



Universität
Zürich^{UZH}



HEP 2013
Stockholm
18-24 July 2013
(info@eps-hep2013.eu)



Studies of electroweak penguin transitions of $b \rightarrow s \mu \mu$

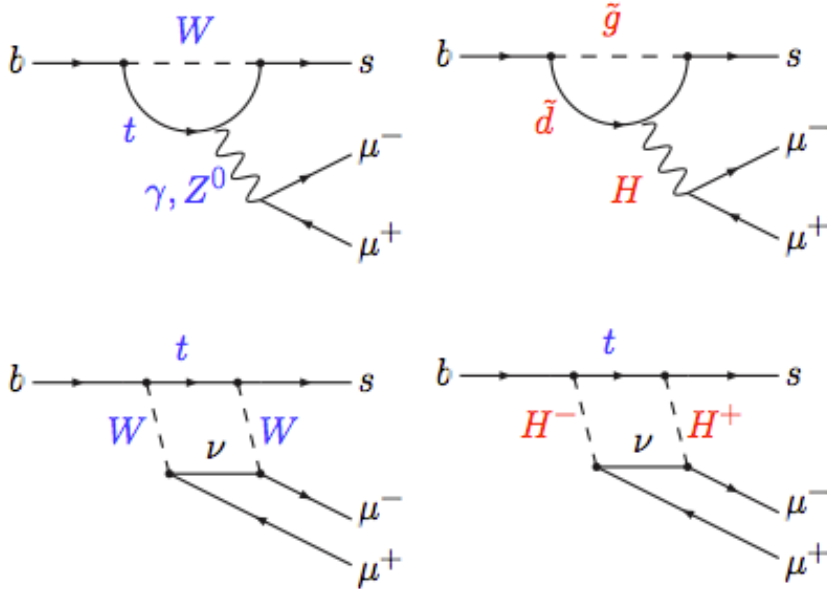
Nicola Serra on behalf of the
LHCb collaboration

Outlook

- $B^0 \rightarrow K^* \mu \mu$:
 - New observables shown here for the first time
 - $B^+ \rightarrow K^+ \mu \mu$:
 - Time-integrated CP violation (new)
 - New resonance structure at low recoil (new)
 - see talk by Giovanni Carboni
 - $B_s \rightarrow \phi \mu \mu$
 - $\Lambda_b \rightarrow \Lambda \mu \mu$
- see talk by Flavio Archilli for $B_s \rightarrow \mu \mu$

Introduction

Rare processes where new physics can enter to modify SM amplitudes



In the SM
contribution from

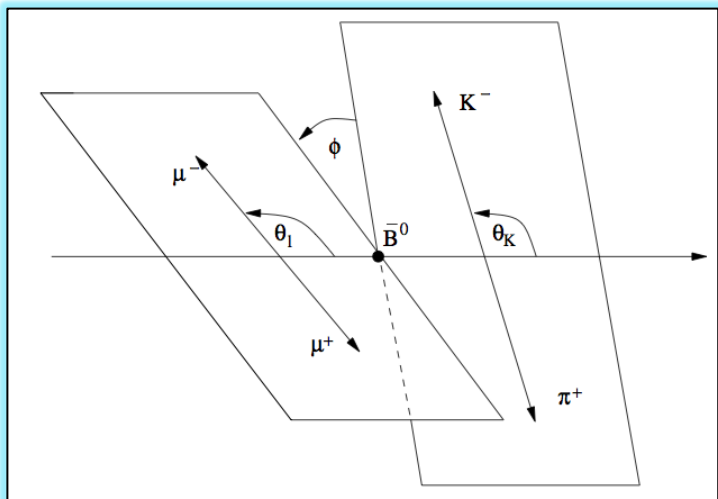
$$\mathcal{O}_7 \sim m_b (\bar{s}_L \sigma_{\mu\nu} b_R) F_{\mu\nu}$$

$$\mathcal{O}_9 \sim (\bar{s}b)_{V-A} (\bar{\ell}\ell)_V$$

$$\mathcal{O}_{10} \sim (\bar{s}b)_{V-A} (\bar{\ell}\ell)_A$$

$$H = \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum (C_i^{SM} + C_i^{NP}) O_i^{SM} + \sum \frac{c}{\Lambda_{NP}} O_{NP}$$

$$B^0 \rightarrow K^*(\rightarrow K^+ \pi^-) \mu \mu$$



The observables F_L and S_i are function of Wilson coefficients and form-factors

Well known observable is the forward-backward asymmetry of the dimuon system: $S_6^s = \frac{3}{4} A_{FB}$

$$\begin{aligned} \frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi} = & \frac{9}{32\pi} \left[\frac{3}{4} (1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4} (1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ & - F_L \cos^2 \theta_K \cos 2\theta_\ell + \\ & S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + \\ & S_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + S_6^s \sin^2 \theta_K \cos \theta_\ell + \\ & S_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + \\ & \left. S_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right] \end{aligned}$$

Altmannshofer et al. (2008)

$$B^0 \rightarrow K^* (\rightarrow K^+ \pi^-) \mu \mu$$

Observables with limited dependence on form-factor uncertainty have been proposed by several authors:

Kruger-Matias (2005), Matias et al. (2012), Egede-Matias-Hurth-Ramon-Reece (2008), Bobeth-Hiller-Van Dyk (2010-11), Beciveric-Schneider (2012)

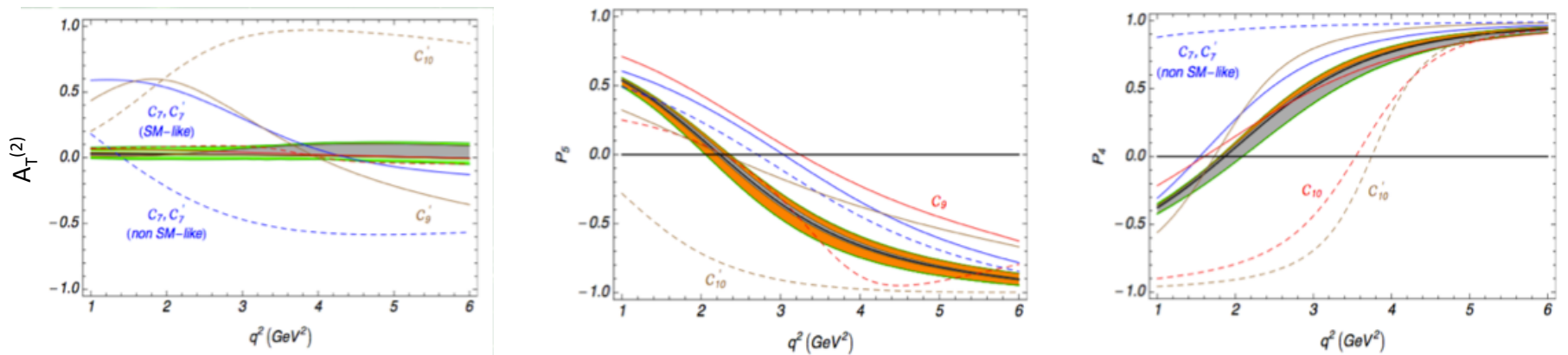
N.D.: There are other observables which are combination of those presented here

$$\begin{aligned} A_T^{(2)} &= \frac{2S_3}{(1 - F_L)} \\ A_T^{Re} &= \frac{S_6}{(1 - F_L)} \\ P'_4 &= \frac{S_4}{\sqrt{(1 - F_L)F_L}} \\ P'_5 &= \frac{S_5}{\sqrt{(1 - F_L)F_L}} \\ P'_6 &= \frac{S_7}{\sqrt{(1 - F_L)F_L}} \\ P'_8 &= \frac{S_8}{\sqrt{(1 - F_L)F_L}} \end{aligned}$$

$$\begin{aligned} \frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi} = & \frac{9}{32\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ & - F_L \cos^2 \theta_K \cos 2\theta_\ell + \frac{1}{2}(1 - F_L) A_T^{(2)} \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \\ & \sqrt{F_L(1 - F_L)} P'_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + \sqrt{F_L(1 - F_L)} P'_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + \\ & (1 - F_L) A_T^{Re} \sin^2 \theta_K \cos \theta_\ell + \sqrt{F_L(1 - F_L)} P'_6 \sin 2\theta_K \sin \theta_\ell \sin \phi + \\ & \left. \sqrt{F_L(1 - F_L)} P'_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + (S/A)_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right] \end{aligned}$$

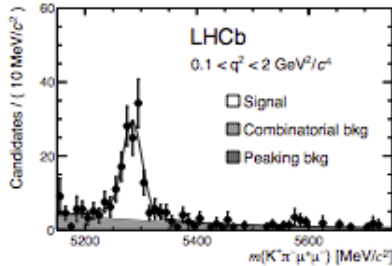
Complementarity

Reproduced from
arXiv:1207.2753

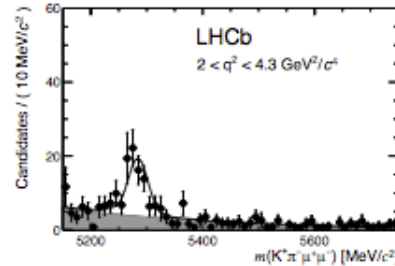


- Important to measure all observables
- Example of complementarity: solid red curve and dashed red curve
- For discussion on the complementarity of S_i see Altmannshofer et al. arXiv:0811.1214
- For discussion of P_i complementarity see J. Matias et al. arXiv:1207.2753

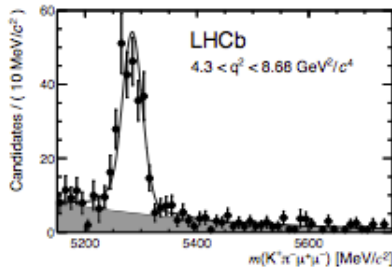
Analysis strategy



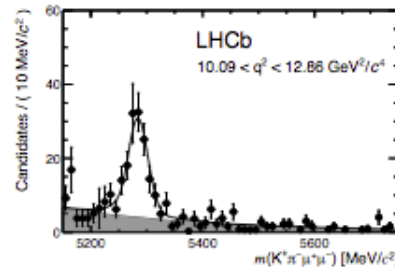
(a) $0.1 < q^2 < 2 \text{ GeV}^2/c^4$



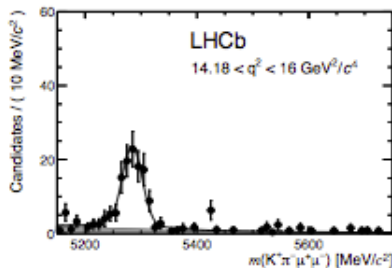
(b) $2.0 < q^2 < 4.3 \text{ GeV}^2/c^4$



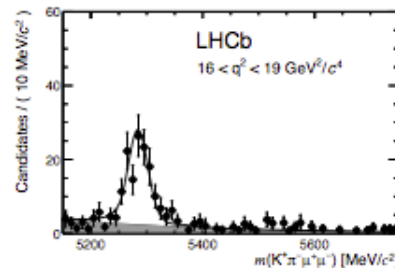
(c) $4.3 < q^2 < 8.68 \text{ GeV}^2/c^4$



