

# Higgs pair production with EFT Modeling at NLO

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On behalf of

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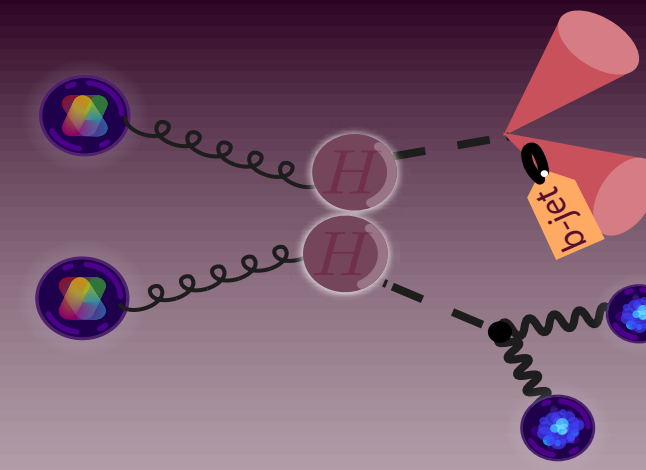
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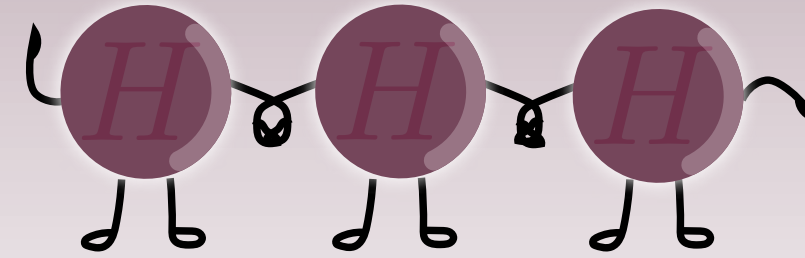
# Why looking for Higgs pairs ?



- The natural “ next-step” after finding the Higgs.

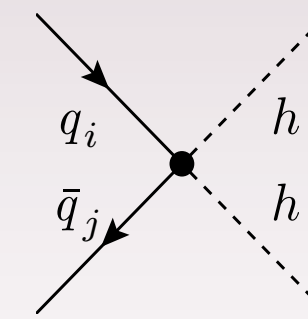
- Essential for probing the trilinear coupling.

S. Di Vita, et al (2017)



- Provides a direct measurement of Higgs coupling to light quarks

[L.A, R. Corral Lopez and R Gröber '19]



- Sensitive to non-linearity in couplings with the Higgs

- Probe some of the Higgs EFT operators



Could we save computational power by reweighting SM MC samples to become EFT ones ?

Both ATLAS and CMS are now trying to implement the reweighting procedure

# HH production in Effective Field Theories

There are 2- famous- EFT's for the Higgs pair process

- Higgs Effective Field Theory - non-linear- a.k.a. Chiral Lagrangian [lonso et al '12 and Buchalla et al '13]

$$\mathcal{L}_{\text{HEFT}} = -m_t \left( c_t \frac{h}{v} + c_{tt} \frac{h^2}{v^2} \right) (\bar{t}_R t_L + h.c.) - c_{hhh} \frac{m_h^2}{2v^2} h^3 + \frac{\alpha_s}{8\pi} \left( c_{ggh} \frac{h}{v} + c_{gghh} \frac{h^2}{v^2} \right) \text{tr}(G_{\mu\nu} G^{\mu\nu})$$

- SM Effective Field Theory, strongly interacting light Higgs (SILH) basis [R. Contino, et al. '13]

$$\mathcal{L}_{\text{SILH}} = \frac{\bar{c}_H}{2v^2} \partial_\mu (H^\dagger H) \partial^\mu (H^\dagger H) - \frac{\bar{c}_6 \lambda}{v^2} (H^\dagger H)^3 + \left( \frac{\bar{c}_u}{v^2} y_u H^\dagger H \bar{q}_L H^c t_R + h.c. \right) + \frac{\bar{c}_g g_s^2}{m_W^2} H^\dagger H \text{tr}(G_{\mu\nu} G^{\mu\nu})$$

- SM Effective Field Theory, Warsaw basis [B. Grzadkowski, et al. '10]

$$\mathcal{L}_{\text{SMEFT}} = C_{H,\square} (H^\dagger H) \square (H^\dagger H) + C_{HD} |(H^\dagger D_\mu H)|^2 + C_H (H^\dagger H)^3 + C_{uH} (H^\dagger H \bar{q}_L H^c t_R + h.c.) \\ + C_{HG} H^\dagger H \text{tr}(G_{\mu\nu} G^{\mu\nu})$$

SILH and Warsaw are the same except for the choice of which operators are removed by E.o.M !

# ATLAS

by

C. Dimitriadi, L. Pereira Sanchez and A. Ferrari

# Reweighting in HEFT

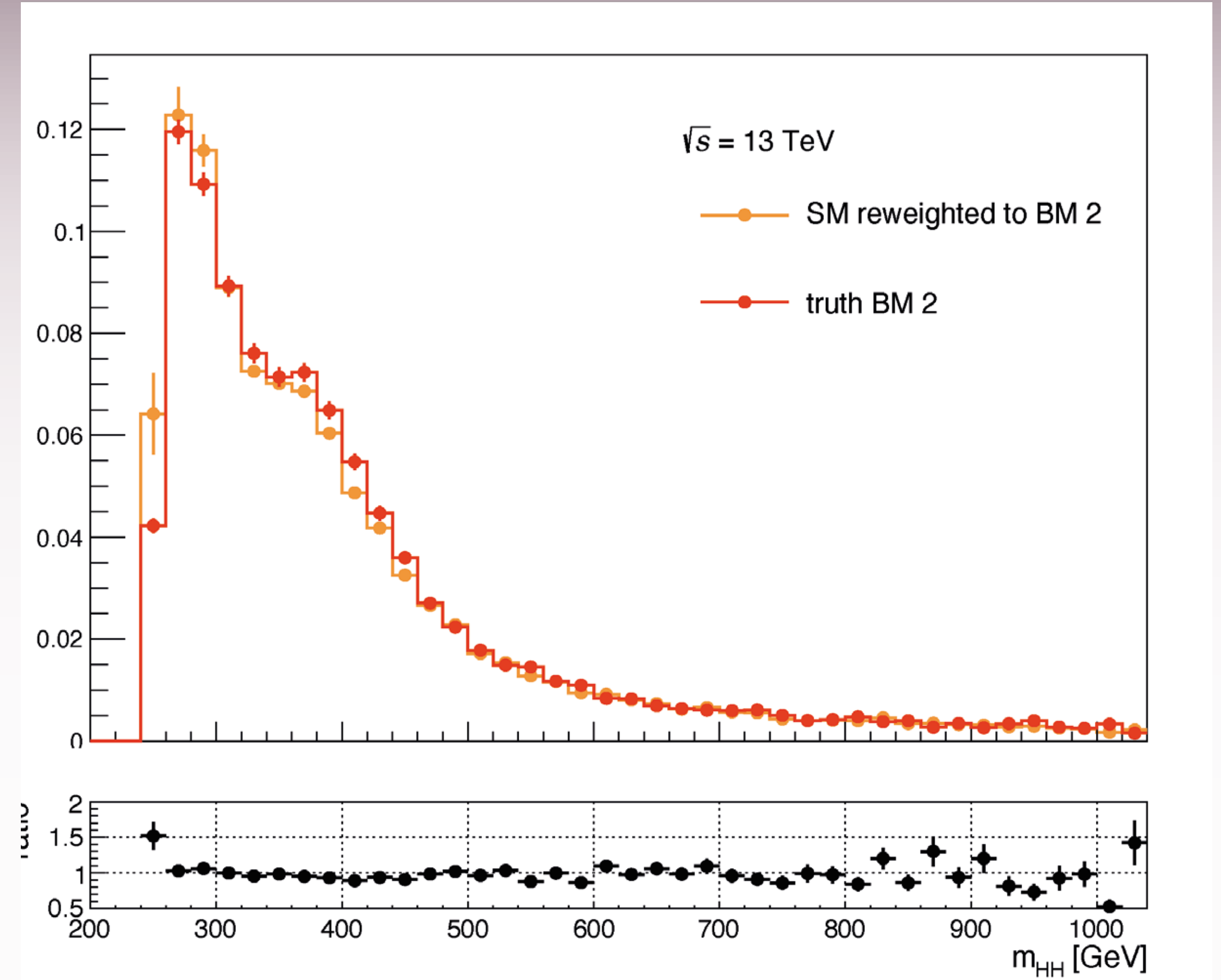
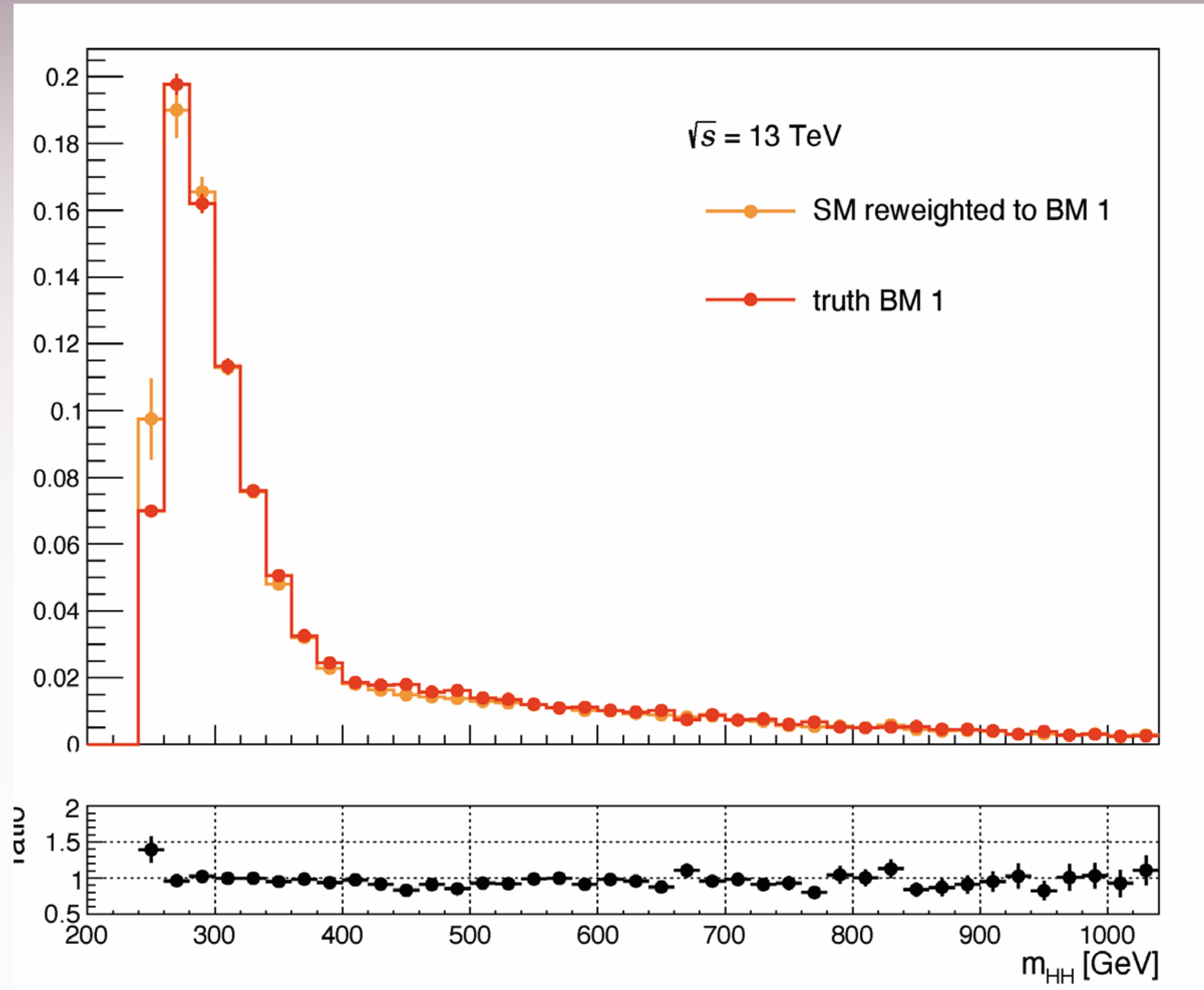
- The weights for the total, and differential cross sections are available for LO and NLO [G. Heinrich, et al. '17]  
[G. Heinrich, et al. '20]
- Reweighting of  $M_{hh}$  distributions were done for 7 benchmarks as a cross-check.

benchmark	ct	chhh	ctt	cggh	cgghh
SM	1	1	0	0	0
1	0.94	3.94	-0.333	0.5	0.333
2	0.61	6.84	0.333	0	-0.333
3	1.05	2.21	-0.333	0.5	0.5
4	0.61	2.79	0.333	-0.5	0.167
5	1.17	3.95	-0.333	0.167	-0.5
6	0.83	5.68	0.333	-0.5	0.333
7	0.94	-0.10	1	0.167	-0.167

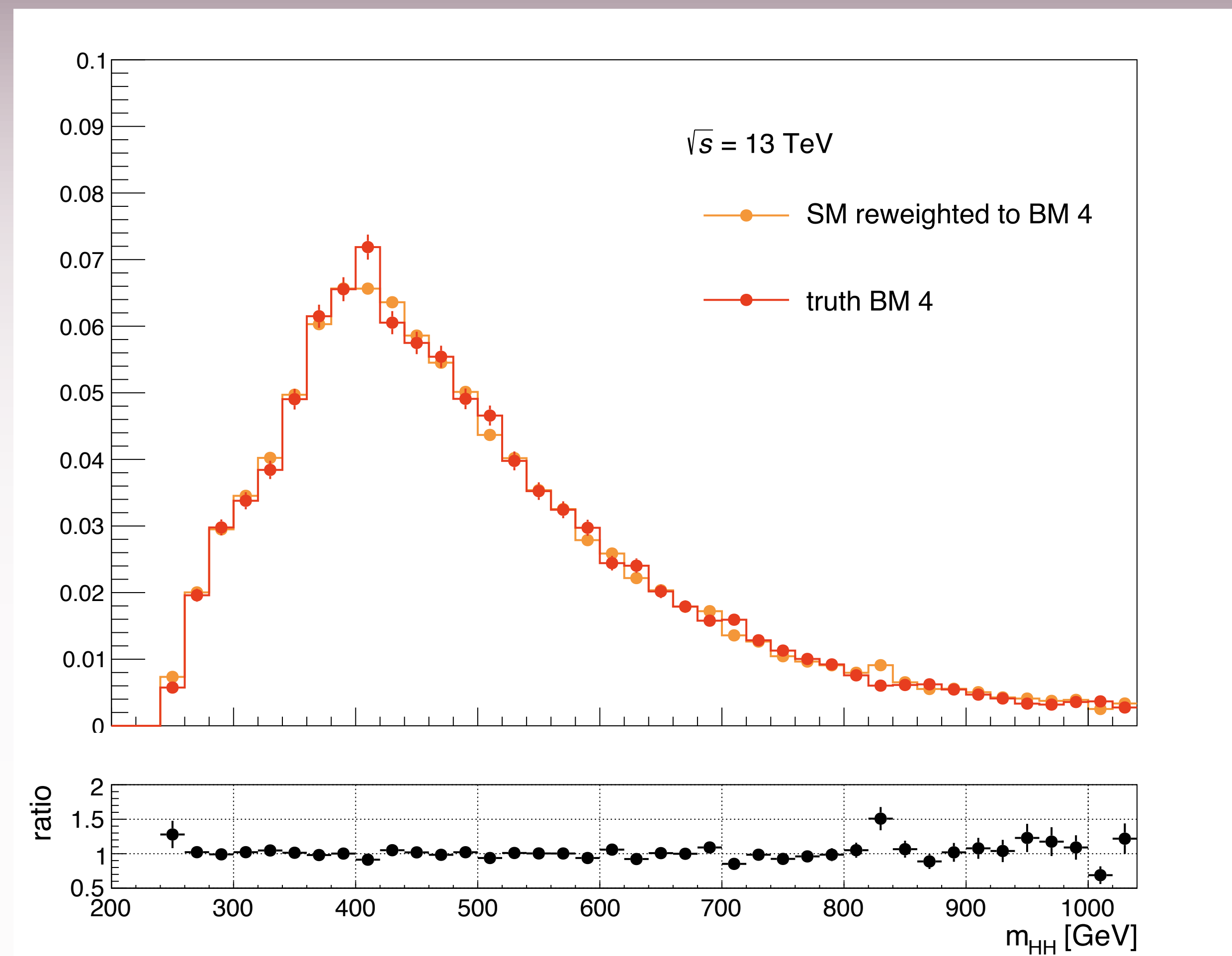
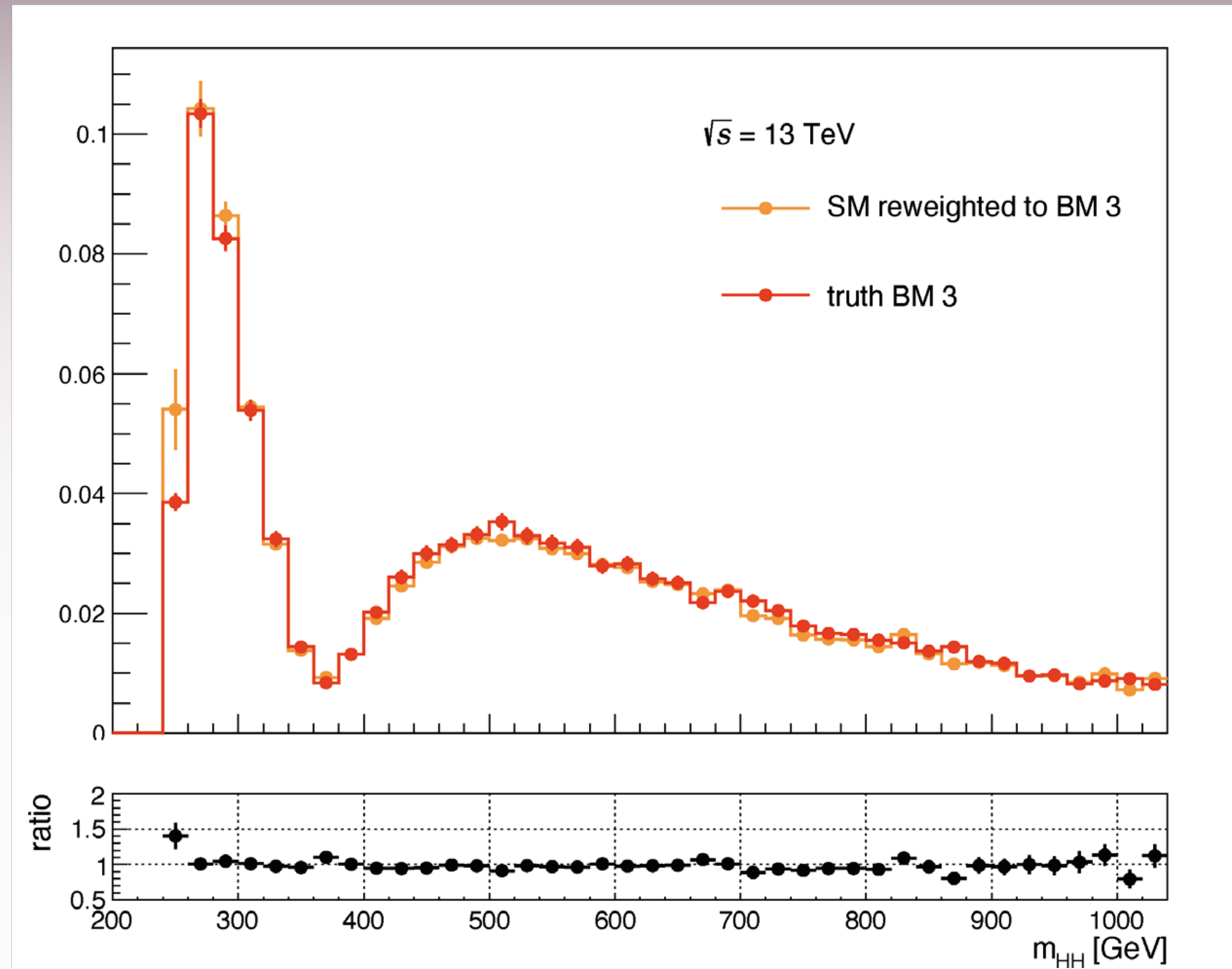
Weights are given by an analytic parametrisation

$$\frac{d\sigma}{dM_{hh}} / \frac{d\sigma^{SM}}{dM_{hh}} = \sum_i A_i C_i$$

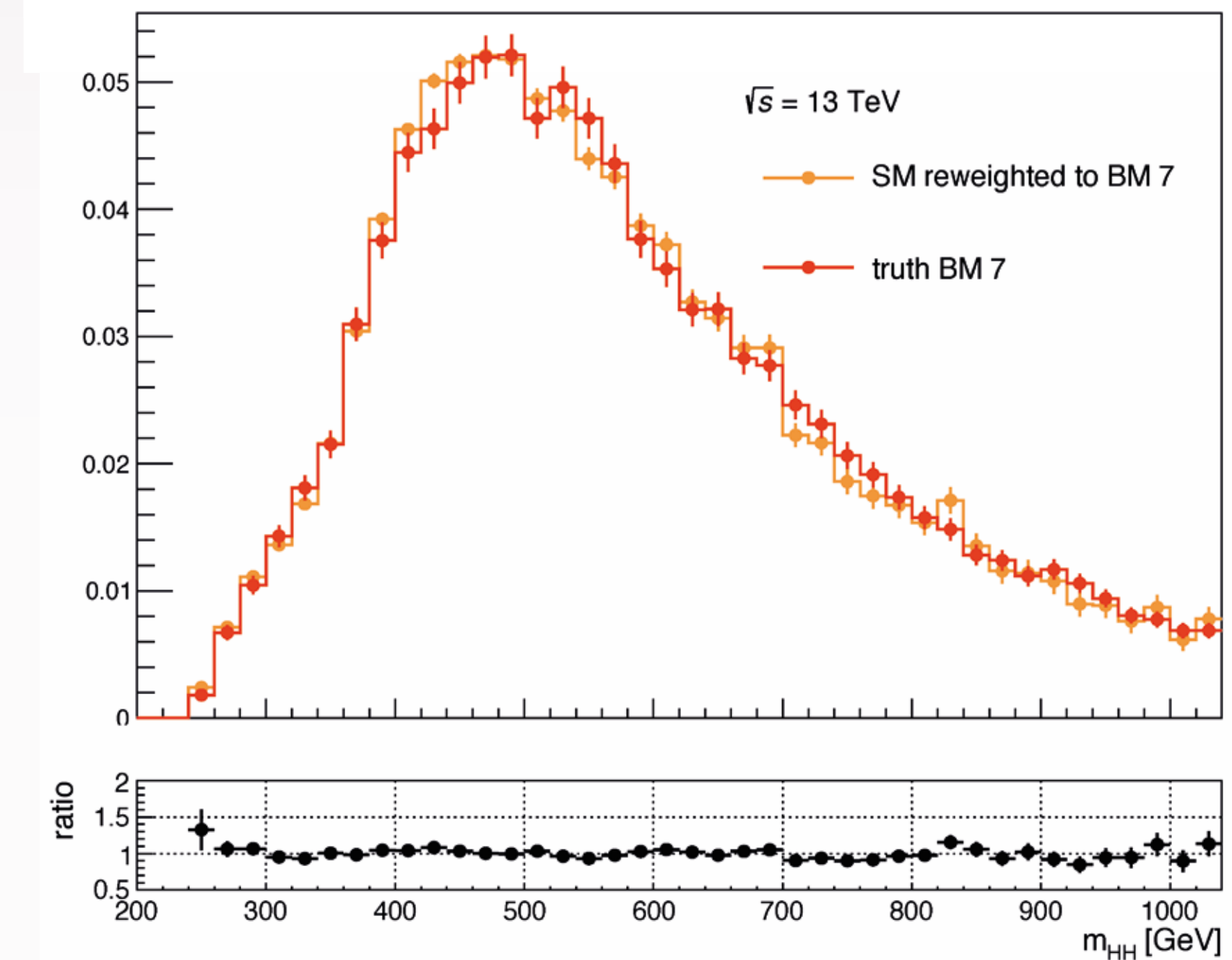
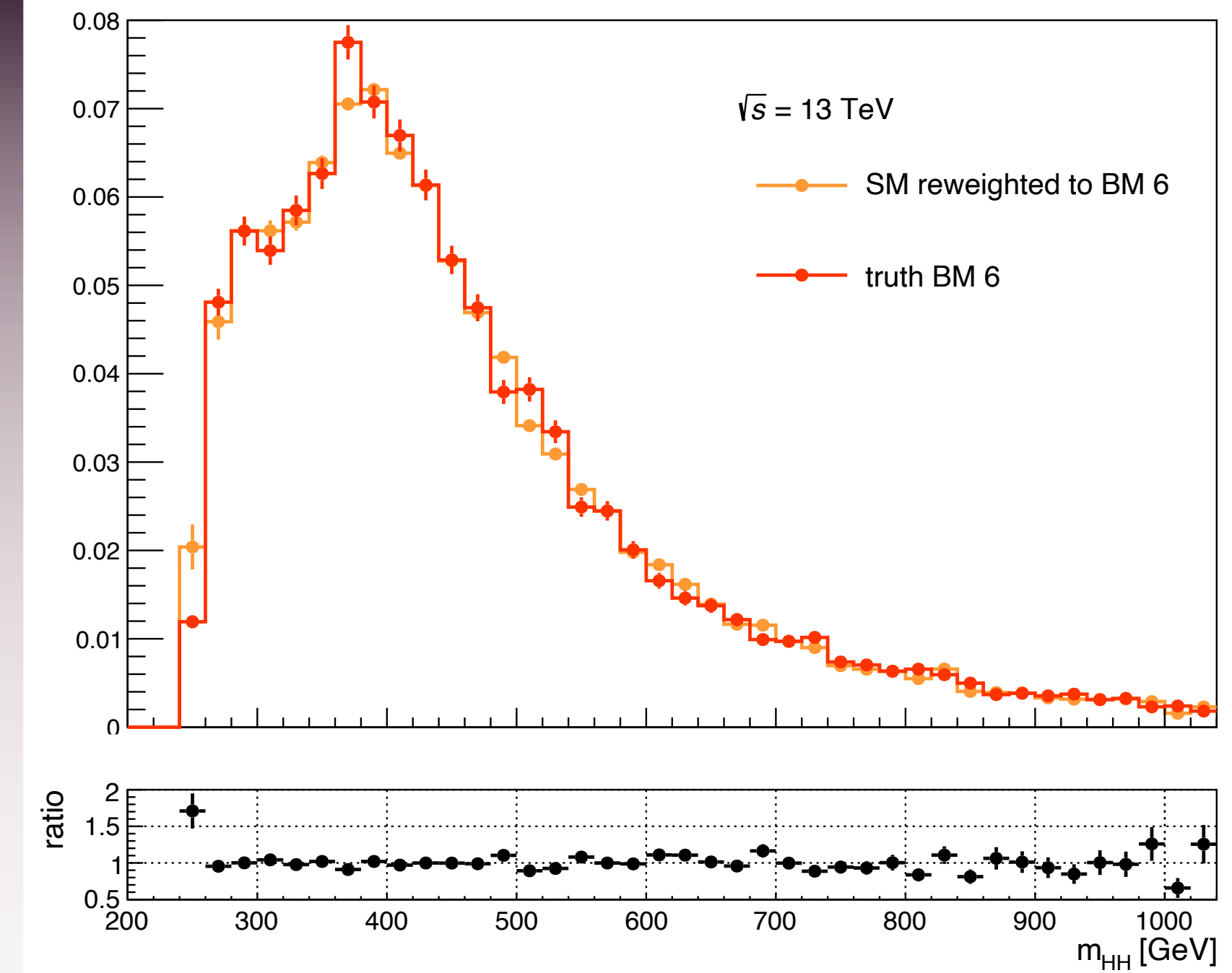
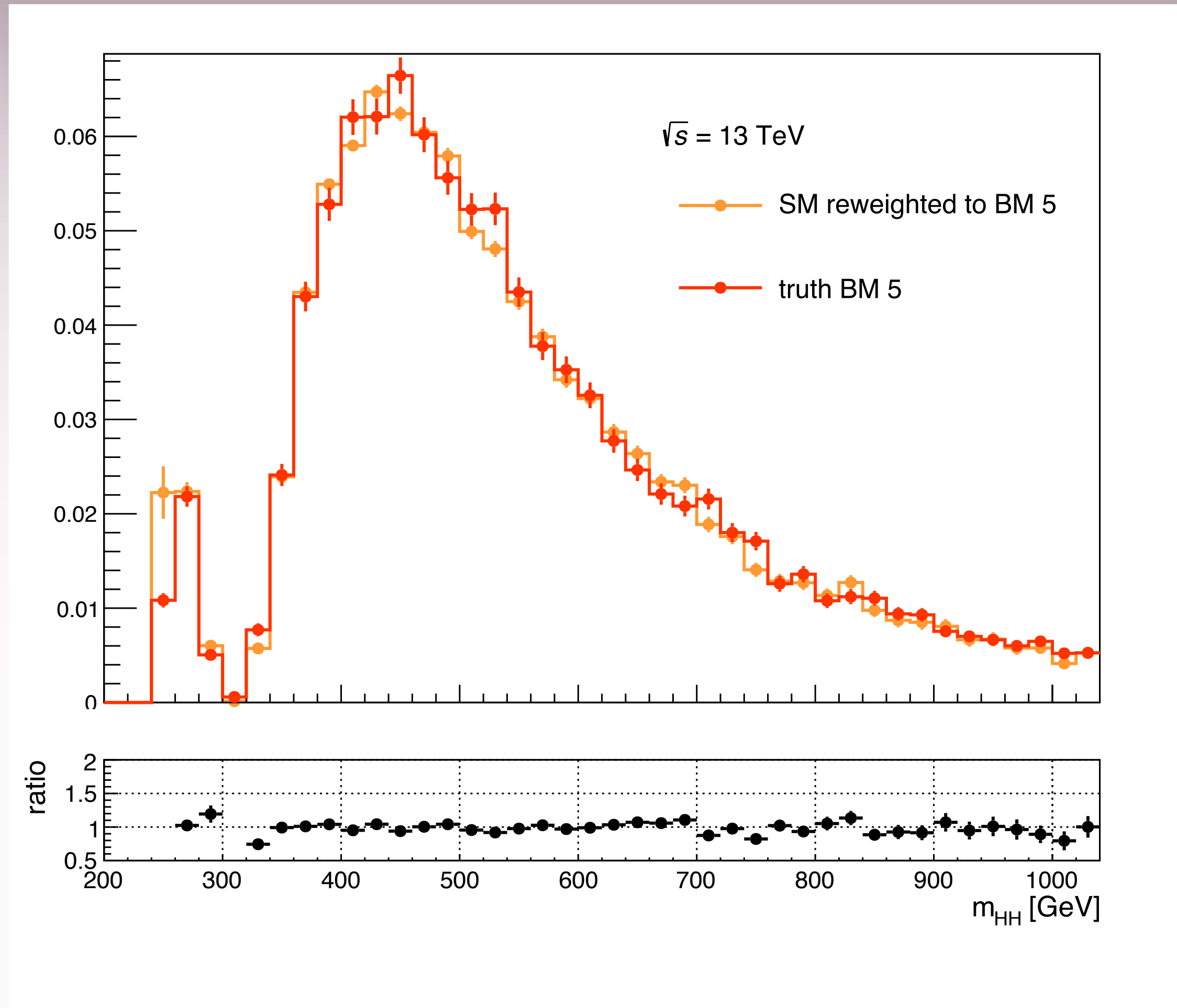
# Reweighting validation @ NLO



