

Single W Boson Photoproduction in p-p and p-A collisions

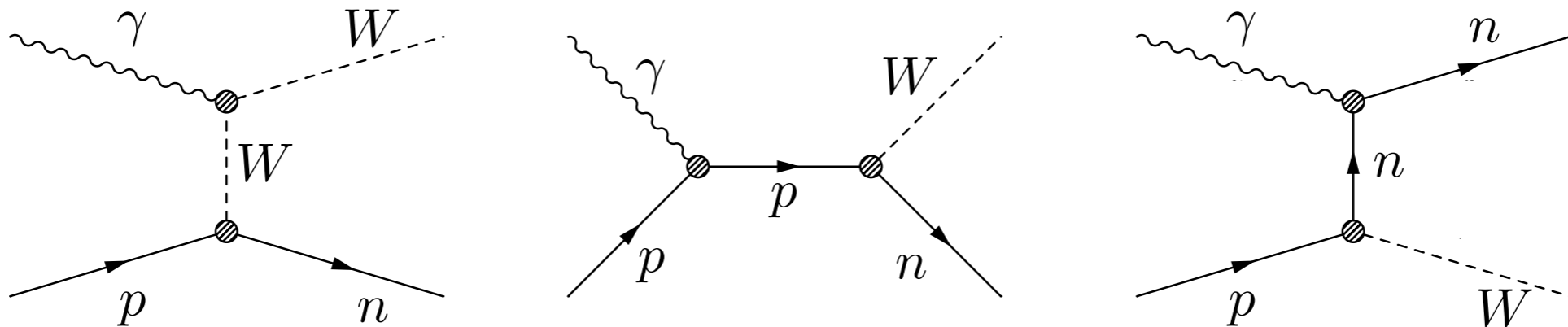
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Motivation

- the couplings of gauge bosons among themselves belong to one of the least tested sectors of electroweak theory
- the photoproduction of single W bosons is a process well-suited to test the $WW\gamma$ coupling
- up to now very low rates for processes involving triple gauge boson coupling (HERA: 3 events for inclusive photoproduction [*Breitweg et al.*, Phys. Lett. B 471, 411 (2000)])
- Can these rates be improved in p-p and p-A collisions at LHC?
- exclusive photoproduction: neutron in forward direction
- in p-p collisions contributions from elastic and inelastic photon spectra

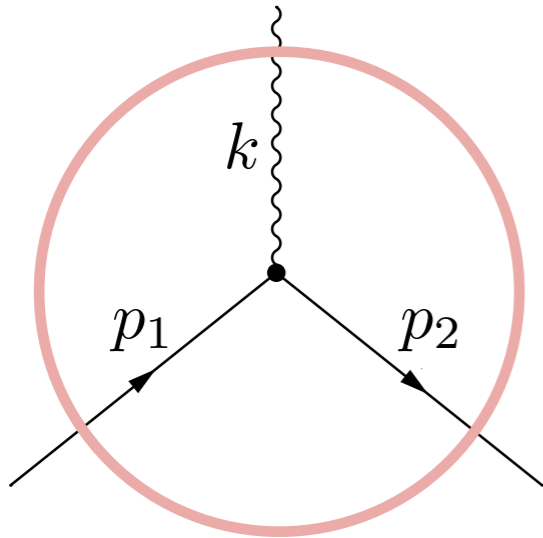
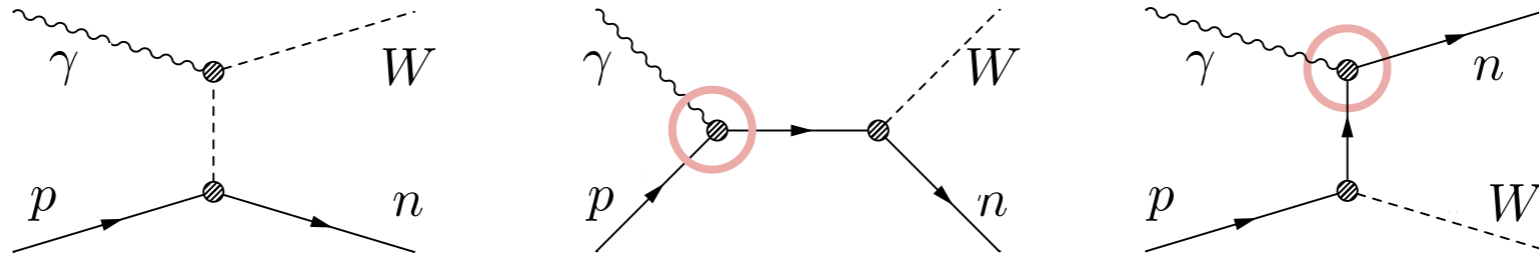
Exclusive Photoproduction of W

- we include three Feynman diagrams in our calculation



- appropriate electromagnetic and weak form factors have to be employed

Electromagnetic NN Vertex



$$\Gamma_{\mu}^{\gamma NN} = -ie \left[F_1^N(k^2) \gamma_{\mu} + i \frac{\kappa_N}{2M} F_2^N(k^2) \sigma_{\mu\nu} (p_2 - p_1)^{\nu} \right]$$

$$F_1^p(0) = F_2^p(0) = F_2^n(0) = 1$$

$$F_1^n(0) = 0$$

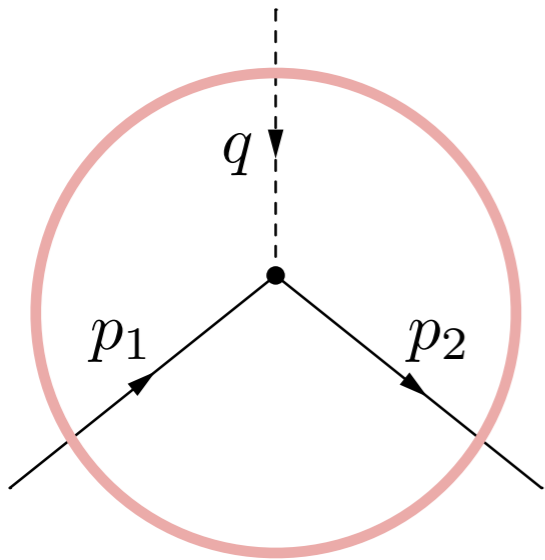
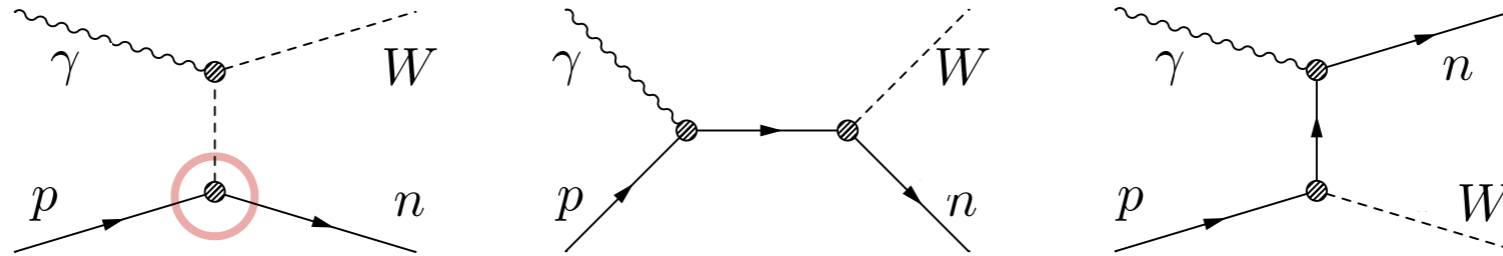
$$\kappa_p = 1.79$$

$$\kappa_n = -1.91$$

$$\sigma_{\mu\nu} = \frac{i}{2} [\gamma_{\mu}, \gamma_{\nu}]$$

Weak NN Vertex

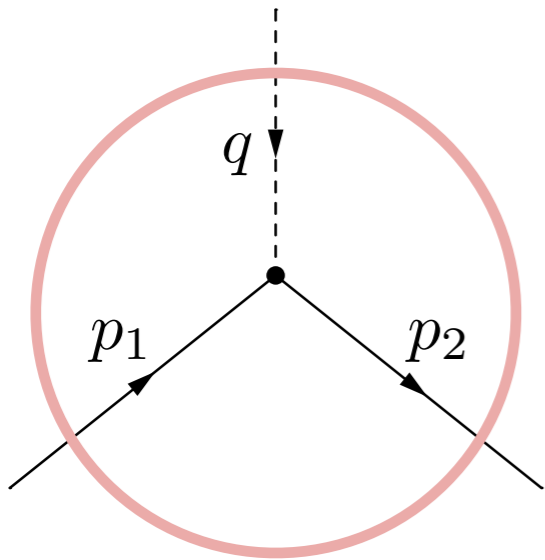
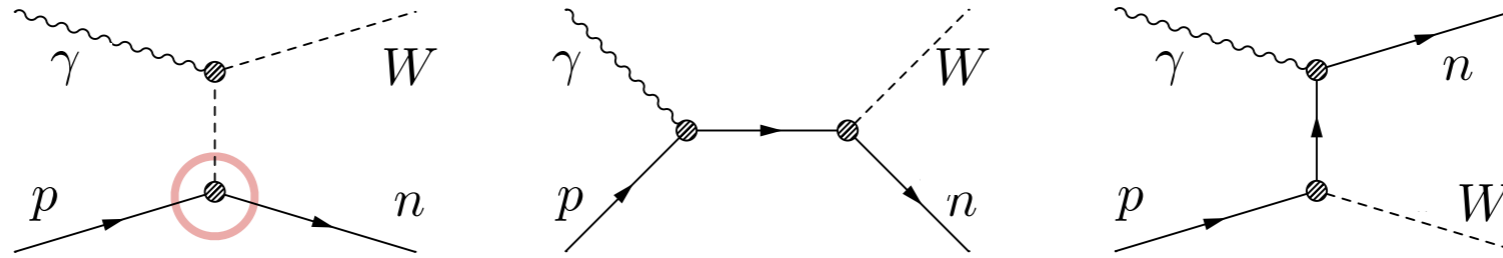
(spacelike region)



$$\Gamma_{\mu}^{WNN} = -ig [F_V(q^2)\gamma_{\mu} + iF_M(q^2)\sigma_{\mu\nu}q^{\nu} + F_S(q^2)q_{\mu} + F_A(q^2)\gamma_{\mu}\gamma_5 + iF_T(q^2)\sigma_{\mu\nu}q^{\nu}\gamma_5 + F_P(q^2)q_{\mu}\gamma_5]$$

Weak NN Vertex

(spacelike region)



$$\Gamma_{\mu}^{WNN} = -ig [F_V(q^2)\gamma_{\mu} + iF_M(q^2)\sigma_{\mu\nu}q^{\nu} + \cancel{F_S(q^2)}q_{\mu} + F_A(q^2)\gamma_{\mu}\gamma_5 + i\cancel{F_T(q^2)}\sigma_{\mu\nu}q^{\nu}\gamma_5 + \cancel{F_P(q^2)}q_{\mu}\gamma_5]$$

- conserved vector current hypothesis:

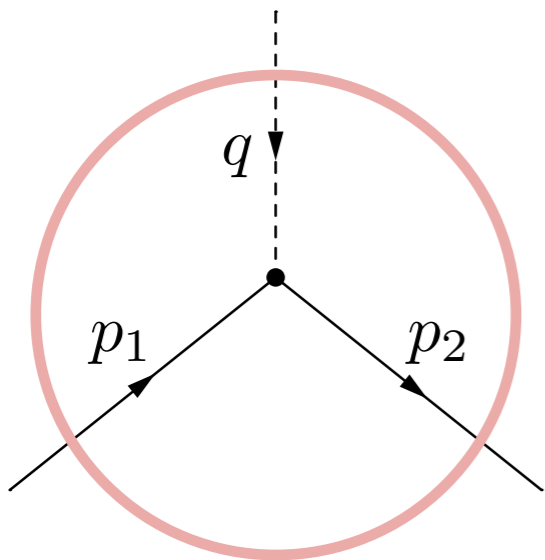
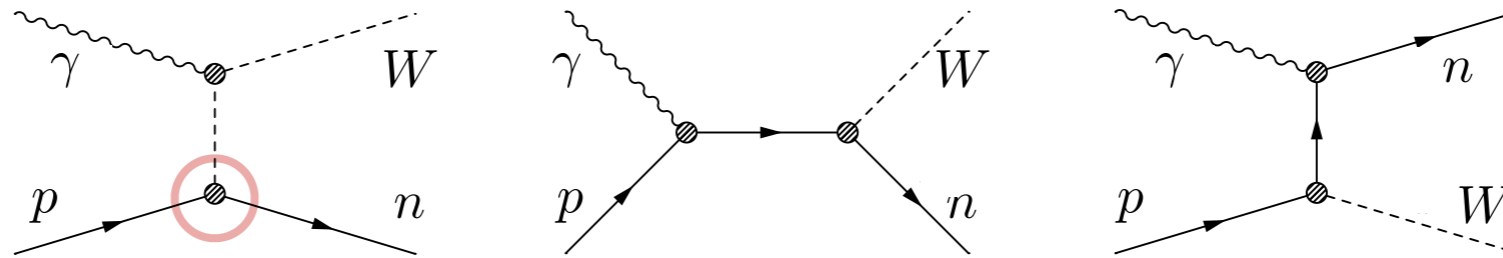
$$F_V(q^2) = F_1^p(q^2) - F_1^n(q^2) = \frac{1 + \tau(1 + (\kappa_p - \kappa_n))}{1 + \tau} G_V(q^2) \quad F_V(0) = 1$$

$$F_M(q^2) = \frac{1}{2M} (\kappa_p F_2^p(q^2) - \kappa_n F_2^n(q^2)) = \frac{1}{2M} \frac{\kappa_p - \kappa_n}{1 + \tau} G_V(q^2)$$

$$G_V(q^2) = \left(1 - \frac{q^2}{m_V^2}\right)^{-2} \quad \tau = \frac{q^2}{4M^2} \quad m_V^2 = 0.71 \text{ GeV}^2$$

