

# The CMS Precision Proton Spectrometer Project for the HL-LHC



Low-x Workshop  
La Biodola, Isola d'Elba

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on behalf of  
The CMS Collaboration



The Compact Muon Solenoid Experiment

## CMS Note

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26 November 2020 (v3, 09 December 2020)

## The CMS Precision Proton Spectrometer at the HL-LHC – Expression of Interest

The CMS Collaboration

### Abstract

The CMS Collaboration intends to pursue the study of central exclusive production (CEP) events,  $pp \rightarrow pXp$ , at the High-Luminosity LHC (HL-LHC) by means of a new near-beam proton spectrometer. In CEP events, the state  $X$  is produced at central rapidities, and the scattered protons do not leave the beam pipe. The kinematics of  $X$  can be fully reconstructed from that of the protons, which gives access to final states otherwise not visible. CEP allows unique sensitivity to physics beyond the standard model, e.g. in the search for anomalous quartic gauge couplings, axion-like particles, and in general new resonances.

CMS has been successfully operating the Precision Proton Spectrometer (PPS) since 2016; PPS started as a joint CMS and TOTEM project, and then evolved into a standard CMS subsystem. The present document outlines the physics interest of a new near-beam proton spectrometer at the HL-LHC, and explores its feasibility and expected performance. The document has been edited by the members of the PPS group and builds on their experience in the construction and operation of PPS.

Discussion with the machine groups has led to the identification of four locations suitable for the installation of movable proton detectors: at 196, 220, 234, and 420 m from the interaction point, on both sides (in this document these locations always imply both sides, unless otherwise noted). The locations at 196, 220, and 234 m can be instrumented with Roman Pot devices similar to the ones presently used. The 420 m location requires a bypass cryostat (which has been developed for other locations in the LHC) and a movable detector vessel approaching the beam from between the two beam pipes.

Acceptance studies indicate that having the beams cross in the vertical plane at the interaction point, as implemented after Long Shutdown 3, is vastly preferable over the present horizontal crossing. This gives access to centrally produced states  $X$  in the mass range 133 GeV–2.7 TeV with the stations at 196, 220, and 234 m. The mass range becomes 43 GeV–2.7 TeV if the 420 m station is included, which makes it possible to study central exclusive production of the 125 GeV Higgs boson. This is a major improvement with respect to the current mass range of 350 GeV–2 TeV.

The radiation background has also been studied. Radiation hardness is required for all components in the tunnel. Service work during short technical stops will not be possible. The irradiation dose rate will be very strongly peaked near the beam. Detectors should therefore be vertically shifted with a

Since 2018:

- HL-LHC Studies based on present-day PPS experience
- regular interactions with machine integration and optics teams.
- presentations to HL-LHC coordination group

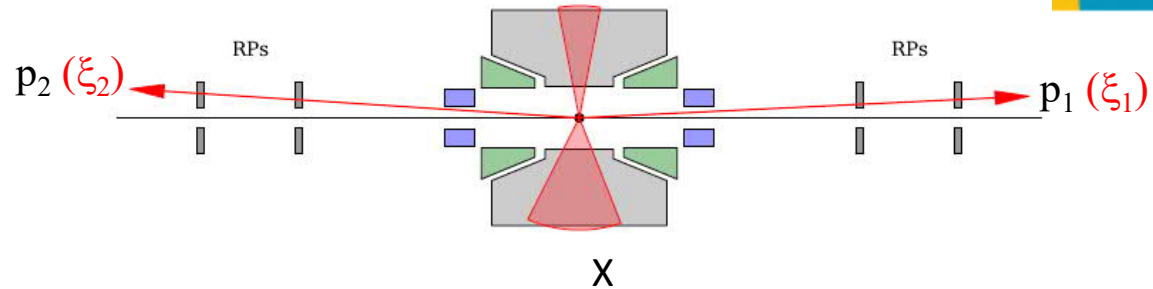
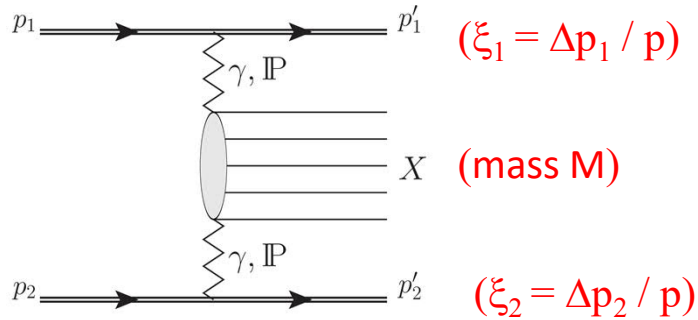
2020:

- September: presentation in the LHCC focus session on CMS PPS / Forward Physics at HL-LHC
- December: publication as CMS NOTE-2020/008, available on CDS and arXiv:

<https://cds.cern.ch/record/2750358>

<http://arxiv.org/abs/2103.02752>

# Central Exclusive Production (Reminder)



## Measurable kinematic variables of the leading protons:

- **Fractional momentum losses ( $\xi_1, \xi_2$ ) via proton tracking**  
 → Reconstruction of mass and rapidity of central system

$$M_X^2 = \xi_1 \xi_2 s$$

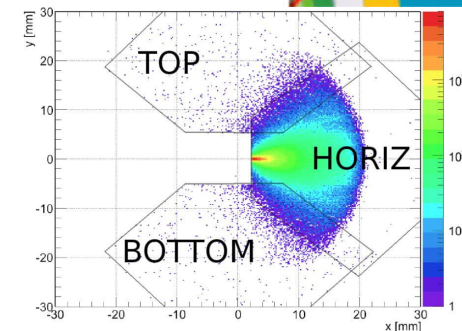
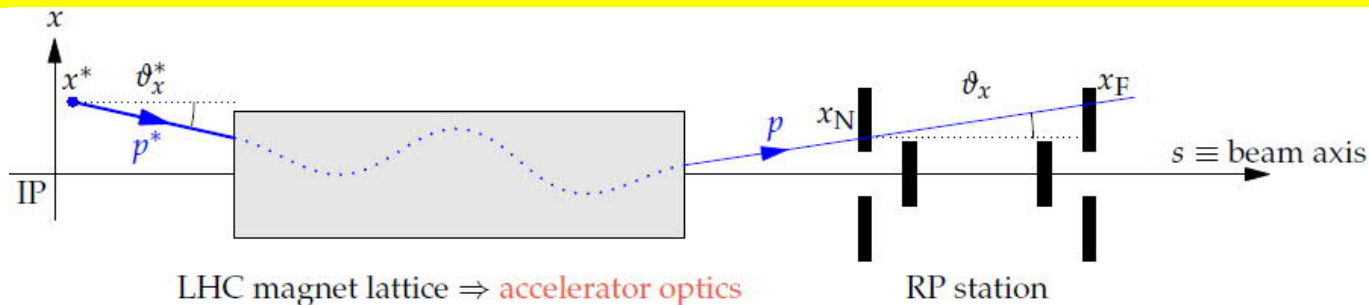
$$y_X = \frac{1}{2} \ln \frac{\xi_1}{\xi_2}$$

- **Transverse momenta ( $\mathbf{p}_{T,1}, \mathbf{p}_{T,2}$ ) via proton tracking**  
 → momentum balance with central system useful for event selection:

$$\mathbf{p}_{T,X} + \mathbf{p}_{T,1} + \mathbf{p}_{T,2} = \mathbf{0}$$

- **Longitudinal vertex position via proton time of flight (ToF)**  
 → important for resolving pileup (up to  $\mu = 200$  at the HL-LHC)

# Proton Measurements in Space and Time



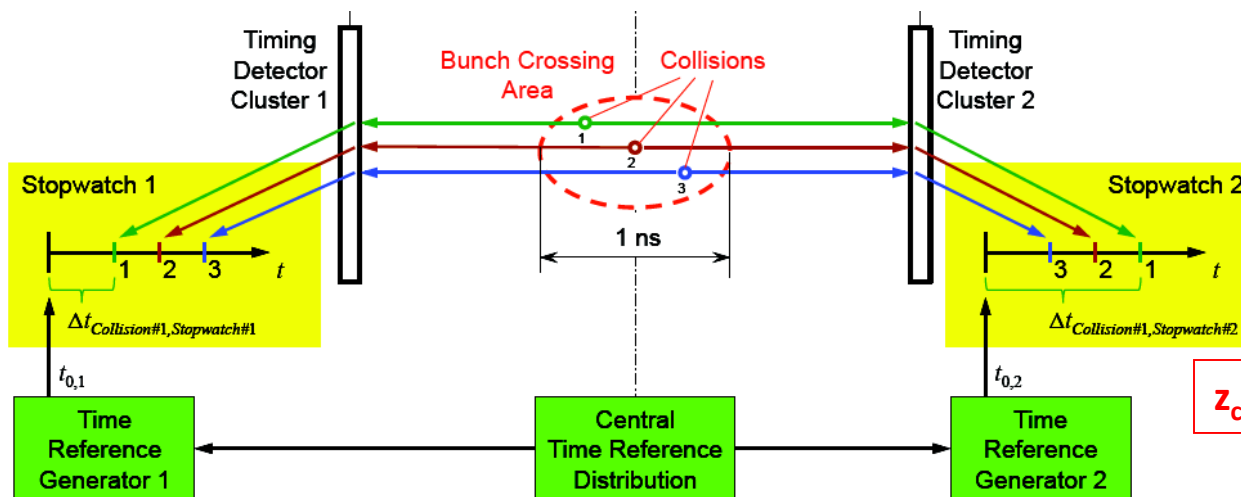
Proton kinematics in RP

$$\begin{pmatrix} x \\ \Theta_x \\ y \\ \Theta_y \\ \Delta p/p \end{pmatrix}_{\text{RP}} = \underbrace{\begin{pmatrix} v_x & L_x & 0 & 0 & D_x \\ v'_x & L'_x & 0 & 0 & D'_x \\ 0 & 0 & v_y & L_y & 0 \\ 0 & 0 & v'_y & L'_y & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix}}_{\text{Product of all lattice element matrices}} \begin{pmatrix} x^* \\ \Theta_x^* \\ y^* \\ \Theta_y^* \\ \Delta p/p \end{pmatrix}_{\text{IP5}}$$

Values at IP5 to be reconstructed

Product of all lattice element matrices

## Longitudinal Vertex Position Measurement via Time-of-Flight Difference



$\mu \leq 200 !$

$$z_{\text{collision}} = \frac{1}{2} (t_{\text{stopwatch 2}} - t_{\text{stopwatch 1}})$$



# Search for Detector Locations (1)



LS3: Long Straight Section in IR5 to be redesigned, all present Roman Pots removed  
→ **new spectrometer to be built**

Objective from physics programme:

Maximise mass acceptance for centrally produced states measured via leading protons.

**Minimum mass:**  $M_{\min} = |\xi|_{\min} \sqrt{s}$ ,  $\xi \equiv \frac{\Delta p_{\text{proton}}}{p_{\text{proton}}} = \frac{x_{\text{track}}}{D_x}$  ← track displacement from beam @ detector  
← dispersion @ detector

$$M_{\min} = \frac{d_{\min}}{D_x} \sqrt{s}$$

Closest approach of detector to beam:  $d_{\min} = (n_{\text{TCT}} + 3) \sigma_x + 0.3 \text{ mm}$  (collimation hierarchy)  
→ look for locations with small  $\frac{\sigma_x}{D_x}$  ( $\sigma_x = \text{beam width}$ )

**Maximum mass:**  $M_{\max} = |\xi|_{\max} \sqrt{s}$  :

determined by the tightest aperture limitation (usually a TCL debris collimator):

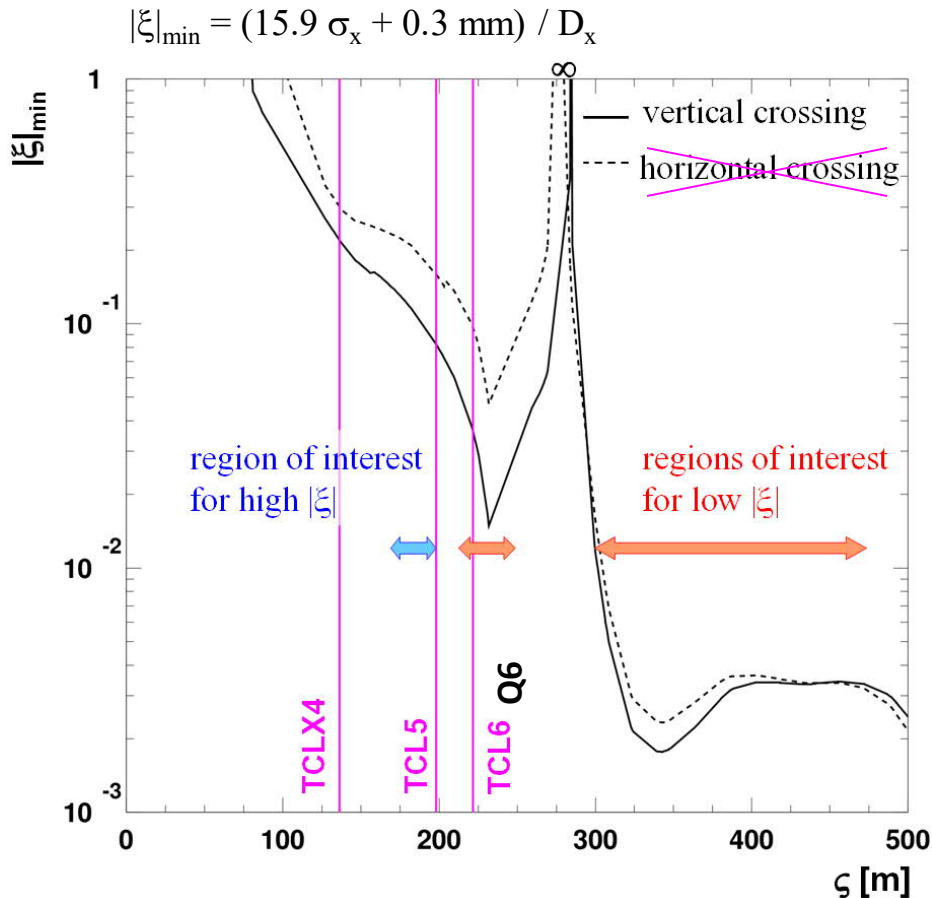
$$M_{\max} = \frac{d_{\text{TCL}}}{D_{\text{TCL}}} \sqrt{s}$$

→ look for locations just before TCL collimators

# Search for Detector Locations (2)



- HL-LHC optics version 1.3
- for (crossing-angle  $\alpha/2, \beta^*$ ) = (250  $\mu$ rad, 15 cm)
- Roman Pots @ (12.9 + 3)  $\sigma$  + 0.3 mm



## Crossing plane in IP5:

Both orientations (horizontal, vertical) studied and discussed in the EoI  
→ strong preference for vertical crossing  
→ CMS request in December 2018

