

# PROBING RAPIDLY OSCILLATING ULTRALIGHT DARK MATTER

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*Super-resolution in Physics Workshop*

**INBAR SAVORAY**

**07.02.21**

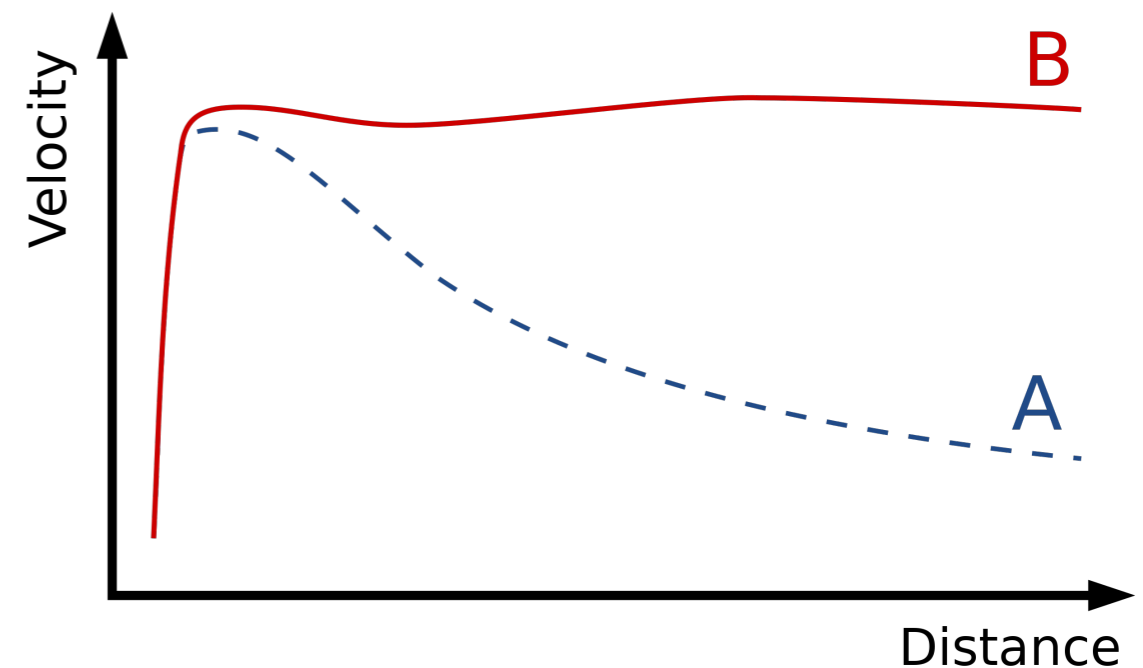
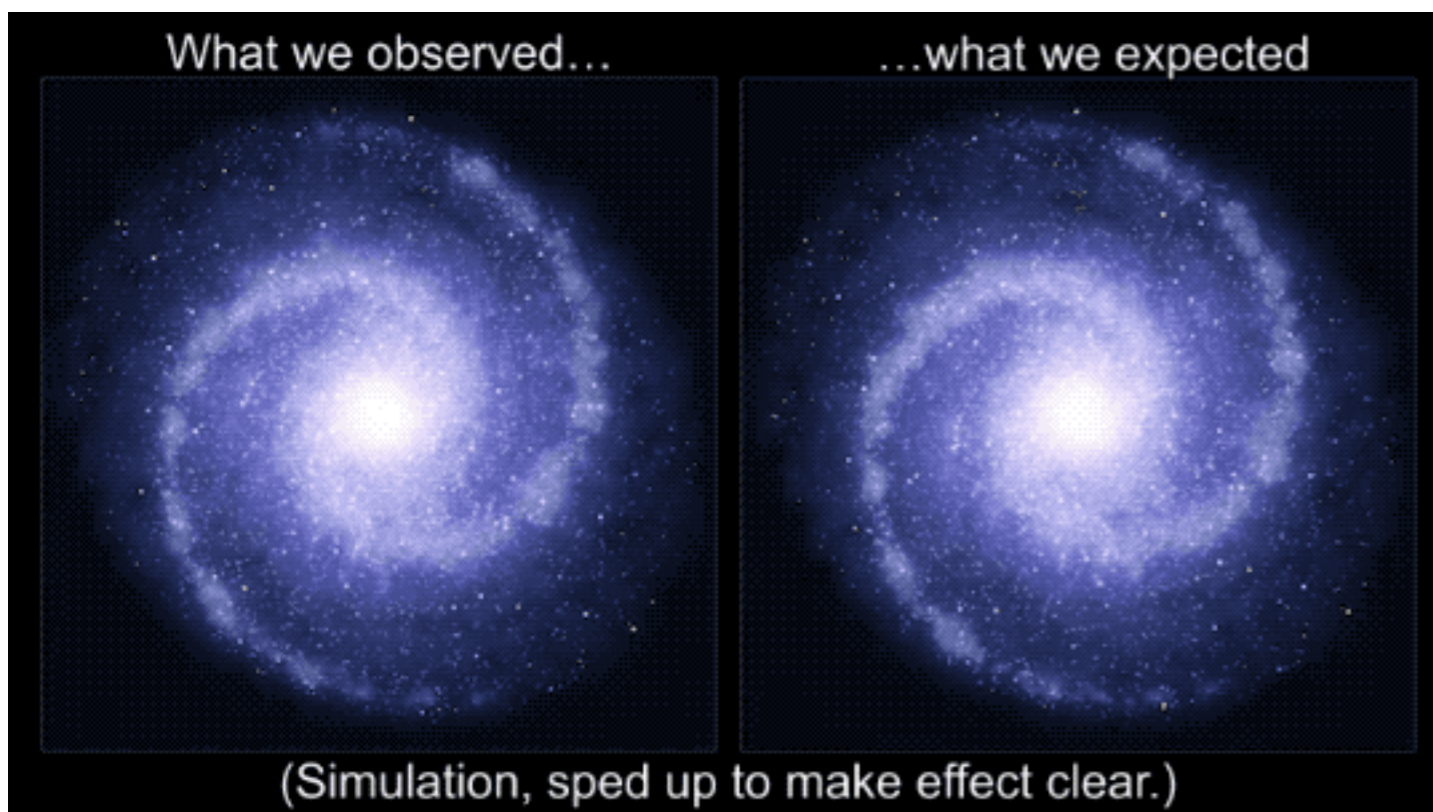
# OUTLINE

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- Dark matter
- Ultralight Dark Matter
- Atomic probes of Ultralight Dark Matter
- Rapidly oscillating Ultralight Dark Matter
- Super-resolution?

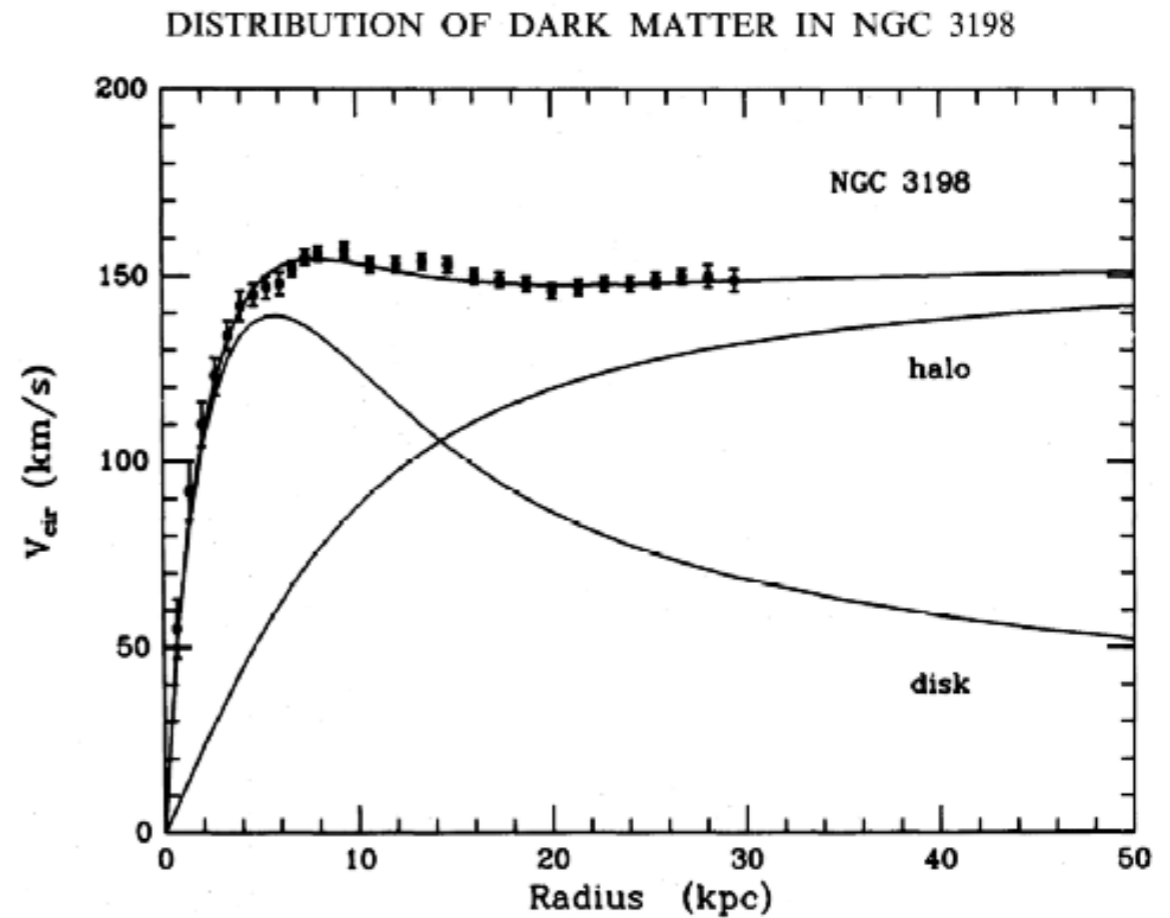
# DARK MATTER

- **The missing mass problem** - the discrepancy between the amount of visible (baryonic) matter, and the expected amount of gravitationally interacting matter in the universe.
- Supported by various astronomical and cosmological observations such as: galaxy rotation curves, structure formation, CMB data and more...