# Reweighting validation @ NLO



# Reweighting validation @ NLO





# CMS

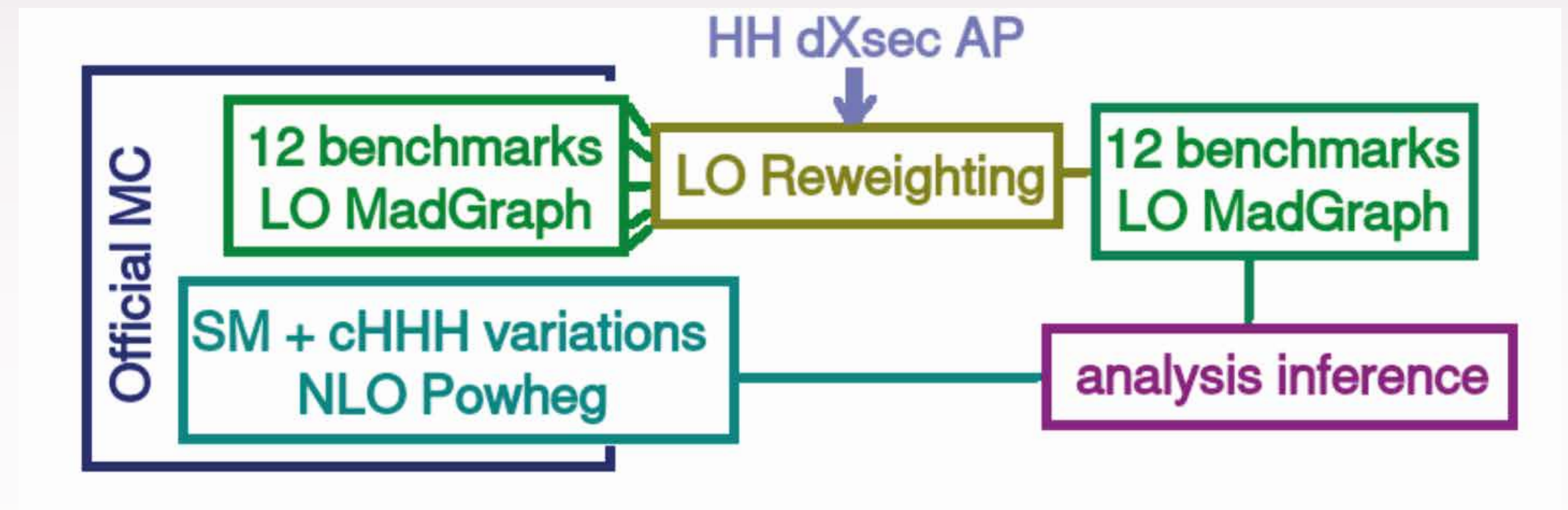
by

P. Mandrik and S. Slabospitskii

# Typical CMS non-resonant ggF HH workflow

at the moment :

- Official MC for representative benchmarks in MadGraph at LO only
- Powheg NLO generator and analytical parametrisation (AP) of ggF HH of differential cross section are available



Can we incorporate NLO reweighting of LO MadGraph/NLOPowheg samples into ongoing analyses and how will it work?

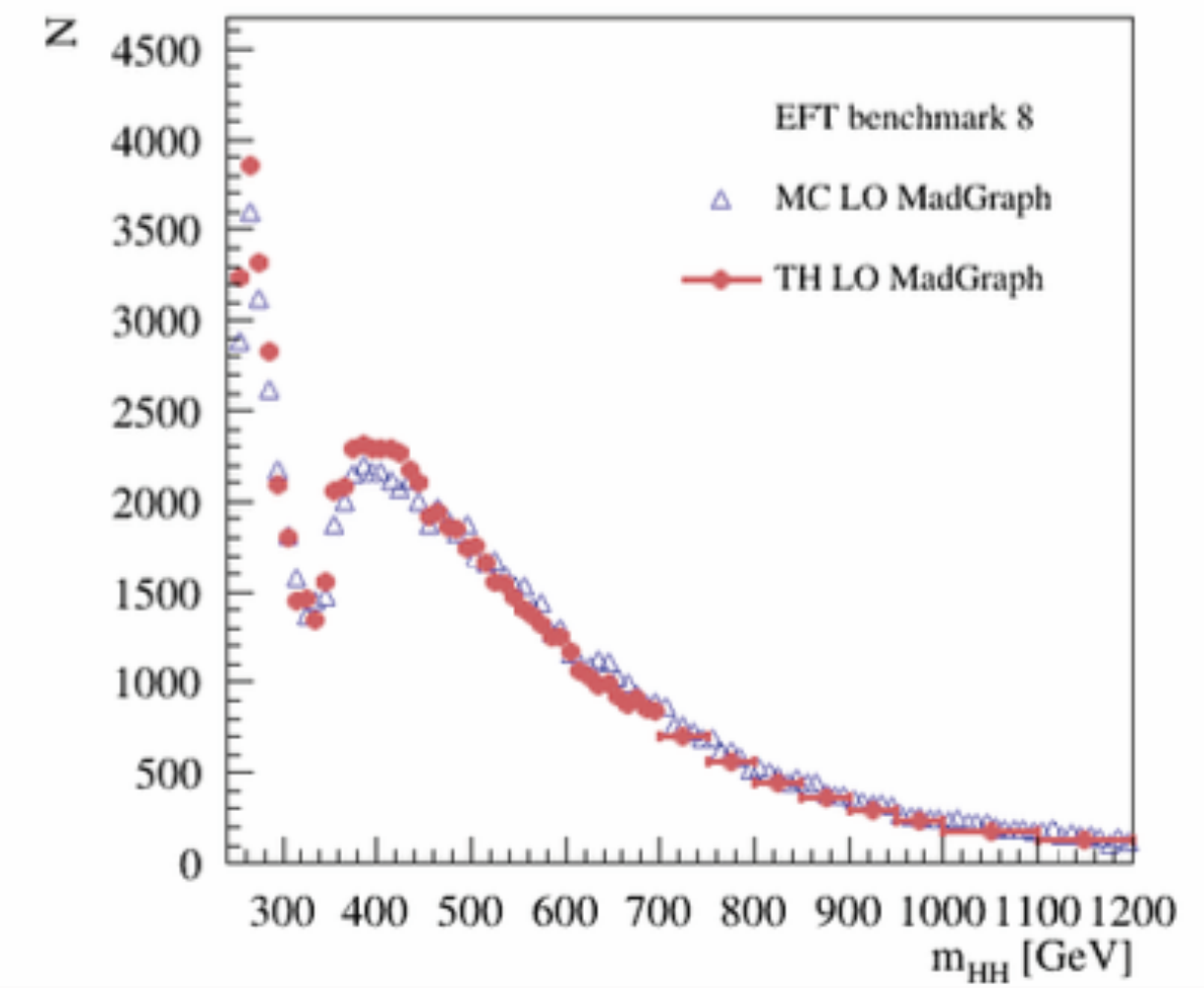
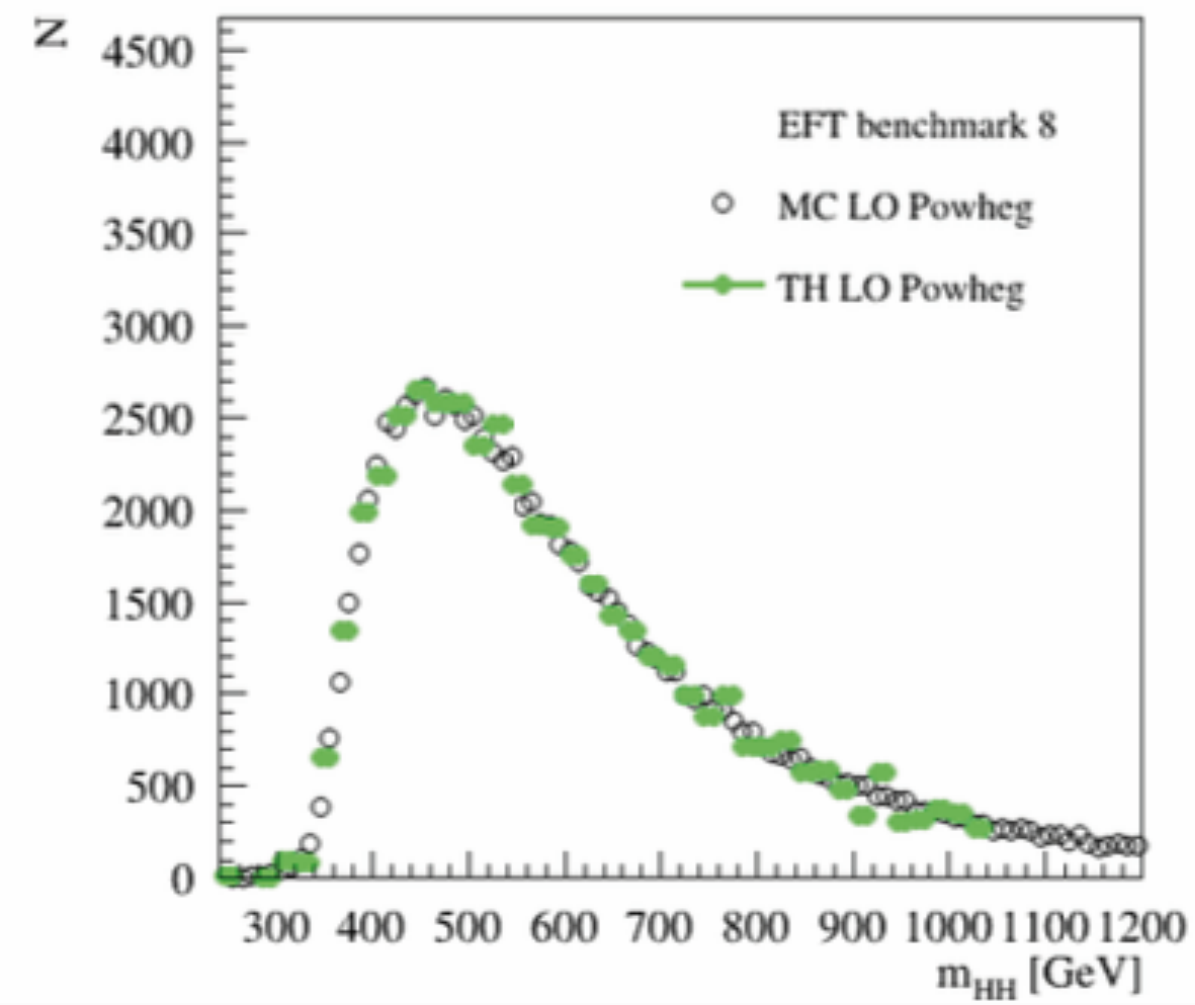
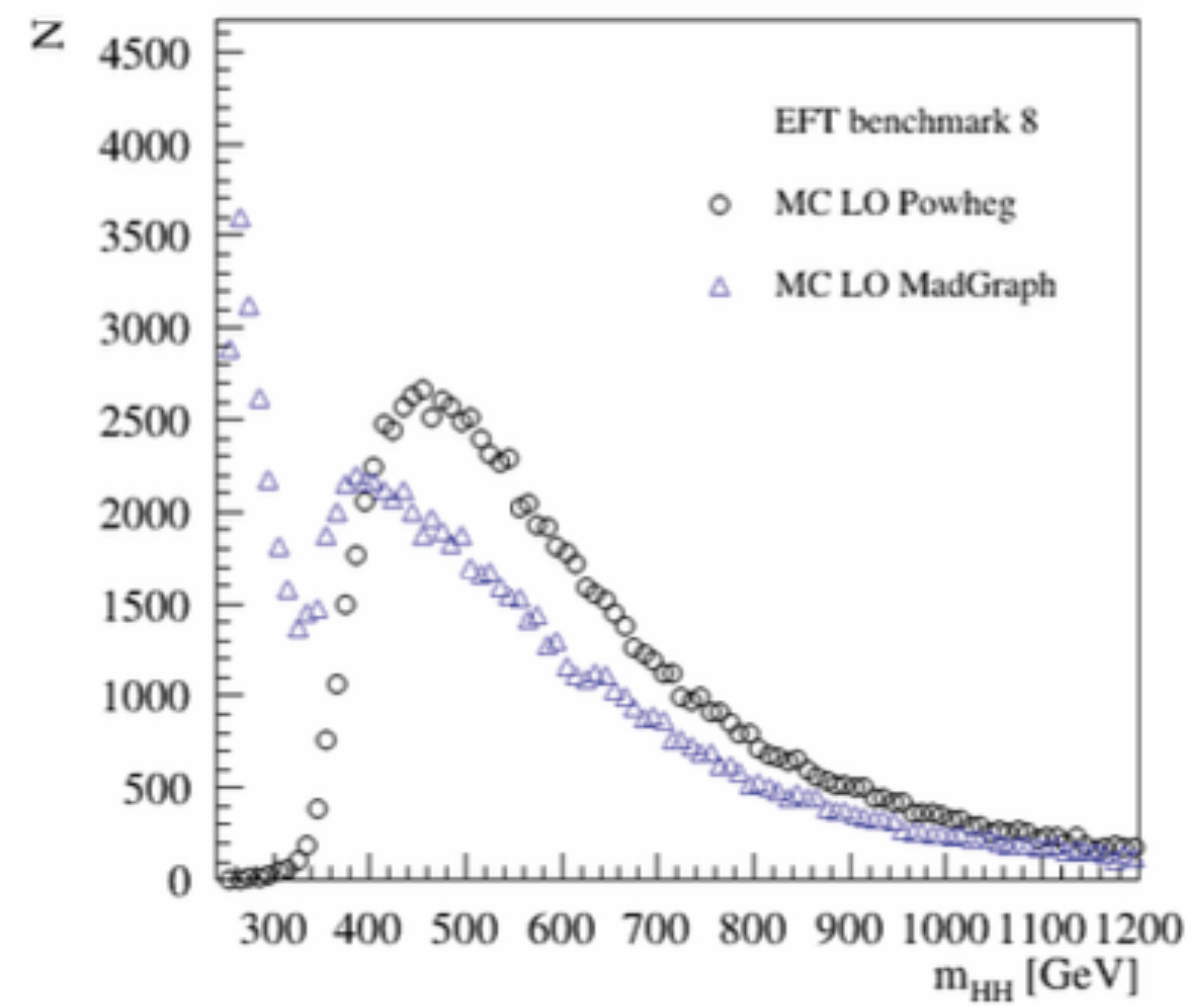
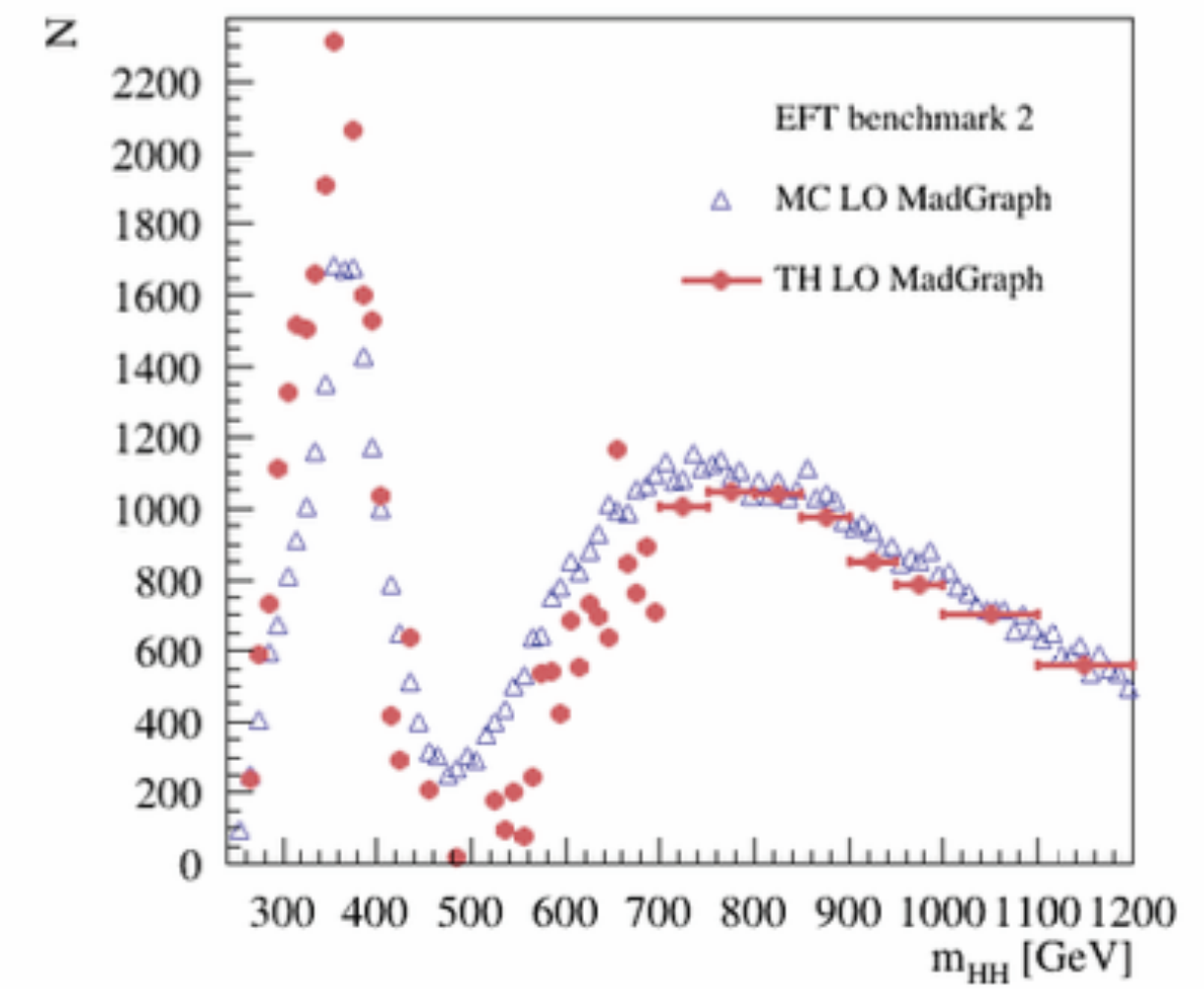
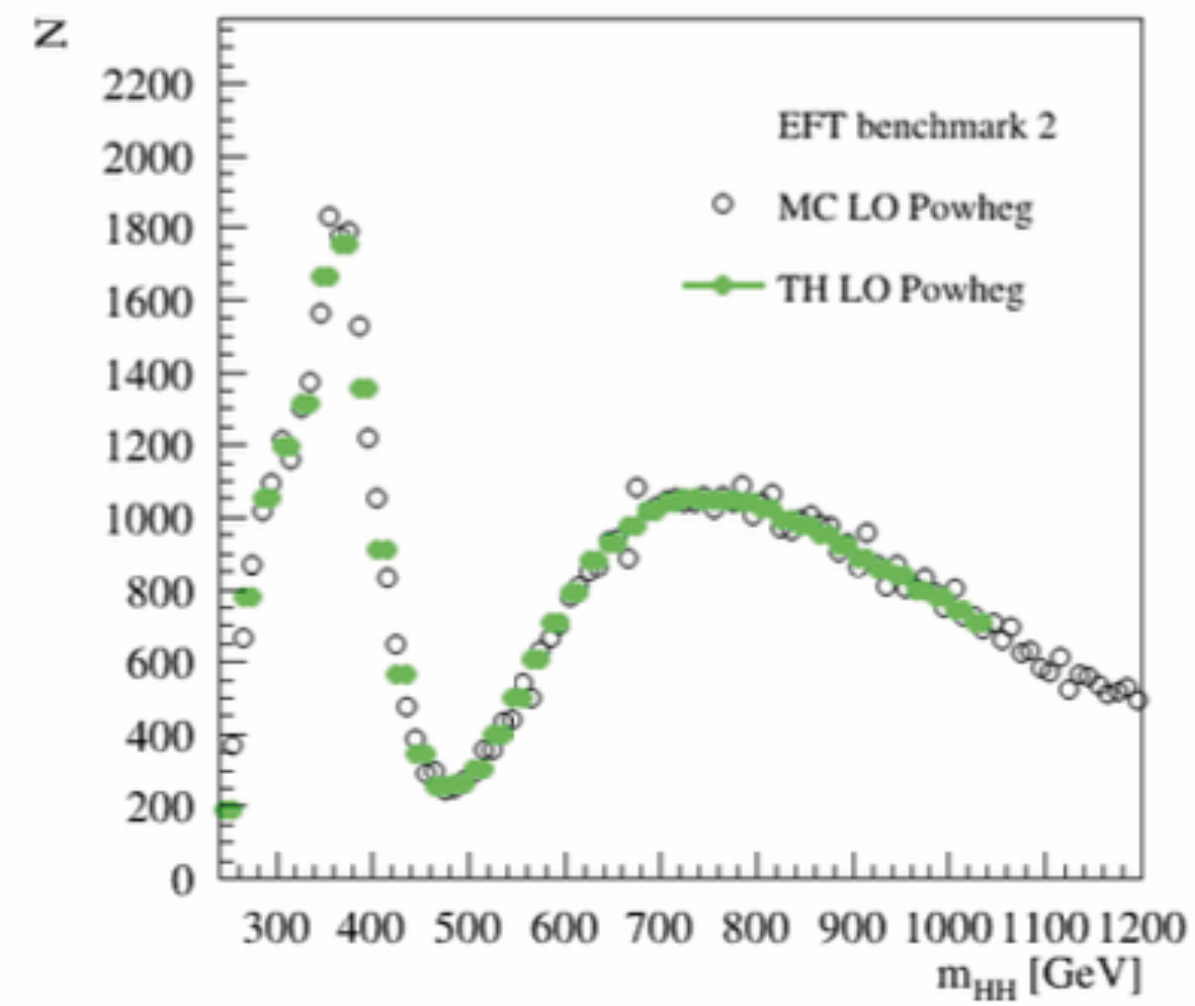
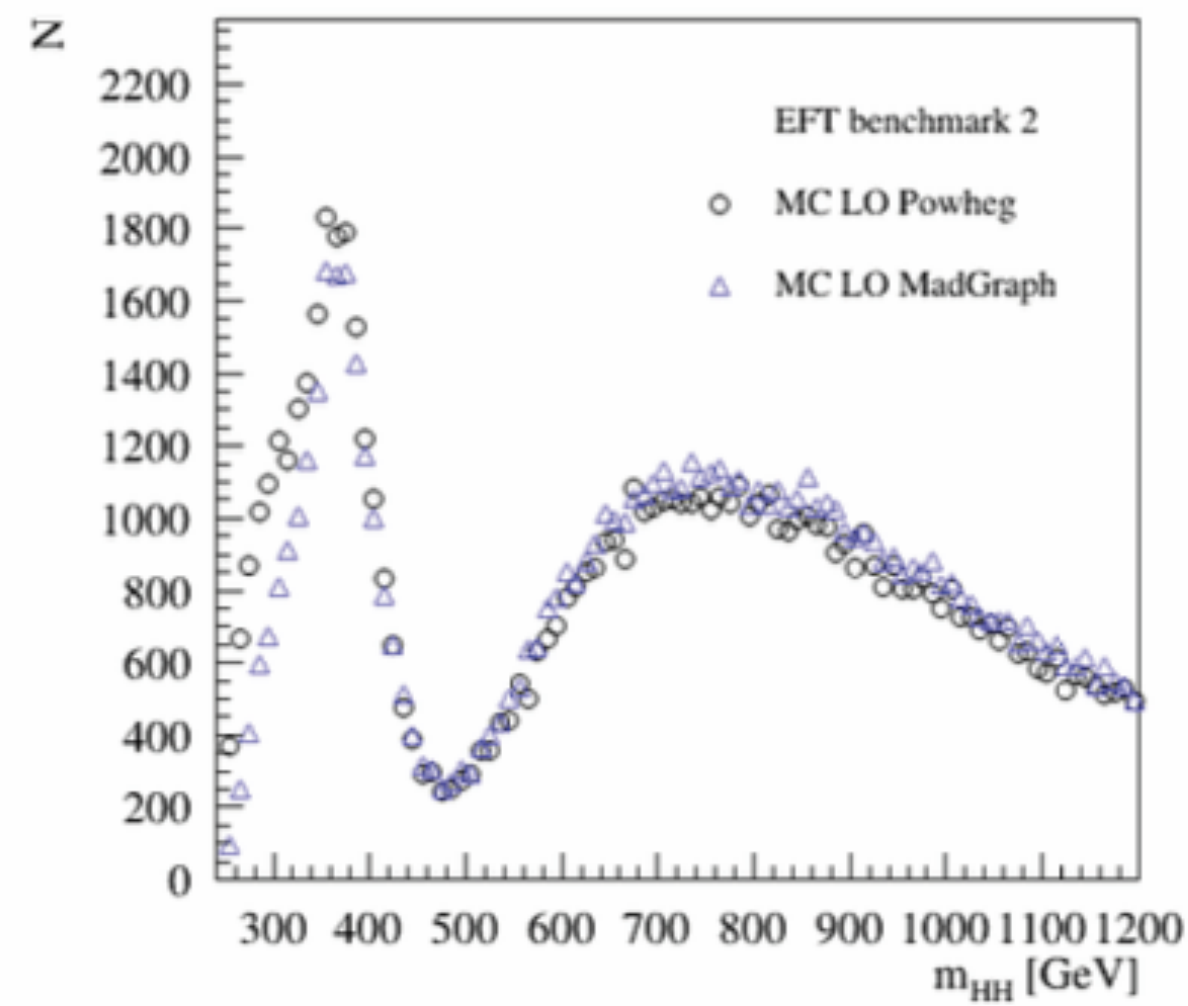
# Input

