



New and conventional bottomonium states



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FPCP2012

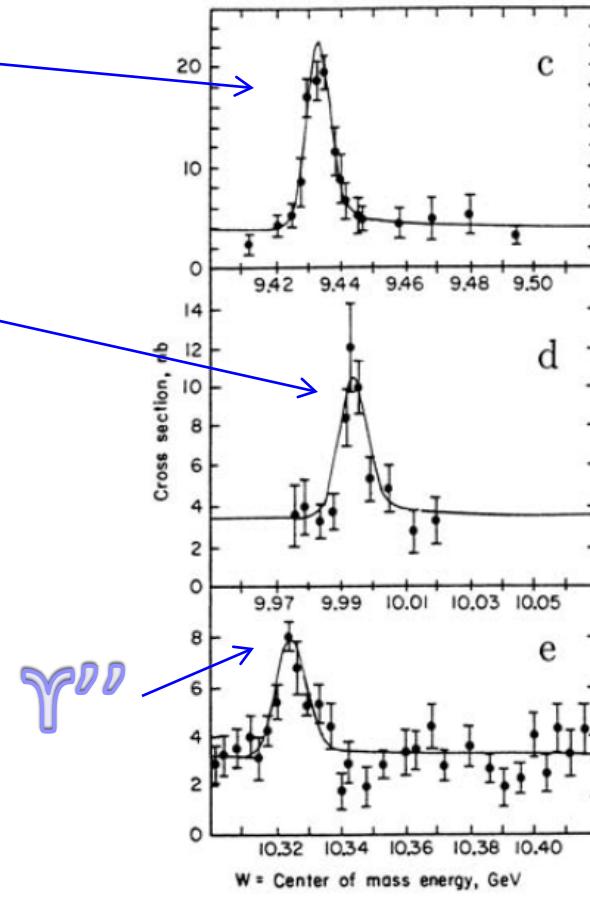
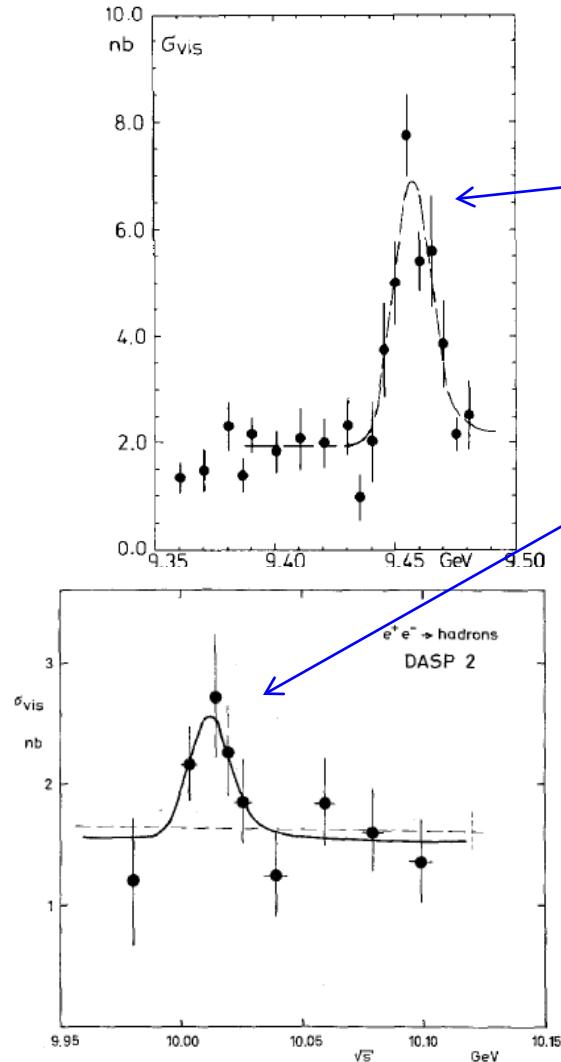


γ discovery in 1978

DORIS ring at DESY

γ, γ' discovered in e^+e^- annihilation.

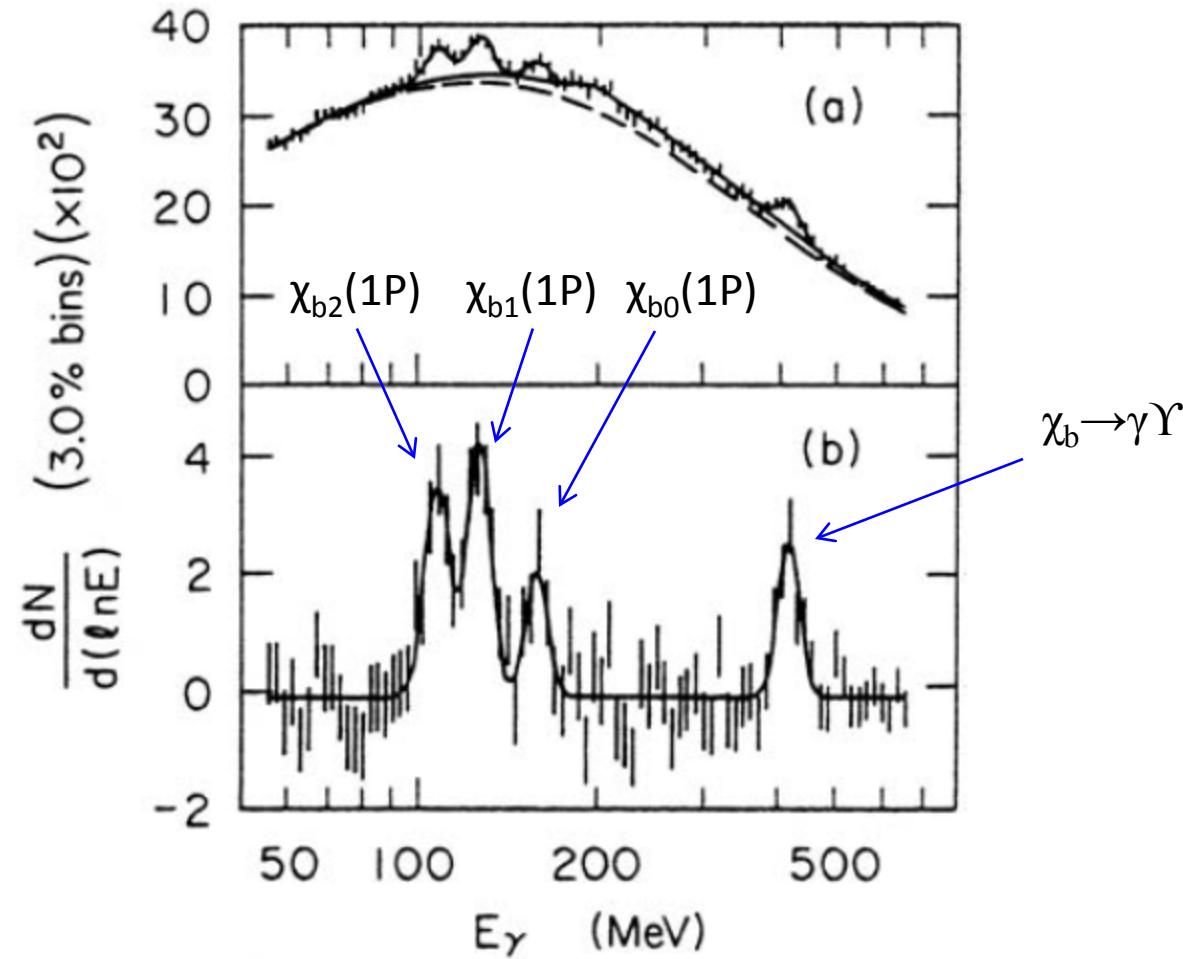
CLEO group at CESR



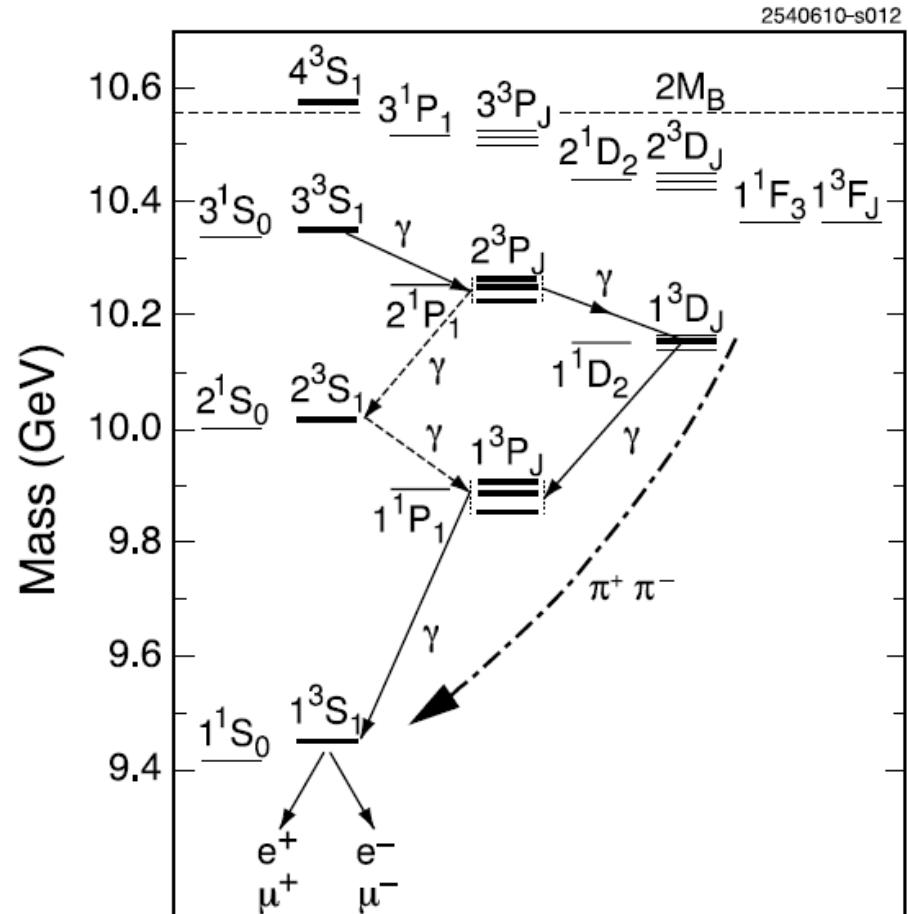
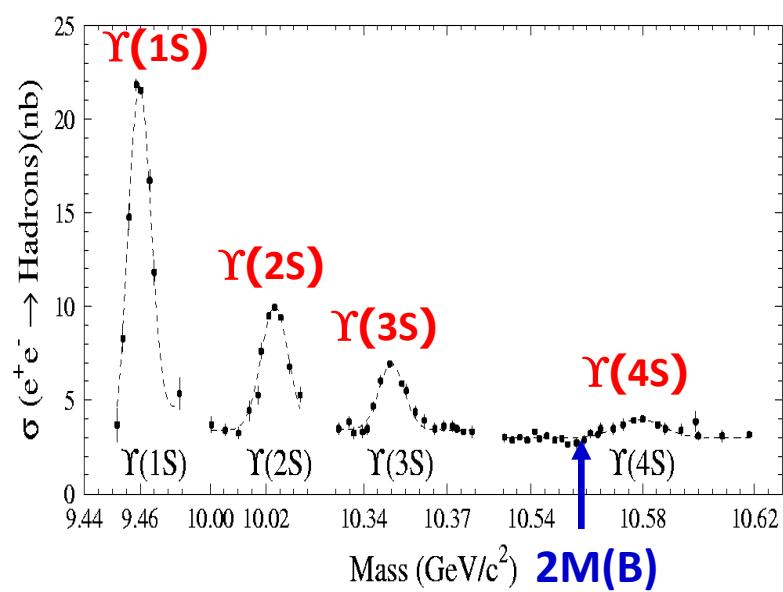
$\chi_b(^3P_{2,1,0})$ observation in 1985

Photon spectrum from Υ' decays.

Crystal Ball Collaboration at DORIS II.



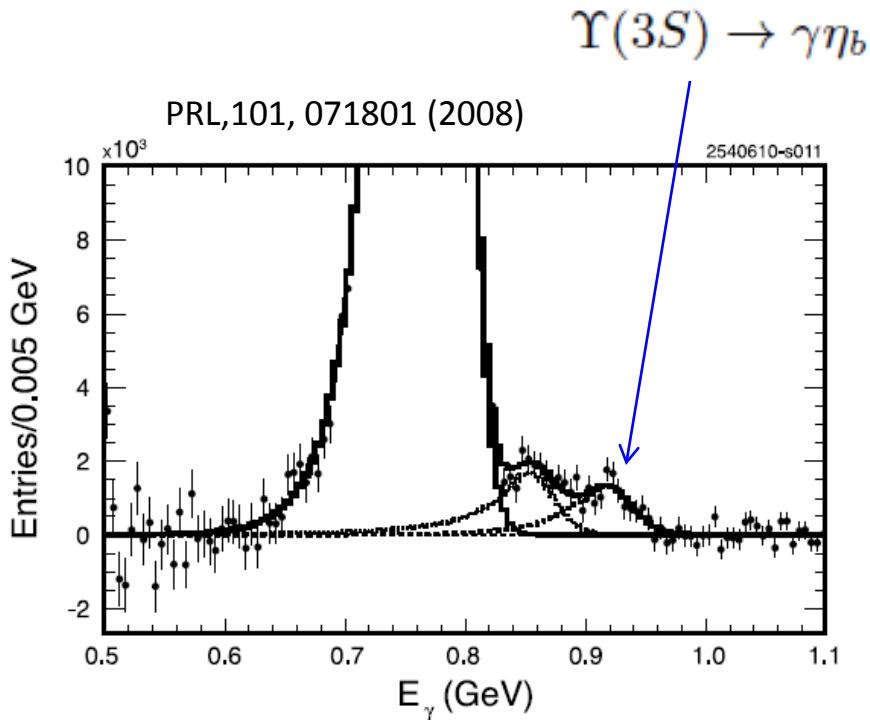
The bottomonium family



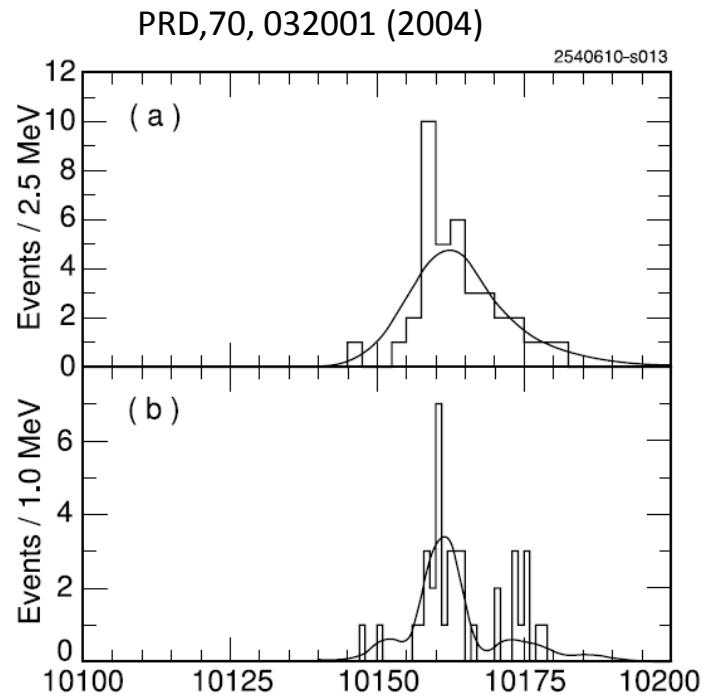
BaBar, CLEO collected large $\Upsilon(3S)$ data because its rich decay trees.