June 2020: machine decision for vertical crossing in IP5

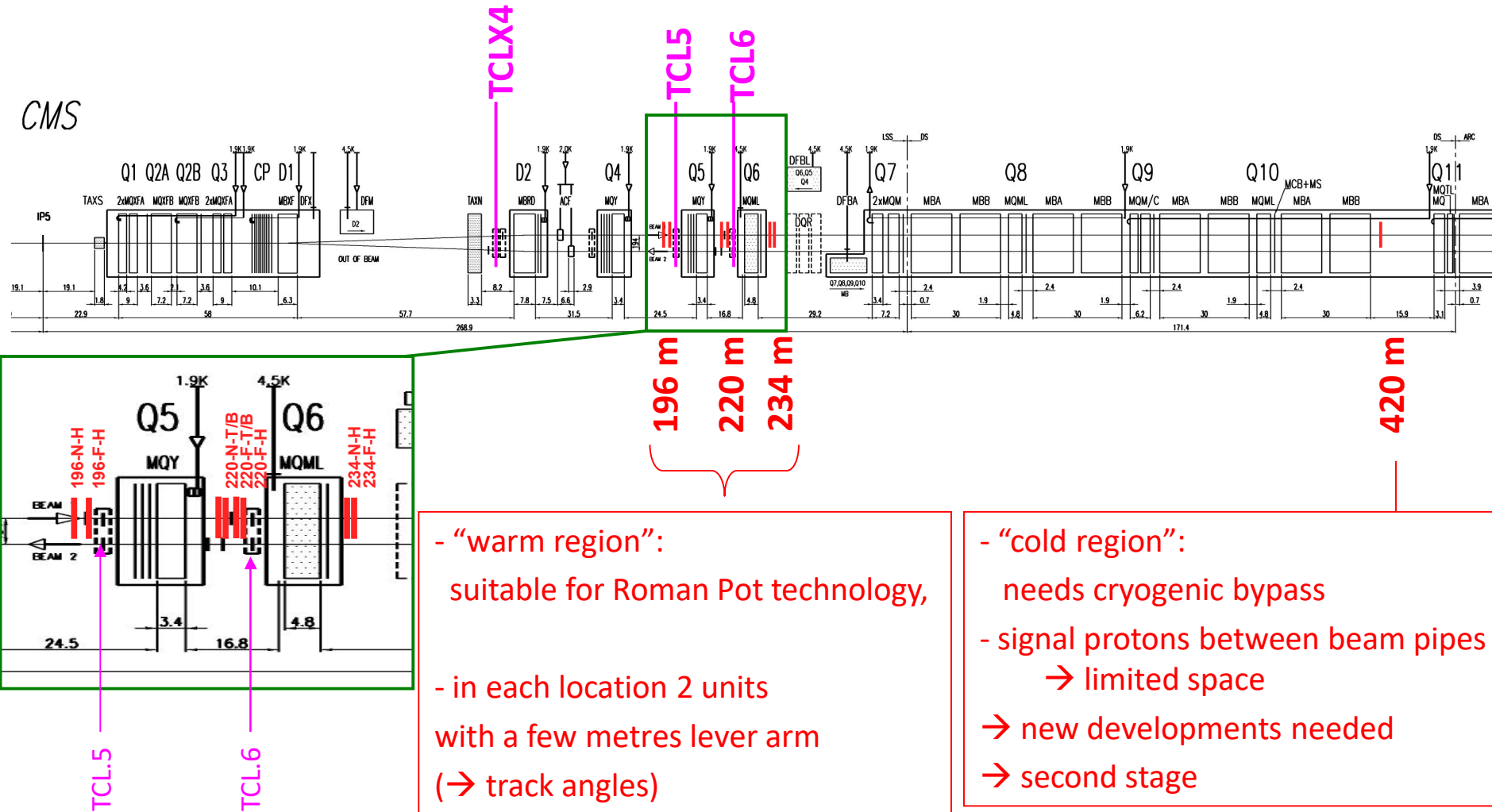
→ All figures in this presentation for vertical crossing



# Layout Overview with Proposed Stations



Free locations identified in discussions with the LHC layout team:



- “warm region”:  
 suitable for Roman Pot technology,  
 - in each location 2 units  
 with a few metres lever arm  
 (→ track angles)

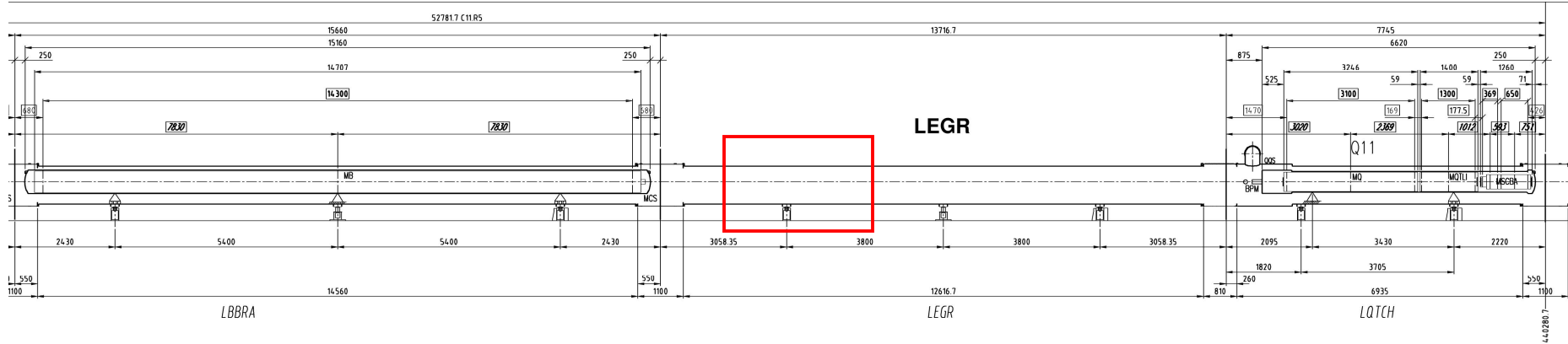
- “cold region”:  
 needs cryogenic bypass  
 - signal protons between beam pipes  
 → limited space  
 → new developments needed  
 → second stage

Feb. 2020: Tentative space reservations at the 28th HL-LHC Coordination Group Meeting

Dec. 2020: Consolidation of reservation: layout drawings with space holders in preparation



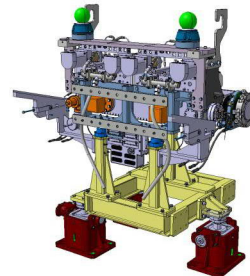
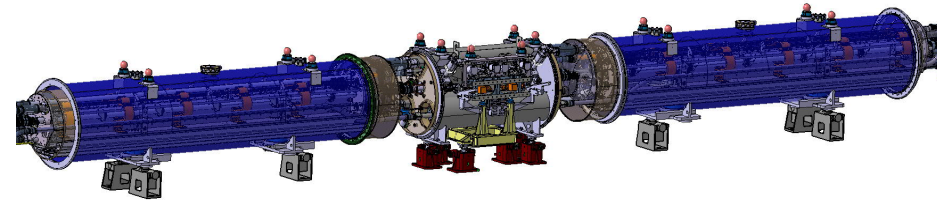
# The 420 m Station



- Region with an empty cryostat (“missing magnet”)
- Signal proton tracks are between the 2 beampipes (positive dispersion)
- Not suitable for present Roman Pot technology → needs special development

## Ideas:

- Reuse connection cryostat from TCLD integration or cryostat designed for the old FP420 project
- Detector vessel options:
  - mini Roman Pot
  - modified TCLD
  - moveable beampipe

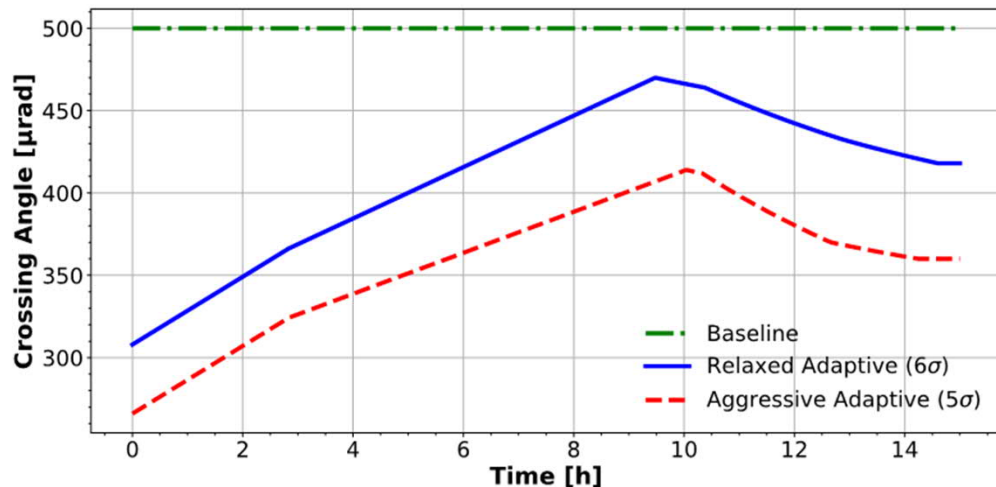


TCLD

# Evolution of Acceptance during a Fill



Luminosity levelling: concurrent variation of crossing-angle ( $\alpha/2$ ) and optics ( $\beta^*$ )



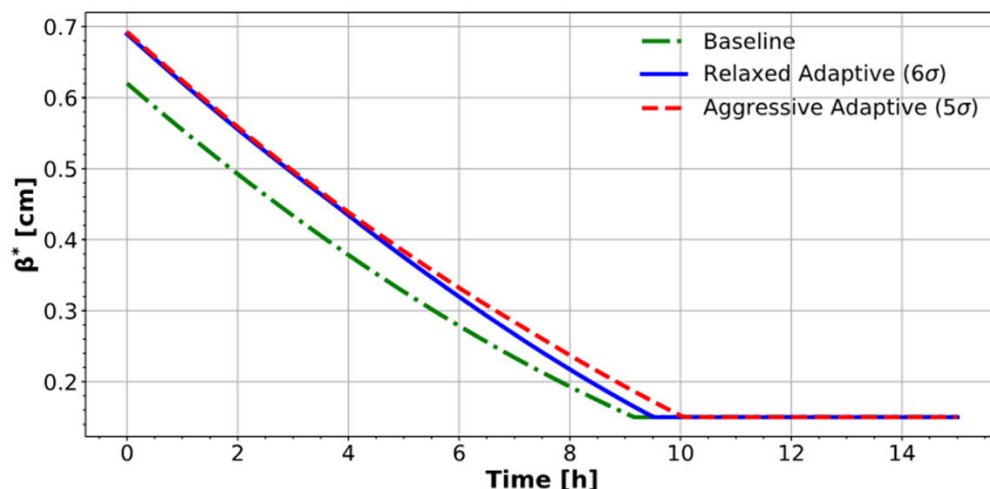
Crossing-angle in IP5 can – in principle – be horizontal (Runs 1 – 3) or vertical (HL-LHC)

Linear dependence of dispersion on X-angle:

$$D_x = D_x(0) - D'_x \alpha_x/2 \quad (\text{X-angle reduces } D_x \text{ !})$$

$$D_y = D'_y \alpha_y/2$$

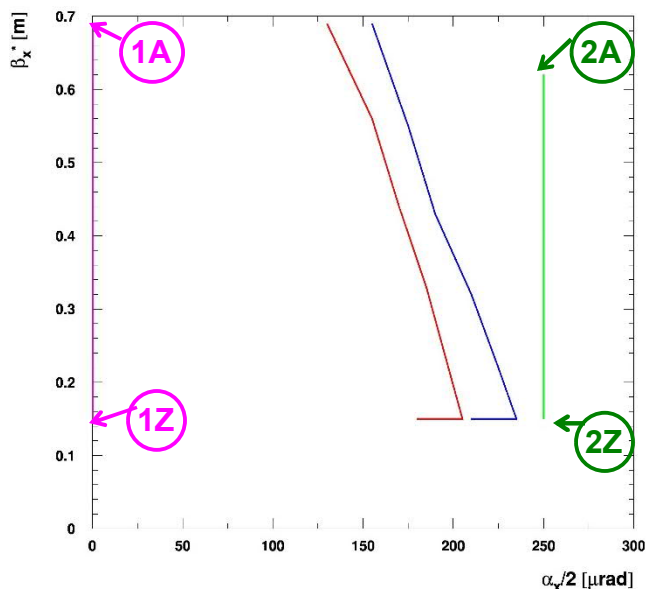
Acceptance depends mainly on  $D_x$ , less on  $D_y$ :  
 → choice of crossing plane very important



$\beta^*$  determines beam width,  
 hence RP distance

All performance parameters to be studied along a “levelling trajectory” in the ( $\alpha_x/2, \beta^*$ ) plane.

# Acceptance in the Mass – Rapidity Plane



Labels (1A), (1Z), (2A), (2Z)  
 = start and end points of **any vertical**  
 and the **simplest horizontal** trajectory

## Note on $p_T$ :

The M-y plot is for proton  $p_{T,1} = p_{T,2} = 0$   
 Fixed non-zero  $p_T$  would shift the contours.

For each point ( $\alpha_x/2, \beta_x^*$ ):

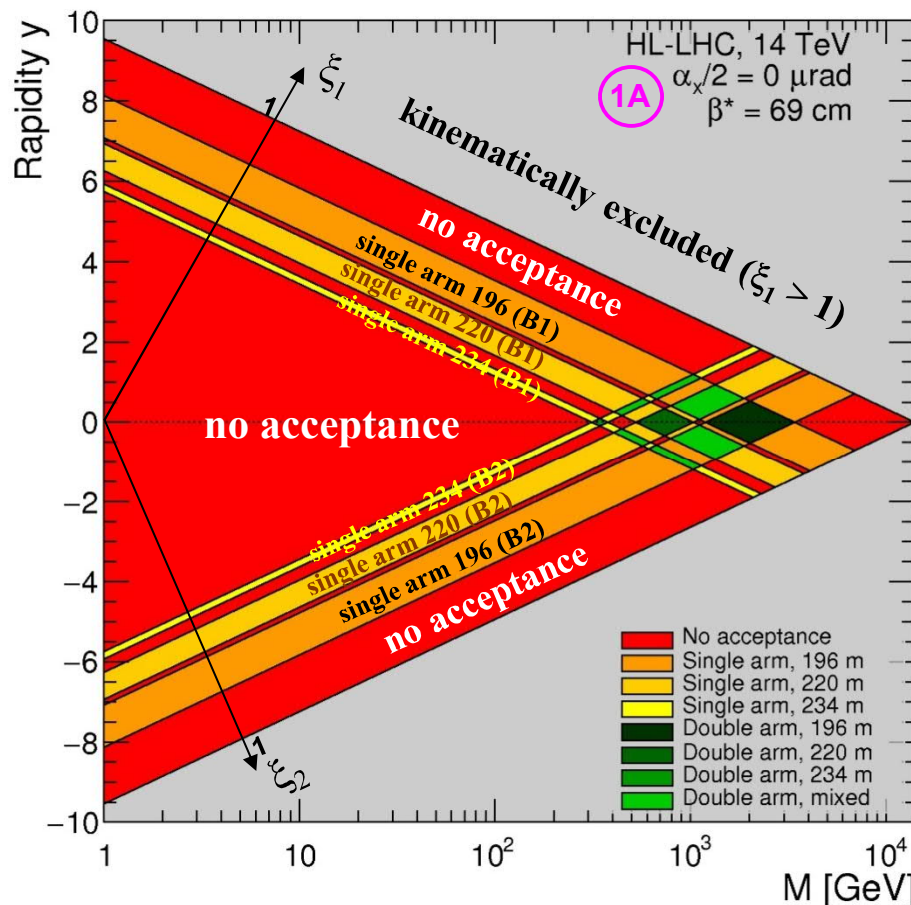
Acceptance for central exclusive events is defined in 2-dim space ( $\xi_1, \xi_2$ )  
 or equivalently – after basis rotation – in (M, y):

