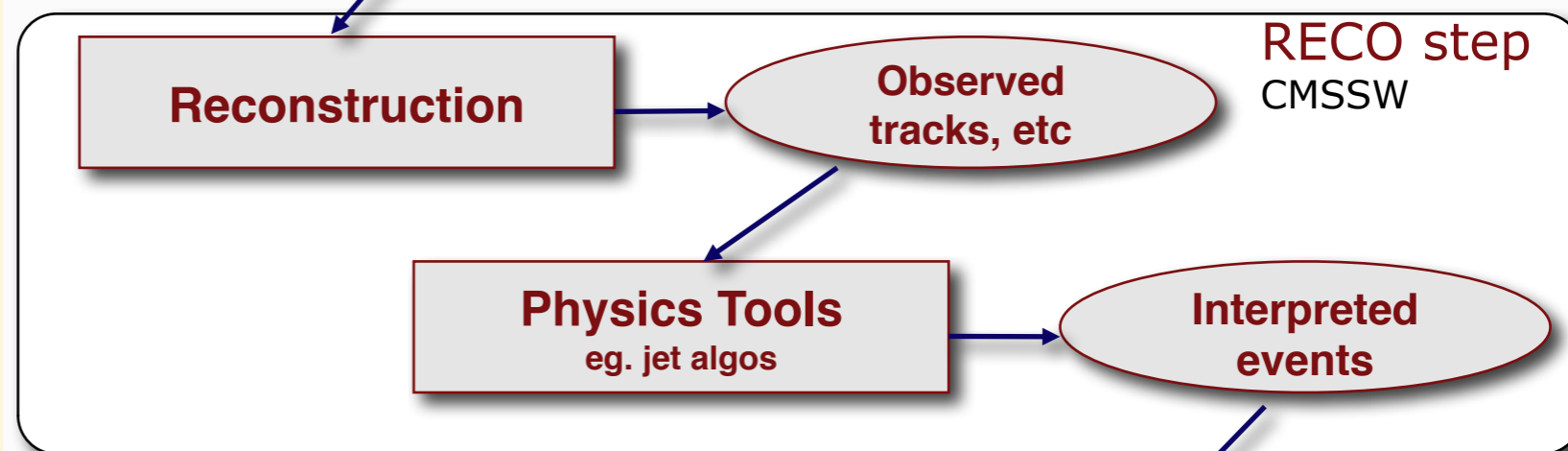
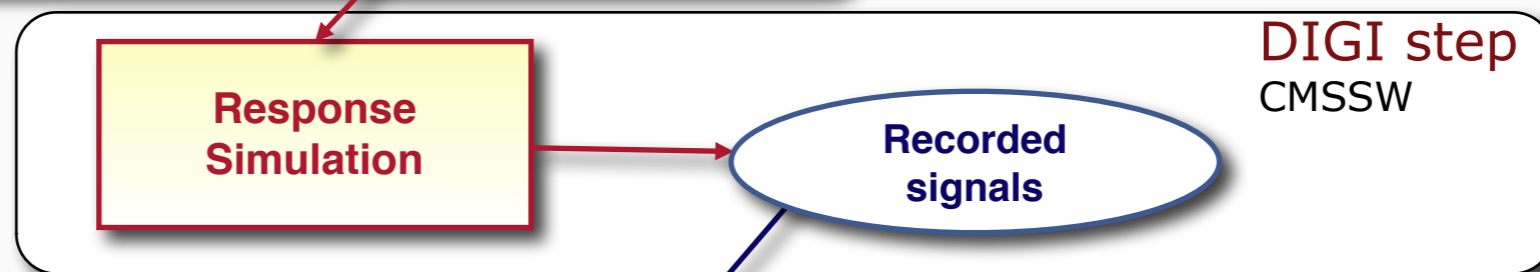