Samples produced privately for SM + 12 EFT benchmarks + extra benchmarks (43 EFT points in total):

- Madgraph LO MC private samples, input configuration cards  
[https://github.com/pmandrik/HH\\_pair/tree/master/generation](https://github.com/pmandrik/HH_pair/tree/master/generation)
- Powheg LO MC private samples, 100K events, generated at full theory, User-Processes-V2/ggHH, input configuration cards (13 TeV)  
[https://github.com/pmandrik/HH\\_pair/tree/master/generation](https://github.com/pmandrik/HH_pair/tree/master/generation)
- Powheg NLO MC private samples, up to 100K events, generated at full theory, User-Processes-V2/ggHH, input configuration cards (13 TeV):  
[https://github.com/pmandrik/HH\\_pair/tree/master/generation](https://github.com/pmandrik/HH_pair/tree/master/generation)
- Analytical parameterization of  $\frac{d\sigma}{dM_{hh}}$  at LO (Madgraph & Powheg based) [A. Carvalho et al '15] MG and [G.Buchalla et al '18] Powheg at NLO (Powheg based) [G.Buchalla et al '18]

C++ implementation of the prediction from analytical parametrisations is available in  
[https://github.com/pmandrik/VSEVA/blob/master/HHWWgg/reweight/reweight\\_HH.C](https://github.com/pmandrik/VSEVA/blob/master/HHWWgg/reweight/reweight_HH.C)

Benchmark	$\kappa_\lambda$	$\kappa_t$	$c_2$	$c_g$	$c_{2g}$
1	7.5	1.0	-1.0	0.0	0.0
2	1.0	1.0	0.5	-0.8	0.6
3	1.0	1.0	-1.5	0.0	-0.8
4	-3.5	1.5	-3.0	0.0	0.0
5	1.0	1.0	0.0	0.8	-1
6	2.4	1.0	0.0	0.2	-0.2
7	5.0	1.0	0.0	0.2	-0.2
8	15.0	1.0	0.0	-1	1
9	1.0	1.0	1.0	-0.6	0.6
10	10.0	1.5	-1.0	0.0	0.0
11	2.4	1.0	0.0	1	-1
12	15.0	1.0	1.0	0.0	0.0
SM	1.0	1.0	0.0	0.0	0.0



# Intermediate conclusions

- Good Powheg/MadGraph and MC/AP agreement for most of the benchmark, but BM-8 for MadGraph vs Powheg and BM-2 for MG vs AP
- CMS official MC samples generated using MadGraph & NLO k-factors from Powheg  
Then reweight to take differences in generators/models prediction

**(MadGraph LO) → (Powheg LO) → (Powheg NLO)**

- Analytical parameterization from MadGraph is defined for LO only
- Analytical parameterization from Powheg is defined up to  $M_{hh} < 1040$  GeV and does not take into account dependence on  $\cos \theta^*$

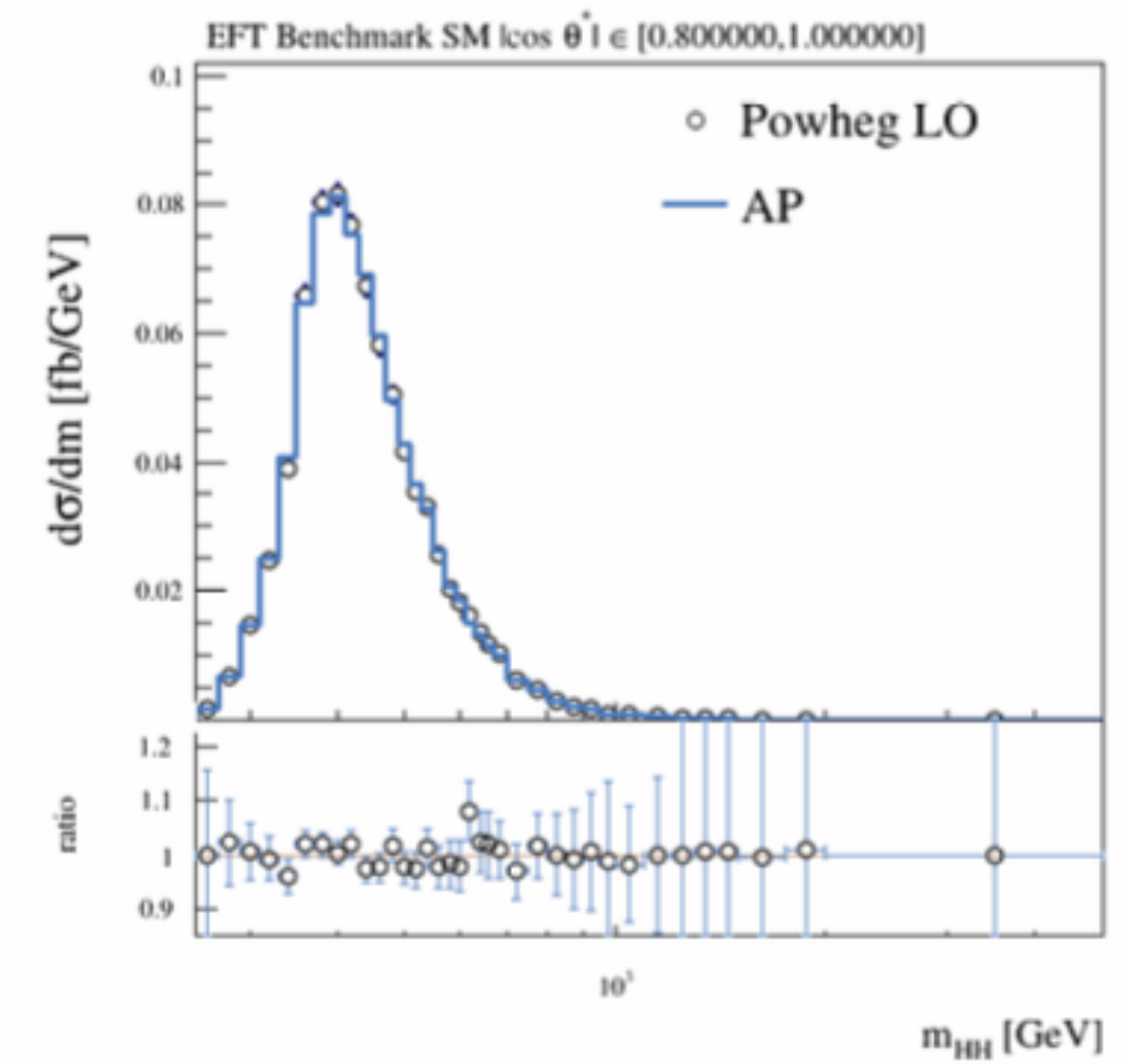
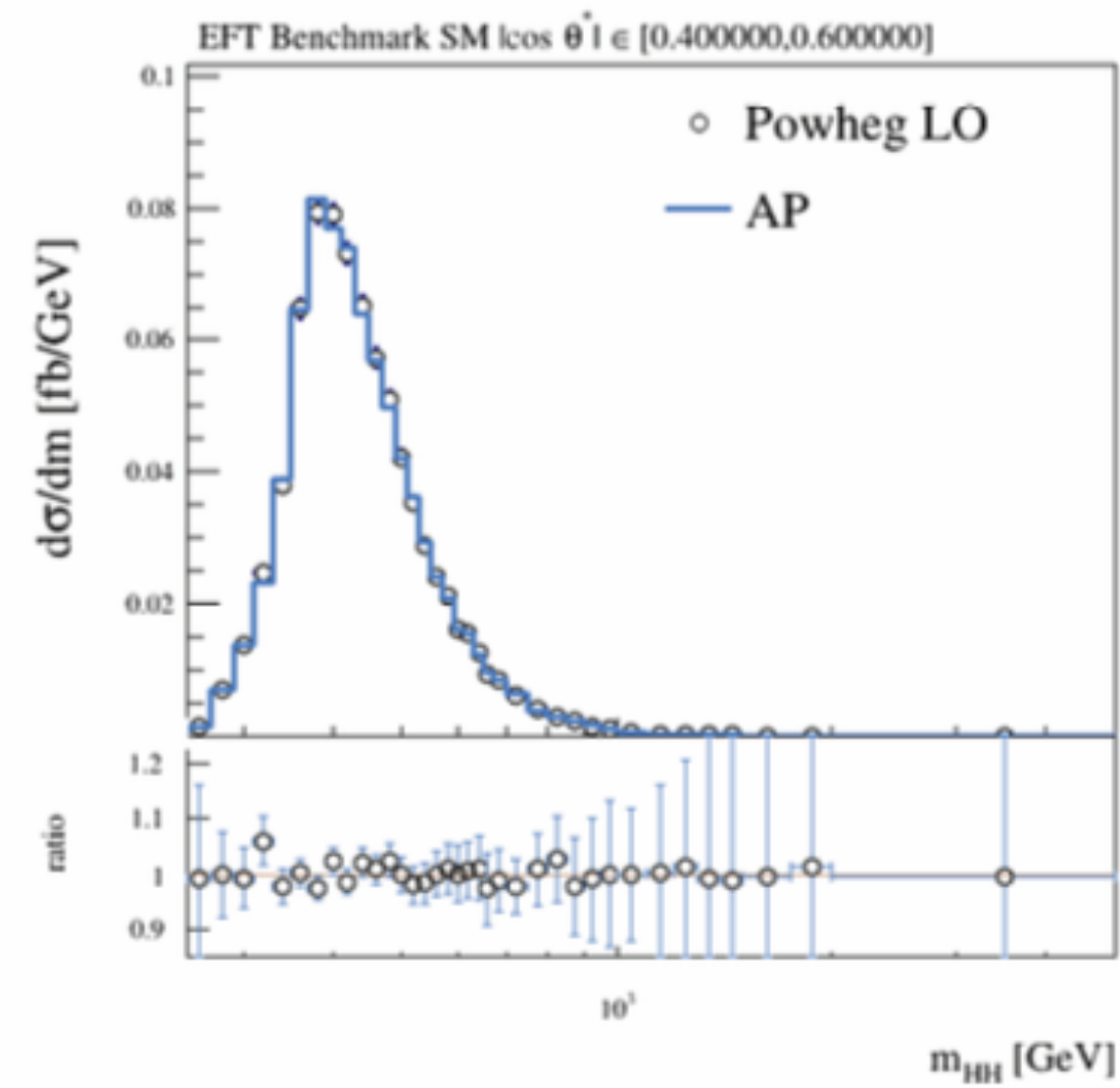
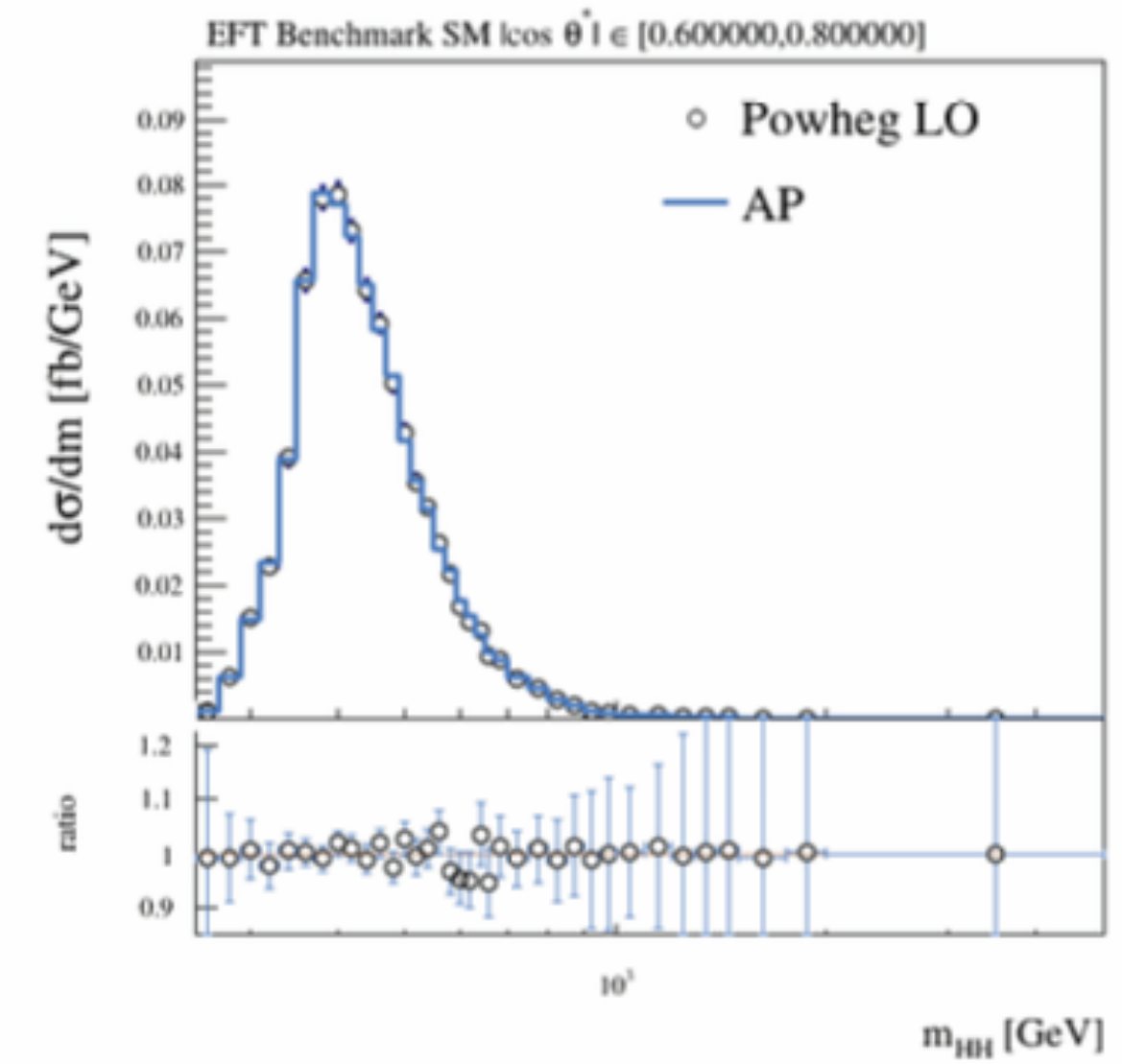
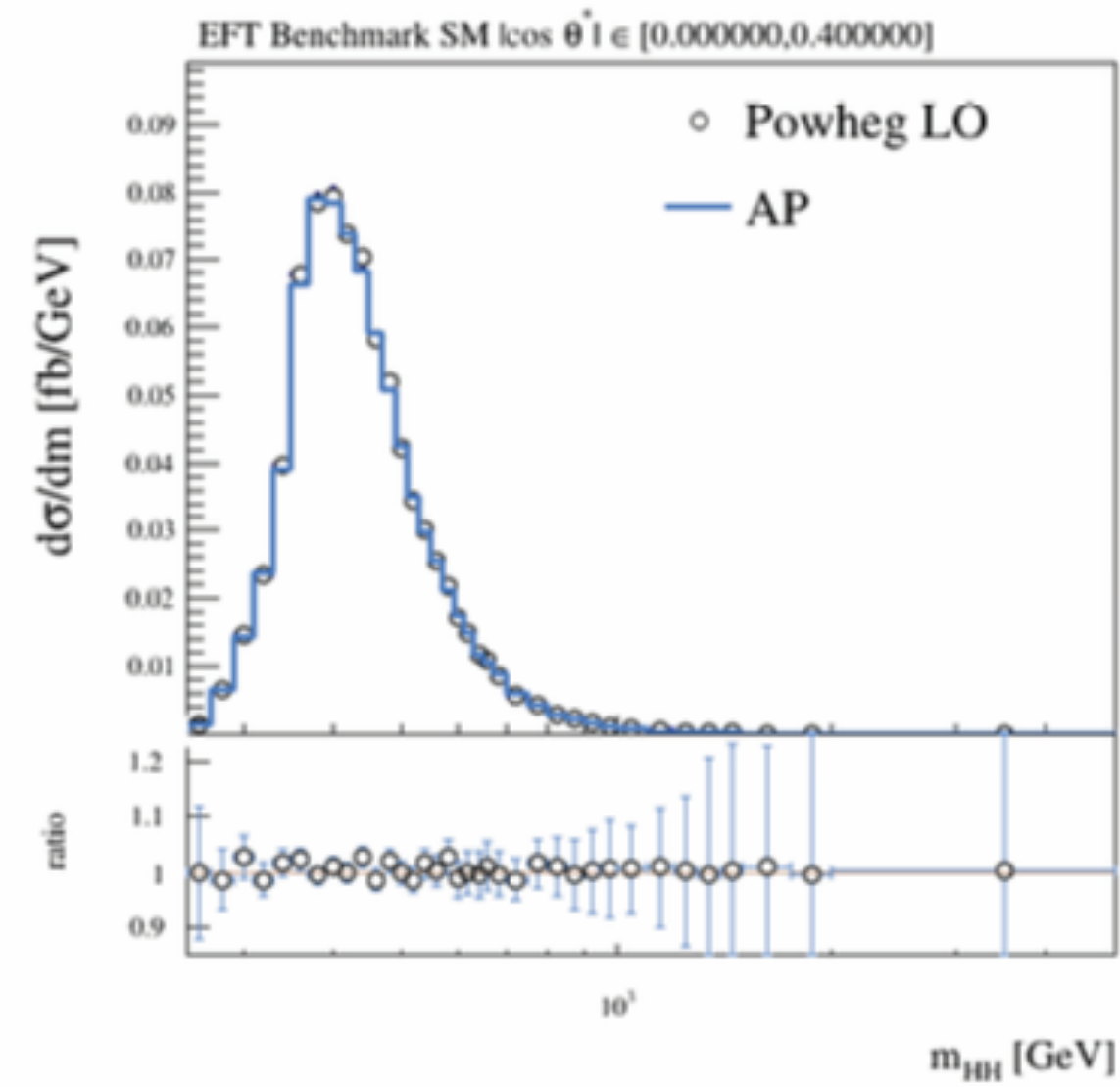
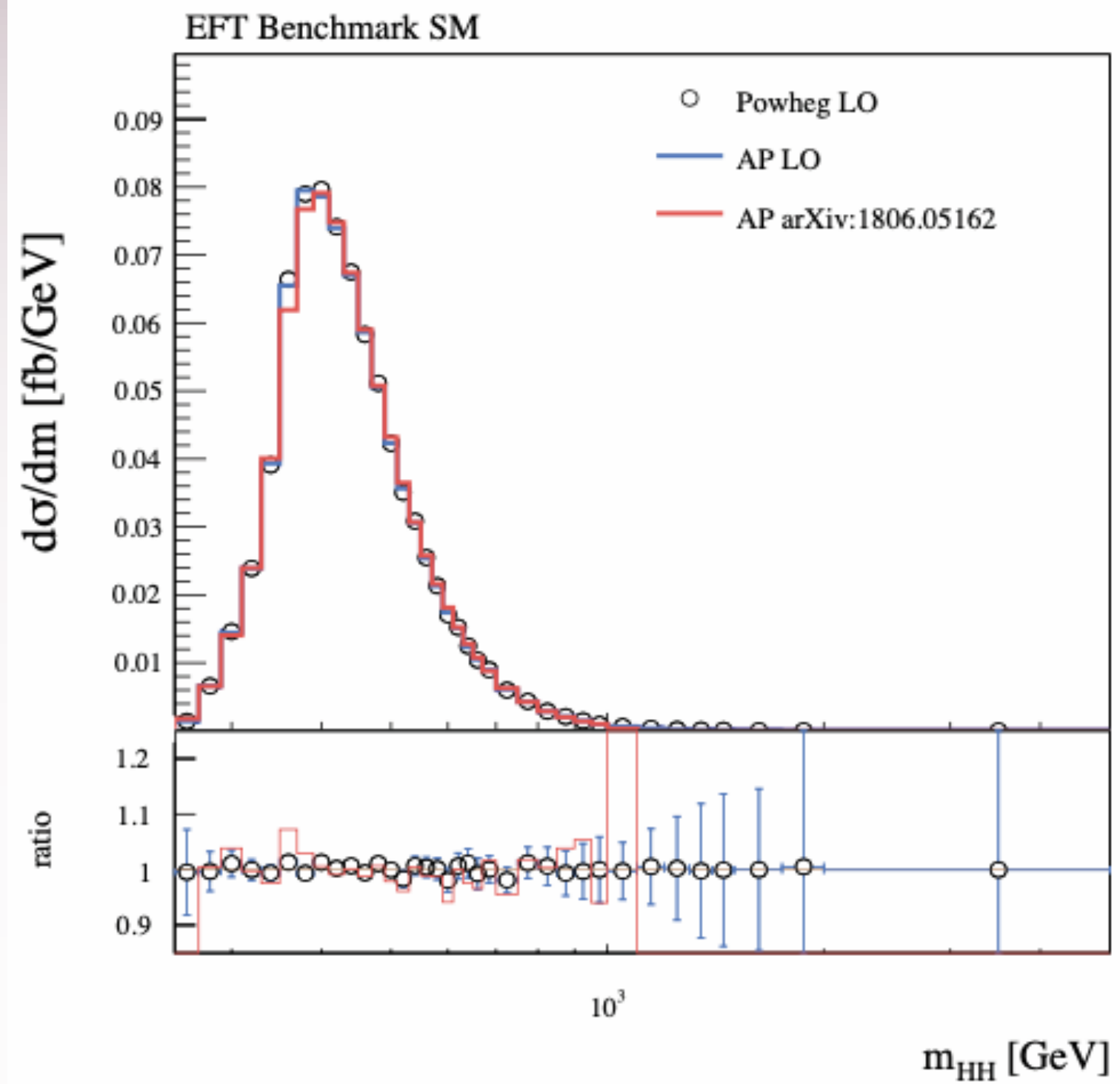
## Solution:

- Derive custom  $A_*$  coefficients using private MC (35/9 benchmarks for fit/validation with  $\approx 10^6$  events) for analytical parameterization:

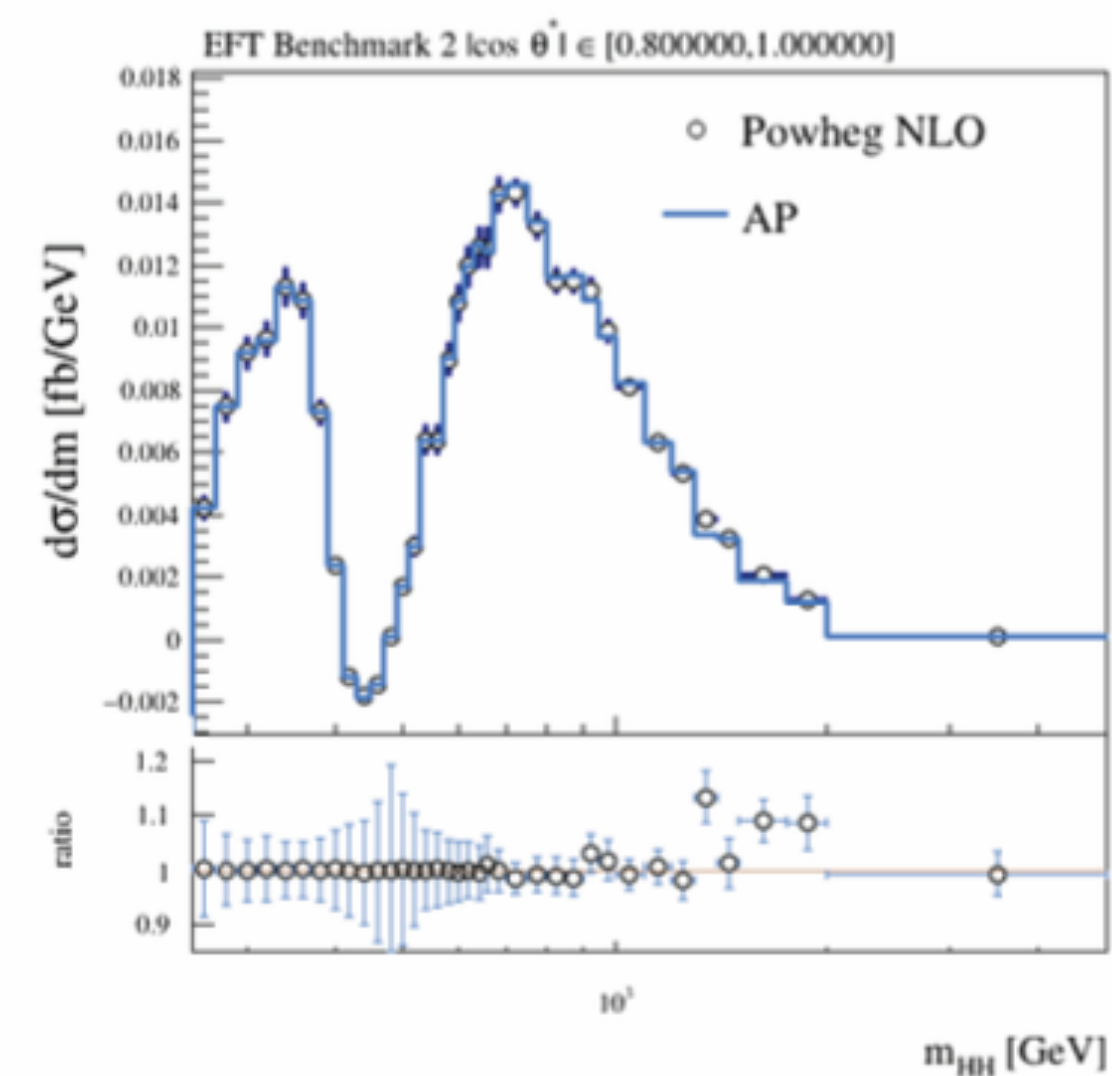
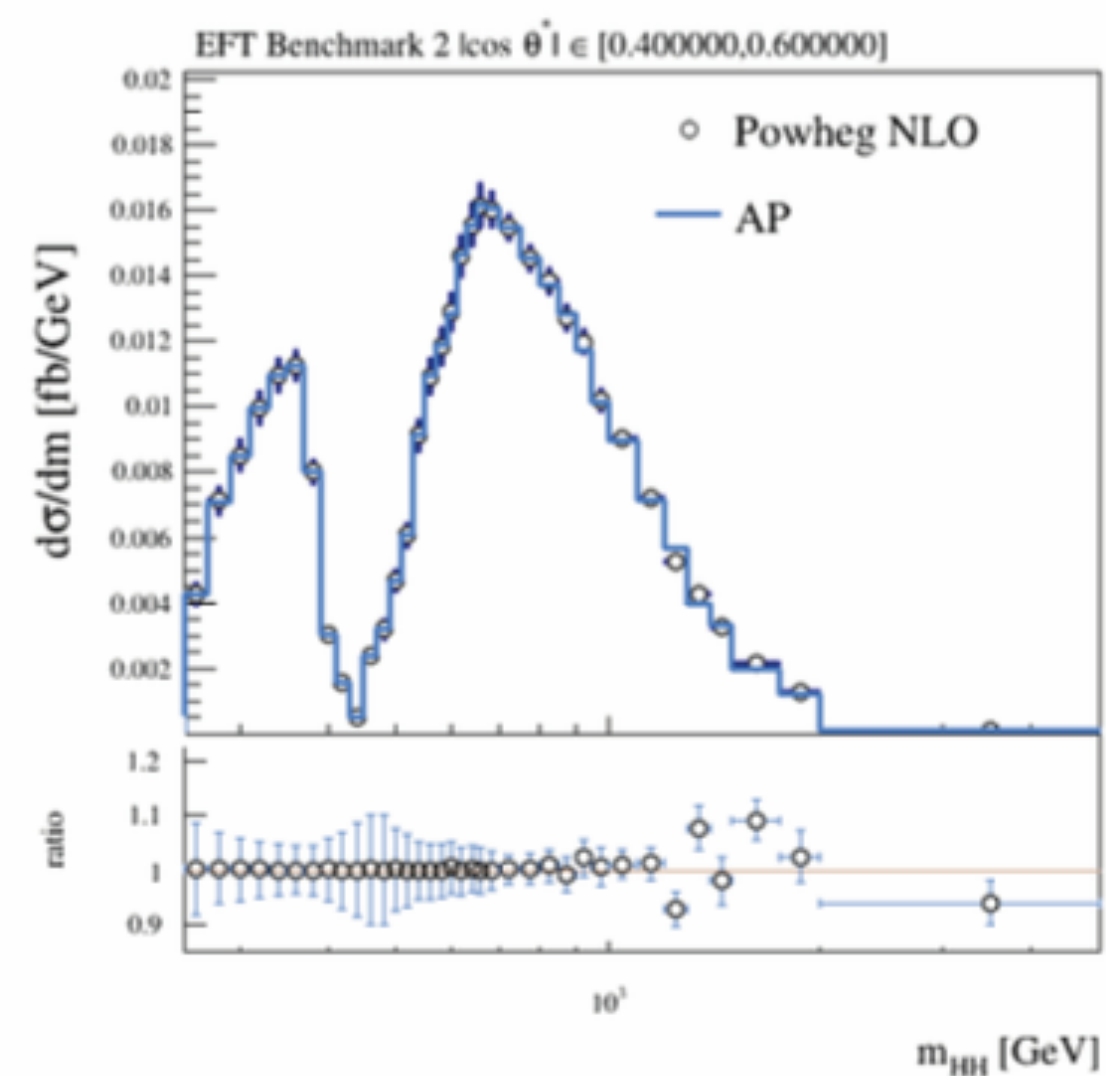
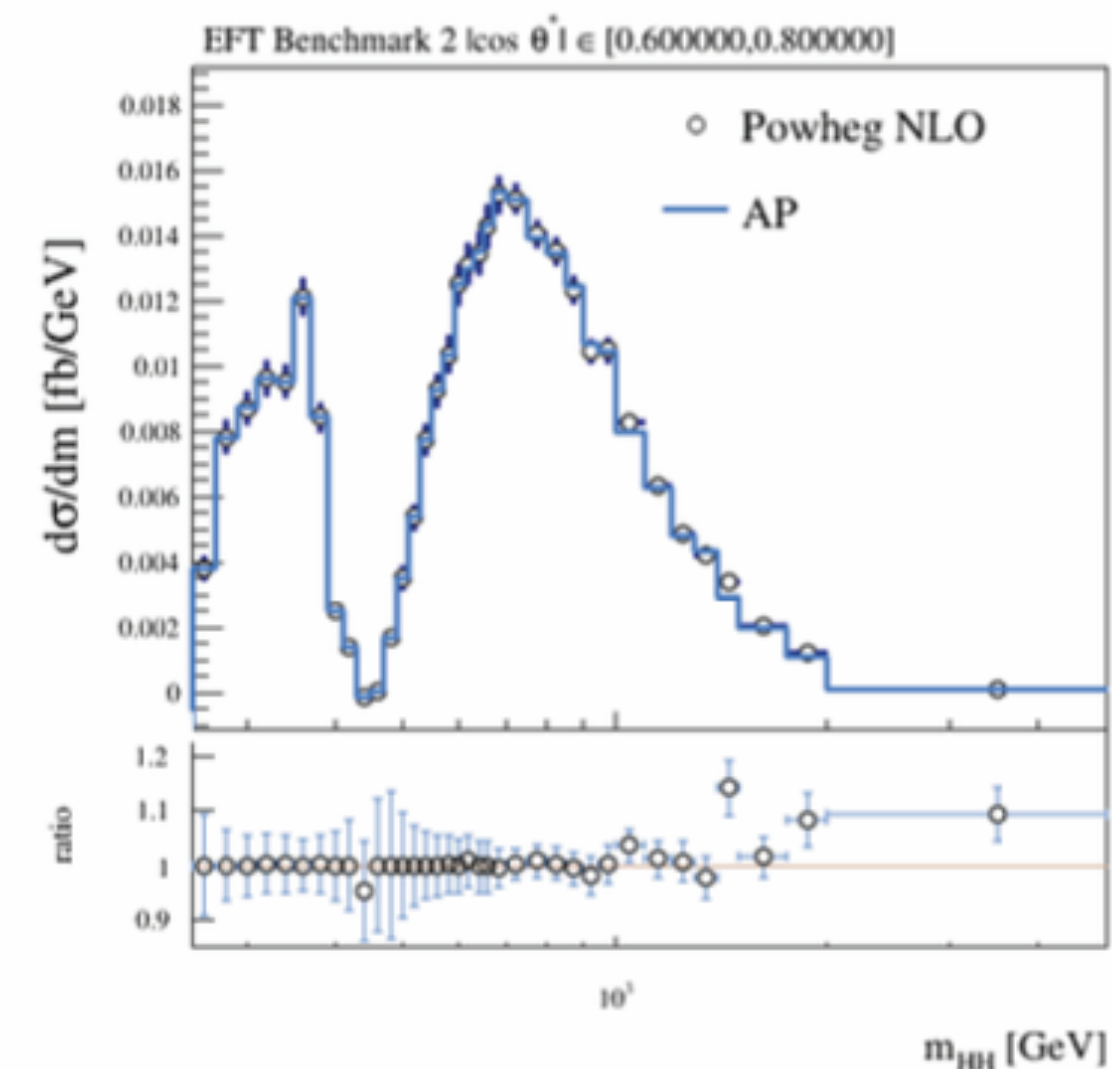
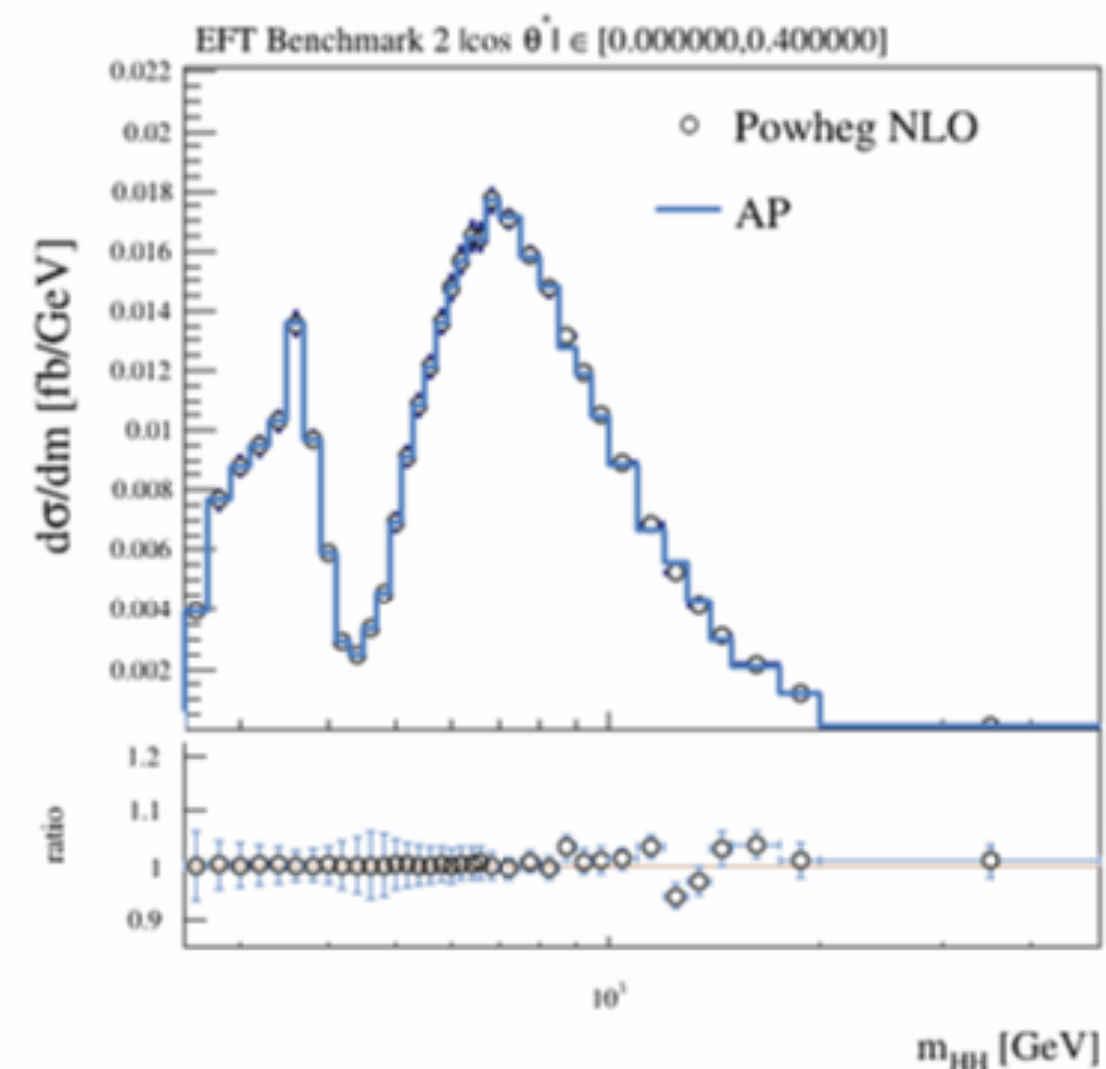
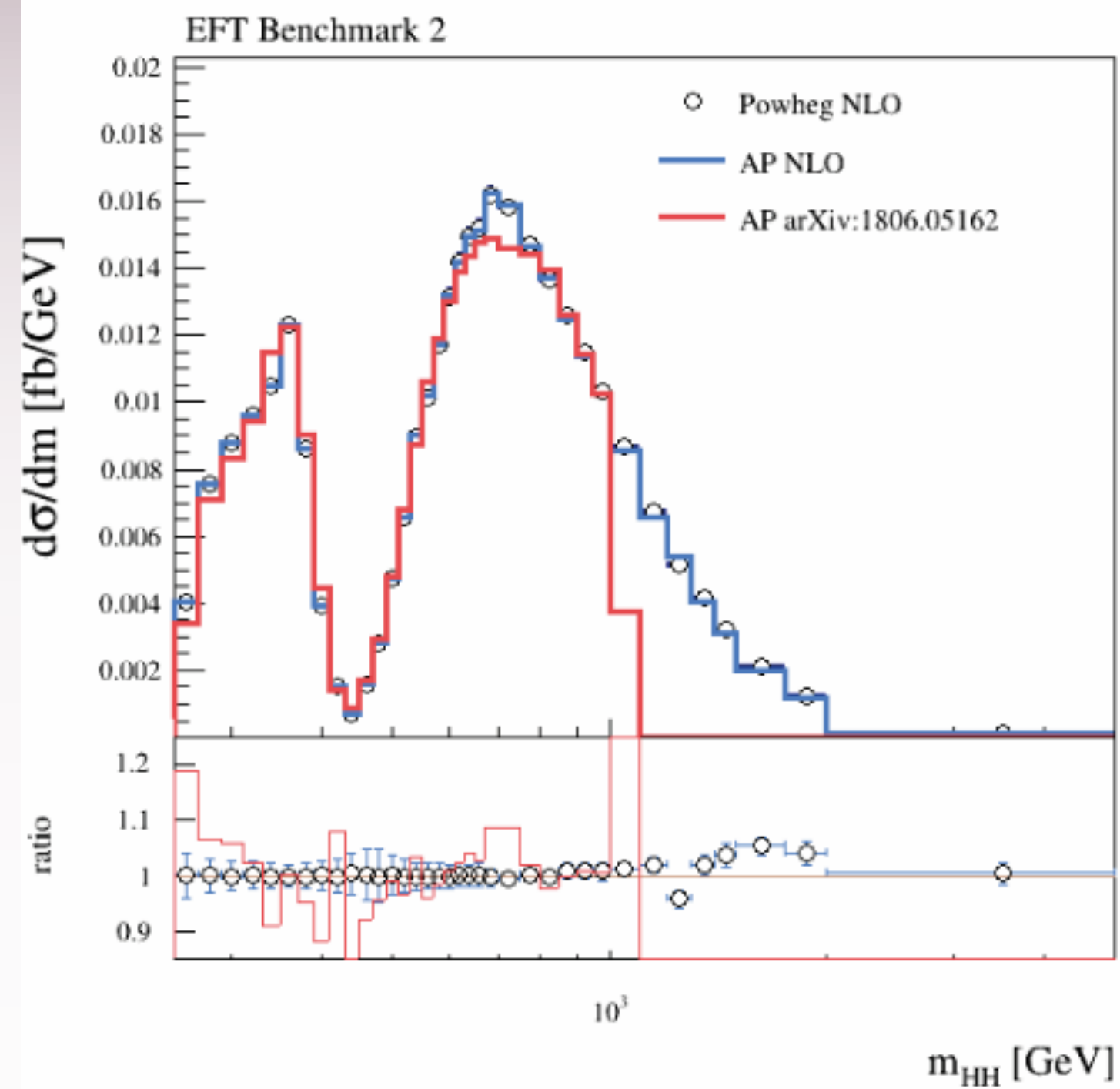
$$\frac{d\sigma}{dM_{hh}d|\cos \theta^*|} = \sum A_i(M_{hh}, |\cos \theta^*|) c_i$$

- To extract the  $A_i(M_{hh}, |\cos \theta^*|)$  coefficients the system of equations is solved by performing a minimisation in ROOT MINUIT.