(d) $10.09 < q^2 < 12.68 \text{ GeV}^2/c^4$



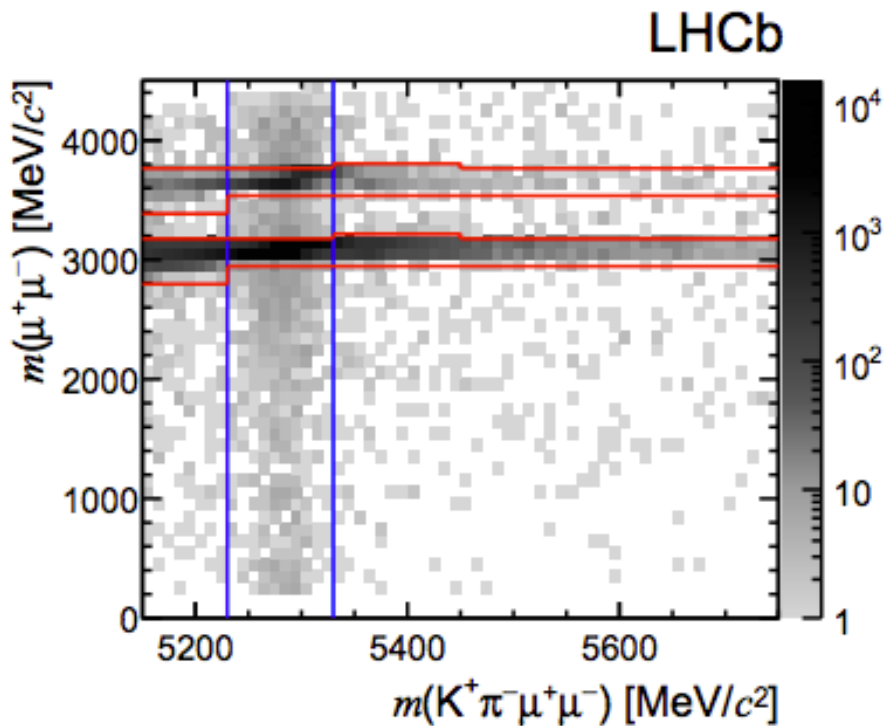
(e) $14.18 < q^2 < 16 \text{ GeV}^2/c^4$



(f) $16 < q^2 < 19 \text{ GeV}^2/c^4$

- Signal selected with a BDT (studied to keep the angular acceptance as flat as possible)
- Acceptance corrected in an event-by-event basis with the MC
- Data/MC agreement checked with control channels, e.g. $B^0 \rightarrow J/\psi K^*$
- Analysis in six bins of q^2 + the region $1 < q^2 < 6 \text{ GeV}^2$ (theory preferred at large recoil)

Analysis strategy II



Vetoed regions of dimuon mass (q) consistent with charmonium resonances

Several other peaking backgrounds examined and reduced to a negligible level

“Folding technique”

By using the transformation $\phi \rightarrow \phi + \pi$ if $\phi < 0$ (arXiv:1304.6325)

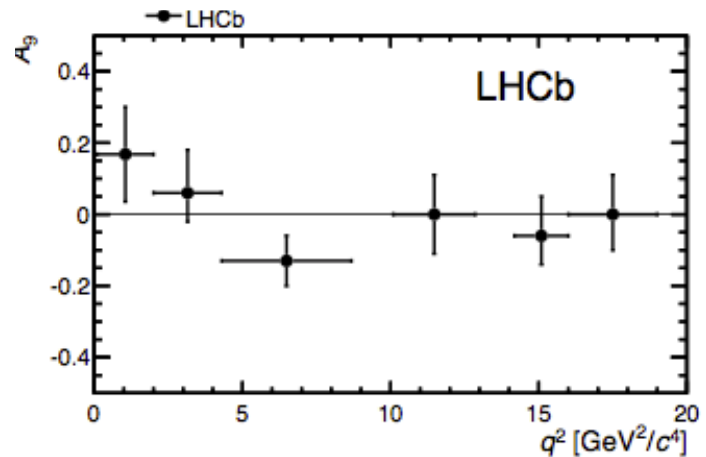
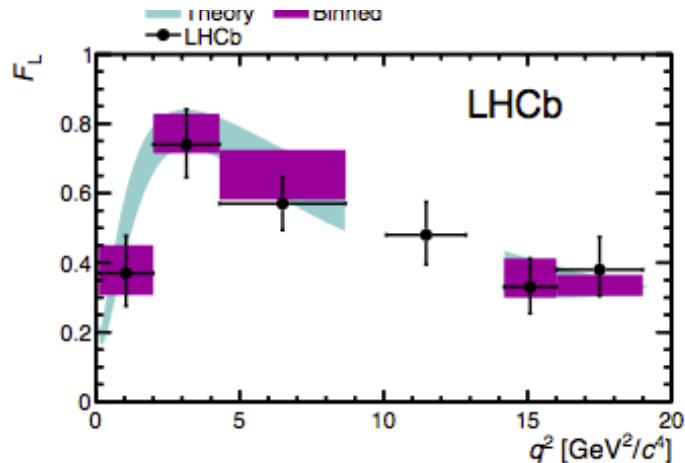
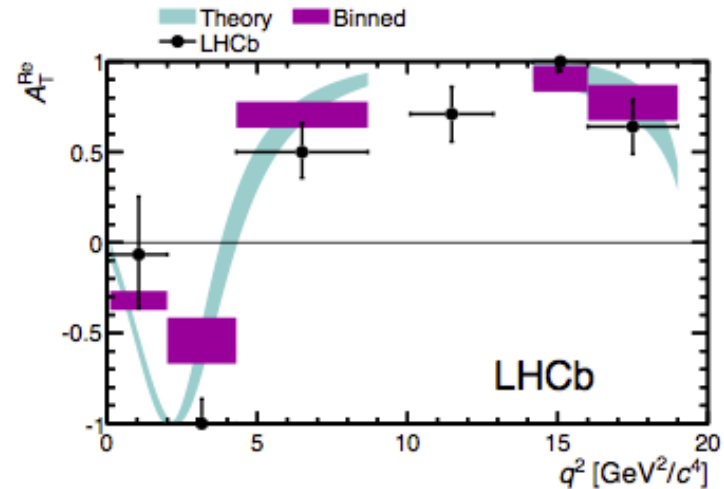
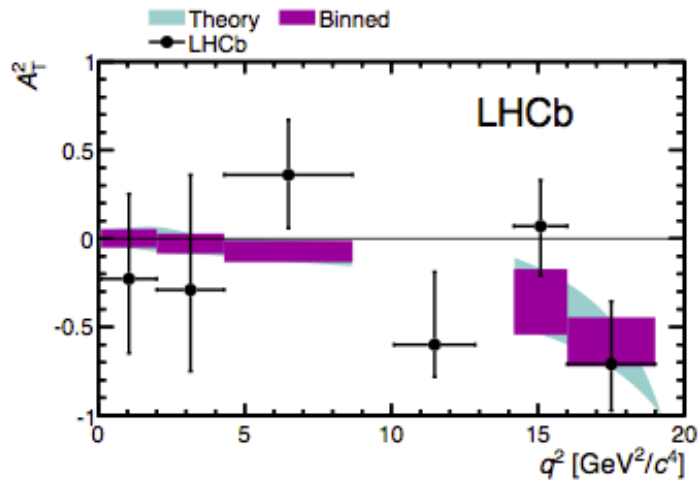
$$\frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi} = \frac{9}{16\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell \right. \\ \left. - F_L \cos^2 \theta_K \cos 2\theta_\ell + \frac{1}{2}(1 - F_L) A_T^{(2)} \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \right. \\ \left. \frac{1}{2}(1 - F_L) A_T^{Re} \sin^2 \theta_K \cos \theta_\ell + (S/A)_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right]$$

Measurement of the other observables with other folding techniques
for P_5' (or equivalently S_5) $\phi \rightarrow -\phi$ (if $\phi < 0$) and $\vartheta_l \rightarrow \pi - \vartheta_l$ (if $\vartheta_l < \pi/2$)

$$\frac{1}{\Gamma} \frac{d^3(\Gamma + \bar{\Gamma})}{d \cos \theta_\ell d \cos \theta_K d\phi} = \frac{9}{8\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - \right. \\ \left. F_L \cos^2 \theta_K \cos 2\theta_\ell + \frac{1}{2}(1 - F_L) A_T^{(2)} \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \right. \\ \left. \sqrt{F_L(1 - F_L)} P_5' \sin 2\theta_K \sin \theta_\ell \cos \phi \right]$$

The other transformations for $P_{4,6,8}'$ (or $S_{4,7,8}$) are in the backup slides

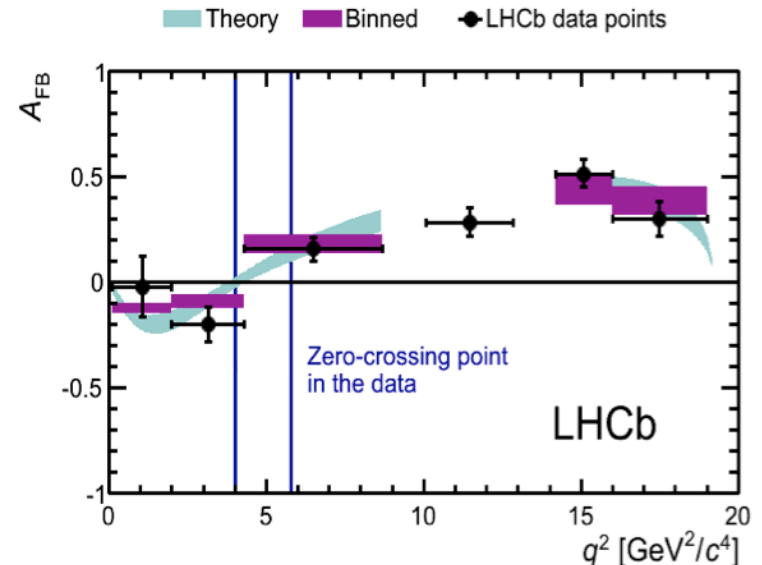
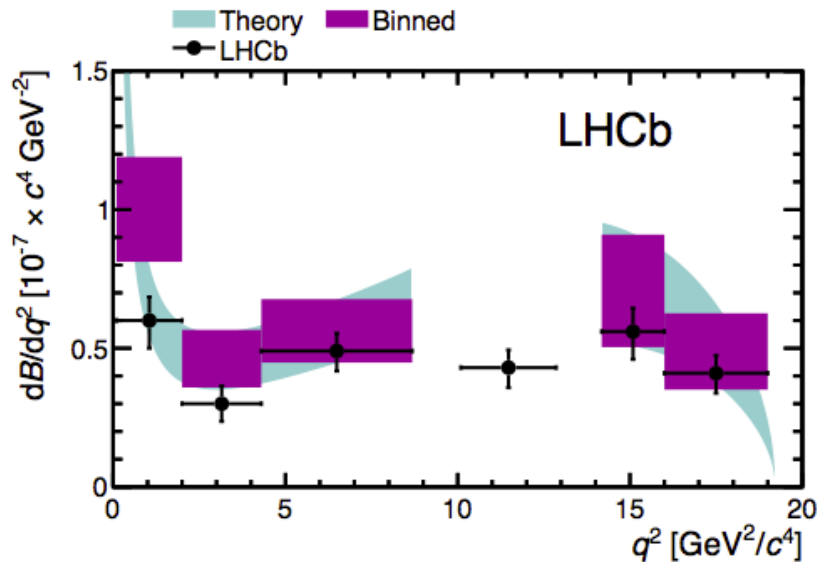
Results



Good agreement with SM predictions
Bobeth-Hiller-Van Dyk (2011); Form-factor from Ball-Zwicky (2005);
Consistent with Matias et al.(2013)

