

Physics Beyond Colliders Annual Workshop 2/3/2021

Detecting dark matter at BASE-CERN

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- 1. Couplings between antiprotons and dark matter
- 2. Dark matter decays in the BASE detectors
- 3. Future perspectives





BSE Measuring magnetic moments

$$\frac{g}{2} = \frac{2\pi v_L}{B_e \mu_N} = \frac{v_L}{v_c}$$

 v_L is the Larmor frequency, the frequency at which rf can most easily flip the spin

To measure v_L , just need a way to identify the antiproton spin state

Detecting changes in the direction or size of μ



The particle has a total magnetic moment μ (spin + orbital)

If we add an inhomogeneous field $B = B_0 + B_2(z^2 - \rho^2/2)$ then there is an additional force $\nabla(\mu \cdot B)$ which contributes to the axial trapping force

This shifts the axial frequency depending on μ , so the spin state can be identified

SE Couplings between dark matter and antiprotons

Measure the coupling $\mathcal{L}_{int} = -\frac{\partial a}{f_a} \bar{\psi} \gamma^{\mu} \gamma^5 \psi$ between ultralight, pseudoscalar ALP relic dark matter and \bar{p} Interaction

$$H_{\text{int}} = -\frac{C_{\bar{p}}a_0}{2f_a}\sin(\omega_a t)\,\vec{\sigma}_{\bar{p}}\cdot\vec{p}_a$$

between the momentum of the axion field \vec{p}_a and the antiproton spin vector $\vec{\sigma}_{\vec{p}}$ oscillating at the axion Compton frequency $\omega_a = m_a c^2 / \hbar$

Should cause characteristic time dependent variation in v_L , by constraining the size of this a- \bar{p} coupling limits extracted



a- \overline{p} coupling limits a natural bi-product of precision CPT tests

The BASE frequency detection system

Measurement of ~1 fA image currents induced in trap electrodes







Resolving single antiproton spin flips requires the highest Q and lowest temperature LC resonant detectors ever built- BASE-CERN is the state of the art

H. Nagahama *et al.*, Rev. Sci. Instrum. **87**, 113305 (2016). 6

-50

50

0



Conversion of Axion-like particles into photons in the detector

Axions can couple to photons via the interaction term $\mathcal{L}_{int} = -\frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$

This modifies Maxwell's equations

$$\begin{aligned} \nabla \cdot \vec{E} &= \rho - g_{a\gamma} \, \vec{B} \cdot \nabla a \\ \nabla \times \vec{B} - \partial_t \vec{E} &= \vec{J} + g_{a\gamma} (\vec{B} \partial_t a - \vec{E} \times \nabla a) \\ \nabla \cdot \vec{B} &= 0 \\ \nabla \times \vec{E} + \partial_t \vec{B} &= 0 \end{aligned}$$

Inside the resonator housing, $d \ll \lambda_a$, and where there is a strong field B_e , the axions source a magnetic field

$$\left|\vec{B}_{a}\right| = \frac{1}{2}rg_{a\gamma}\left|\vec{B}_{0}\right|\sqrt{\rho_{a}\hbar c}$$





f(v,Q,q) is a lorentzian line-shape function proportional to Re{Z} e_n is the equivalent input noise of the amplifier

 κ is the coupling constant

Q is the resonator Q-factor

 N_T is the number of turns

l is the length of the <u>toriod</u> along the magnet B field

- r₁ is the inner radius of the toroid
- r_2 is the outer radius

 $g_{a\gamma}$ is the coupling constant

- B is the static magnetic field
- ho_a is the dark matter density

How to measure T_z ?

J. A. Devlin et al., Phys. Rev. Lett. 126, 041301 (2021). 8

SE A quantum "Boltzmann" thermometer



Trapped antiproton's axial motion reaches thermal equilibrium with the detector- can use it as a "quantum" sensor

3.

Magnetic inhomogneity gives axial frequency shift proportional to $\mu \propto n_-$



Knowing strength of inhomogeneity, can determine T_z





We use maximum likelihood estimation to estimate $g_{a\gamma}$, and then decide if we have a discovery, or set limits



J. A. Devlin et al., Phys. Rev. Lett. 126, 041301 (2021); J. W. Foster et al. Phys. Rev. D 97 (2018). 10

SE Next step: frequency tuning

Cryogenic adjustable capacitance with no loss of Q already developed



SE Future potential

Small detector(s): 5 cm long, 5 cm diameter

Immediately realizable with BASE technology today, 6-9 months assembly time $\left(\frac{e_n}{\kappa} = 2.4 \text{ nV}/\sqrt{\text{Hz}}, 7 \text{ T}, 1 \text{ year acquisition}\right)$

Possible in the short term with detector RnD work using BASE know-how $(\frac{e_n}{\kappa} = 2.4 \text{ nV}/\sqrt{\text{Hz}}, 7 \text{ T}, 1 \text{ year acquisition, use multiple}$ higher frequency coils without excessive Q-loss)

Large detector: 80 cm long, 14 cm diameter

Optimistic projection, "normal" Penning trap magnet $\left(\frac{e_n}{r} = 2.4 \text{ nV}/\sqrt{\text{Hz}}, 7 \text{ T}, 1 \text{ year acquisition, 10 mK, Q=200,000}\right)$

Pushing the sensitivity further

- Much large detector volumes- in discussion with RADES/babyIAXO
- Colder detectors- laser cooled resonators?
 - Lower noise amplifiers particle assisted readout?



Input from other CERN experts appreciated! e.g. NbTi coating or HTS tape for large cavities.

SE Takeaway message-unique possibilities at BASE

- 1. Long-term investigations of **antiproton**-relic **dark matter** coupling
- 2. Transfer **state-of-the-art technologies** developed for high precision antimatter measurements to detect relic dark matter particles
- 3. BASE Penning traps at CERN as an interface between **quantum measurement** techniques and large detectors for weak EM radiation



• **Mainz:** Measurement of the magnetic moment of the proton, implementation of new





- **CERN-AD:** Measurement of the magnetic moment of the antiproton and proton/antiproton q/m ratio
- Hannover/PTB: QLEDS-laser cooling project, new technologies







Institutes: RIKEN, MPIK, CERN, University of Mainz, Tokyo University, GSI Darmstadt, University of Hannover, PTB Braunschweig