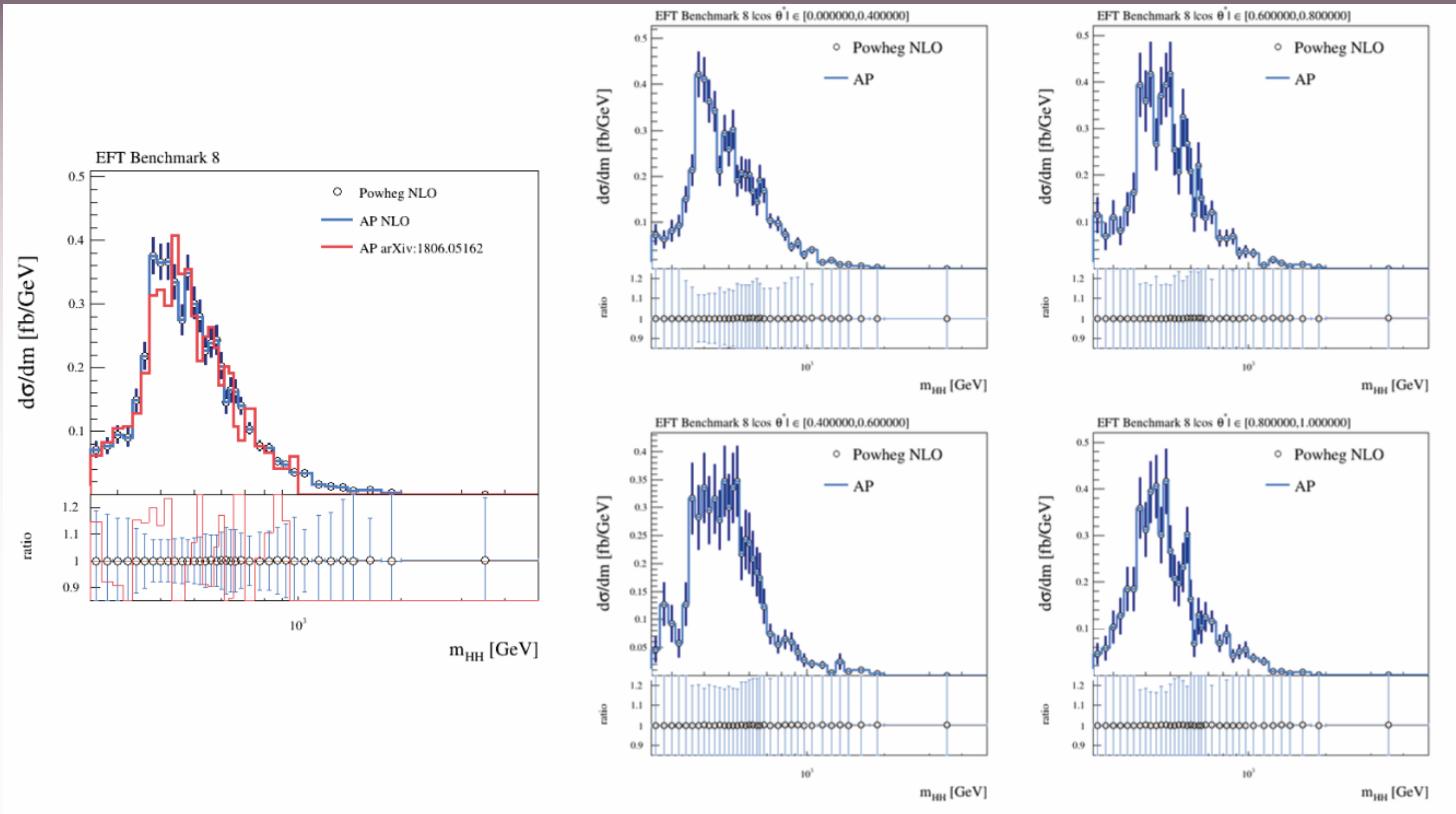
# AP @ LO results



# AP @ LO results



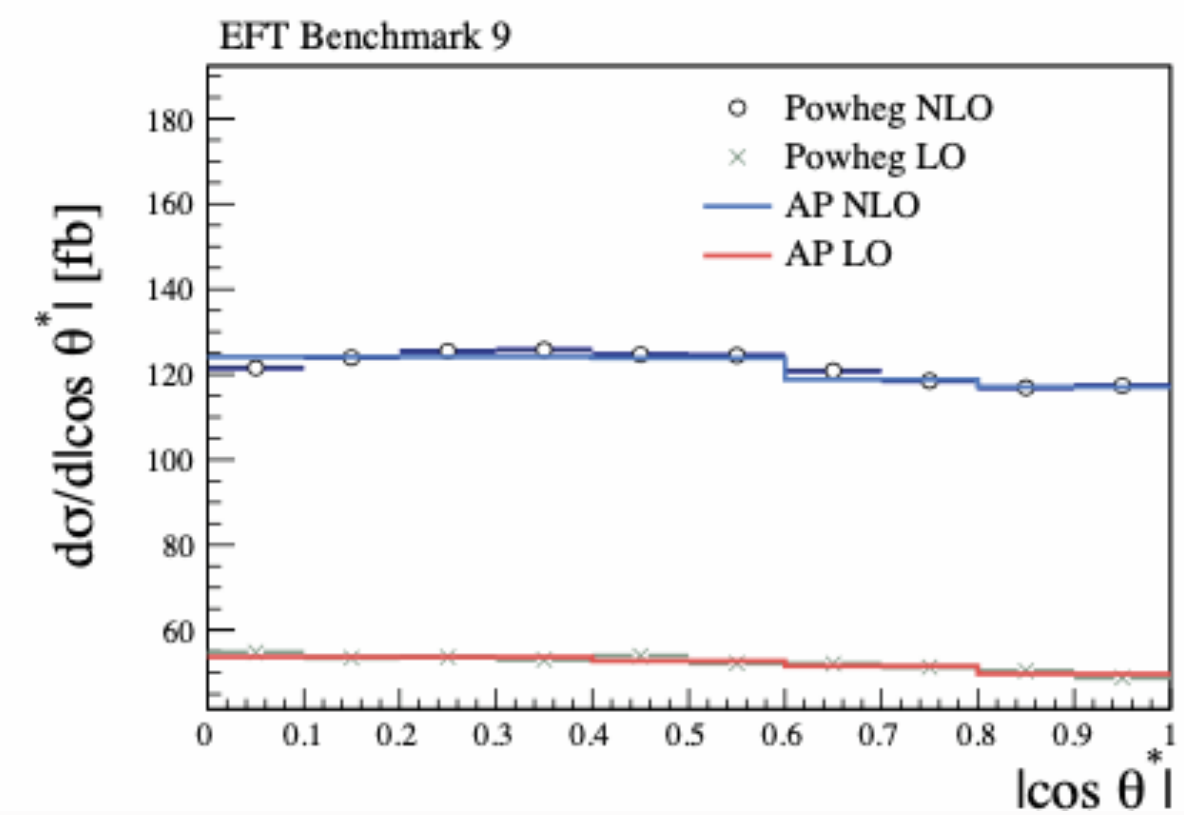
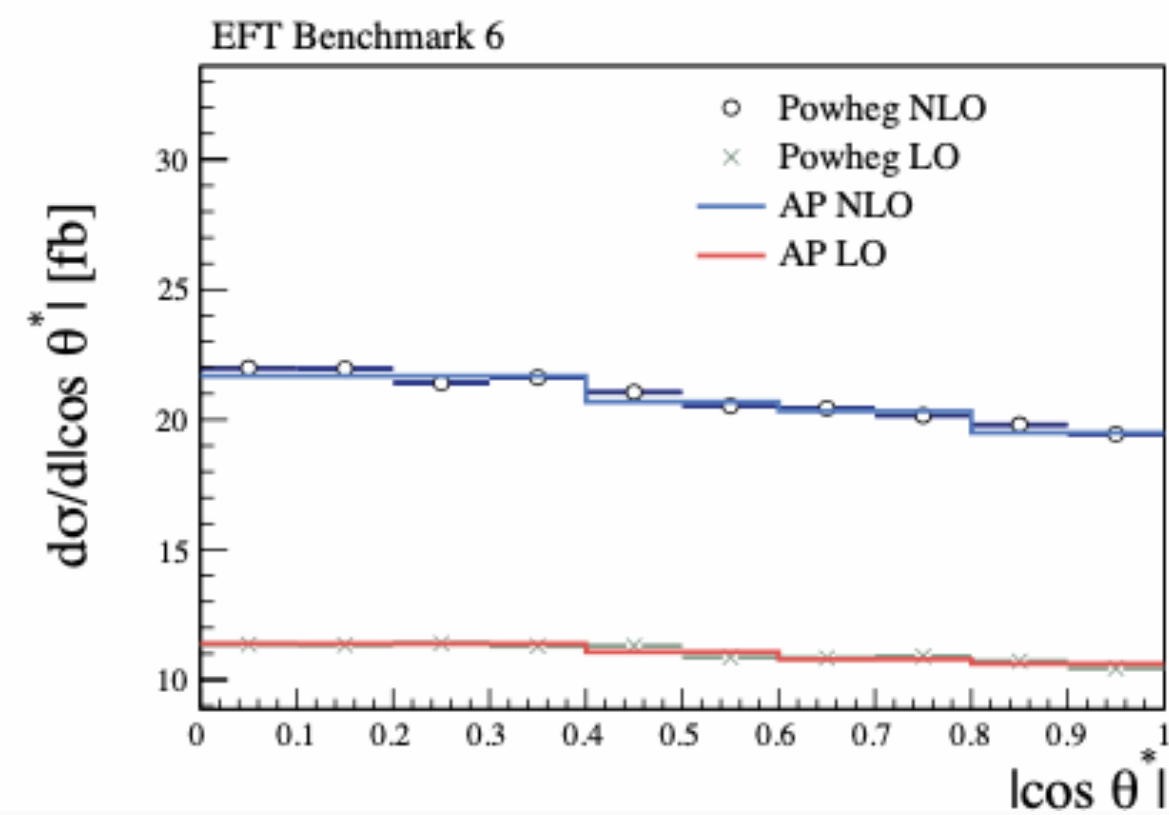
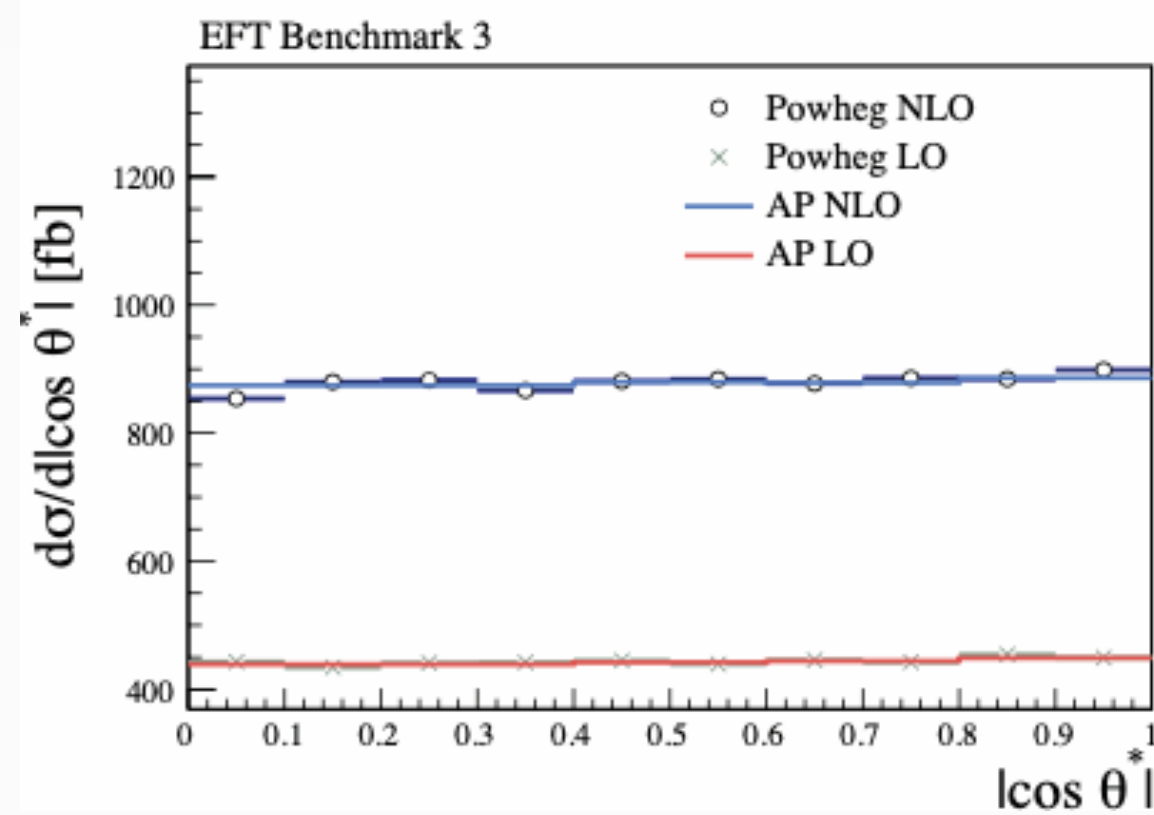
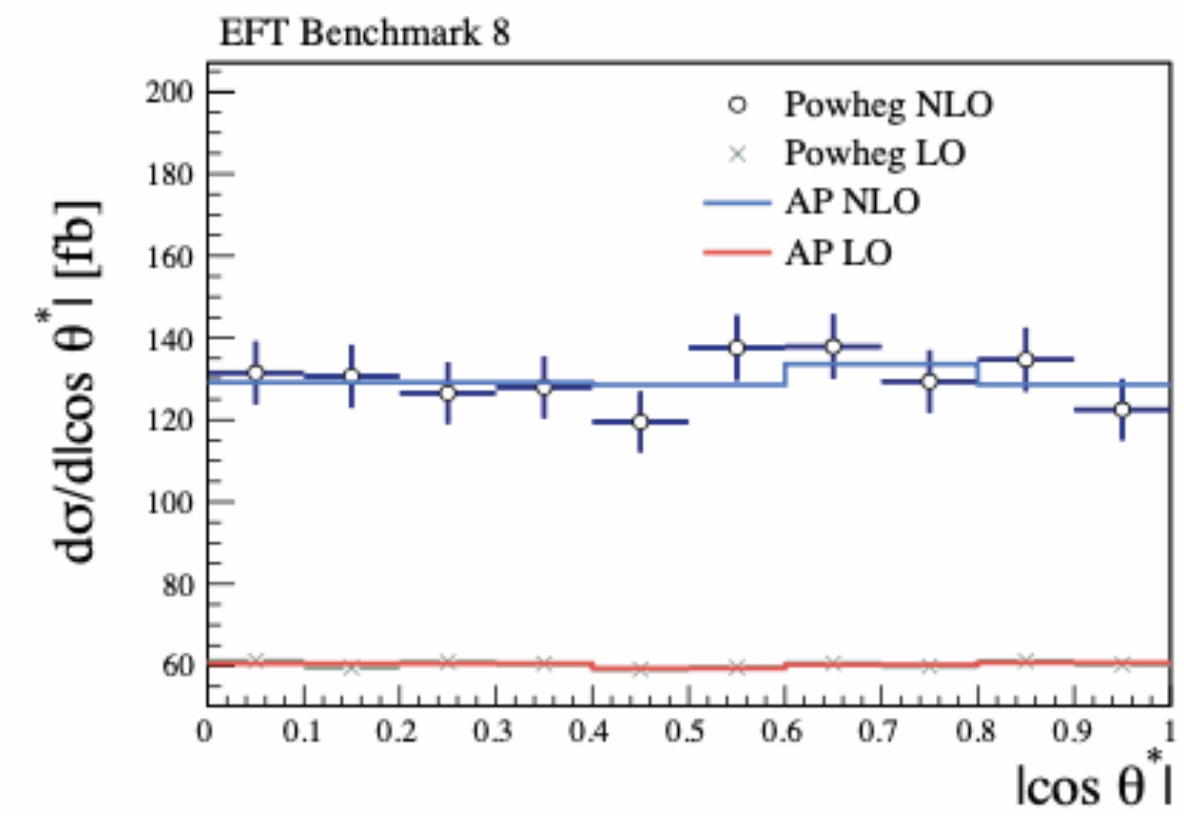
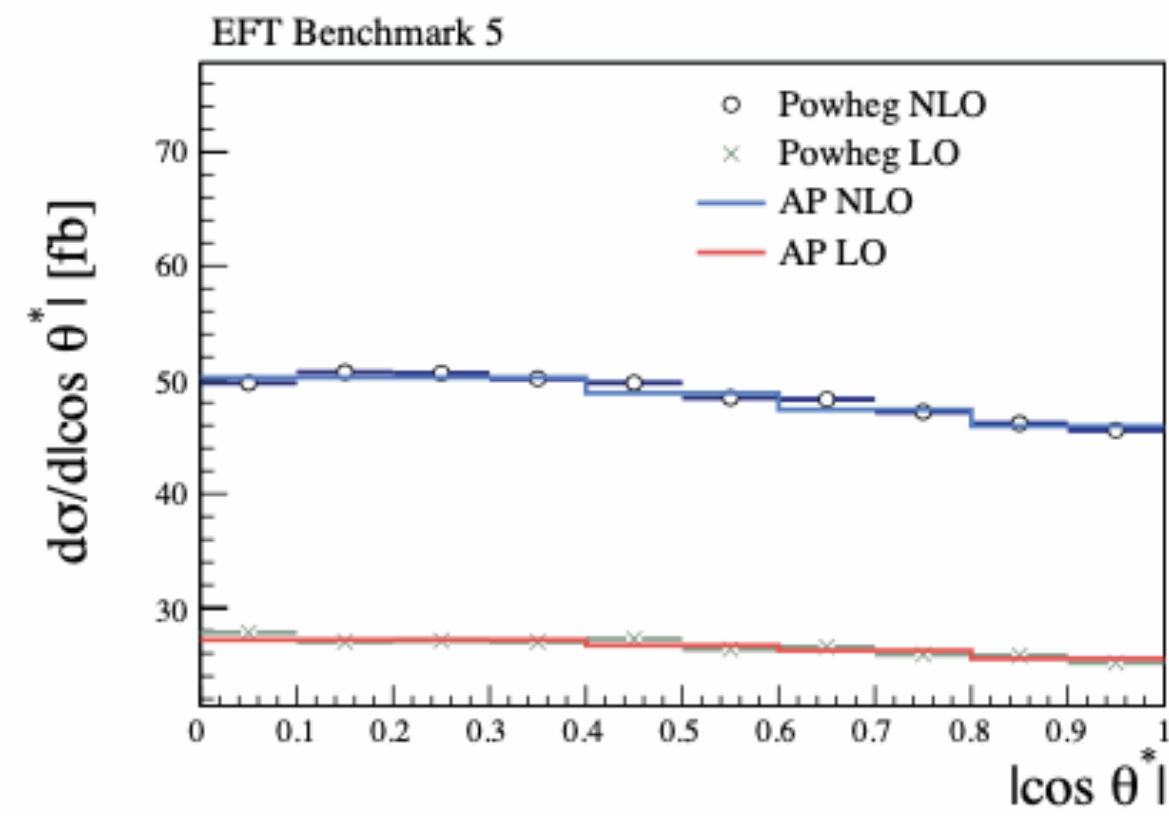
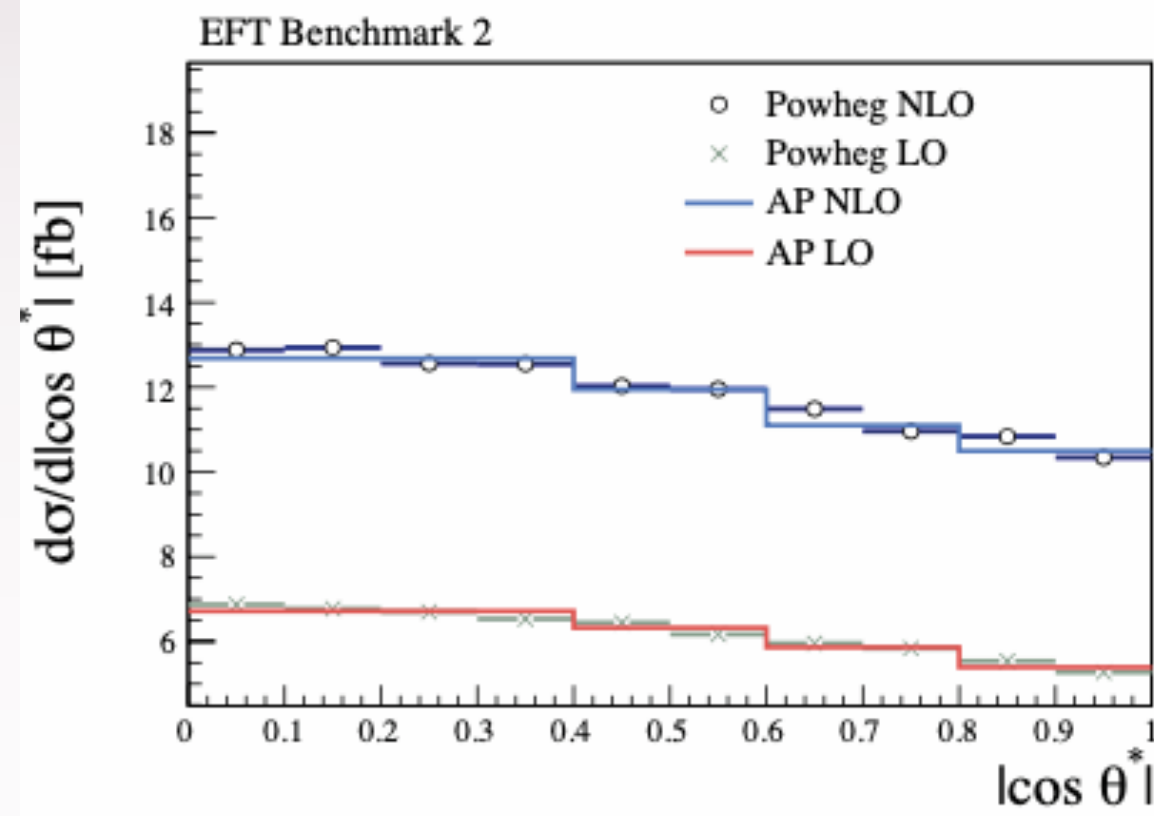
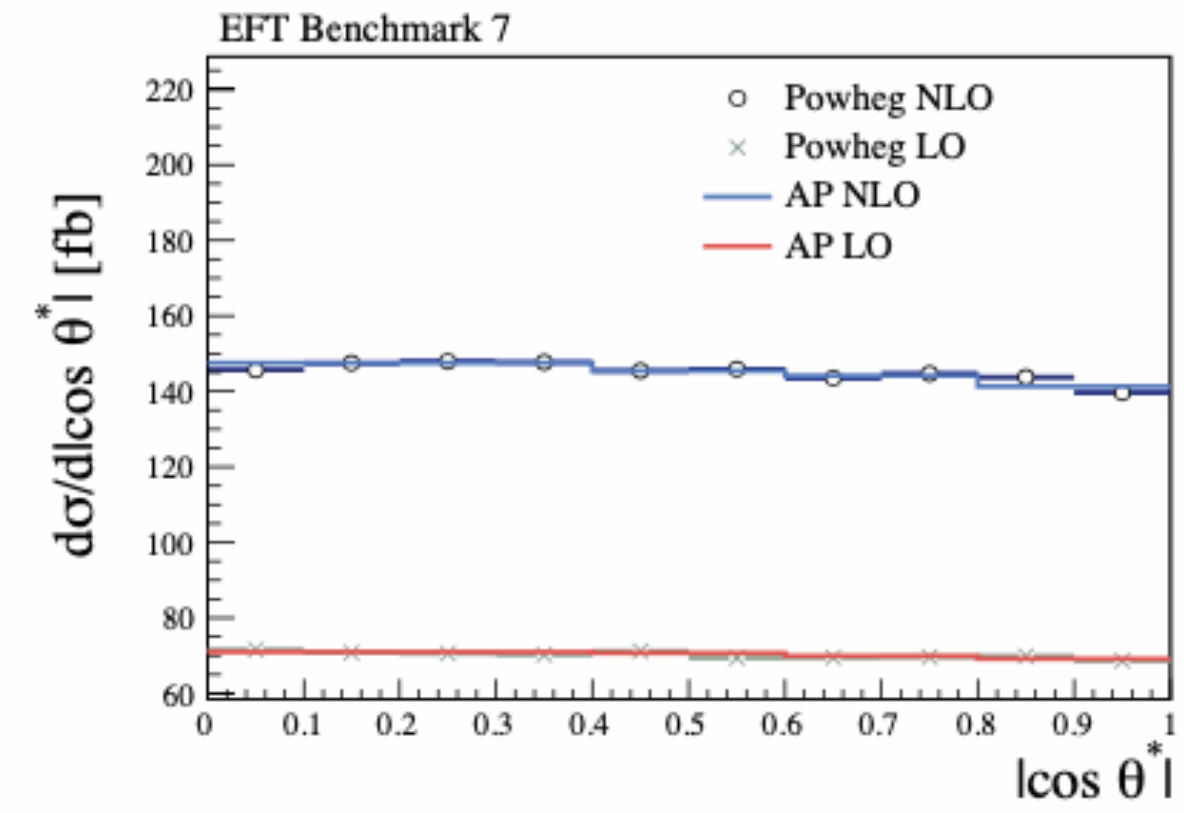
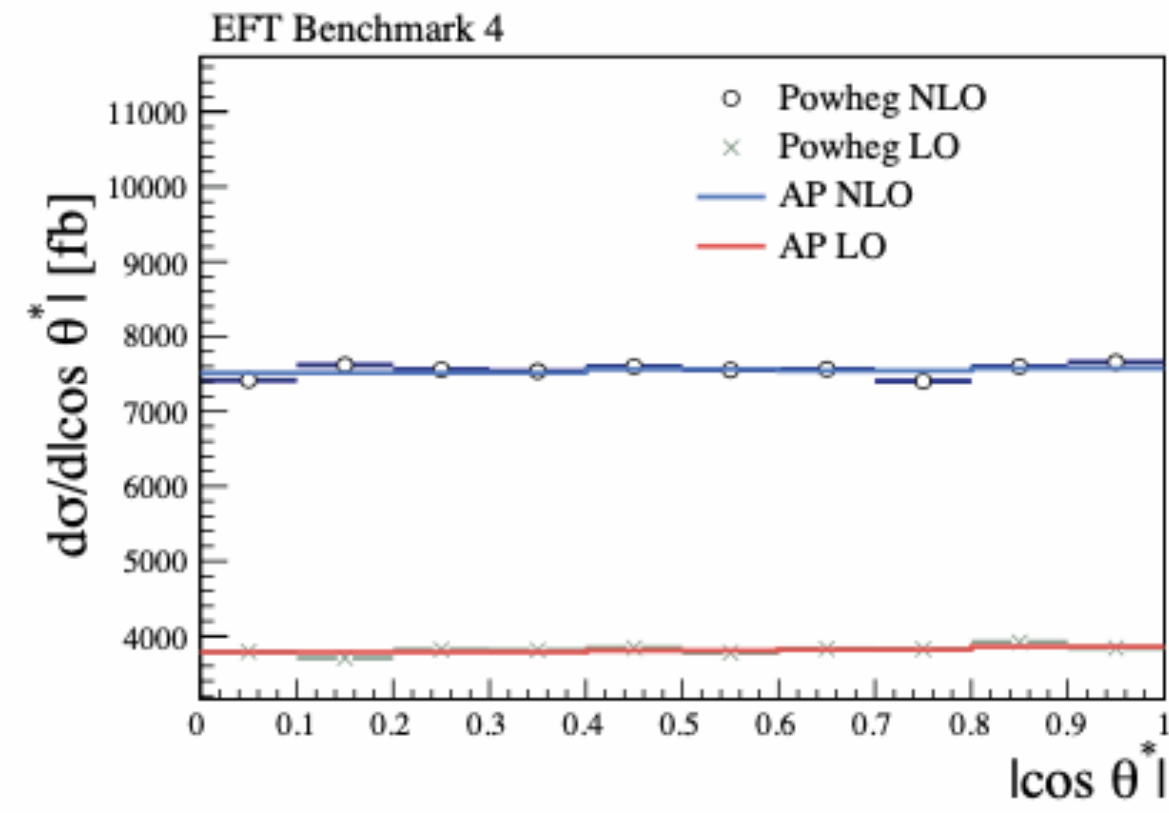
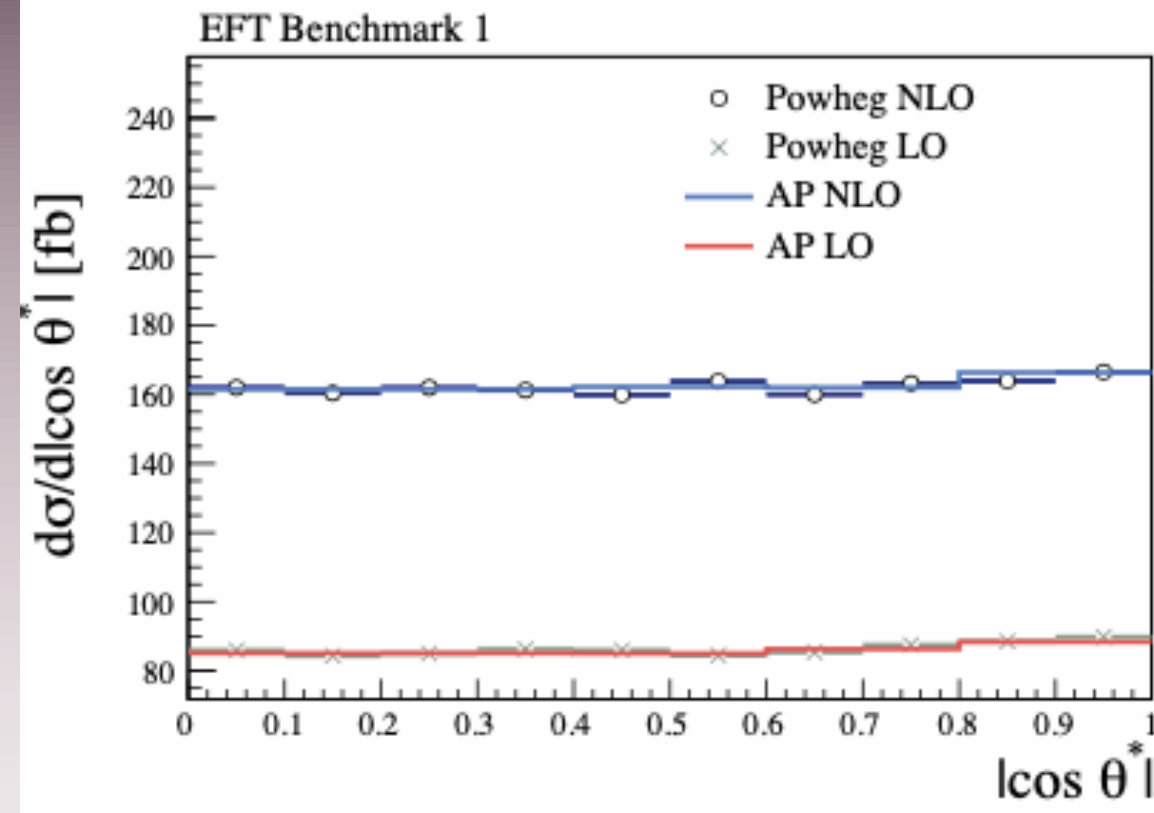
# Benchmark 8, a special one