Weak NN Vertex

(spacelike region)



$$\begin{aligned}
 \Gamma_{\mu}^{WNN} = & -ig [F_V(q^2)\gamma_{\mu} + iF_M(q^2)\sigma_{\mu\nu}q^{\nu} + \cancel{F_S(q^2)}q_{\mu} \\
 & + F_A(q^2)\gamma_{\mu}\gamma_5 + i\cancel{F_T}(q^2)\sigma_{\mu\nu}q^{\nu}\gamma_5 \\
 & + \cancel{F_P}(q^2)q_{\mu}\gamma_5]
 \end{aligned}$$

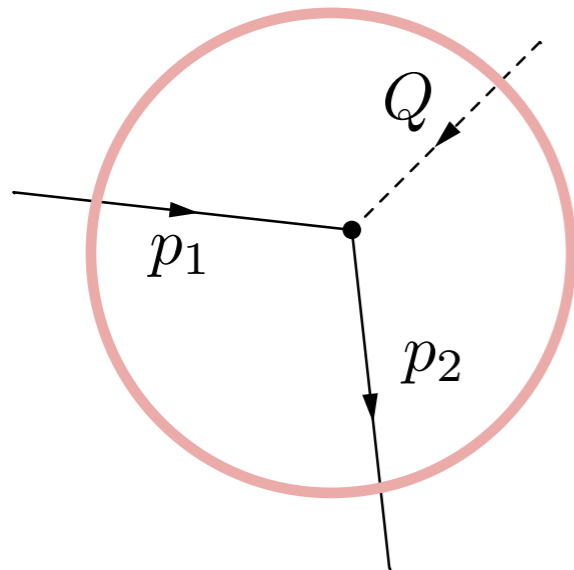
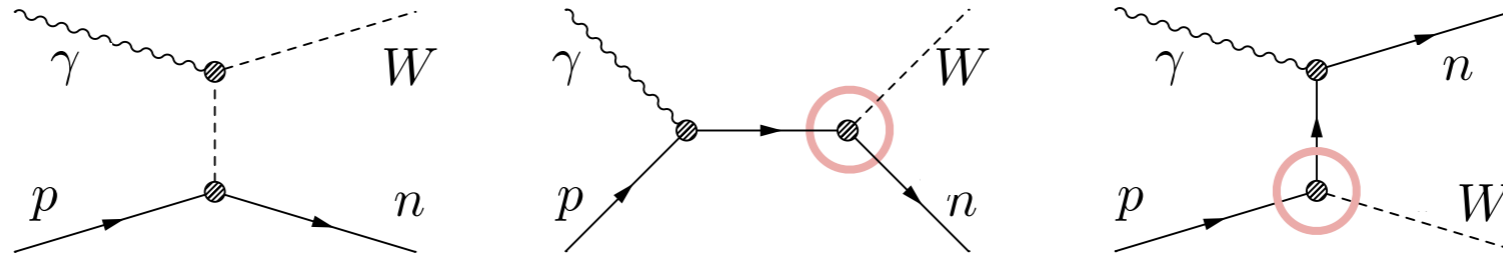
- by analogy with vector form factors, the axial vector form factor $F_A(q^2)$ is usually parameterized as a dipole form

$$F_A(q^2) = F_A(0) \left(1 + \frac{q^2}{m_A^2}\right)^{-2}$$

with $F_A(0) = -1.26$ and $m_A = 0.95 \text{ GeV}$

Weak NN Vertex

(timelike region)

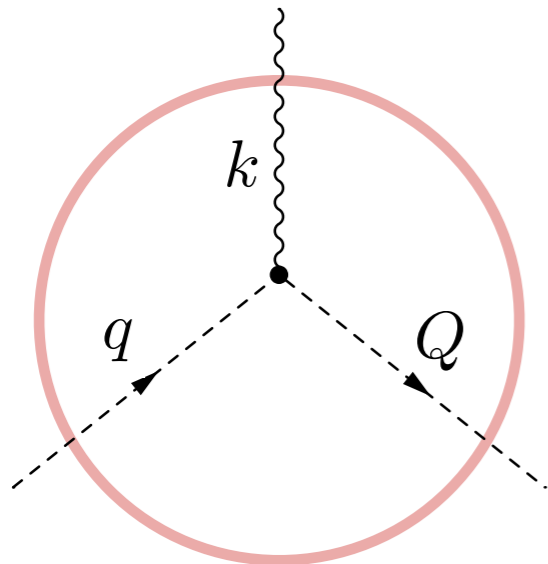
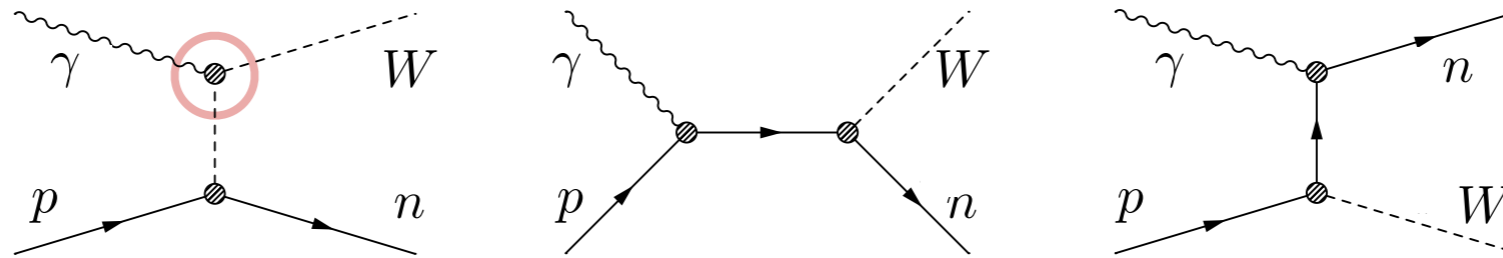


$$\Gamma_{\mu}^{WNN} = -ig [F_V(Q^2)\gamma_{\mu} + iF_M(Q^2)\sigma_{\mu\nu}Q^{\nu} + F_A(Q^2)\gamma_{\mu}\gamma_5]$$

here: $Q^2 = M_W^2$

- *Fearing et al.* [Phys. Rev. D5, 158 & 177 (1972)] and *Kallianpur* [Phys. Rev. D34, 3343 (1986)] use two different choices of timelike weak form factors:
 - ▶ constant timelike form factors
 - ▶ dipole timelike form factors

Triple Gauge Boson Coupling



$$\Gamma_{\mu\alpha\beta}^{WW\gamma} = ie \left\{ F_1^W(k^2) [(2Q - k)_\mu g_{\alpha\beta} - Q_\alpha g_{\mu\beta} - (Q - k)_\beta g_{\mu\alpha}] + \kappa_W F_2^W(k^2) [k_\beta g_{\mu\beta} - k_\alpha g_{\mu\beta}] \right\}$$

$$F_1^W(0) = F_2^W(0) = 1$$

- Standard Model Coupling recovered for $\kappa_W = 1$

Gauge Invariance

$$\mathcal{M}_{fi} = -ige \epsilon_\mu M^{\mu\beta} \epsilon_\beta^W$$

$$M^{\mu\beta} = M_1^{\mu\beta} + M_2^{\mu\beta} + M_3^{\mu\beta}$$

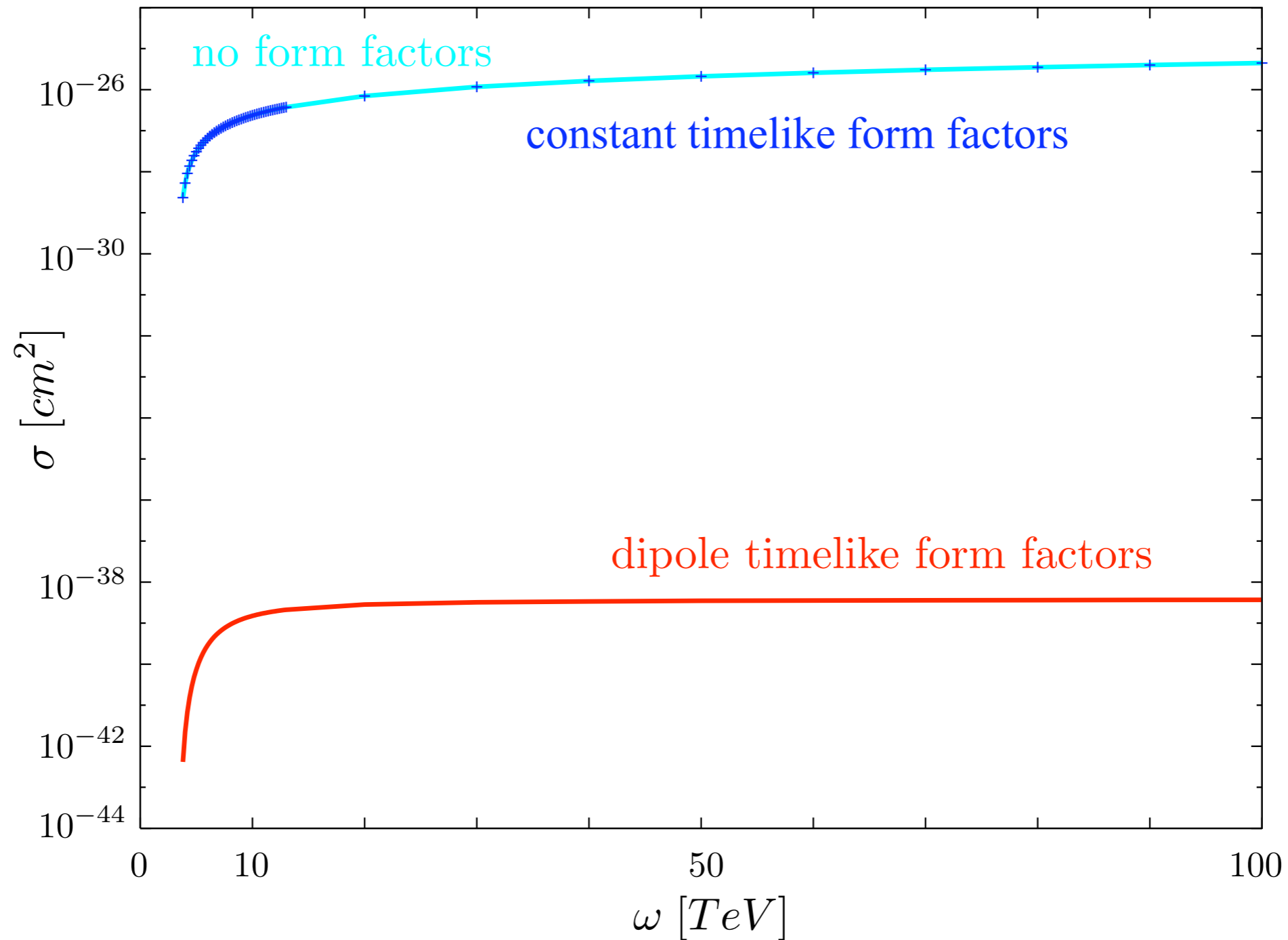
- employing weak vertices between offshell states results in loss of explicit gauge invariance

$$k_\mu M^{\mu\beta} \epsilon_\beta^W \neq 0$$

- *Fearing et al.* [Phys. Rev. D 5, 158 (1972)] introduce a technique how to maintain gauge invariance by adding a term $\Delta M^{\mu\beta}$
- $\Delta M^{\mu\beta}$ has to fulfill certain requirements:
 - ▶ it should cancel the extra terms arising
 - ▶ it should not contain new singularities in the physical region
 - ▶ it should satisfy $\Delta M \ll M$