Old strategy: γ lineshape

BaBar's η_b discovery

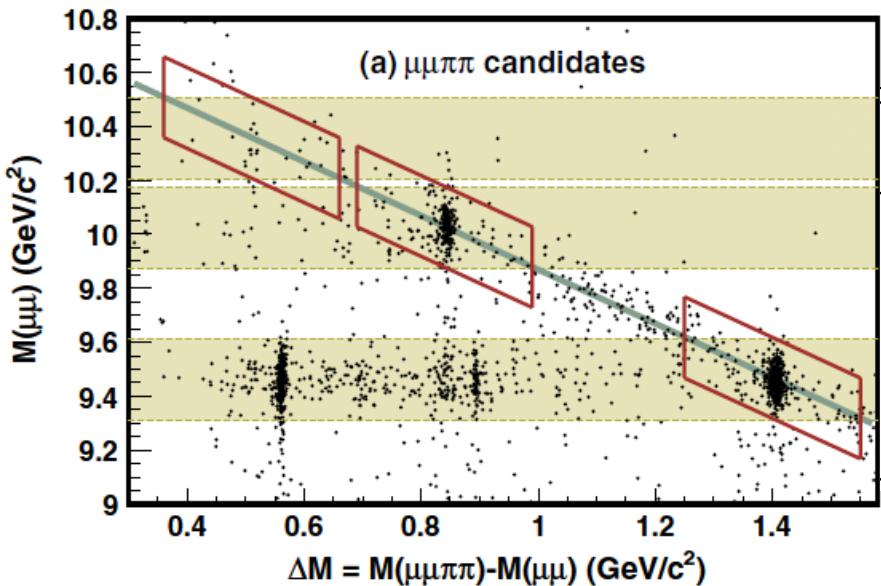


CLEO's $\Upsilon(1D)$ discovery



Recoil mass of two soft γ from the
 $\Upsilon(3S) \rightarrow \gamma\gamma\gamma l^+l^-$ events.

$\Upsilon(5S)$ is just different



$e^+e^- \rightarrow \Upsilon(1S/2S/3S) \pi^+\pi^-$

Process	$\Gamma_{\Upsilon(1S)\pi^+\pi^-}$
$\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.0060 MeV
$\Upsilon(3S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.0009 MeV
$\Upsilon(4S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.0019 MeV
$\Upsilon(10860) \rightarrow \Upsilon(1S)\pi^+\pi^-$	0.59 MeV

$\sim \times 100$



$\Upsilon(5S)$ result: (21.7 fb⁻¹)
 Belle, K.F. Chen et al.
 PRL 100, 112001 (2008)

Process	σ (pb)	Γ (MeV)
$\Upsilon(1S)\pi^+\pi^-$	$1.61 \pm 0.10 \pm 0.12$	$0.59 \pm 0.04 \pm 0.09$
$\Upsilon(2S)\pi^+\pi^-$	$2.35 \pm 0.19 \pm 0.32$	$0.85 \pm 0.07 \pm 0.16$
$\Upsilon(3S)\pi^+\pi^-$	$1.44^{+0.55}_{-0.45} \pm 0.19$	$0.52^{+0.20}_{-0.17} \pm 0.10$
$\Upsilon(1S)K^+K^-$	$0.185^{+0.048}_{-0.041} \pm 0.028$	$0.067^{+0.017}_{-0.015} \pm 0.013$

So Belle took large data at $\Upsilon(5S)$ resonance. Now 121.4 fb⁻¹ total.

The $\pi^+\pi^-$ transition of $\Upsilon(5S)$

Missing mass technique



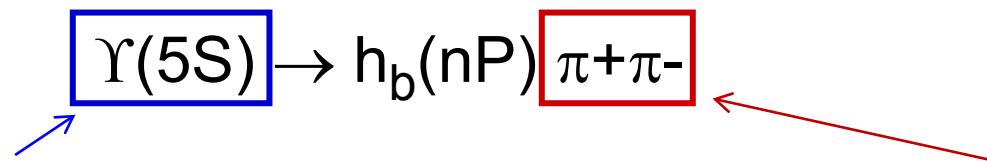
Good for study of unknown particles

Full reconstruction, when the decays of daughter particles are known.

$$\Upsilon(5S) \rightarrow \Upsilon(1S,2S,3S) \pi^+\pi^-$$

$$\Upsilon(nS) \rightarrow \mu^+ \mu^-$$

Missing mass technique:



$P(5S)$ from accelerator information

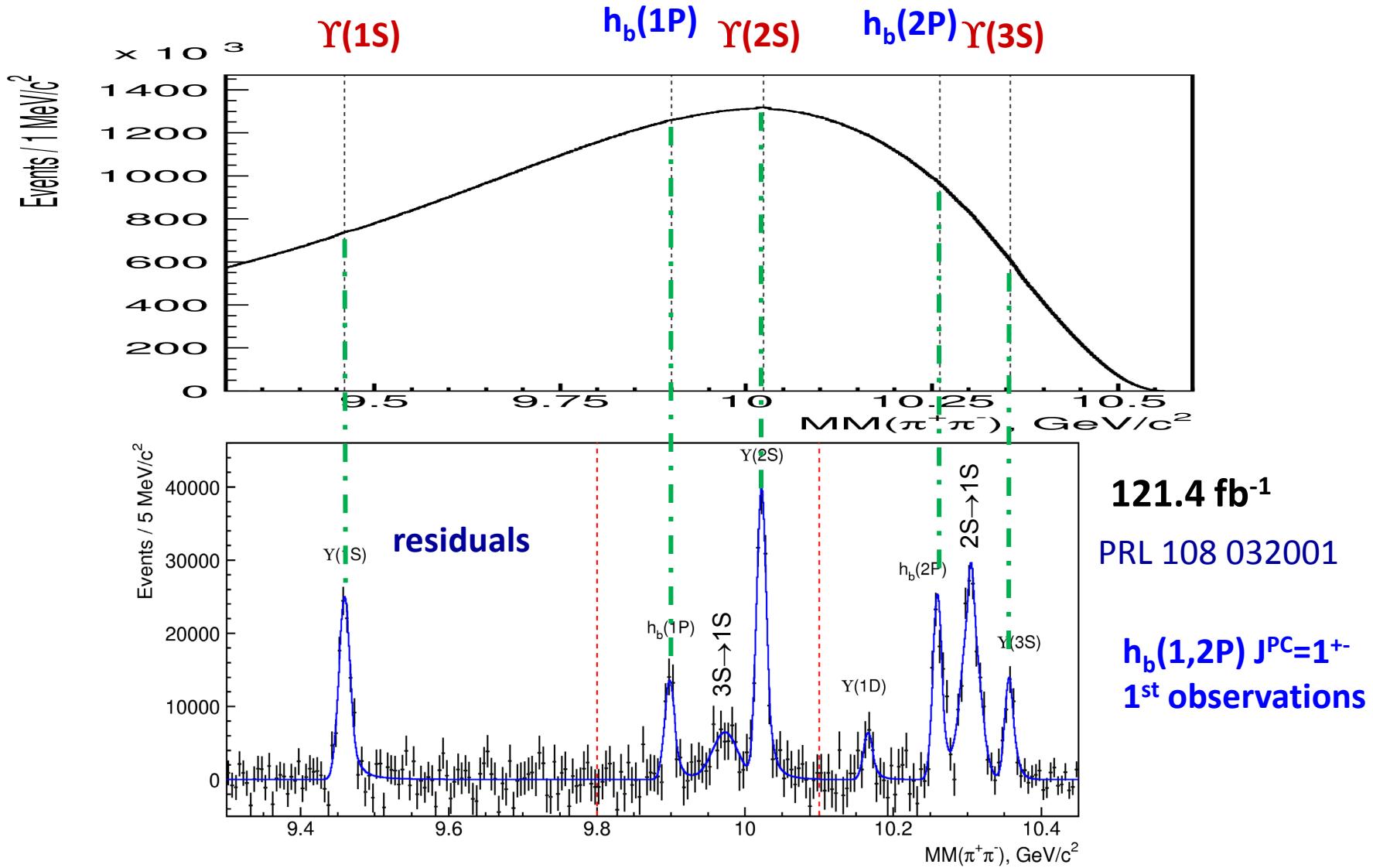
$P(\pi\pi)$ is reconstructed.

$$M(h_b) = \sqrt{(P(5S) - P(\pi^+\pi^-))^2} \equiv \textcolor{red}{MM}(\pi^+\pi^-)$$

missing mass of $\pi^+\pi^-$

Since we don't need to reconstruct h_b decays...

Results of $\pi^+\pi^-$ missing mass study



$h_b(1,2P)$ mass

$(b\bar{b})$: $S=0$ $L=1$ $J^{PC}=1^{+-}$

Expected mass: center-of-gravity
 $\approx (M\chi_{b0} + 3 M\chi_{b1} + 5 M\chi_{b2}) / 9$

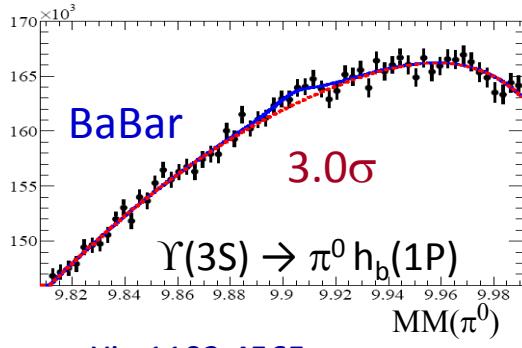
$\Delta M_{HF} \Rightarrow$ test of hyperfine interaction

Deviations from expected masses

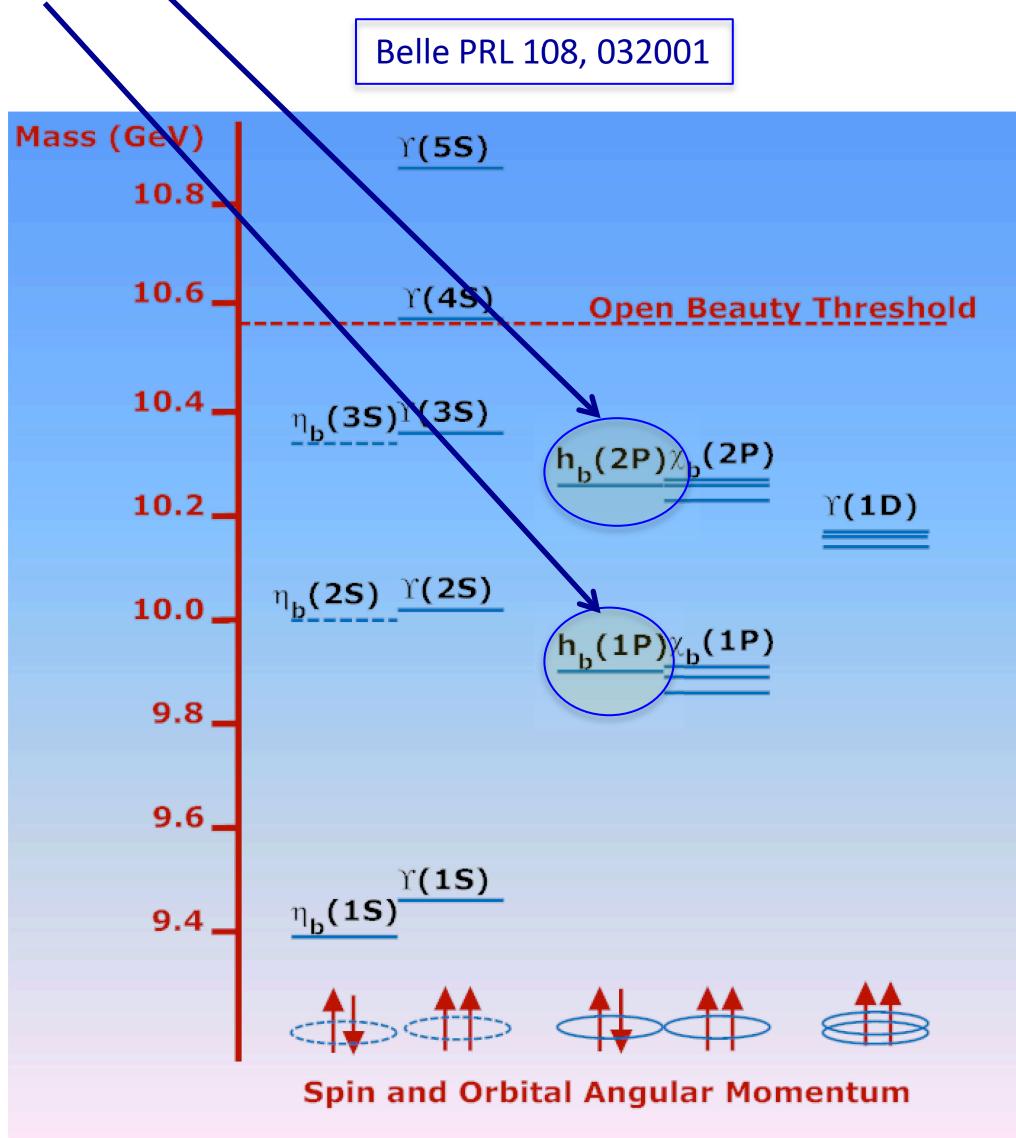
$$\left. \begin{array}{ll} h_b(1P) & (1.7 \pm 1.5) \text{ MeV}/c^2 \\ h_b(2P) & (0.5^{+1.6}_{-1.2}) \text{ MeV}/c^2 \end{array} \right\}$$

Agrees with expectations

Previous search



arXiv:1102.4565



Large $\Upsilon(5S) \rightarrow h_b(nP) \pi^+ \pi^-$ rate

$S=1$ for $\Upsilon(5S)$; $S=0$ for h_b

$$\frac{\Gamma(\Upsilon(5S) \rightarrow h_b(nP) \pi^+ \pi^-)}{\Gamma(\Upsilon(5S) \rightarrow \Upsilon(2S) \pi^+ \pi^-)} = \begin{cases} 0.407 \pm 0.079_{-0.076}^{+0.043} & h_b(1P) \\ 0.78 \pm 0.09_{-0.10}^{+0.22} & h_b(2P) \end{cases}$$

↑
No spin-flip

Spin-flip

Process with spin-flip of heavy quark is not suppressed

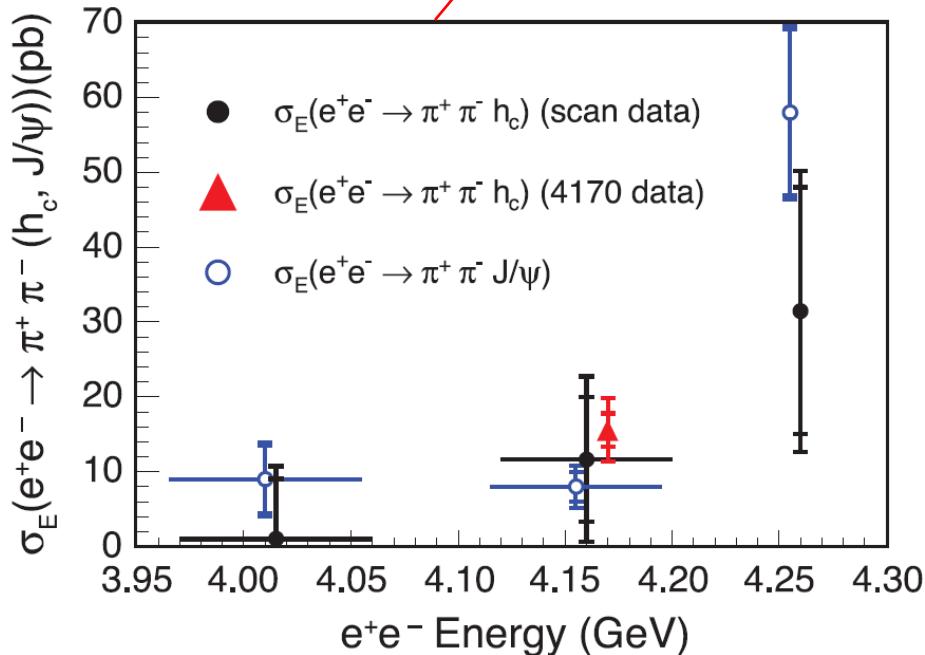
⇒ Mechanism of $\Upsilon(5S) \rightarrow h_b(nP) \pi^+ \pi^-$ decay is exotic

CLEO-c's $e^+e^- \rightarrow \pi^+\pi^- h_c$ cross section

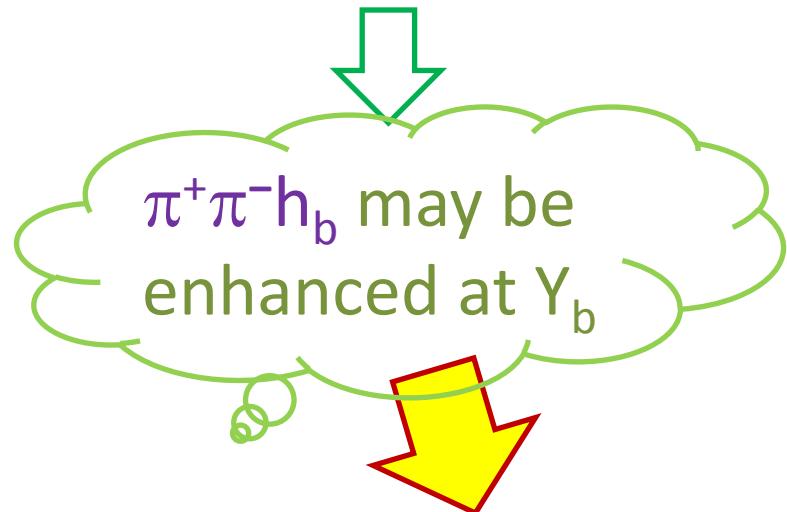
An interesting comparison!