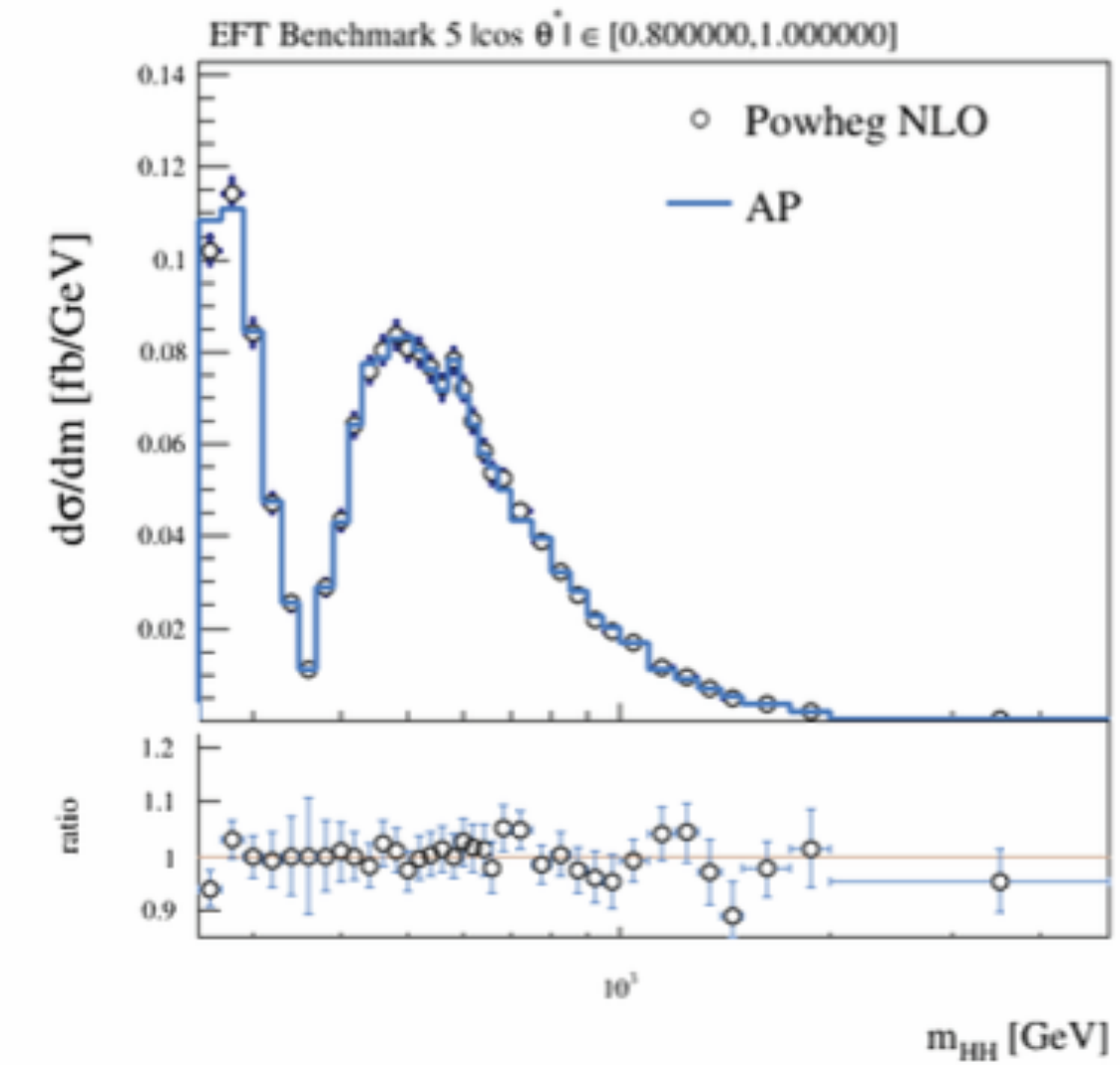
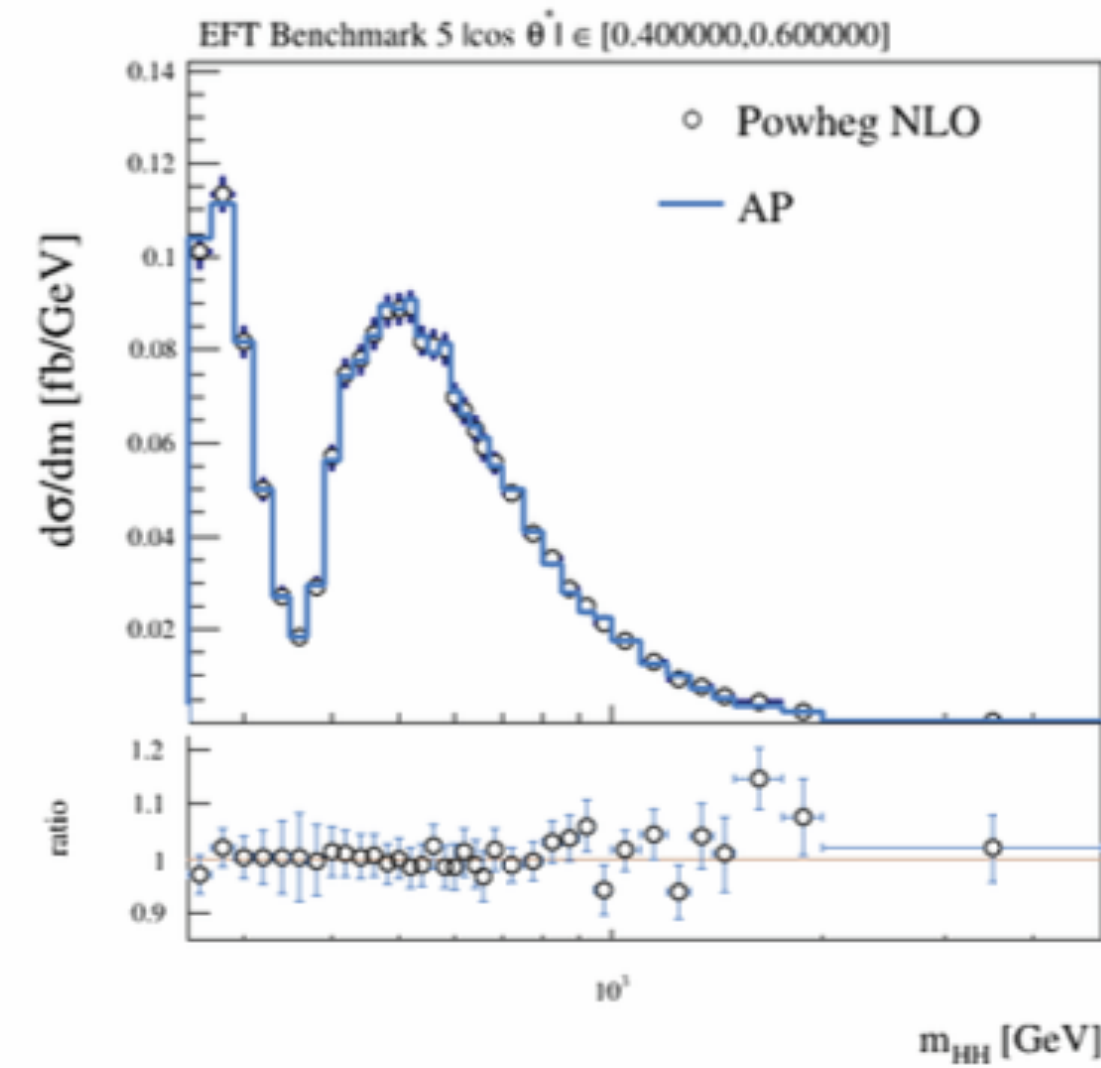
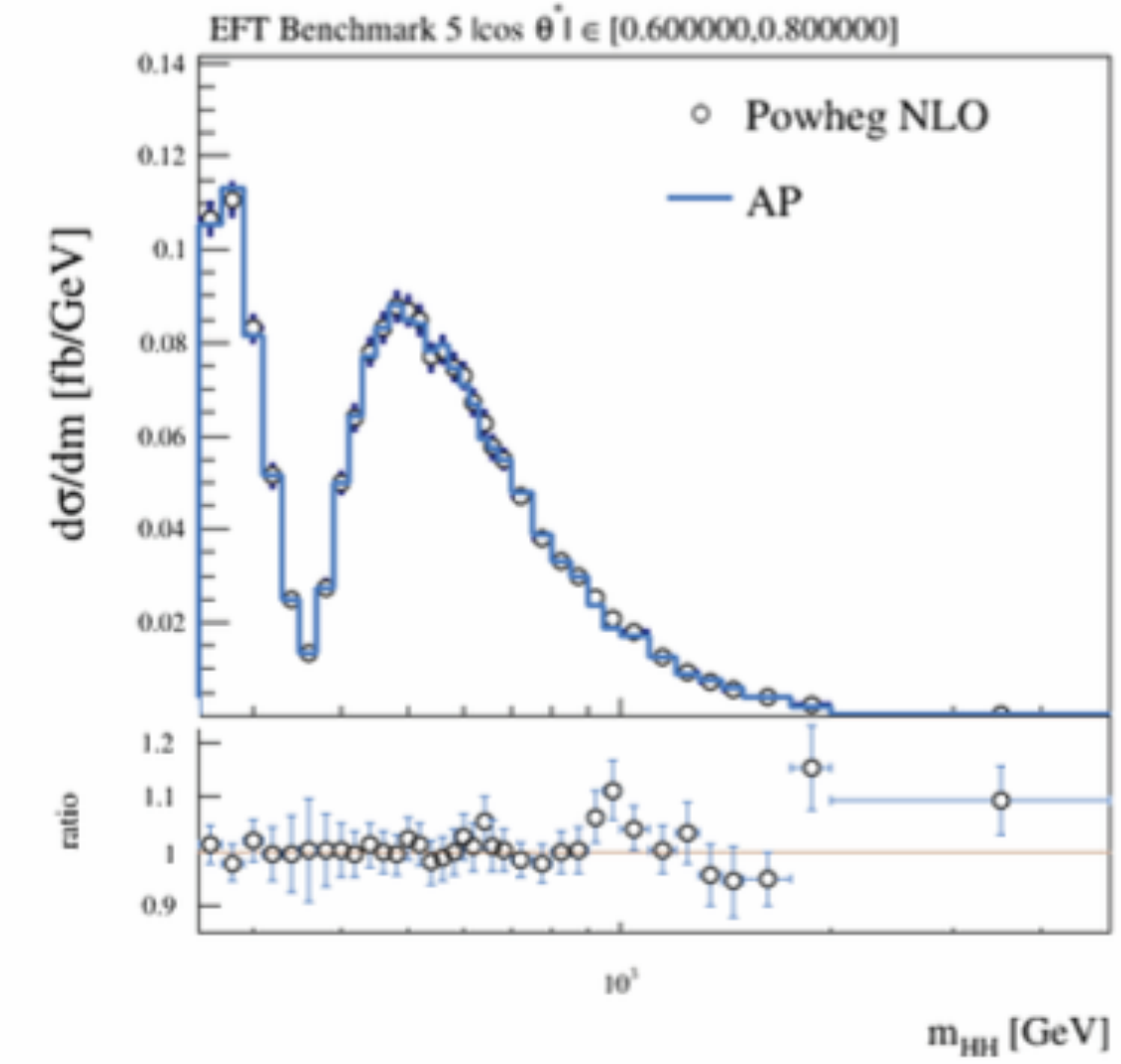
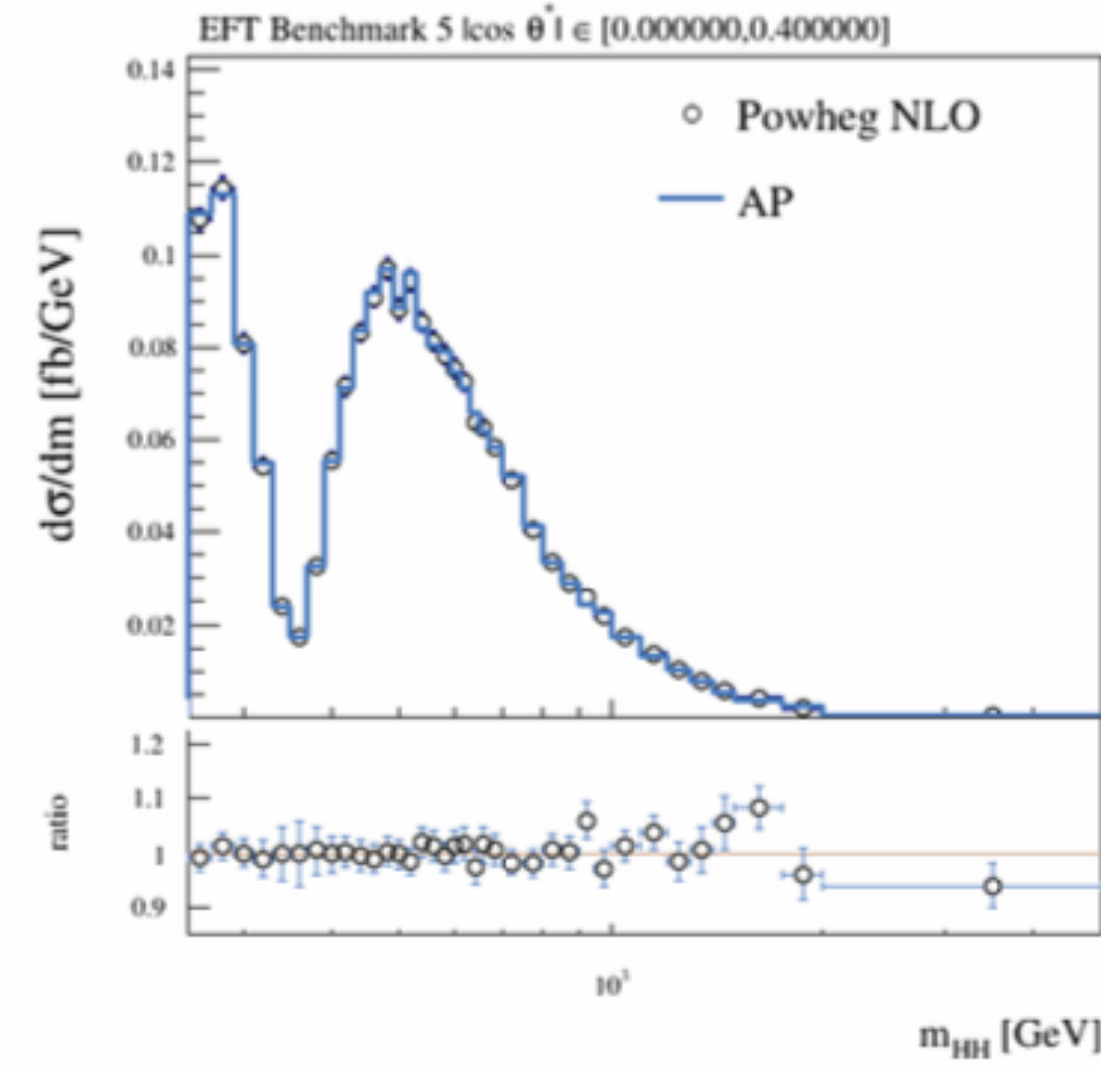
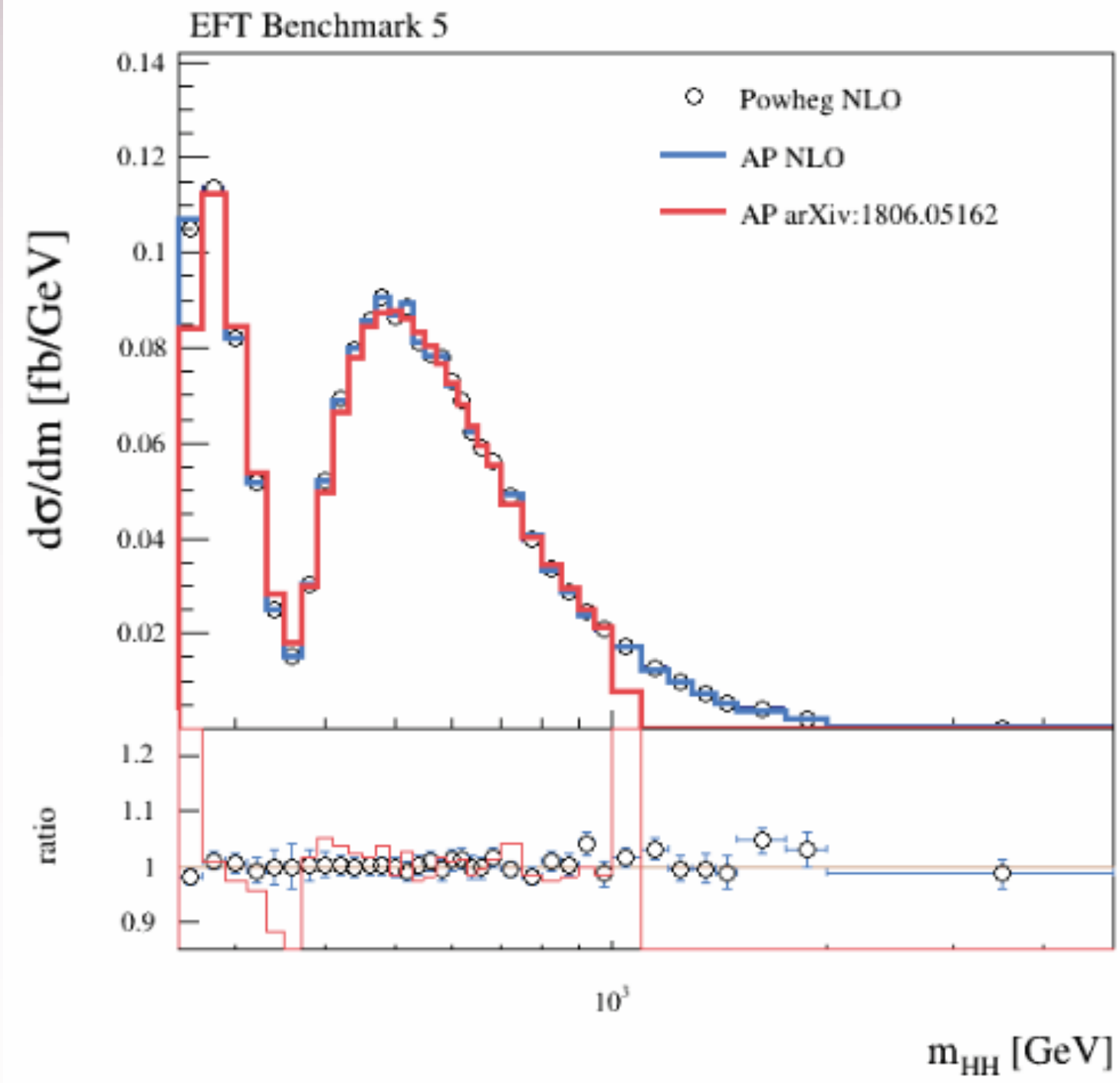
- In addition, generation of events for this benchmark at least  $\approx 10$  times slower



# NLO results



# NLO results



# Theory

by

L. Alasfar, R. Gröber and G. Heinrich

# Translation between the EFT's

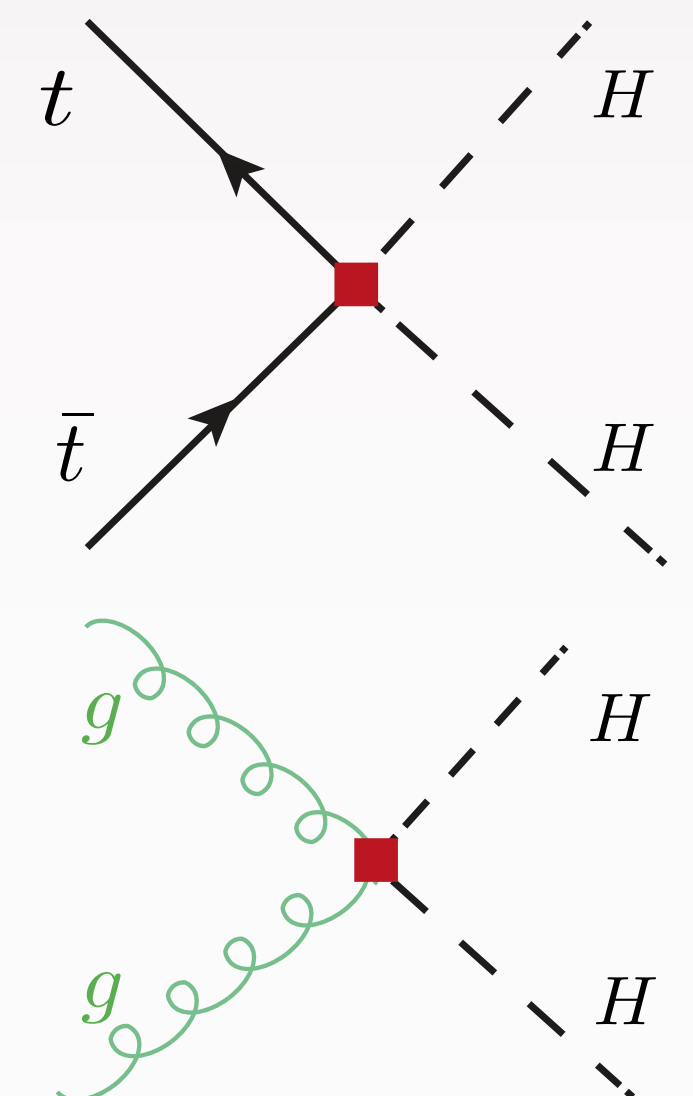
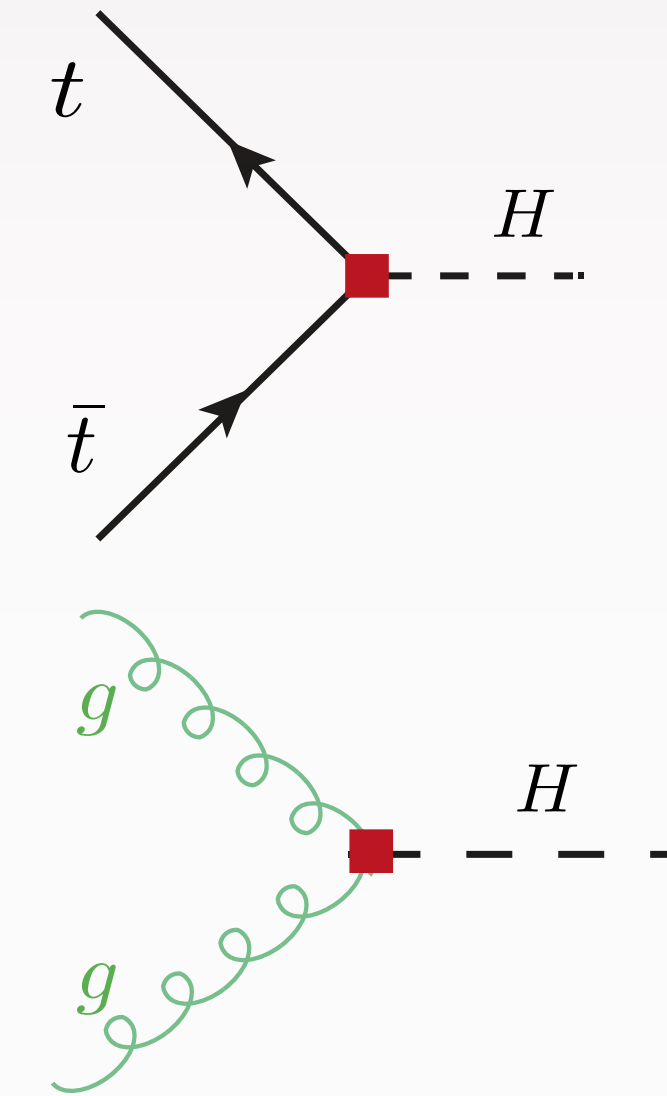
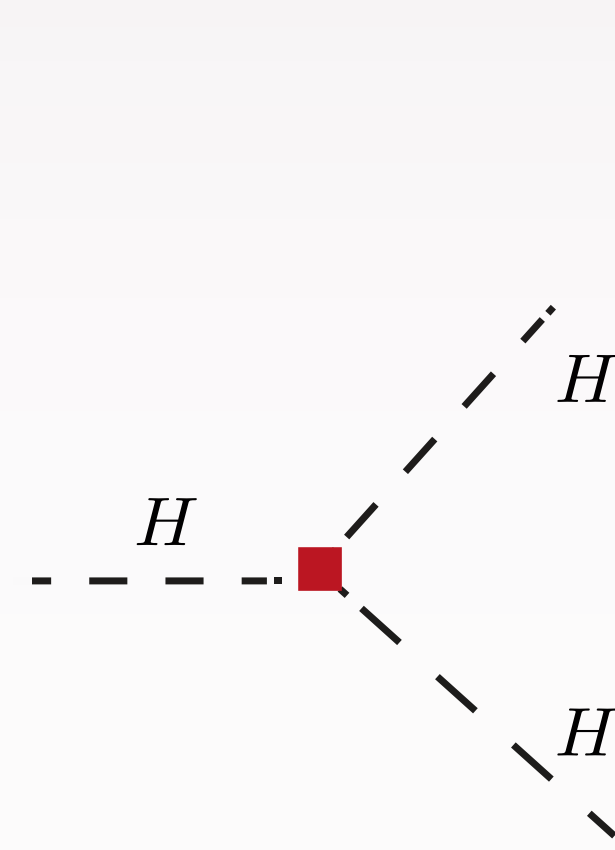
Equiv. use field redefinition

$$h \rightarrow h + C_{H,kin} \left( h + h^2/v + h^3/(3v^2) \right)$$

One can use the equations of motions for the Higgs to eliminate some of the Warsaw basis operators

We define the Wilson coefficient for Warsaw basis  $C_{H,kin} = \left( C_{H,\square} - \frac{1}{4} C_{HD} \right) v^2$