# DARK MATTER

- **The missing mass problem** - the discrepancy between the amount of visible (baryonic) matter, and the expected amount of gravitationally interacting matter in the universe.
- Supported by various astronomical and cosmological observations such as: galaxy rotation curves, structure formation, CMB data and more...
- **Dark matter** -
  - a “new” type of matter
  - has gravitational interactions, but only very weak (or no) interactions with light.



# ULTRALIGHT SCALAR DARK MATTER

- DM number density in the galaxy  $n_\phi = \rho_{DM}/m_\phi$
- De-Broglie length  $\lambda_{DB} = h/(m_\phi v)$
- Ultralight DM ( $m < 0.1$  eV) - many particles within a De-Broglie cube

$$N = n \cdot \lambda_{dB}^3 = \frac{\rho_{DM}}{m^4 v^3} \approx \left( \frac{\rho_{DM}}{0.3 \text{ GeV}/\text{cm}^3} \right) \left( \frac{0.1 \text{ eV}}{m} \right)^4$$

- Cannot be treated as individual particles, but rather as a **macroscopic coherent field** (much like a laser/BEC/superfluid)
- Ultralight scalar DM is a field the coherently oscillating

$$\langle \phi(t) \rangle \simeq \frac{\sqrt{2\rho_{DM}}}{m_\phi} \cos(m_\phi t)$$

# ULTRALIGHT DM COUPLING TO THE SM

- The DM field can be searched for through its interactions with “ordinary” (Standard Model) particles.
- e.g. *Higgs-portal* models, in which the DM couples to the SM through its mixing with the Higgs, inheriting its interactions.

$$g_{\phi i} = g_{hi} \sin \theta$$



# ULTRALIGHT DM COUPLING TO THE SM

- If the DM field is coupled to electrons and/or photons

$$\mathcal{L}_{int} \supset -g_{\phi e} m_e \phi \bar{e} e + \frac{g_{\phi\gamma}}{4} \phi F^{\mu\nu} F_{\mu\nu}$$

- The DM background field would appear to “give mass” to the electrons (similarly to the Higgs), and/or modify the field strength of EM fields.
- Temporal oscillations of  $m_e, \alpha$  at a frequency determined by the mass of the field.

$$\delta m_e = g_{\phi e} m_e \langle \phi(t, \vec{x}) \rangle, \quad \delta \alpha = g_{\phi\gamma} \langle \phi(t, \vec{x}) \rangle \alpha^{SM} \quad \langle \phi(t, \vec{x}) \rangle \simeq \frac{\sqrt{2\rho_{DM}}}{m_\phi} \cos(m_\phi t),$$

# TEMPORAL VARIATIONS OF FUNDAMENTAL CONSTANTS

- Oscillations of the DM field induce oscillations of  $m_e, \alpha$

$$\delta m_e = g_{\phi e} m_e \langle \phi(t, \vec{x}) \rangle, \quad \delta \alpha = g_{\phi \gamma} \langle \phi(t, \vec{x}) \rangle \alpha^{SM} \quad \langle \phi(t, \vec{x}) \rangle \simeq \frac{\sqrt{2\rho_{DM}}}{m_\phi} \cos(m_\phi t),$$

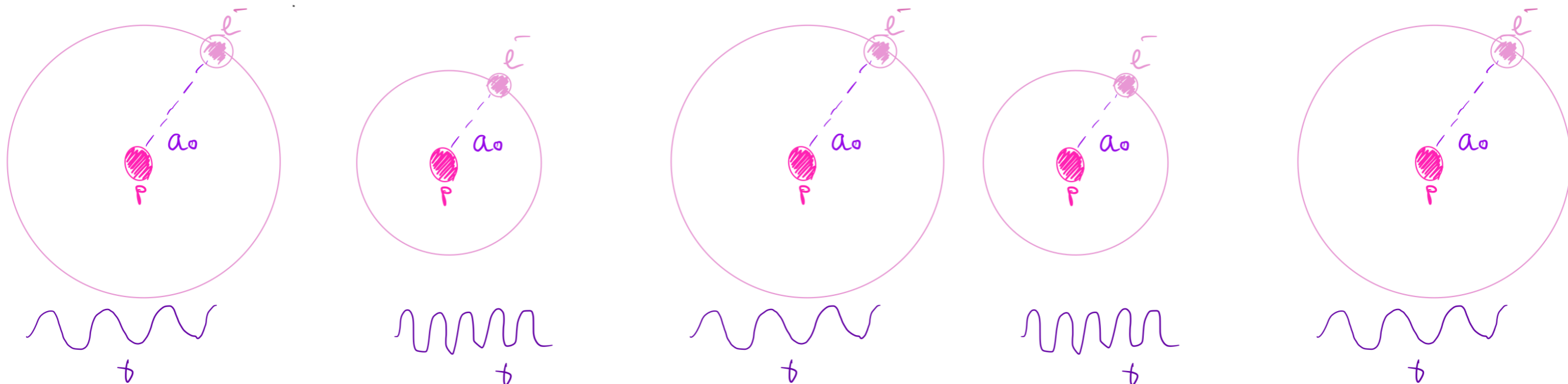
- Temporal oscillations of atomic frequencies

Y. V. Stadnik and V. V. Flambaum, [1412.7801]

- Rydberg energy levels  $R_\infty \propto \alpha^2 m_e$

- cavity modes frequencies  $a_0 \propto (\alpha m_e)^{-1}$

- H hyperfine transition  $\omega_{HF} \propto \alpha^4 m_e^2$



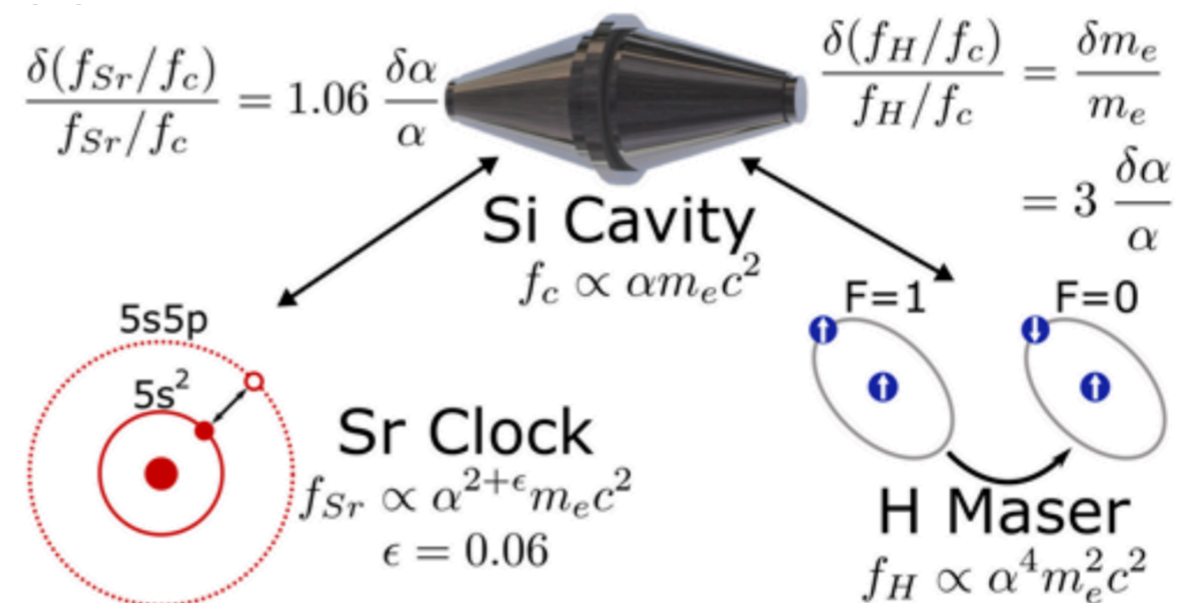


# SEARCHING FOR ULDM WITH ATOMIC CLOCKS

➤ Goal - measure oscillations of fundamental constants induced by ULDM.