$$M^2 = \xi_1 \xi_2 s$$

$$\ln \frac{M}{\sqrt{s}} = \frac{1}{2} (\ln \xi_1 + \ln \xi_2)$$

$$y = \frac{1}{2} \ln \frac{\xi_1}{\xi_2}$$

$$y = \frac{1}{2} (\ln \xi_1 - \ln \xi_2)$$



# Acceptance in Mass – Rapidity Plane

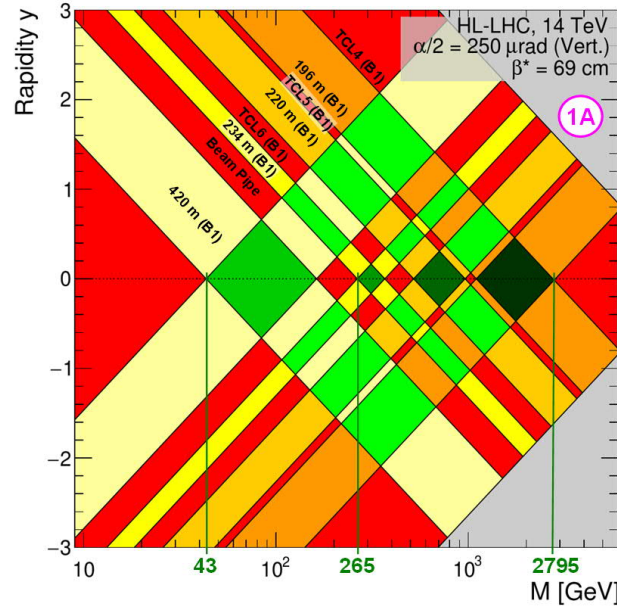


## Vertical crossing in IP5

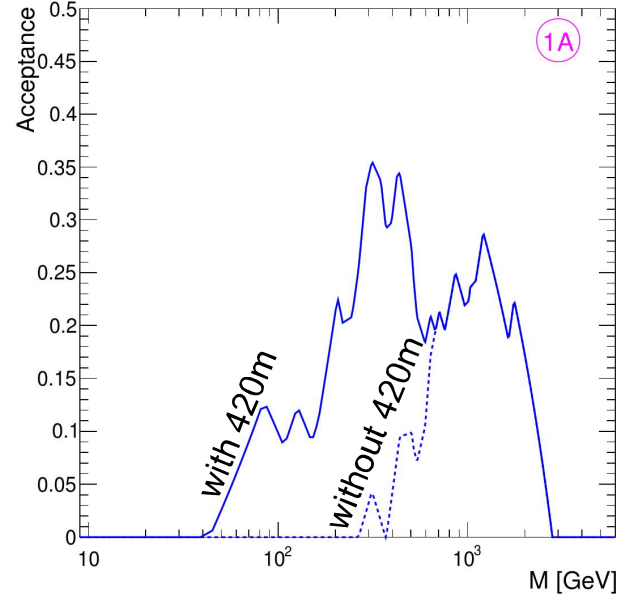
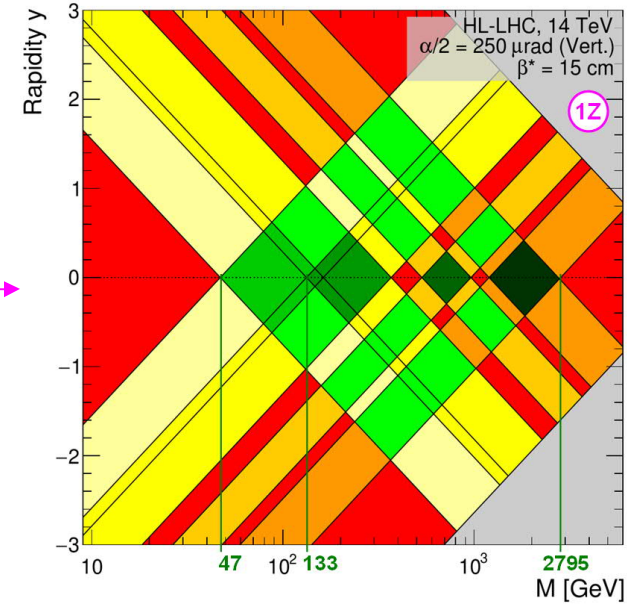
$$\ln \frac{M}{\sqrt{s}} = \frac{1}{2} (\ln \xi_1 + \ln \xi_2)$$

$$y = \frac{1}{2} (\ln \xi_1 - \ln \xi_2)$$

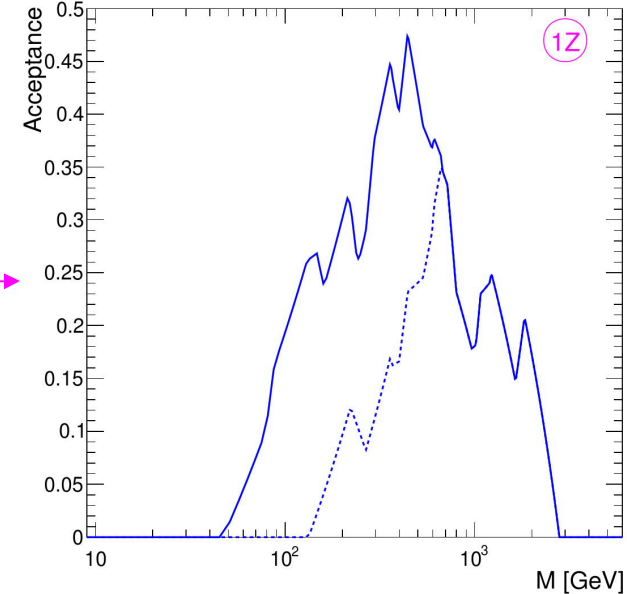
Projection on M  
assuming flat  
y distribution  
(in kinematically  
allowed region)



Fill →



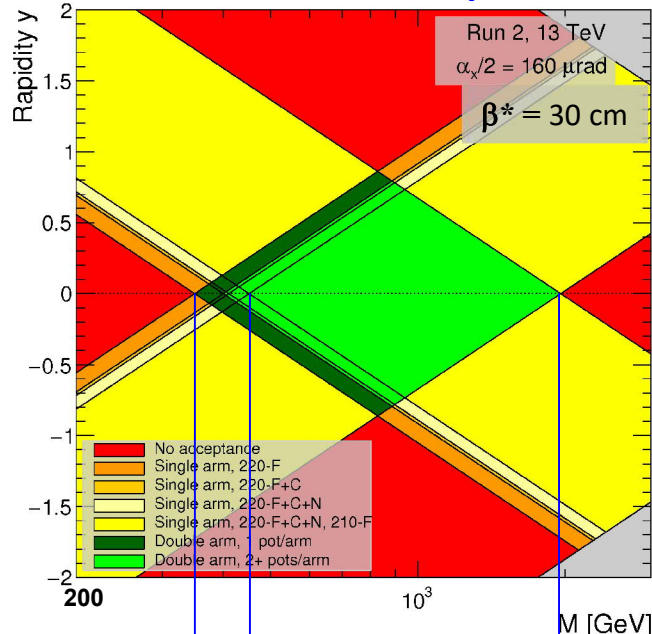
Fill →



# Comparison Mass-Rapidity Acceptance Run 2 / HL-LHC



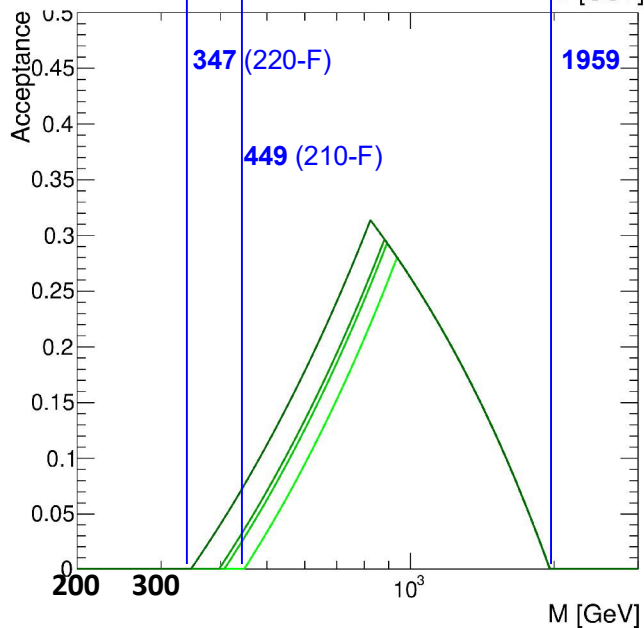
## 2018 Nominal Optics



## HL-LHC (vertical crossing):

without 420 m: 133 GeV – 2.7 TeV  
 with 420 m: 43 GeV – 2.7 TeV

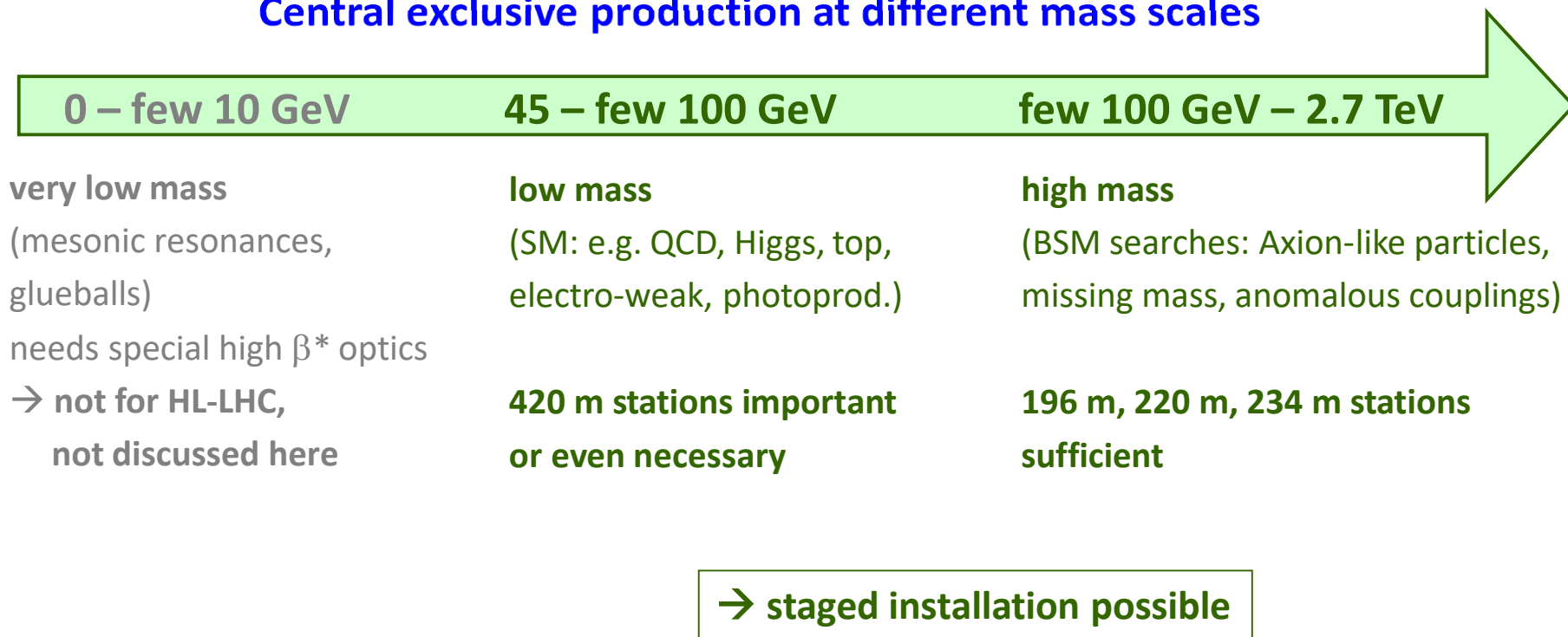
Physics programme allows  
 a staged installation (420 m later)



The physics section anticipates:

- integrated lumi: 300 fb<sup>-1</sup>/year [nominal ÷ ultimate peak: (5 ÷ 7.5) x 10<sup>34</sup> cm<sup>-2</sup> s<sup>-1</sup>]
- pileup multiplicity: nominal ÷ ultimate = 140 ÷ 200
- detector stations at +- 196 m, +-220 m, +-234 m, (+-420 m)
- vertical crossing-angle in IP5

## Central exclusive production at different mass scales



Most recent physics talk: Michael Pitt @ LHC Forward Physics Meeting (March 2021):  
<https://indico.cern.ch/event/955960/>



# Physics Examples

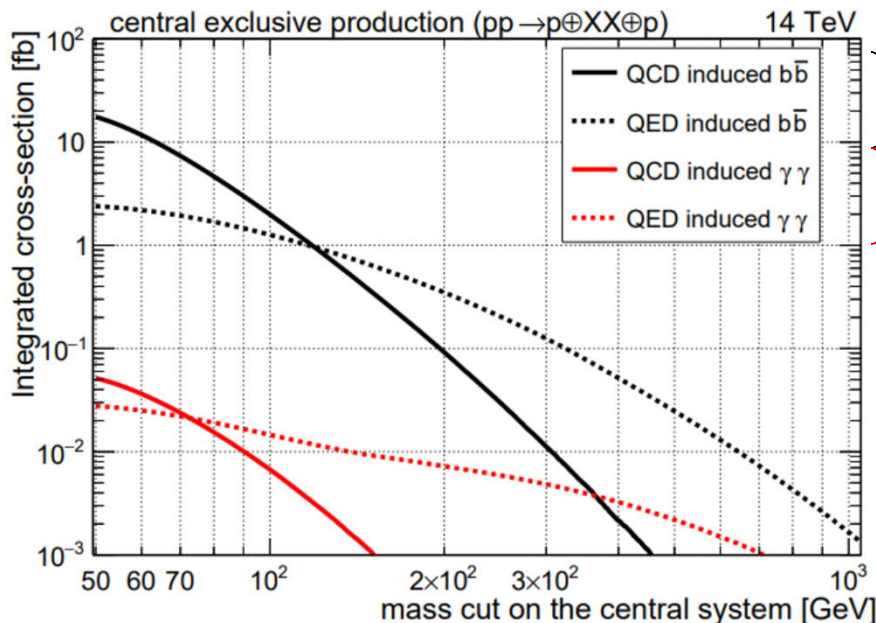


Low masses: QCD dominant

→ study exclusive jj  
(screening effects; bb as Higgs backg.)

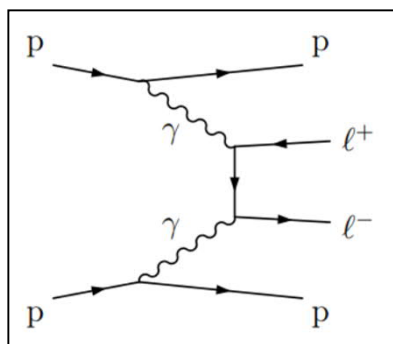
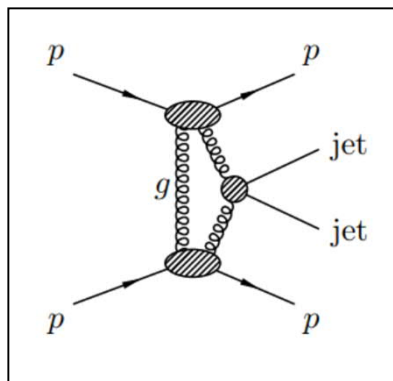
High masses: QED dominant

→ study heavy objects, anomalies



FPMC generator

SuperChic\_v4 generator



Fiducial cross-sections (2-arm tagged) of CEP SM processes @  $\sqrt{s} = 14$  TeV with and without the 420 m stations

[FPMC generator,  $p_T > 20$  GeV for all objects generated, survival prob. = 3% (QCD) and 90% (QED)]