$\bullet \sigma(h_c \pi^+\pi^-) \cong \sigma(J/\psi \pi^+\pi^-)$
 $\bullet \Gamma(Y(4260) \rightarrow J/\psi \pi^+\pi^-) > 508 \text{ keV@ 90\% C.L.}$
⇒ Large!! X. H. Mo et al., Phys. Lett. B 640 (2006) 182

PRL107, 041803 (2011)
arXiv:1102.3424



$\pi^+\pi^- h_c$ production is enhanced at $Y(4260)$



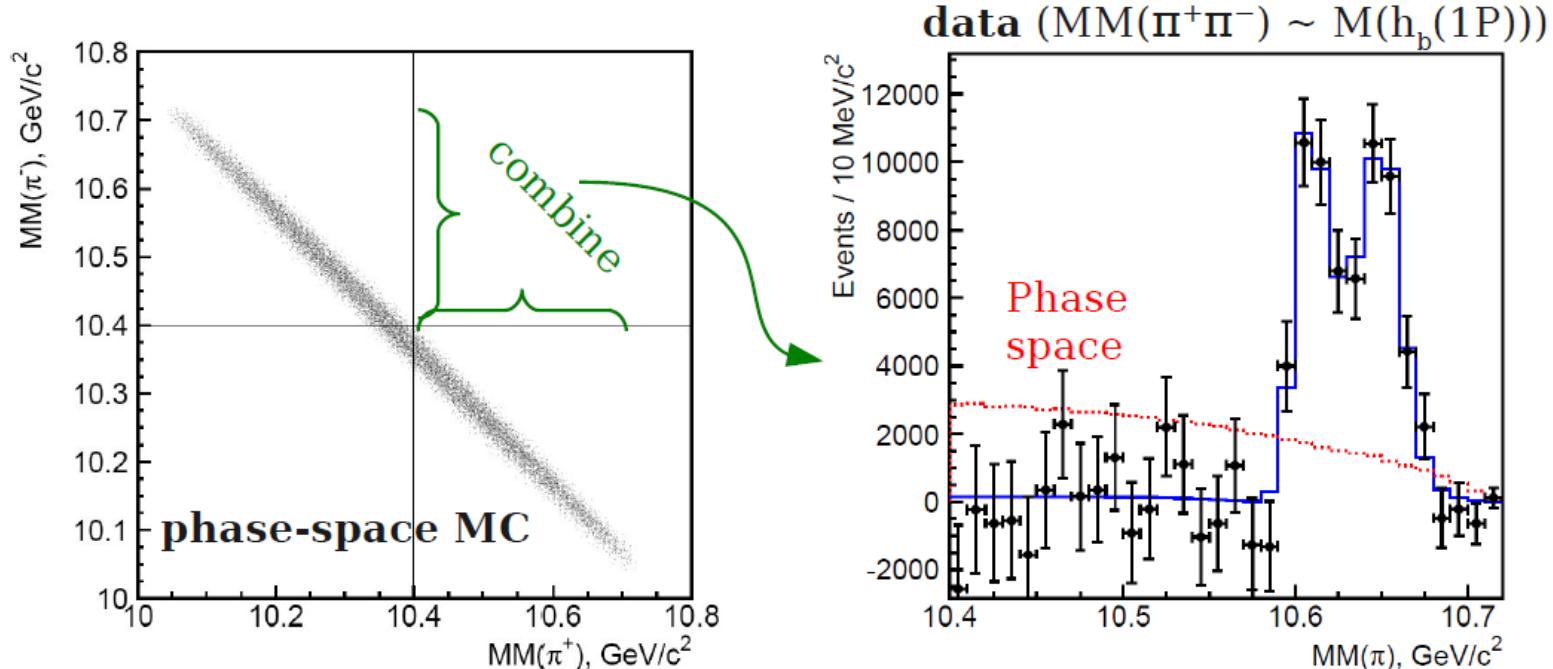
Substructure of $\pi^+\pi^- h_c$ in $Y(4260)$ has **not** been studied.

Substructure of $\pi^+\pi^- h_b$: Enough statistics in Belle's $Y(5S)$ data.

Original Motivation
for h_b search

Resonant substructure of $h_b(nP)\pi^+\pi^-$

Now look at the missing mass of a single pion π^- , then $MM(\pi^-) = M(h_b \pi^+)$



Fit Function: $|BW(s, M_1, \Gamma_1) + ae^{i\phi} BW(s, M_2, \Gamma_2) + be^{i\psi}|^2 \frac{qp}{\sqrt{s}}$

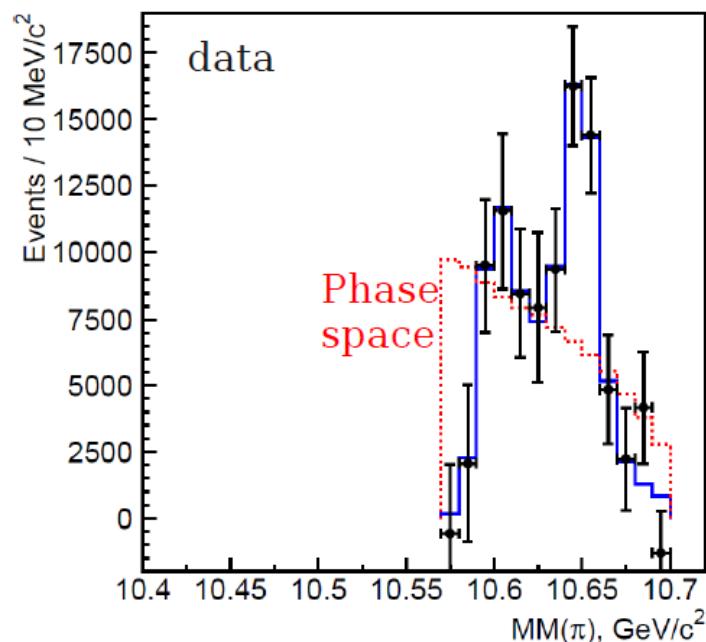
Significances with systematic:
 1 resonance v.s. 0: **6.6σ**
 2 resonances v.s. 0: **16σ**

Similar analysis for $h_b(2P)\pi^+\pi^-$

Significances: For $h_b(2P)$ yields in MM(π^-) bins, the allowed phase space is smaller.

1 resonance v.s. 0: **1.9 σ**

2 resonances v.s. 0: **4.7 σ**



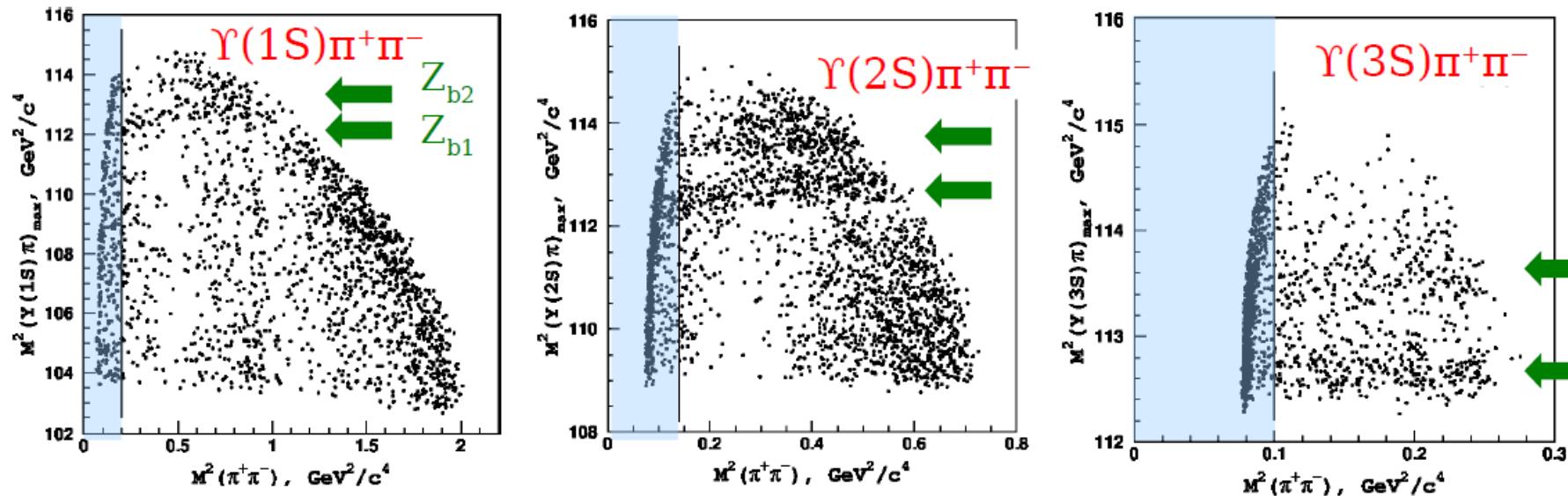
	$h_b(1P)\pi^+\pi^-$	$h_b(2P)\pi^+\pi^-$
$M_{Z_{b1}}$	$10605.1 \pm 2.2^{+3.0}_{-1.0}$	$10596 \pm 7^{+5}_{-2}$
$\Gamma_{Z_{b1}}$	$11.4^{+4.5+2.1}_{-3.9-1.2}$	16^{+16+13}_{-10-4}
$M_{Z_{b2}}$	$10654.5 \pm 2.5^{+1.0}_{-1.9}$	$10651 \pm 4 \pm 2$
$\Gamma_{Z_{b2}}$	$11.4^{+5.4+2.1}_{-4.7-5.7}$	12^{+11+8}_{-9-2}
a	$1.8^{+1.0+0.1}_{-0.7-0.5}$	$1.3^{+3.1+0.4}_{-1.1-0.7}$
ϕ	188^{+44+4}_{-58-9}	$255^{+56+12}_{-72-183}$
NR	~ 0	~ 0

Mass/width in MeV(/c²), phase in degrees

- Good agreement between $h_b(1P)$ and $h_b(2P)$.
- Non resonant part ~ 0 : Nearly all produced through

$$\Upsilon(5S) \rightarrow Z^\pm \pi^\mp \rightarrow h_b(nP)\pi^\pm \pi^\mp$$

Dalitz Plots of $\Upsilon(5S) \rightarrow \Upsilon(nS) \pi^+ \pi^-$



Exclude high background from photon conversions

$$S(s_1, s_2) = |A_{Z_{b1}} + A_{Z_{b2}} + A_{\text{NR}} + A_{f_0(980)} + A_{f_2(1270)}|^2$$

$$s_i = M_{\Upsilon\pi_i}^2$$

Isospin symmetry

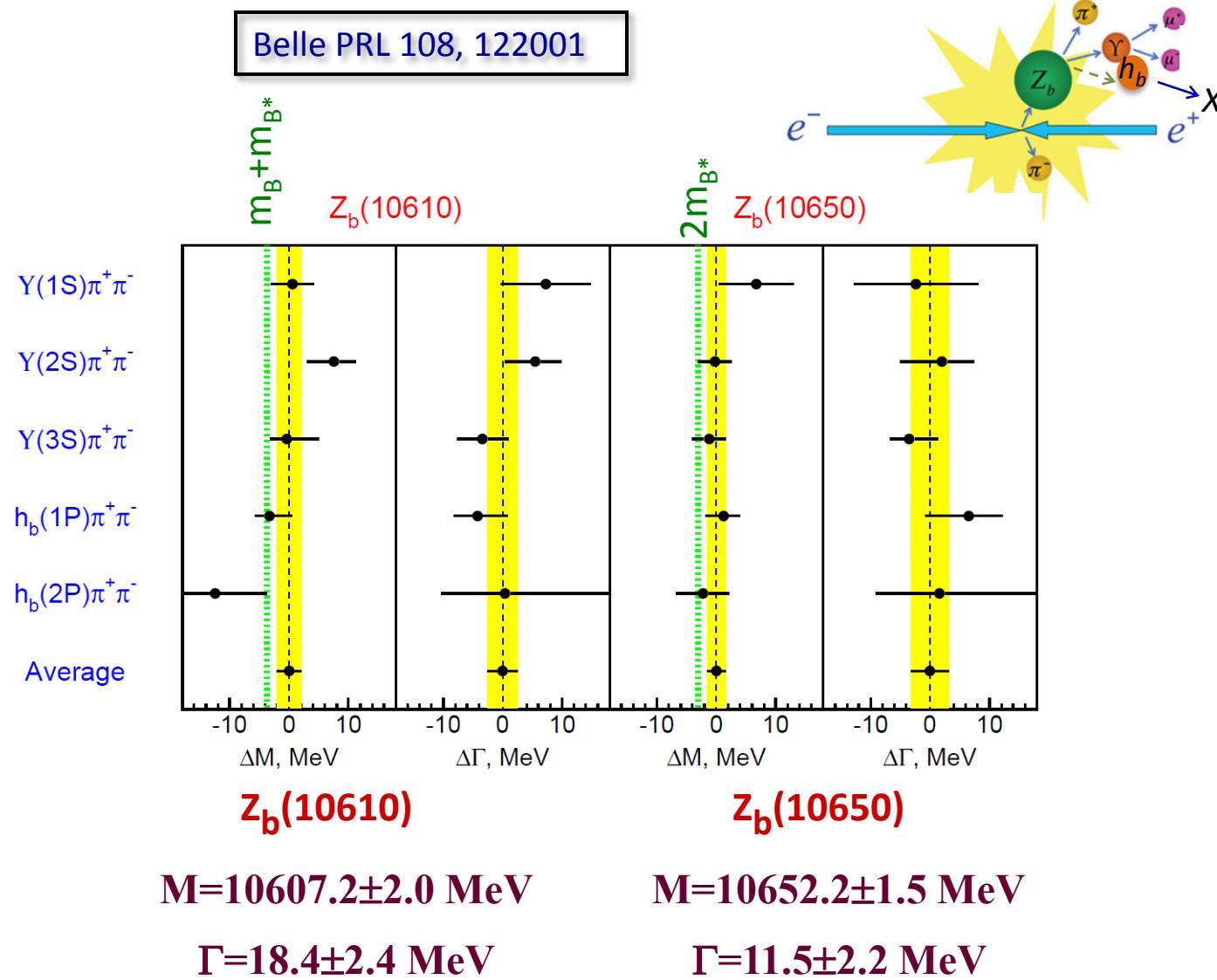
$$A_{Z_{bk}} = a_k e^{i\phi_k} \left(\frac{\sqrt{M_k \Gamma_k}}{M_k^2 - s_1 + iM_k \Gamma_k} + \frac{\sqrt{M_k \Gamma_k}}{M_k^2 - s_2 + iM_k \Gamma_k} \right)$$