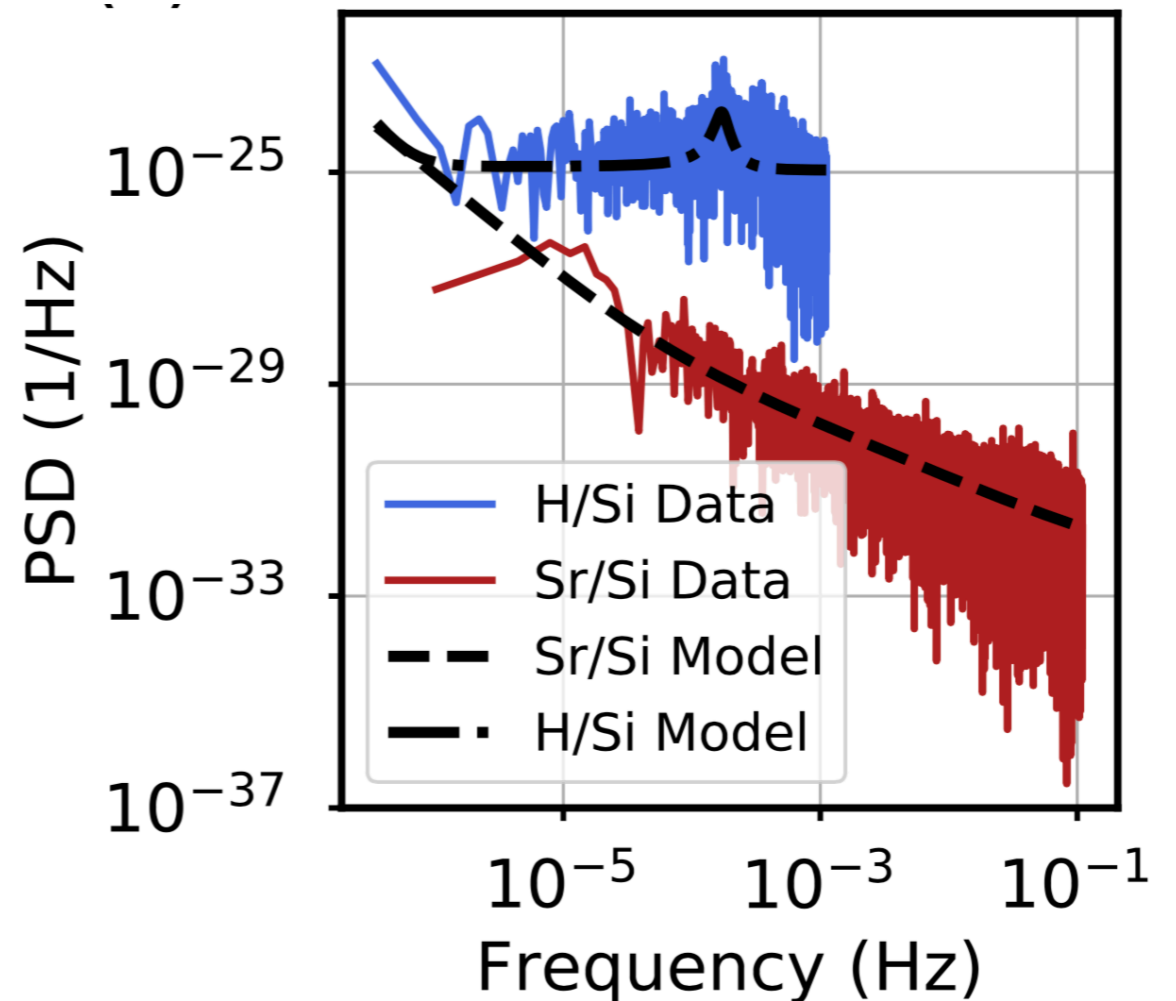
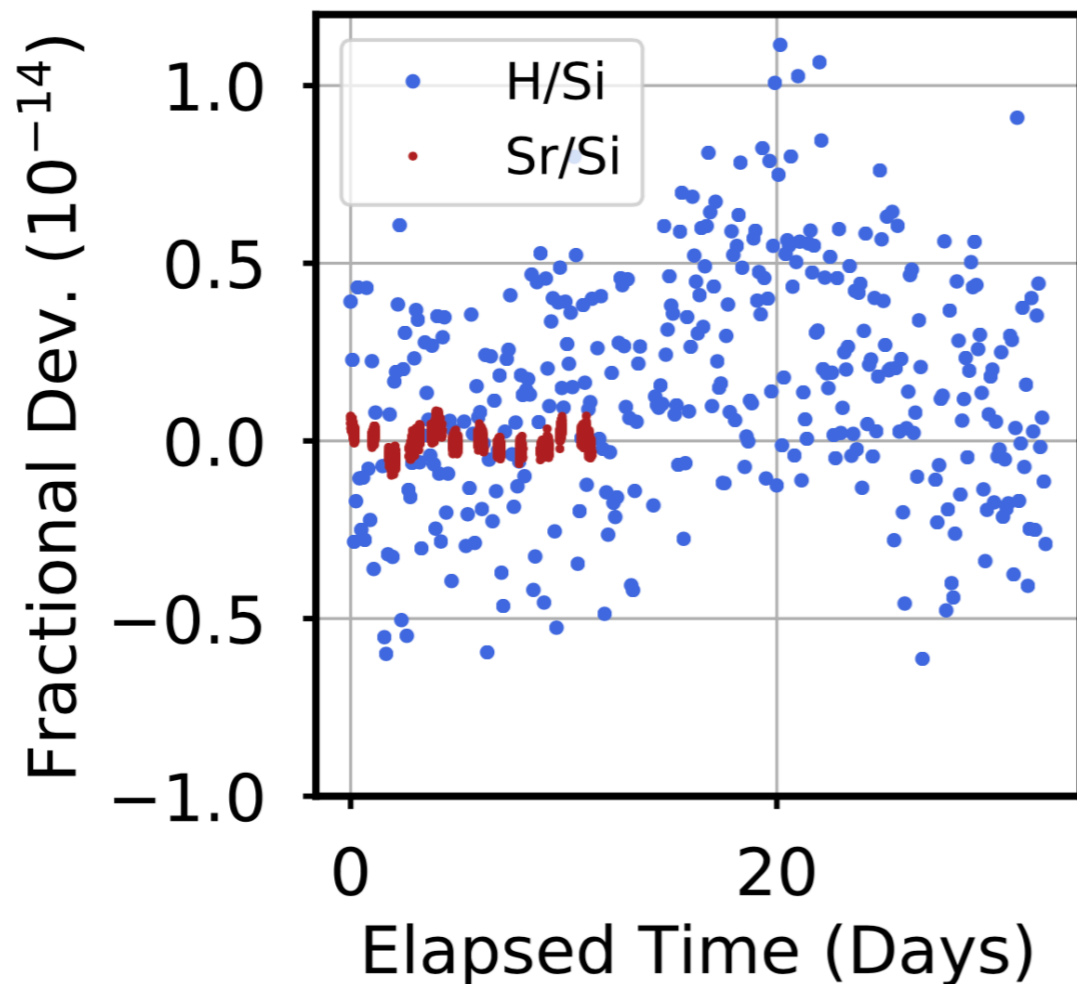
1. Compare two atomic clocks with clock frequencies that depend differently on  $\alpha, m_e$

$$\frac{\delta(f_A/f_B)}{f_A/f_B} = \frac{\delta f_A}{f_A} - \frac{\delta f_B}{f_B} \approx \Delta n_{A,B}^{\alpha} \frac{\delta \alpha}{\alpha} + \Delta n_{A,B}^{m_e} \frac{\delta m_e}{m_e}$$