Process	fiducial cross section [fb]			
	all stations		w/o 420	
	IP - IP	$\gamma - \gamma$	IP - IP	$\gamma - \gamma$
jj	$\mathcal{O}(10^6)$	60	$\mathcal{O}(10^4)$	2
$W^+W^-$	—	37	—	15
$\mu\mu$	—	46	—	1.3
$t\bar{t}$	—	0.15	—	0.1
H	0.6	0.07	0	0
$\gamma\gamma$	—	0.02	—	0.003

needs 420 m stations

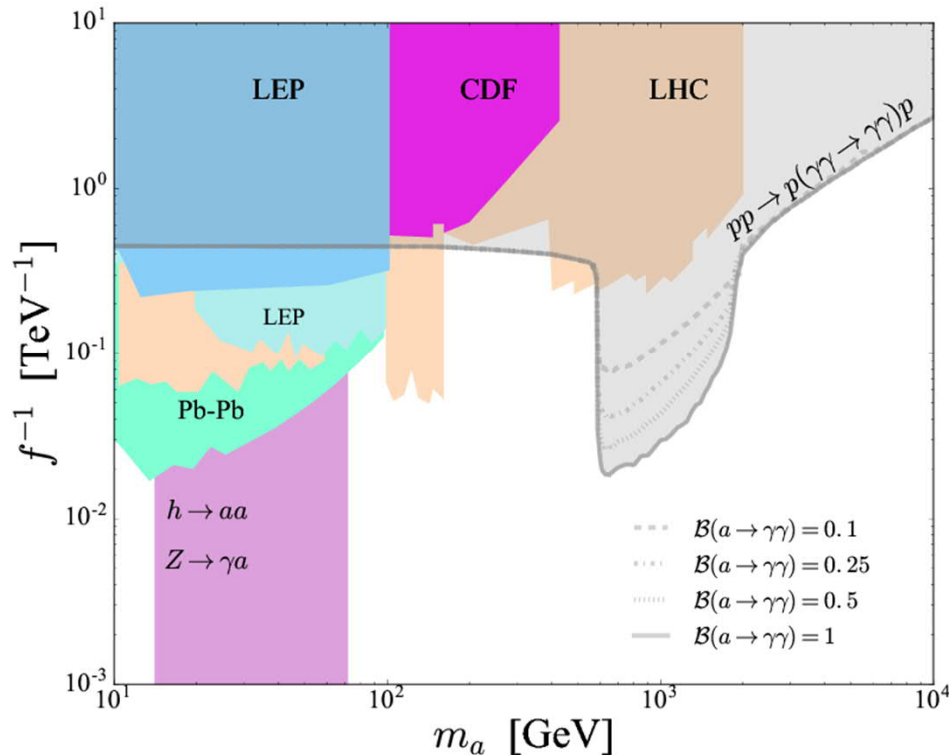




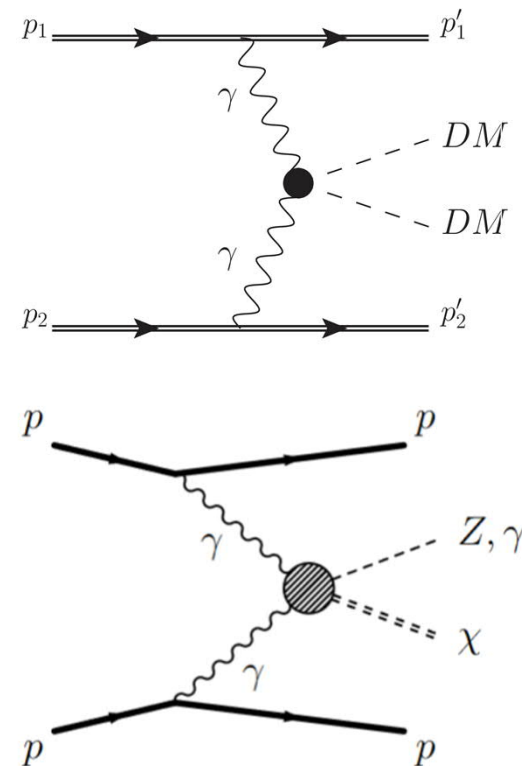
## Search for Axion-Like Particles

via  $\gamma\gamma \rightarrow a \rightarrow \gamma\gamma$

95% CL exclusion regions



## Search for invisible particles (“missing mass”)



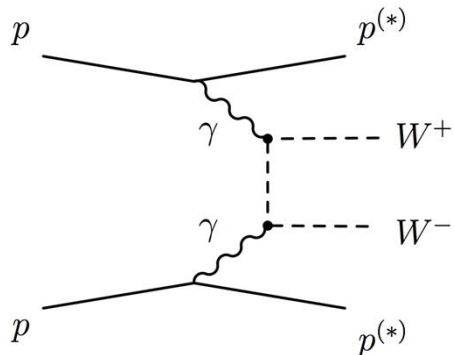
Total central mass measured via protons !

Light grey shaded: PPS @ LHC for 300 fb $^{-1}$

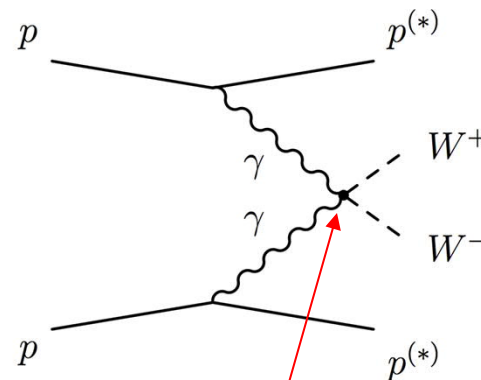
# Physics Examples: Anomalous Gauge Couplings ( $\gamma WW$ )



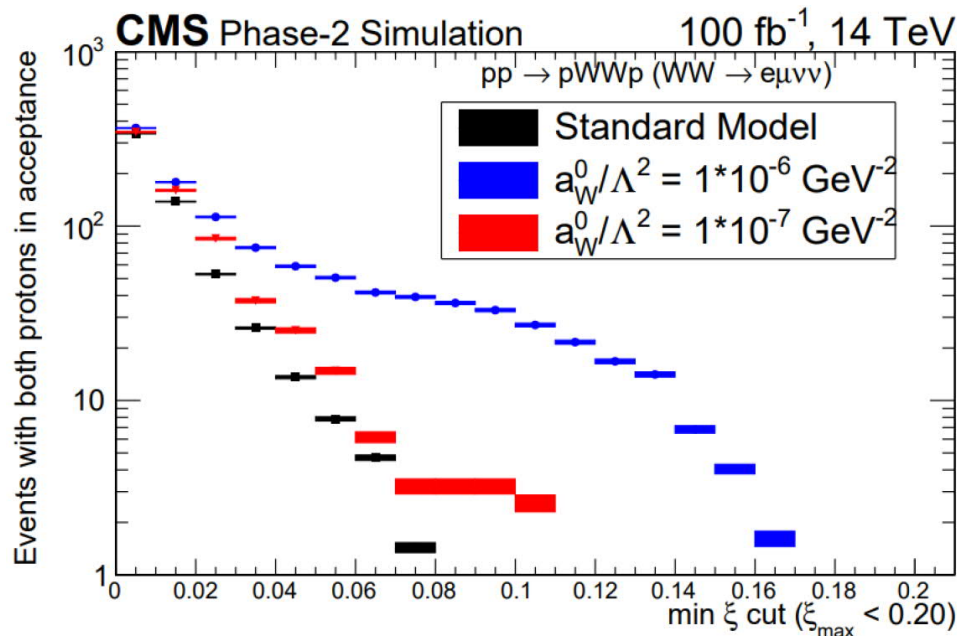
Triple Gauge Couplings



Quartic Gauge Coupling



anomalous contributions (AQGC)?



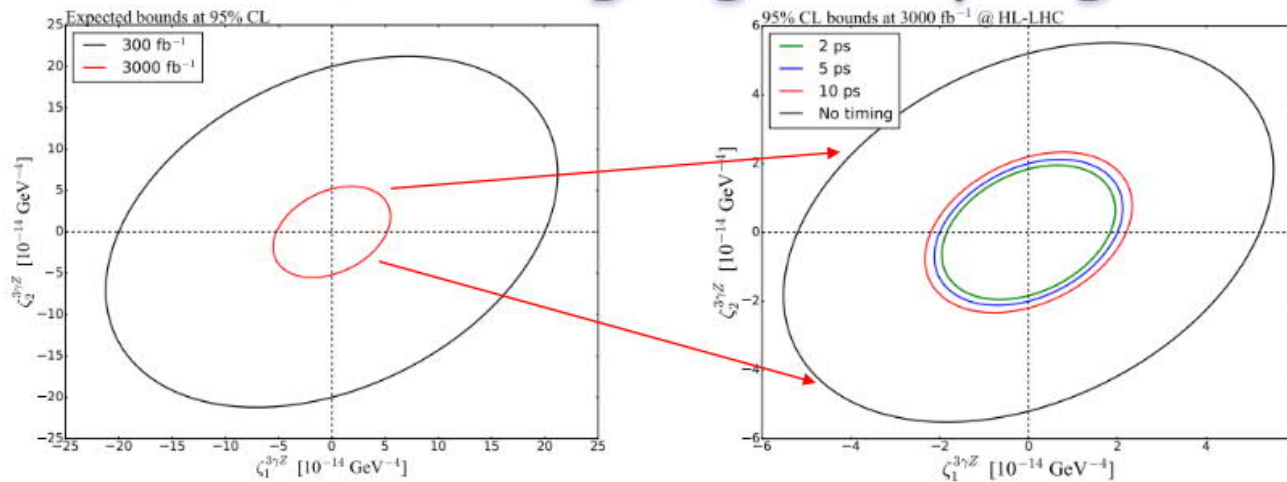
(generator level)

Deviations due to AQGC  
expected to be visible at high masses

# Physics Examples: Anomalous Gauge Couplings ( $\gamma\gamma Z$ )



## Anomalous gauge couplings



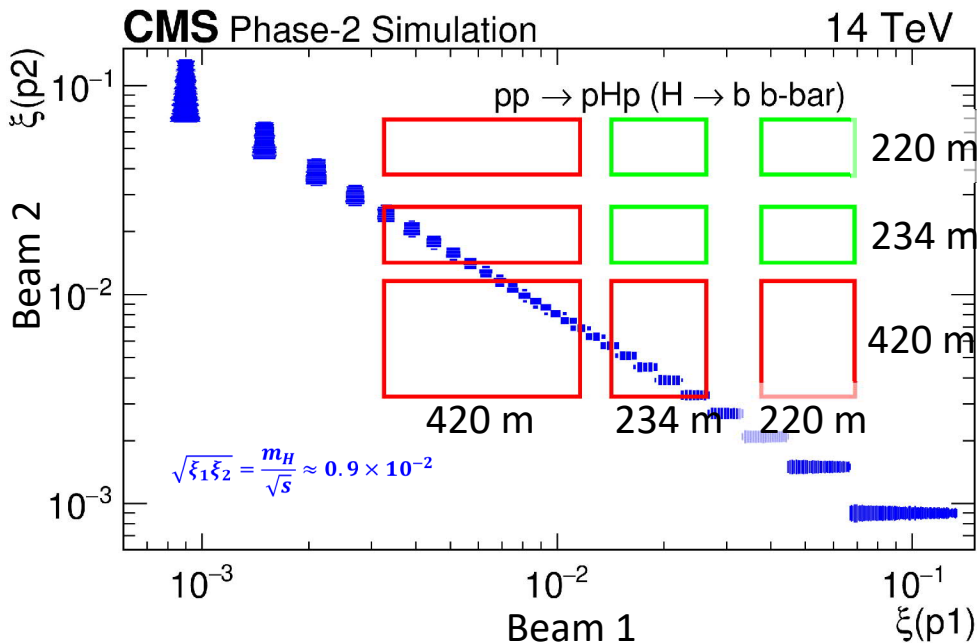
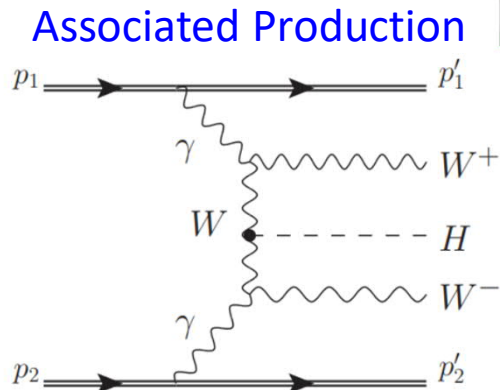
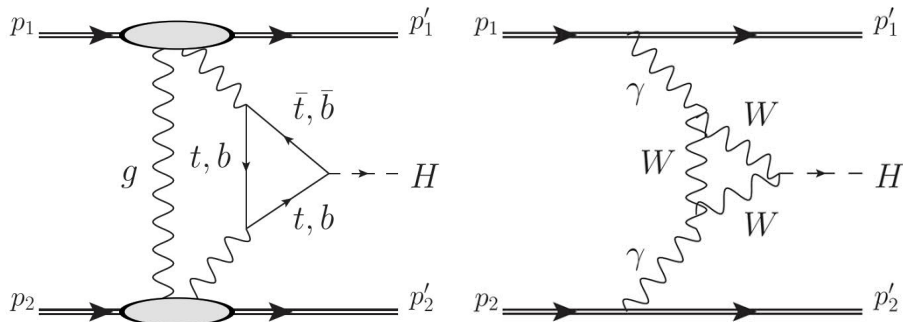
$\langle\mu\rangle=200$ , Lumi=3000fb $^{-1}$

Figure 6: Expected bounds on the anomalous  $\gamma\gamma Z$  couplings at 95% CL with 300 fb $^{-1}$  and 3000 fb $^{-1}$  at the HL-LHC without time-of-flight measurement (left). Expected bounds at 95% CL for timing resolutions of  $\delta t = 2, 5, 10$  ps at the HL-LHC (right). Figure from Ref. [60].

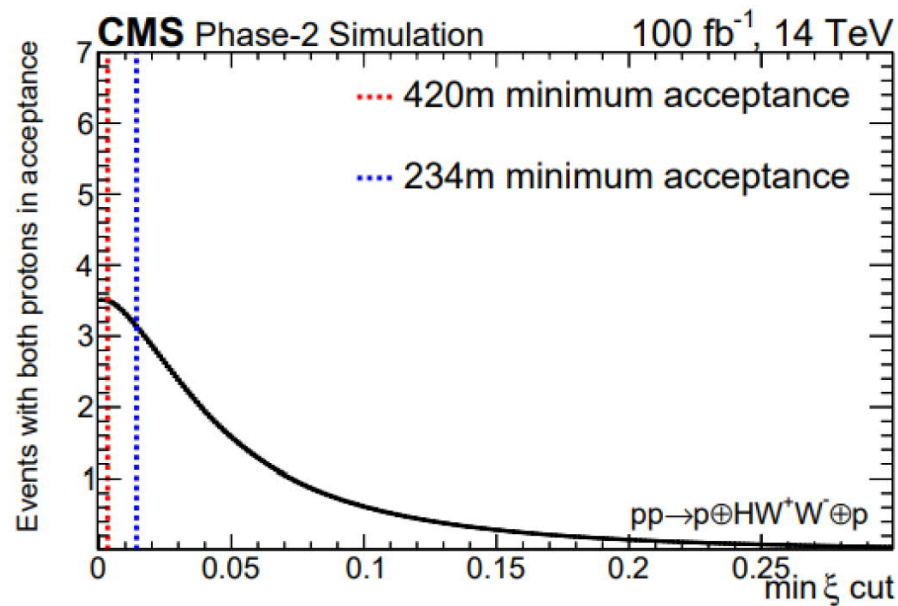
[60] HL-LHC and HE-LHC Working Group Collaboration, “Standard Model Physics at the HL-LHC and HE-LHC”, CERN-LPCC-2018-03, arXiv:1902.04070 .

