

Beam-beam effects in High Energy LHC

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Thanks to O. Dominguez and F. Zimmermann

HE-LHC

- $E = 16.5 \text{ TeV}$
- Radiation damping time is 2h in the transverse emittance unit.

Updated parameter list for LHC energy upgrade at 33 TeV centre-of-mass energy

	nominal LHC	LHC energy upgrade
beam energy [TeV]	7	16.5
dipole field [T]	8.33	20
dipole coil aperture [mm]	56	40
beam half aperture [cm]	2.2 (x), 1.8 (y)	1.3
#bunches	2808	1404
bunch population [10^{11}]	1.15	1.29
initial transverse normalized emittance [μm]	3.75	3.75, 1.84
initial longitudinal emittance [eVs]	2.5	4.0
number of IPs contributing to tune shift	3	2
initial total beam-beam tune shift	0.01	0.01 (x & y)
maximum total beam-beam tune shift	0.01	0.01
RF voltage [MV]	16	32
rms bunch length [cm]	7.55	6.5
rms momentum spread [10^{-4}]	1.13	0.9
IP beta function [m]	0.55	1 (x), 0.43 (y)
initial rms IP spot size [μm]	16.7	14.6 (x), 6.3 (y)
full crossing angle [μrad]	285 ($9.5 \sigma_{x,y}$)	175 ($12 \sigma_{x0}$)
Piwinski angle	0.65	0.39
geometric luminosity loss from crossing	0.84	0.93
stored beam energy [MJ]	362	478.5
SR power per ring [kW]	3.6	62.3
dipole SR heat load dW/ds [W/m/aperture]	0.21	3.64
energy loss per turn [keV]	6.7	207.1
critical photon energy	44	576
longitudinal SR emittance damping time [h]	12.9	0.98
horizontal SR emittance damping time [h]	25.8	1.97
initial longitudinal IBS emittance rise time [h]	61	64
initial horizontal IBS emittance rise time [h]	80	80
initial vertical IBS emittance rise time [h] ($\kappa=0.2$)	~400	~400
<i>note: IBS rise times > SR damping times</i>		
events per crossing	19	76
initial luminosity [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	1.0	2.0
peak luminosity [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	1.0	2.0
beam lifetime [h]	46	12.6
integrated luminosity over 10 h [fb^{-1}]	0.3	0.5

Introduction

- Interplay of the synchrotron radiation, intra-beam scattering and beam-beam interaction.
- Coherent beam-beam effects in HE-LHC.
- Incoherent emittance growth in HE-LHC.
- Effects of IP optics parameters in the coherent and incoherent phenomena.

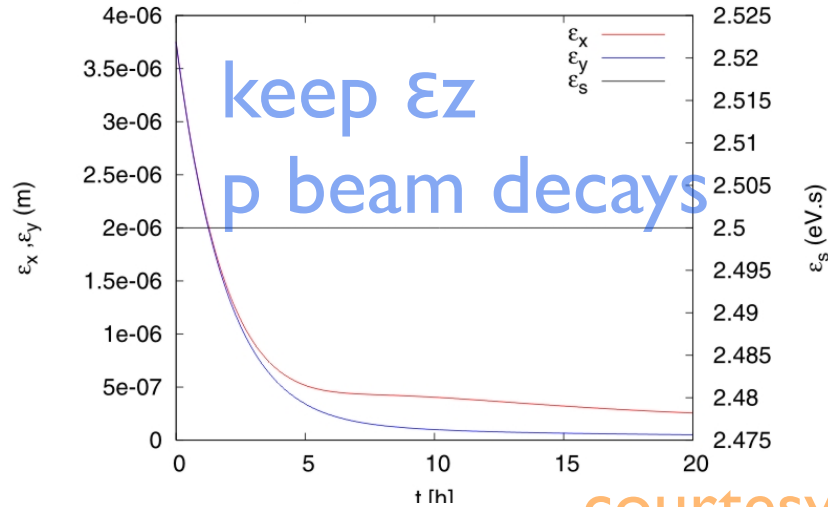
Beam size evolution with radiation damping and IBS