$$A_{\text{NR}} = c_1 + c_2 M_{\pi\pi}^2$$

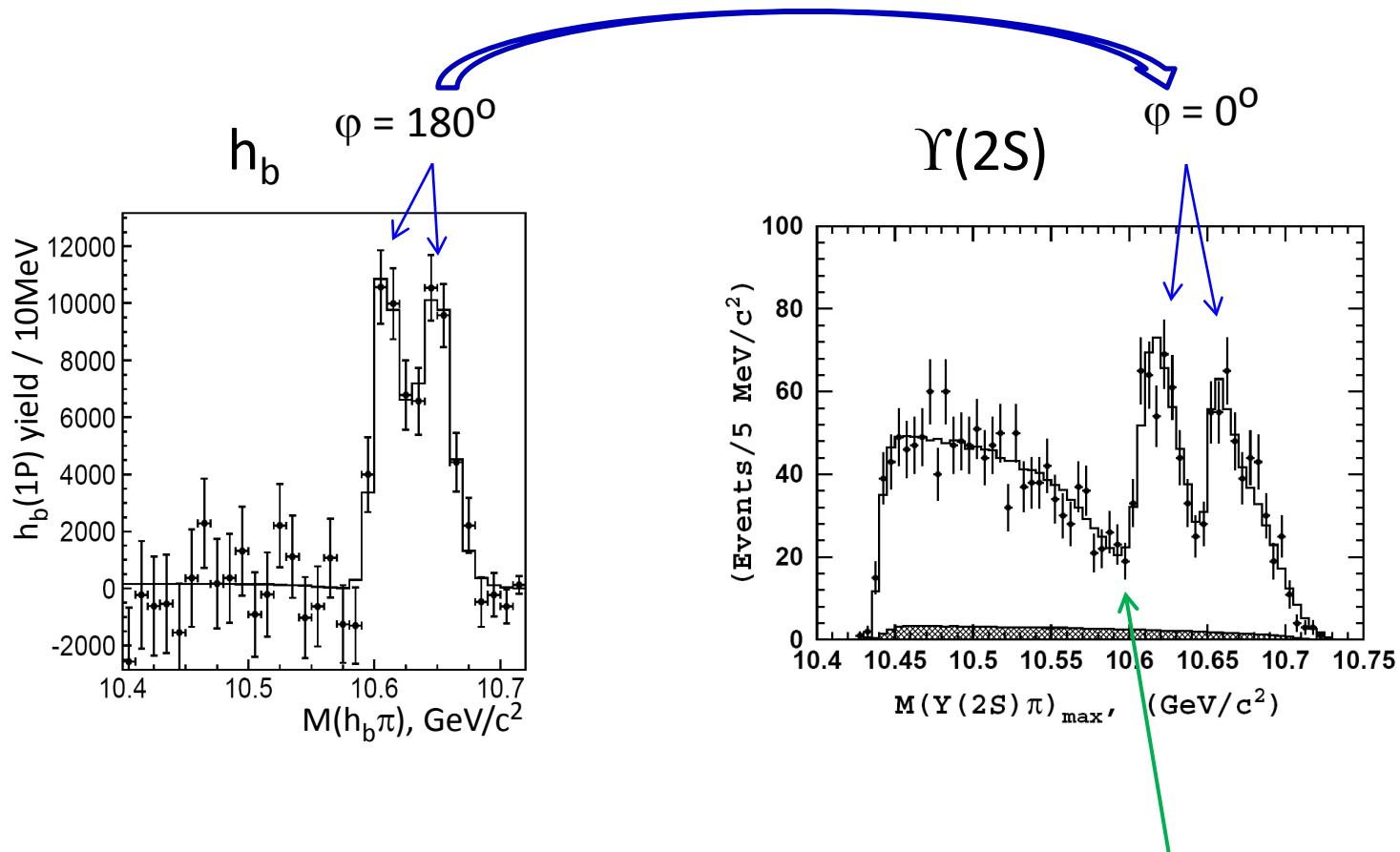
A. Voloshin, PRD74,054022 (2006)
Prog. Part. Nucl. Phys. 61,455 (2008)

Flatté D-Wave
 Breit-Wigner

Summary of parameter measurements



Projections of $\Upsilon(5S) \rightarrow [\Upsilon(nS), h_b] \pi^+ \pi^-$



Dip due to destructive interference with non-resonant amplitude.
 \Rightarrow Information on J^P of Z_b ?

Introduction to η_b

h_b has large expected radiative η_b decay

$h_b(1P) \rightarrow ggg$ (57%), $\eta_b(1S)\gamma$ (41%), γgg (2%)

Godfrey & Rosner, PRD66 014012 (2002)

$h_b(2P) \rightarrow ggg$ (63%), $\eta_b(1S)\gamma$ (13%), $\eta_b(2S)\gamma$ (19%), γgg (2%)

Large $h_b(mP)$ samples give opportunity to study
 $\eta_b(nS)$ states.

Reminder: Experimental status of η_b

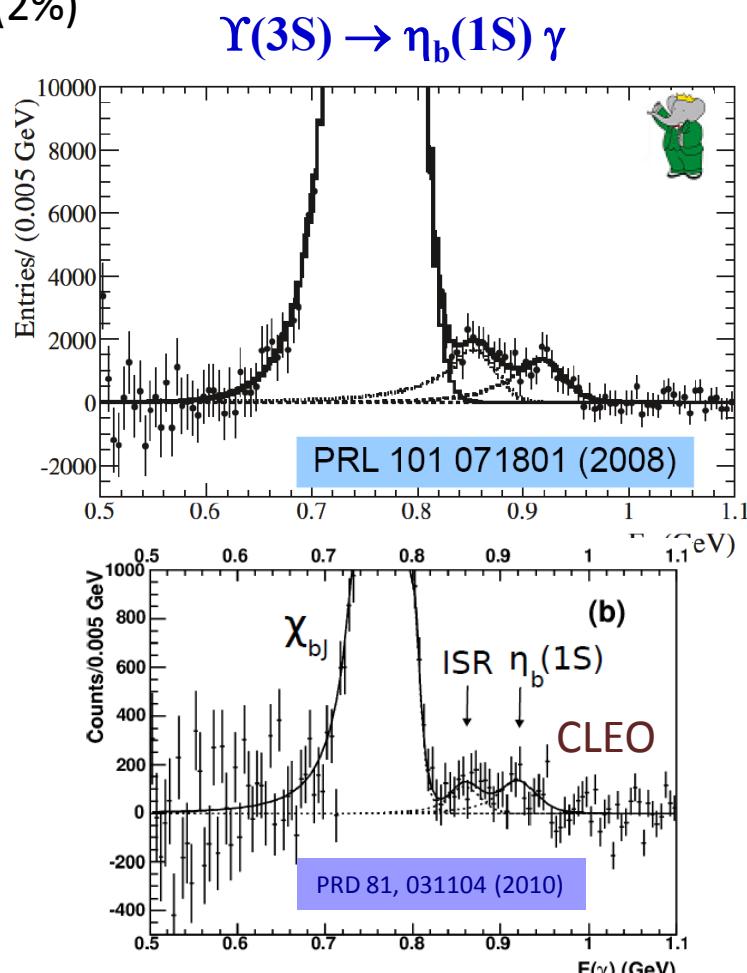
$$M[\eta_b(1S)] = 9390.9 \pm 2.8 \text{ MeV} \text{ (BaBar + CLEO)}$$

$$M[\Upsilon(1S)] - M[\eta_b(1S)] = 69.3 \pm 2.8 \text{ MeV}$$

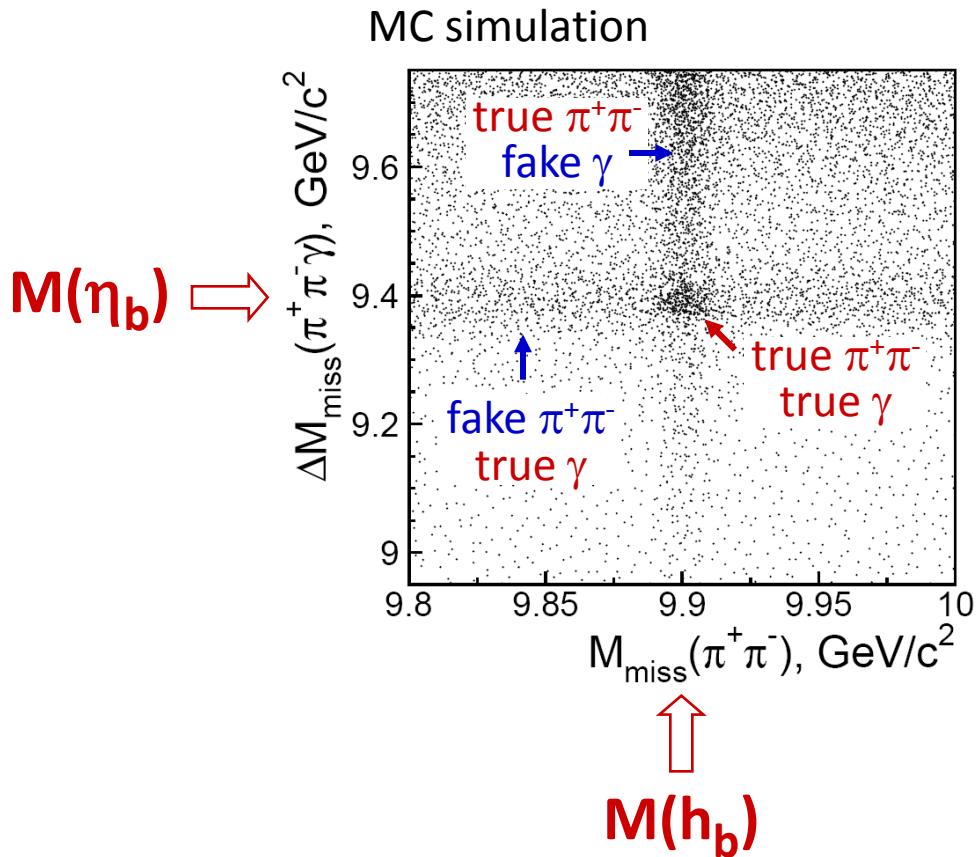
pNRQCD: 41 ± 14 MeV Lattice: 60 ± 8 MeV

Kniehl et al., PRL92,242001(2004) Meinel, PRD82,114502(2010)

Width of $\eta_b(1S)$: no information



Method of signal extraction



$$\Delta M_{\text{miss}}(\pi^+\pi^-\gamma) \equiv M_{\text{miss}}(\pi^+\pi^-\gamma) - M_{\text{miss}}(\pi^+\pi^-) + M[h_b]$$

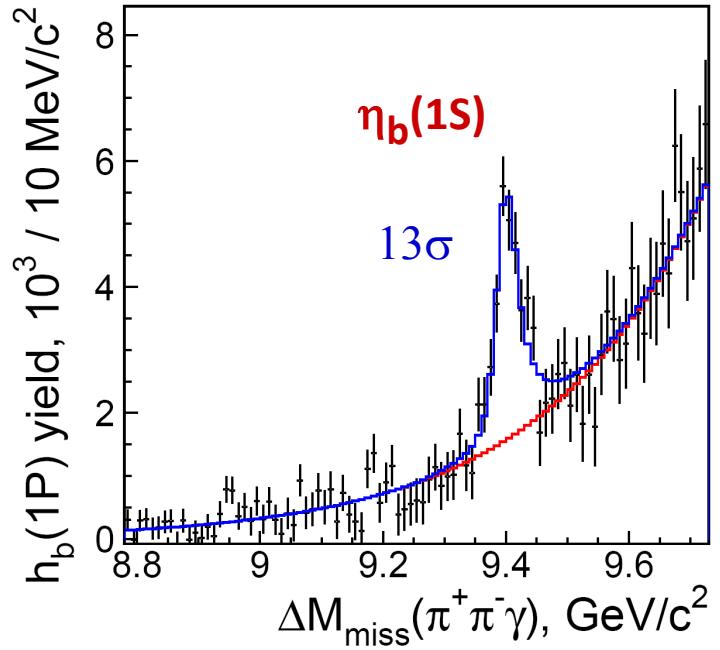
- rectangular bg bands
- no correlation

Approach:

fit $M_{\text{miss}}(\pi^+\pi^-)$ spectra
in $\Delta M_{\text{miss}}(\pi^+\pi^-\gamma)$ bins

Observation of $h_b(1P) \rightarrow \eta_b(1S) \gamma$

N.Brambilla et al., Eur.Phys.J.
C71(2011) 1534 (arXiv:1010.5827)

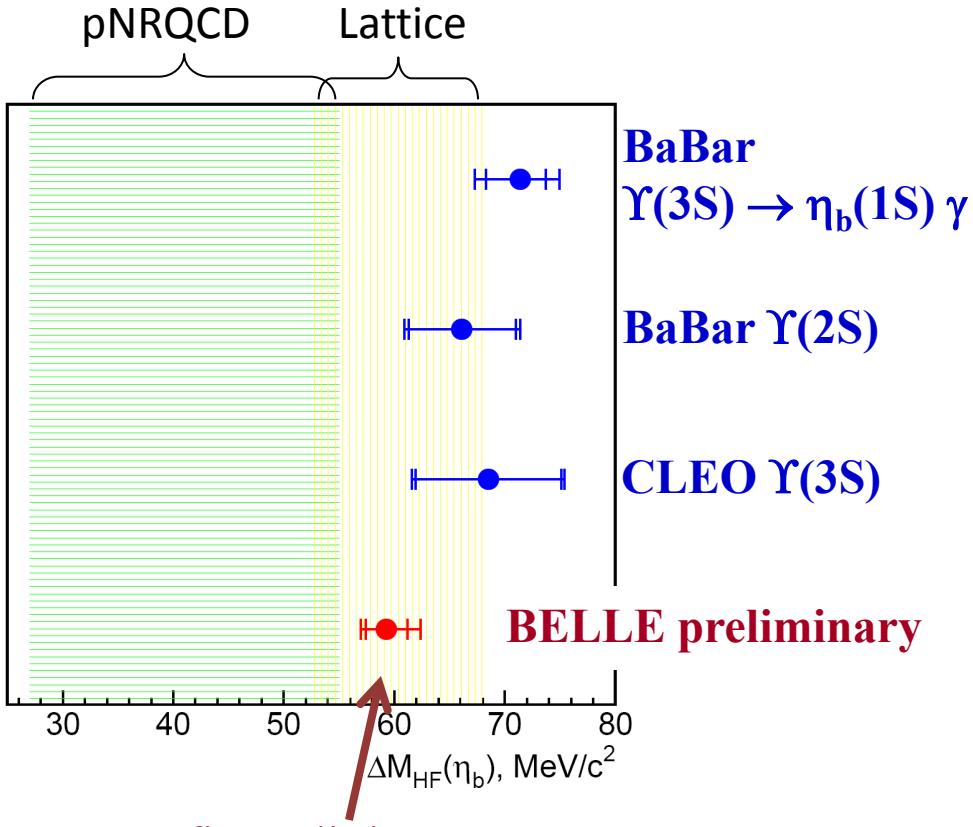


$$N[\eta_b(1S)] = (21.9 \pm 2.0^{+5.6}_{-1.7}) \cdot 10^3$$

$$M[\eta_b(1S)] = (9401.0 \pm 1.9^{+1.4}_{-2.4}) \text{ MeV}/c^2$$

$$\Gamma[\eta_b(1S)] = (12.4^{+5.5+11.5}_{-4.6-3.4}) \text{ MeV}$$

$$\text{Godfrey \& Rosner : BF} = 41\% \quad \mathcal{B}[h_b(1P) \rightarrow \eta_b(1S)\gamma] = (49.8 \pm 6.8^{+10.9}_{-5.2})\%$$



$$\Delta M_{\text{HF}}[\eta_b(1S)] = 59.3 \pm 1.9^{+2.4}_{-1.4} \text{ MeV}/c^2$$