- $\gamma\gamma Z$  coupling can be probed in  $\gamma\gamma \rightarrow Z\gamma$  channel search.
- Sensitivity is improved with timing detectors

# Physics Examples: Higgs Boson



- needs 420 m stations
- 60 events per  $100 \text{ fb}^{-1}$   
(all decay channels, no experimental cuts)

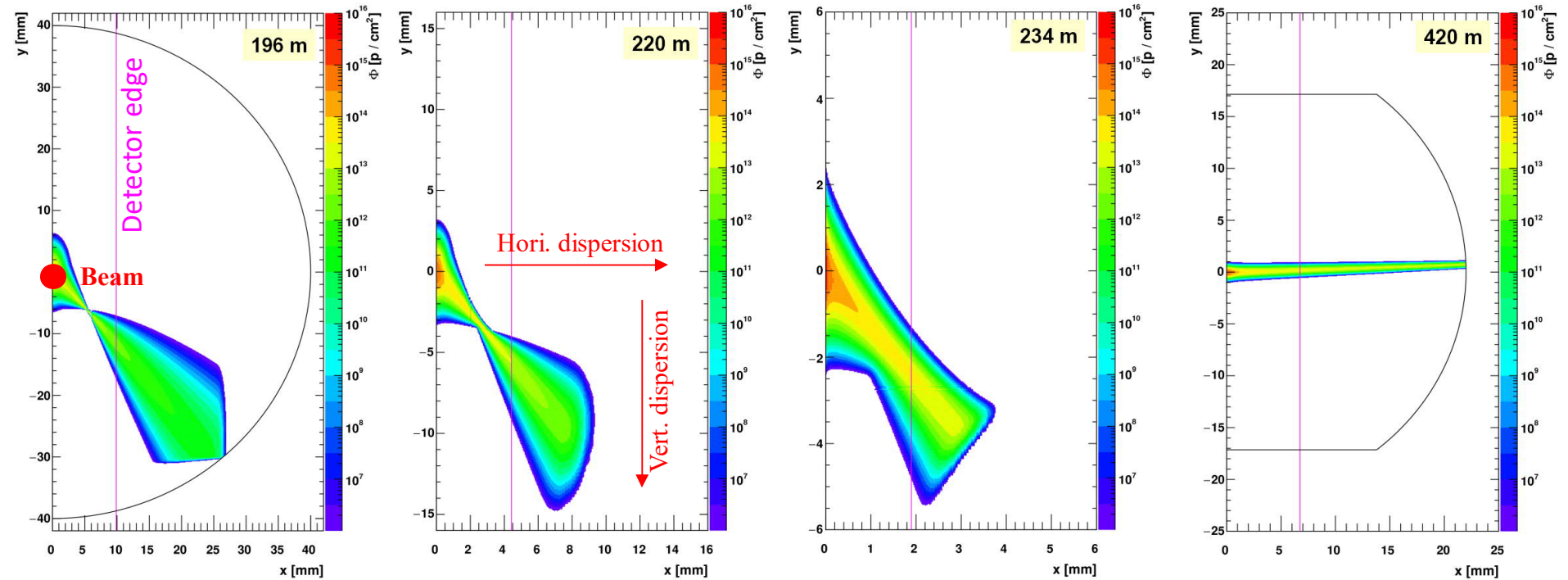


- feasible without 420 m stations  
(large total central mass!)
- 3 events per  $100 \text{ fb}^{-1}$   
(all decay channels, no experimental cuts)

# Detector Requirements: Fluence Maps in Detector Planes



Calculation for Single Diffractive protons (dominant background)



- strongly peaked irradiation
- **Signal distributions similar:**
  - required sensitive areas much larger than at present
  - needs larger thin window (challenge!) but present pot size is appropriate
- occupancy
  - impact on segmentation

Peak fluence after 1 year

Station	$x_{\text{peak}}$ [mm]	$y_{\text{peak}}$ [mm]	$\Phi$ [p/cm <sup>2</sup> ] (300 fb <sup>-1</sup> )
196 m	9.9	-11.6	$5.4 (5.7) \times 10^{14}$
220 m	4.5	-5.7	$2.9 (3.0) \times 10^{15}$
234 m	2.3	-2.7	$1.4 (1.3) \times 10^{16}$
420 m	6.8	0.2	$6.0 (6.0) \times 10^{15}$



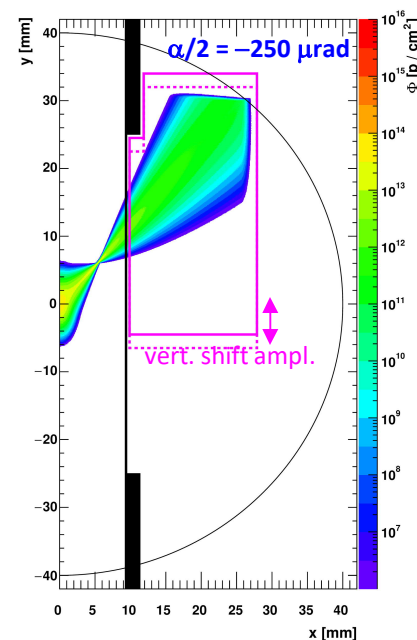
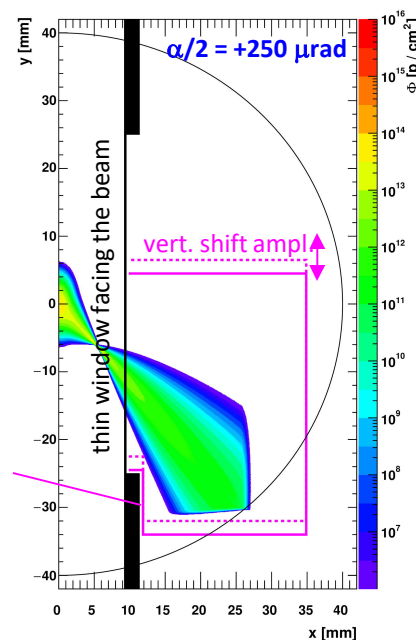
# Design Considerations



- **Detector Area:**

- from hit maps
- extension for vertical shifts to dilute irradiation peak  
(calculated from hit maps and radiation hardness of detectors and electronics)
- detectors covering one polarity of vertical crossing-angle  
but housings for both polarities to allow for annual X-angle flips  
→ replace detectors once a year

196 m station:  
Hitmap, beam pipe,  
thin window and  
detector area:



- **Detector Segmentation:**

- spatial resolution (most crucial for tracking)
  - occupancy (most crucial for timing):
    - \* hit maps: different in the 4 stations
    - \* deadtime
    - \* acceptable level of proton pileup
- different for tracking and timing



Ideas based on present-day PPS experience and other CMS developments.

Presently explored:

## 1. Tracking:

3D silicon pixel detectors:

- used by PPS and CMS tracker in Runs 2 & 3
- improved HL-LHC developments for CMS tracker (sensors and electronics)

## 2. Timing:

– Diamond:

- own developments by TOTEM+PPS, operating in Runs 2 & 3 (very small areas)
- for equipping larger areas: new electronics developments needed

– Ultra-Fast Silicon Detectors (UFSD a.k.a LGAD) from CMS MTD-ETL

(Mip Timing Detector – Endcap Timing Layer)

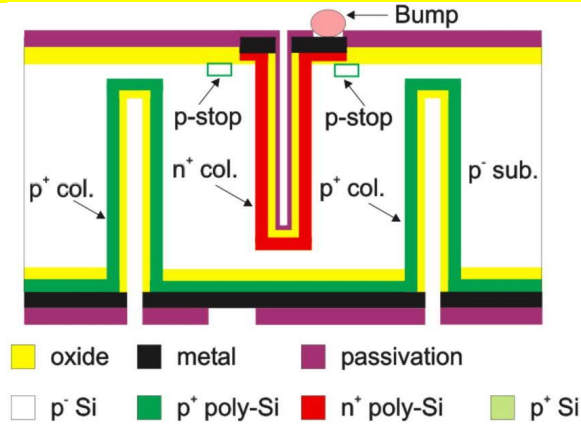
- maximal use of synergy with ETL system (sensors & ETROC electronics)
- some adaptations in segmentation and area needed

No separate pots for tracking and timing

→ combined detector packages (about 10 timing + 6 tracking planes)



# 3D Pixel Silicon Tracking Detectors



- **Spatial resolution:**

Present sensors ( $150 \times 100 \mu\text{m}^2$  pixels, inclination of  $18^\circ$  w.r.t. beam):  $\sigma_x = 25 \mu\text{m}$

- **“Slim Edge” technology:**  $200 \mu\text{m}$  edge  $\rightarrow$   $50 \mu\text{m}$  insensitive margin

- **Radiation hardness:**

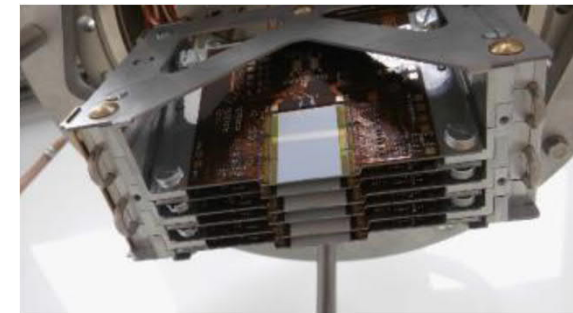
- **sensors:** present generation:  $5 \times 10^{15} \text{ p/cm}^2$   
future (CMS tracker phase 2):  $2 \times 10^{16} \text{ p/cm}^2$   
 $\rightarrow$  sufficient

- **electronics chip bonded to sensor:** degraded in the irradiation peak, fine elsewhere  
 $\rightarrow$  move detector vertically by  $0.5 \text{ mm}$  after  $20 \text{ fb}^{-1}$  to displace the peak

- Run 2: manually in short technical stops

- Run 3 and HL-LHC: remotely controlled piezo-electric motors

At HL-LHC use new pixel readout chips being developed for CMS



- **Possible use for timing:**

Recent tests yield time resolutions of  $20 - 30 \text{ ps}$

# Diamond Timing Detectors



Presently used in PPS: “Double Diamonds” (DD)

(2 layers of scCVD diamond → 1 amplifier)

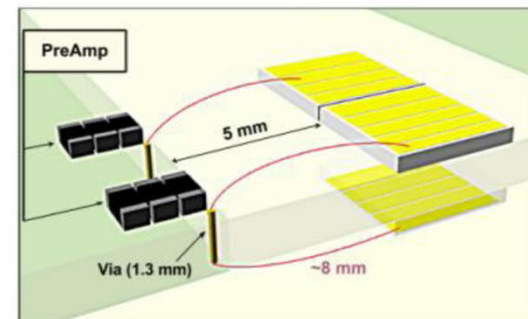
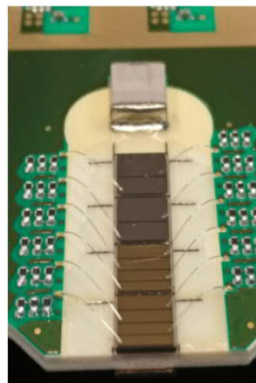
[single crystal chemical vapour deposited]

Time resolution per DD plane:

ideal conditions (testbeam):

**50 ps**

LHC Run 2 (non-perfect readout & biasing): 100 – 150 ps



With present technology: expected achievable time resolution ~ 50 – 60 ps / plane

→ with ~ 10 planes per spectrometer arm: **15 – 20 ps / arm**

Radiation hardness:

after  $100 \text{ fb}^{-1}$  ( $5 \times 10^{15} \text{ p/cm}^2$ ):

- mainly deterioration of electronics,
- only minor deterioration of the **sensors in the tiny irradiation peak near the beam**

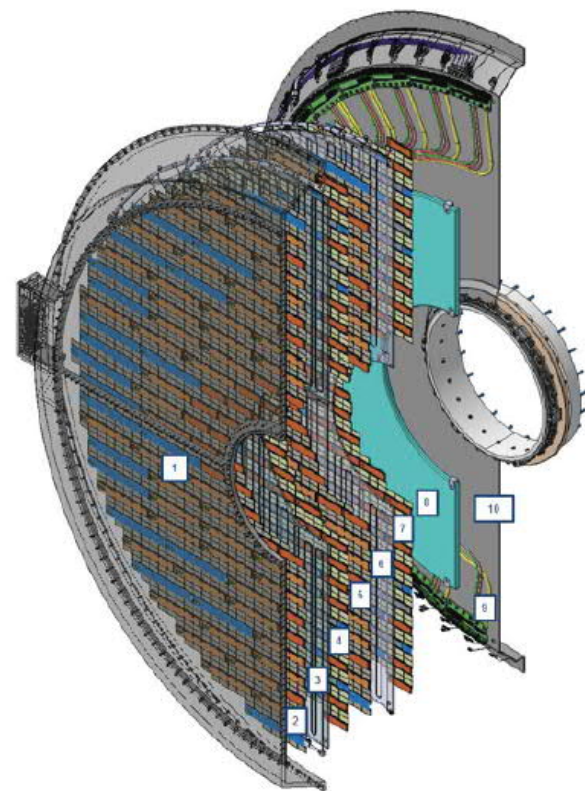


## ETL detector

The ETL detector will be located on the front of the CE covering  $1.6 < |\eta| < 3.0$

Two disks per endcap provide 2 hits per track

- ETL module contains LGAD sensors bump bonded to ETROC ASICs
  - Basic unit: Module
    - $2 \times 4 \text{ cm}^2$  LGAD bump-bonded to 2 ETROC ASICs
    - Mounted on the two sides of cooling plates
- LV, BV, and readout provided by service hybrids
  - Using IpGBT, GBT-SCA and bPOL ASICs
- Each side of each disk is populated with modules and service hybrids
- Time resolution
  - Sensor contribution:  $\sim 30\text{-}40\text{ps}$
  - Electronics contribution  $\sim 30\text{ps}$
  - $50\text{ps}$  per hit  $\rightarrow 35\text{ps}$  per track
- Surface  $\sim 14 \text{ m}^2$ ;  $\sim 8.5 \text{ M}$  channels
- Nominal fluence:  $1.6 \times 10^{15} n_{\text{eq}}/\text{cm}^2$  (@  $3000 \text{ fb}^{-1}$ )



# Conclusions and Outlook



- CMS Expression of Interest for a new PPS spectrometer at HL-LHC:  
extend central production to lower cross-sections and wider mass range

- 4 relevant locations on both sides of IP5 :
  - just before TCL5 (~ 196 m): high masses
  - just before TCL6 (~ 220 m): intermed. masses
  - just after Q6 (~ 234 m): lower masses
  - 420 m: lowest masses

Station	$M_{\min}$ [GeV] @ $y = 0$	$M_{\max}$ [GeV] @ $y = 0$
196 m	1100.87–1197.80	2754.27
220 m	519.89–533.18	962.70
234 m	264.96–132.80	368.11
420 m	43.38–47.04	162.66

- Detector technologies presently studied:
  - Tracking: 3d silicon pixel detectors
  - Time of Flight (to resolve pileup with multiplicity  $\mu \leq 200$ ):
    - Diamond detectors (like present PPS)
    - UFSD (LGAT) from CMS MTD-ETL

- Next step: TDR(s):

First priority: detector vessel (warm stations), machine integration, services

Staged approach: 420 m station in a second step

Physics studies ongoing



The End.