O. Dominguez and F. Zimmermann

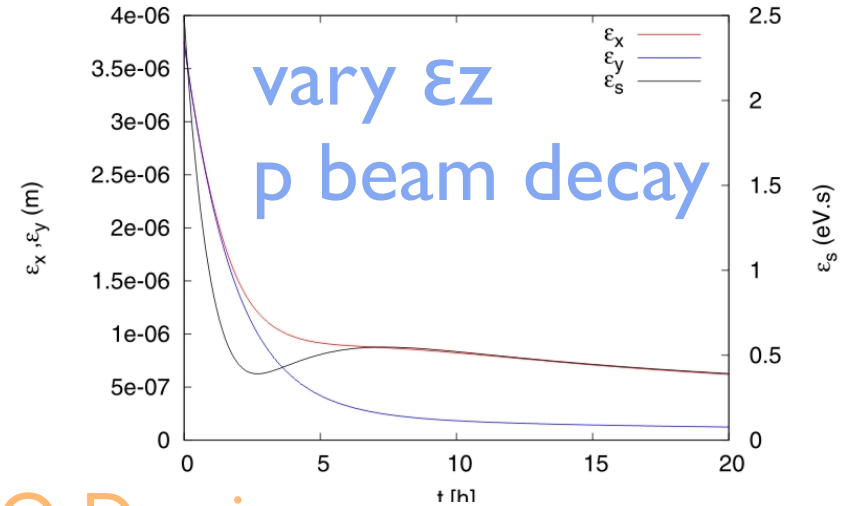
Initial

$\epsilon_x, \epsilon_y = 0.21, 0.10$ nm,
 $\xi_x, \xi_y = 0.0053, 0.0053$

ϵ_x, ϵ_y and ϵ_s vs time with burning off

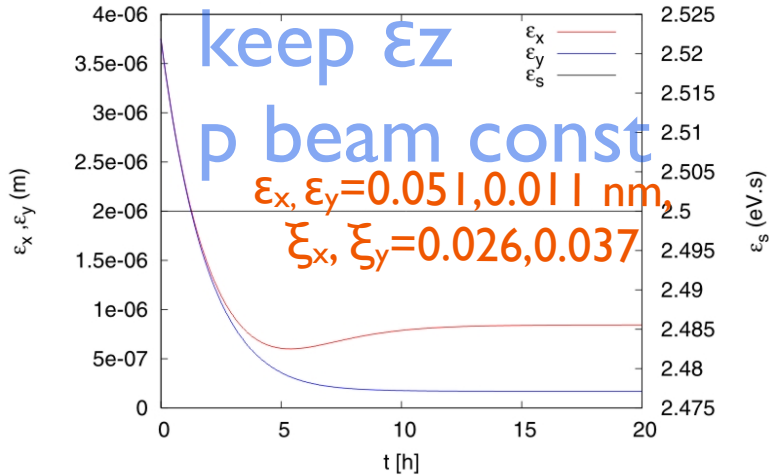


ϵ_x, ϵ_y and ϵ_s vs time with burning off

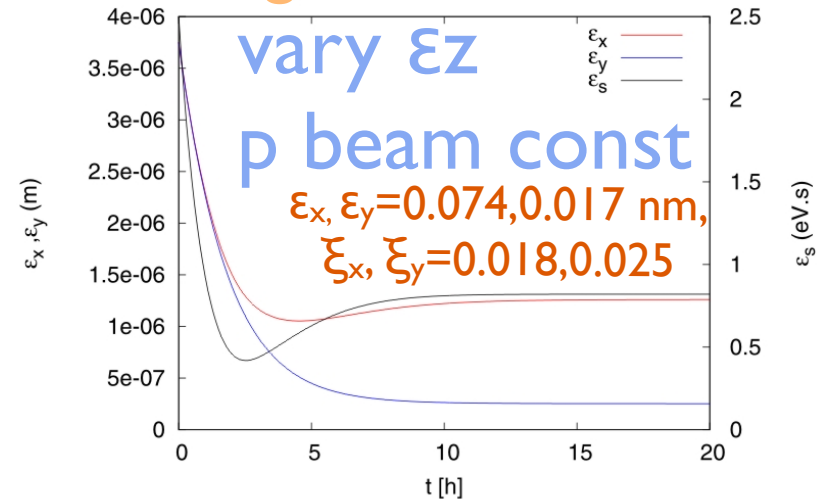


courtesy of O. Dominguez

ϵ_x, ϵ_y and ϵ_s vs time without burning off

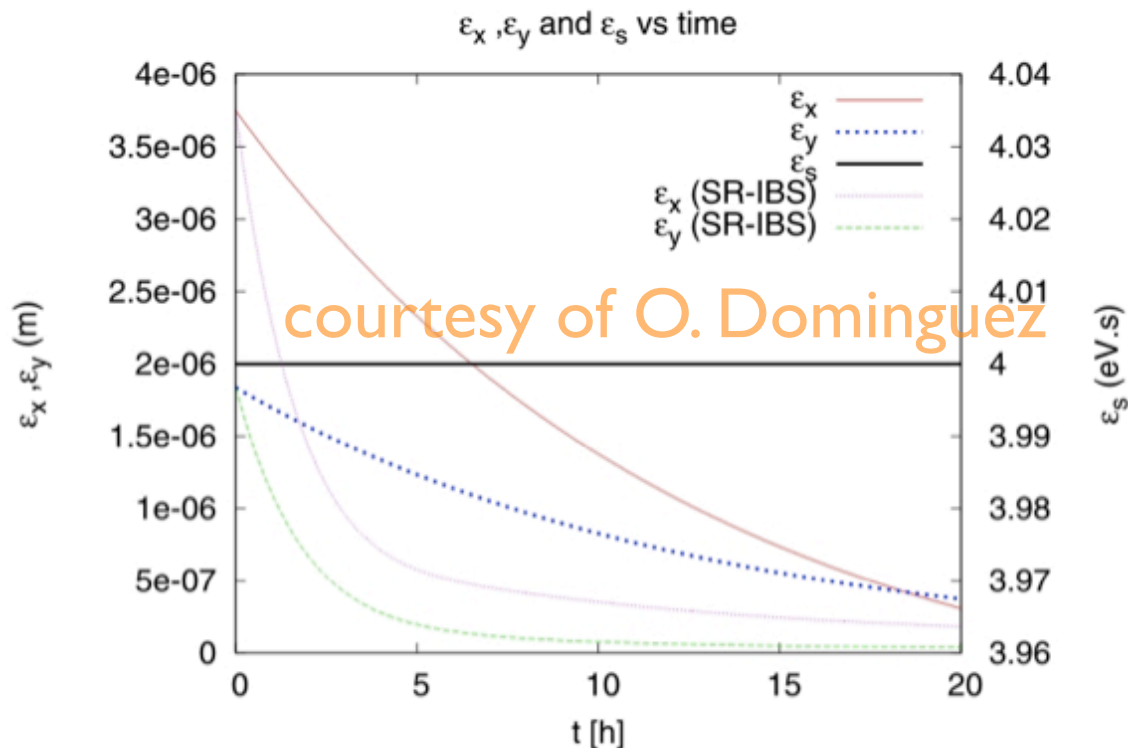


ϵ_x, ϵ_y and ϵ_s vs time without burning off



Control transverse excitation

- Emittance is controllable by applying external fluctuation (kicker) with keeping the beam-beam parameter. Conservative case.
- We study higher beam-beam parameter utilizing the emittance radiation damping.



keep $\epsilon_z, \xi_x, \xi_y,$
 ρ beam decays

courtesy of O. Dominguez

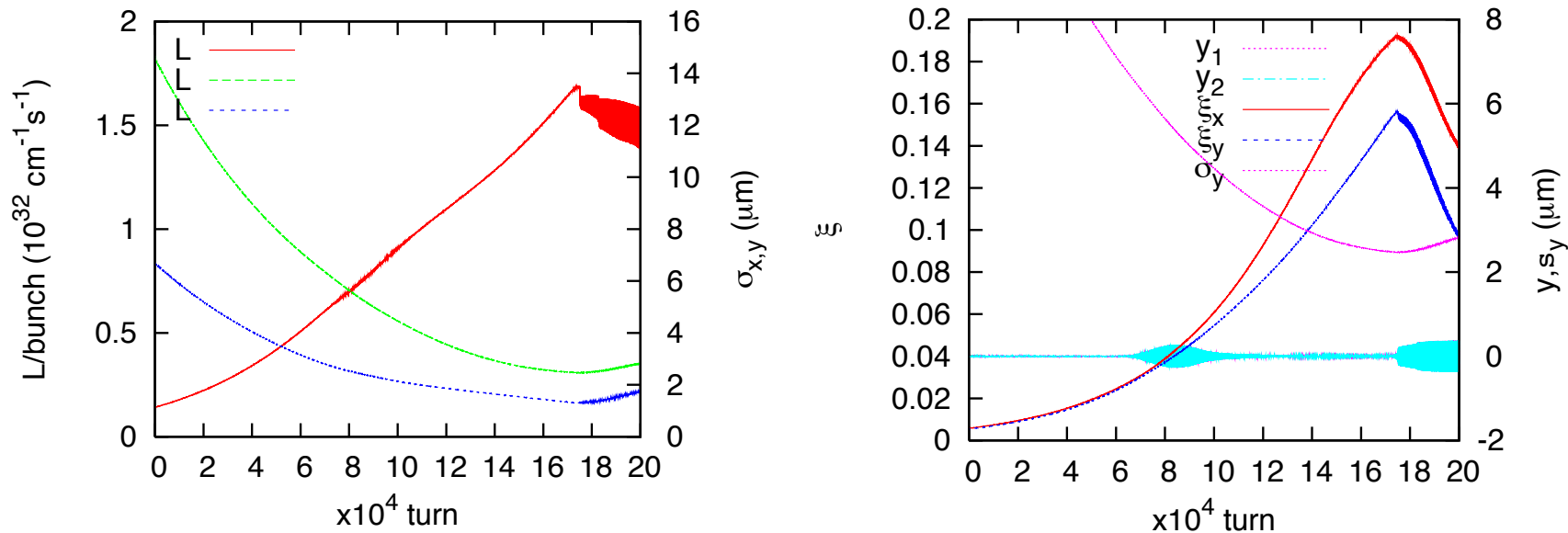
Coherent and incoherent effects of the beam-beam interaction

- Coherent effect is studied by the strong-strong simulation. Single IP and 2D sim.
- Incoherent effect has been studied by the weak-strong simulation. Two IP, 3D sim., crossing angle are taken into account.

Study of coherent effect in HE-LHC

- Strong-strong simulation with a code BBSS. **Single IP.Turn is regarded as collision occurrence for 2IP.**
- Considering simulation time, the radiation damping time has to be chosen faster than actual value.
- Simulations are performed with 2000 and 200 times faster damping time.
- We expect that extrapolation gives information.