Results

LHCb collaboration (1fb⁻¹), arXiv:1304.6325



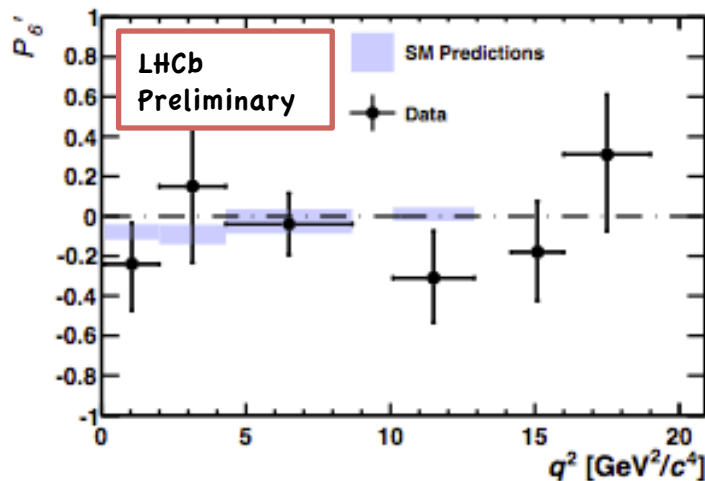
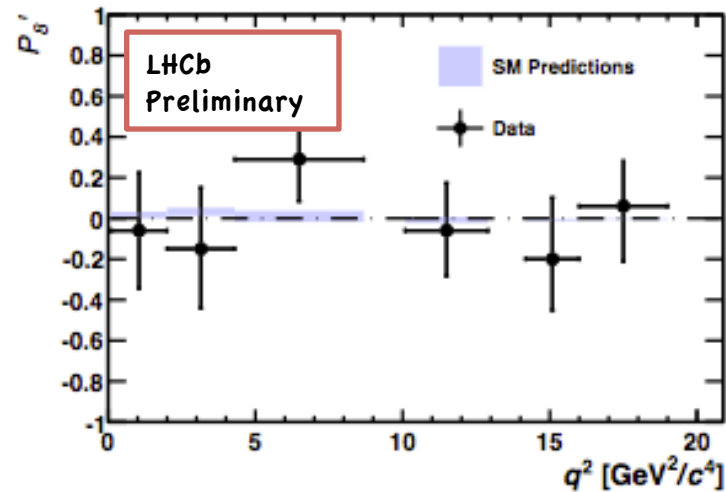
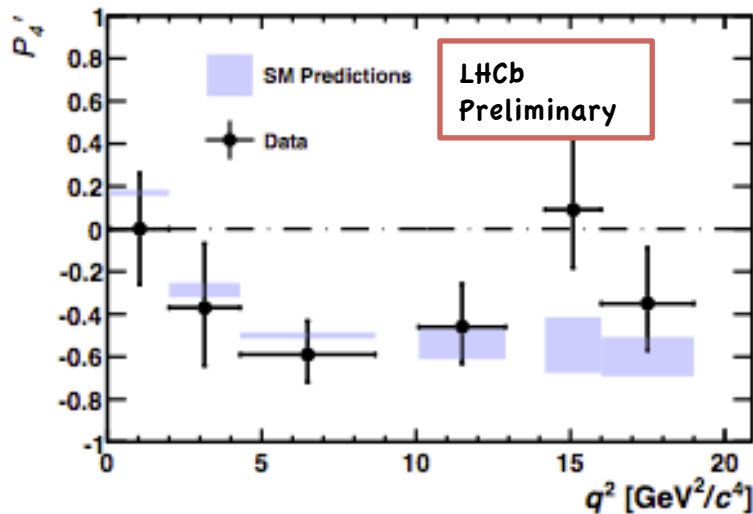
Good agreement with SM predictions

First measurement of the zero-crossing point: $q_0^2 = 4.9 \pm 0.9$ GeV²/c⁴

Results for new observables

NEW

LHCb collaboration (1fb⁻¹), LHCb-PAPER-2013-037

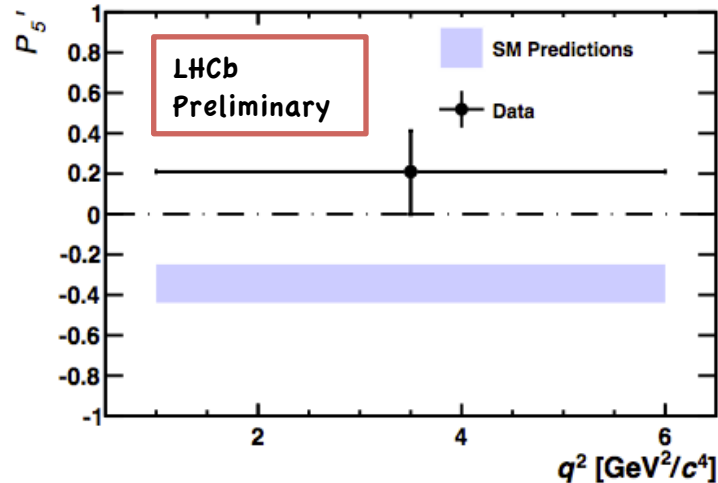
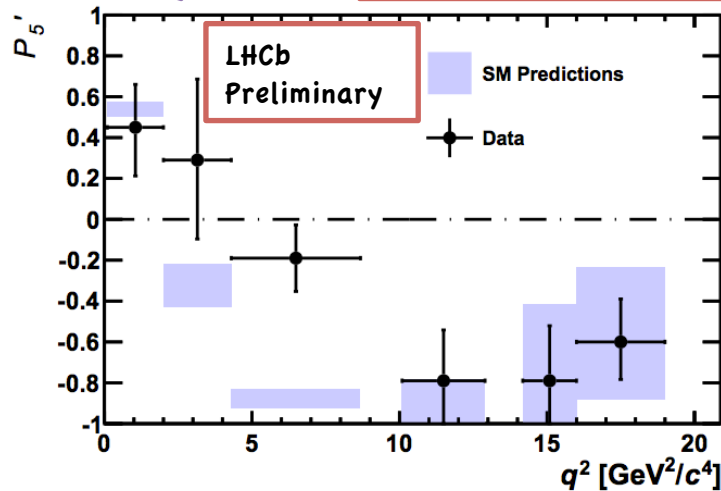


Good agreement with SM predictions (from J.Matias et al. arXiv:1303.5794)

Results for new observables

NEW

LHCb collaboration (1fb⁻¹), LHCb-PAPER-2013-037



- Discrepancy with respect to SM predictions (arXiv:1303.5794) at low q^2
- 3.7 sigma discrepancy in the region $4.3 < q^2 < 8.68$ GeV²/c⁴
- 0.5% probability (2.8 sigma) to observe such a deviation considering 24 independent measurements)
- 2.5 sigma discrepancy in the region $1.0 < q^2 < 6.0$ GeV²/c⁴

N.B.: Jaeger-Camelich (arXiv:1212.2263) have predictions in the region $1.0 < q^2 < 6.0$ GeV²/c⁴ with much larger theoretical error and small shift in the central value (QCD factorization breaking + $c\bar{c}$ loop)

Results for new observables

LHCb Preliminary

$q^2 [\text{GeV}^2/c^4]$	P'_4	P'_5	P'_6	P'_8
0.10 – 2.00	$0.00^{+0.26}_{-0.26} \pm 0.03$	$0.45^{+0.19}_{-0.22} \pm 0.09$	$-0.24^{+0.19}_{-0.22} \pm 0.05$	$-0.06^{+0.28}_{-0.28} \pm 0.02$
2.00 – 4.30	$-0.37^{+0.29}_{-0.26} \pm 0.08$	$0.29^{+0.39}_{-0.38} \pm 0.07$	$0.15^{+0.36}_{-0.38} \pm 0.05$	$-0.15^{+0.29}_{-0.28} \pm 0.07$
4.30 – 8.68	$-0.59^{+0.15}_{-0.12} \pm 0.05$	$-0.19^{+0.16}_{-0.16} \pm 0.03$	$-0.04^{+0.15}_{-0.15} \pm 0.05$	$0.29^{+0.17}_{-0.19} \pm 0.03$
10.09 – 12.90	$-0.46^{+0.20}_{-0.17} \pm 0.03$	$-0.79^{+0.16}_{-0.19} \pm 0.19$	$-0.31^{+0.23}_{-0.22} \pm 0.05$	$-0.06^{+0.23}_{-0.22} \pm 0.02$
14.18 – 16.00	$0.09^{+0.35}_{-0.27} \pm 0.04$	$-0.79^{+0.20}_{-0.13} \pm 0.18$	$-0.18^{+0.25}_{-0.24} \pm 0.03$	$-0.20^{+0.30}_{-0.25} \pm 0.03$
16.00 – 19.00	$-0.35^{+0.26}_{-0.22} \pm 0.03$	$-0.60^{+0.19}_{-0.16} \pm 0.09$	$0.31^{+0.38}_{-0.37} \pm 0.10$	$0.06^{+0.26}_{-0.27} \pm 0.03$
1.00 – 6.00	$-0.29^{+0.18}_{-0.16} \pm 0.03$	$0.21^{+0.20}_{-0.21} \pm 0.03$	$-0.18^{+0.21}_{-0.21} \pm 0.03$	$0.23^{+0.18}_{-0.19} \pm 0.02$

$q^2 [\text{GeV}^2/c^4]$	S_4	S_5	S_7	S_8
0.10 – 2.00	$-0.01^{+0.12}_{-0.12} \pm 0.03$	$0.22^{+0.09}_{-0.10} \pm 0.04$	$-0.12^{+0.11}_{-0.11} \pm 0.03$	$-0.04^{+0.13}_{-0.12} \pm 0.01$
2.00 – 4.30	$-0.14^{+0.13}_{-0.12} \pm 0.03$	$0.11^{+0.14}_{-0.13} \pm 0.03$	$0.06^{+0.15}_{-0.15} \pm 0.02$	$-0.05^{+0.12}_{-0.12} \pm 0.02$
4.30 – 8.68	$-0.29^{+0.06}_{-0.06} \pm 0.02$	$-0.09^{+0.08}_{-0.08} \pm 0.01$	$-0.03^{+0.07}_{-0.08} \pm 0.04$	$0.13^{+0.08}_{-0.08} \pm 0.01$
10.09 – 12.90	$-0.22^{+0.09}_{-0.08} \pm 0.02$	$-0.40^{+0.08}_{-0.10} \pm 0.10$	$-0.17^{+0.11}_{-0.12} \pm 0.03$	$-0.03^{+0.10}_{-0.10} \pm 0.01$
14.18 – 16.00	$0.05^{+0.14}_{-0.08} \pm 0.01$	$-0.38^{+0.10}_{-0.09} \pm 0.09$	$-0.08^{+0.13}_{-0.14} \pm 0.01$	$-0.10^{+0.13}_{-0.12} \pm 0.02$
16.00 – 19.00	$-0.16^{+0.11}_{-0.09} \pm 0.01$	$-0.29^{+0.09}_{-0.08} \pm 0.04$	$0.15^{+0.16}_{-0.15} \pm 0.03$	$0.03^{+0.12}_{-0.12} \pm 0.02$

$B^+ \rightarrow K^+ \mu \mu$

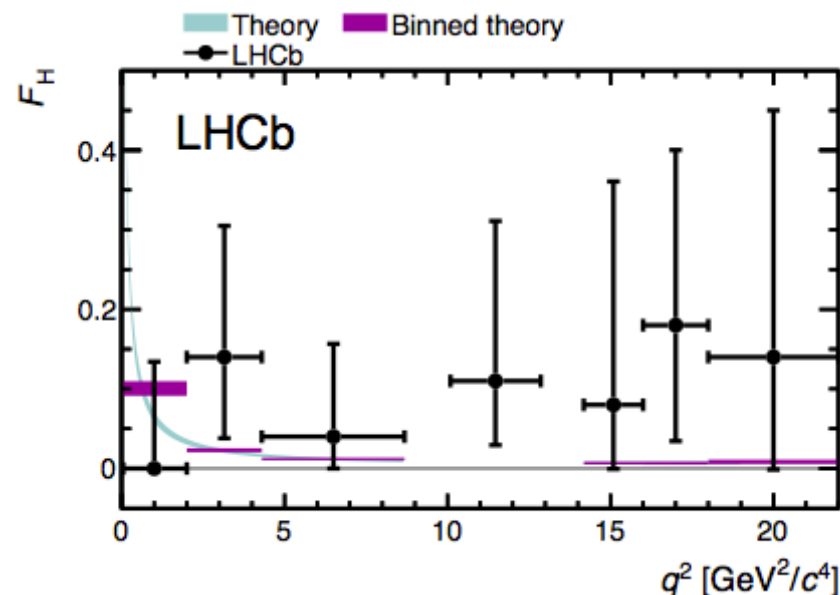
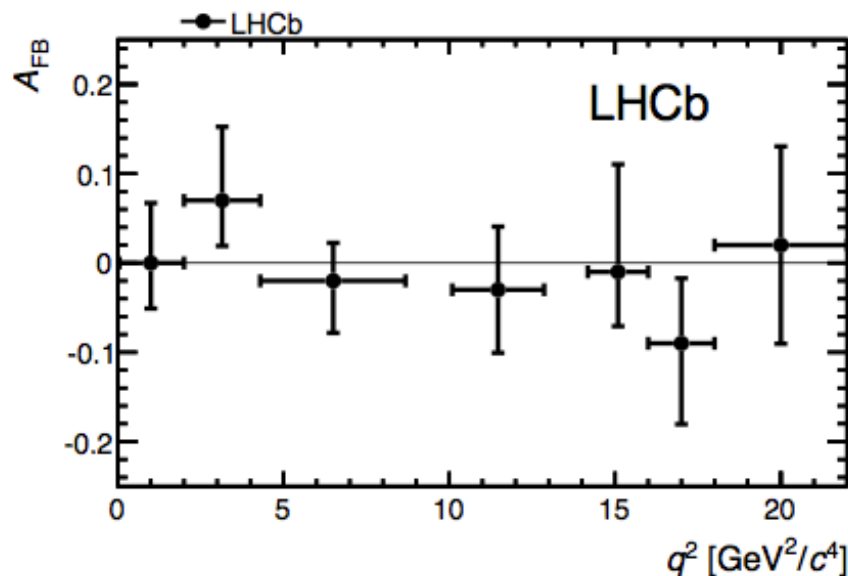
LHCb collaboration (1fb^{-1})
JHEP 02 (2013) 105