# Appendix



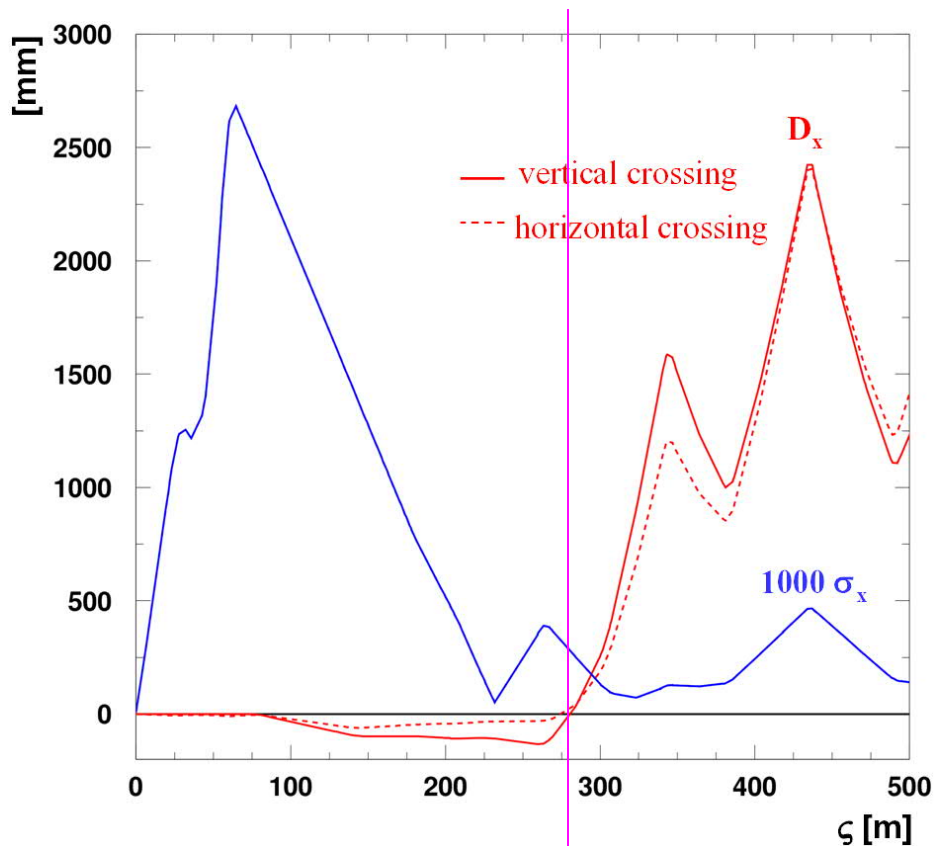


Parameter	Nominal LHC (design report)	HL-LHC (standard)	HL-LHC (BCMS)#
Beam energy in collision [TeV]	7	7	7
Particles per bunch, $N$ [ $10^{11}$ ]	1.15	2.2	2.2
Number of bunches per beam	2808	2748	2604
Number of collisions in IP1 and IP5*	2808	2736	2592
$N_{\text{tot}}$ [ $10^{14}$ ]	3.2	6.0	5.7
Beam current [A]	0.58	1.09	1.03
Crossing angle in IP1 and IP5 [ $\mu\text{rad}$ ]	285	510	510
Minimum normalized long-range beam-beam separation [ $\sigma$ ]	9.4	12.5	12.5
Minimum $\beta^*$ [m]	0.55	0.2	0.2
$\epsilon_n$ [ $\mu\text{m}$ ]	3.75	2.50	2.50
$\epsilon_L$ [eVs]	2.50	2.50	2.50
R.M.S. energy spread [0.0001]	1.13	1.08	1.08
R.M.S. bunch length [cm]	7.55	8.1	8.1
IBS horizontal in collision [h]	80→106	18.8	18.8
IBS longitudinal in collision [h]	61→60	20.6	20.6
Piwinski parameter	0.65	2.5	2.5
Total reduction factor $R_0$ without crab cavities at min. $\beta^*$	0.836	0.369	0.369
Total reduction factor $R_1$ with crab cavities at min. $\beta^*$	(0.981)	0.715	0.715
Beam-beam tune shift/IP	0.0031	0.01	0.01
Peak luminosity without crab cavities $L_{\text{peak}}$ [ $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ]	1.00	6.52	6.18
Peak luminosity with crab cavities $L_{\text{peak}} \times R_1/R_0$ [ $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ]	(1.18)	12.6	11.9
Events/crossing without levelling and without crab cavities	27	172	172
Levelled luminosity for $\mu = 140$ [ $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ]	-	5.32 <sup>†</sup>	5.02 <sup>†</sup>
Events/crossing $\mu$ (with levelling and crab cavities) <sup>‡</sup>	27	140	140
Maximum line density of pile-up events during fill [events/mm]	0.21	1.3	1.3
Levelling time [h] (assuming no emittance growth) <sup>‡</sup>	-	5.23	5.23
Number of collisions in IP2/IP8	2808	2452/2524**	2288/2396**
$N$ at injection [ $10^{11}$ ] <sup>††</sup>	1.20	2.30	2.30
Maximum number of bunches per injection	288	288	288
$N_{\text{tot}}$ /injection [ $10^{13}$ ]	3.46	6.62	6.62
$\epsilon_n$ at SPS extraction [ $\mu\text{m}$ ] <sup>‡‡</sup>	3.50	2.00	<2.00***

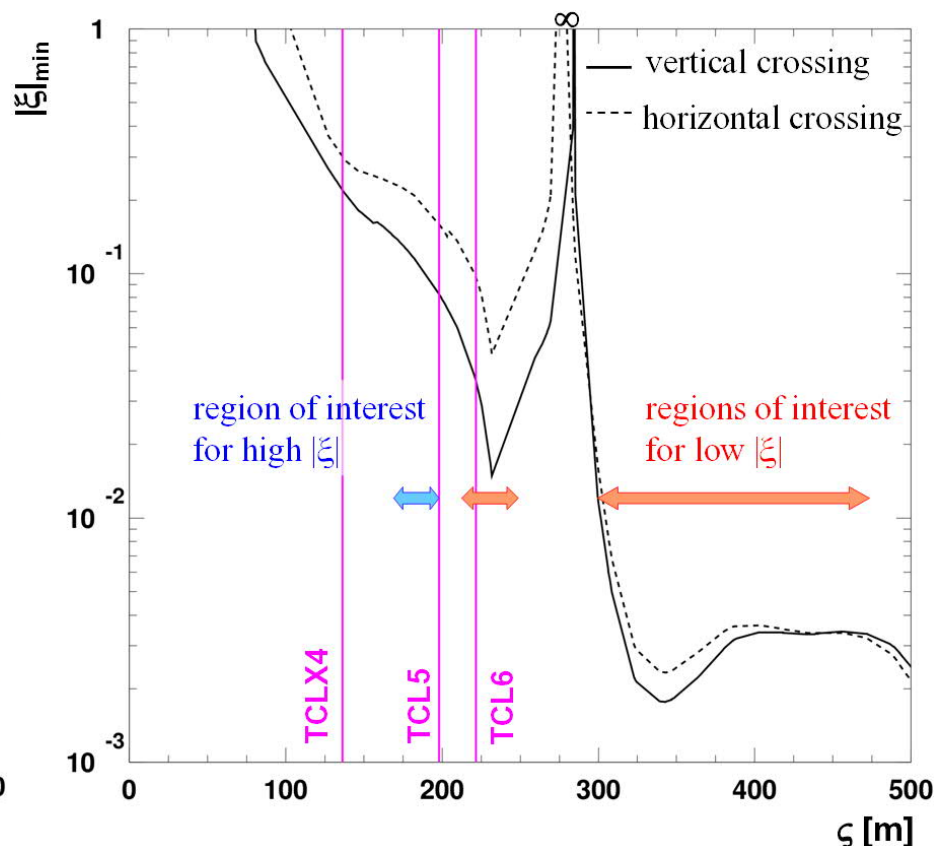
# Search for Detector Locations (2)



- HL-LHC optics version 1.3
- for (crossing-angle  $\alpha/2$ ,  $\beta^*$ ) = (250  $\mu$ rad, 15 cm)
- Roman Pots @ (12.9 + 3)  $\sigma$  + 0.3 mm



$$|\xi|_{\min} = (15.9 \sigma_x + 0.3 \text{ mm}) / D_x$$

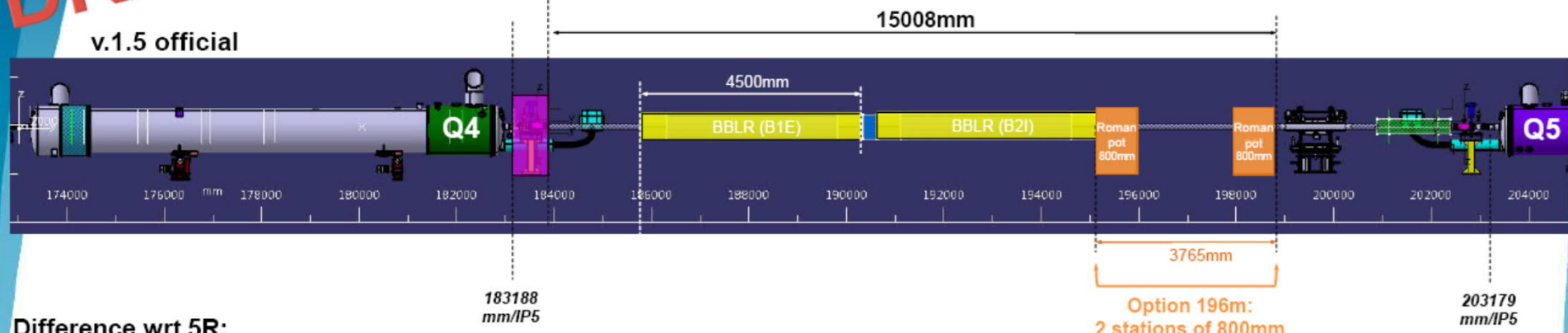


for  $s > \sim 270$  m :  $D_x > 0$   
 → protons with momentum loss are between the beam pipes  
 → no standard Roman Pot possible → needs new technology  
**Free only around 420 m.**

# DRAFT

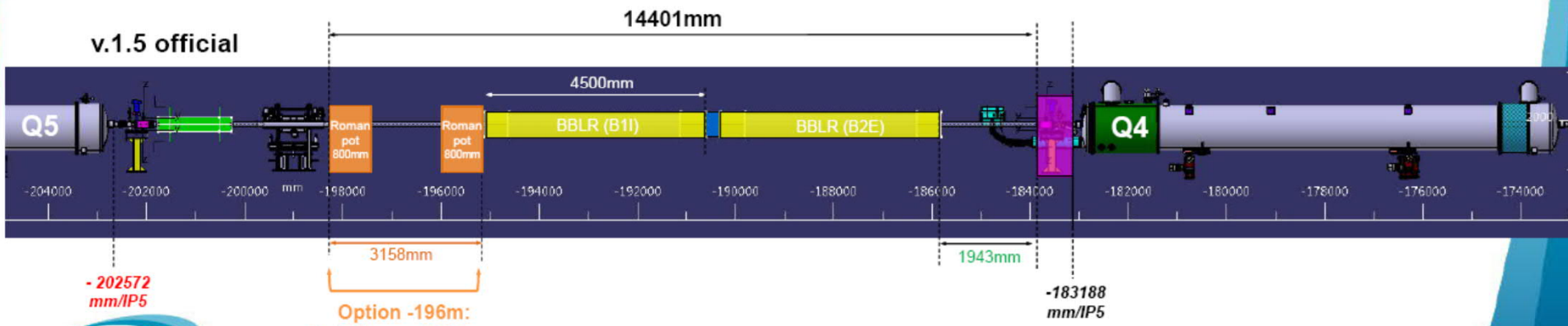
## Region of interest: 180-200m/IP5 → 196m

### LSS5R - MACHINE LAYOUT v.1.5



Difference wrt 5R:  
 $203179 - 202572 = 607\text{mm}$   
There is 607mm less space  
in 5L = 1R !!