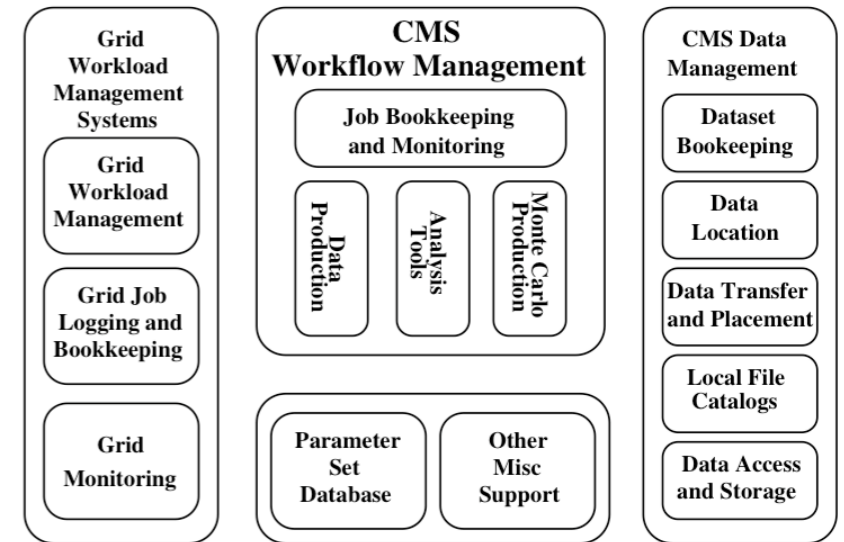
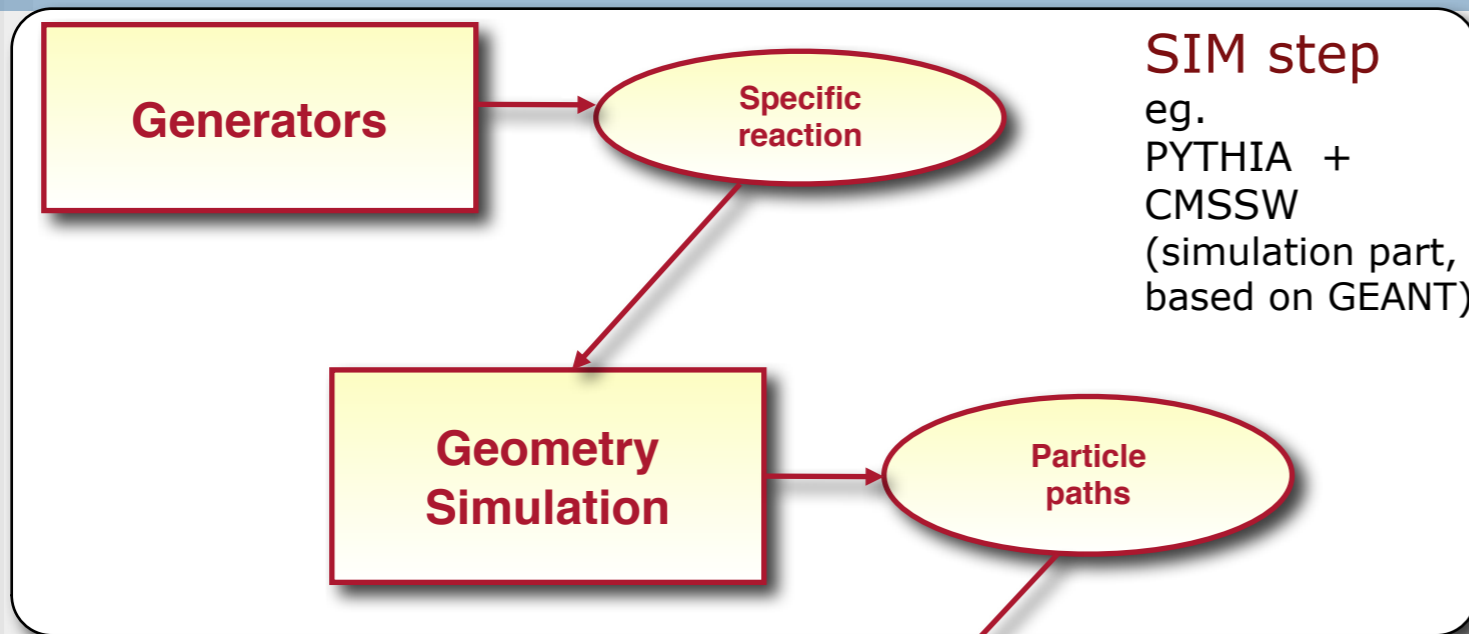