# SEARCHING FOR ULDM WITH ATOMIC CLOCKS

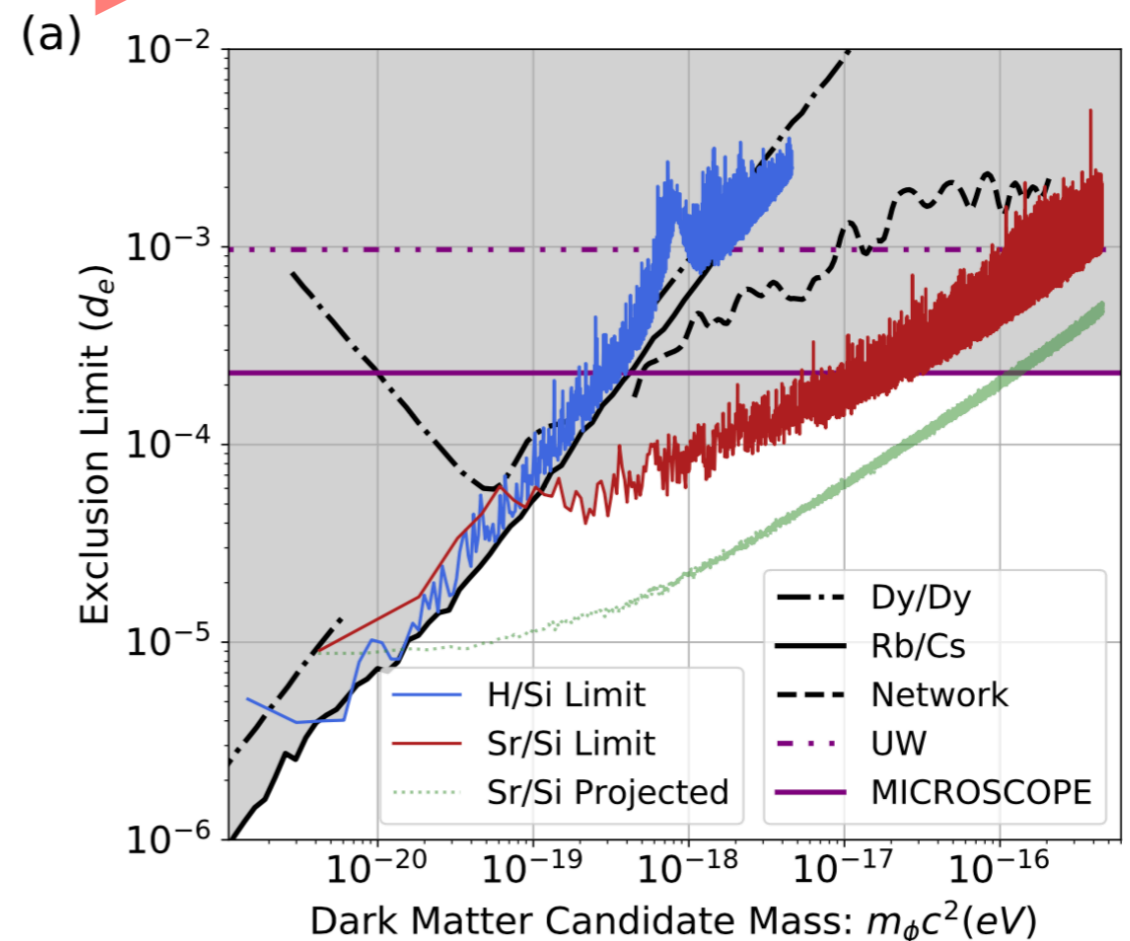
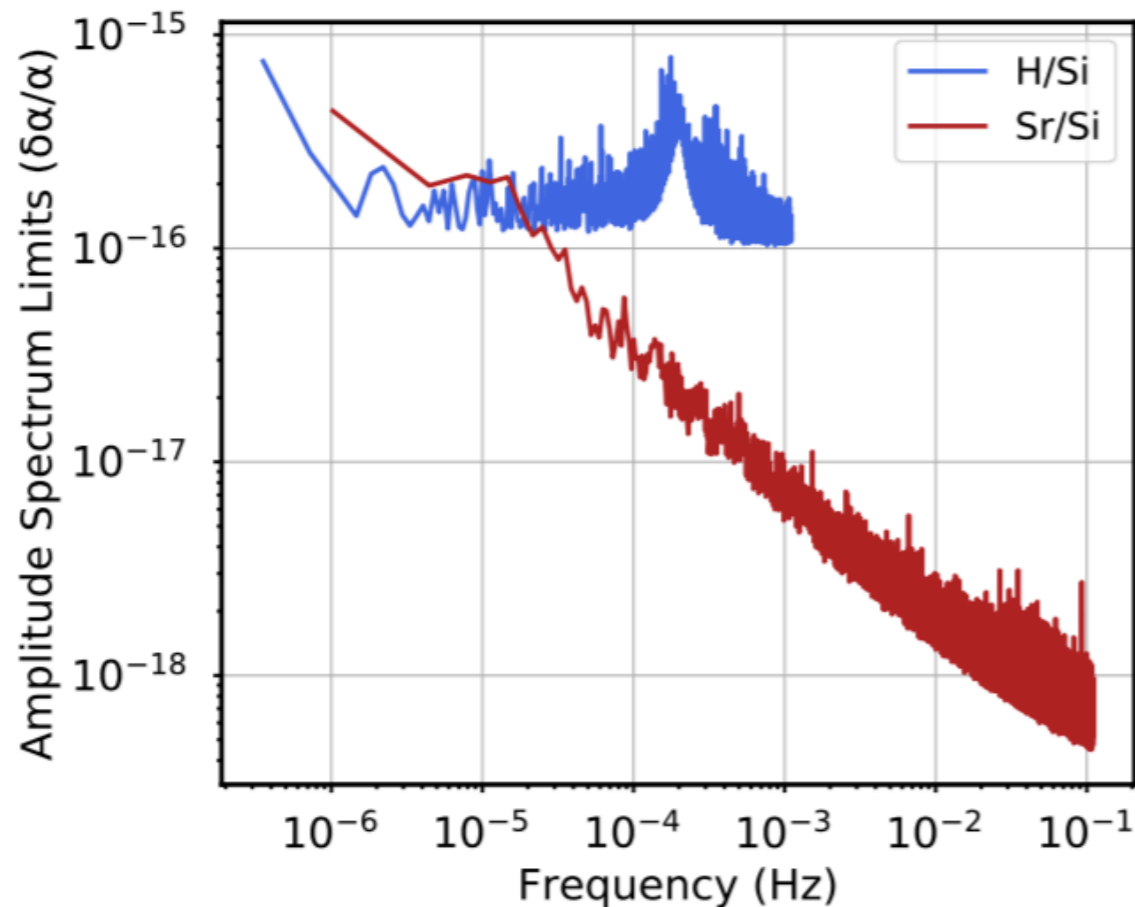
2. Sample the deviations  $\frac{\delta(f_A/f_B)}{f_A/f_B}$  over time
3. Obtain the power spectrum (periodogram)
4. Compare to known noise model (e.g. clock stability)



# SEARCHING FOR ULDM WITH ATOMIC CLOCKS

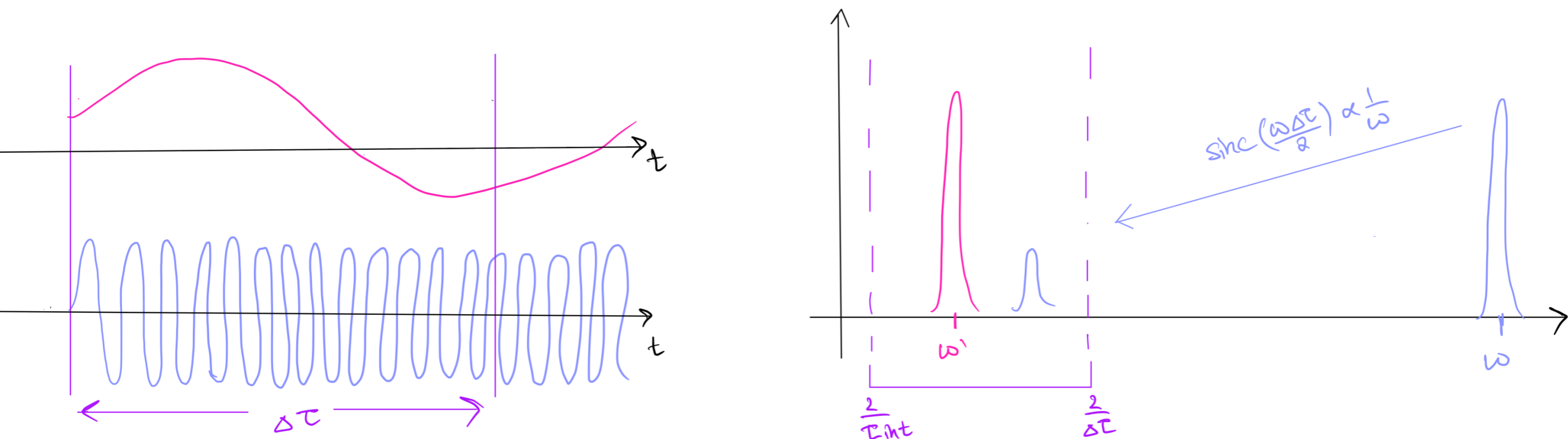
- Bounds on variations of fundamental constants translated to bounds on DM couplings to photons and to electrons.

$$\frac{\delta\alpha}{\alpha}(\omega = m_\phi) = - \frac{d_e}{m_\phi} \frac{\sqrt{2\rho_{DM}}}{M_{pl}}$$



# RAPIDLY OSCILLATING DARK MATTER

- “Heavy” ultralight DM  $m_\phi \gtrsim 10^{-12} \text{ eV}, \nu_\phi \gtrsim 100 \text{ Hz}$
- Theoretically motivated - relaxion DM. A. Banerjee et al., [1810.01889].
- **Blind spot** for atomic probes of temporal variations of fundamental constants due to rapid oscillations.
- Long averaging times ( $> 1 \text{ s}$ ) significantly lower the sensitivity to signals at frequencies above the Nyquist rate ( $> 1 \text{ Hz}$ )