$$M^{\mu\beta} \longrightarrow M^{\mu\beta} + \Delta M^{\mu\beta}$$

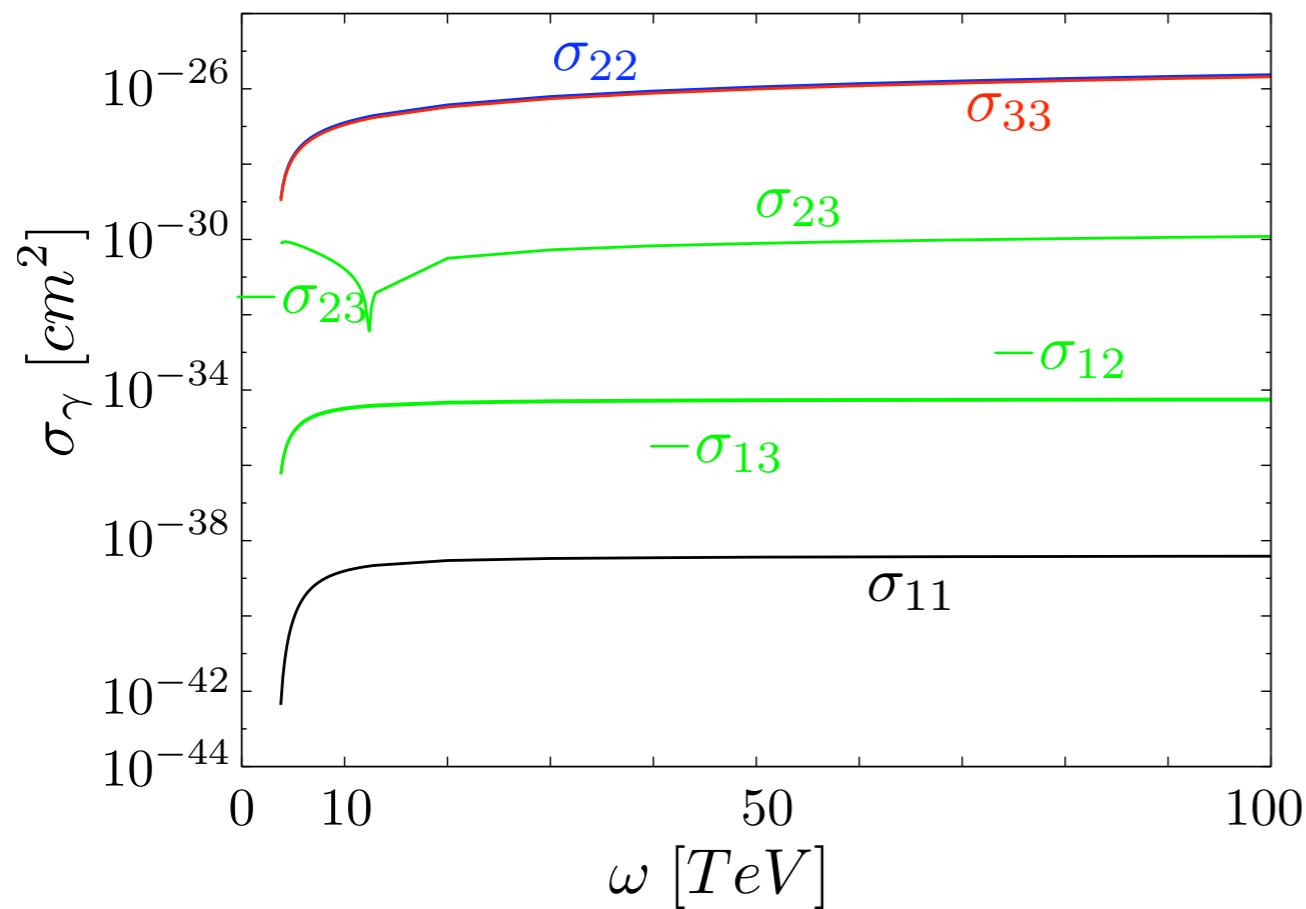
Results (Photoproduction)



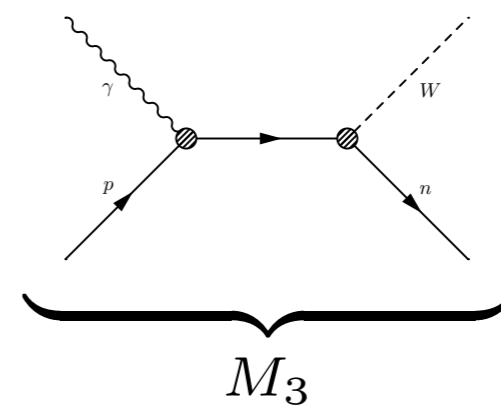
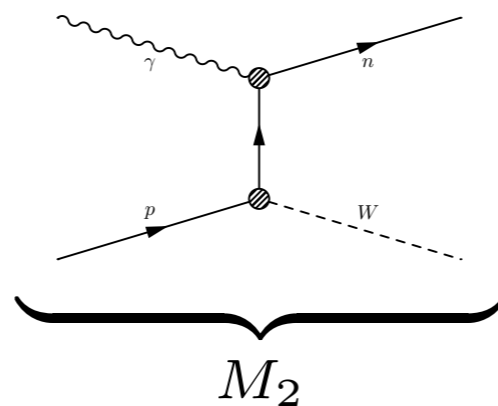
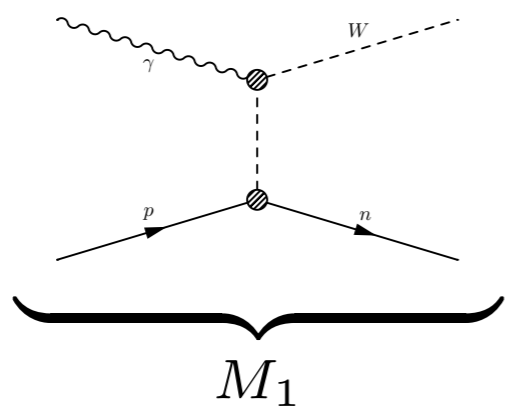
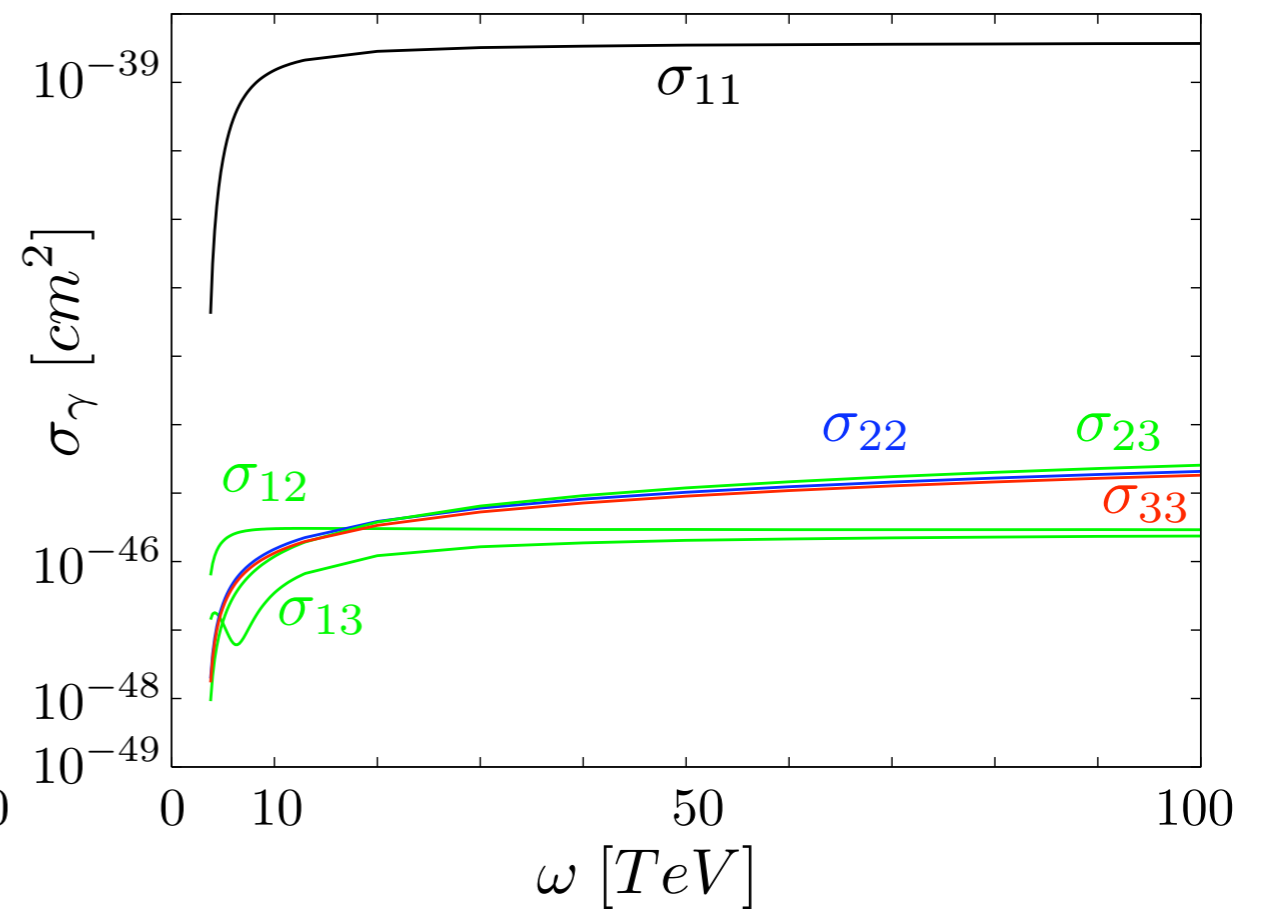
- total photoproduction cross section as function of photon energy without form factors and with constant and dipole weak timelike form factors

Results (Photoproduction)