First observation of the $\eta_b(2S)$

$h_b(2P) \rightarrow \gamma \eta_b(2S)$: same method

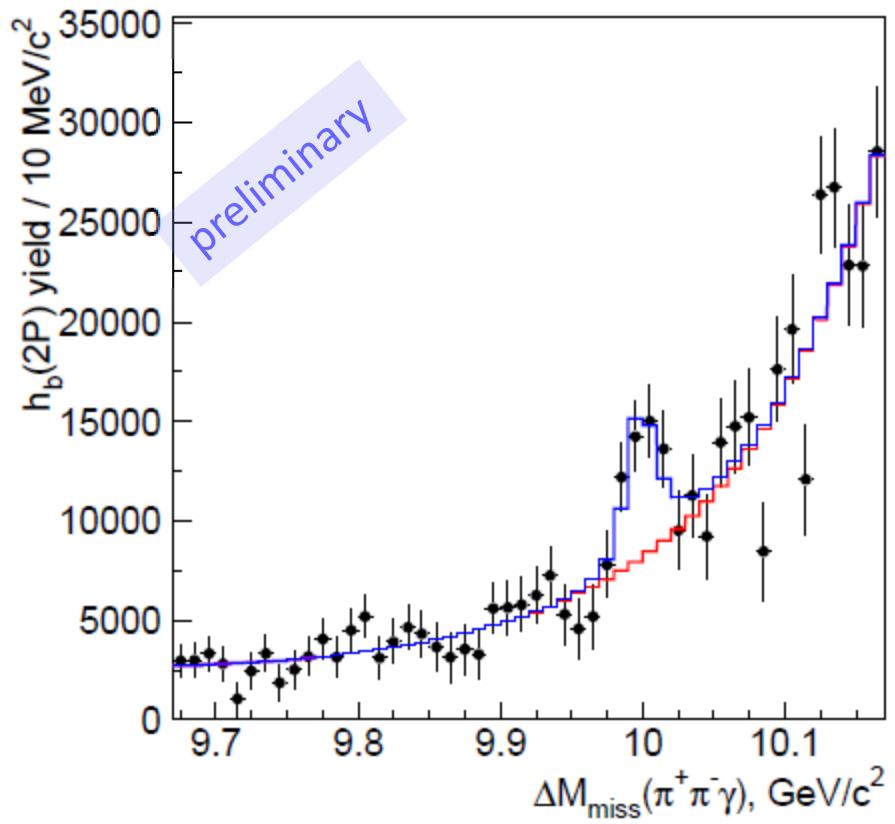
$$M[\eta_b(2S)] = 9999.0 \pm 3.5^{+2.8}_{-1.9} \text{ MeV}$$

$$Bf[h_b(2P) \rightarrow \gamma \eta_b(2S)] = 47.5 \pm 10.5^{+6.8}_{-7.7} \%$$

$$\Delta M_{\text{hfs}}(2S) = 24.3 \pm 3.5^{+2.8}_{-1.9} \text{ MeV}$$

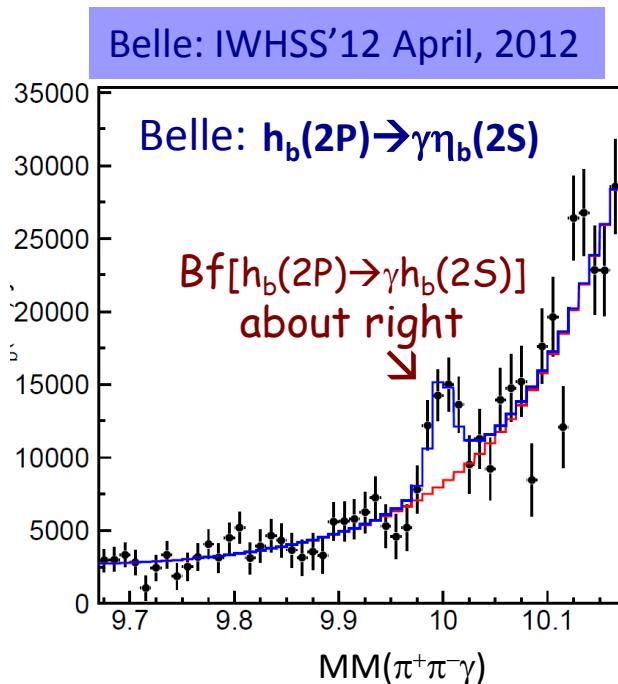
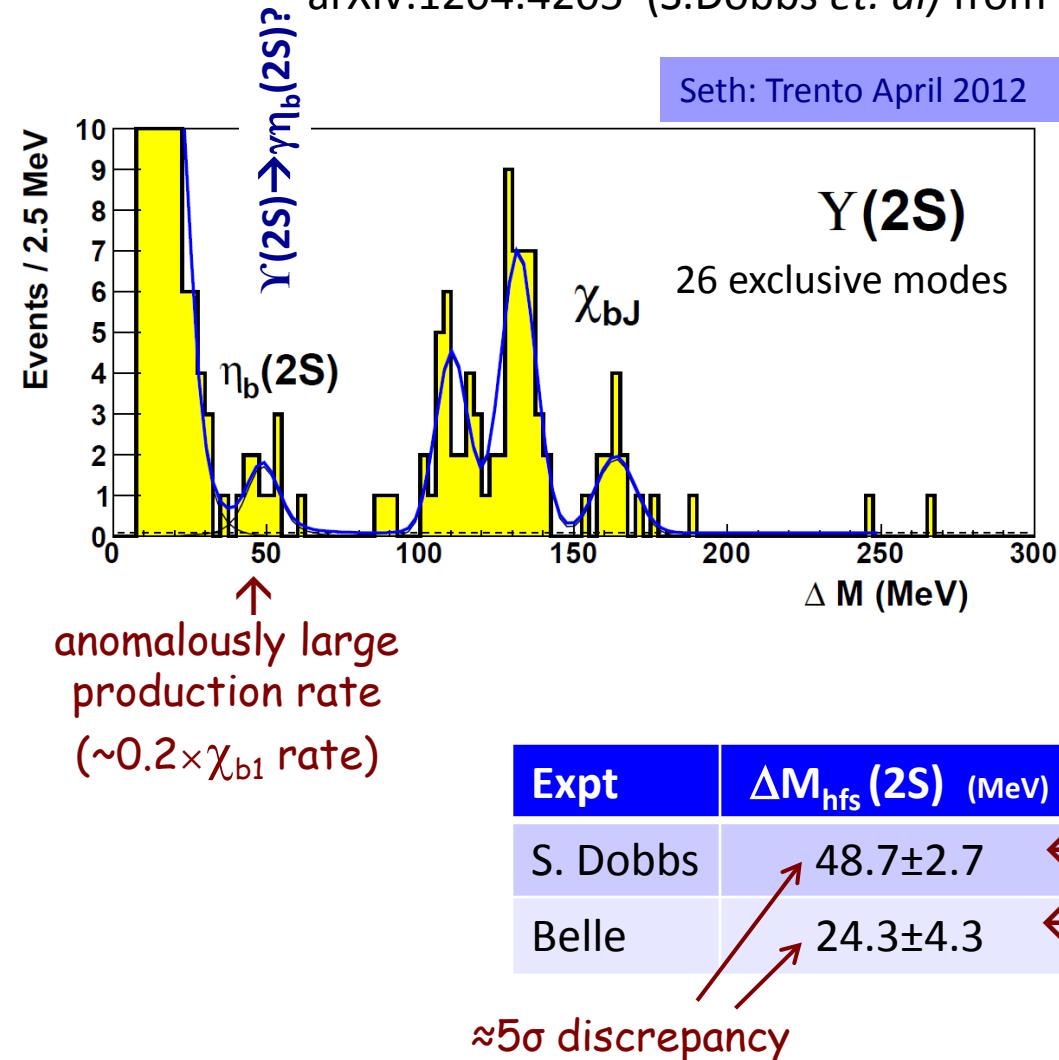
Lattice Meinl PRD82,114502(2010)

$$\Delta M_{HF} = 23.5 \pm 4.7 \text{ MeV}$$

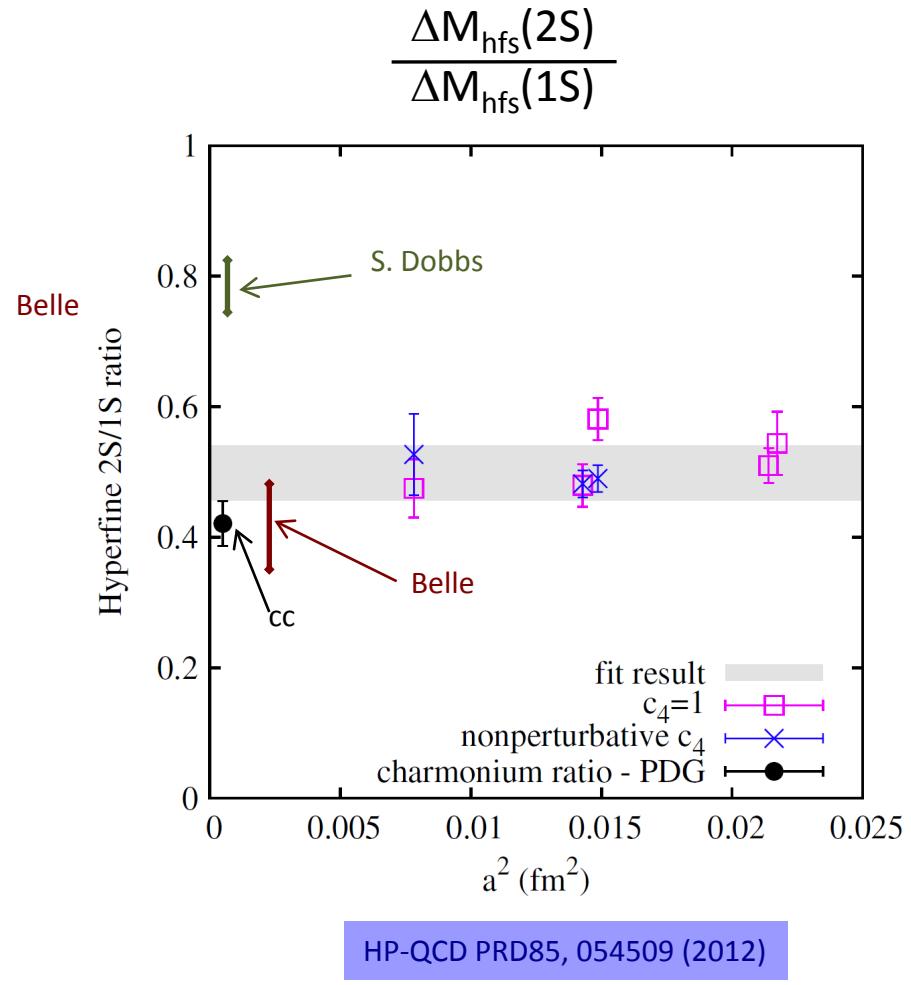
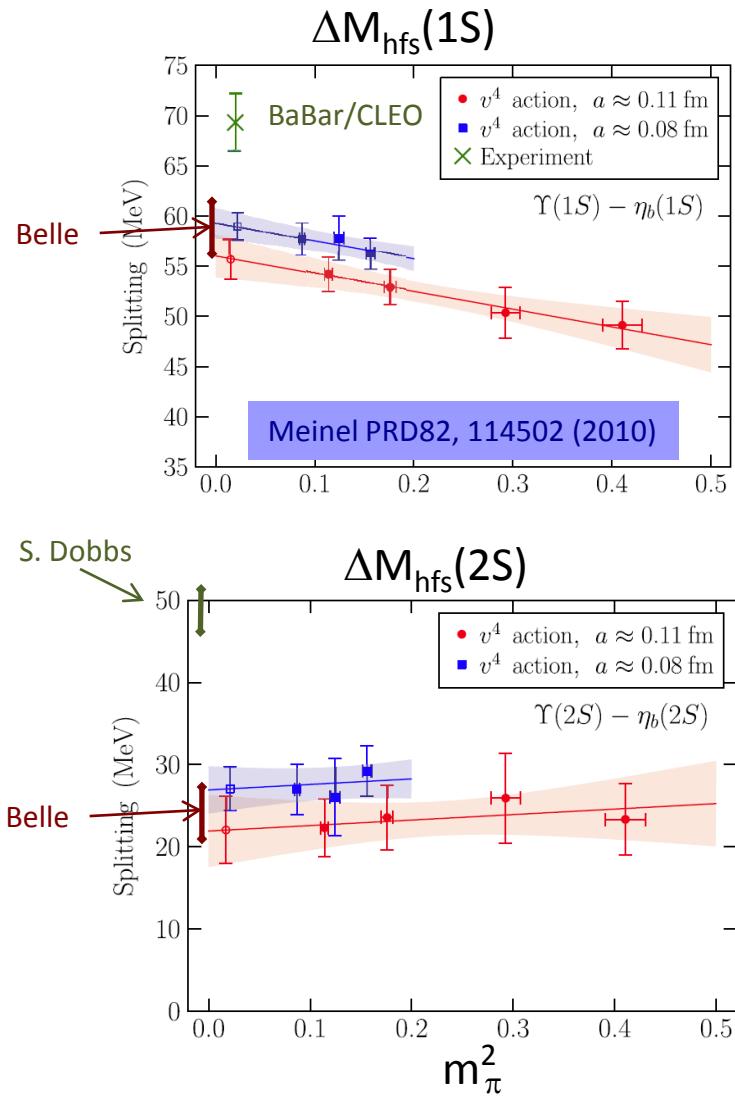


Comparison: $\eta_b(2S)$ “signals”

arXiv:1204.4205 (S.Dobbs *et. al.*) from CLEO data



LQCD predictions for $\Delta M_{\text{hfs}}(1,2S)$

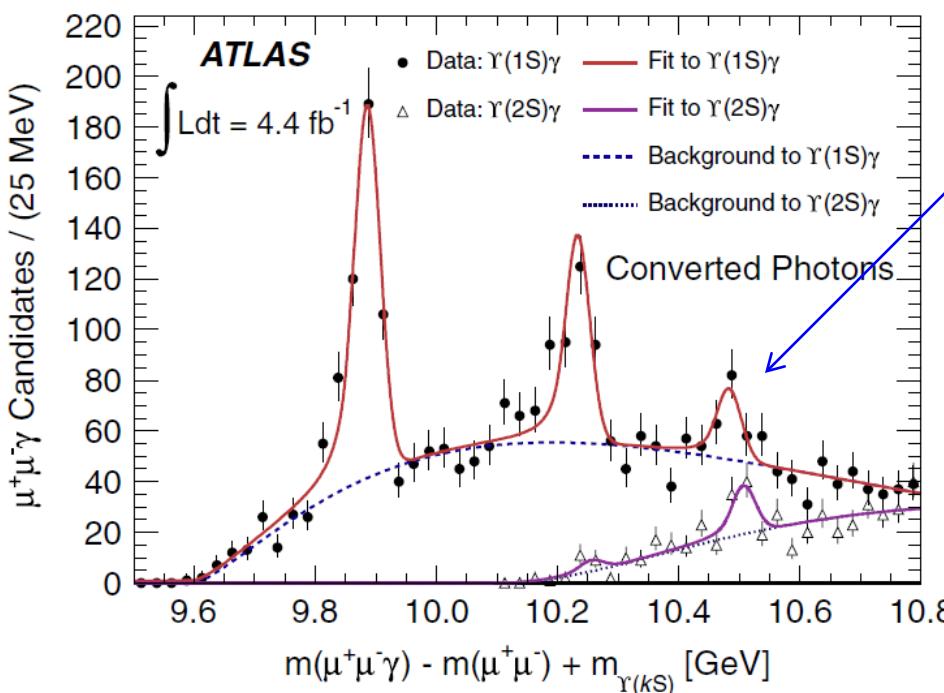


Observation of $\chi_b(3P)$

Converted photon is used for mass determination

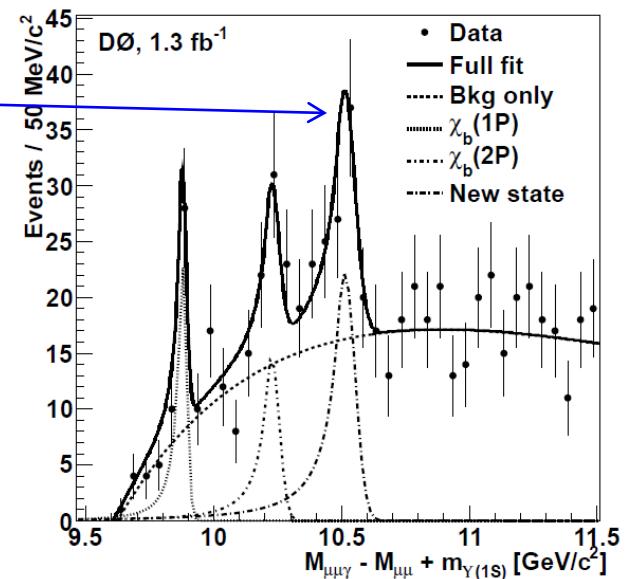
$\chi_b(3P) \rightarrow \gamma(1,2S)\gamma$ first by ATLAS:

PRL 108,152001



Confirmed by D0:

arXiv:1203.6034



$10.530 \pm 0.005 \pm 0.009 \text{ GeV}$ ATLAS

Spin-weighted mass for $\chi_b(3P)$: $10.551 \pm 0.014 \pm 0.017 \text{ GeV}$ D0

10.5439 (Potential model, arXiv: 1201.4096)

$\gamma(nS) \rightarrow \eta \gamma(mS)$ Status

$\gamma(nS) \rightarrow \eta \gamma(mS)$: E1M2
transition, spin-flip

M. B. Voloshin, Prog. Part. Nucl. Phys. 61, 455-511 (2008).
Y.-P. Kuang, Front. Phys. China 1, 19-37 (2006).

QCD multipole formalism:

- spin-flip amplitude scale as $1/m_b$.
- η transition suppressed compared to $\pi\pi$ transition.

Scale from $\psi' \rightarrow \eta J/\psi$

$$B[Y(2S) \rightarrow Y(1S)\eta] \sim 8 \times 10^{-4}$$

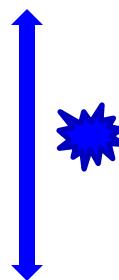
Experiment:

$$(2.39 \pm 0.31 \pm 0.14) \times 10^{-4}$$

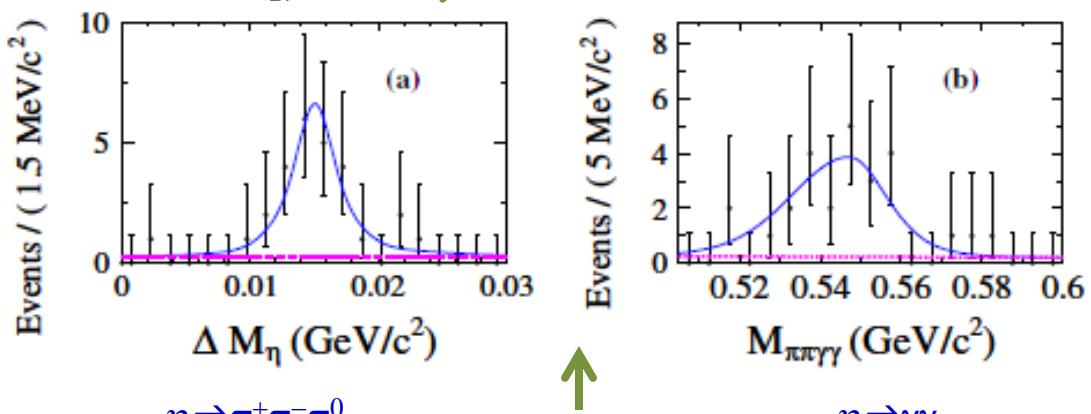


$$(2.1 \pm 0.7 \pm 0.4) \times 10^{-4}$$

(CLEO, PRL 101, 192001)



$$\begin{aligned} B[Y(4S) \rightarrow Y(1S)\eta] &\approx 2.5 \times \\ B[Y(4S) \rightarrow Y(1S)\pi^+\pi^-] & \\ (\text{BaBar, PRD 78, 112002}) & \end{aligned}$$

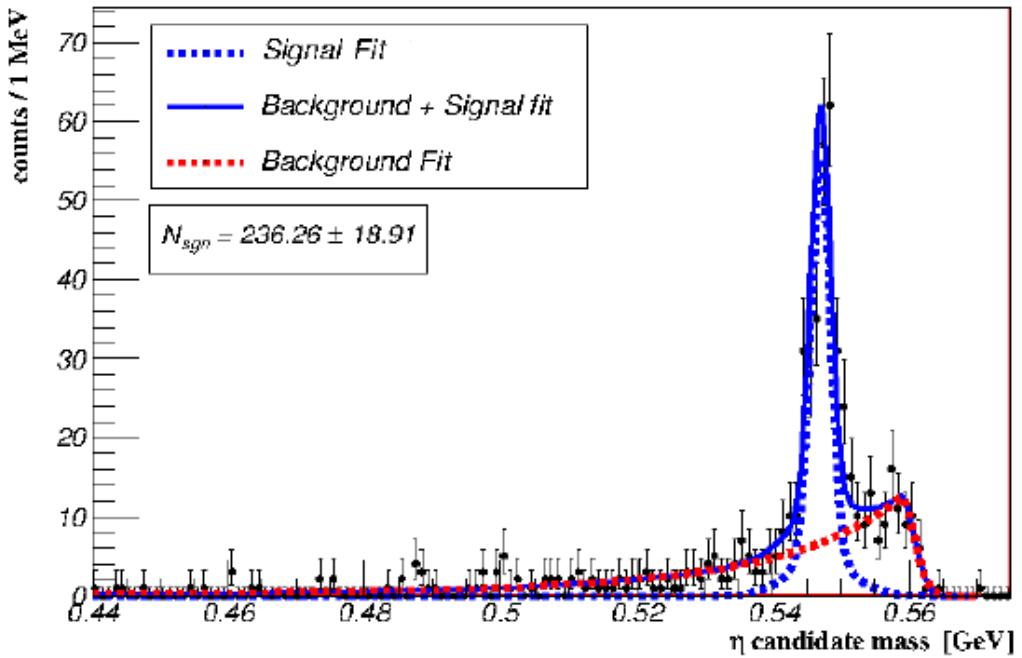


Phys. Rev. D 84 092003

$$B[Y(3S) \rightarrow Y(1S)\eta] < 1.0 \times 10^{-4}$$

$\gamma(2S) \rightarrow \eta, \pi^0 \gamma(1S)$

Preliminary

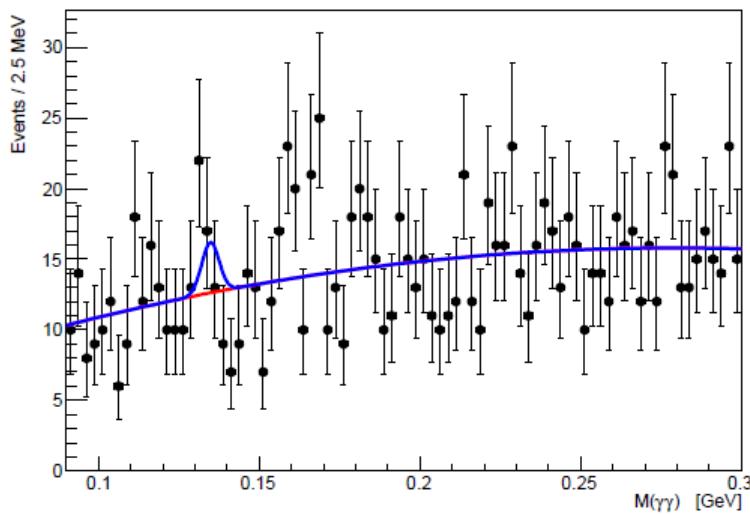


$$\begin{aligned} B[\Upsilon(2S) \rightarrow \Upsilon(1S)\eta] \\ = (3.41 \pm 0.28 \pm 0.35) \times 10^{-4} \end{aligned}$$

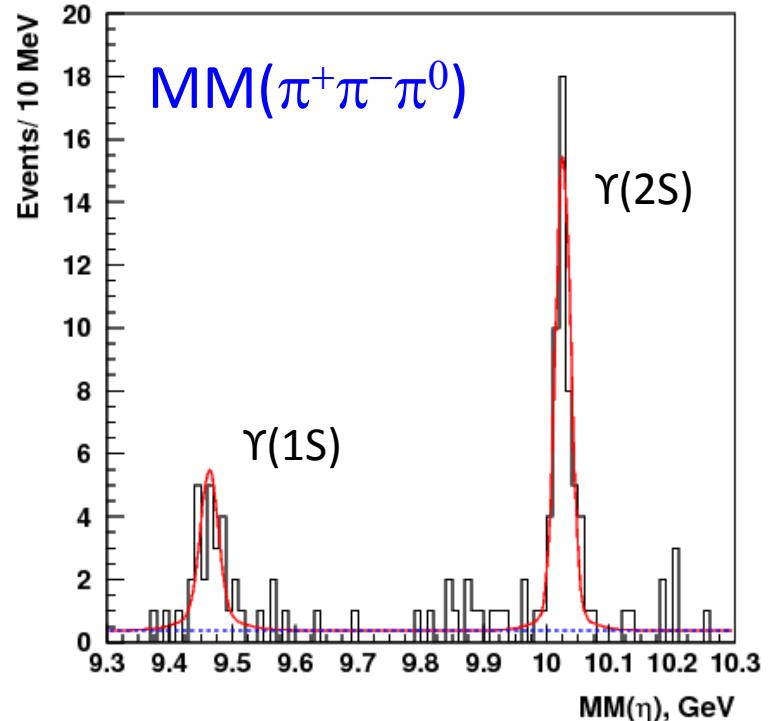
Still < half of theory prediction

$B(\Upsilon(2S) \rightarrow \pi^0 \Upsilon(1S)) < 0.43 \times 10^{-4}$ (90% CL)
 $\sim 6 \times 10^{-4}$ (Theory)

Scale from $\psi' \rightarrow \eta J/\psi$



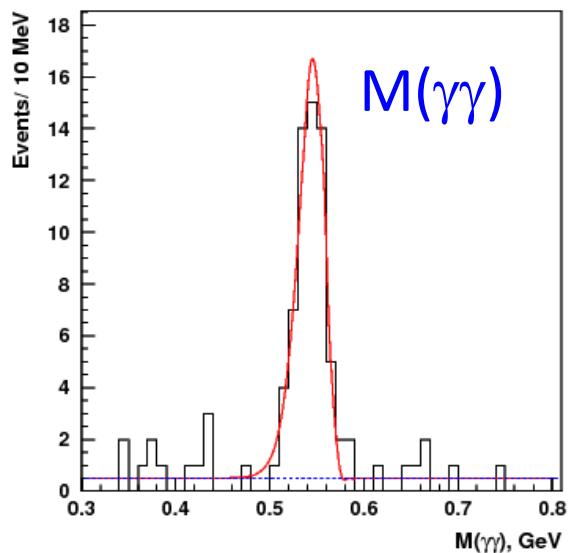
Observation of $\Upsilon(5S) \rightarrow \Upsilon(1,2S)\eta$



- Three modes:

- $\Upsilon(1,2S)[\mu^+\mu^-] \eta[\pi^+\pi^-\pi^0]$
- $\Upsilon(2S)[\Upsilon(1S)\pi^+\pi^-] \eta[\gamma\gamma]$
- $\Upsilon(1S)[\mu^+\mu^-] \eta'[\eta\pi^+\pi^-]$

preliminary

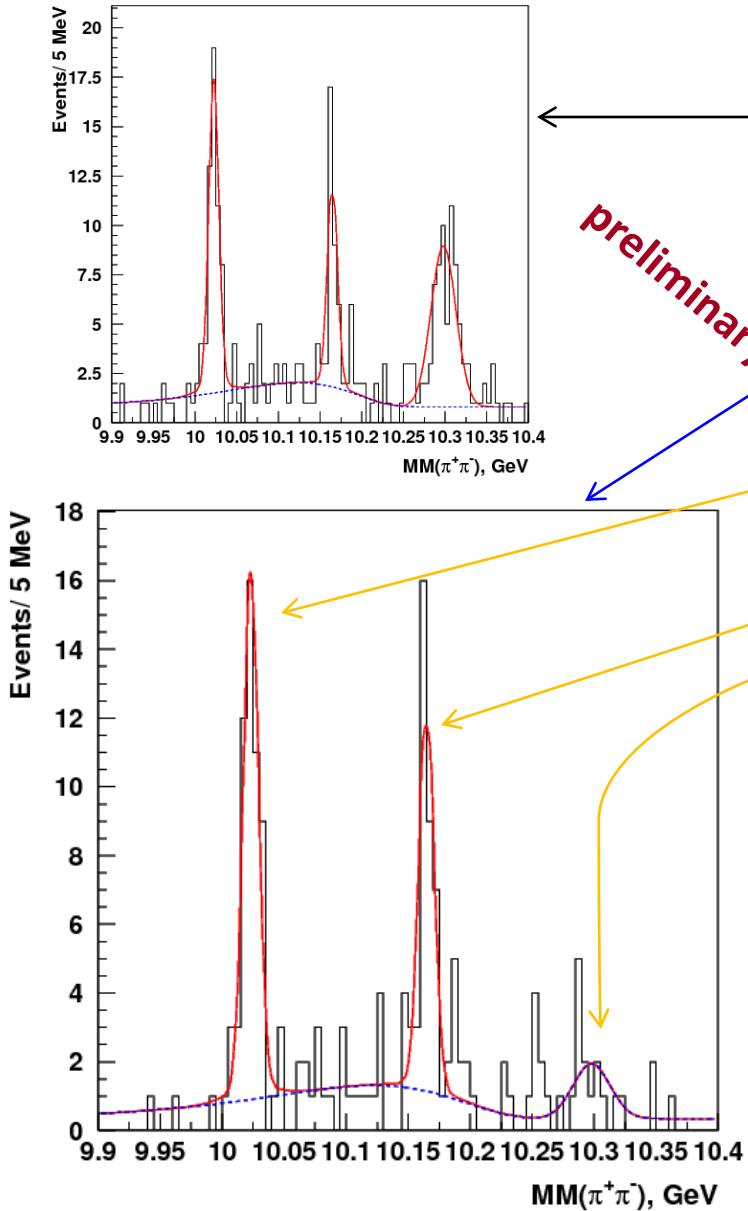


$$B[\Upsilon(5S) \rightarrow \Upsilon(1S)\eta] = (7.3 \pm 1.6 \pm 0.8) \cdot 10^{-4}$$

$$B[\Upsilon(5S) \rightarrow \Upsilon(2S)\eta] = (38 \pm 4 \pm 5) \cdot 10^{-4}$$

$$B[\Upsilon(5S) \rightarrow \Upsilon(1S)\eta'] < 1.2 \cdot 10^{-4}$$

Observation of $\Upsilon(5S) \rightarrow \Upsilon(1D)\pi^+\pi^-$



- $\Upsilon(1S)[\mu^+\mu^-]\pi^+\pi^-\gamma\gamma$ final state
- After $\chi_b(1P) \rightarrow Y\gamma$ selection
- Three peaks in $MM(\pi^+\pi^-)$:
 - $\Upsilon(2S)\pi^+\pi^-$
 - $\Upsilon(1D)\pi^+\pi^-$
 - $\Upsilon(2S)[Y\pi^+\pi^-]\eta[\gamma\gamma]$ reflection

statistical significance 9σ

$$B[\Upsilon(5S) \rightarrow \Upsilon(2S)\pi^+\pi^-] = (7.5 \pm 1.1 \pm 0.8) \times 10^{-3} \text{ (cross check)}$$

$$B[\Upsilon(5S) \rightarrow \Upsilon(1D)\pi^+\pi^-] B[\Upsilon(1D) \rightarrow \chi_b(1P)\gamma \rightarrow \Upsilon(1S)\gamma\gamma] \\ = (2.0 \pm 0.4 \pm 0.3) \times 10^{-4}$$

P.S. : An evidence of $\Upsilon(1D)$
was seen in inclusive
 $MM(\pi^+\pi^-)$ spectra

Summary

Belle's large $\Upsilon(5S)$ data gives unexpectedly new results.