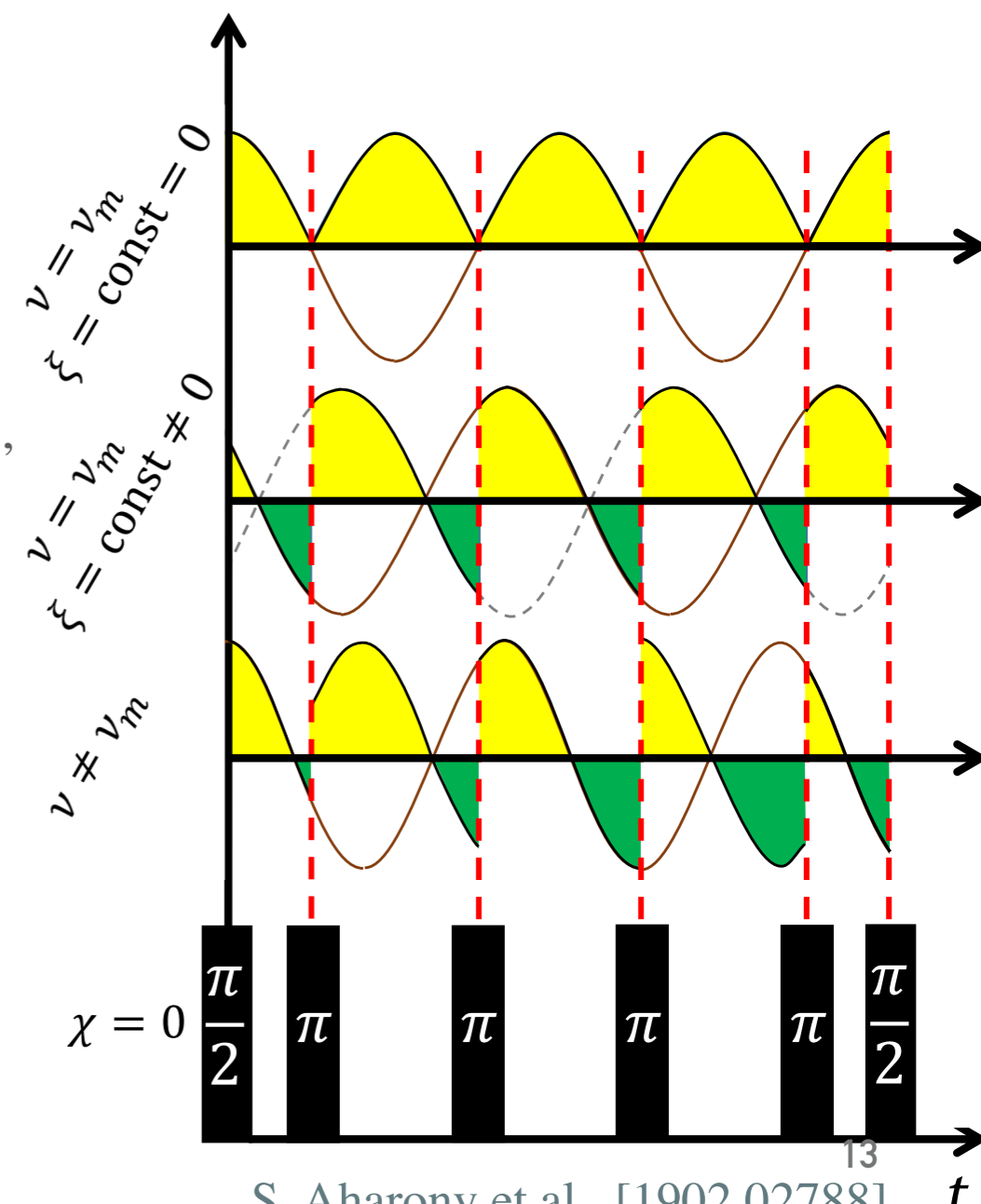
# OBSERVING RAPIDLY OSCILLATING DARK MATTER

- **Dynamical Decoupling** - echo pulses applied at  $\nu_m \gg \nu_s$
- Effectively - amplitude modulation of the signal by an alternating sign

$$\psi(\tau, n, \xi) = \int_{-\infty}^{\infty} G(t, \tau, n) \Delta f(t, \nu, \xi) dt.$$

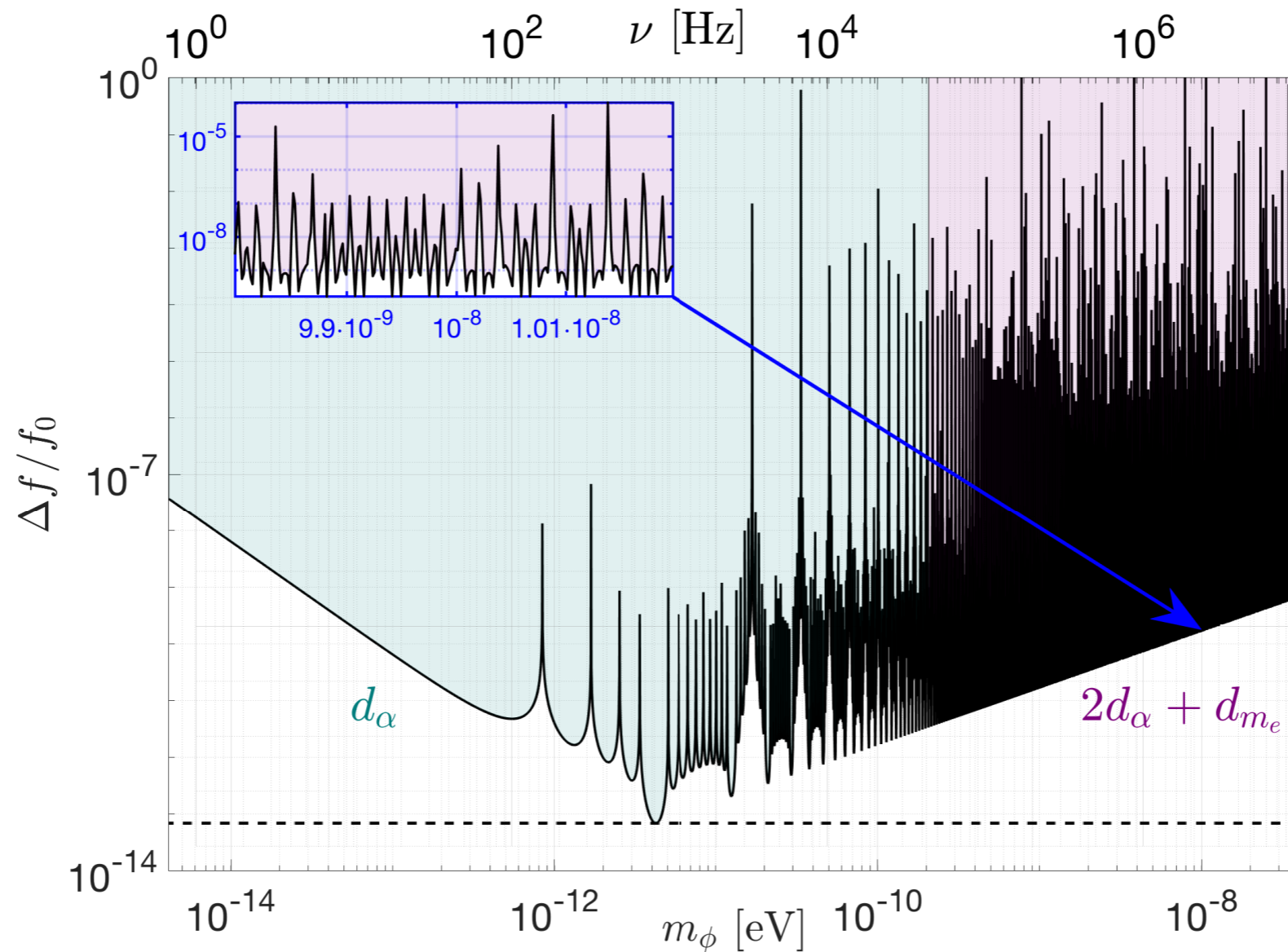
$$G(t, \tau, n) = \text{rect}\left(\frac{t}{2n\tau}\right) \times \left[ \Theta(t) + 2 \sum_{k=1}^{\infty} (-1)^k \Theta\left(t - (2k-1)\tau\right) \right],$$

- Thus, effects at the modulation frequency are accumulated, allows for long interrogation times.



# PROBING RAPIDLY OSCILLATING DM USING DD

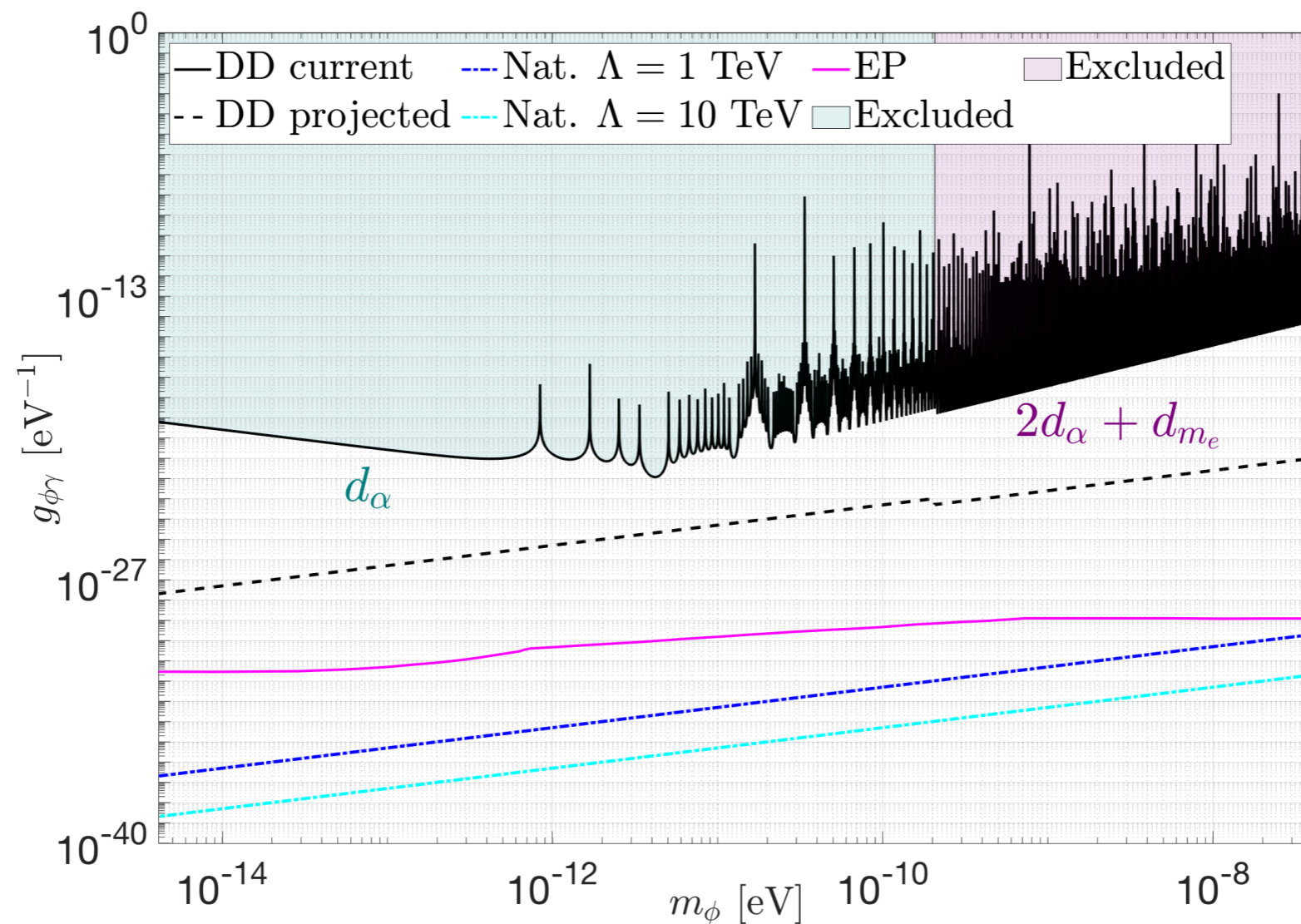
$$\frac{\Delta f(\nu)}{f_0} = \frac{\delta f_{\text{ion}}(\nu) - \delta f_{\text{laser}}(\nu)}{f_0} = (2 - F(\nu)) \frac{\delta\alpha}{\alpha} + (1 - F(\nu)) \frac{\delta m_e}{m_e}$$