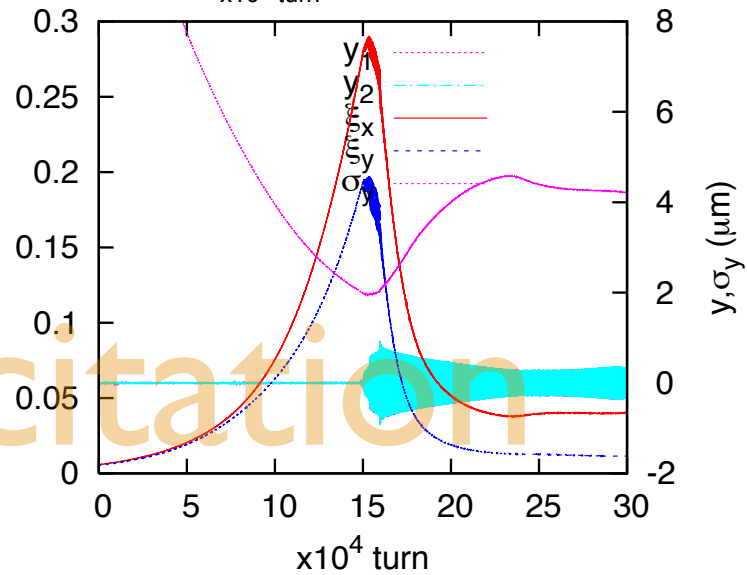
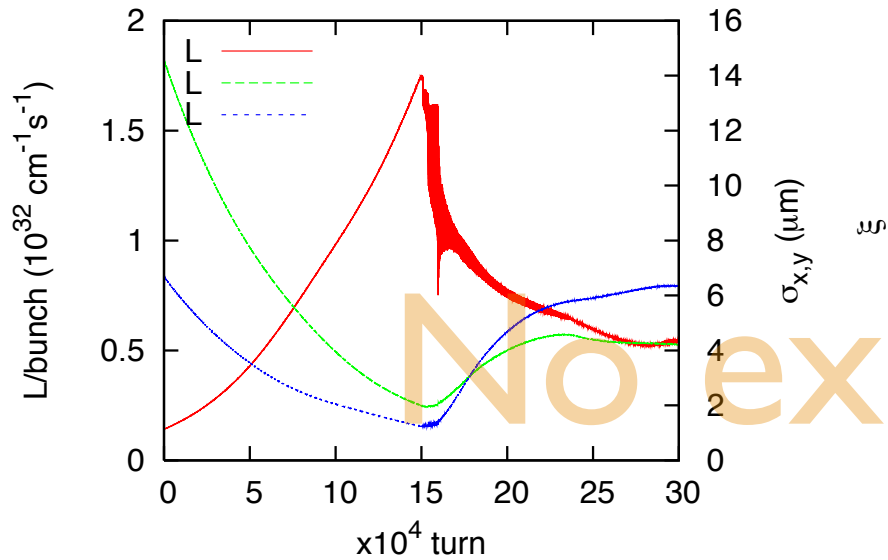
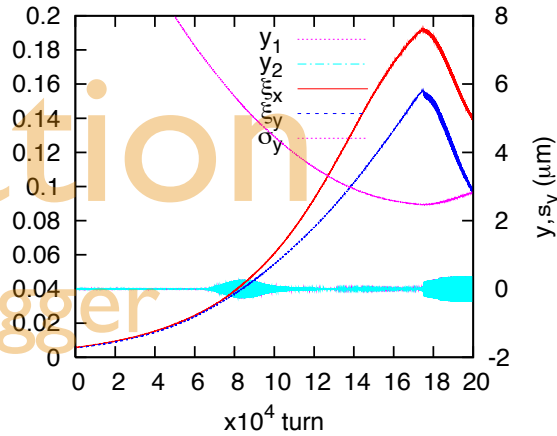
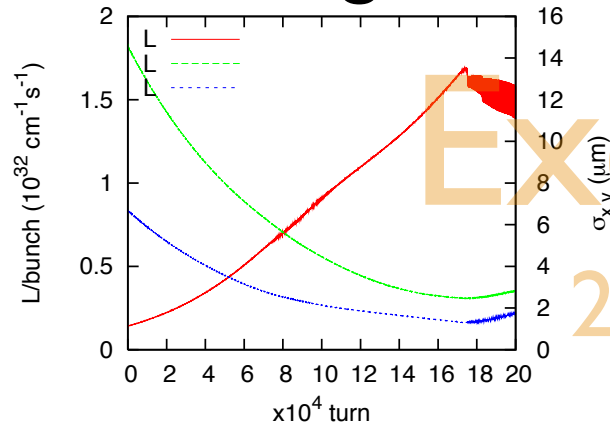
Assume 2000x faster damping time 200 times bigger excitation



- Dipole oscillation limit the luminosity. The beam-beam parameter is very high $\xi \sim 0.15$.
- The dipole oscillation was seen at $\xi > 0.05$ in a flat beam such as lepton colliders for $\sigma_z \ll \beta_y$.

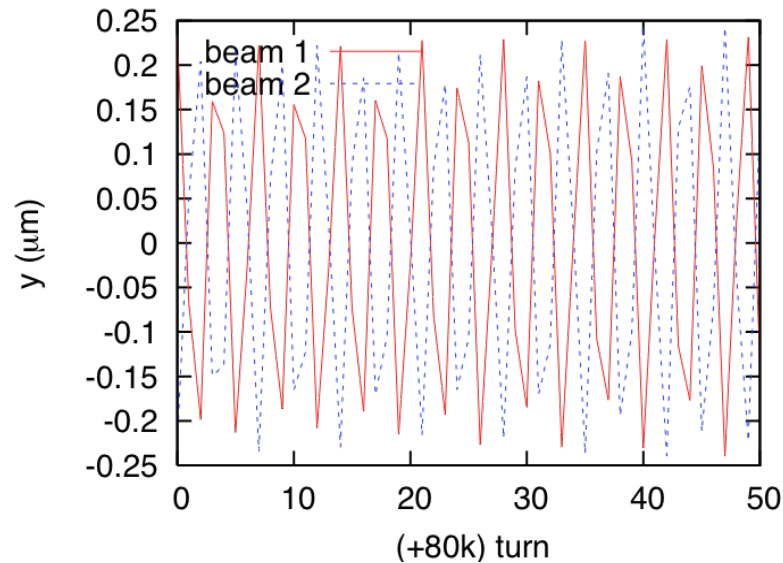
Excitation ON/OFF

- No big difference.