Inclusive $\gamma, \pi^0, \pi^+\pi^-$, spectrum is the main method in discovery of new states.

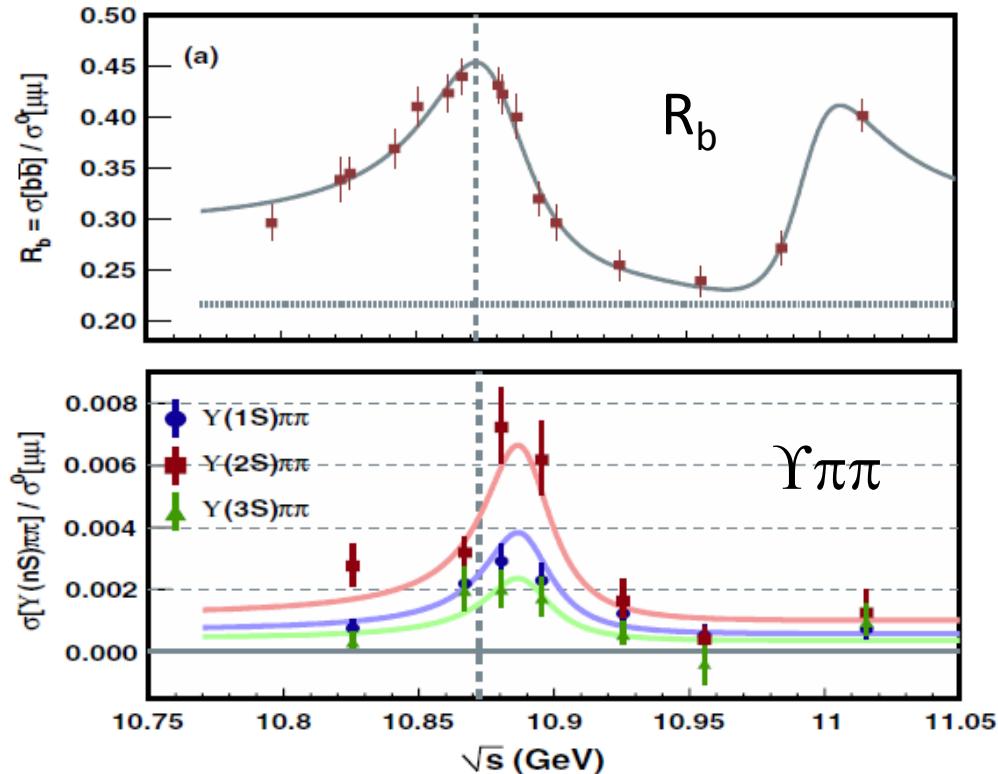
More results from exclusive reconstruction method.



- Observation of $h_b(1P), h_b(2P), Z_{b1}^+, Z_{b2}^+$.
 - Exclusive method in final states $\Upsilon(nS)\pi^+\pi^-$.
 - Charm counterpart in $\Upsilon(4260)$ decay?
- Observation of $h_b(1P) \rightarrow \eta_b(1S)\gamma$.
 - **Observation of $\eta_b(2S)$** from Belle and CLEO (S. Dobbs).
 - Belle: $\Delta M_{hfs}(2S) = 24.3 \pm 4.3$ MeV $Bf[h_b(2P) \rightarrow \gamma\eta_b(2S)] = 47 \pm 13\%$
 - S. Dobbs (CLEO): $\Delta M_{hfs}(2S) = 48.7 \pm 2.7$ MeV, no BF info.
 - Disagreement on the $\Delta M_{hfs}(2S)$.
- Observation of $\chi_b(3P)$ from ATLAS and D0.
- Observation of $\Upsilon(5S) \rightarrow \Upsilon(1,2S)\eta$. First measurement.
- Observation of $\Upsilon(5S) \rightarrow \Upsilon(1D)\pi^+\pi^-$. First measurement.
- Measurement of $\Upsilon(2S) \rightarrow \eta\Upsilon(1S)$ from Belle. Still smaller than theory.

BACKUP

(2) $\Upsilon(nS)\pi\pi$ production shape



Belle, Phys. Rev., D82, 091106R (2010)

Energy scan \Rightarrow
mean of shapes of $\sigma(\Upsilon\pi\pi)$
and hadronic cross section R_b
shapes differ by 2.0σ

$$\mu_{\Upsilon\pi\pi} - \mu_{10860} = 9 \pm 4 \text{ MeV}/c^2$$

$$\Upsilon(1S)\pi^+\pi^- \sigma_{\text{peak}} (\text{pb}) \quad 2.78^{+0.42}_{-0.34} \pm 0.23$$

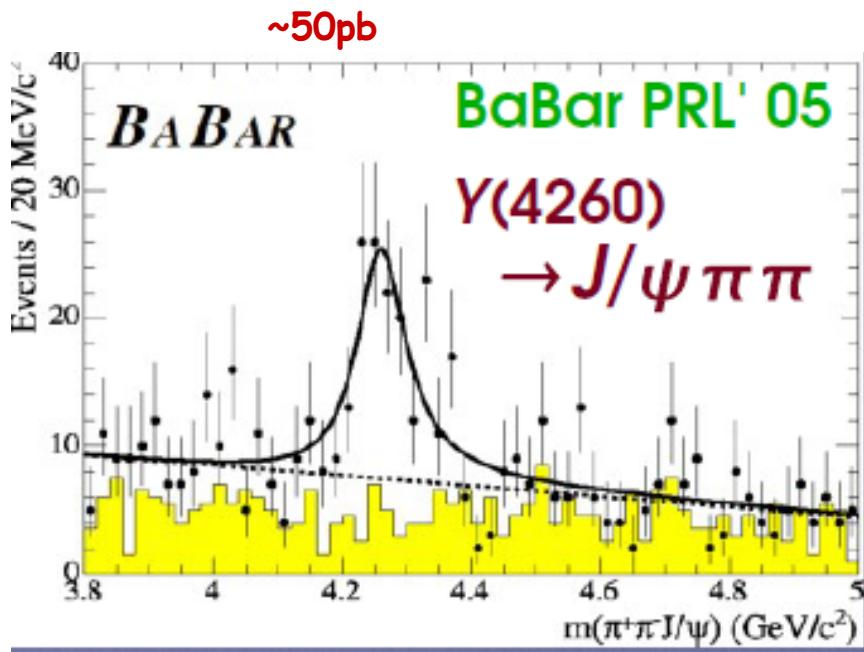
$$\Upsilon(2S)\pi^+\pi^- \sigma_{\text{peak}} (\text{pb}) \quad 4.82^{+0.77}_{-0.62} \pm 0.66$$

$$\Upsilon(3S)\pi^+\pi^- \sigma_{\text{peak}} (\text{pb}) \quad 1.71^{+0.35}_{-0.31} \pm 0.24$$

$$\mu \text{ (MeV}/c^2) \quad 10888.4^{+2.7}_{-2.6} \pm 1.2$$

$$\Gamma \text{ (MeV}/c^2) \quad 30.7^{+8.3}_{-7.0} \pm 3.1$$

Reminder of $\Upsilon(4260)$



- Observed from initial state radiation events: $e^+e^- \rightarrow \gamma J/\psi \pi\pi$.
- Mass far above $D\bar{D}^{(*)}$ threshold.
- Width ~ 100 MeV (quite narrow).
- $J^{PC} = 1^{--}$ from production mechanism .

$$e^+e^- \rightarrow \gamma(\pi^+\pi^-J/\psi)$$



BaBar, Phys. Rev. Lett. **95**, 142001 (2005).
arXiv:0808.1543v2.
Belle, Phys. Rev. Lett. **99**, 182004 (2007).

$$e^+e^- \rightarrow \pi^+\pi^-J/\psi$$

$$e^+e^- \rightarrow \pi^0\pi^0J/\psi$$



CLEO, Phys. Rev. Lett. **96**, 162003 (2006).

$\Upsilon(4260)$ notes

- Decay to $D\bar{D}^{(*)}$ not seen.
- Small coupling to e^+e^- ,
 $BR(J/\psi\pi^+\pi^-) \times \Gamma(e^+e^-)$
 $= 7.5 \pm 0.9 \pm 0.8$ eV. BaBar, arXiv:0808.1543

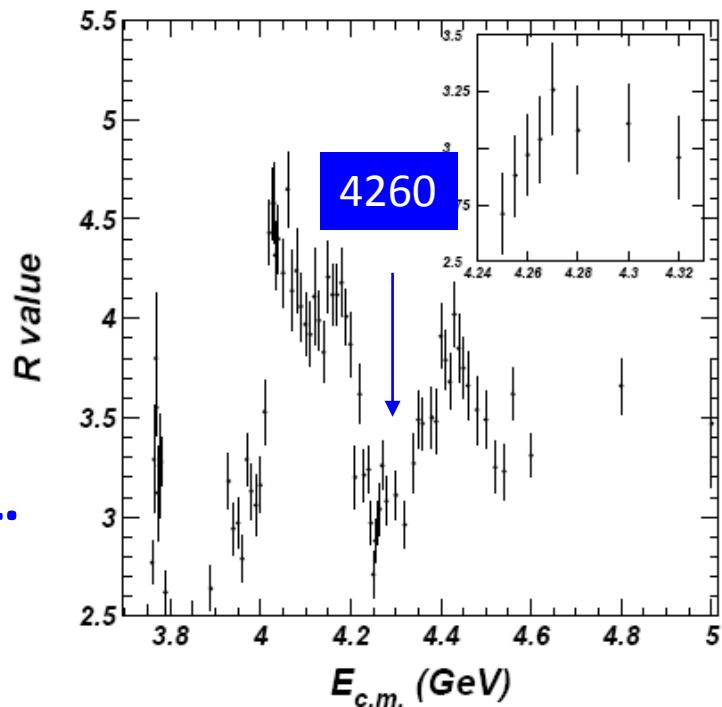
- Small coupling to $p\bar{p}$
 $BR(\Upsilon(4260) \rightarrow p\bar{p}) /$
 $BR(\Upsilon(4260) \rightarrow J/\psi\pi^+\pi^-) < 0.13$.

BaBar, Phys.Rev.D73,012005(2006)

- $\Gamma(\Upsilon \rightarrow J/\psi\pi^+\pi^-) > 508$ keV @ 90% C.L.
 \Rightarrow Large !!

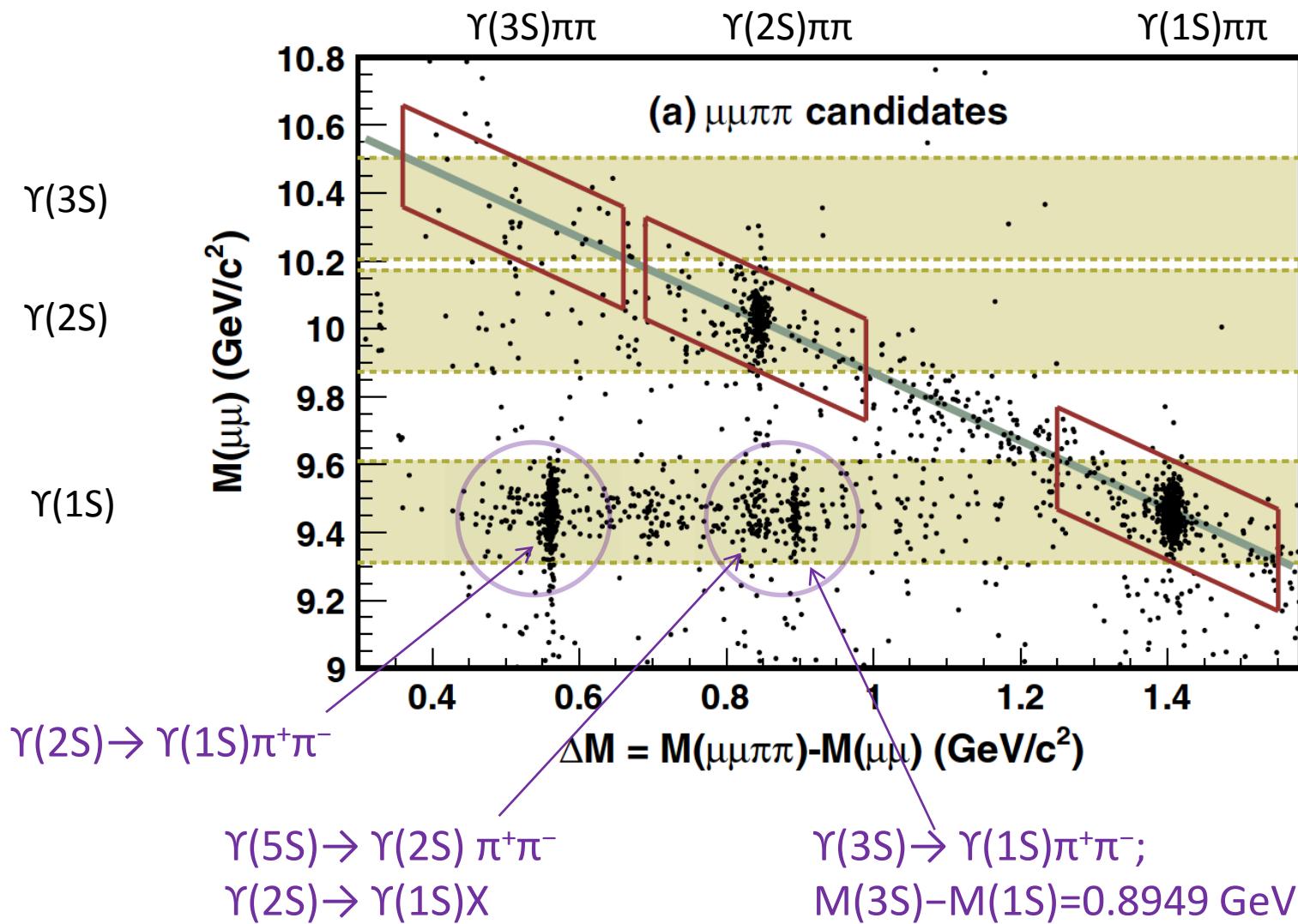
X. H. Mo et al., Phys. Lett. B 640 (2006) 182

BES, Phys. Rev. Lett. 88, 101802 (2002)
At a dip in $\sigma(e^+e^-) \rightarrow \text{hadrons}$