First model independent bounds on variations of  $\alpha, m_e$  up to MHz frequencies.

# PROBING RAPIDLY OSCILLATING DM USING DD

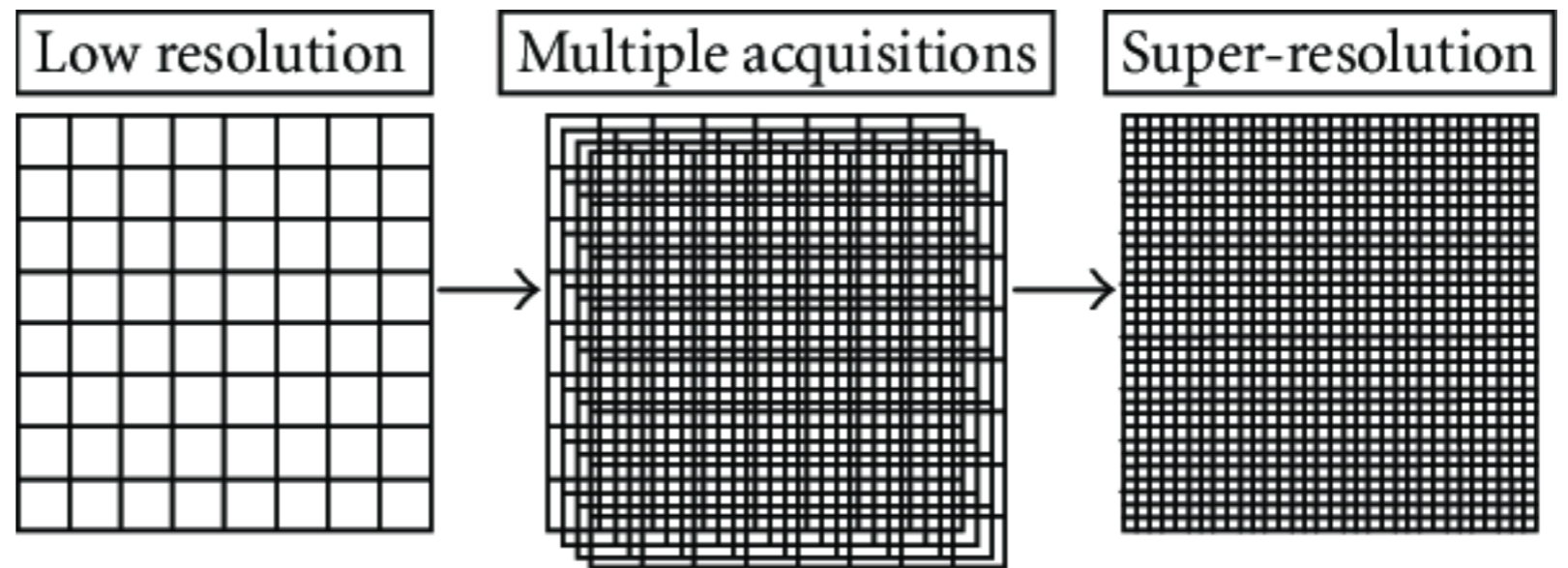
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First model independent bounds on variations of  $\alpha, m_e$  up to MHz frequencies.

# CAN WE DO BETTER?

- DD - resonance search, requires scanning over modulation frequency.
- Can we apply super-resolution/sub-Nyquist sampling techniques to enhance the detectability of rapidly oscillating DM?





# CAN WE USE SUPER-RESOLUTION/ SUB-NYQUIST SAMPLING?

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- Open questions for discussion:
  - Can we use super-resolution techniques to enhance the detectability of rapidly oscillating DM signals?
    - Single probe - different sampling method?
    - Many probes - sub-pixel shifts?
  - Are there better techniques to search for a sinusoidal signal with unknown frequency?
  - What are the experimental factors limiting the sensitivity of super-resolution techniques?

**THANK YOU!**

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# BACKUP SLIDES

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# ULTRALIGHT SCALAR DARK MATTER

- ▶ Ultralight DM ( $m < 0.1$  eV) - large occupation number in the galaxy

$$N = n \cdot \lambda_{dB}^3 = \frac{\rho_{DM}}{m^4 v^3} \approx \left( \frac{\rho_{DM}}{0.3 \text{ GeV/cm}^3} \right) \left( \frac{0.1 \text{ eV}}{m} \right)^4$$

- ▶ Can be described as a macroscopic classical field (much like a laser/BEC/superfluid)

$$\phi(t) = A \cos(m_\phi t)$$

- ▶ Energy density

$$\rho_{DM} = T_{00} \approx \frac{1}{2} m_\phi^2 A^2$$

- ▶ Ultralight scalar DM is the coherently oscillating as

$$\langle \phi(t) \rangle \simeq \frac{\sqrt{2\rho_{DM}}}{m_\phi} \cos(m_\phi t)$$

# RAPIDLY OSCILLATING DARK MATTER

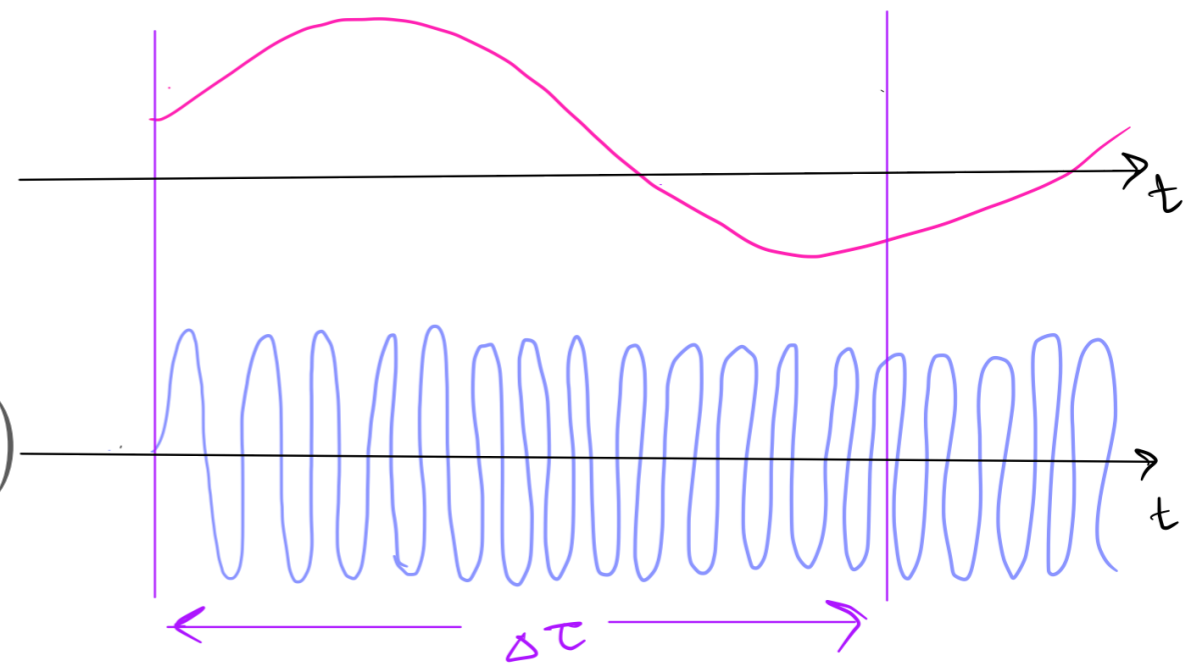
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$$\frac{\delta(f_A/f_B)}{f_A/f_B} = \frac{1}{\Delta\tau} \frac{g_{\phi\gamma} \sqrt{2\rho_{\text{DM}}} \Delta n_{A,B}}{m_\phi} \int_{t_n - \frac{\Delta\tau}{2}}^{t_n + \frac{\Delta\tau}{2}} \sin(m_\phi t) dt$$