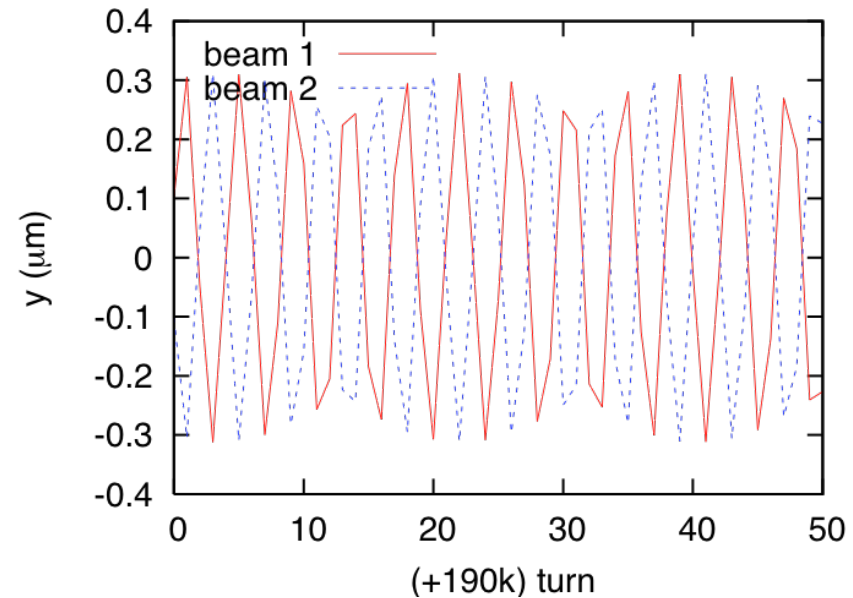
Dipole oscillation

- π mode



~ 80000 turn

$\times 2000 = 1.6 \times 10^8 = 4$ hour

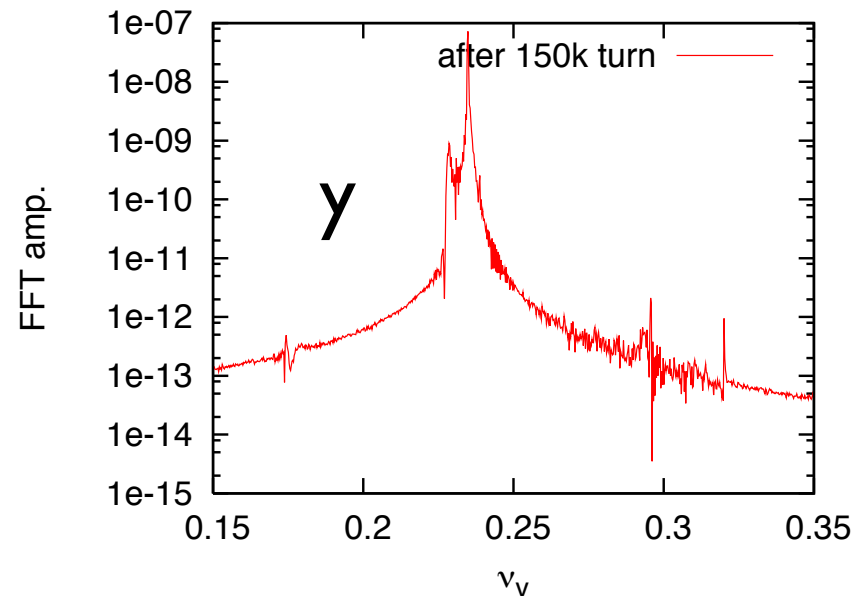
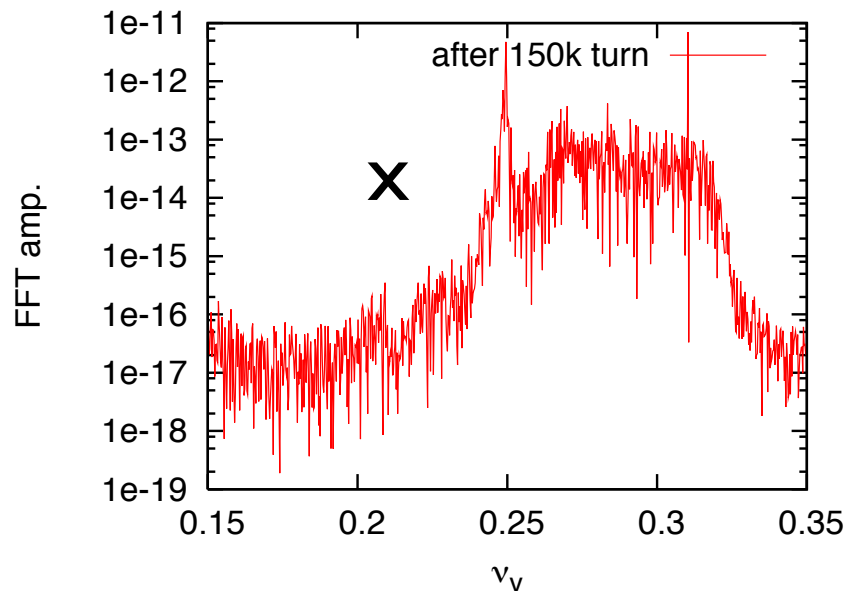
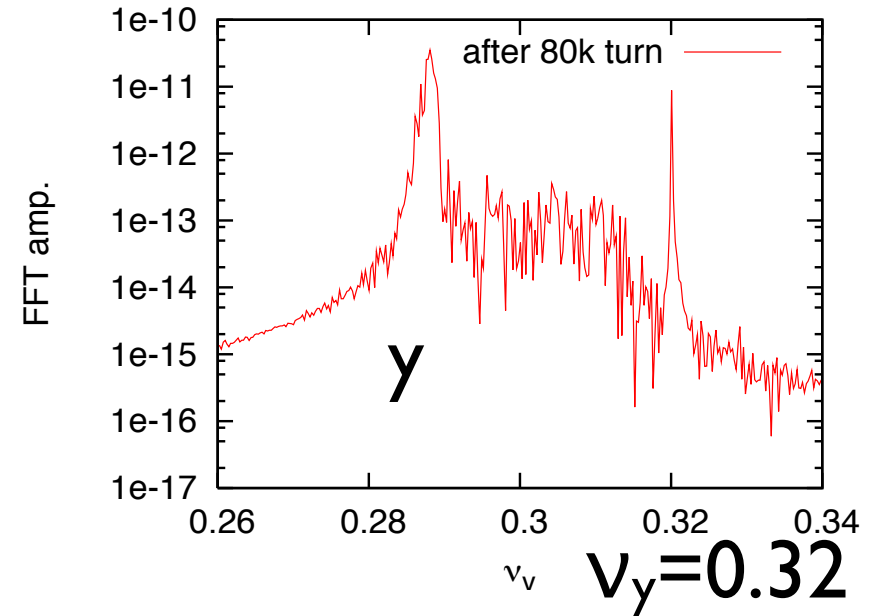
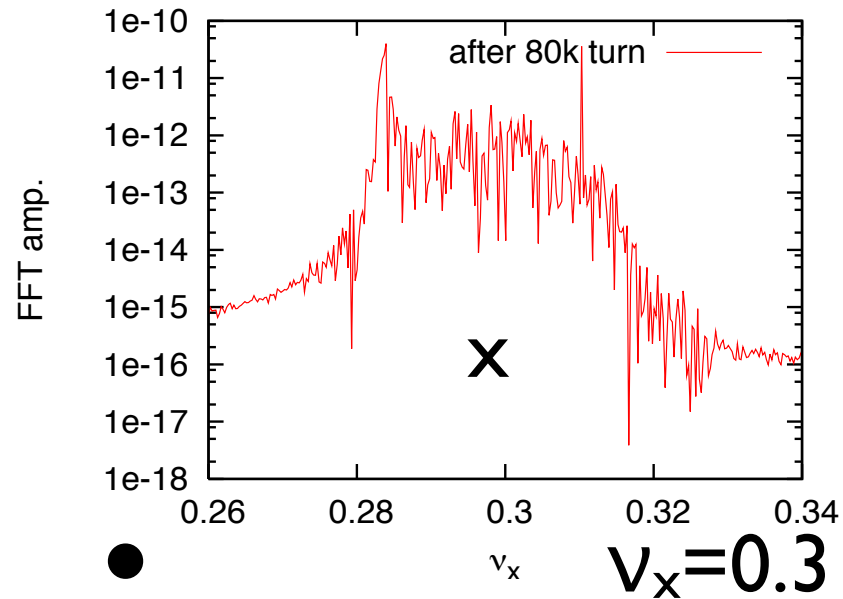


~ 190000 turn

$\times 2000 = 3.8 \times 10^8 = 9$ hour

Not realized by IBS

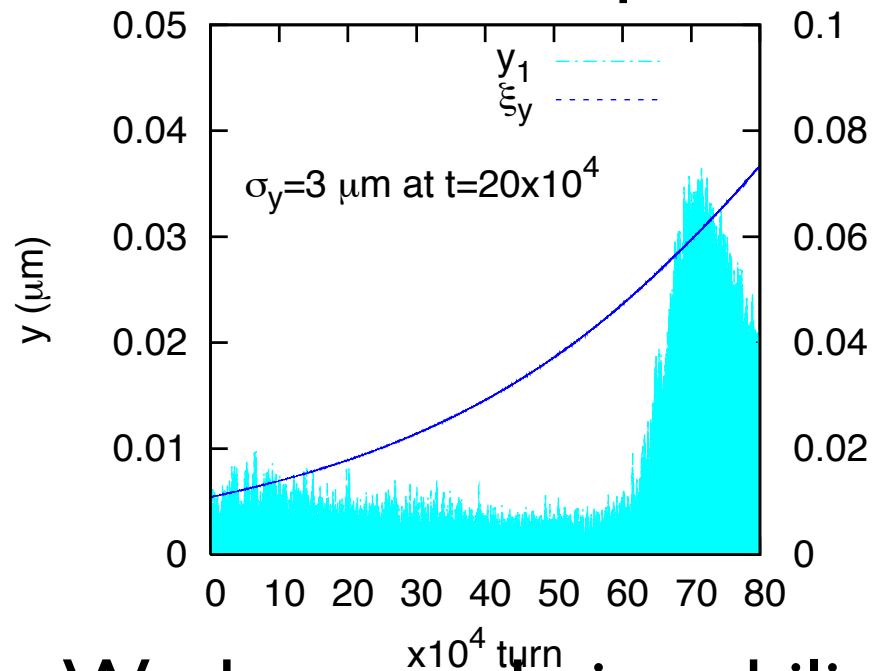
Fourier spectra for the instability



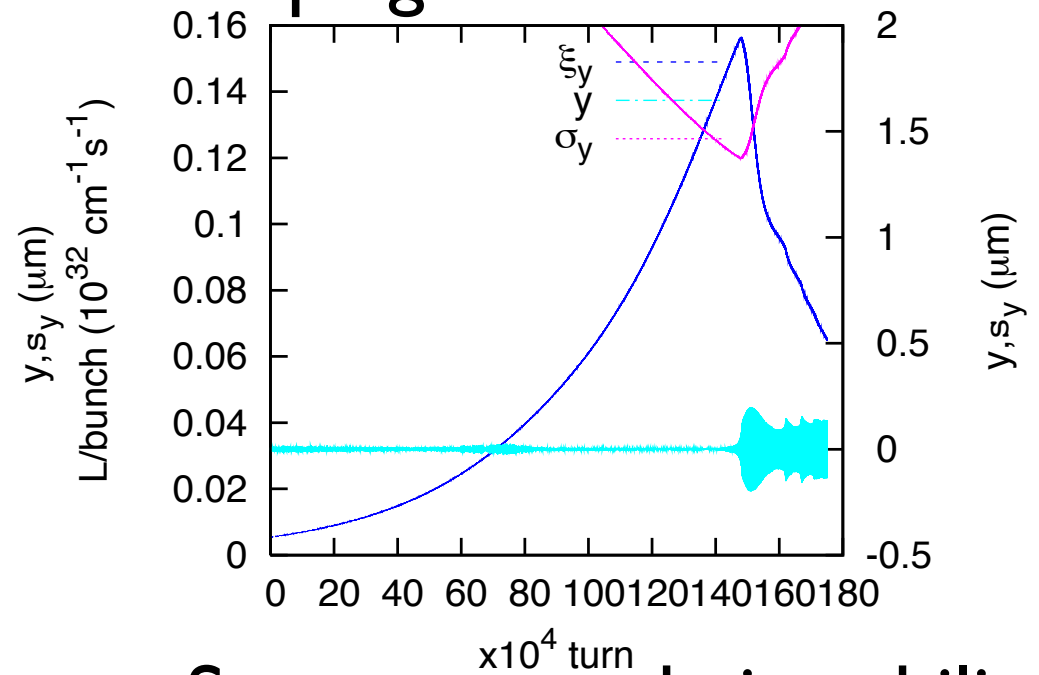
200x faster damping

20 times bigger excitation

- Excitation ON/OFF
- No remarkable difference. Coherent instability does not depend on the damping time.