Decay described by a single angle ϑ_l

$$\frac{1}{\Gamma} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{d\cos\theta_l} = \frac{3}{4}(1 - F_H)(1 - \cos^2\theta_l) + \frac{1}{2}F_H + A_{\text{FB}} \cos\theta_l$$

F_H and A_{FB} sensitive to θ_p and θ_s

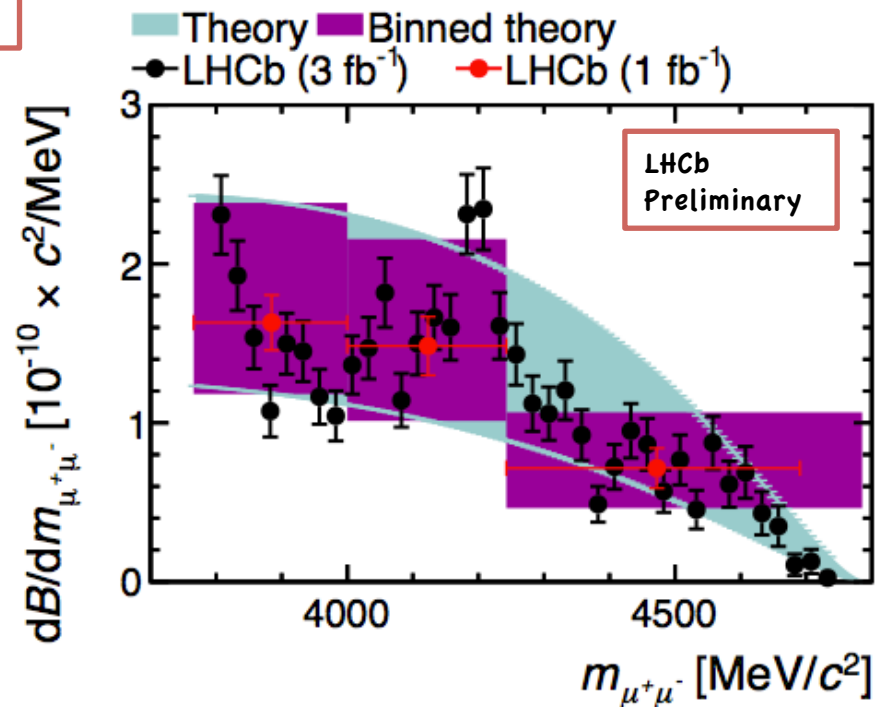
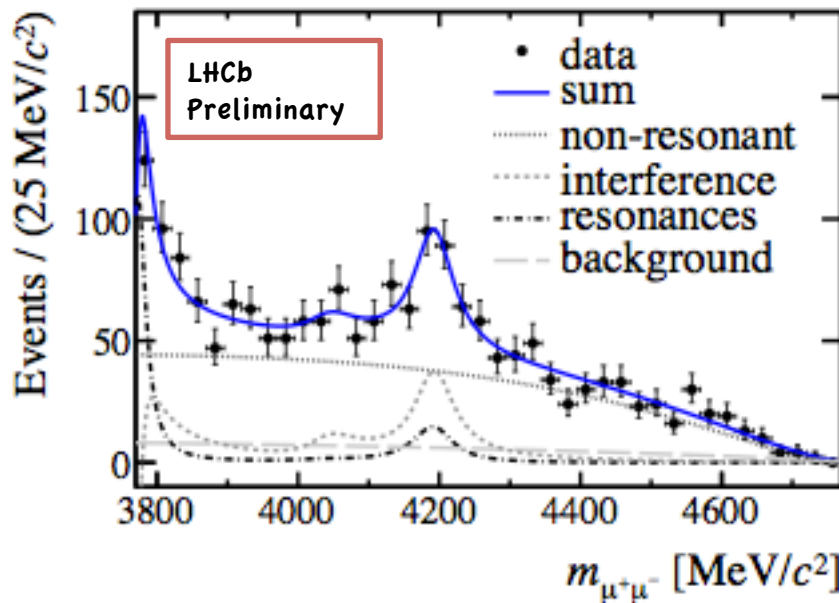
In the SM both observables are about zero



New resonance structure

NEW

LHCb collaboration (3fb^{-1})
LHCb-PAPER-2013-039



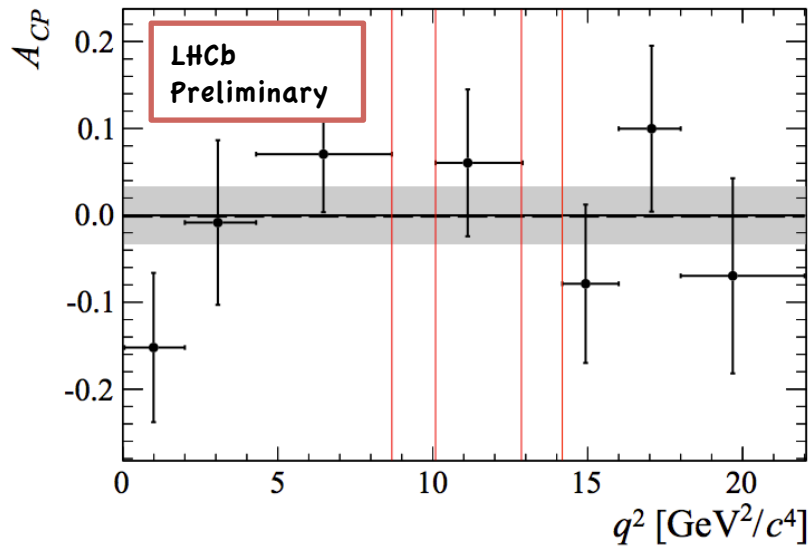
Observation of a resonant structure ($\psi(4160)$) at high q^2
First observation of $B^+ \rightarrow K^+ \psi(4160)$ and of $\psi(4160) \rightarrow \mu\mu$
Unexpected result $\rightarrow$ large influence on EWP at low recoil

See talk by Giovanni Carboni

Time-integrated CP Violation

NEW

LHCb collaboration (1fb⁻¹)
LHCb-PAPER-2013-043



Production and detection asymmetry taken into account by using the $B^+ \rightarrow J/\psi K^+$ decay

$$A_{CP}(K^+ \mu \mu) = A_{\text{RAW}}(K^+ \mu \mu) - A_{\text{RAW}}(J/\psi K^+)$$

Left-right detector asymmetry removed averaging magnet polarities

LHCb Preliminary:

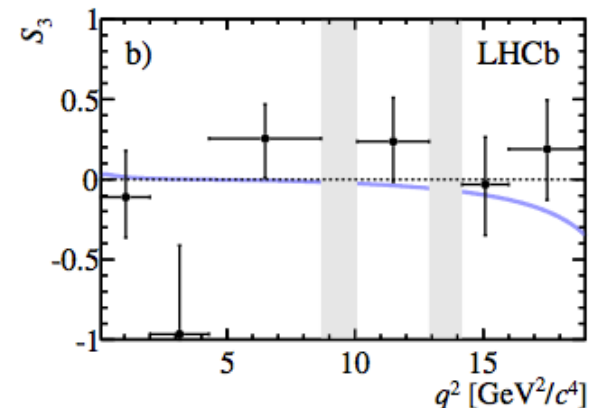
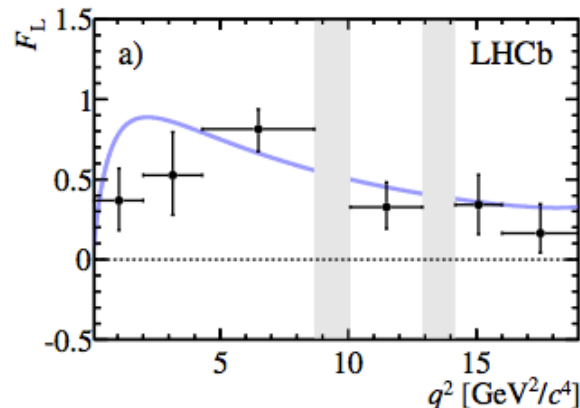
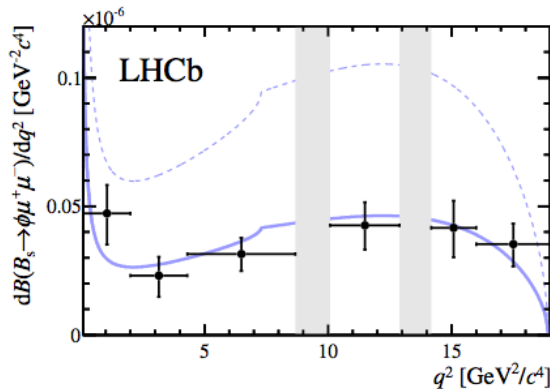
$$A_{CP}(K^+ \mu \mu) = 0.000 \pm 0.033(\text{stat}) \pm 0.005(\text{syst}) \pm 0.007(\text{norm})$$

$B_s^0 \rightarrow \phi \mu \mu$

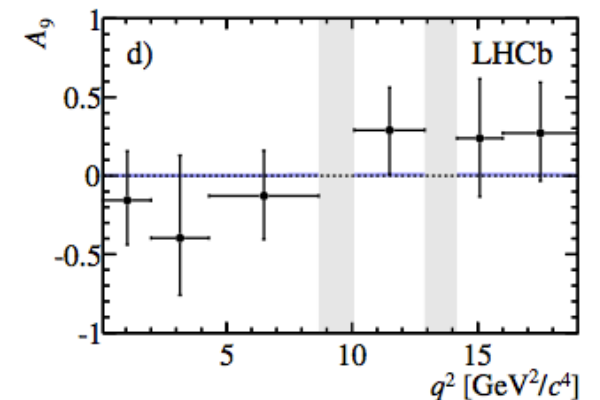
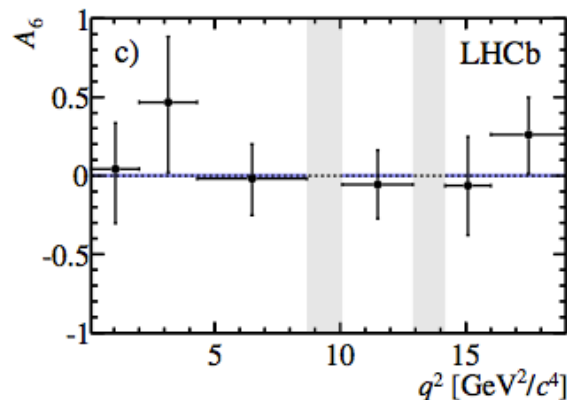
LHCb collaboration (1fb⁻¹)
arXiv:1305.2168

Similar to $B^0 \rightarrow K^* \mu \mu$, but not self tagging.
for certain angular terms only A are accessible!