$$= \frac{g\sqrt{2\rho_{\text{DM}}} \Delta n_{A,B}}{m_\phi} \text{sinc}\left(\frac{m_\phi \Delta\tau}{2}\right) \sin(m_\phi t_n)$$

additional  $\frac{1}{m_\phi}$  suppression

for  $m_\phi \gg \frac{2}{\Delta\tau}$

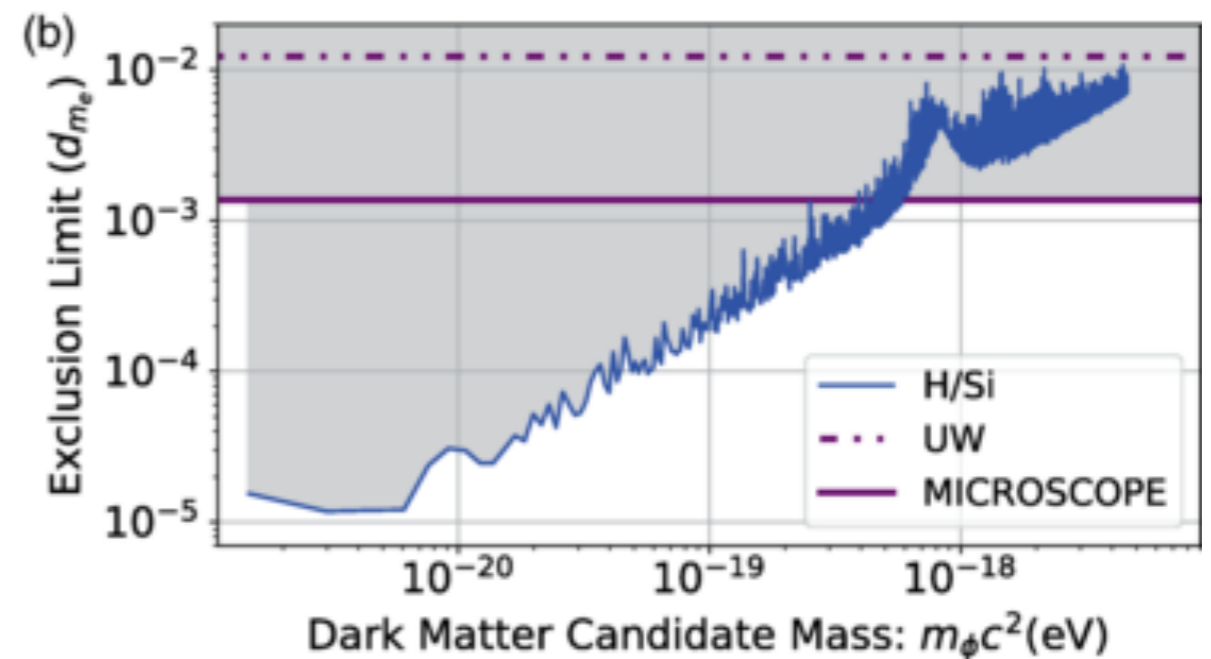
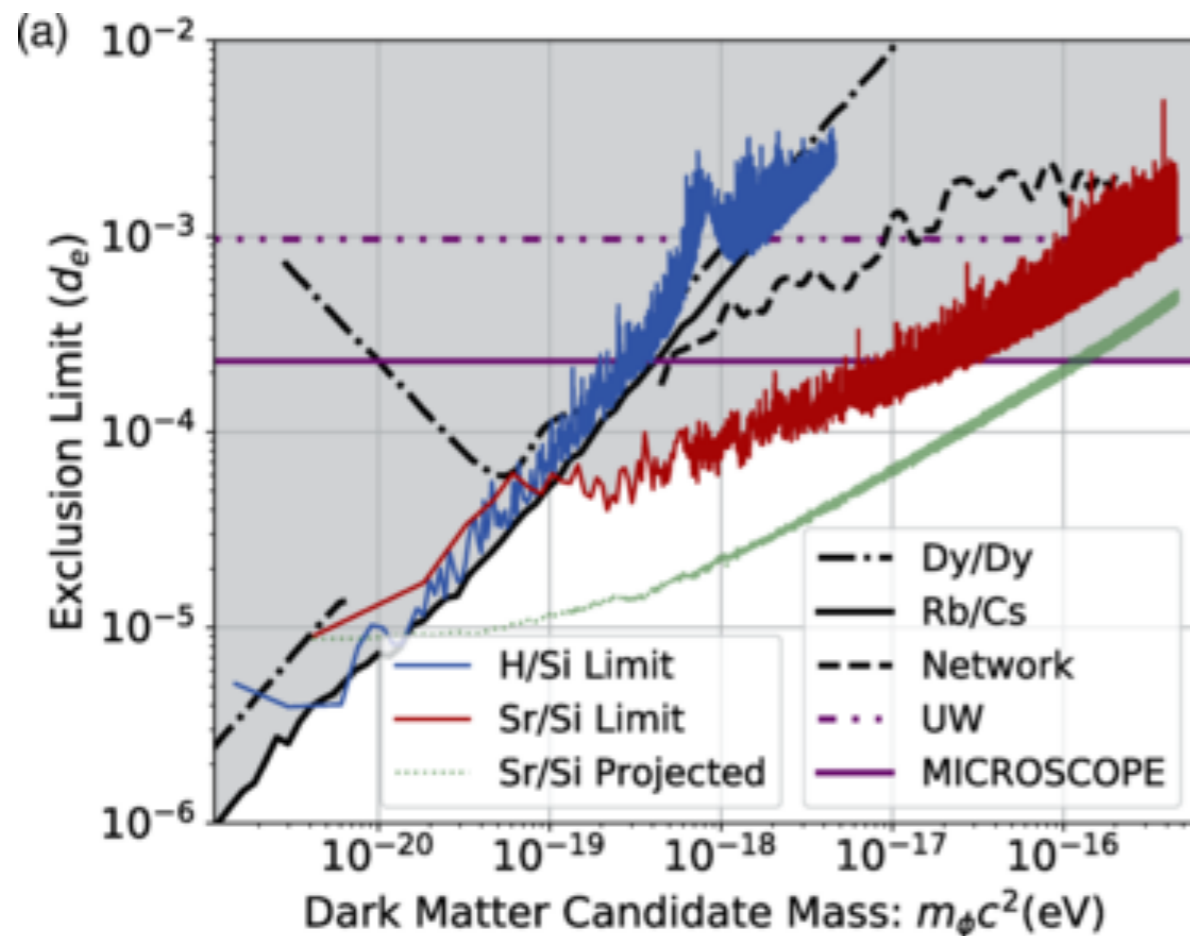


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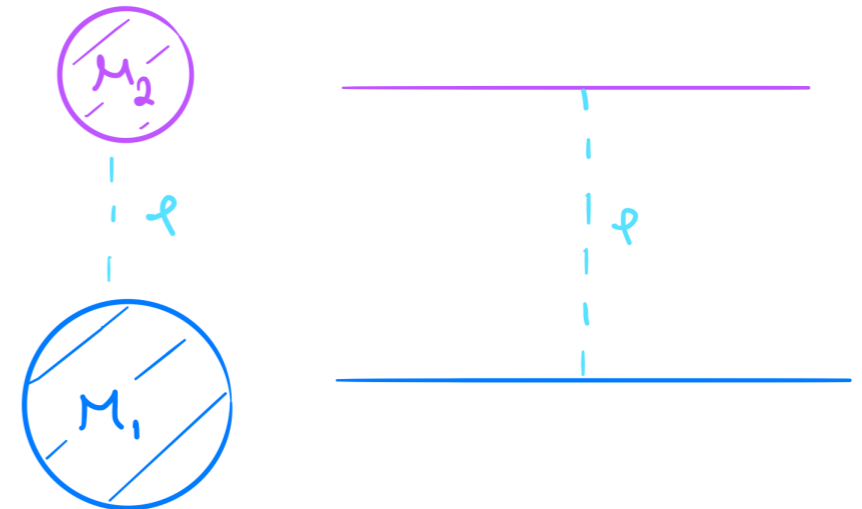
$$\frac{\delta\alpha}{\alpha}(\omega = m_\phi) = -\frac{d_e}{M_{pl}} \frac{\sqrt{2\rho_{DM}}}{m_\phi}$$

$$\frac{\delta m_e}{m_e}(\omega = m_\phi) = \frac{d_{m_e}}{M_{pl}} \frac{\sqrt{2\rho_{DM}}}{m_\phi}$$



