



Quark-mass effects in POWHEG and Hres results

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work done in collaboration with: E. Bagnaschi and G. Degrandi

POWHEG formulation and the role of h fact

$$d\sigma^{\text{NLO+PS}} = d\Phi_B \bar{B}^s(\Phi_B) \left[\Delta^s(p_\perp^{\min}) + d\Phi_{R|B} \frac{R^s(\Phi_R)}{B(\Phi_B)} \Delta^s(p_T(\Phi)) \right] + d\Phi_R R^f(\Phi_R) + d\Phi_R R_{reg}(\Phi_R)$$

$$\bar{B}^s = B(\Phi_B) + \left[V(\Phi_B) + \int d\Phi_{R|B} R^s(\Phi_{R|B}) \right]$$

$R = R_{reg} + R_{div}$ is the sum of all the real emission squared matrix elements,
with a regular (divergent) behaviour in the collinear limit

R^s enters in the Sudakov form factor $\Delta^s(p_T(\Phi))$

$$R^s = \frac{h^2}{h^2 + p_T^2} R_{div} \quad R^f = \frac{p_T^2}{h^2 + p_T^2} R_{div}$$

at low p_T , the damping factor $\rightarrow 1$, R_{div} tends to its collinear approximation,
at large p_T , the damping factor $\rightarrow 0$ and suppresses R_{div} in the Sudakov and in the square bracket

the scale h fixes the upper limit for the Sudakov form factor to play a role,
effectively is the upper limit for the inclusion of multiple parton emissions

the total cross section does NOT depend on the value of h

A proposal to treat quark-mass effects with POWHEG

- In the following identity the square bracket is a correction to the first, only-top, term because of the yukawa suppression of the bottom coupling

$$|\mathcal{M}(t + b)|^2 = |\mathcal{M}(t)|^2 + [|\mathcal{M}(t + b)|^2 - |\mathcal{M}(t)|^2]$$

- The first term contains the full top-quark squared amplitude;
the square bracket contains the top-bottom interference and the bottom squared amplitude

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- The total cross section is independent of the choice of h
→ the total cross section, including quark-mass effects, can be written as

$$\sigma(t + b) = \sigma(t, h = m_H/1.2) + [\sigma(t + b, h = m_b) - \sigma(t, h = m_b)]$$