See Altmannshofer et al. (2008) and Bobeth-Hiller-Piranishvili (2008)



- Total BR lower than predictions
- Good agreement for the angular observables

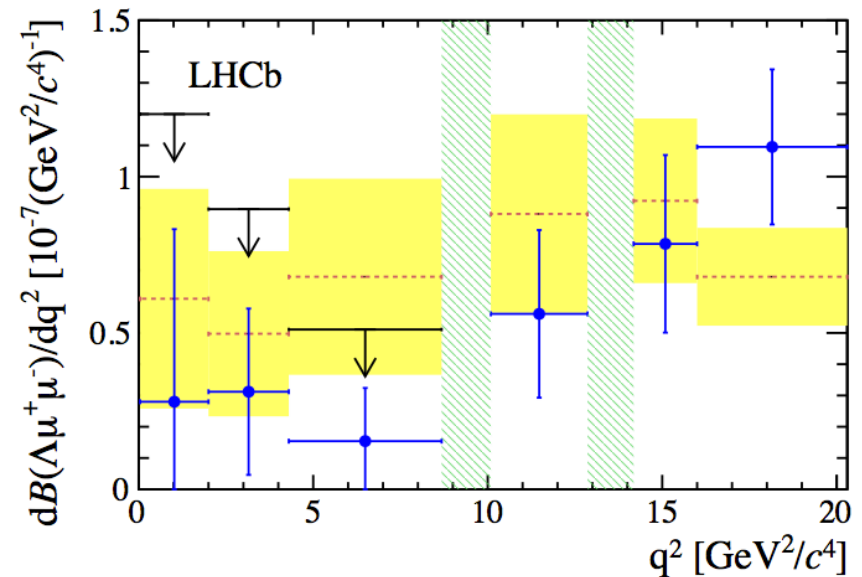
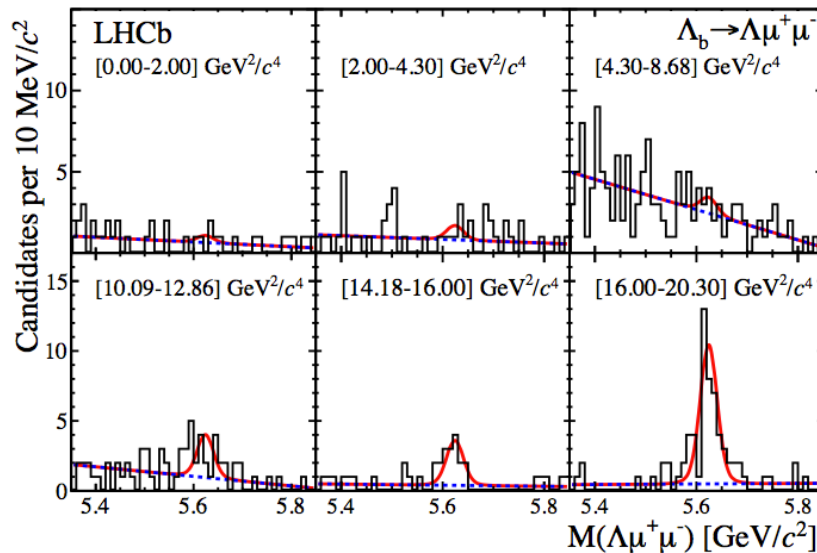


$\Lambda_b \rightarrow \Lambda \mu \mu$

LHCb collaboration (1fb⁻¹), arXiv:1306.2577

- Additional complexity/observables wrt mesons
- First step measuring the yield
- Decay $\Lambda_b \rightarrow \Lambda J/\psi$ used as a control channel

(Hiller et al. 2007 and ref therein)



Predictions from Detmold et al. (2012)

Observed about 80 events in the whole q^2 region

$$\text{BR}(\Lambda_b \rightarrow \Lambda \mu \mu) = (0.96 \pm 0.16(\text{stat}) \pm 0.13(\text{syst}) \pm 0.21(\text{norm})) \times 10^{-6}$$

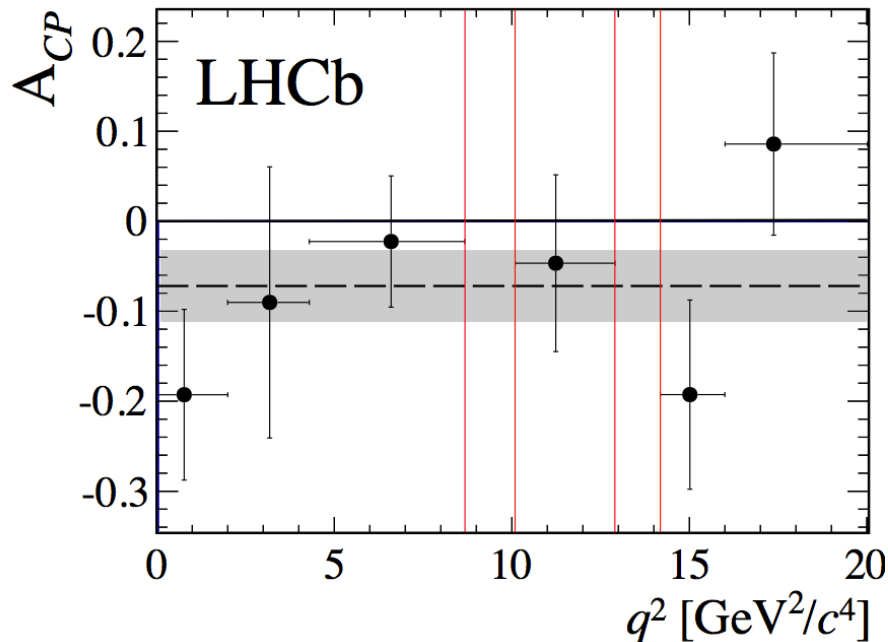
Conclusions

- LHCb has new results on rare electroweak penguins
- Generally good agreement with the SM, apart for a local discrepancy in the low q^2 region in the observable P_5' in the decay $B^0 \rightarrow K^* \mu \mu$
- Most of the results are obtained with 1fb^{-1} ... stay tuned for the 3fb^{-1} analyses

Backup slides

Time-integrated CP Violation

Phys. Rev. Lett. 110 (2013) 031801



$$A_{CP}(K^*\mu\mu) = -0.072 \pm 0.040 \pm 0.005$$

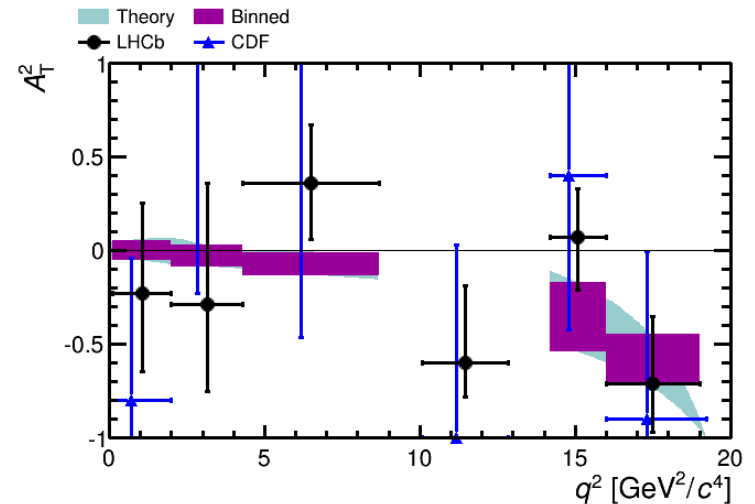
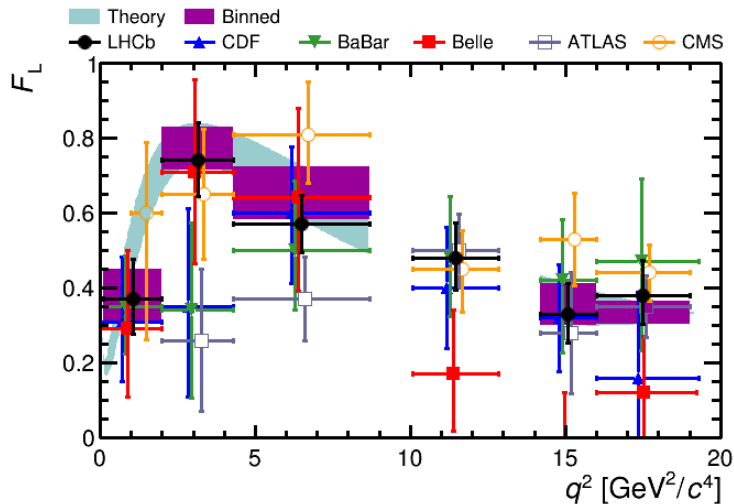
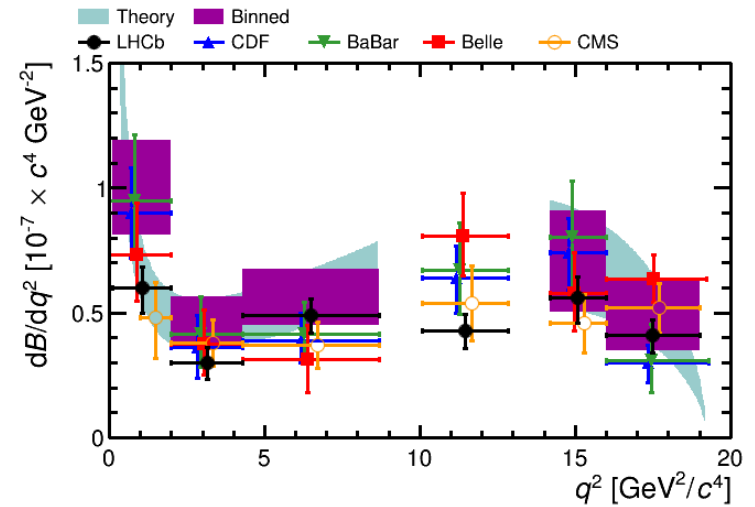
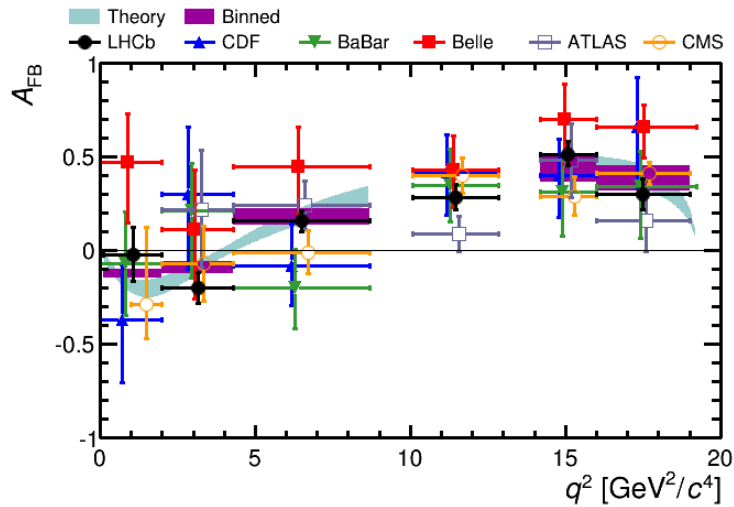
Production and detection asymmetry taken into account by using the $B^0 \rightarrow J/\psi K^*$ decay

$$A_{CP}(K^*\mu\mu) \approx A_{RAW}(K^*\mu\mu) - A_{RAW}(J/\psi K^*)$$

Kinematic difference taken into account by reweighting

Left-right detector asymmetry removed averaging magnet polarities

Comparison with other experiments

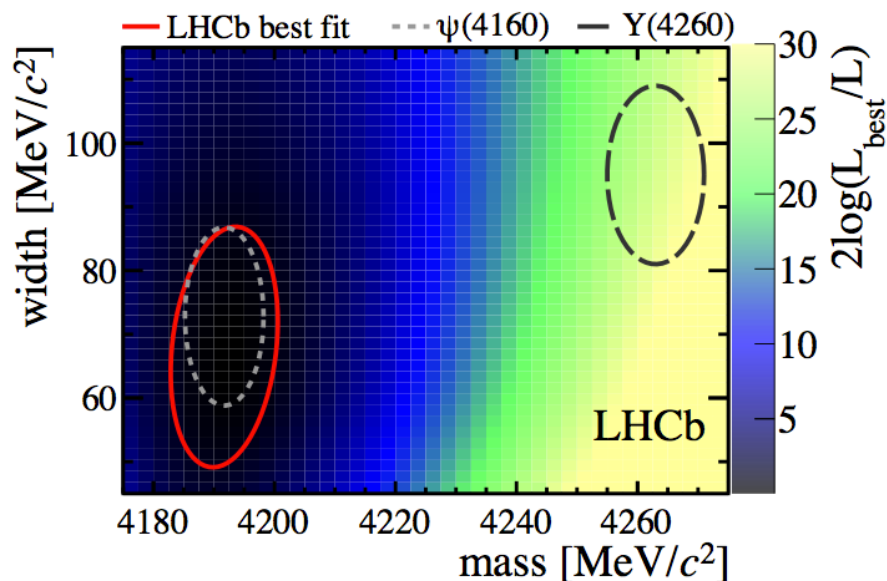


New resonance structure

$$\mathcal{P}_{\text{sig}} \propto P(m_{\mu^+\mu^-}) |\mathcal{A}|^2 f^2(m_{\mu^+\mu^-}^2),$$