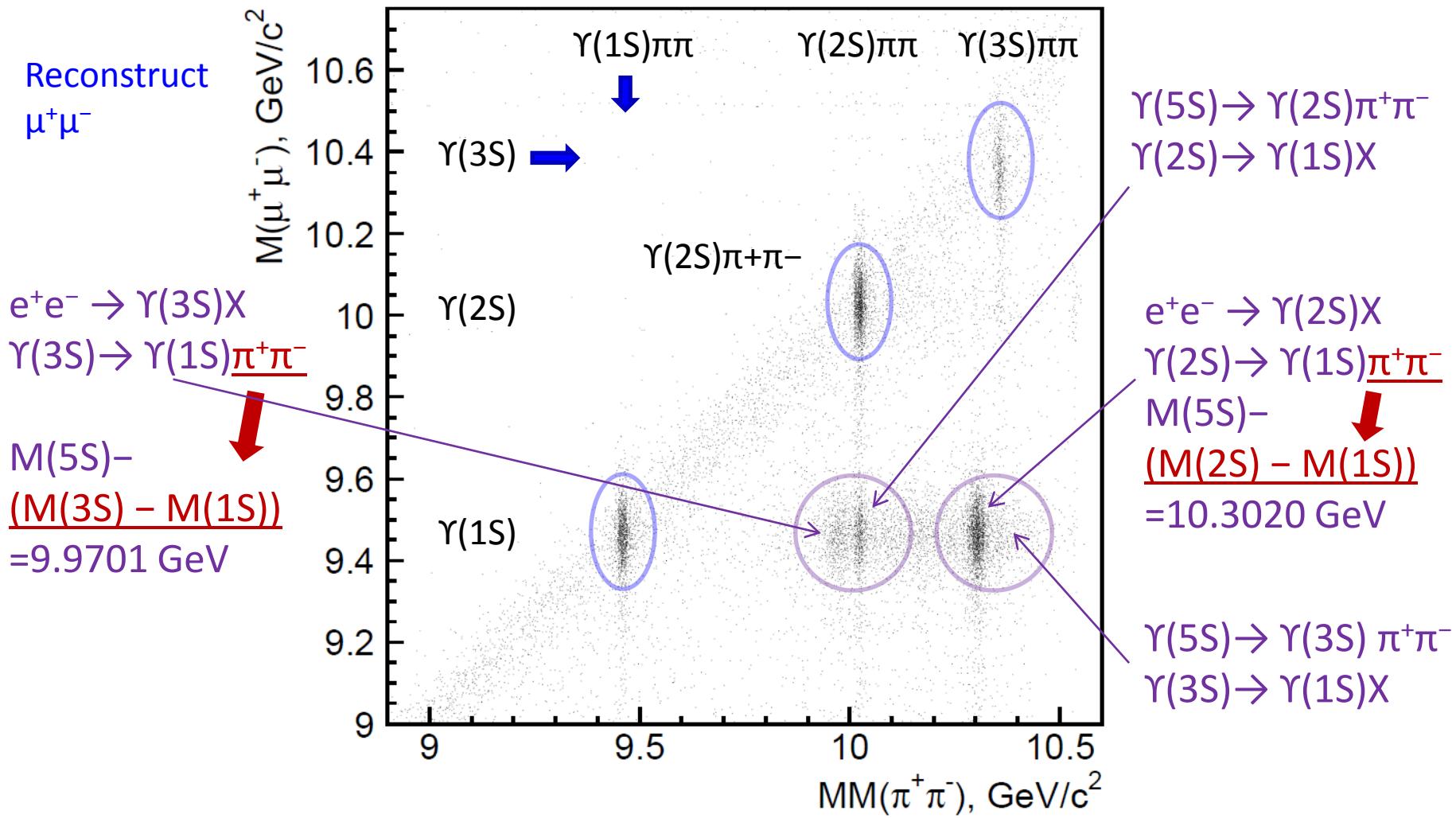


In b-sector, may be there also exists a counterpart Υ_b which has also large h_b branching fraction.

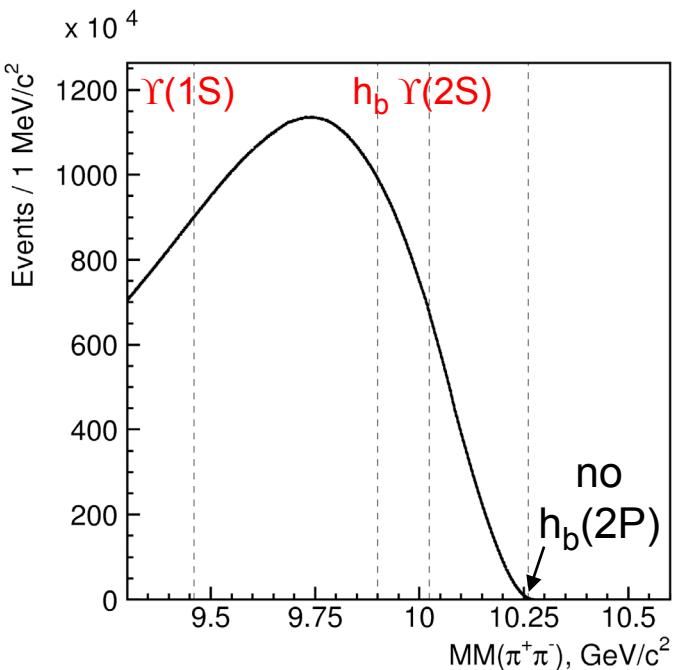
Exclusive (full) reconstruction



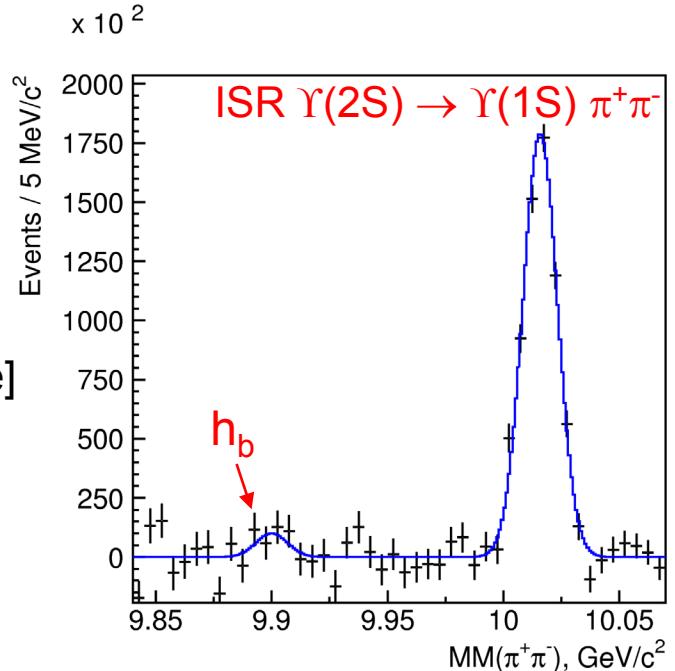
Missing mass technique: Calibration sources in $\Upsilon(5S)$ events



h_b in $\Upsilon(4S)$ data



$L = 711 \text{ fb}^{-1}$
[$\times 6$ $\Upsilon(5\text{S})$ sample]



No significant signal of $h_b(1\text{P})$: $(34 \pm 20) \times 10^3$ (1.7σ)

$$\frac{\sigma[e^+e^- \rightarrow h_b(1\text{P}) \pi^+\pi^-] @ \Upsilon(4\text{S})}{\sigma[e^+e^- \rightarrow h_b(1\text{P}) \pi^+\pi^-] @ \Upsilon(5\text{S})} < 0.28 \text{ at } 90\% \text{ C.L.}$$

⇒ $\Upsilon(4\text{S})$ does not show anomalous properties

Angular Analysis

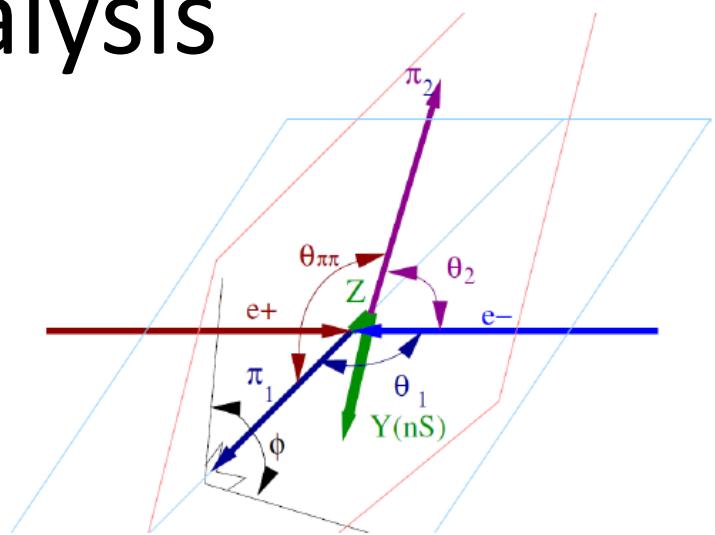
Z_b velocity is very small ($\beta < 0.02$) \Rightarrow measure all pion momenta in the c.m. frame.

$$\Upsilon(5S) \rightarrow Z_b(\Upsilon(nS), h_b + \boldsymbol{\pi}_1) \boldsymbol{\pi}_2$$

$$\theta_1 = \angle(\boldsymbol{\pi}_1, e^+)$$

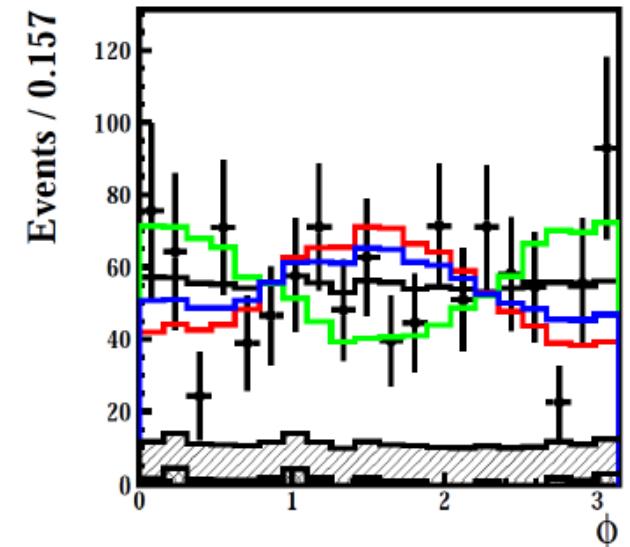
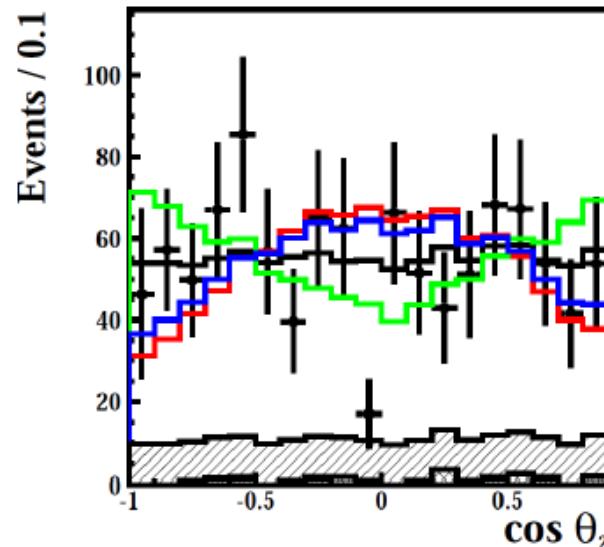
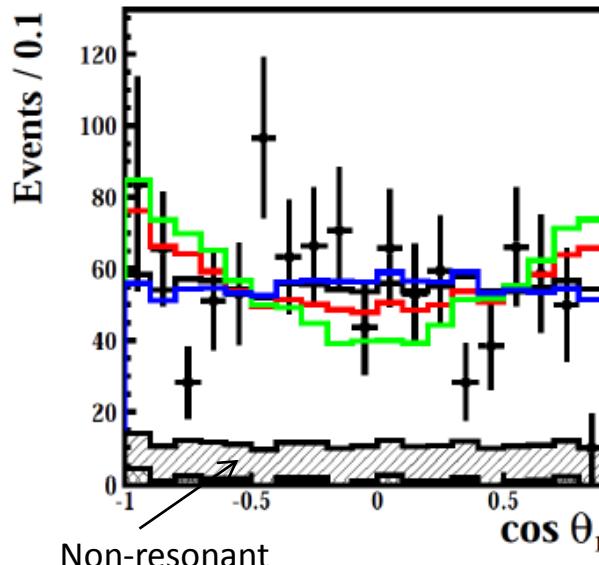
$$\theta_2 = \angle(\boldsymbol{\pi}_2, e^+)$$

$$\phi = \angle(\text{plane}(\boldsymbol{\pi}_1, e^+), \text{plane}(\boldsymbol{\pi}_1, \boldsymbol{\pi}_2))$$



$$J^P = 1^+, 1^-, 2^+, 2^-$$

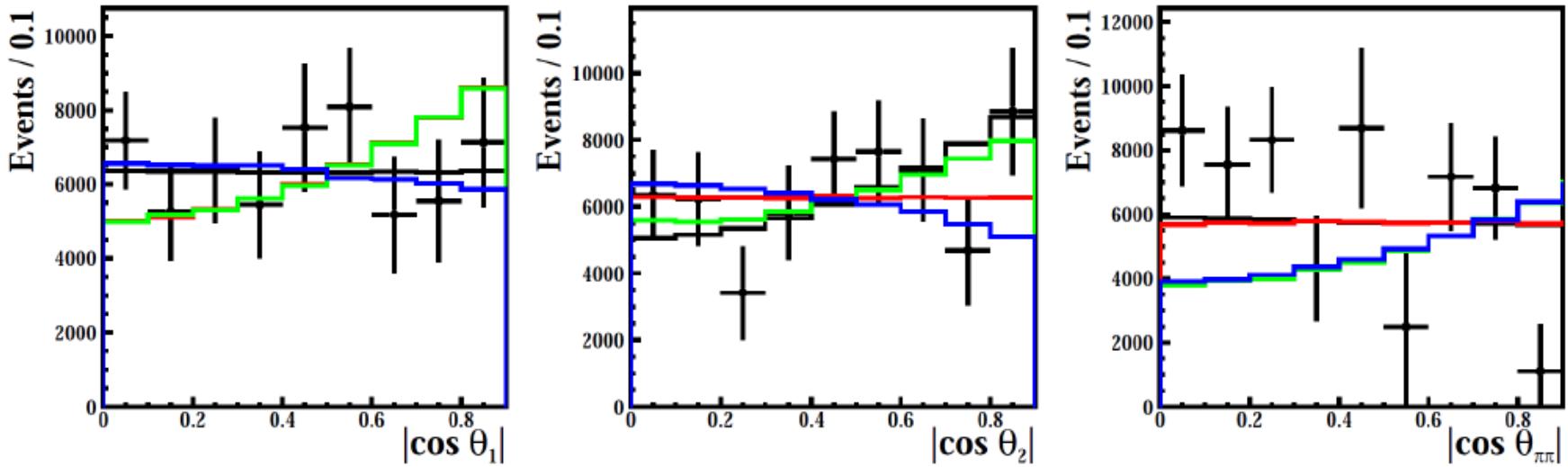
Example: $\Upsilon(5S) \rightarrow Z_b(10610) \boldsymbol{\pi}_2 \rightarrow \Upsilon(2S) \boldsymbol{\pi}_1 \boldsymbol{\pi}_2$



Angular Analysis (h_b final state)

$J^P = 1^+, 1^-, 2^+, 2^-$

Example: $\Upsilon(5S) \rightarrow Z_b(10610) \pi_2 \rightarrow h_b(1P) \pi_1 \pi_2$



Probabilities at which other J^P hypotheses are disfavored with respect to 1^+

J^P	$Z_b(10610)$			$Z_b(10650)$		
	$\Upsilon(2S)\pi^+\pi^-$	$\Upsilon(3S)\pi^+\pi^-$	$h_b(1P)\pi^+\pi^-$	$\Upsilon(2S)\pi^+\pi^-$	$\Upsilon(3S)\pi^+\pi^-$	$h_b(1P)\pi^+\pi^-$
1^-	3.6σ	0.3σ	0.3σ	3.7σ	2.6σ	2.7σ
2^+	4.3σ	3.5σ	4.3σ	4.4σ	2.7σ	2.1σ
2^-	2.7σ	2.8σ		2.9σ	2.6σ	

$J^P=1^+$ is favored!

$\Upsilon(5S) \rightarrow \Upsilon(1S)\eta'$