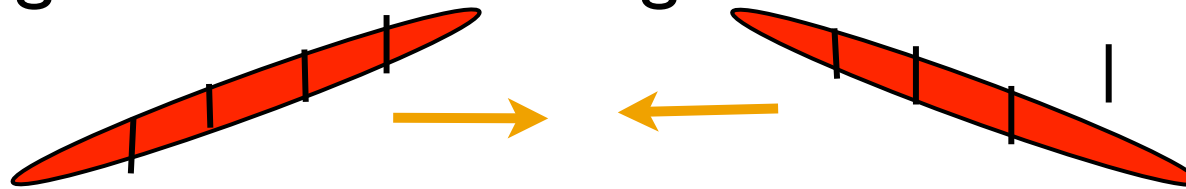
Weak π -mode instability
at 8×10^5 turn \sim 4 hour.



Strong π -mode instability
at 15×10^5 turn \sim 8 hour

Incoherent emittance growth

- 2 IP, phase difference π (pessimistic case)
- crossing angle, 5 slices along z.
- Turning off the radiation damping, luminosity degradation is investigated.

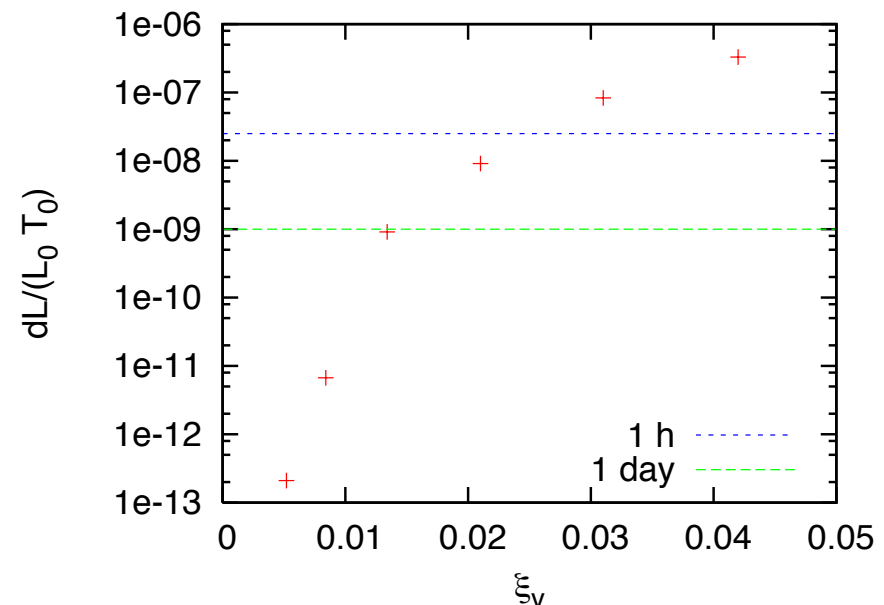
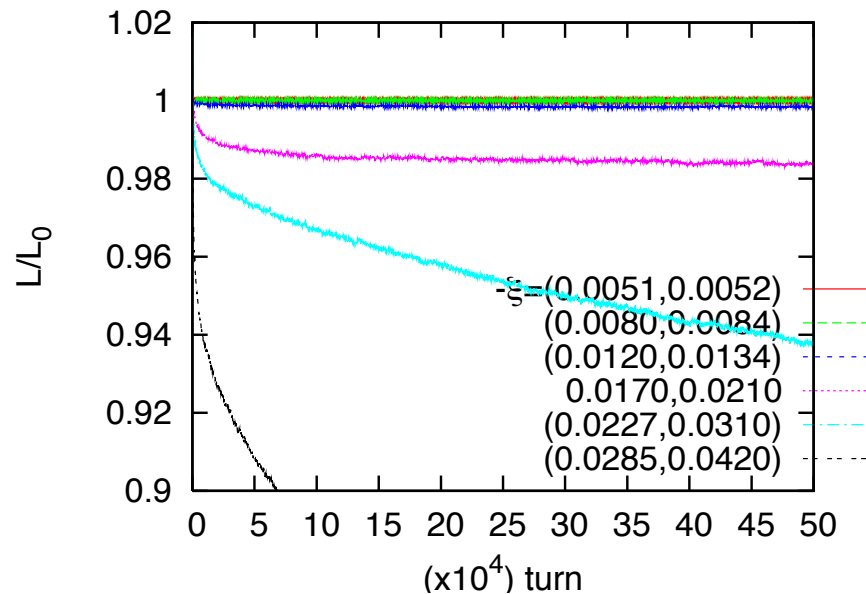


t (h)	ϵ_x (nm)	ϵ_y (nm)	ξ_x (/IP)	ξ_y (/IP)
0	0.21	0.1	0.0051	0.0052
1	0.13	0.062	0.0080	0.0084
2	0.076	0.037	0.012	0.013
3	0.046	0.022	0.017	0.021
4	0.027	0.014	0.023	0.031
5	0.016	0.0097	0.029	0.042

Study 6 cases, which are $t=0-5$ hour after the injection without IBS.

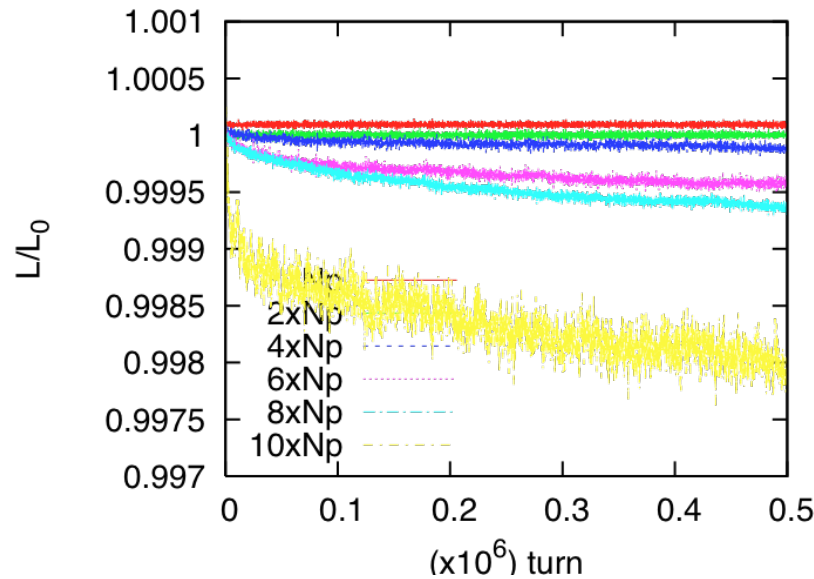
Luminosity degradation due to incoherent emittance growth

- The beam-beam parameter is 0.0134/IP for 1 day and 0.025/IP for 1 hour life time.
- This result can be also applied for the present LHC.