### LSS5L - MACHINE LAYOUT v.1.5



18/11/2020

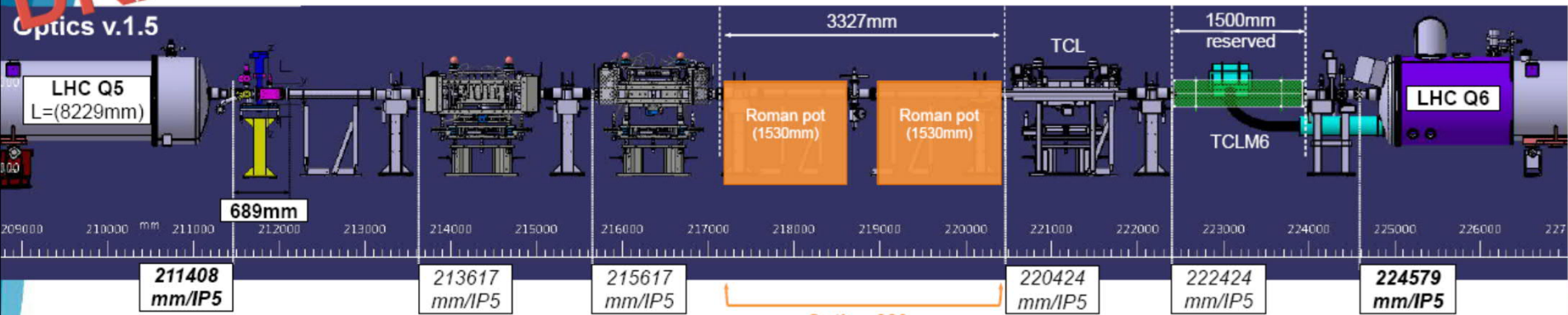
M. Gonzalez, P. Fessia



**DRAFT**

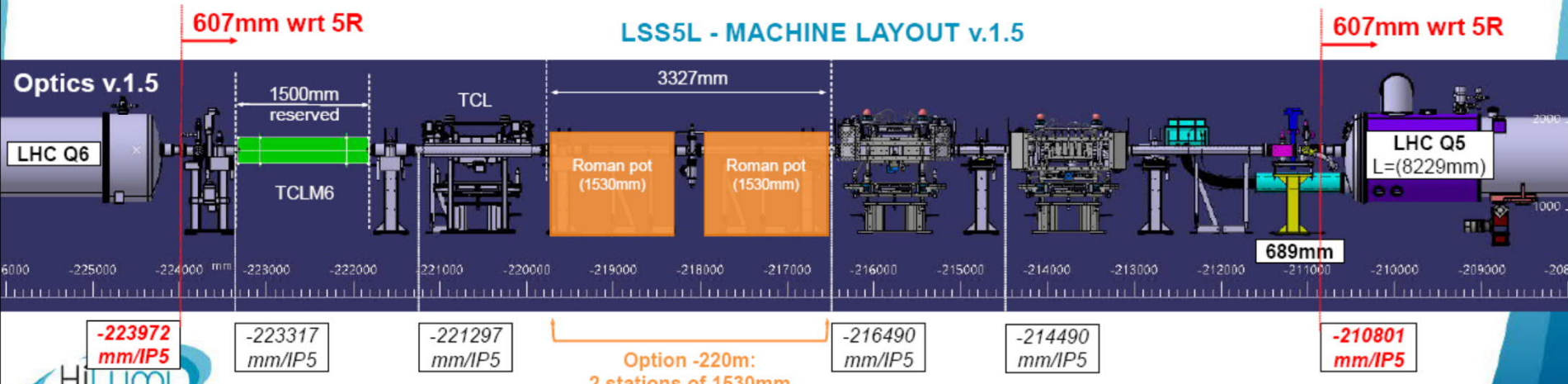
Region of interest: 210-250m/IP5 → 220m

LSS5R - MACHINE LAYOUT v.1.5



All components are shifted by 607mm towards the IP in 5L = 1R !! Wrt 5R

LSS5L - MACHINE LAYOUT v.1.5



18/11/2020

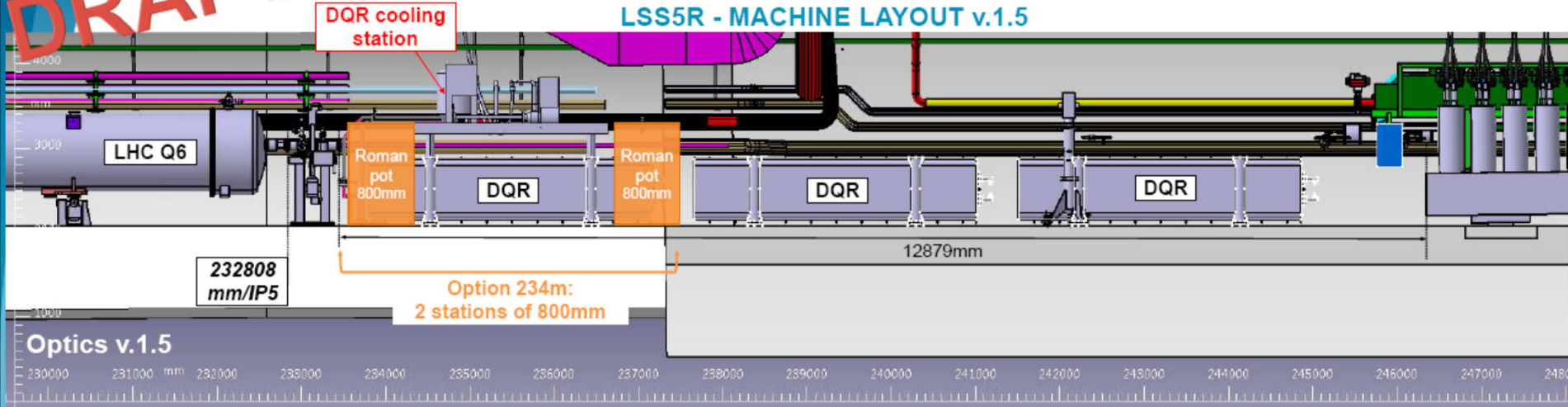
M. Gonzalez, P. Fessia

Only in this station: vertical units for alignment and optics calibration

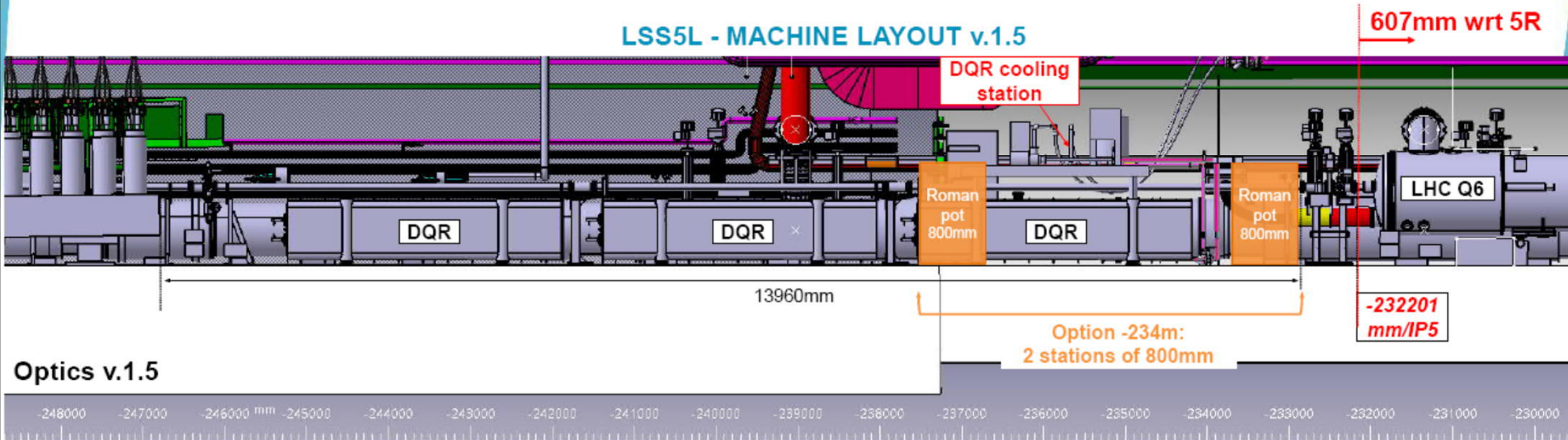
**DRAFT**

Region of interest: 210-250m/IP5 → 234m

LSS5R - MACHINE LAYOUT v.1.5



LSS5L - MACHINE LAYOUT v.1.5



# XRP Insertion Distance vs. $\beta^*$



Assume insertion rule:  $d_{\text{XRP}} = (n_{\text{TCT}} + 3)\sigma_{\text{XRP}} + 0.3 \text{ mm}$

Collimation scheme presently foreseen:

$$d_{\text{TCT}} = \text{const.} \rightarrow n_{\text{TCT}}(\beta^*) = n_{\text{TCT}}(\beta_0^*) \sqrt{\frac{\beta^*}{\beta_0^*}}$$

$$\sigma_{\text{XRP}} = \sqrt{\frac{\varepsilon_n \beta_{\text{XRP}}}{\gamma}}$$

We need  $\beta_{\text{XRP}}(\beta^*)$  !

ATS invariance of optical functions:  $v_{\text{XRP}} = \sqrt{\frac{\beta_{\text{XRP}}(\beta^*)}{\beta^*}} \cos \mu_{\text{XRP}}(\beta^*)$  : magnification independent of  $\beta^*$

$L_{\text{XRP}} = \sqrt{\beta_{\text{XRP}}(\beta^*) \beta^*} \sin \mu_{\text{XRP}}(\beta^*)$  : eff. length independent of  $\beta^*$

$$\Rightarrow \left\{ \begin{array}{l} \tan \mu_{\text{XRP}}(\beta^*) = \frac{L_{\text{XRP}}}{v_{\text{XRP}}} \frac{1}{\beta^*} \\ \beta_{\text{XRP}}(\beta^*) = \frac{L_{\text{XRP}} v_{\text{XRP}}}{\sin \mu_{\text{XRP}}(\beta^*) \cos \mu_{\text{XRP}}(\beta^*)} \end{array} \right\} \Rightarrow \beta_{\text{XRP}}(\beta^*) = v_{\text{XRP}}^2 \beta^* + \frac{L_{\text{XRP}}^2}{\beta^*}$$

$$\sigma_{\text{XRP}} = \sqrt{\frac{\varepsilon_n}{\gamma} \left( v_{\text{XRP}}^2 \beta^* + \frac{L_{\text{XRP}}^2}{\beta^*} \right)}$$

$$d_{\text{XRP}} = \left( n_{\text{TCT}}(\beta_0^*) \sqrt{\frac{\beta^*}{\beta_0^*}} + 3 \right) \sqrt{\frac{\varepsilon_n}{\gamma} \left( v_{\text{XRP}}^2 \beta^* + \frac{L_{\text{XRP}}^2}{\beta^*} \right)} + 0.3 \text{ mm}$$

# Mass Acceptance Calculation



Calculate mass limits:  $M_{\min/\max} = \xi_{\min/\max} \sqrt{s}$  in  $(\alpha/2, \beta^*)$  plane  
( for symmetric optics in Beam 1 / Beam 2 with  $\xi_{1 \min/\max} = \xi_{2 \min/\max}$  )

Cannot simulate every  $(\alpha/2, \beta^*)$  point  $\rightarrow$  analytical approach:

gap + insensitive XRP detector margin

$$M_{\min} = \xi_{\min} \sqrt{s} \text{ with } \xi_{\min} = \frac{d_{\text{XRP}}(\beta^*) + \delta}{D_{x,\text{XRP}}(\frac{\alpha_x}{2}, \xi_{\min})} \text{ resolved for } \xi_{\min}$$

$$M_{\max} = \xi_{\max} \sqrt{s} = \frac{d_A}{D_A(\frac{\alpha}{2}, \xi_{\max})} \sqrt{s}$$

$d_{\text{XRP}}$ : detector distance from beam centre:  
analytical expression depending on  
TCT collimator settings  
and optics properties

$D_{\text{XRP}}$ : hori. dispersion @ detector location,  
parametrisation in  $(\alpha/2, \xi)$  from MAD-X

Based on full aperture study

$d_A$ : aperture limitation (hori. or vert.) upstream,  
in most cases: TCLs

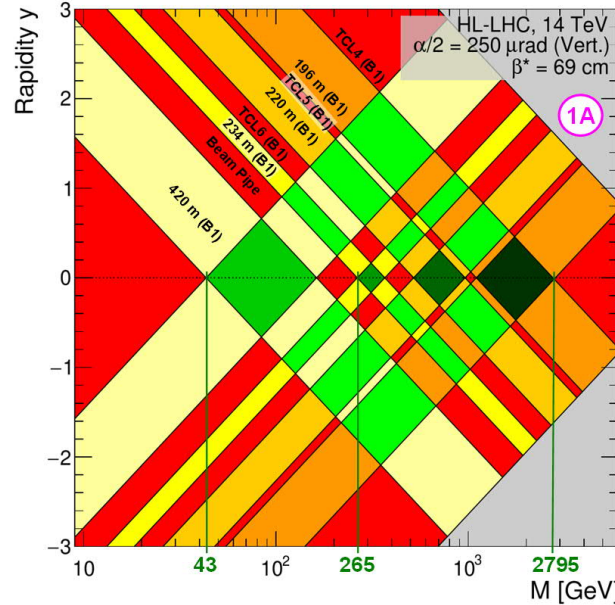
$D_A$ : dispersion (hori. or vert.) @ aperture limit.,  
parametrisation in  $(\alpha/2, \xi)$  from MAD-X



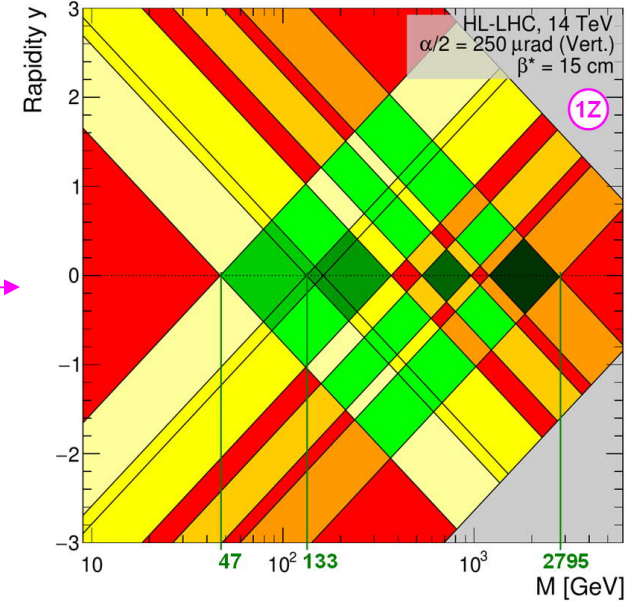
# Acceptance in Mass – Rapidity Plane



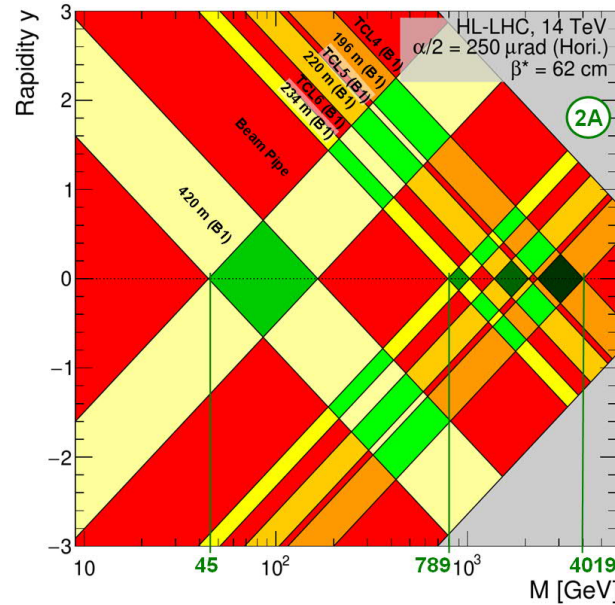
Vertical crossing in IP5



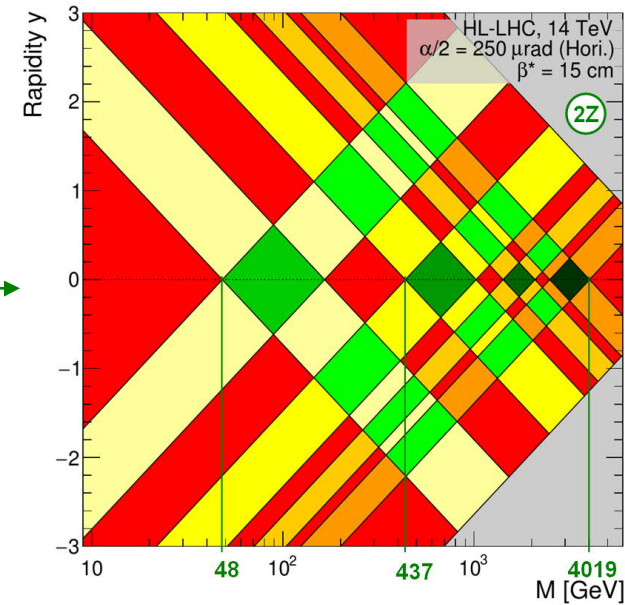
Fill →



Horizontal crossing in IP5



Fill →



Large gaps !

$$\ln \frac{M}{\sqrt{s}} = \frac{1}{2} (\ln \xi_1 + \ln \xi_2)$$

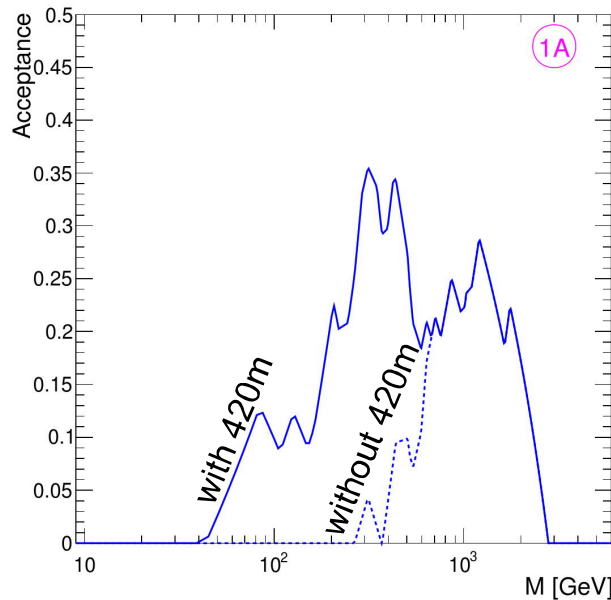
$$y = \frac{1}{2} (\ln \xi_1 - \ln \xi_2)$$

# Mass Acceptance

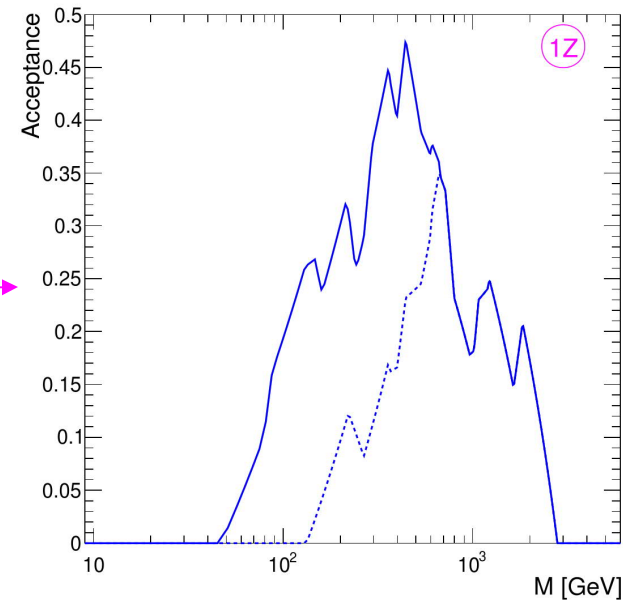
projections on M axis assuming flat rapidity distribution