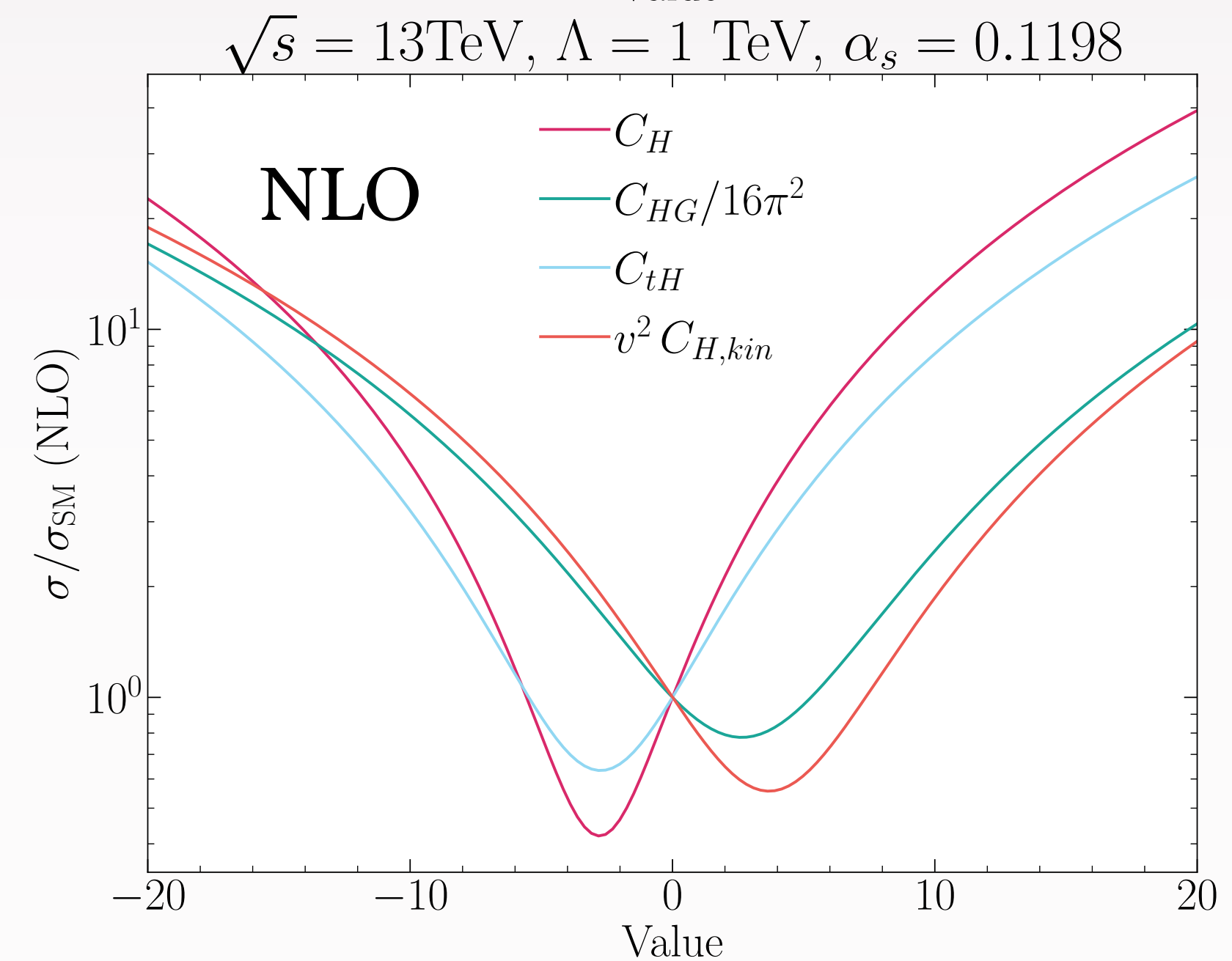
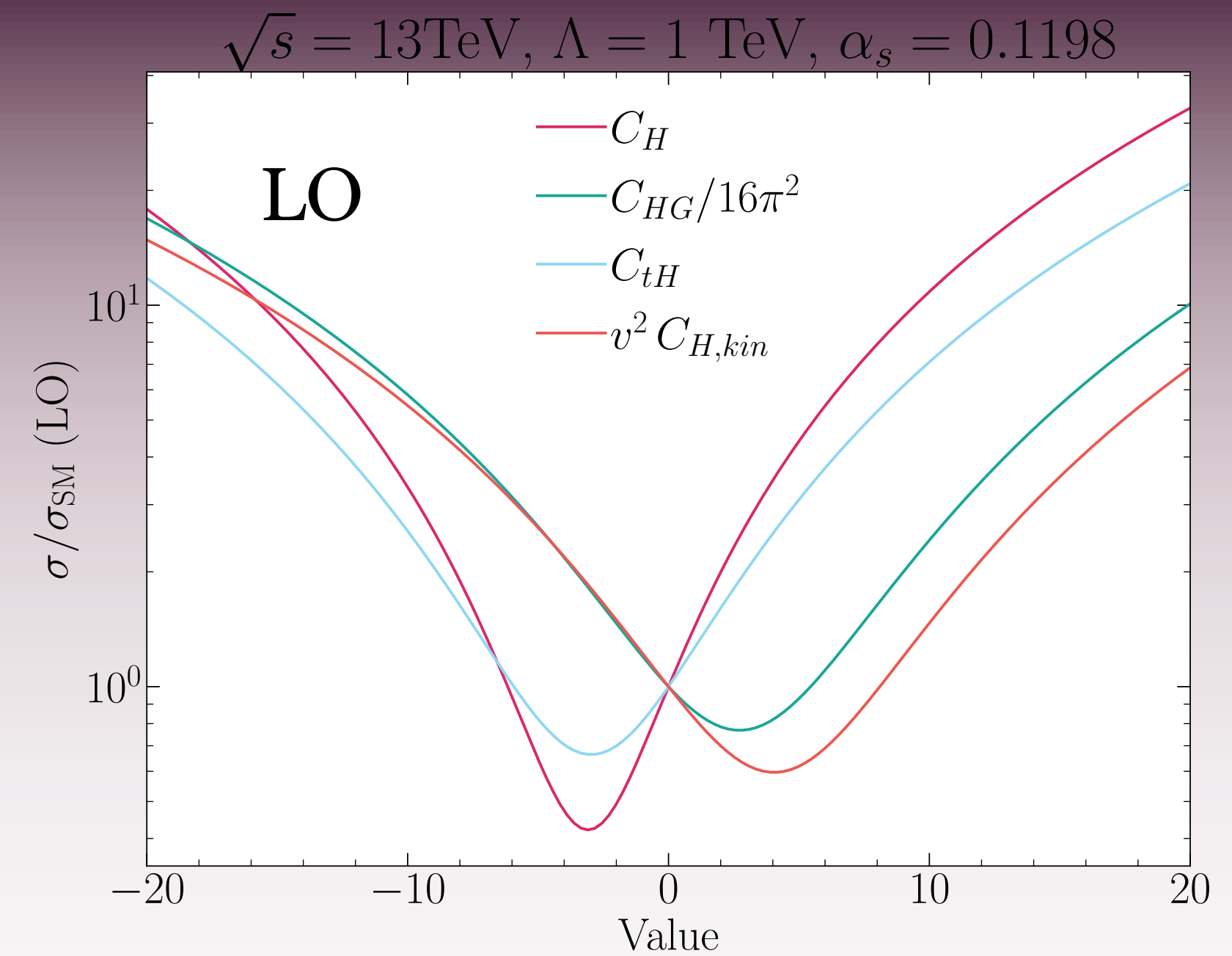
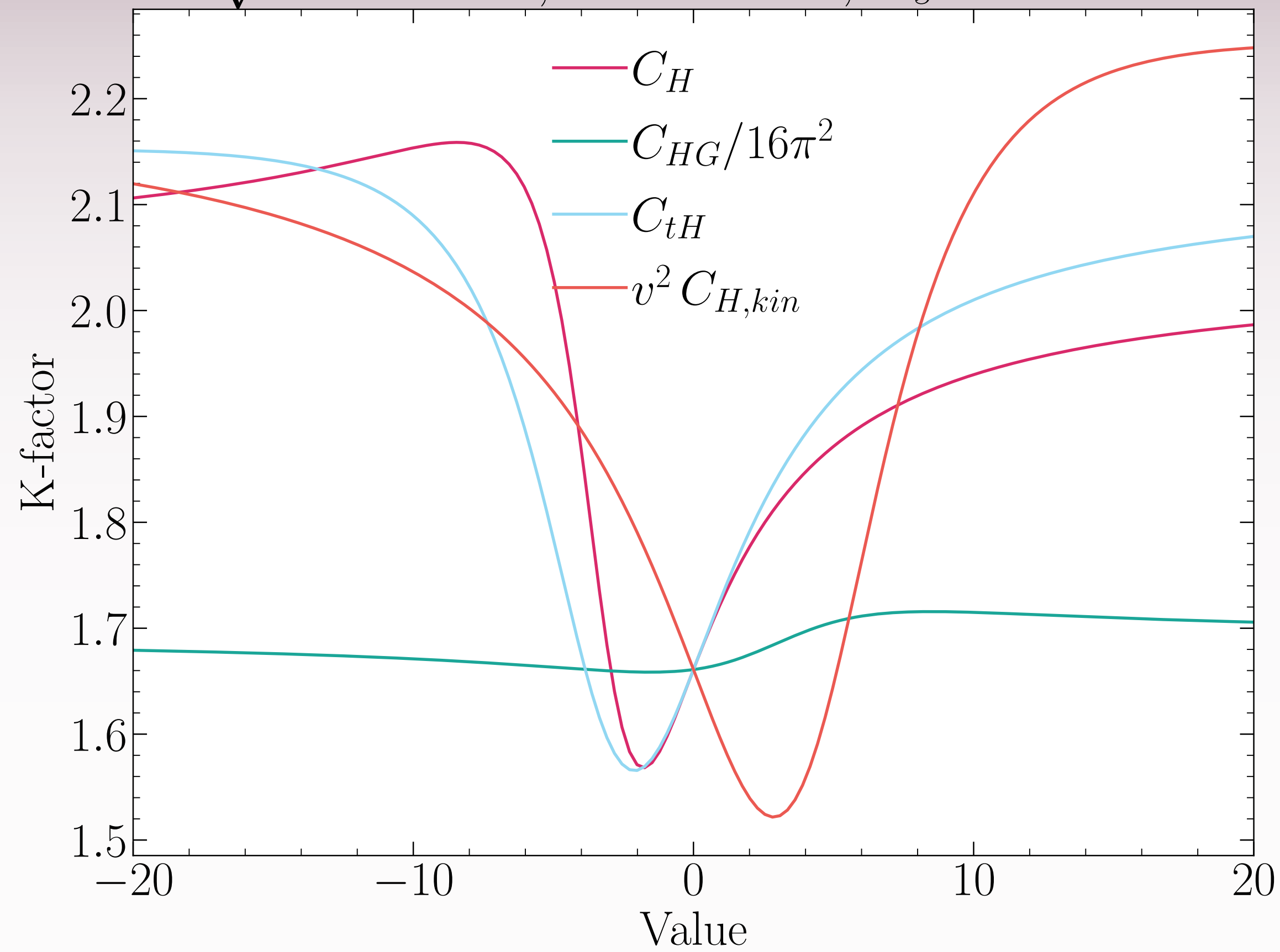
HEFT	SILH	Warsaw
$c_{hhh}$	$1 + \bar{c}_6 - \frac{3}{2} \bar{c}_H$	$1 - 2 \frac{v^4}{m_h^2} C_H + 3 C_{H,kin}$
$c_t$	$1 - \frac{\bar{c}_H}{2} - \bar{c}_u$	$1 + C_{H,kin} - C_{uH} \frac{v^3}{\sqrt{2} m_t}$
$c_{tt}$	$-\left( \frac{3}{2} \bar{c}_u + \frac{\bar{c}_H}{2} \right)$	$-C_{uH} \frac{3v^3}{2\sqrt{2} m_t} + C_{H,kin}$
$c_{ggh}$	$\frac{128\pi^2}{g_2^2} \bar{c}_g$	$\frac{8\pi}{\alpha_s} v^2 C_{HG}$
$c_{gghh}$	$\frac{64\pi^2}{g_2^2} \bar{c}_g$	$\frac{4\pi}{\alpha_s} v^2 C_{HG}$



# Cross-sections and K-factors

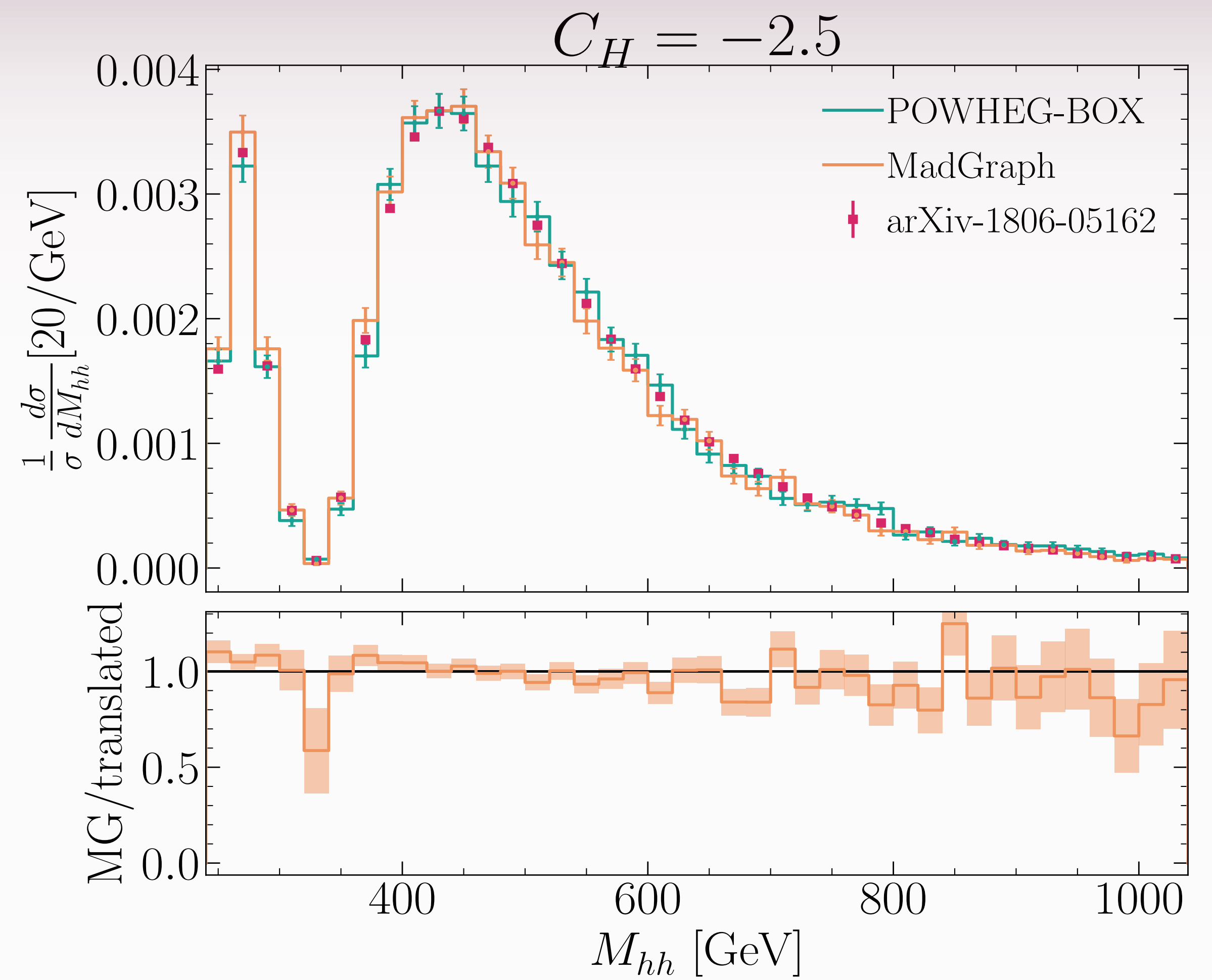
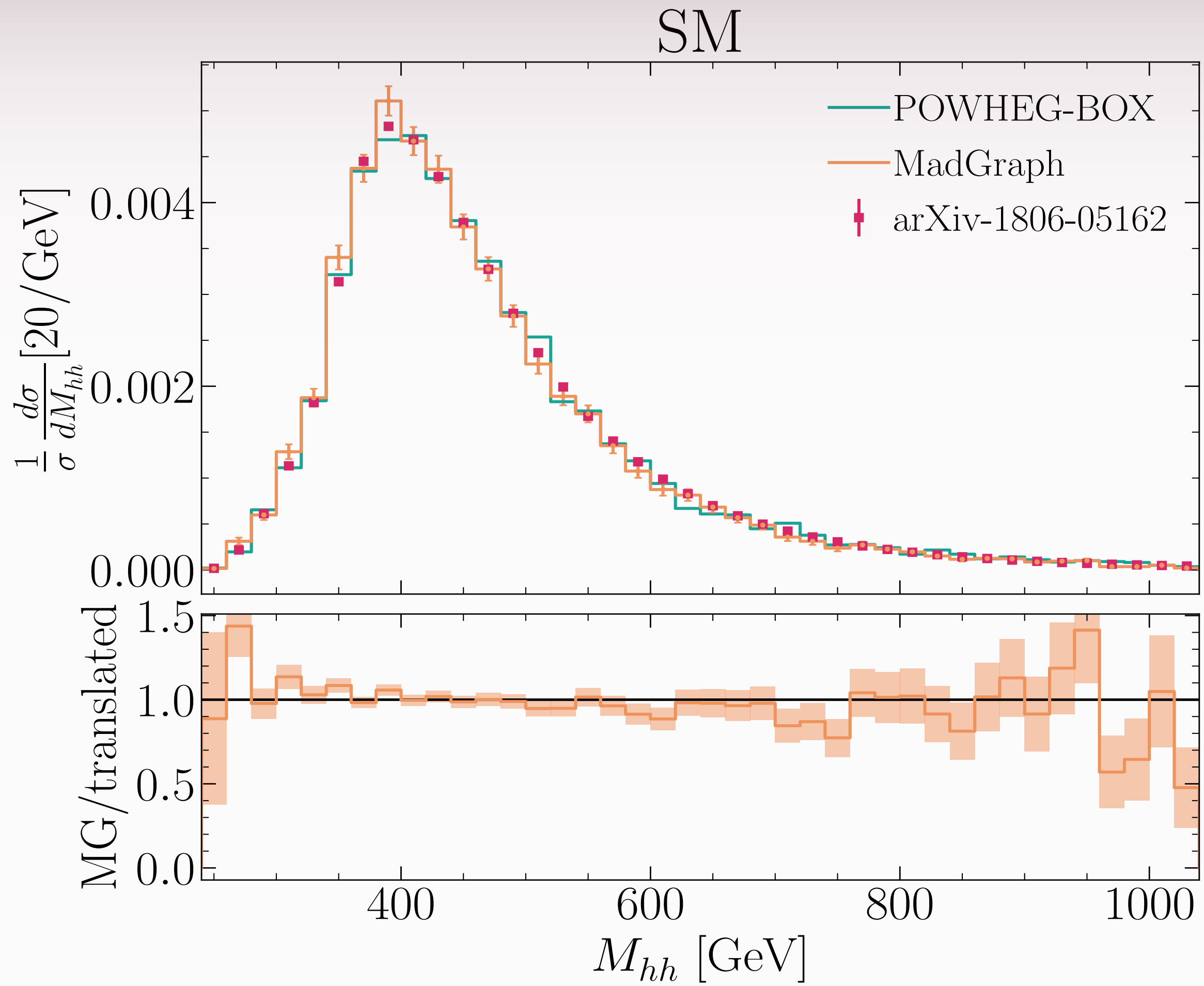
We redefine the SMEFT operators such that  $C \rightarrow \frac{C}{\Lambda^2}$

$\sqrt{s} = 13\text{TeV}, \Lambda = 1\text{ TeV}, \alpha_s = 0.1198$

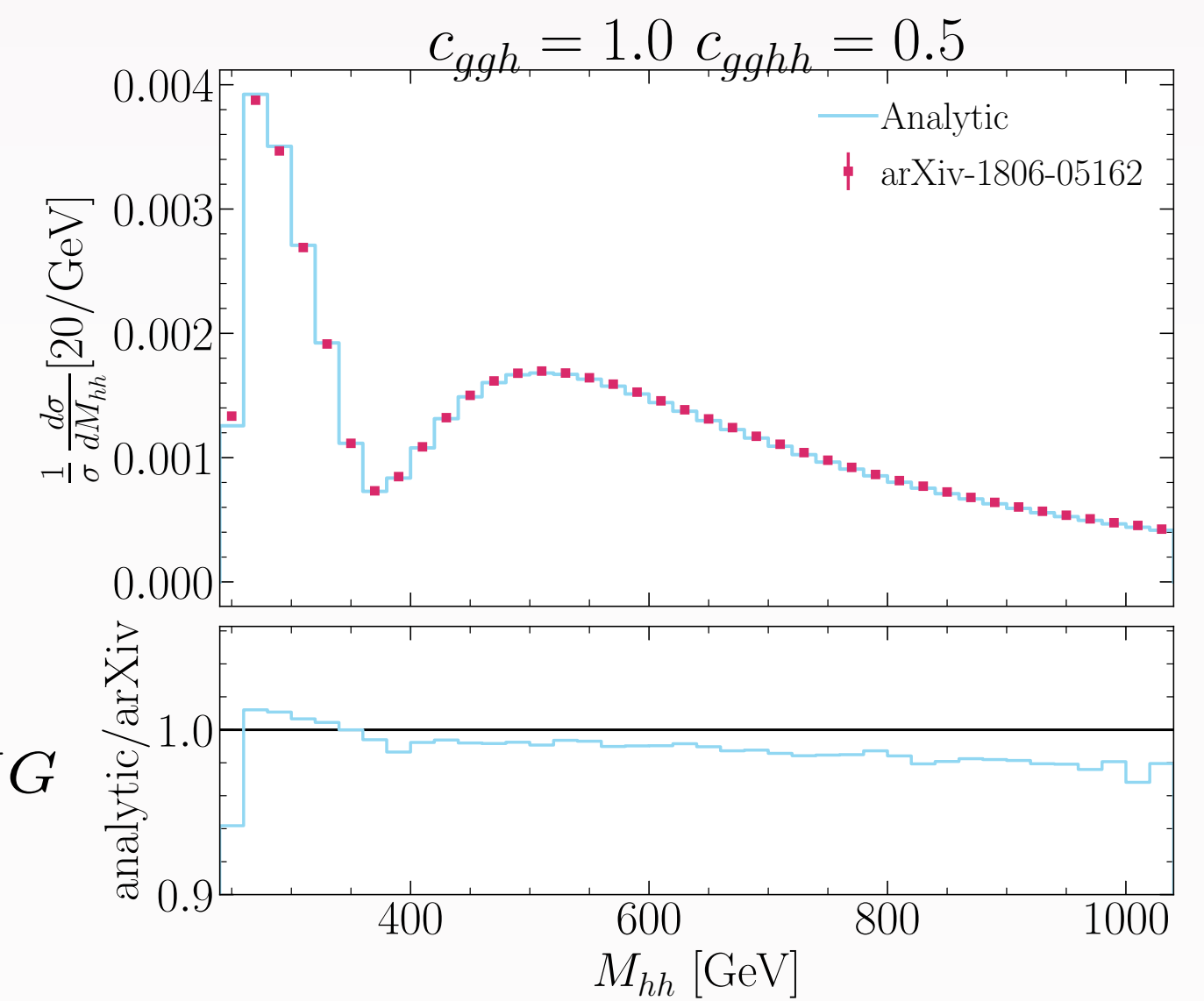
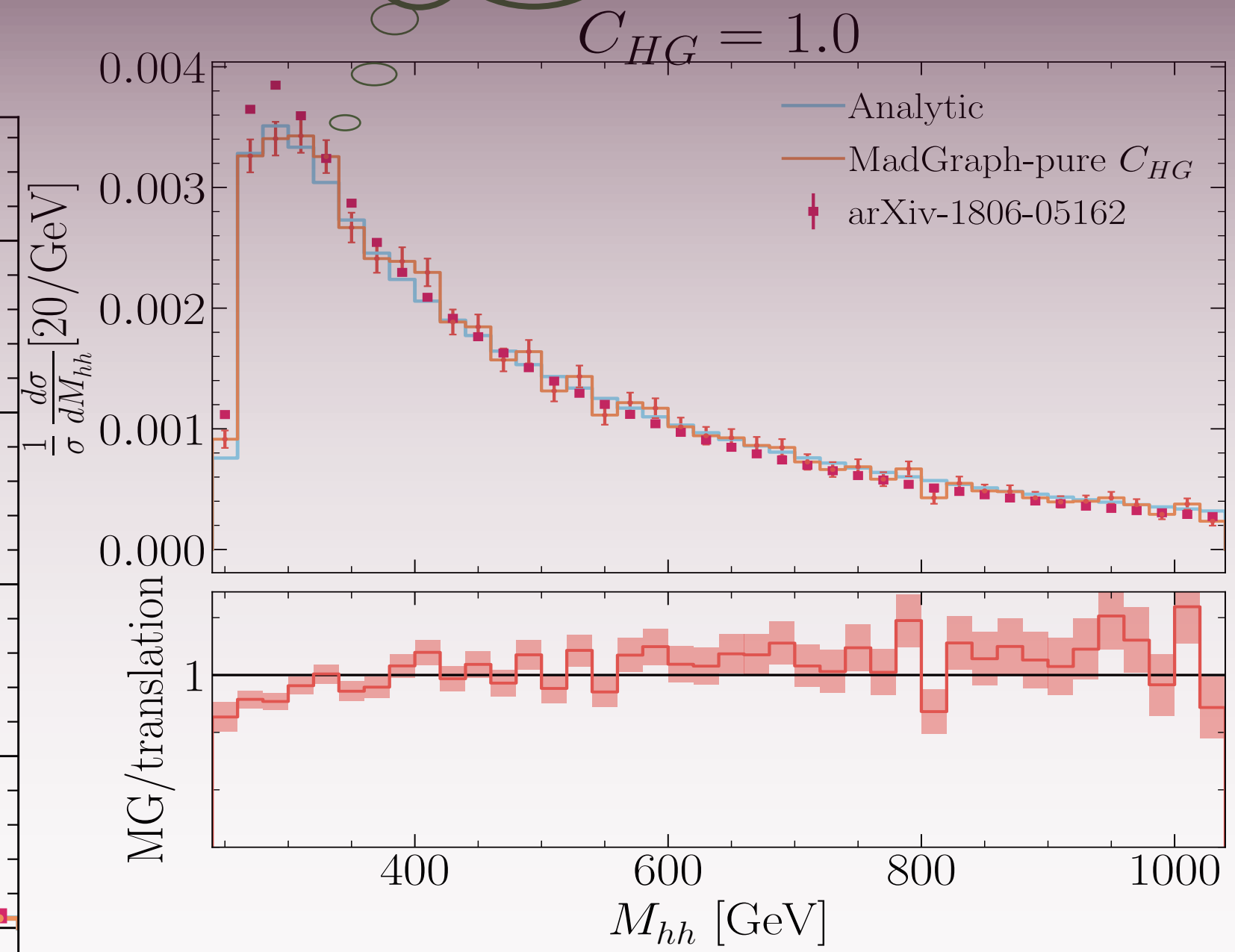
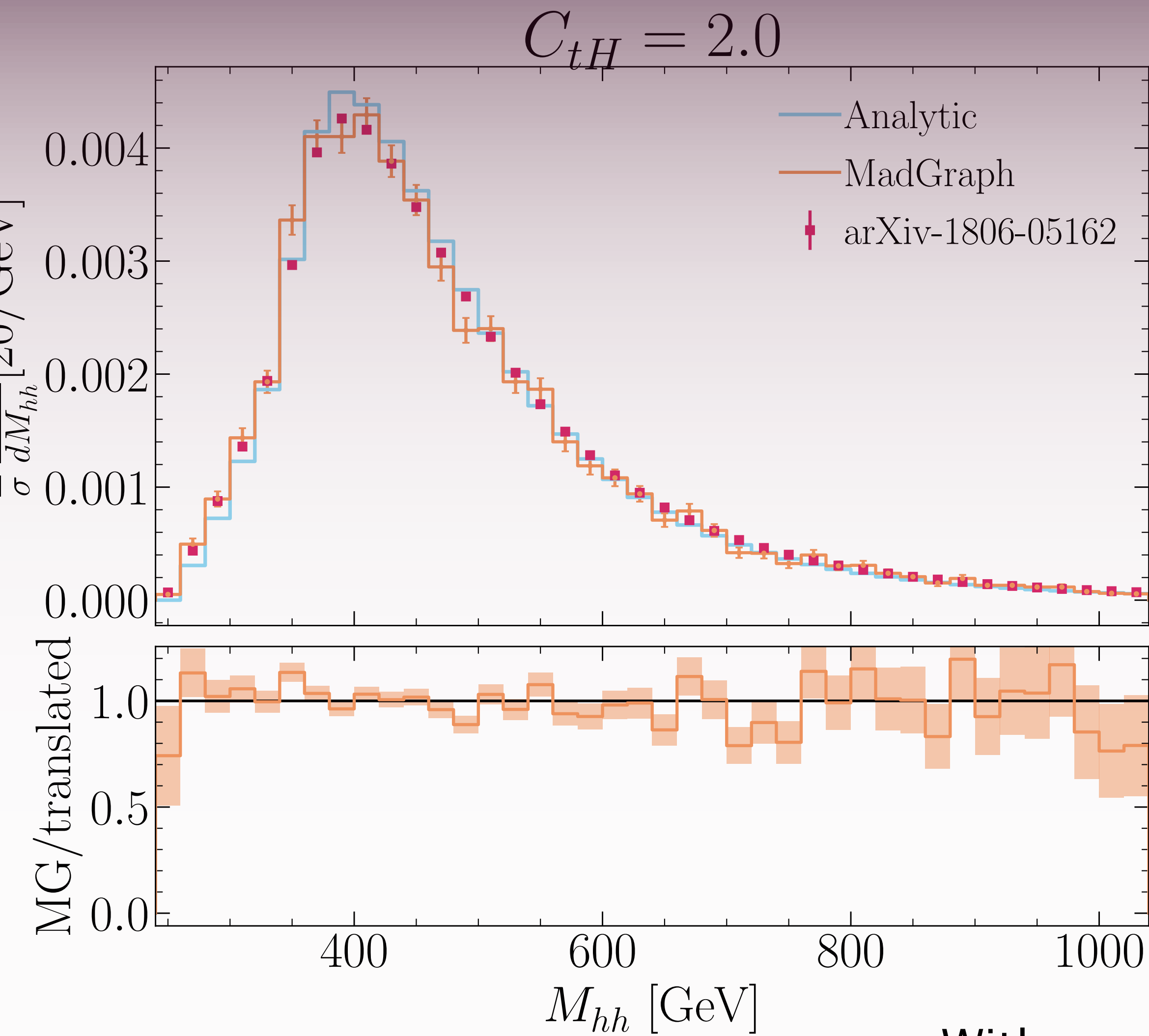