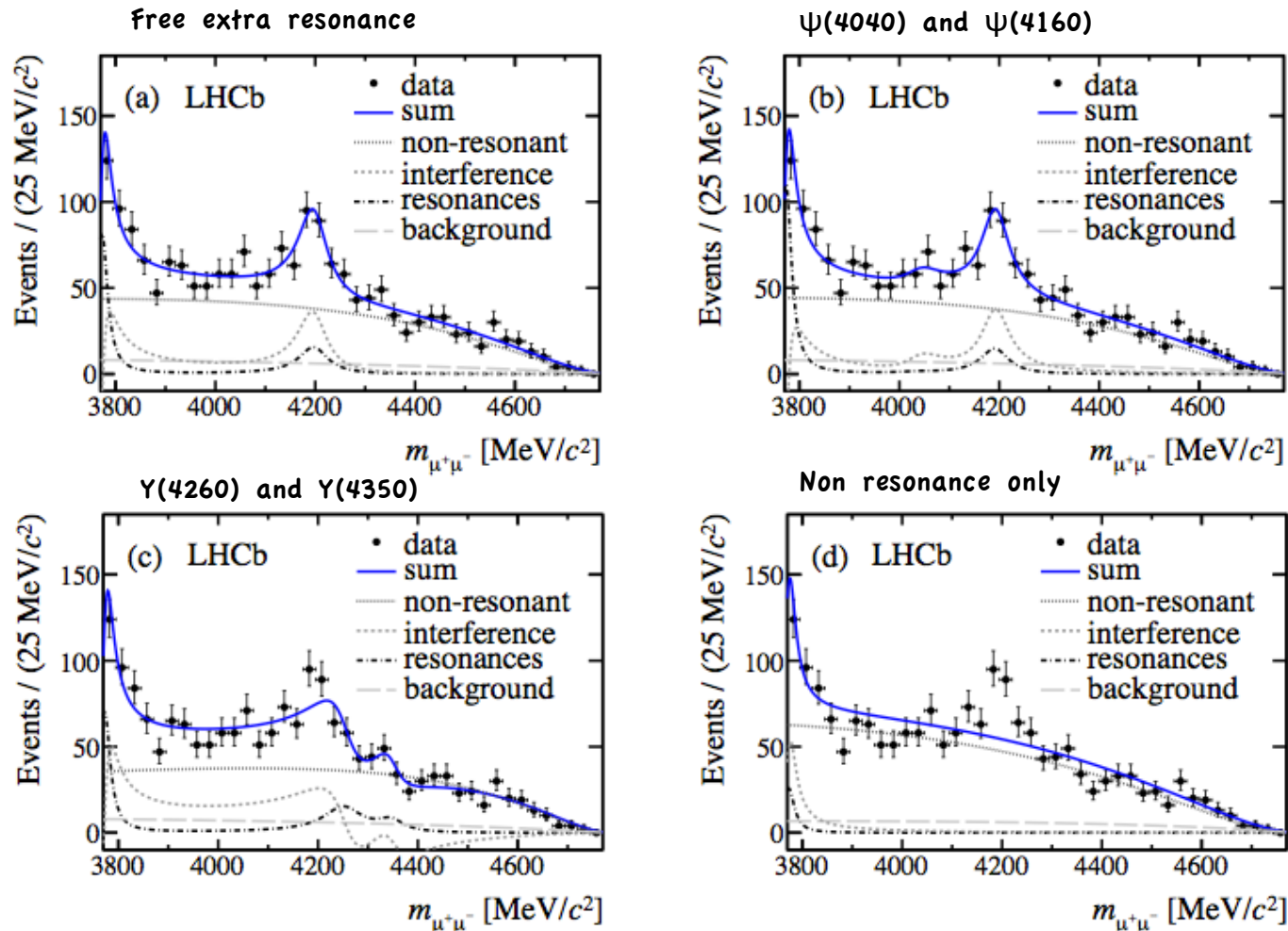
$$|\mathcal{A}|^2 = |A_{\text{nr}}^{\text{V}} + \sum_k e^{i\delta_k} A_{\text{r}}^k|^2 + |A_{\text{nr}}^{\text{AV}}|^2,$$

	Unconstrained	$\psi(4160)$
$\mathcal{B}[\times 10^{-9}]$	$3.9^{+0.7}_{-0.6}$	$3.5^{+0.9}_{-0.8}$
Mass [MeV/c ²]	4191^{+9}_{-8}	4190 ± 5
Width [MeV/c ²]	65^{+22}_{-16}	66 ± 12
Phase [rad]	-1.7 ± 0.3	-1.8 ± 0.3

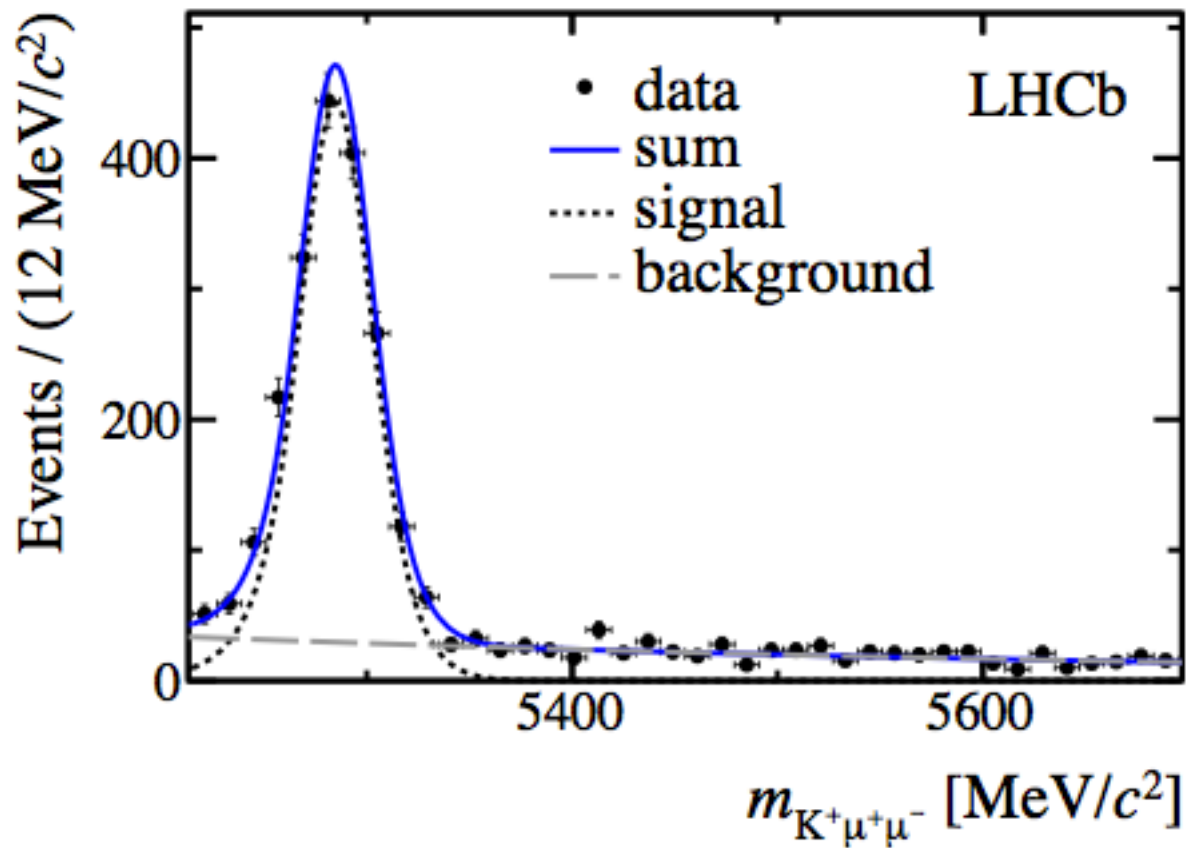


- **Significance of the new structure exceed 6 sigmas**
- **Compatible with hypothesis of $\psi(4160)$**
- **Hypothesis of $Y(4260)$ disfavored at 4 sigmas**