constant timelike form factors

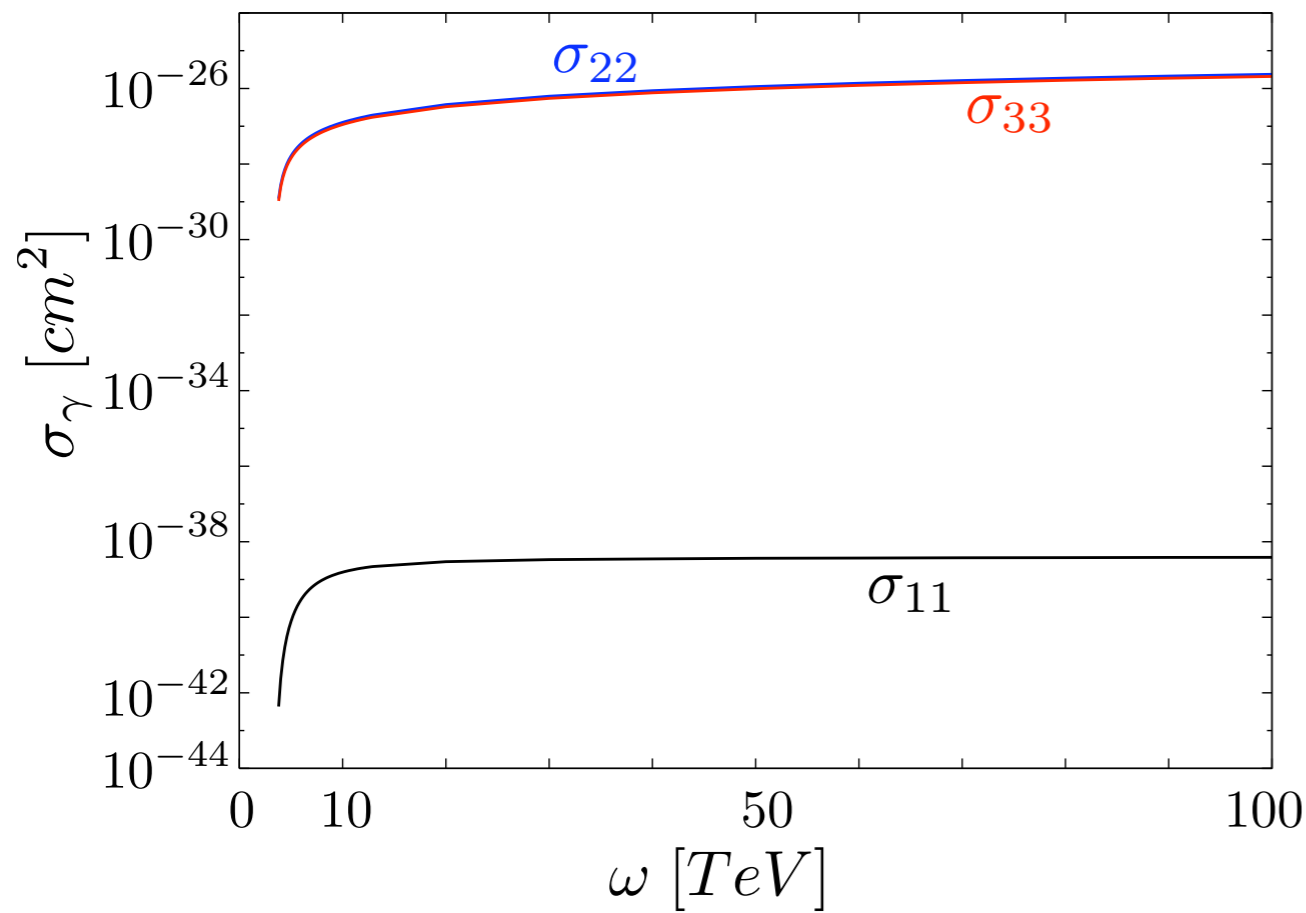


dipole timelike form factors

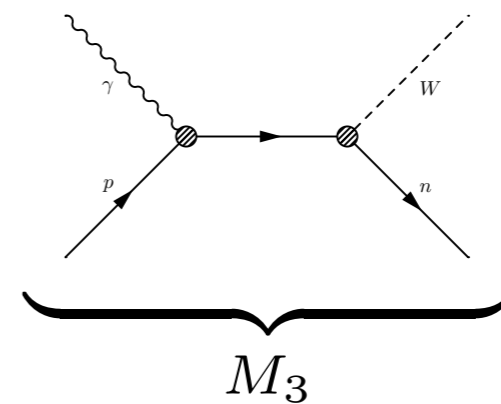
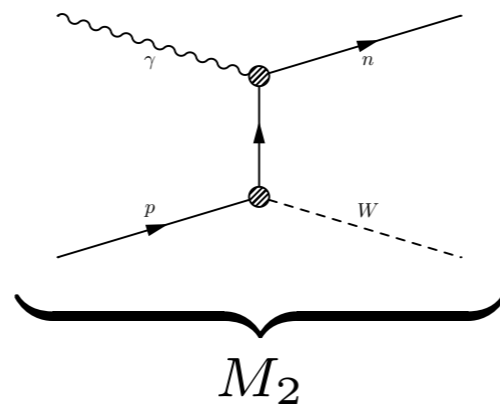
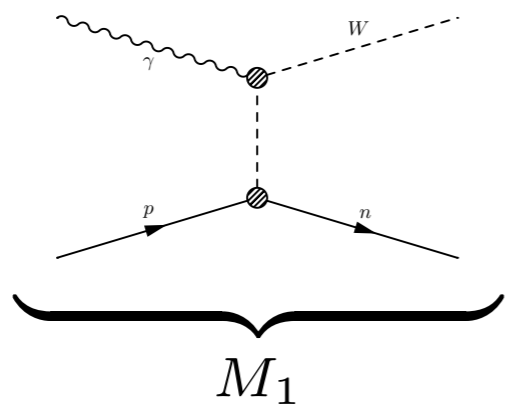
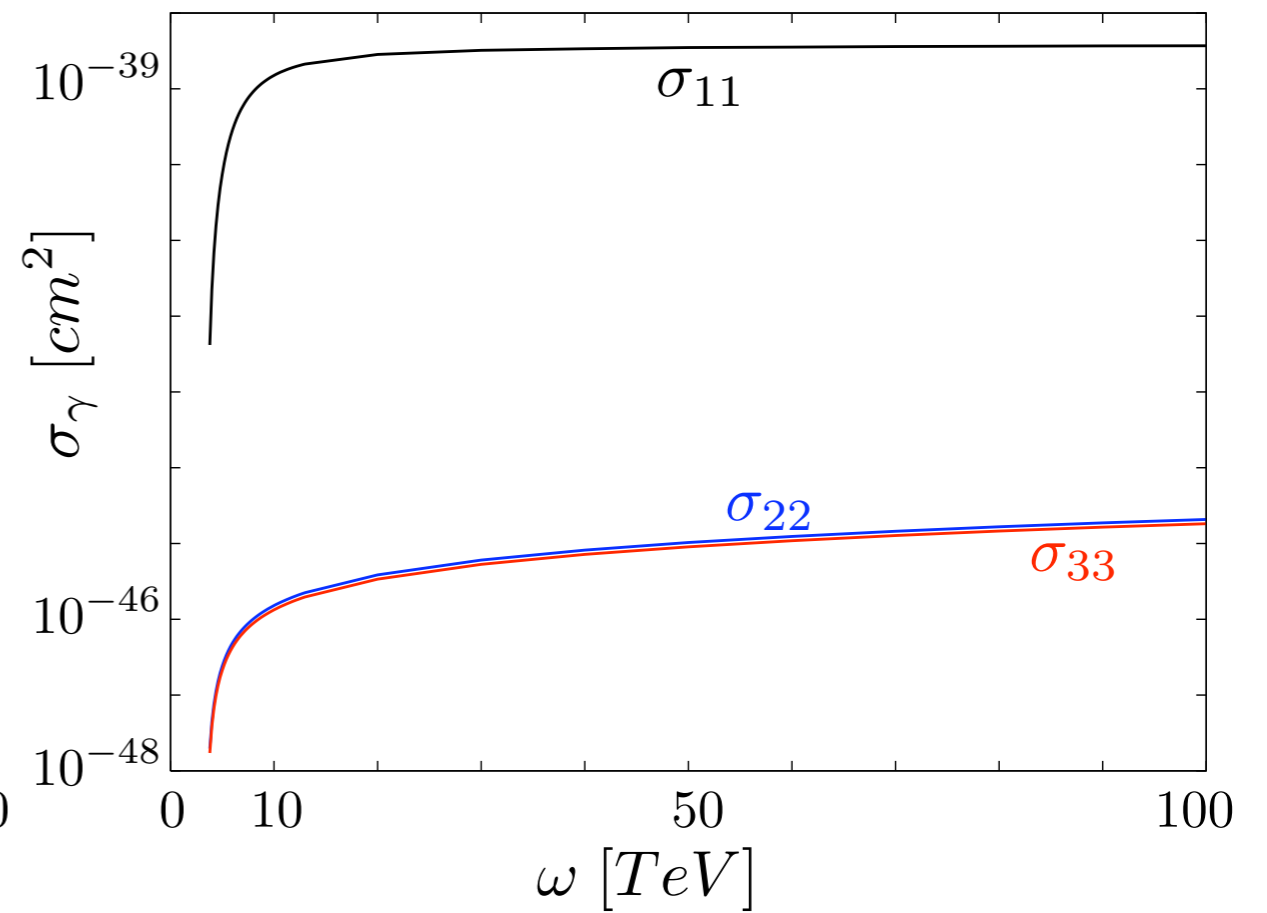


Results (Photoproduction)

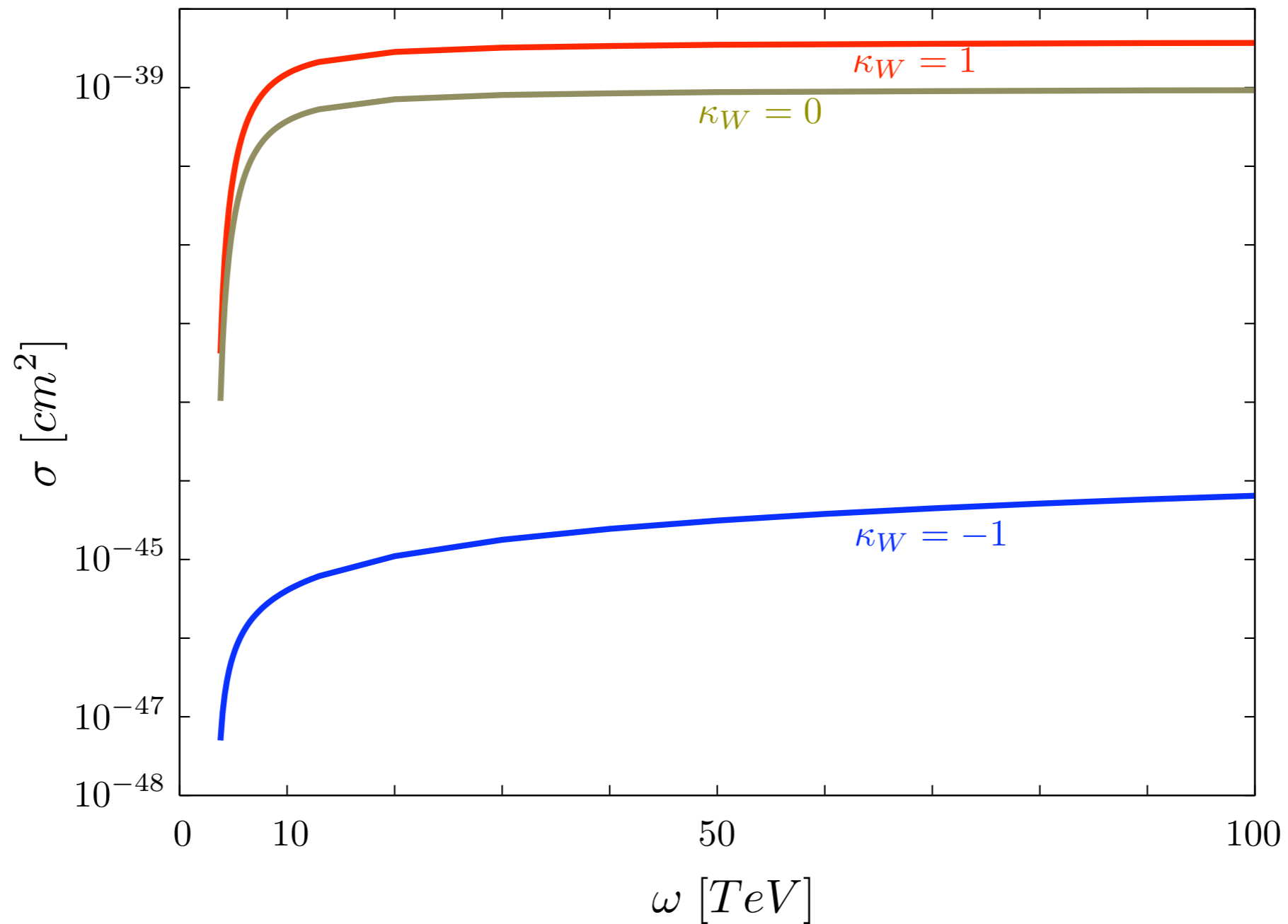
constant timelike form factors



dipole timelike form factors



Results (Photoproduction)



- total photoproduction cross section as function of photon energy for different values of κ_W

Equivalent Photon Approximation

- impact parameter larger than the extension of the nucleus/proton
- approximated equivalent photon spectrum for ion of radius R and mass M_A

$$f_{\gamma|A}(u) \approx \frac{2Z^2\alpha}{\pi} \ln \left(\frac{1}{uRM_A} \right) \quad \omega = u \cdot E_A$$

- approximate proton spectrum can be derived in the same way as ion spectrum, with R being the charge radius of the proton

$$f_{\gamma|p}(u) \approx \frac{\alpha}{\pi} \ln \left(\frac{0.71 \text{ GeV}^2}{u^2 M_p^2} \right) \quad \omega = u \cdot E_p$$

$$\sigma = \int \frac{du}{u} f_{\gamma|p/A}(u) \sigma_\gamma$$

Equivalent Photon Approximation (proton, elastic)

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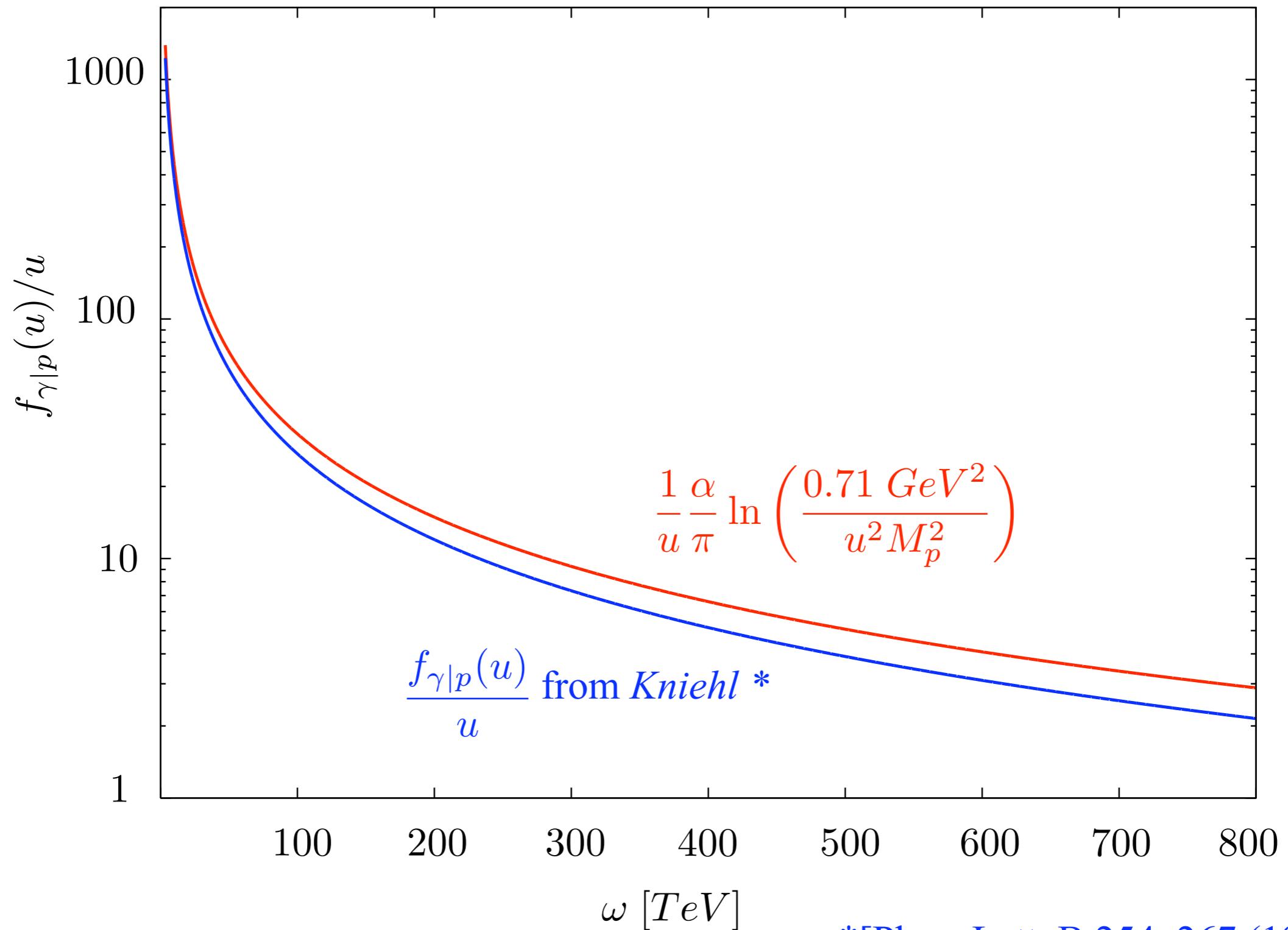
- parametrization from *Kniehl* [Phys. Lett. B 254, 267 (1991)]

$$f_{\gamma|p}(u) = \frac{\alpha}{2\pi} u \left[c_1 y \ln \left(1 + \frac{c_2}{z} \right) - (y + c_3) \ln \left(1 - \frac{1}{z} \right) \right. \\ \left. + \frac{c_4}{z-1} + \frac{c_5 y + c_6}{z} + \frac{c_7 y + c_8}{z^2} + \frac{c_9 y + c_{10}}{z^3} \right]$$

$$y = \frac{1}{2} - \frac{2}{u} + \frac{2}{u^2} \quad z = 1 + \frac{M_p^2}{0.71 \text{ GeV}^2} \frac{u^2}{1-u}$$