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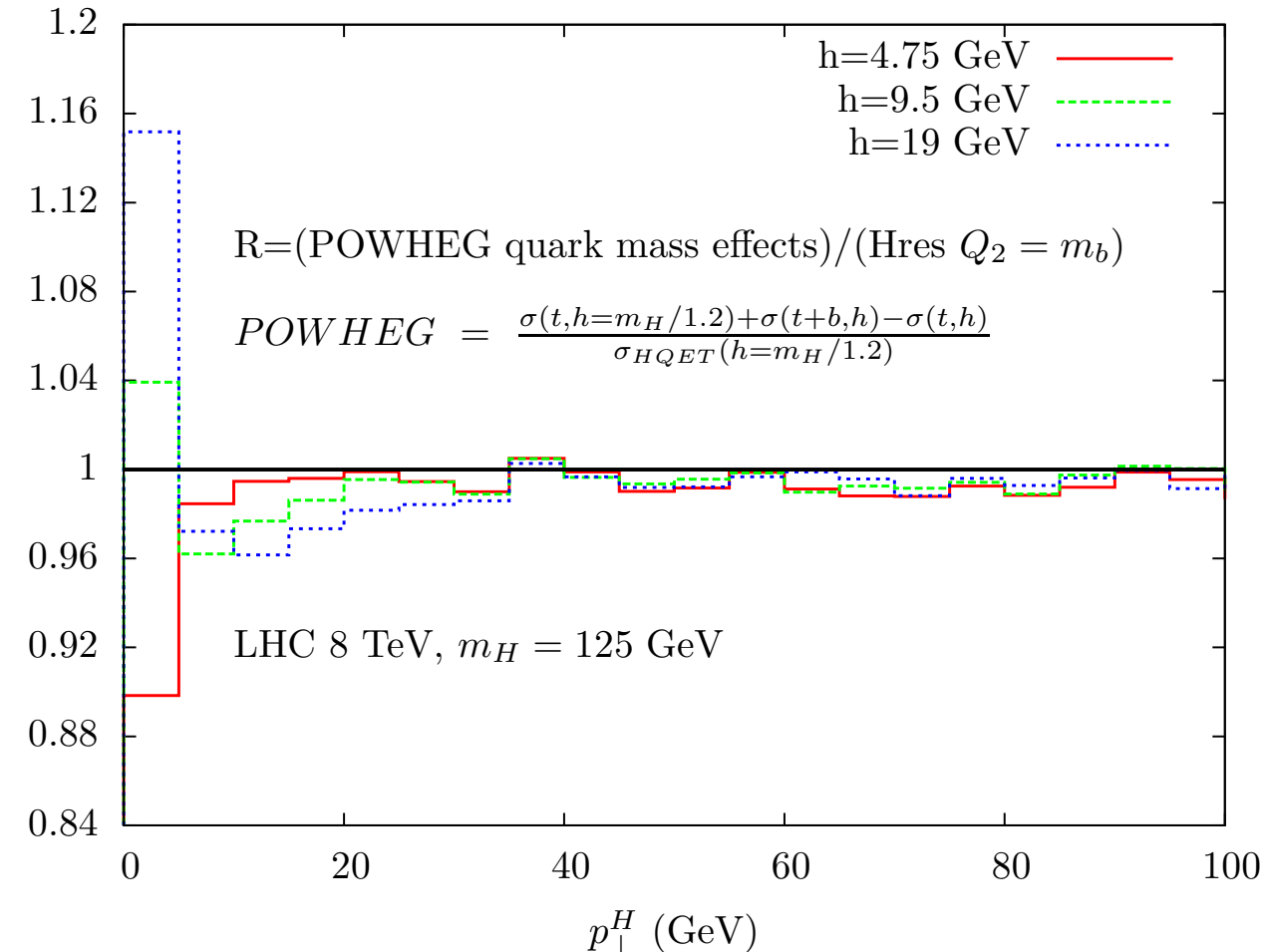
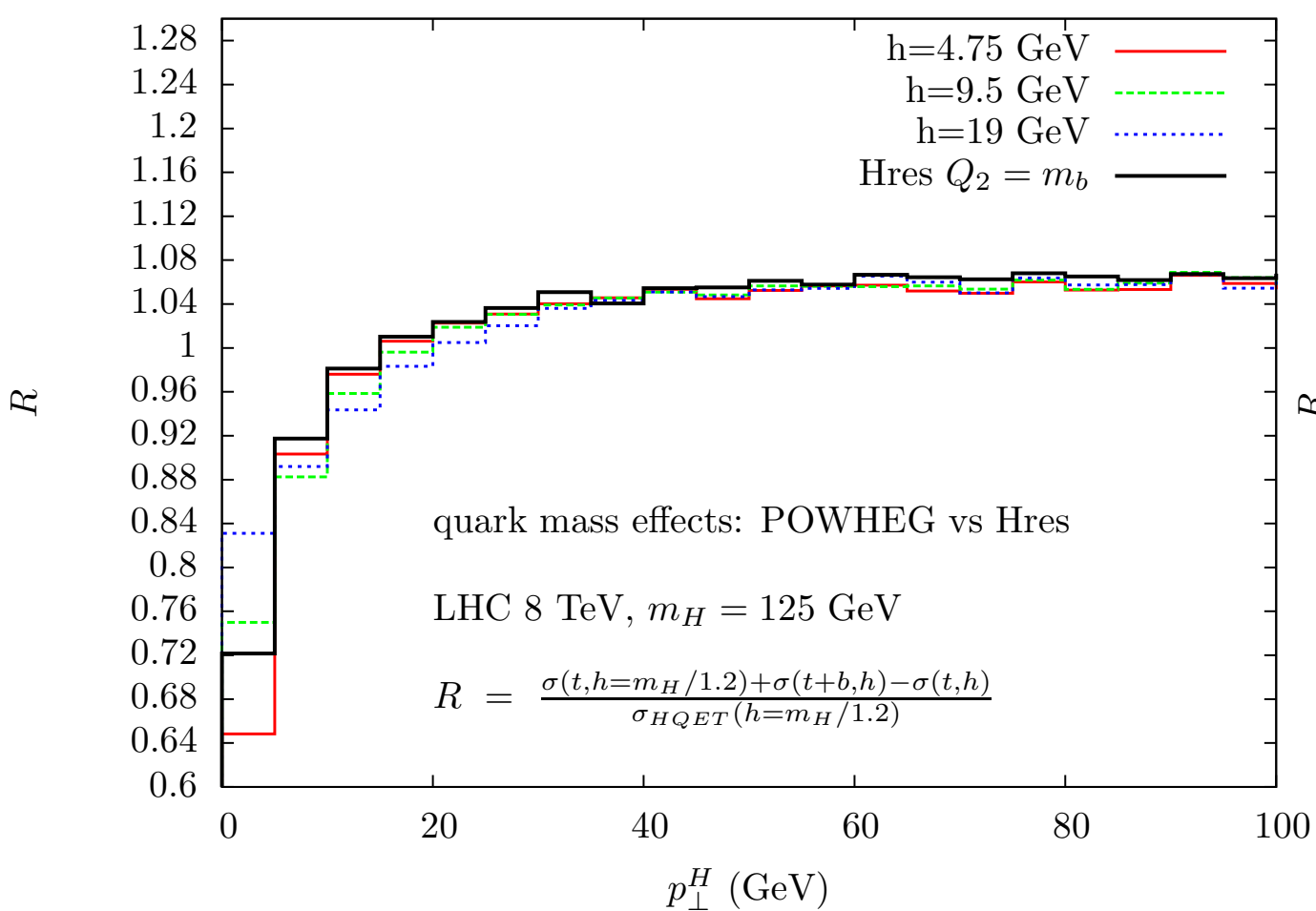
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- Since the first term depends only on the top quark, a sensible choice is $h = m_H/1.2$
- Since the square bracket contains the top-bottom interference and the bottom squared amplitude, but no pure top-quark contribution, a sensible choice is $h = m_b$
- We propose to use the above formula also for the differential distributions