# TESTS OF FIFTH-FORCE AND EP

- New scalar mediates a force (interaction) between neutral masses



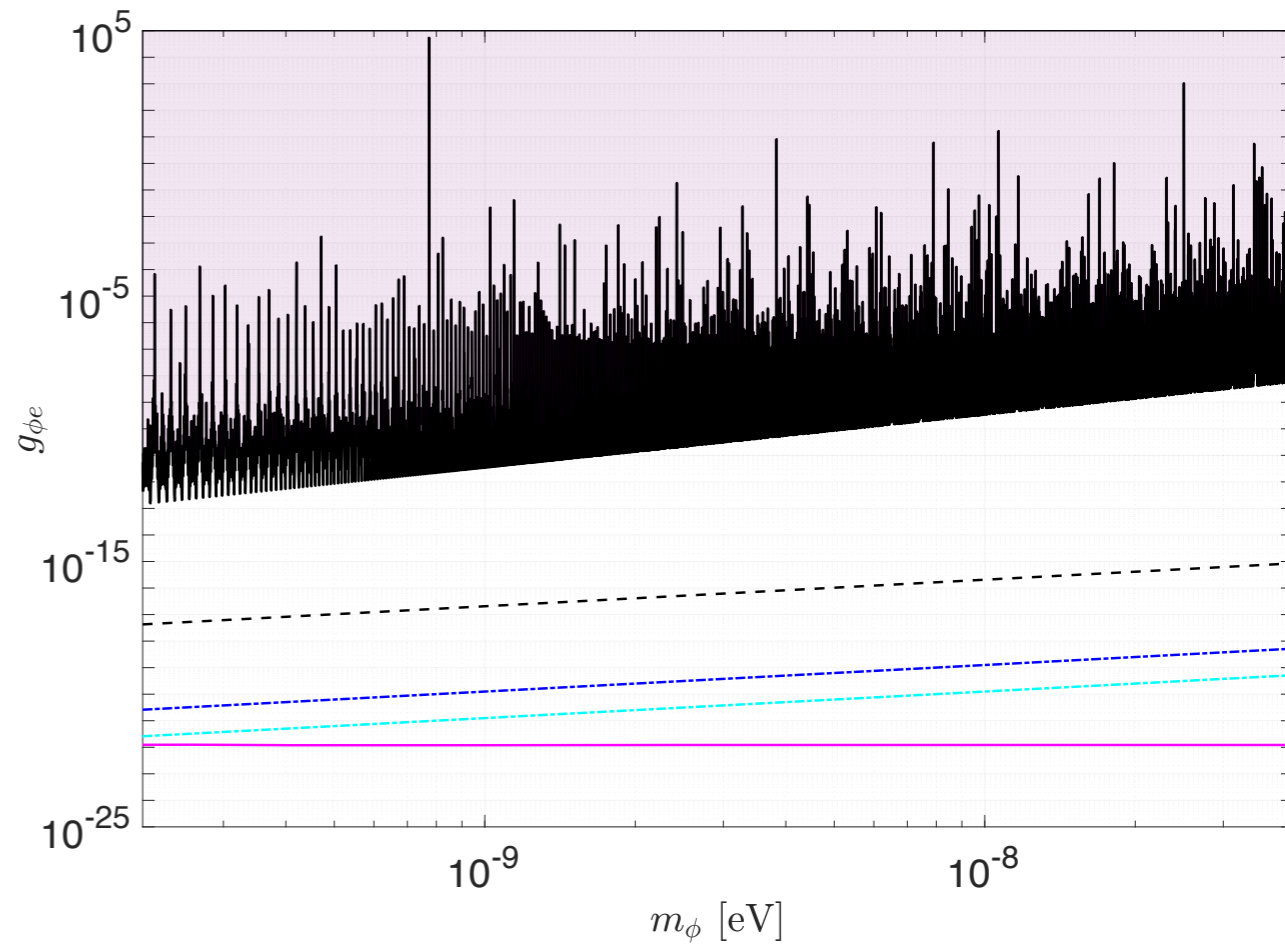
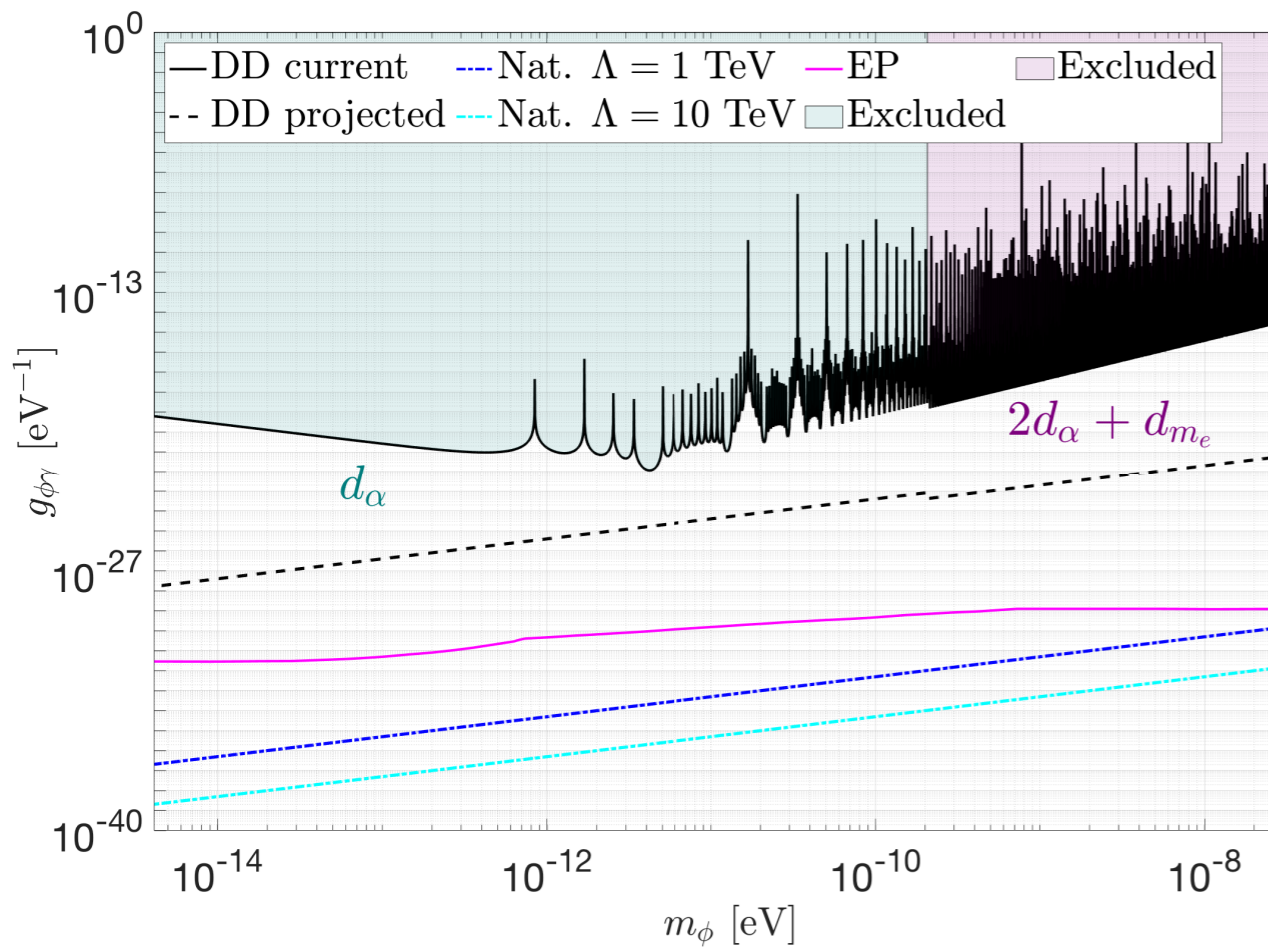
- Searches for deviations from gravity
  - massive force carrier - Yukawa potential

$$V_{\phi} \propto \frac{e^{-m_{\phi}r}}{r}$$

- Non-universal couplings to masses of different compositions - violation of the universality of free-fall (Einstein's Equivalence Principle)

$$\eta = \frac{a_A - a_B}{a_A + a_B} \propto \frac{1}{G} \left( g_{\phi e} Q_{m_e}^C + g_{\phi \gamma} Q_e^C \right) \left( g_{\phi e} \left( Q_{m_e}^A - Q_{m_e}^B \right) + g_{\phi \gamma} \left( Q_e^A - Q_e^B \right) \right)$$

# COMPARISON – FIFTH FORCE VS. FUND. CONSTANTS





# QUADRATICALLY COUPLED DM

- Higher order interactions of the DM candidate with the SM

$$\mathcal{L} = g_{\phi^2 e} \phi^2 m_e \bar{e} e$$

- Variations of  $m_e$

$$\frac{\delta m_e}{m_e} = g_{\phi^2 e} \frac{\rho_{DM}}{m_\phi^2} \cos(2m_\phi t)$$

- **5F/EP** - gains DM density sensitivity, different r-dependance

$$V(r) \propto \phi^2 \sim \frac{\rho_{DM}}{m_\phi^2} \left( 1 - s_A \frac{GM_A}{r} \right)^2$$

$$\eta = \frac{a_A - a_B}{a_A + a_B} \propto \frac{\rho_{DM}}{Gm_\phi^2} \left( g_{\phi^2 e} Q_{m_e}^C + g_{\phi^2 \gamma} Q_e^C \right) \left( g_{\phi^2 e} \left( Q_{m_e}^A - Q_{m_e}^B \right) + g_{\phi^2 \gamma} \left( Q_e^A - Q_e^B \right) \right)$$

# COMPARISON – QUADRATICALLY COUPLED DM