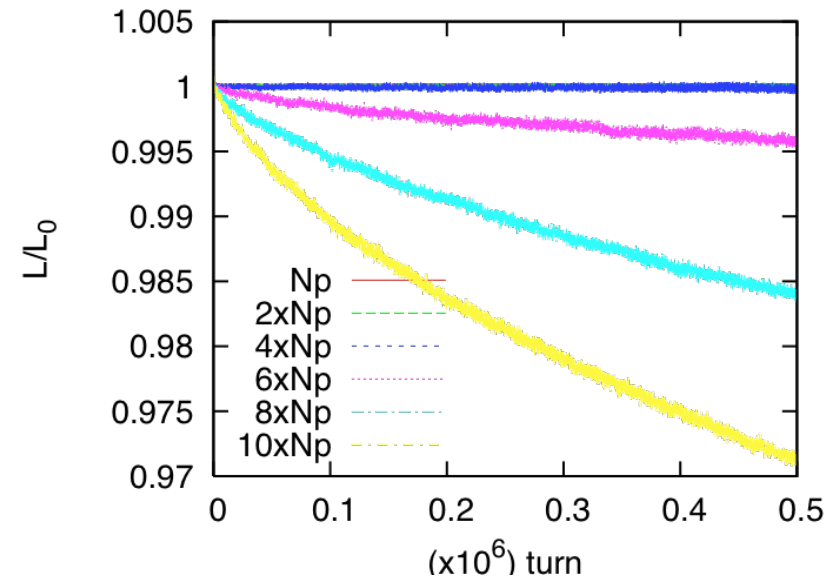


Effect of crossing angle

- 2 IP
- $\xi=0.0034/\text{IP}$ for N_p .



zero crossing angle



$\theta_H \sigma_z / \sigma_x = 0.6$

Luminosity degradation without crossing angle is very small.

Tentative result for HE-LHC

- A weak coherent instability arises at $\xi \sim 0.03$, but disappears for further turn evolutions.
- A strong coherent instability arises at $\xi \sim 0.15$ and degrades the luminosity. This condition is not realized by IBS.
- Incoherent growth is dominant compared to the coherent instability.
- Incoherent growth time is $0.0134/\text{IP}$ for 1 day and $0.025/\text{IP}$ for 1 hour life time. It is possible to be $\sim 0.02/\text{IP}$ considering radiation damping time 2 h.
- Since the geometrical beam-beam parameter limit of IBS is $0.03/\text{IP}$, the beam size should be controlled to be $0.02/\text{IP}$ using an external noise.

x-y coupling for HE-LHC beam beam performance

- x-y coupling affects the beam-beam performance essentially in KEKB. It is effect of beam-beam dynamics, but not that of geometric.
- How x-y coupling affects LHC performance.
- Model with 2000 time faster damping to study coherent instability.
- Weak strong simulation to study incoherent growth.

x-y coupling at IP

- Parametrization of x-y coupling

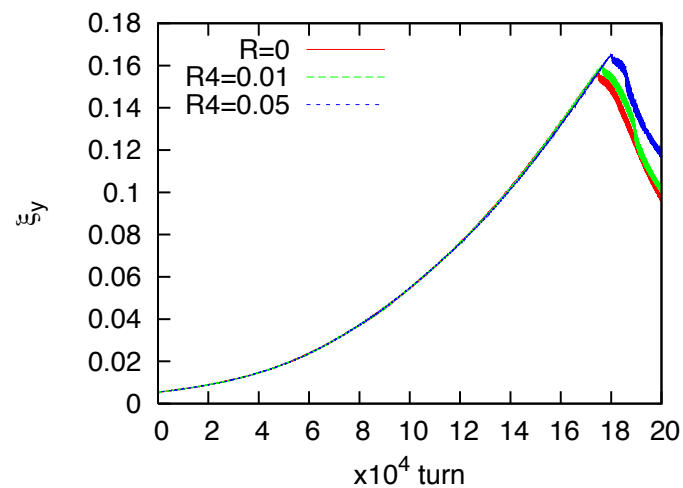
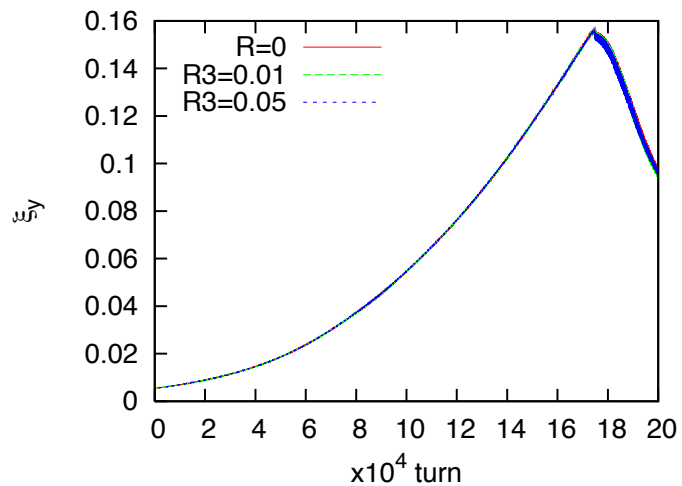
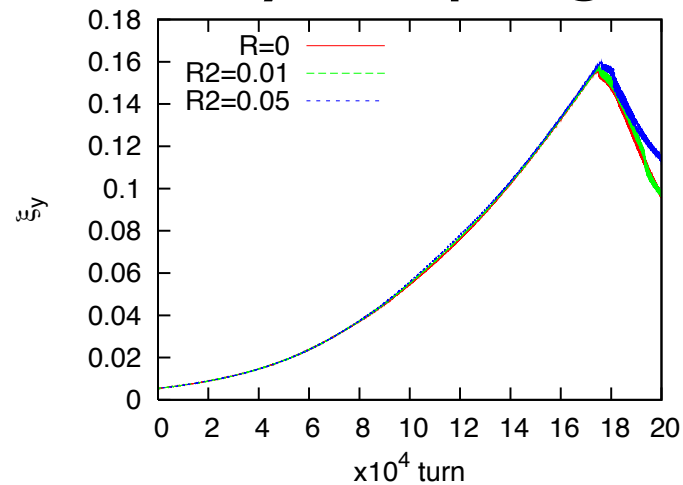
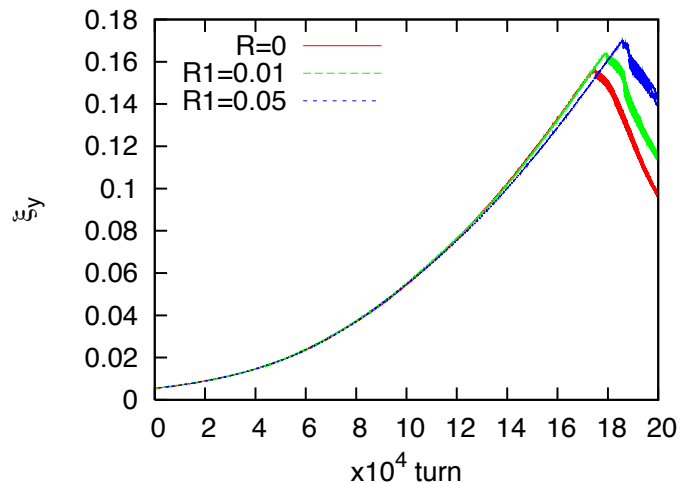
$$\begin{pmatrix} x \\ x' \\ y \\ y' \end{pmatrix} = RB \begin{pmatrix} X \\ X' \\ Y \\ Y' \end{pmatrix} \quad R = \begin{pmatrix} r_0 & 0 & r_4 & -r_2 \\ 0 & r_0 & -r_3 & r_1 \\ -r_1 & -r_2 & r_0 & 0 \\ -r_3 & -r_4 & 0 & r_0 \end{pmatrix}$$

$$B = \begin{pmatrix} \sqrt{\beta_x} & 0 & 0 & 0 \\ -\alpha_x/\sqrt{\beta_x} & 1/\sqrt{\beta_x} & 0 & 0 \\ 0 & 0 & \sqrt{\beta_y} & 0 \\ 0 & 0 & -\alpha_y/\sqrt{\beta_y} & 1/\sqrt{\beta_y} \end{pmatrix}$$