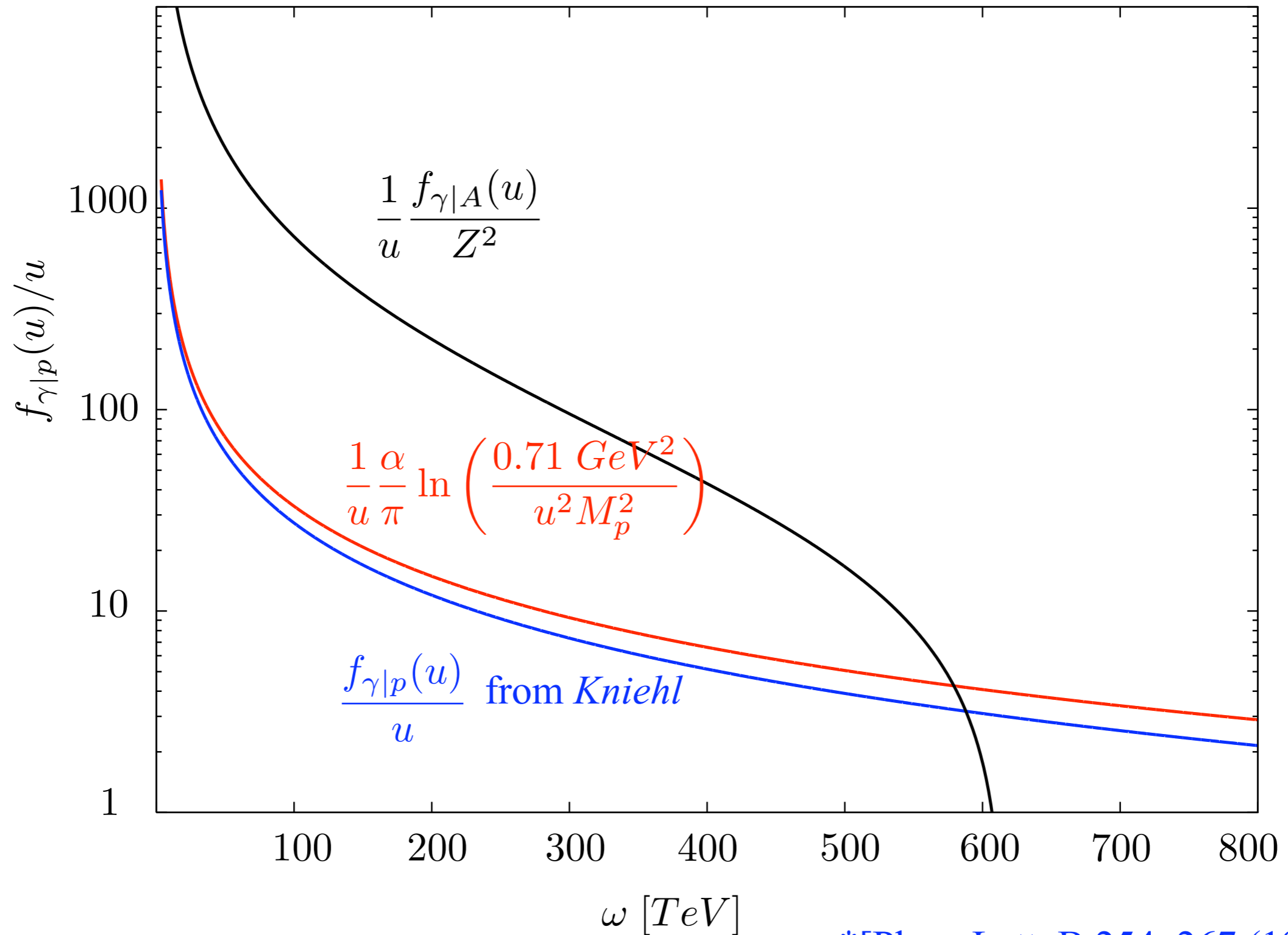
$$\sigma = \int \frac{du}{u} f_{\gamma|p}(u) \sigma_\gamma \quad \omega = u \cdot E_p$$

Equivalent Photon Approximation (proton, elastic)



*[Phys. Lett. B 254, 267 (1991)]

Equivalent Photon Approximation (elastic)



*[Phys. Lett. B 254, 267 (1991)]

Results (p-p and p-A collisions)

- total cross sections:
 - ▶ p-p collisions: $\sim 5 \cdot 10^{-40} \text{ cm}^2$
 - ▶ Pb-p collisions: $\sim 9 \cdot 10^{-37} \text{ cm}^2$

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- luminosities:

- ▶ p-p collisions: $L_{pp} \approx 10^{29} \dots 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

- ▶ Pb-p collisions: $L_{pA} \approx 10^{29} \dots \text{ cm}^{-2} \text{ s}^{-1}$

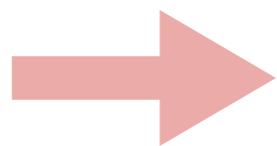
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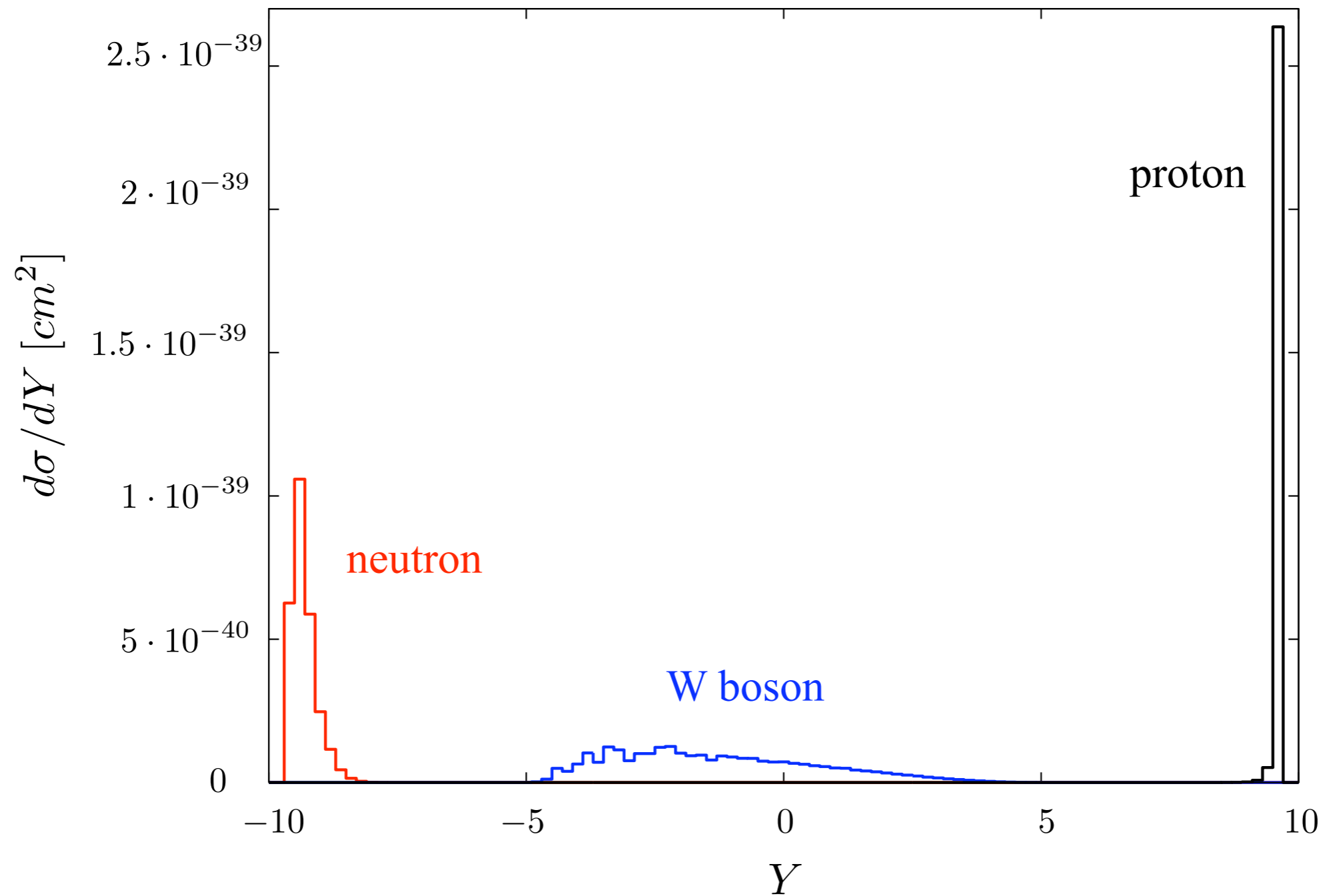
- ▶ p-p collisions: $L_{pp} \approx 10^{29} \dots 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- ▶ Pb-p collisions: $L_{pA} \approx 10^{29} \dots \text{ cm}^{-2} \text{ s}^{-1}$



$5 \cdot 10^{-4} \dots 50 \text{ events}/10^7 \text{ s}$ p-p collisions

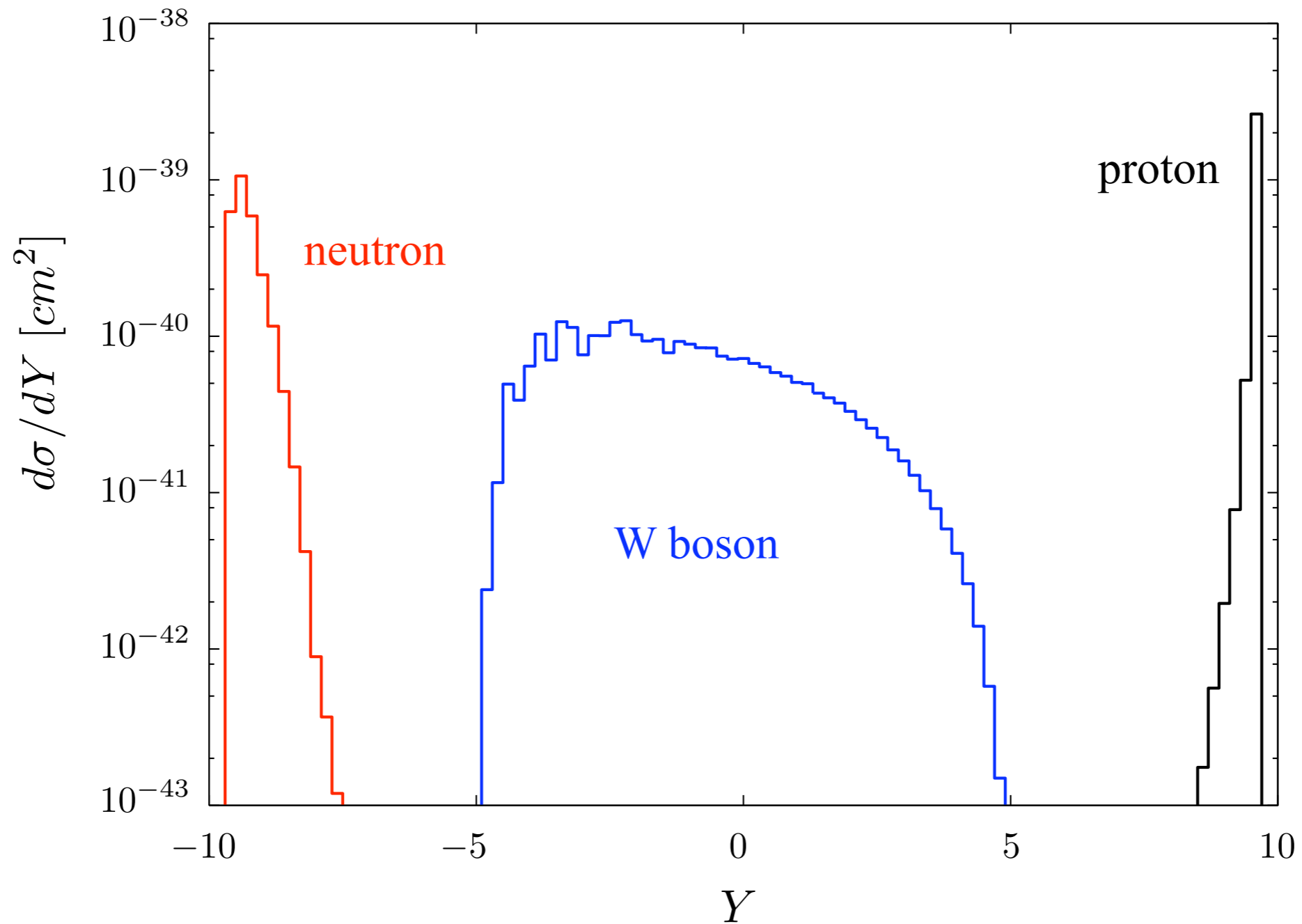
at least $0.1 \text{ events}/10^6 \text{ s}$ Pb-p collisions