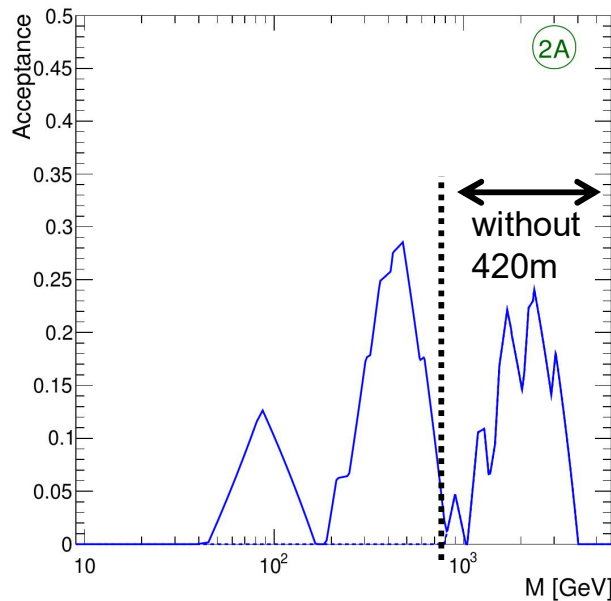
Vertical crossing in IP5



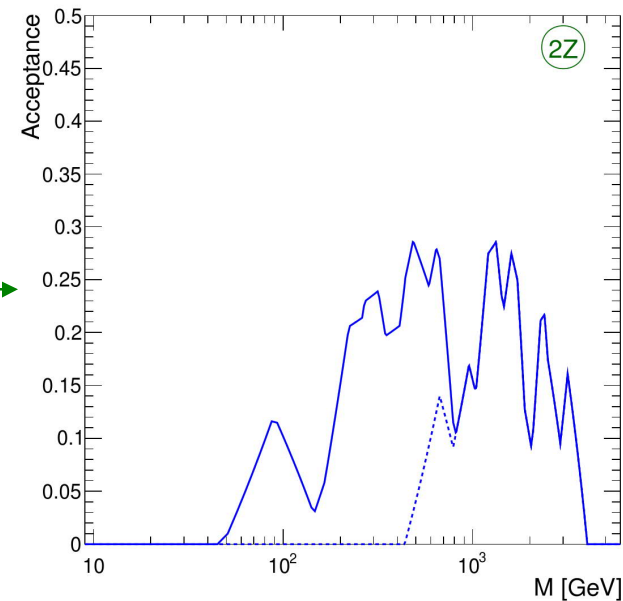
Fill



Horizontal crossing in IP5



Fill

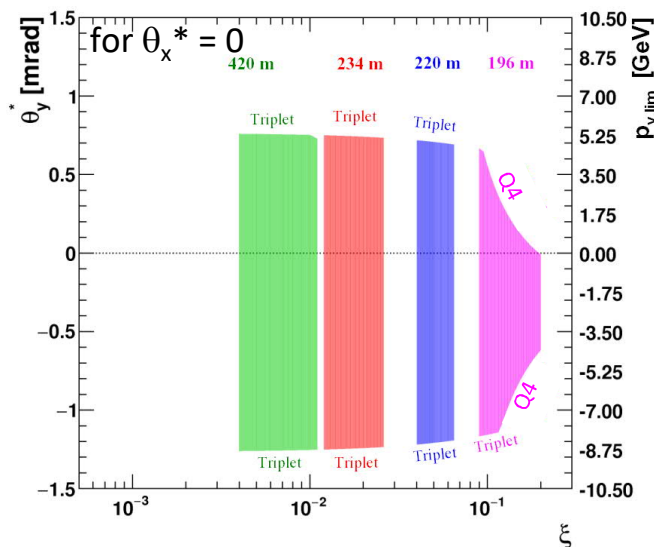
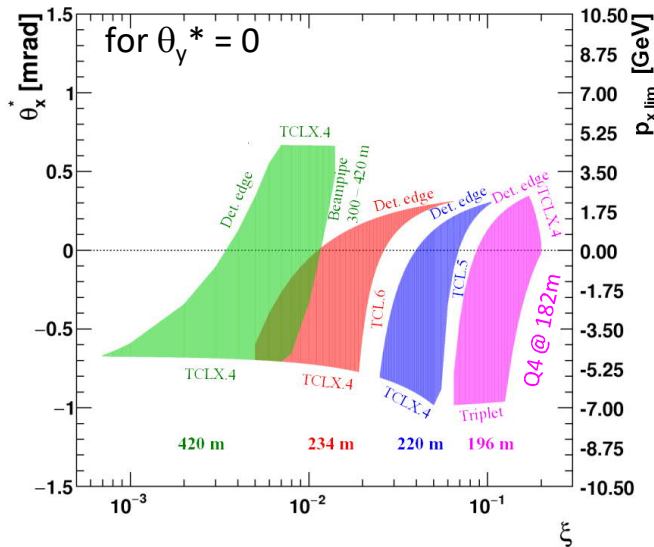


# Single-Arm Proton Acceptance (Non-Zero $p_T$ )



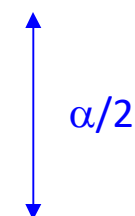
Acceptance numbers in EoI are taking into account all aperture limitations upstream of XRPs but assuming full instrumentation of the scoring planes.

## Physics Runs

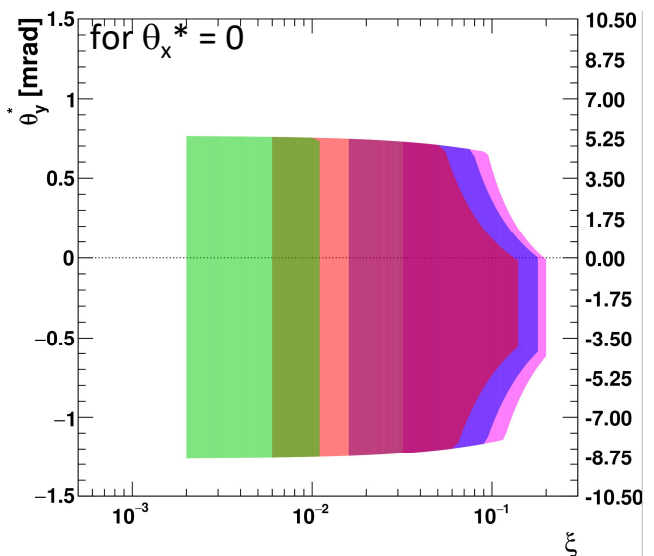
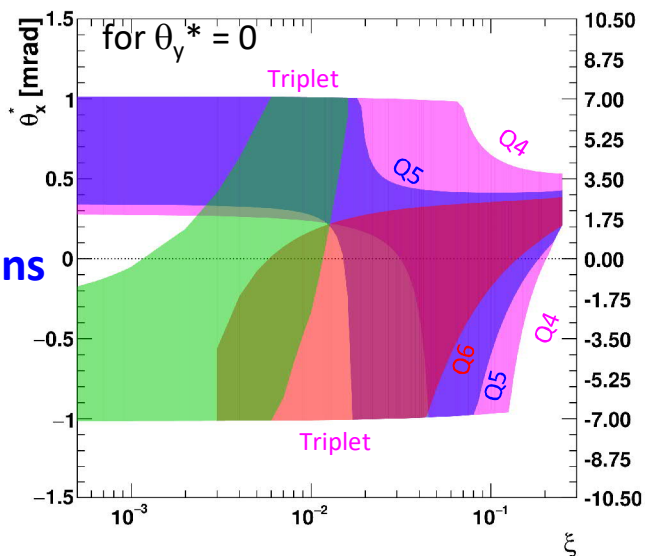


These plots are for  $\alpha/2 = +250 \mu\text{rad}$

Continuous variation of  $\alpha/2$  during the fill shifts the  $\theta_{y^*}$  acceptance blocks (approx.) up and down



## Calibration Runs



- RPs at  $5 \sigma$  from beam
- no TCLs





# ToF and Vertex Resolution

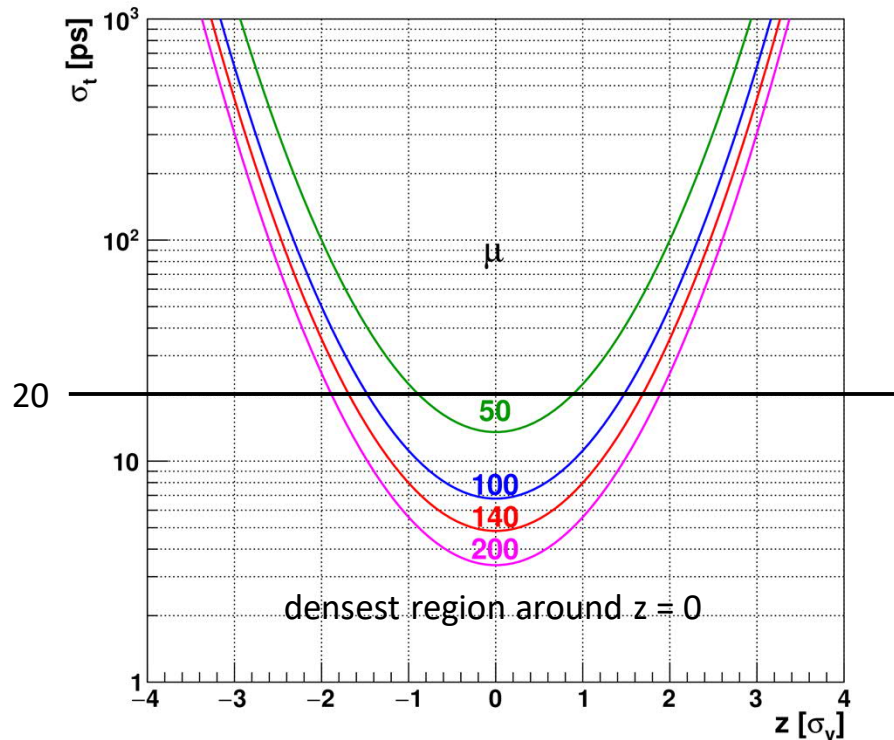


With present technology: expected time resolution  $\sim 50 - 60$  ps / plane

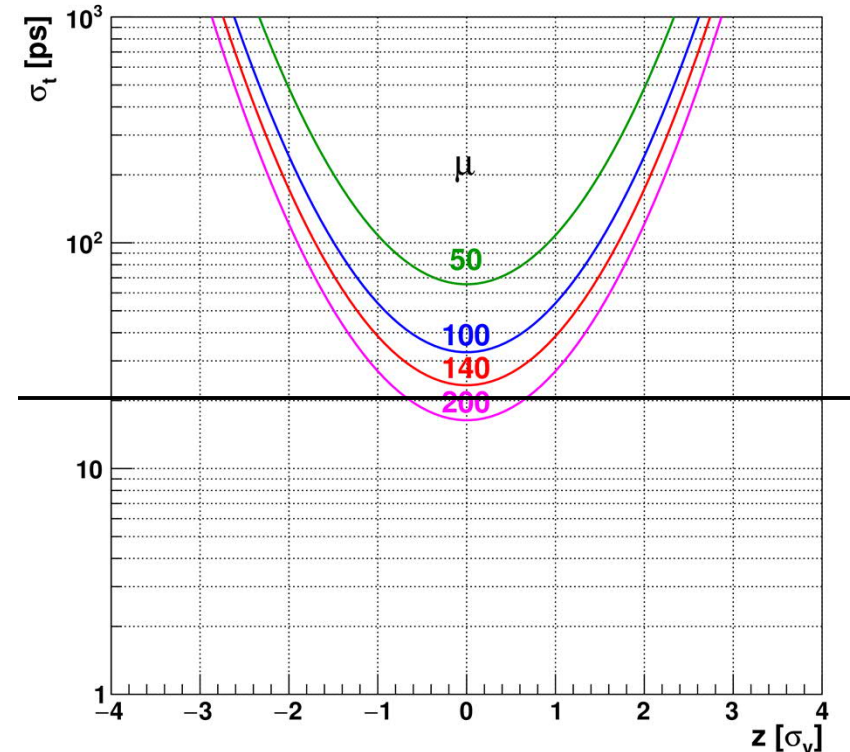
→ with  $\sim 10$  planes per spectrometer arm: 15 – 20 ps / arm

Required time resolution per arm to resolve mean vertex distance:

## PPS alone



## Combined with MTD timing



For  $\mu = 140$  (200):

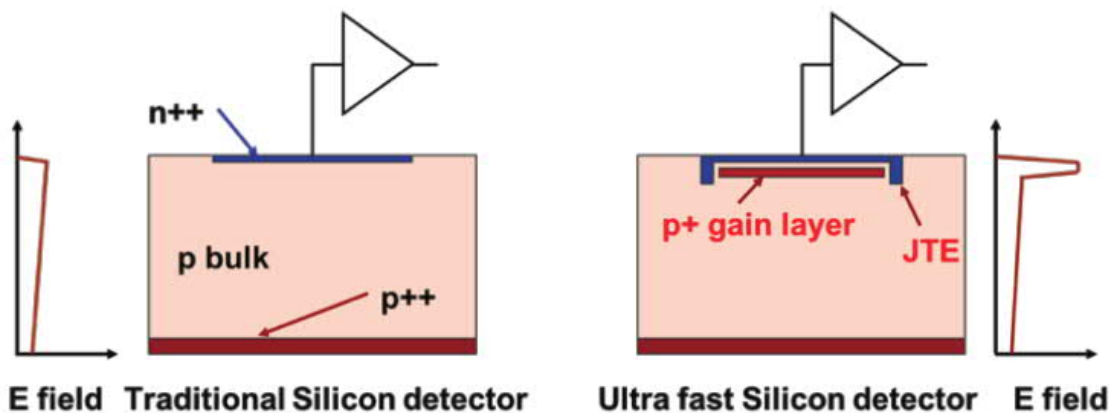
resolves vertices only outside  $1.7$  ( $1.9$ )  $\sigma$

But: event topology selections reduce eligible vertices!

to be studied in detail

# Low gain avalanche diodes

- The LGAD sensors, as proposed and first manufactured by CNM
- High field obtained by adding an extra doping layer
- $E \sim 300$  kV/cm, close to breakdown voltage
- Gain is the key ingredient to good time resolution

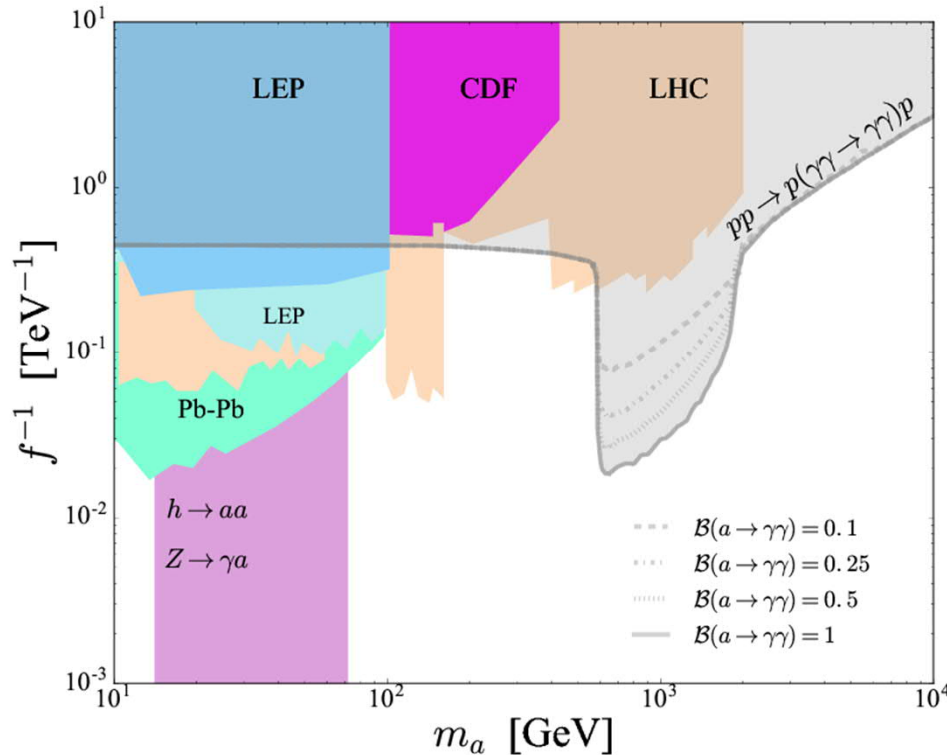




## Search for Axion-Like Particles

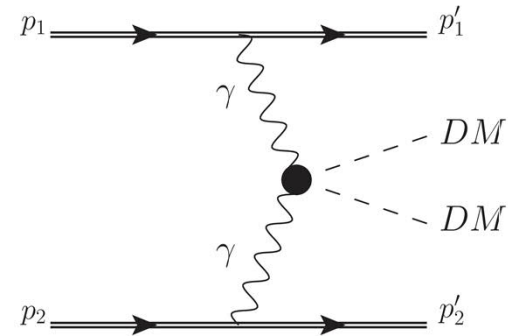
via  $\gamma\gamma \rightarrow a \rightarrow \gamma\gamma$

95% CL exclusion regions



Light grey shaded: PPS @ LHC for 300 fb<sup>-1</sup>

## Search for invisible particles (“missing mass”)

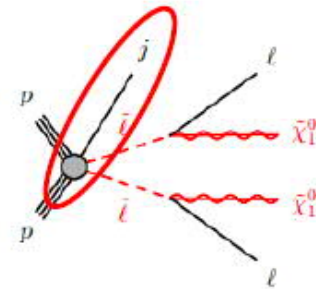


Example: SUSY searches in compressed mass scenario

$$pp \rightarrow \tilde{\ell}\tilde{\ell} \rightarrow \ell\ell\tilde{\chi}_1^0\tilde{\chi}_1^0$$

**Conventional search:** need ISR jets to boost neutralinos  $\rightarrow$  high missing  $E_T$

[PRD 101 (2020), 052005]



**Central production:**

measure  $m_{\tilde{\ell}\tilde{\ell}}$  via protons!

[JHEP 1904, 010 (2019); PRL 123 (2019) 141801]