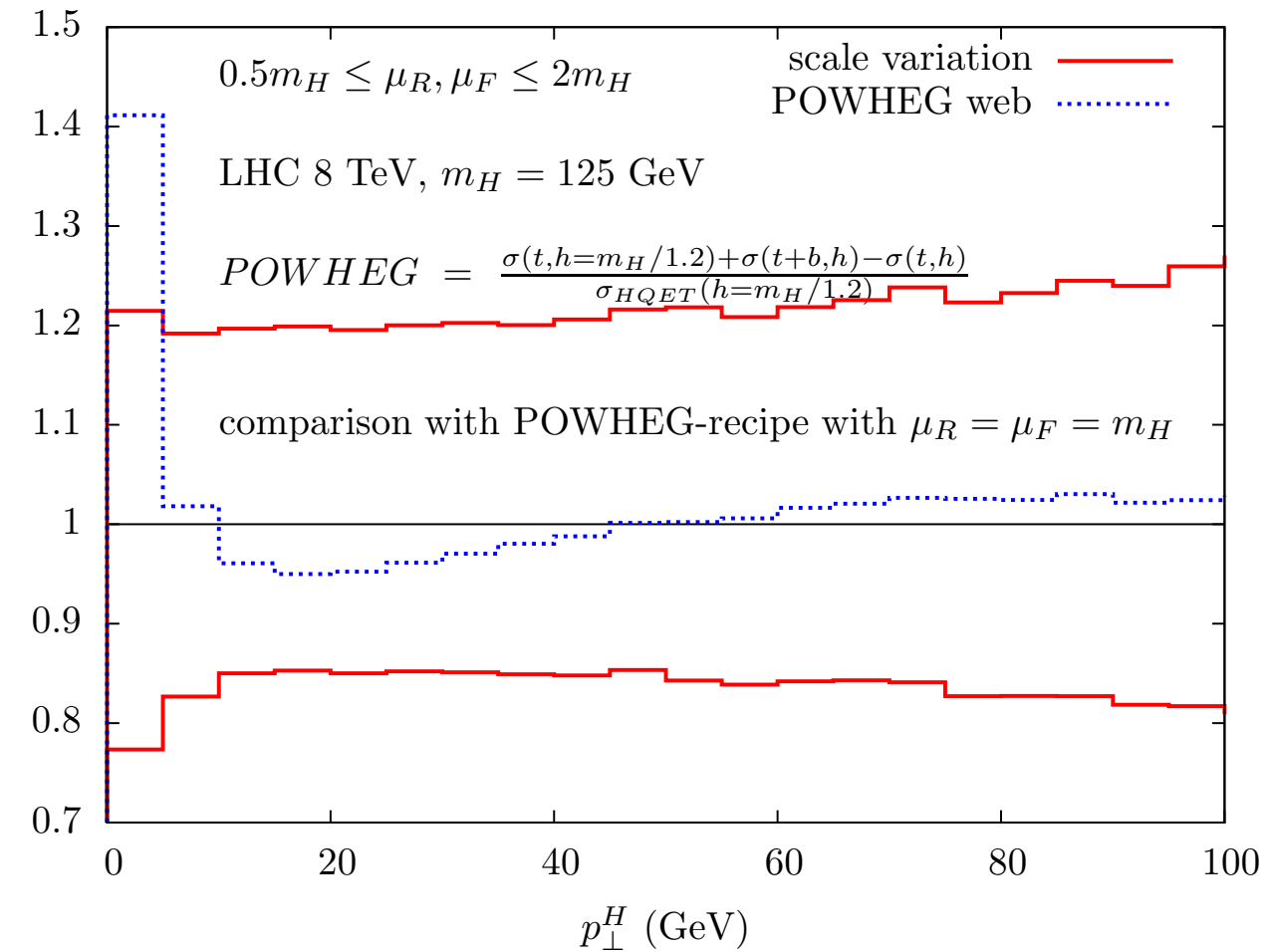
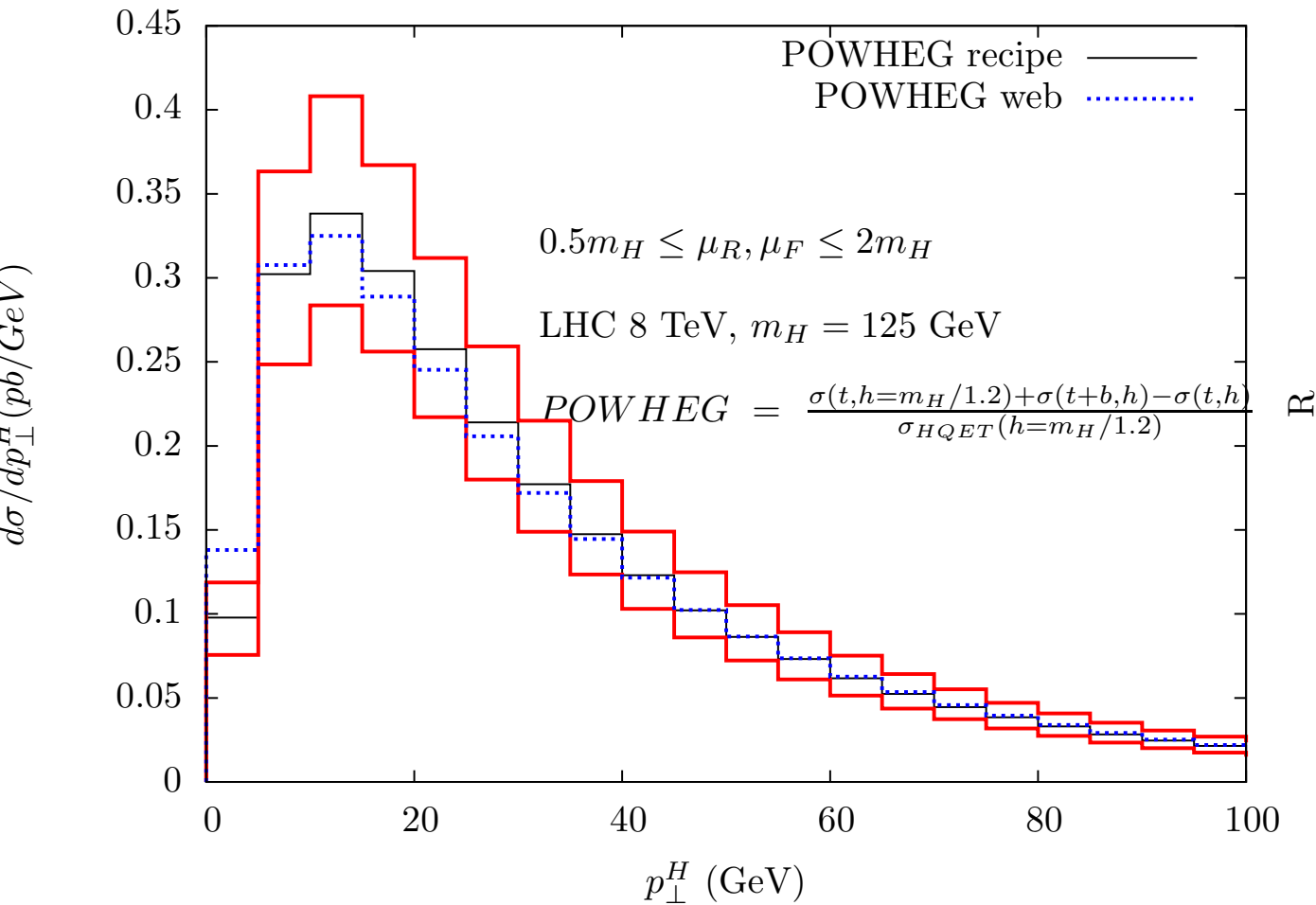
Numerical comparison with Hres

- Hres results (arXiv:1306.4581) kindly provided by M. Grazzini



- Significant suppression due to bottom mass effects in the first two bins, rather flat and positive corrections above 30 GeV
- Agreement with Hres, choosing $h=4.75$ GeV, better than 2% level, with the exception of the first bin
- The statistical accuracy and the bin size can still be improved

Scale variation (preliminary)



- Canonical renormalization and factorization scale variation (red) computed with the new recipe
- Comparison with the present quark-mass-effect POWHEG version in the POWHEG-box (blue)

Conclusions

- The use of $hfact$ to control the range where multiple parton emissions plays a role allows to treat in a different way top and bottom parts of the amplitude
- A simple combination of 3 POWHEG runs reproduces quite accurately the LO+NLL Hres calculation
- under discussion: merging the 3 contributions in one single code