Partitioning

- there can be event stores between individual components

Why this structure

- flexibility**, have different versions of pieces
- efficient** for repeated studies don't have to start all over again if some improvement in later stage
- Manageability**
large programs, hard to build, understand, debug, maintain, ...



Partitioning

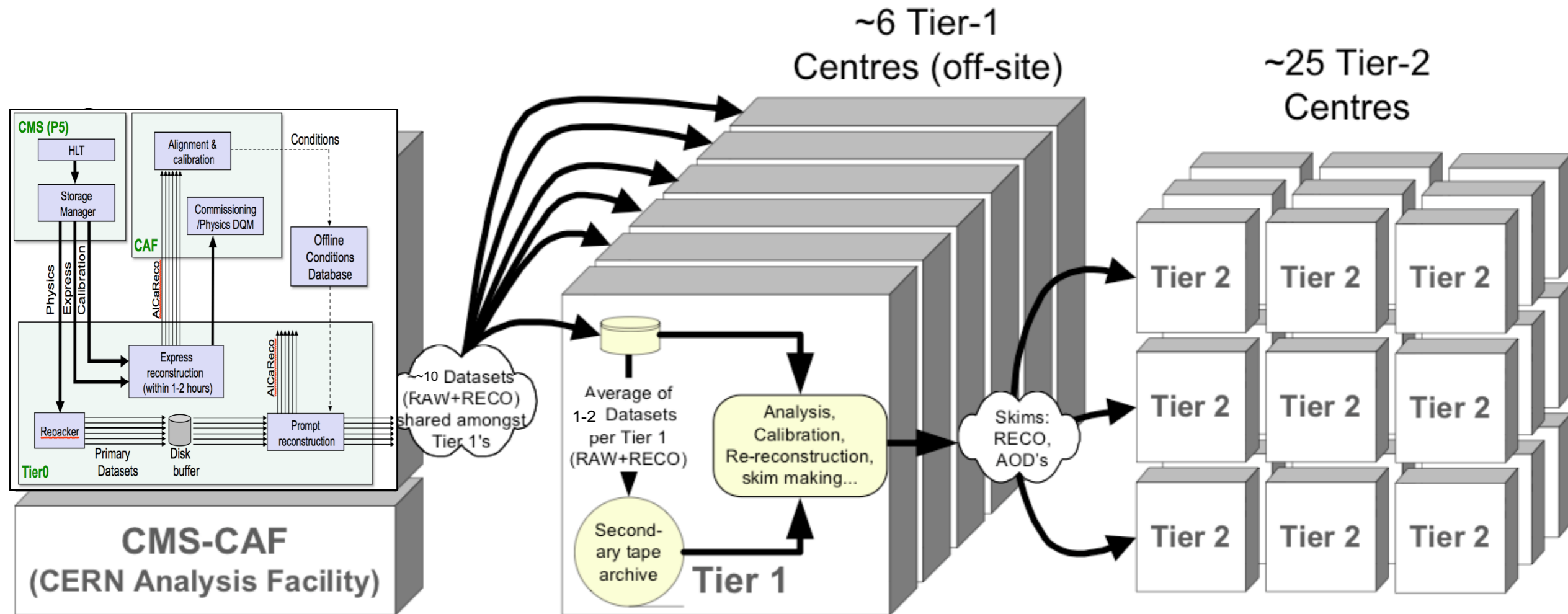
- there can be event stores between individual components

Why this structure

- flexibility**, have different versions of pieces
- efficient** for repeated studies don't have to start all over again if some improvement in later stage
- Manageability**
large programs, hard to build, understand, debug, maintain, ...