These plots are obtained from the translation from HEFT  
 from [G. Buchalla, et al. '18] arXiv:1806.05162 (hep-ph)

# Validation plots (LO only)



# Validation plots (LO only)



With modified definition of  $C_{HG}$

# Conclusion

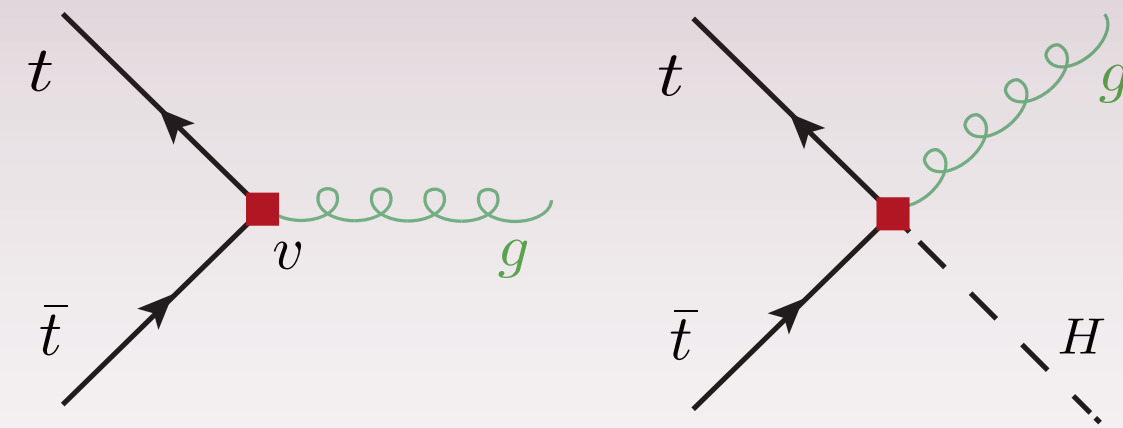
- Reweighting MC samples (LO and NLO) was achieved for HEFT.
- Low statistics in the 1st bin is noticed to affect some benchmarks.  
Moreover, some benchmarks show disagreement between MG and Powheg derived AP's.
- Custom coefficients of the AP of the differential NLO and LO Powheg cross sections as a function of  $M_{hh}$  and  $\cos \theta^*$  are derived, allowing to increase reweighting range from  $M_{hh} < 1040$  up to  $M_{hh} < 5000$  GeV and take into account  $\cos \theta^*$  dependence.
  - Predictions for the SM and 12 EFT benchmarks are in a good agreement with MC  $\implies$  suitable for the analysis of 12 benchmarks
  - Constraints on NLO  $A_{16}$ - $A_{23}$  coefficients are limited by statistical fluctuations  $\implies$  EFT parameter scan with care
- The issue with HEFT to Warsaw for  $C_{HG}$ , where the translation depends on  $\alpha_s$  can be avoided by redefining the SMEFT operator.



# Outlook

- To include or not to include the chromomagnetic operator, in SMEFT

$$\hat{O}_{tG} := y_t g_s (\bar{q}_L \sigma^{\mu\nu} T_A t_R) H^c G_{\mu\nu}^A$$



- NLO SMEFT event generator, via POWHEG-BOX
- Better statistics for the 1st bin and NLO coefficients.

Backup

# Reminder

For the following EFT Lagrangian:

$$\mathcal{L}_{ggF} = -\kappa_\lambda \lambda_{HHH}^{SM} v H^3 - \frac{m_t}{v} (\kappa_t H + \frac{\kappa_2}{v} H^2) (\bar{t}_L t_R + h.c.) + \frac{\alpha_S}{12\pi v} (\kappa_g H - \frac{\kappa_{2g}}{2v} H^2) G_{\mu\nu}^a G^{a,\mu\nu}$$

the cross section for ggF HH production in  $pp$  can be parameterized as a polynomial of the  $\kappa$ 's:

$$\begin{aligned}\sigma_{LO} = & A_1 \kappa_t^4 + A_2 \kappa_2^2 + A_3 \kappa_t^2 \kappa_\lambda^2 + A_4 \kappa_g^2 \kappa_\lambda^2 + A_5 \kappa_{2g}^2 + A_6 \kappa_2 \kappa_t^2 + A_7 \kappa_t^3 \kappa_\lambda \\ & + A_8 \kappa_2 \kappa_t \kappa_\lambda + A_9 \kappa_2 \kappa_g \kappa_\lambda + A_{10} \kappa_2 \kappa_{2g} + A_{11} \kappa_t^2 \kappa_g \kappa_\lambda + A_{12} \kappa_t^2 \kappa_{2g} \\ & + A_{13} \kappa_t \kappa_\lambda^2 \kappa_g + A_{14} \kappa_t \kappa_\lambda \kappa_{2g} + A_{15} \kappa_g \kappa_\lambda \kappa_{2g} .\end{aligned}$$

$$\begin{aligned}\sigma_{NLO} = & \dots + A_{16} \kappa_t^3 \kappa_g + A_{17} \kappa_t \kappa_2 \kappa_g + A_{18} \kappa_t \kappa_g^2 \kappa_\lambda + A_{19} \kappa_t \kappa_g \kappa_{2g} \\ & + A_{20} \kappa_t^2 \kappa_g^2 + A_{21} \kappa_2 \kappa_g^2 + A_{22} \kappa_g^3 \kappa_\lambda + A_{23} \kappa_g^2 \kappa_{2g} .\end{aligned}$$

to take into account differences in the shapes a binned parametrisation of the cross section, for example, as a function of  $m_{HH}$  can be used:

$$\frac{d\sigma}{dm_{HH}} = \sum A_i(m_{HH}) c_i$$

To extract the  $A_i(m_{HH}, |\cos \theta^*|)$  coefficients the system of equations is solved by performing a minimisation in ROOT MINUIT. In every bin  $\mathbf{b} = (m_{HH}, |\cos \theta^*|)$  :

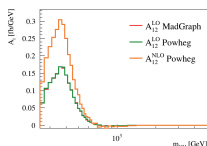
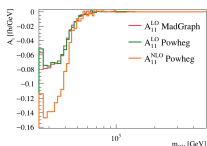
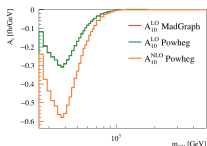
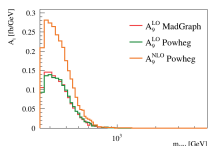
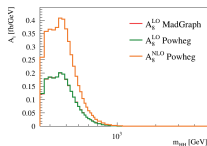
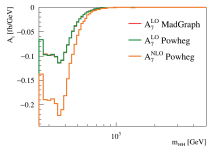
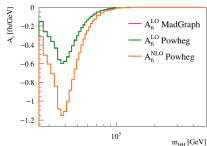
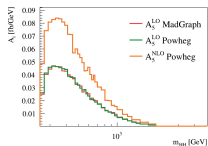
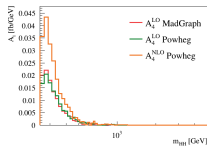
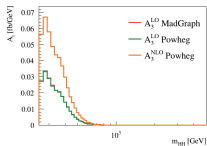
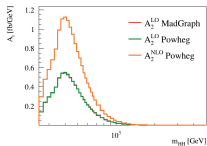
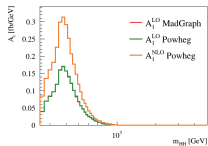
Edges in  $m_{HH} = [ 250, 270, 290, 310, 330, 350, 370,$   
 $390, 410, 430, 450, 470, 490, 510,$   
 $530, 550, 570, 590, 610, 630, 650,$   
 $670, 700, 750, 800, 850, 900, 950,$   
 $1000, 1100, 1200, 1300, 1400, 1500, 1750,$   
 $2000, 5000 ] \text{ Gev}$

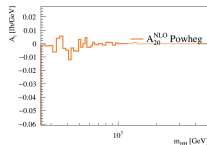
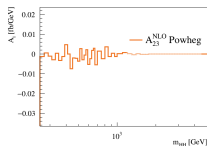
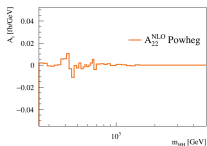
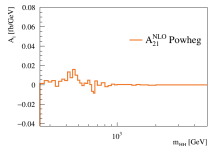
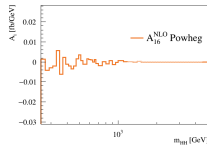
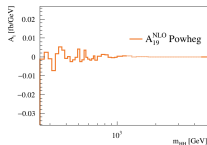
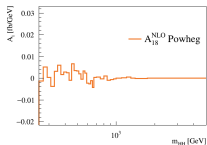
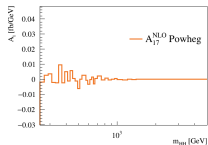
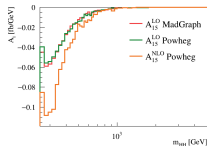
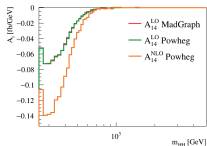
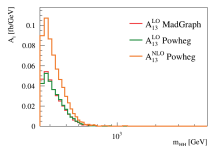
Edges in  $|\cos \theta^*| = [ 0, 0.4, 0.6, 0.8, 1 ]$

following function is minimized :

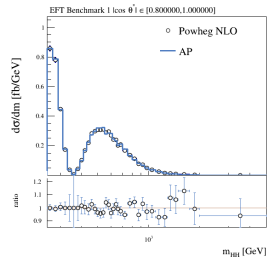
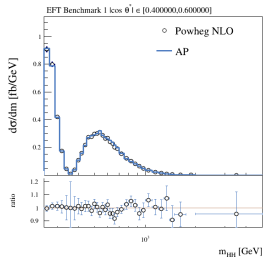
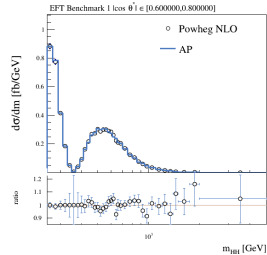
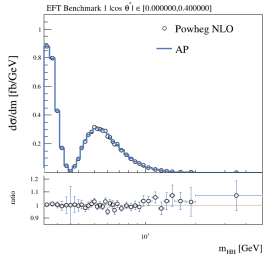
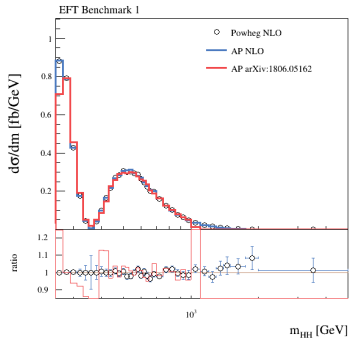
$$f(\mathbf{b}) = \sum_i^{N_{samples}} \left( \frac{\frac{d\sigma_i^{MC}(\mathbf{b})}{d\mathbf{b}} - \frac{d\sigma_i^{AP}(\mathbf{b})}{d\mathbf{b}}}{0.1 \cdot \frac{d\sigma_i^{MC}(\mathbf{b})}{d\mathbf{b}}} \right)^2$$

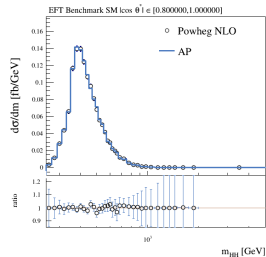
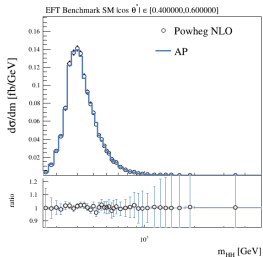
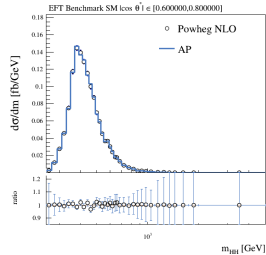
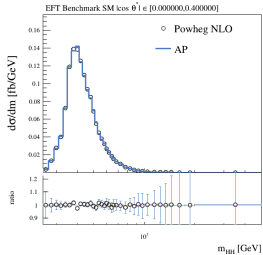
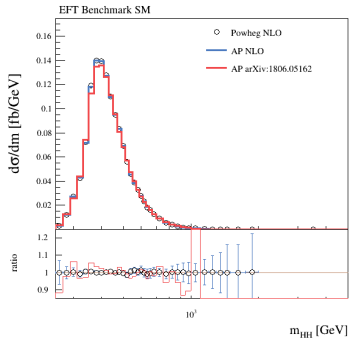
# $A_i$ from the fit, explicitly



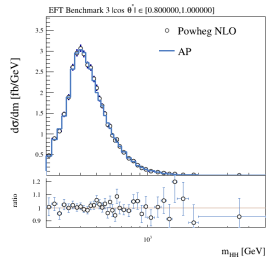
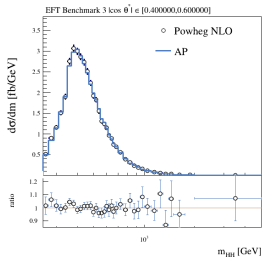
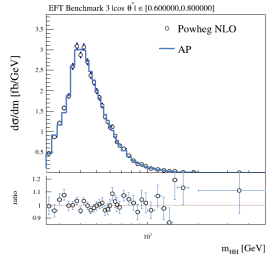
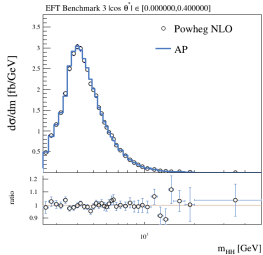
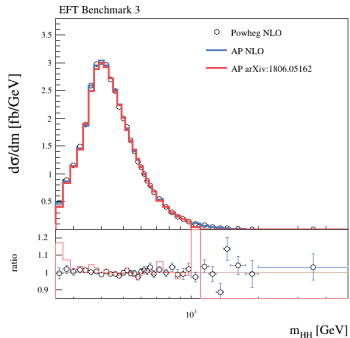


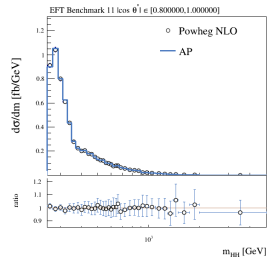
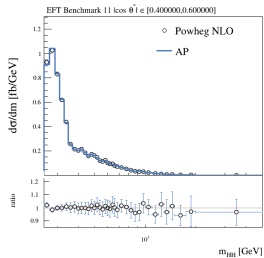
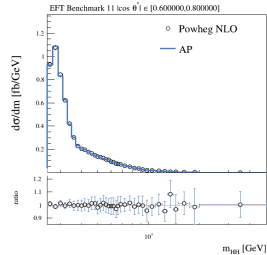
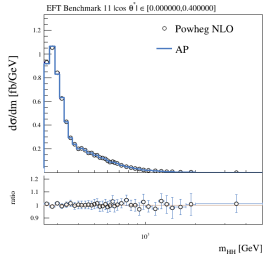
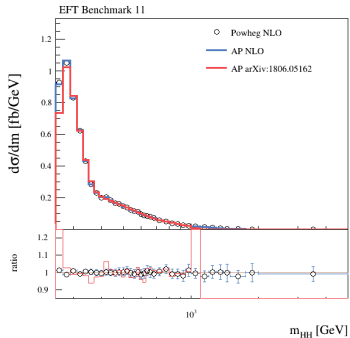
Constraints on NLO  $A_{16}$ - $A_{23}$  coefficients are limited by statistical fluctuations!

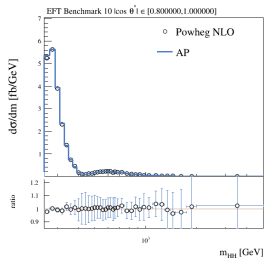
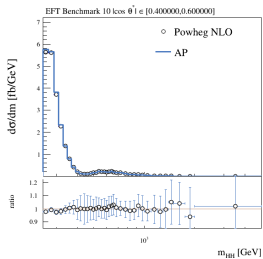
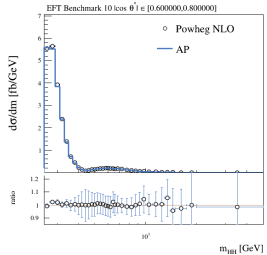
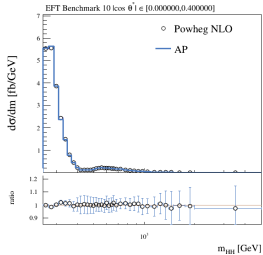
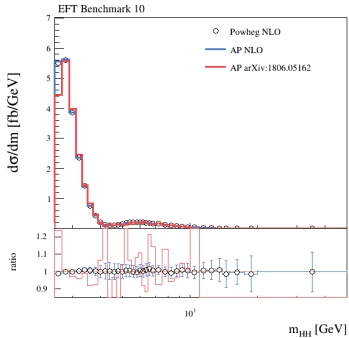


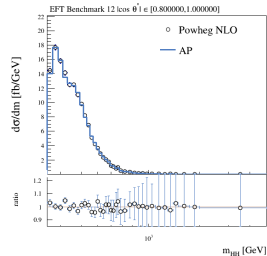
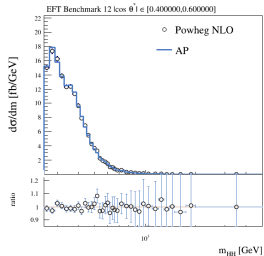
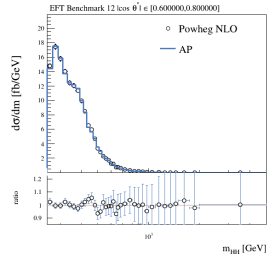
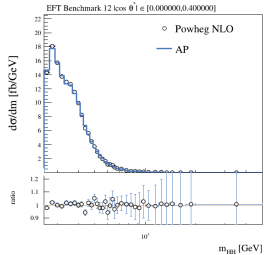
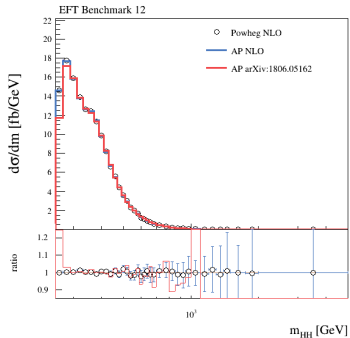


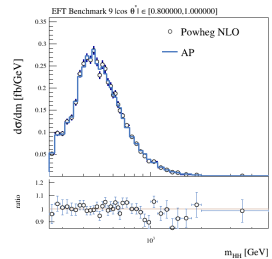
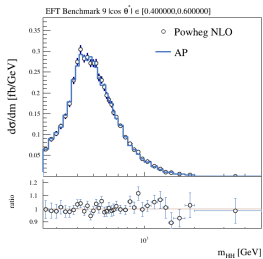
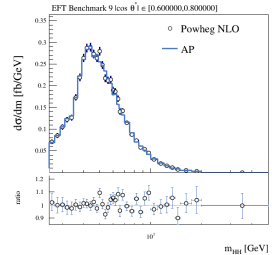
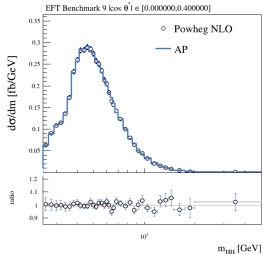
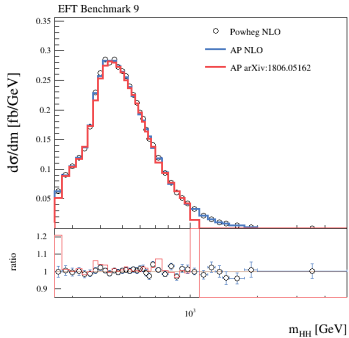


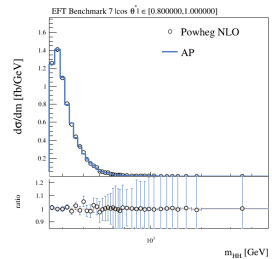
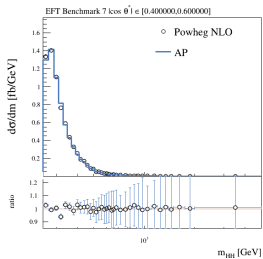
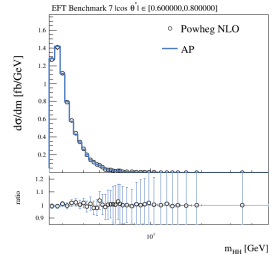
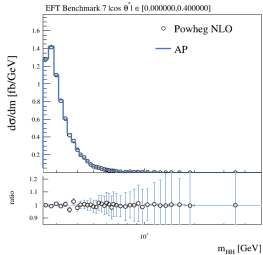
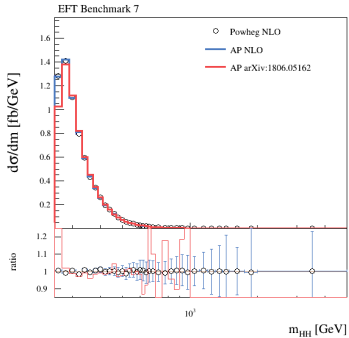


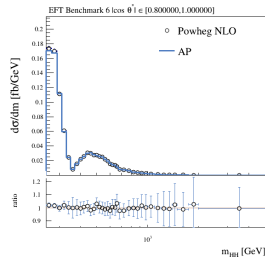
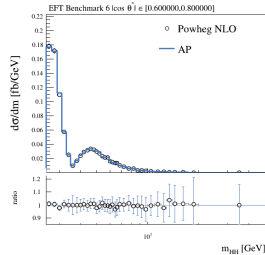
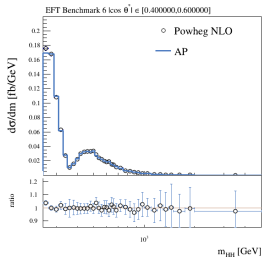
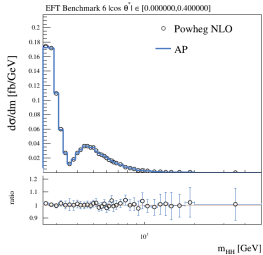
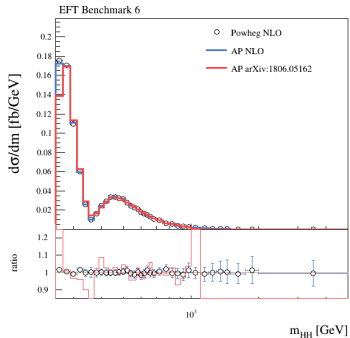


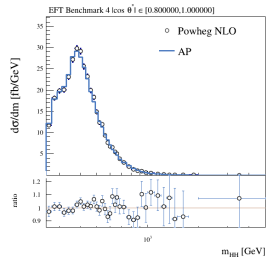
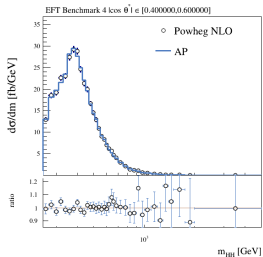
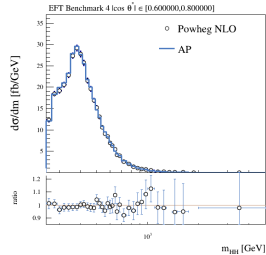
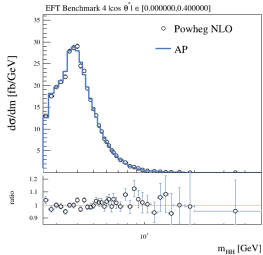
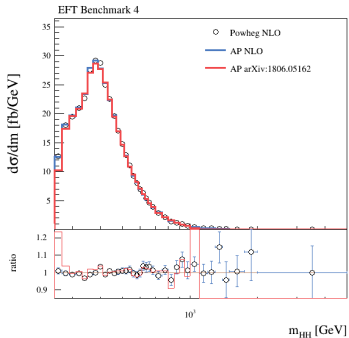




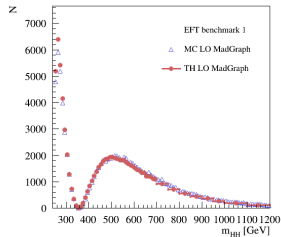
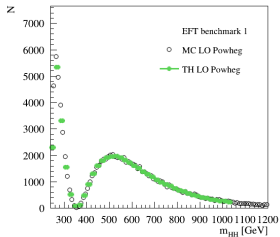
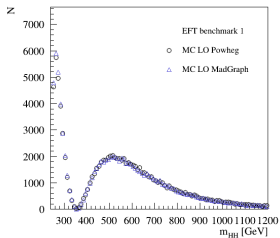
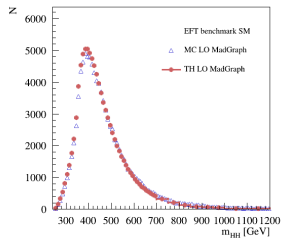
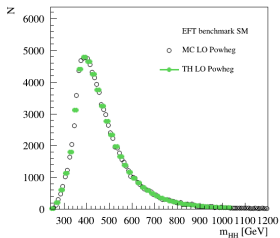
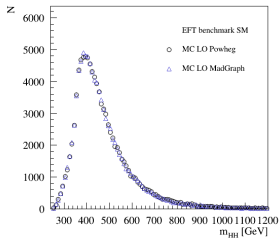












# $M_{HH}$ @ LO