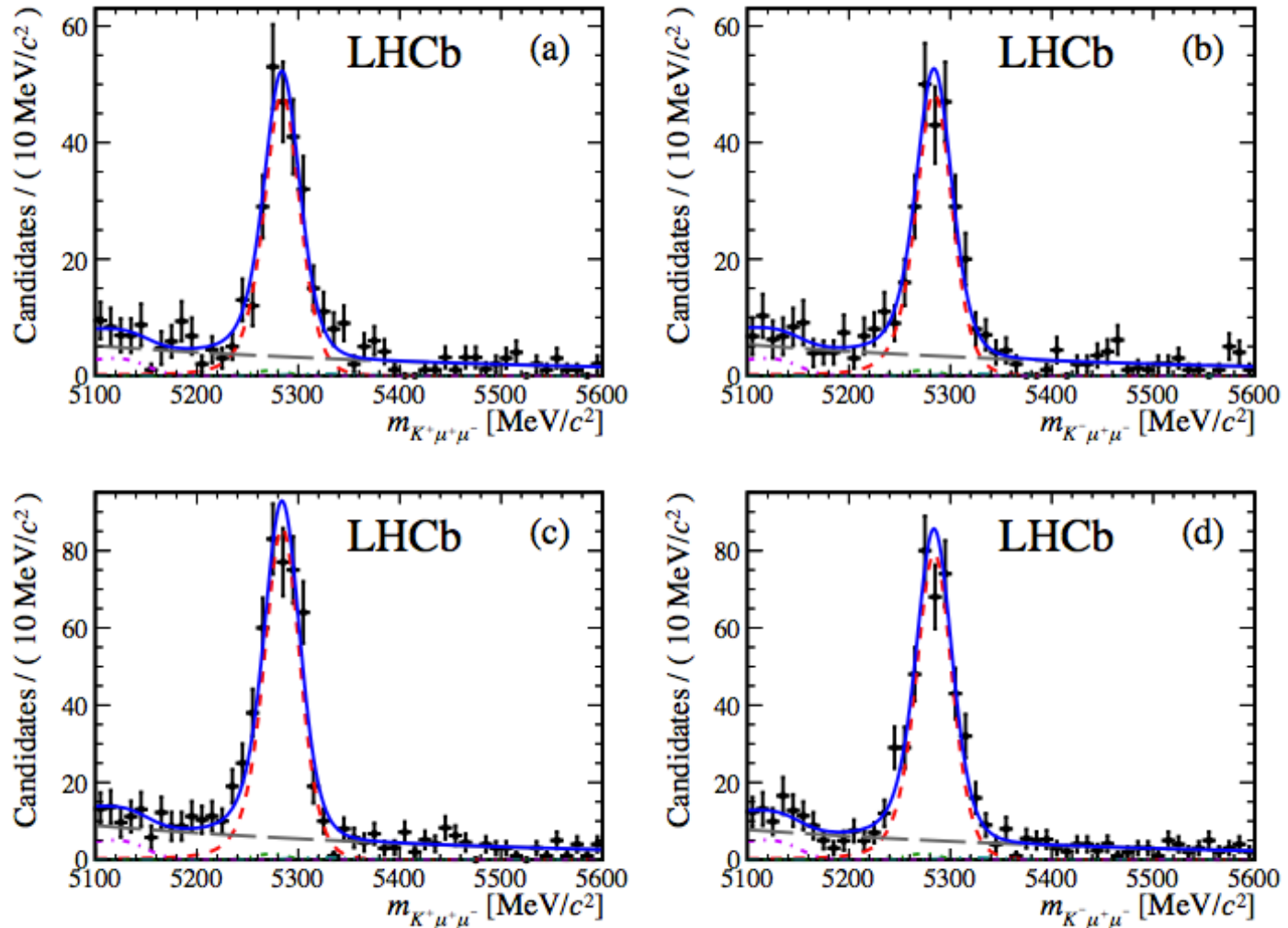
Fit with various hypotheses



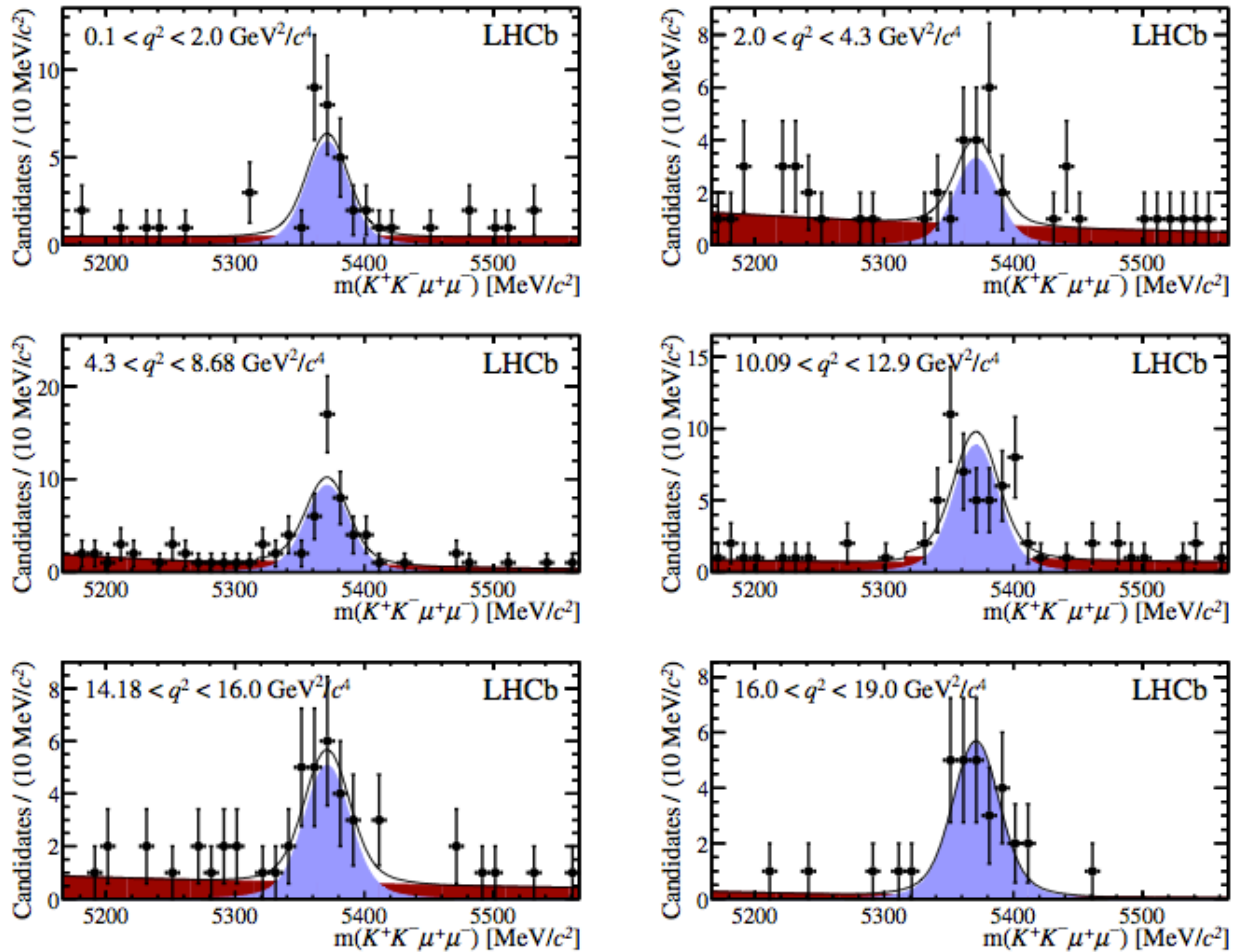
B-mass fit



$$A_{CP} (B^+ \rightarrow K^+ \mu \mu)$$



$$B_s^0 \rightarrow \phi \mu \mu$$



$$B_s^0 \rightarrow \phi \mu \mu$$

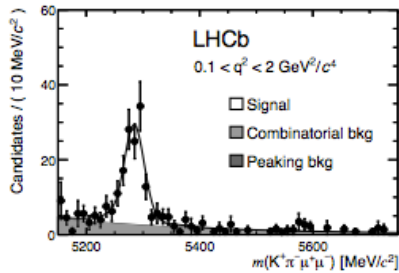
$$\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{dq^2 d\cos\theta_\ell d\cos\theta_K d\Phi} = \frac{9}{32\pi} \left[S_1^s \sin^2 \theta_K + S_1^c \cos^2 \theta_K \right. \\
+ S_2^s \sin^2 \theta_K \cos 2\theta_\ell + S_2^c \cos^2 \theta_K \cos 2\theta_\ell \\
+ S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\Phi + S_4 \sin 2\theta_K \sin 2\theta_\ell \cos \Phi \\
+ A_5 \sin 2\theta_K \sin \theta_\ell \cos \Phi + A_6 \sin^2 \theta_K \cos \theta_\ell \\
+ S_7 \sin 2\theta_K \sin \theta_\ell \sin \Phi + A_8 \sin 2\theta_K \sin 2\theta_\ell \sin \Phi \\
\left. + A_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\Phi \right]$$

Projecting over the three angles separately:

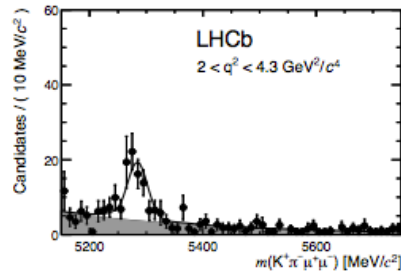
$$\frac{1}{d\Gamma/dq^2} \frac{d^2\Gamma}{dq^2 d\cos\theta_K} = \frac{3}{4}(1 - F_L)(1 - \cos^2 \theta_K) + \frac{3}{2}F_L \cos^2 \theta_K, \\
\frac{1}{d\Gamma/dq^2} \frac{d^2\Gamma}{dq^2 d\cos\theta_\ell} = \frac{3}{8}(1 - F_L)(1 + \cos^2 \theta_\ell) + \frac{3}{4}F_L(1 - \cos^2 \theta_\ell) + \frac{3}{4}A_6 \cos \theta_\ell, \\
\frac{1}{d\Gamma/dq^2} \frac{d^2\Gamma}{dq^2 d\Phi} = \frac{1}{2\pi} + \frac{1}{2\pi}S_3 \cos 2\Phi + \frac{1}{2\pi}A_9 \sin 2\Phi,$$

New observables $B^0 \rightarrow K^* \mu \mu$

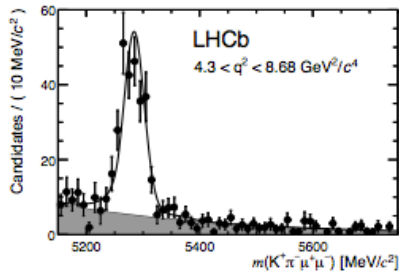
B-mass plots



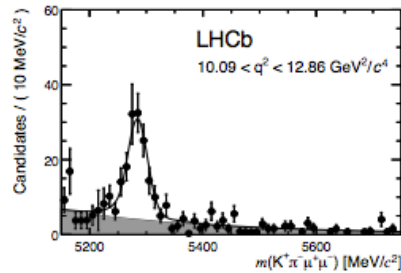
(a) $0.1 < q^2 < 2 \text{ GeV}^2/c^4$



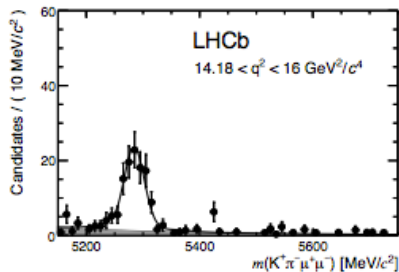
(b) $2.0 < q^2 < 4.3 \text{ GeV}^2/c^4$



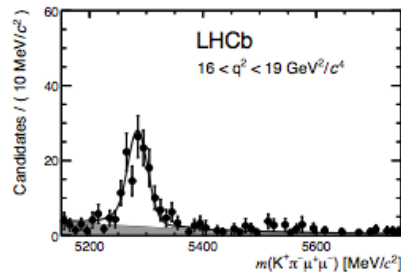
(c) $4.3 < q^2 < 8.68 \text{ GeV}^2/c^4$



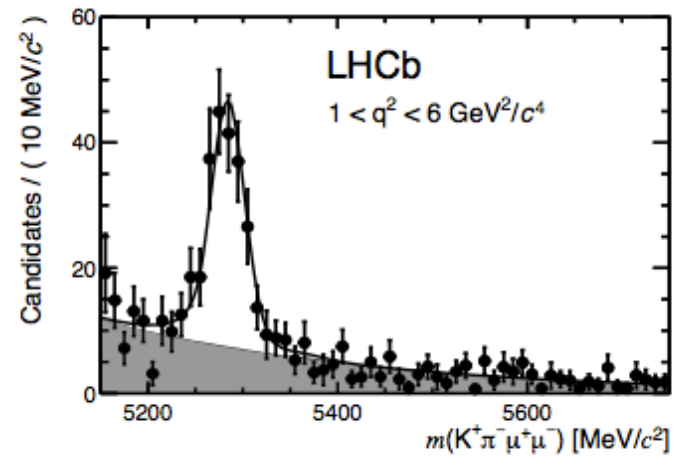
(d) $10.09 < q^2 < 12.68 \text{ GeV}^2/c^4$