Results (p-p collision)



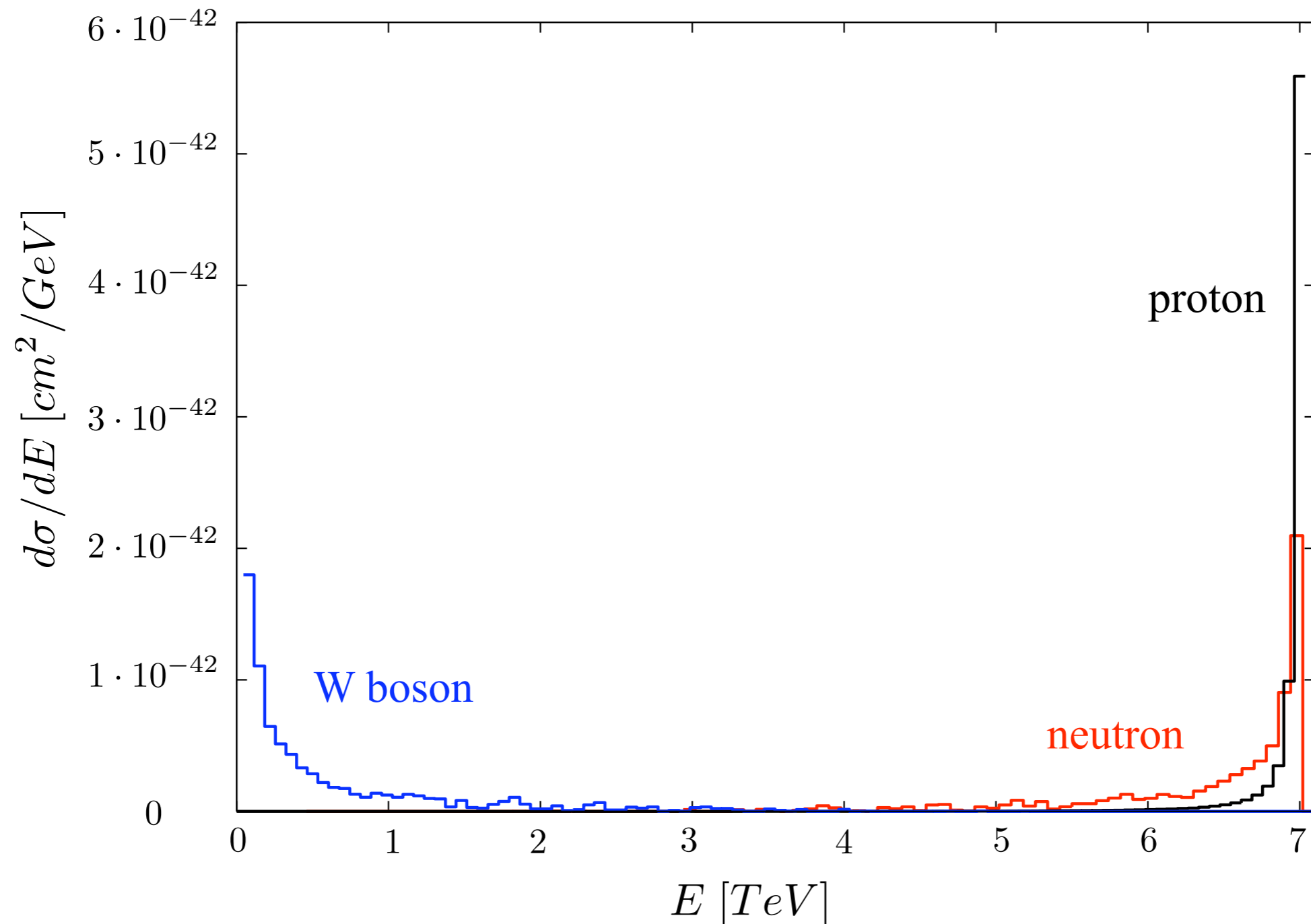
- differential cross section as function of rapidity in p-p collision

Results (p-p collision)



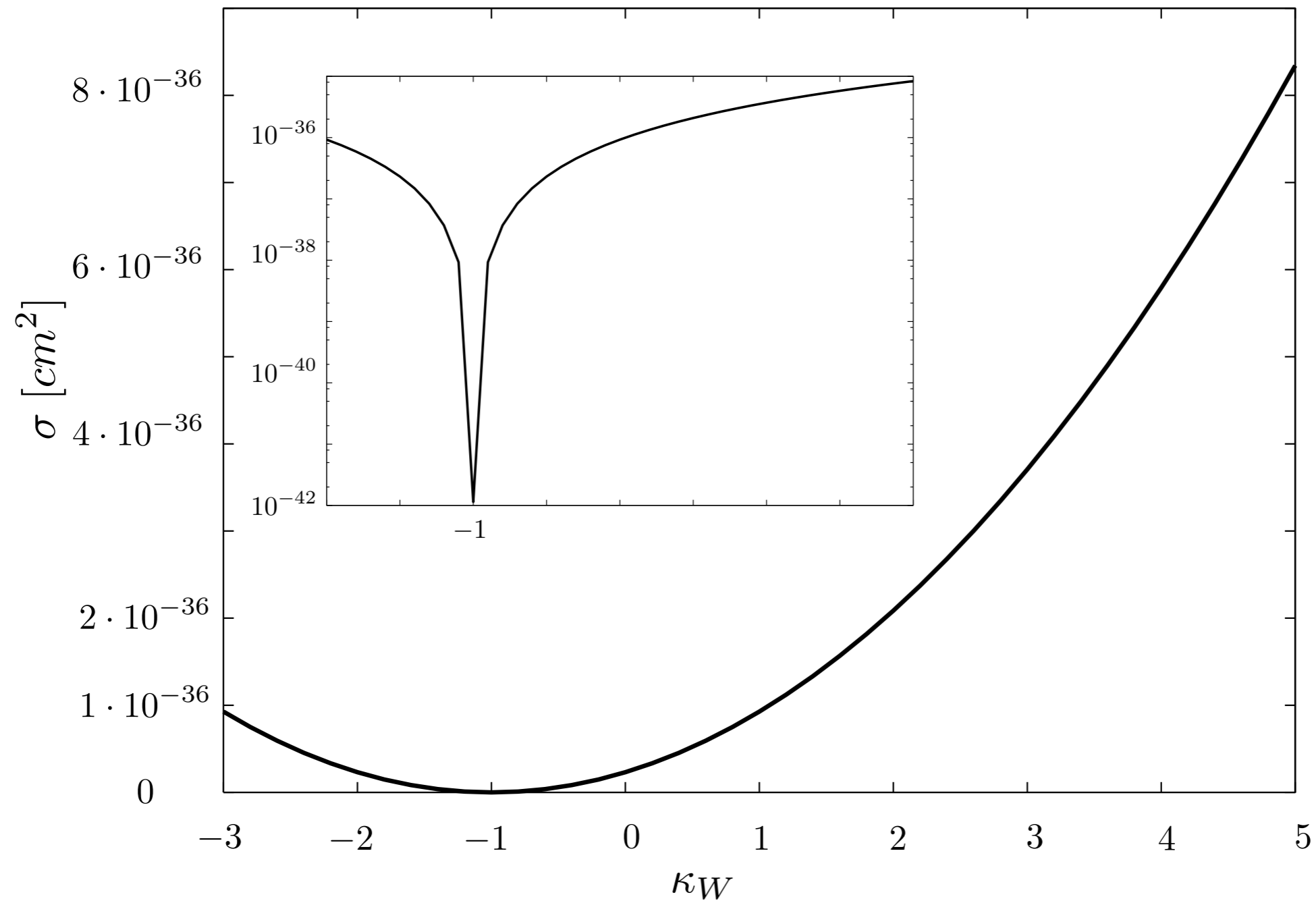
- differential cross section as function of rapidity in p-p collision

Results (p-p collision)



- differential cross section as function of energies in p-p collision, incident particles at 7 TeV

Results (p-A collision)



- total cross section as function of κ_W in Pb-p collision

Equivalent Photon Approximation (proton, inelastic)

- for large momentum transfers Q^2 of the photon, the proton should be regarded as a collection of partons, which radiate as pointlike particles
- for simplicity, we neglect the dependence of the parton distribution functions on Q^2

$$\sigma = \int dx \int du \sum_{q_i} e_i^2 f_{q_i|p}(x, Q_{av}^2) f_{\gamma|q_i}(u) \sigma_\gamma$$

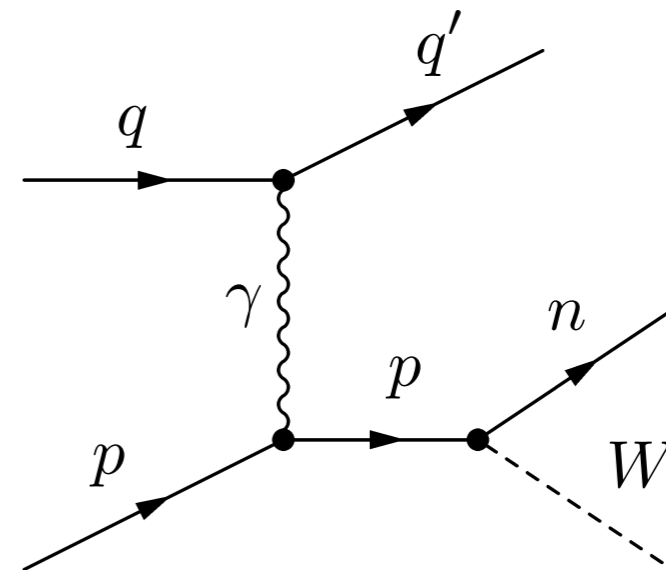
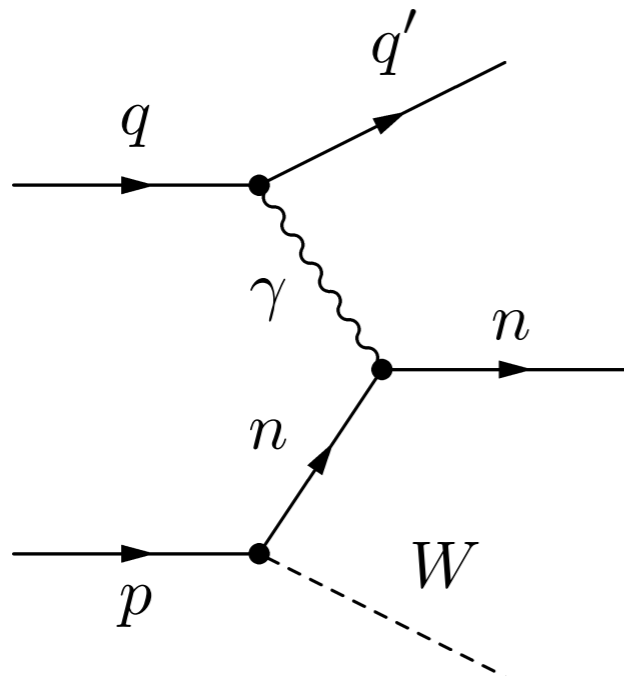
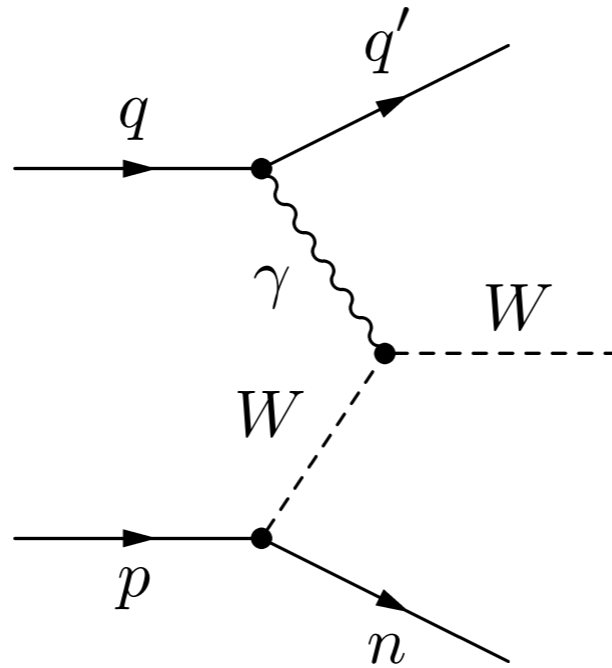
$$f_{\gamma|q} = \frac{\alpha}{2\pi} \frac{1 + (1-u)^2}{u} \ln \left(\frac{Q_{max}^2}{Q_{min}^2} \right) \quad \omega = x \cdot u \cdot E_p$$