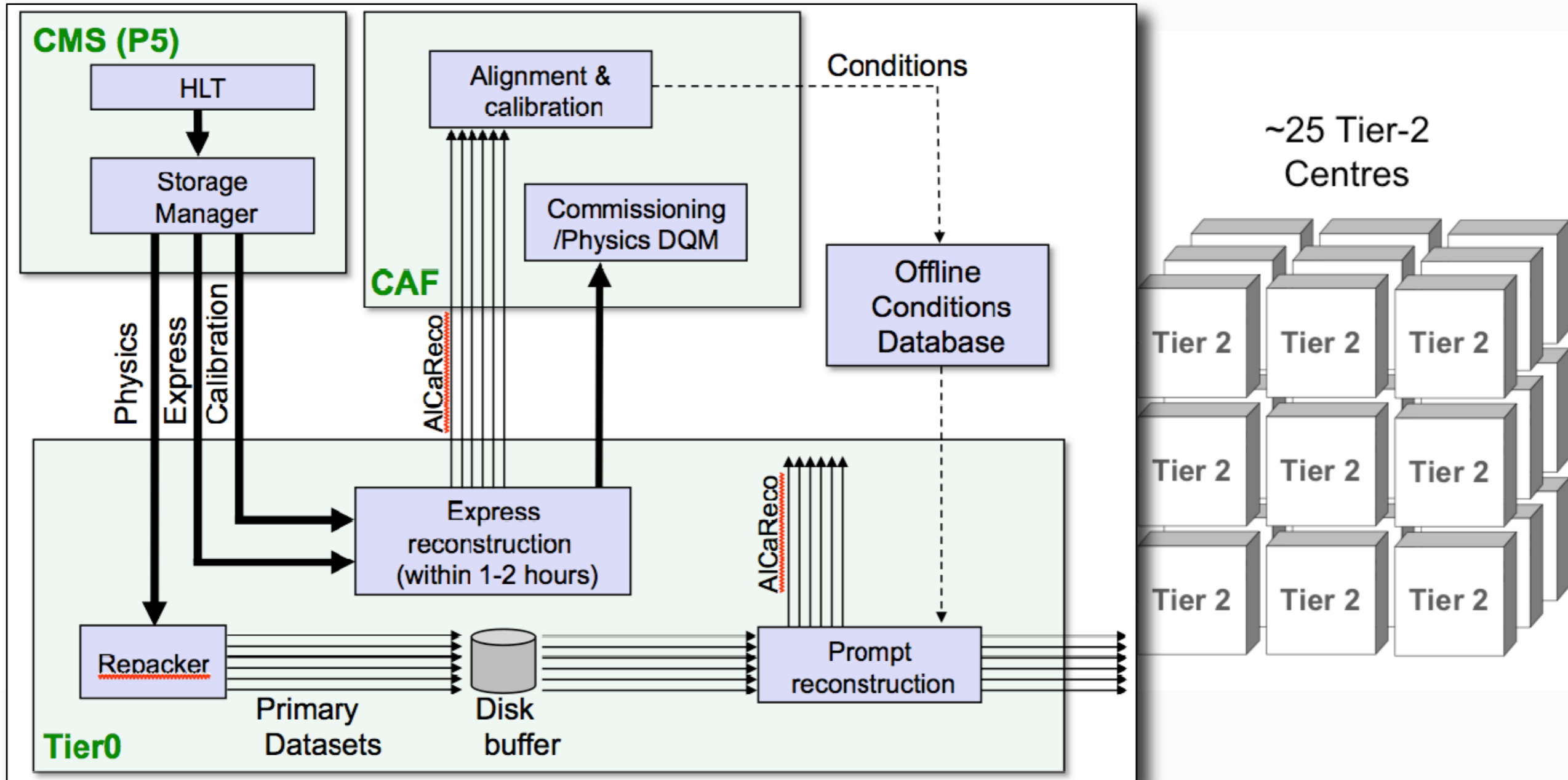


CMS Computing Model: Data Flow



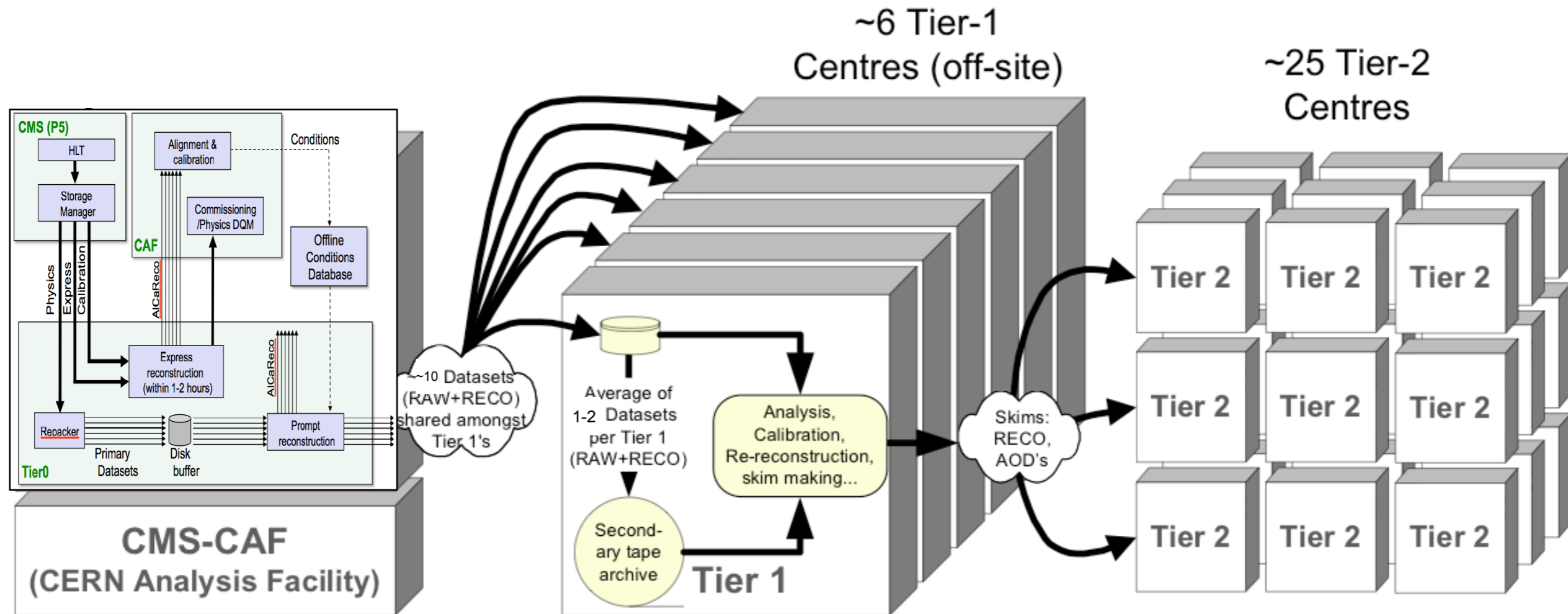


CMS Computing Model: Data Flow





CMS Computing Model: Data Flow



Summary

“Doing something ordinary is a waste of time” (Madonna)

- Details on track fitting, Kalman filters
- Secondary Vertex finding
- Alignment
- Particle Identification
- Calibration techniques, “in-situ” methods
- Particle/Energy flow
- Trigger menus, their studies
- more details on parameter fitting, eg. lifetime and mass measurements
- how to estimate systematic errors
- Databases, persistent data storage
- Programming languages in use (F77, C, C++, JAVA, ...)
-

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **Reconstruction and Analysis**
is how we get from raw data to physics papers
- On your way

- **Reconstruction and Analysis**
is how we get from raw data to physics papers
- **On your way**
 - first you have too much information → reduce

- **Reconstruction and Analysis**
is how we get from raw data to physics papers
- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **What makes it hard, but also exciting**

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **What makes it hard, but also exciting**
 - many many cross checks

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **What makes it hard, but also exciting**
 - many many cross checks
 - more cross checks

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **What makes it hard, but also exciting**
 - many many cross checks
 - more cross checks
 - sometimes some “art” involved

- **Reconstruction and Analysis**
is how we get from raw data to physics papers

- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses

- **What makes it hard, but also exciting**
 - many many cross checks
 - more cross checks
 - sometimes some “art” involved
 - tuning, evolutionary improvement

- **Reconstruction and Analysis**
is how we get from raw data to physics papers
- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses
- **What makes it hard, but also exciting**
 - many many cross checks
 - more cross checks
 - sometimes some “art” involved
 - tuning, evolutionary improvement
- **Even to me it is often a miracle that we can generate wonderful results from these complicated instruments!**

- **Reconstruction and Analysis**
is how we get from raw data to physics papers
- **On your way**
 - first you have too much information → reduce
 - sometimes too little information or little prior knowledge
 - make hypotheses
- **What makes it hard, but also exciting**
 - many many cross checks
 - more cross checks
 - sometimes some “art” involved
 - tuning, evolutionary improvement
- **Even to me it is often a miracle that we can generate wonderful results from these complicated instruments!**