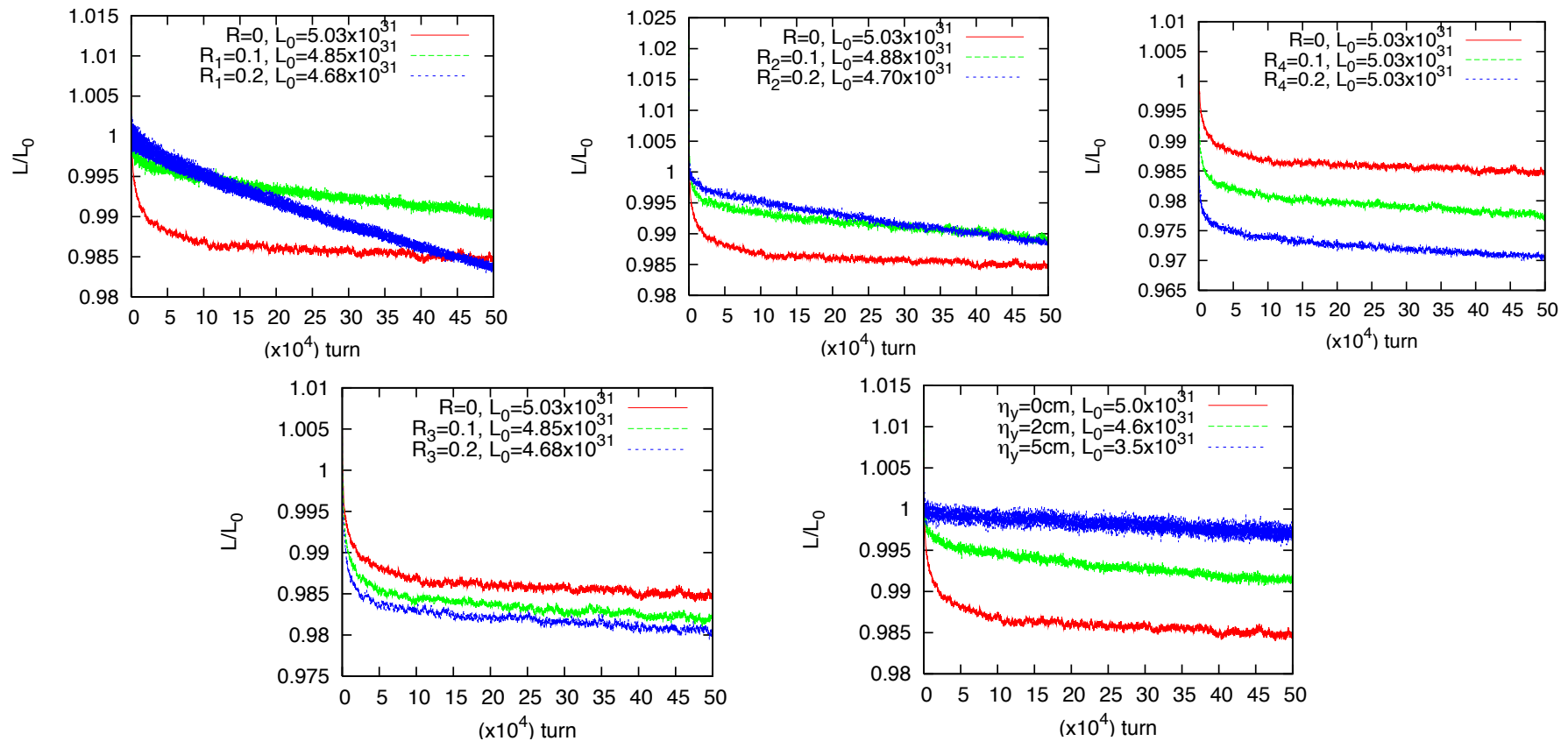
- r_i 's are a kind of Twiss parameter. They are function of s .
- Horizontal beam-beam kick induces vertical kick.

Effect of x-y coupling in coherent instability

- No remarkable effect for x-y coupling



Effect of x-y coupling for incoherent emittance growth



- $\xi/IP=(0.017,0.021)$
- Clear degradation is seen only for $R1,R2=0.2$.
- The sensitivity is 10^2 - 10^3 looser.

Summary

- Coherent instabilities are seen in the simulation, but does not seem serious.
- Incoherent emittance growth is weak for the design beam-beam parameter.
- Higher beam-beam parameter $\xi_{\text{tot}} > 0.03$ is possible in the simulation.
- IP optics error did not affect the beam-beam performance, different from KEKB. The reason may be due to the round beam.
- External noises should be taken care.
- Effect of nonlinear optics should be studied.

Beam-beam limit in proton colliders

Study of beam noise, PAC07

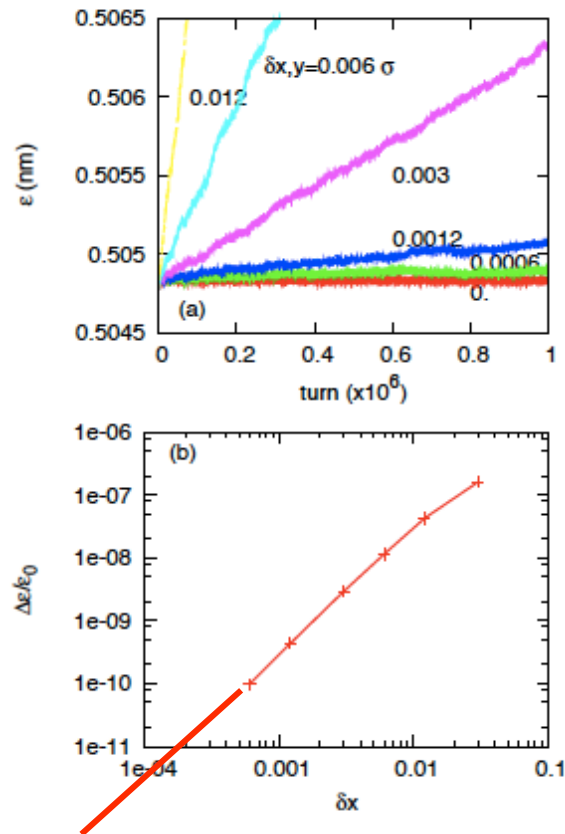
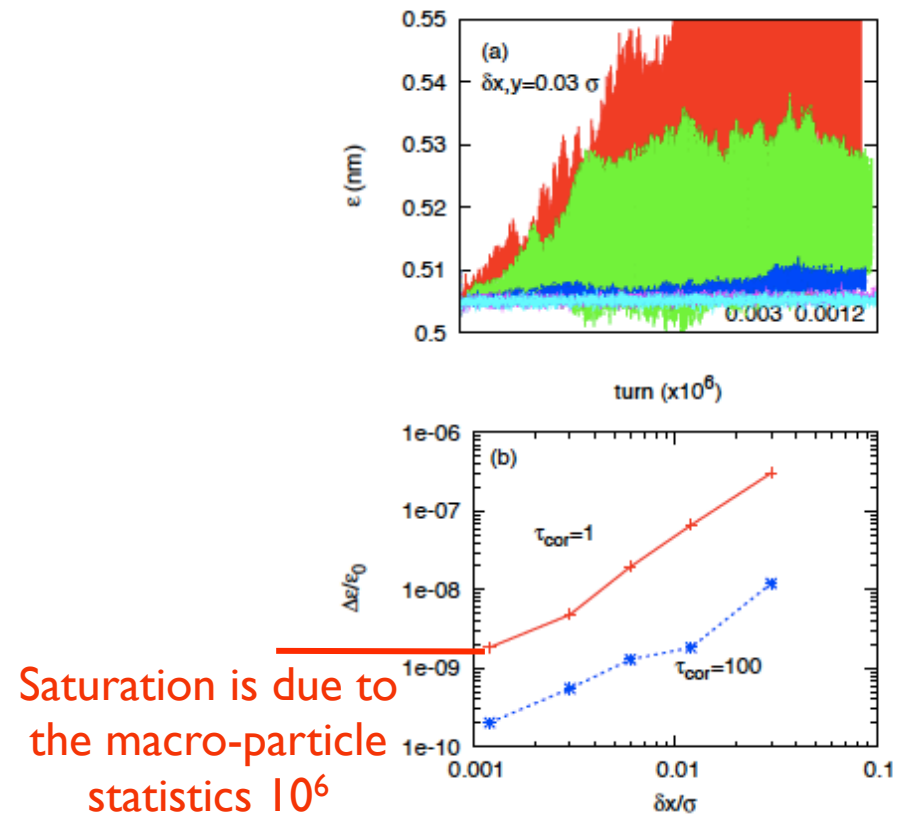


Figure 1: Emittance growth due to noise given by a weak-strong simulation. Plots (a) and (b) depict the evolution of emittance for various noise amplitude and their emittance growth rate, respectively.