$$Q_{min}^2 = 1 \text{ GeV}^2$$

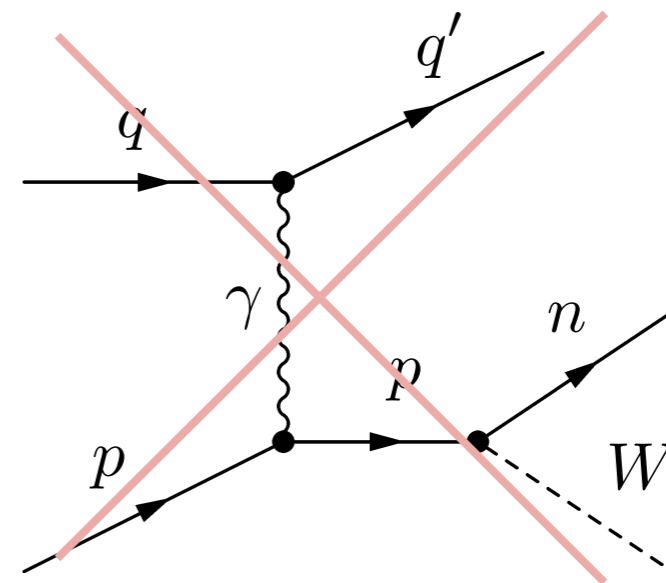
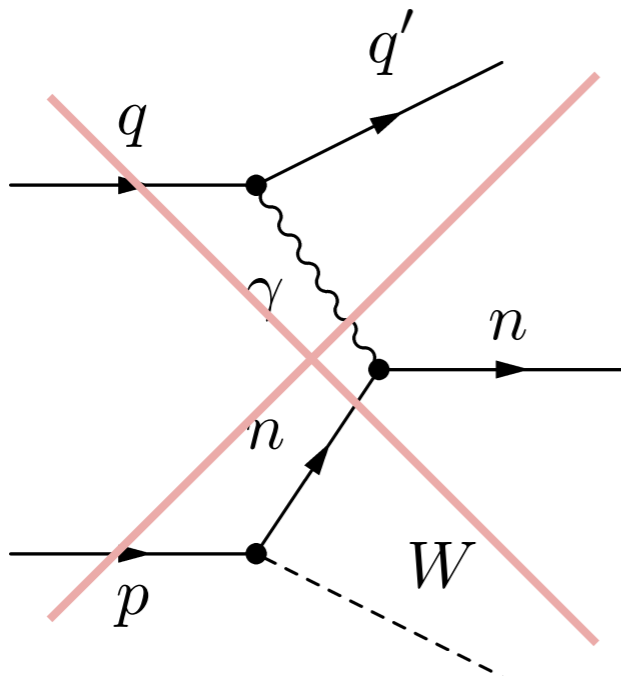
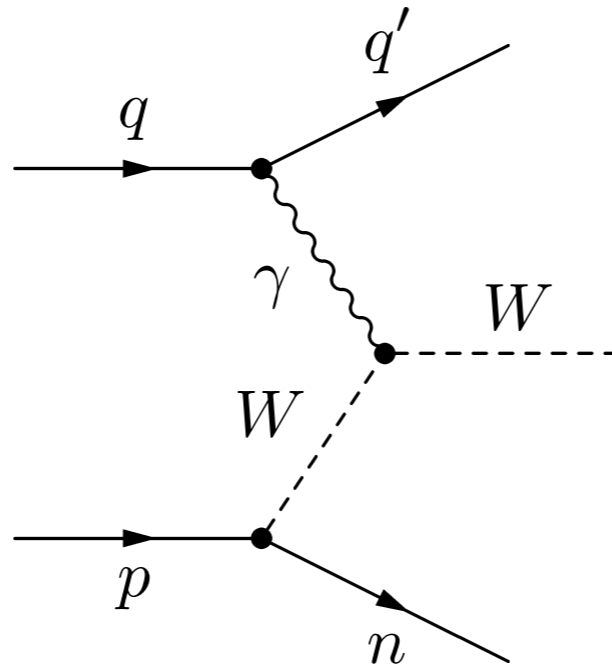
$$Q_{max}^2 = M_W^2$$

Equivalent Photon Approximation

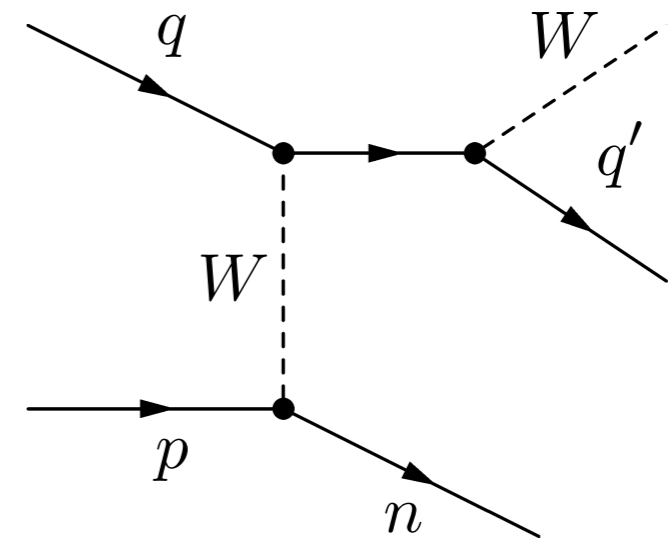
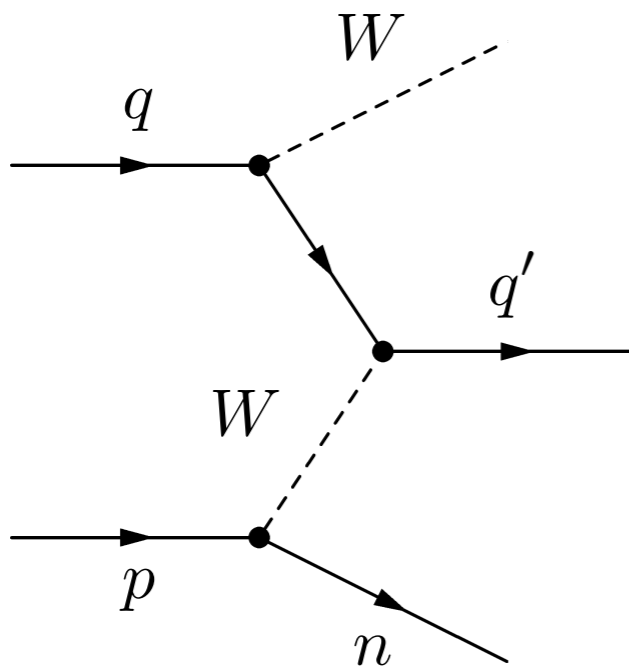
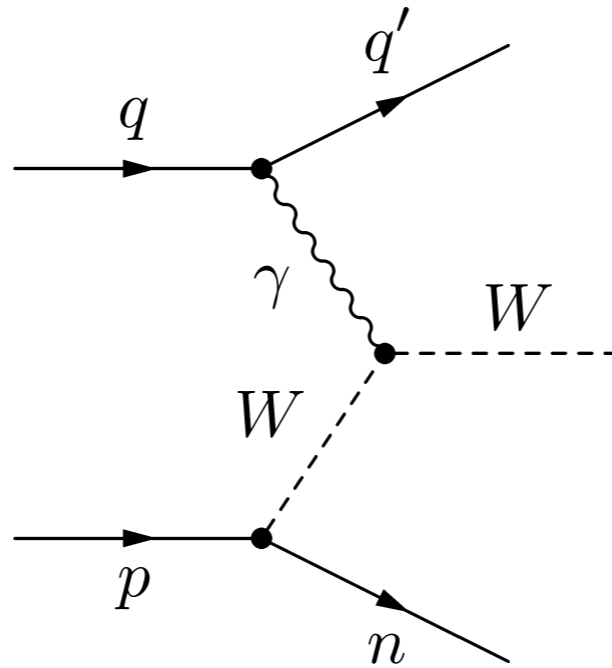
(proton, inelastic)



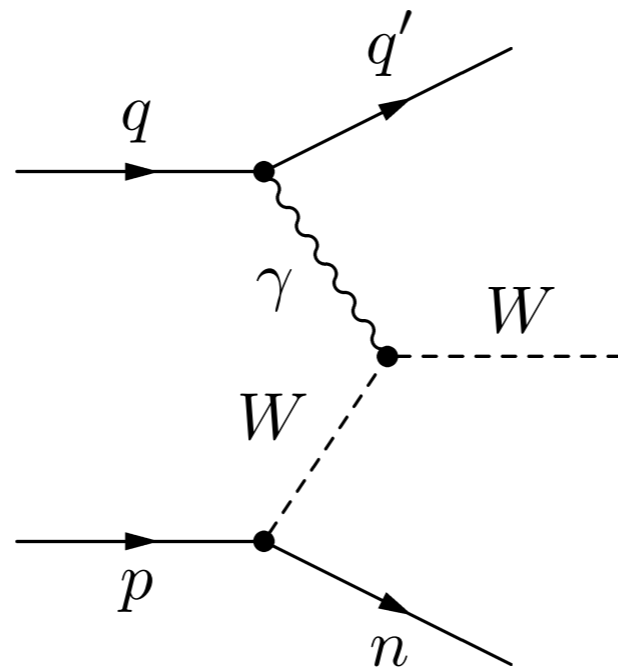
Equivalent Photon Approximation (proton, inelastic)



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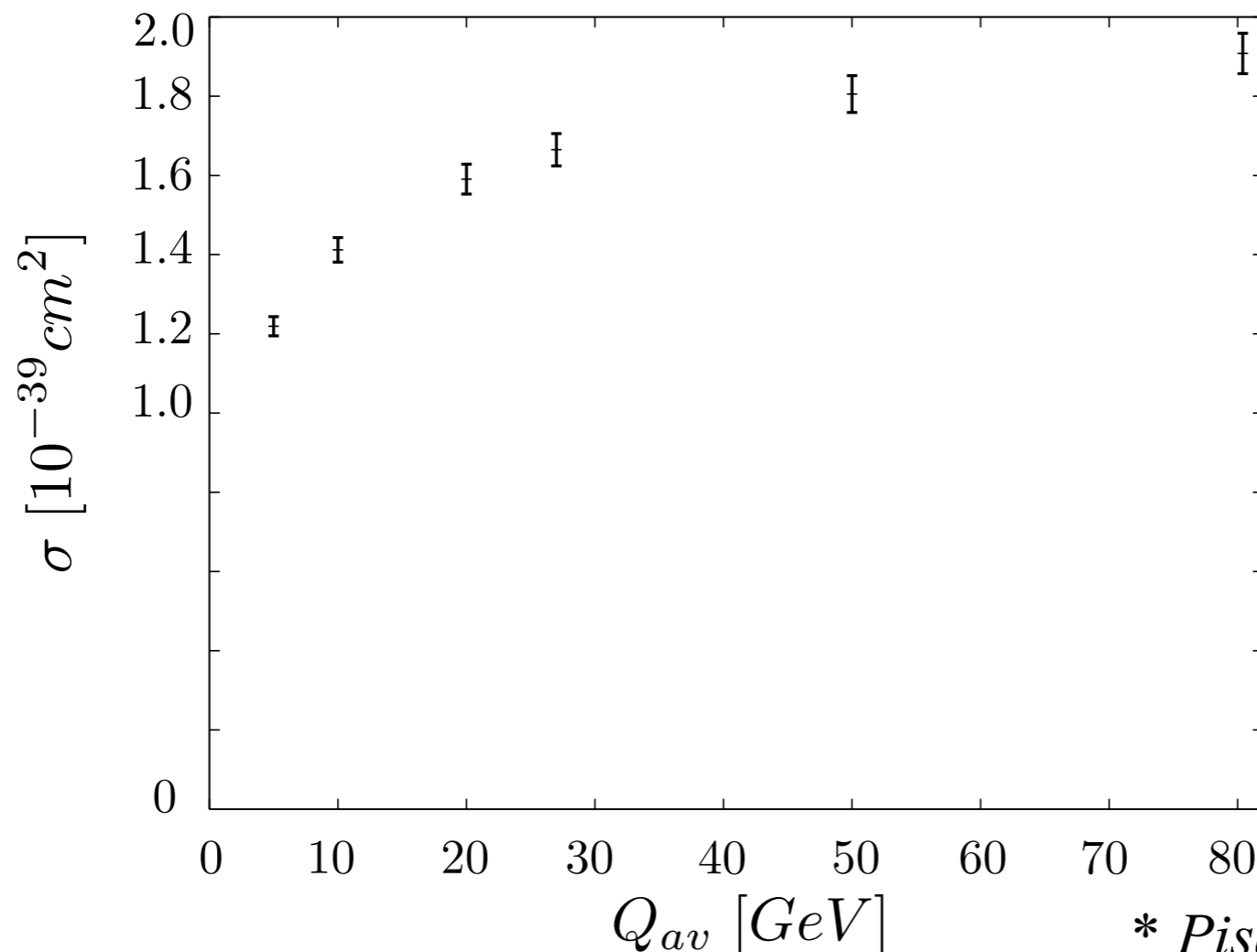
Results (p-p collisions)