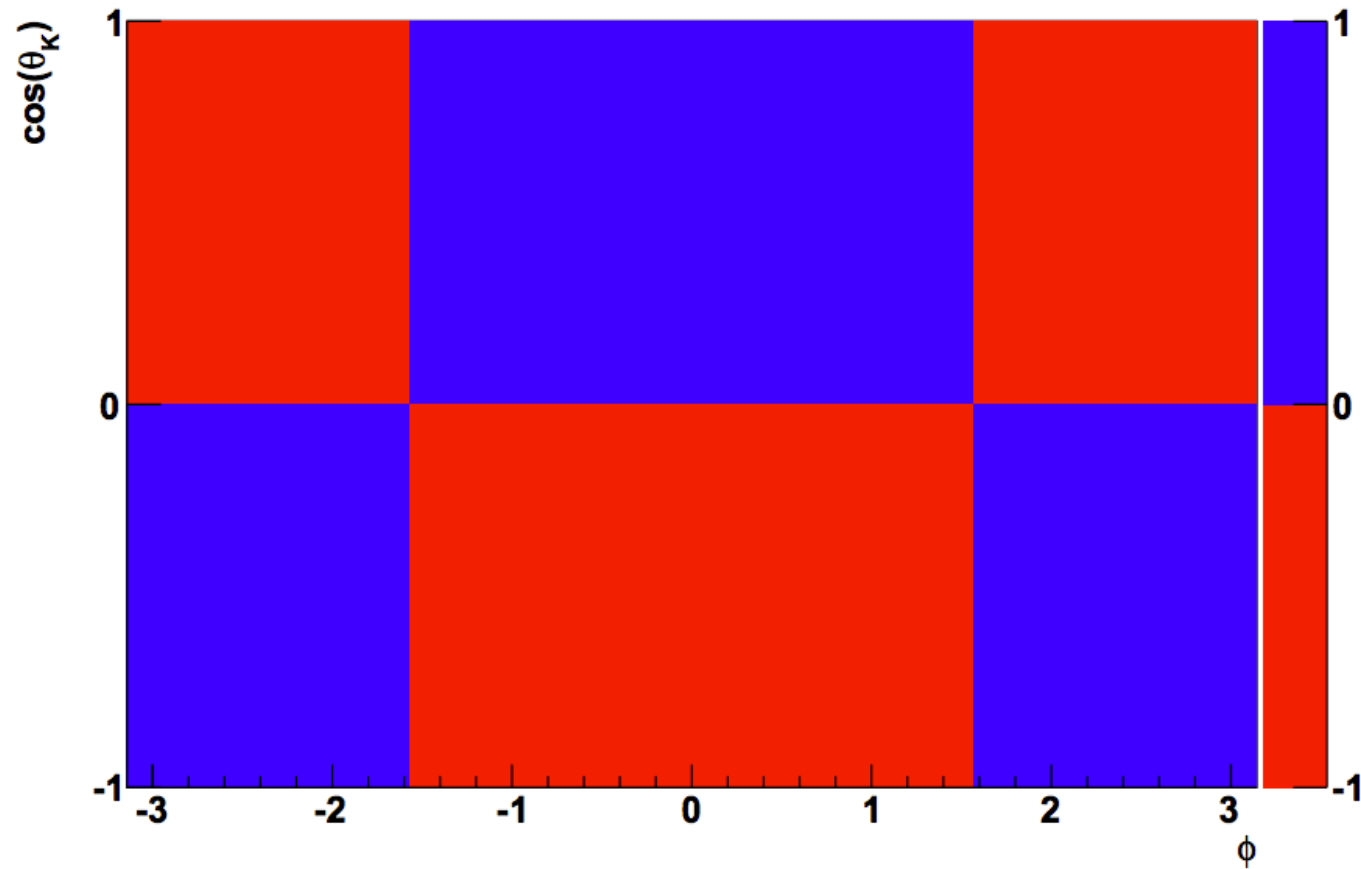
(e) $14.18 < q^2 < 16 \text{ GeV}^2/c^4$



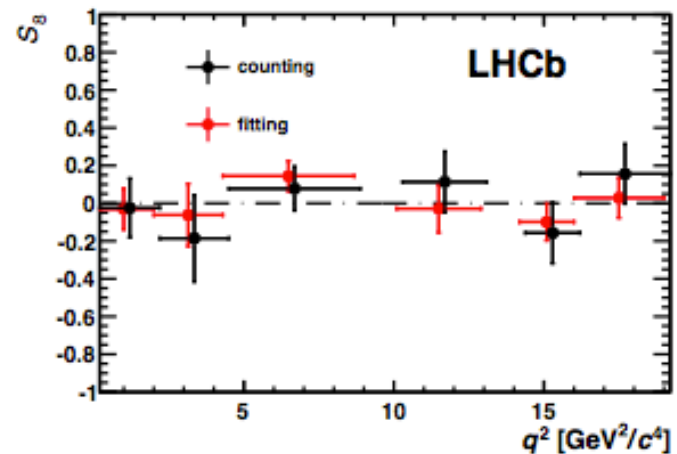
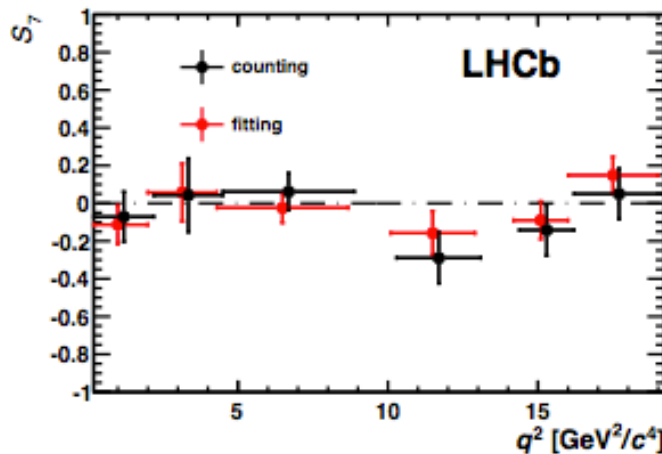
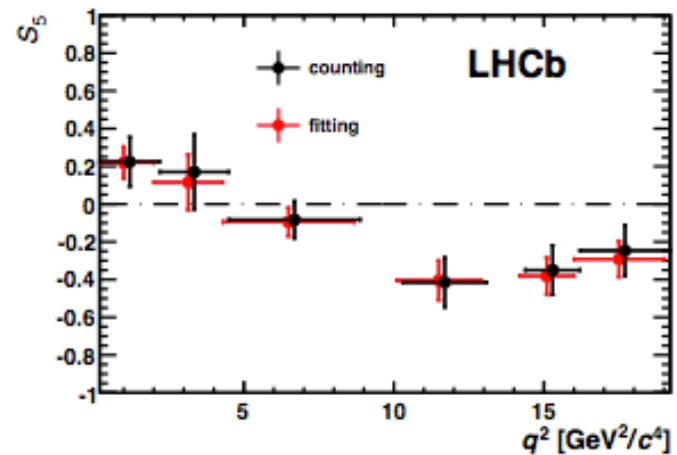
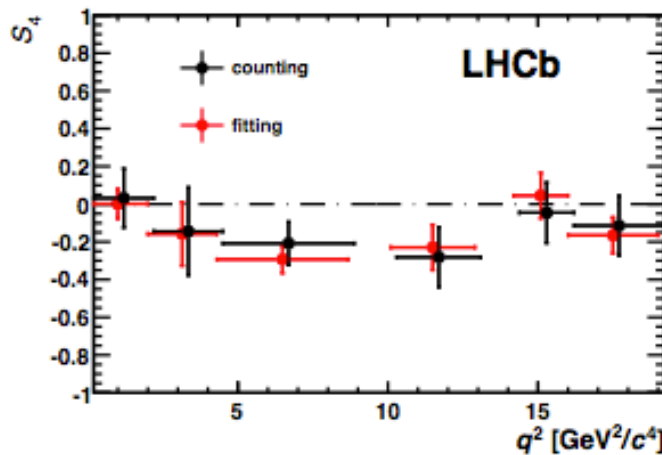
(f) $16 < q^2 < 19 \text{ GeV}^2/c^4$



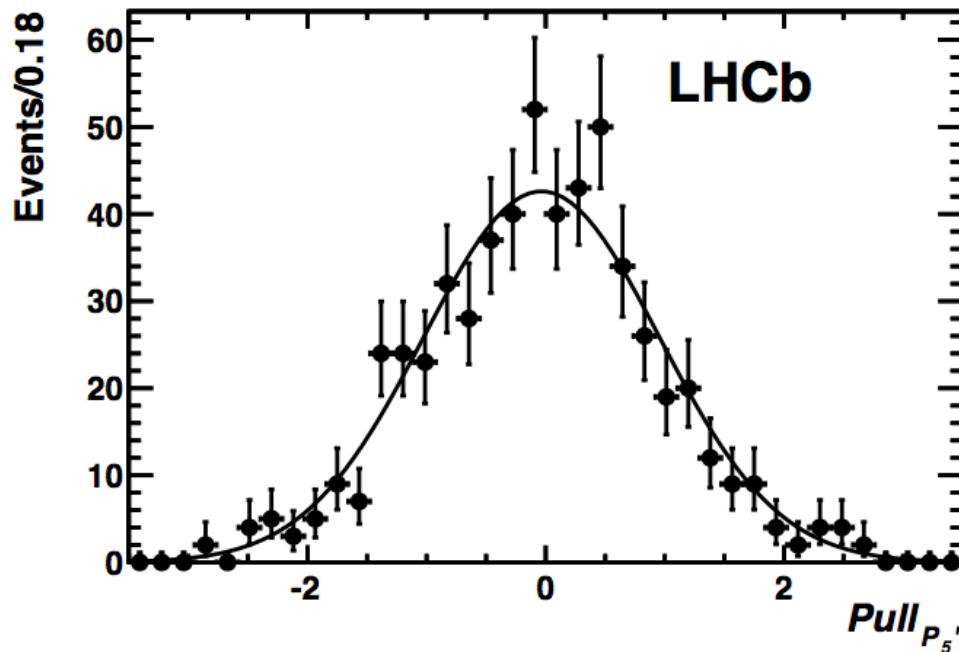
S_5 counting



Counting VS fitting



X-checks with $B^0 \rightarrow J/\psi K^*$



- Procedure repeated with the full sample of $B^0 \rightarrow J/\psi K^*$
- Procedure repeated chopping off the full sample in about 300 sub-samples

Foldings

$$P'_4, S_4: \begin{cases} \phi \rightarrow -\phi & \text{for } \phi < 0 \\ \phi \rightarrow \pi - \phi & \text{for } \theta_\ell > \pi/2 \\ \theta_\ell \rightarrow \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2, \end{cases}$$

$$P'_5, S_5: \begin{cases} \phi \rightarrow -\phi & \text{for } \phi < 0 \\ \theta_\ell \rightarrow \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2, \end{cases}$$

$$P'_6, S_7: \begin{cases} \phi \rightarrow \pi - \phi & \text{for } \phi > \pi/2 \\ \phi \rightarrow -\pi - \phi & \text{for } \phi < -\pi/2 \\ \theta_\ell \rightarrow \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2, \end{cases}$$

$$P'_8, S_8: \begin{cases} \phi \rightarrow \pi - \phi & \text{for } \phi > \pi/2 \\ \phi \rightarrow -\pi - \phi & \text{for } \phi < -\pi/2 \\ \theta_K \rightarrow \pi - \theta_K & \text{for } \theta_\ell > \pi/2 \\ \theta_\ell \rightarrow \pi - \theta_\ell & \text{for } \theta_\ell > \pi/2. \end{cases}$$

Pdfs to fit

$$Pdf = \frac{9}{8\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \sqrt{F_L(1 - F_L)} P'_5 \sin 2\theta_K \sin \theta_\ell \cos \phi \right]$$

$$Pdf = \frac{9}{8\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \sqrt{F_L(1 - F_L)} P'_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi \right]$$

$$Pdf = \frac{9}{8\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \sqrt{F_L(1 - F_L)} P'_6 \sin 2\theta_K \sin \theta_\ell \sin \phi \right]$$

$$Pdf = \frac{9}{8\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_\ell - F_L \cos^2 \theta_K \cos 2\theta_\ell + S_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + \sqrt{F_L(1 - F_L)} P'_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi \right]$$

S-wave

Differential decay rate

$$(1 - F_S)W_P + \frac{9}{32\pi} (W_S + W_{SP})$$

S-wave

$$W_S = \frac{2}{3} F_S \sin^2 \theta_\ell$$

Interference terms

$$\begin{aligned} & \frac{4}{3} A_S \sin^2 \theta_\ell \cos \theta_K + A_S^{(4)} \sin \theta_K \sin 2\theta_\ell \cos \phi + \\ & A_S^{(5)} \sin \theta_K \sin \theta_\ell \cos \phi + A_S^{(7)} \sin \theta_K \sin \theta_\ell \sin \phi \\ & + A_S^{(8)} \sin \theta_K \sin 2\theta_\ell \sin \phi . \end{aligned} \quad (9)$$