Saturation is due to the macro-particle statistics 10^6

Figure 2: Emittance growth due to the fluctuation given by the strong-strong simulation. t_{cor} is the correlation time (in turns) of the fluctuation.

Beam-beam limit in proton colliders

Strong-strong or weak-strong, EPAC08

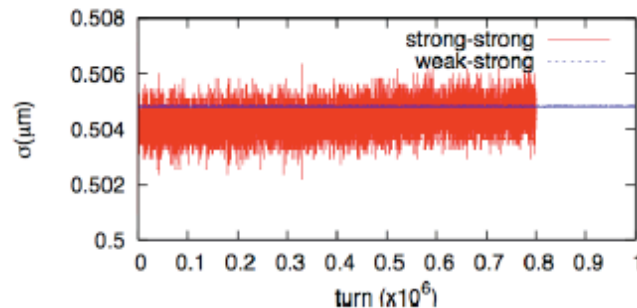


Figure 1: Emittance growth in weak-strong (blue) and strong-strong (red) simulations.

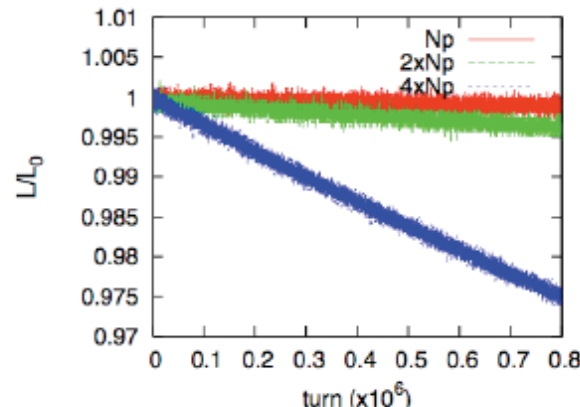


Figure 2: Luminosity decrement given by the strong-strong simulation for nominal, twice, and four times bunch populations in LHC.

$$dL/L(SS) = 10^{-9}/\text{turn}$$

$$dL/L(WS) = 10^{-14}/\text{turn}$$

This growth, which is worse than that of W.S, is fake. The weak-strong simulation is more feasible than the strong-strong one.

Beam-beam limit in proton colliders

Strong-strong or **weak-strong**, EPAC08

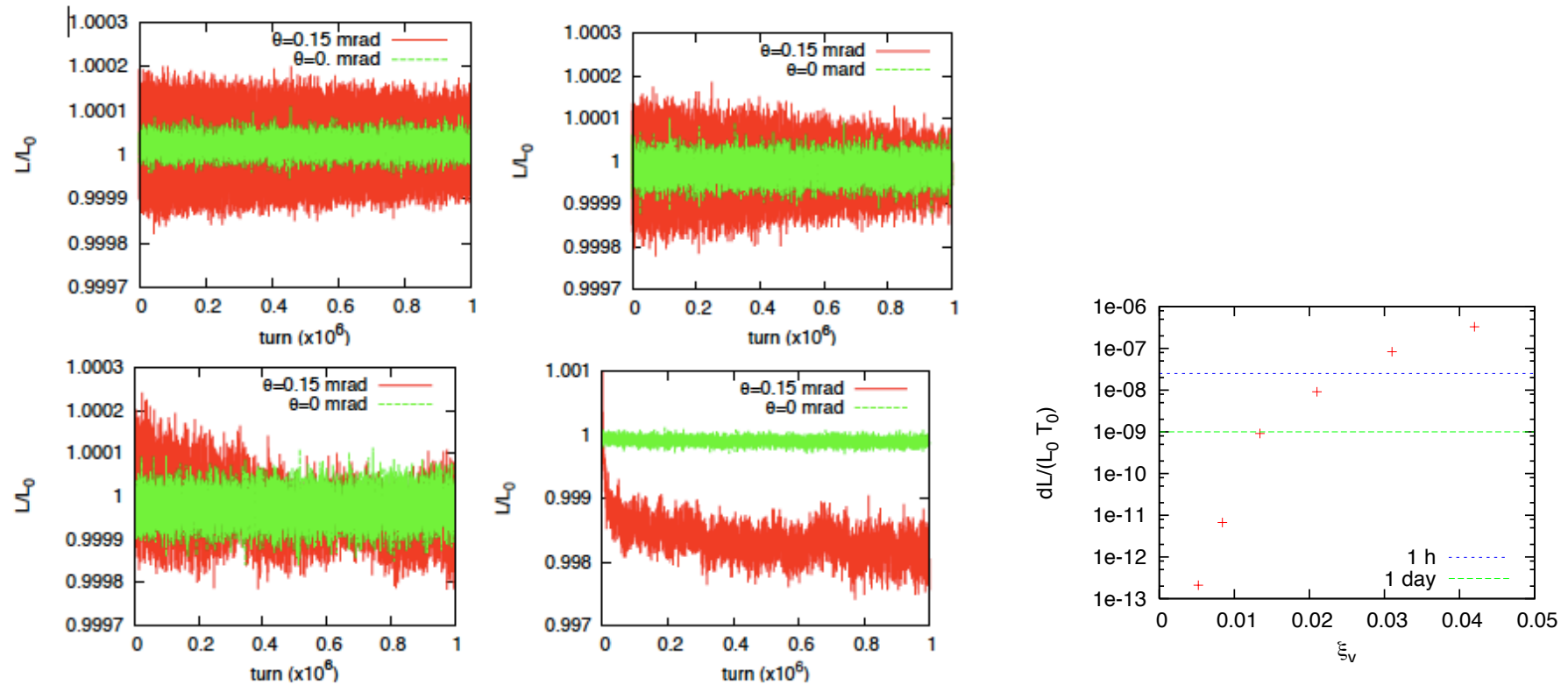
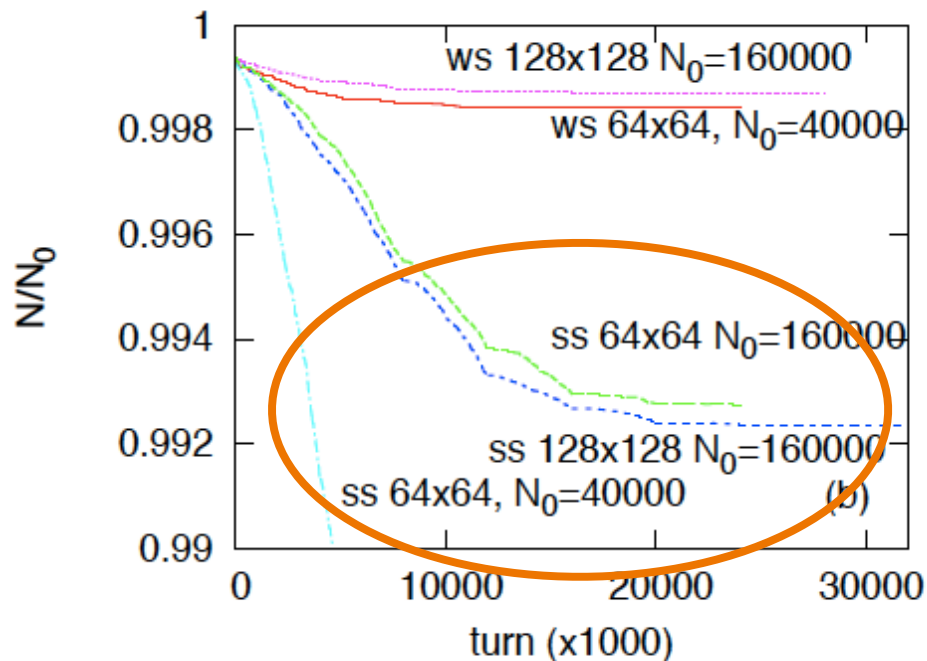


Figure 6: Luminosity decrement at the nominal LHC collision for the bunch populations, $N_p=1.15 \times 10^{11}$ (nominal), $2 \times N_p$, $4 \times N_p$ and $6 \times N_p$. Two lines plot luminosity evolution with crossing angle 0.15 mrad and 0 mrad.

Fake beam loss in Space charge simulation for J-PARC (PAC07)

- PIC simulation like S.S. beam-beam or frozen.



W.S with realistic beam distribution
(NonGaussian)

These losses should be fake too.

The numerical noise in electron machines < radiation excitation. S.S is feasible.

**Thank you for your
attention**