- total cross sections:

- ▶ p-p collisions, elastic: $\sim 5 \cdot 10^{-40} \text{ cm}^2$

- ▶ p-p collisions, inelastic: $\sim 1.7 \cdot 10^{-39} \text{ cm}^2$

$$Q_{av}^{2*} = \frac{Q_{max}^2 - Q_{min}^2}{\log Q_{max}^2 - \log Q_{min}^2}$$

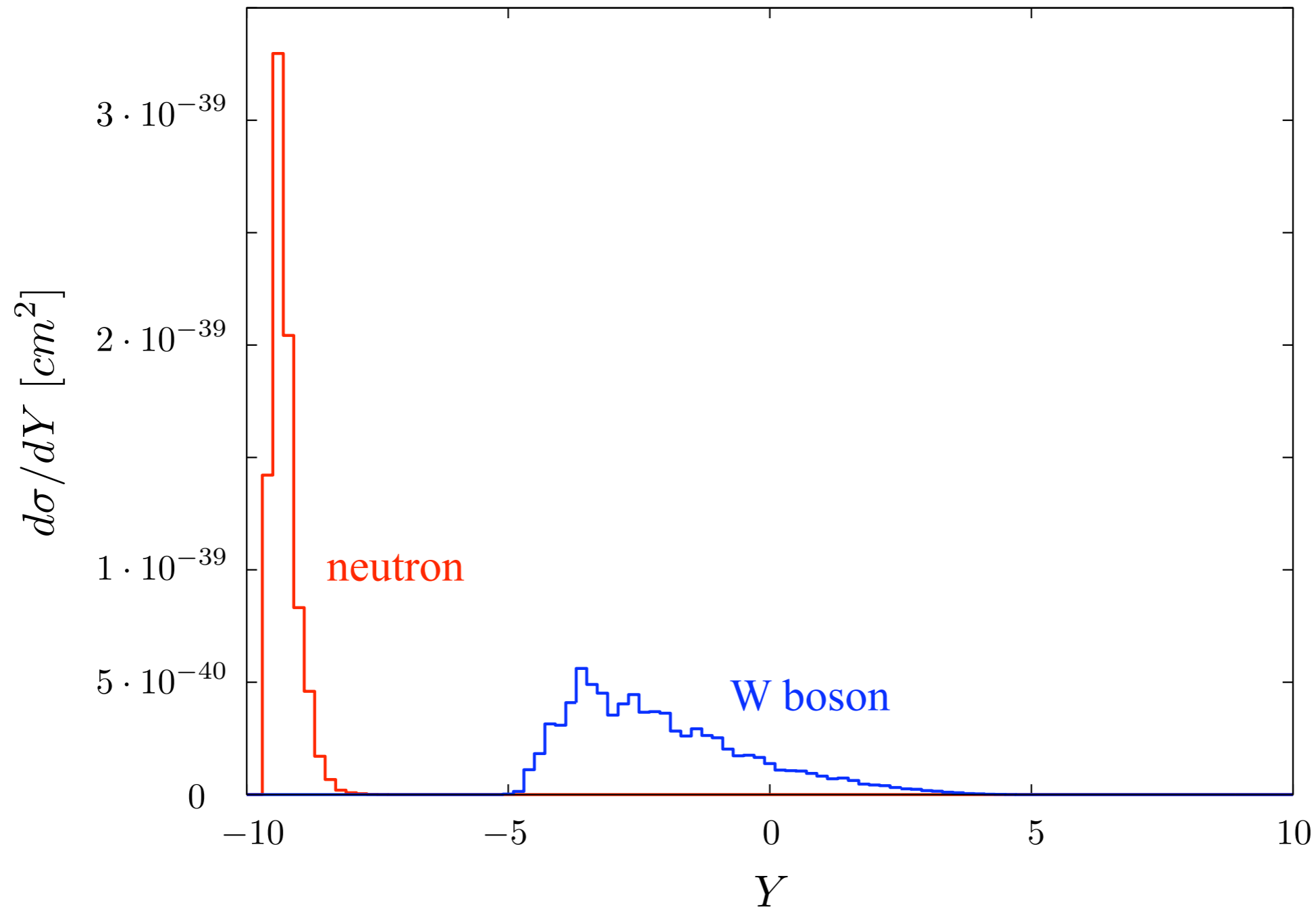


$$(1.2 \dots 1.9) \cdot 10^{-39} \text{ cm}^2$$

$$1 \text{ GeV} \leq Q_{av} \leq M_W$$

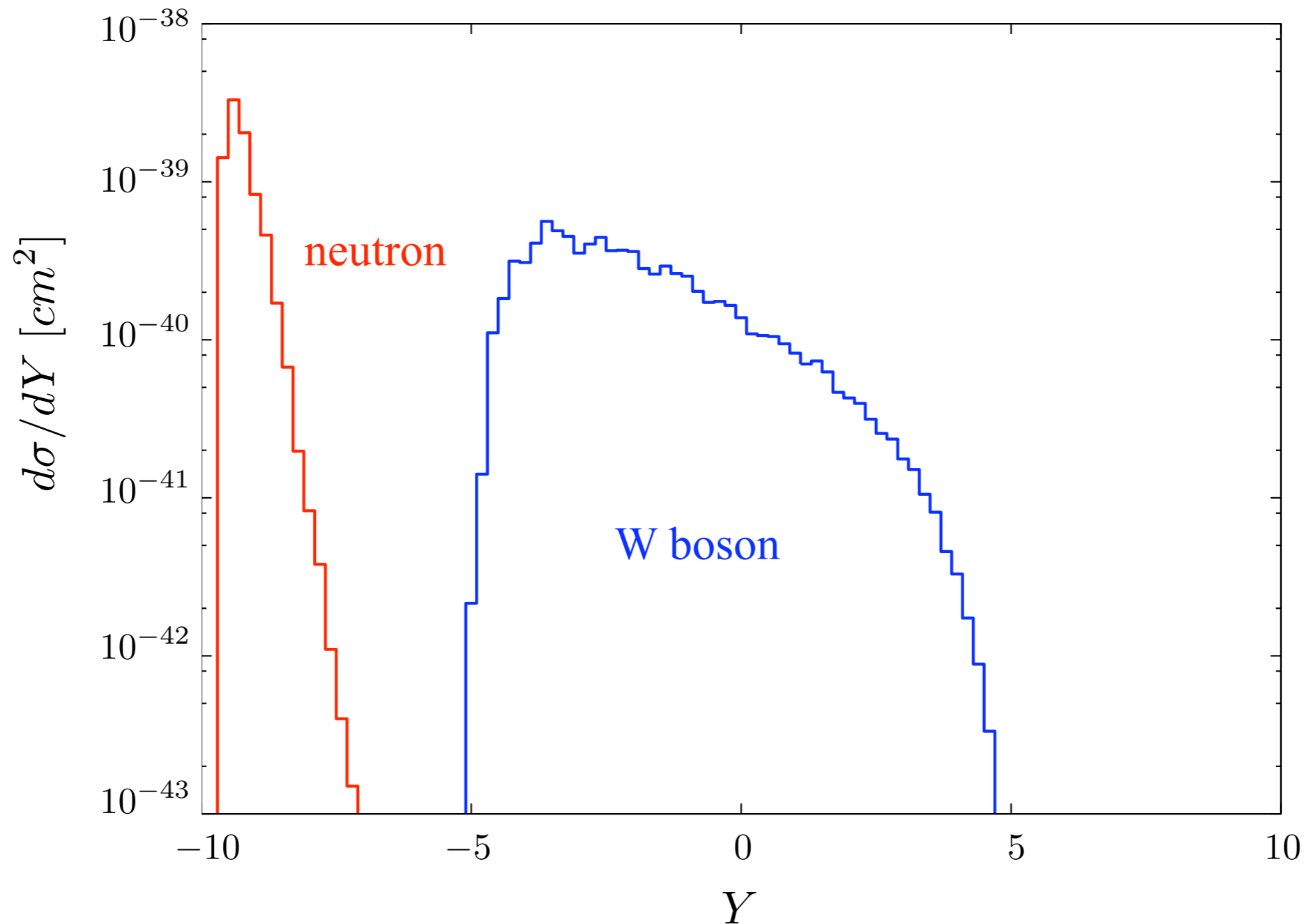
* *Pisano* [Eur. Phys. J. C 38, 79 (2004)]

Results (p-p collisions)



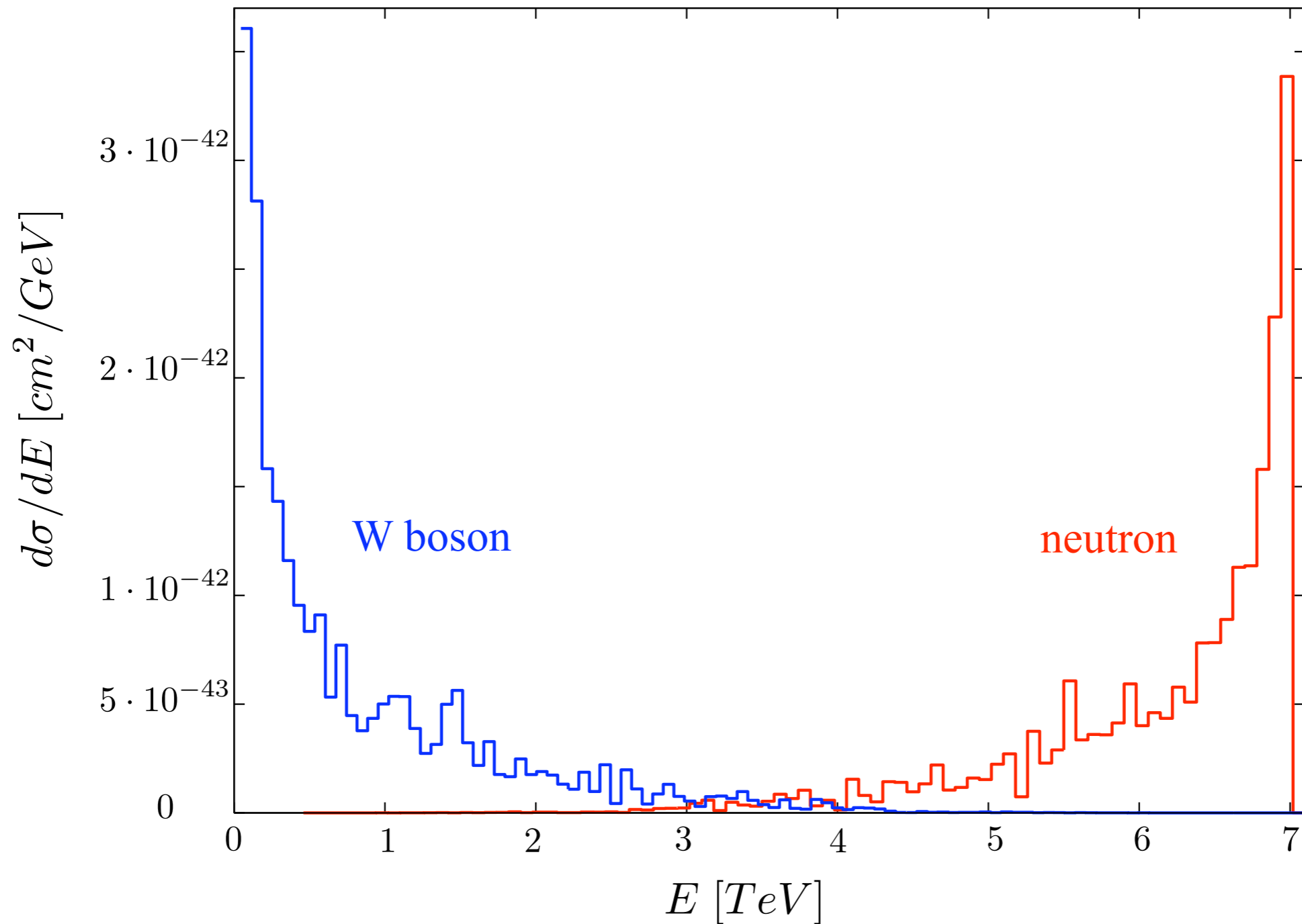
- differential cross section as function of rapidity in p-p collision (inelastic EPA)

Results (p-p collisions)



- differential cross section as function of rapidity in p-p collision (inelastic EPA)

Results (p-p collision)



- differential cross section as function of energies in p-p collision, incident particles at 7 TeV

Conclusion

- we give an estimate of the total cross section for exclusive single W boson production in p-p and p-Pb collisions
- in p-p collisions two possibilities: elastic and inelastic EPA
- for elastic EPA & for timelike form factors which fall off:
 - ▶ total cross section is sensitive to the anomalous magnetic moment of the W boson
 - ▶ differential and total cross sections do not depend on the choice of the form factors in the timelike region

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- in p-p collisions two possibilities: elastic and inelastic EPA
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Open Questions

- feasibility of measuring this process
- Is it possible to distinguish single W boson production from other processes?
- inelastic EPA: two further diagrams exist, their contribution needs to be checked

