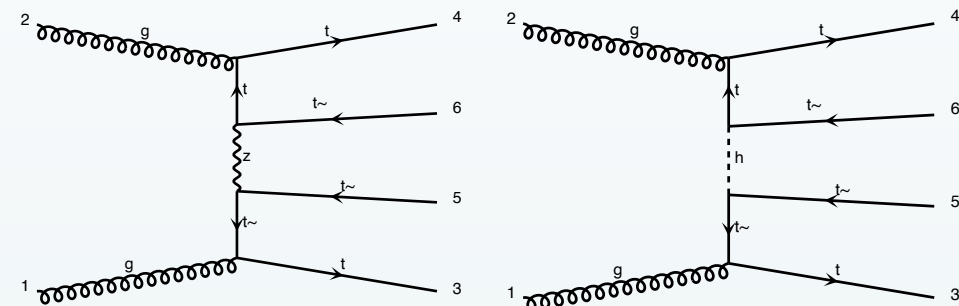
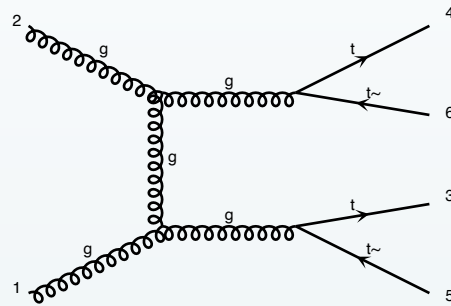


# SM predictions for 4-top production

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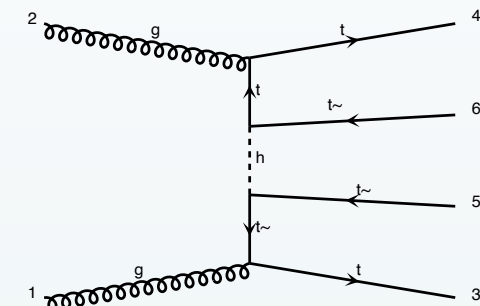
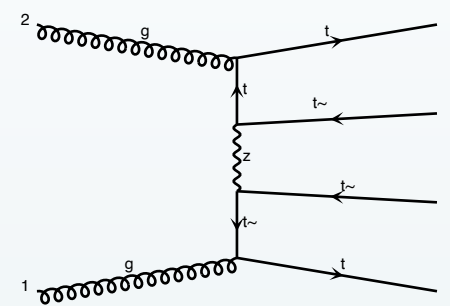
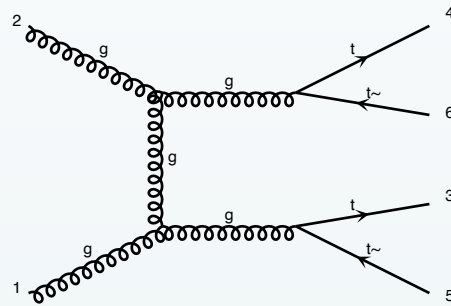
# Anatomy of four-tops



- Small cross section at the LHC,  $\sim 5\text{-}15 \text{ fb @ } 13 \text{ TeV}$
- Dominated by gluon fusion
- "Subleading orders" can be significant: sensitive to top Yukawa

Leading order	X-section LHC-13	scale dependence
QCD <sup>2</sup>	6.8 fb	+73% -40%
EW <sup>2</sup>	2.6 fb	+45% -30%
QCD*EW	-1.8 fb	+62% -36%

# "Subleading orders"

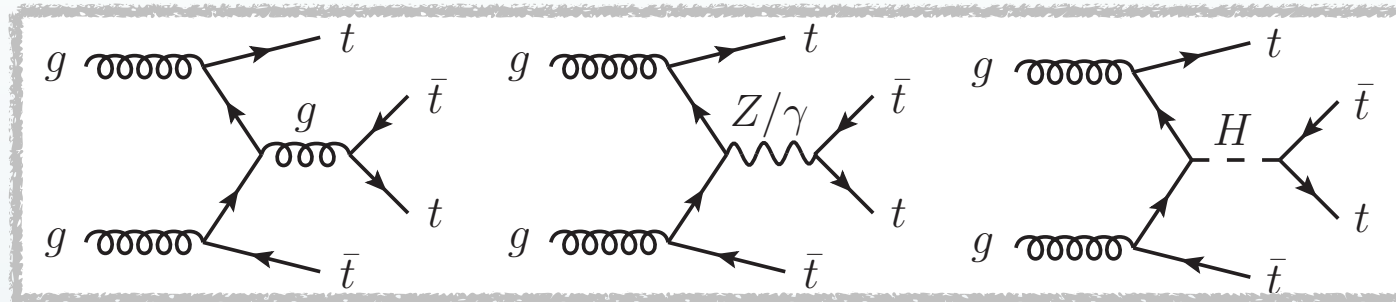


- Large "subleading" contributions due to scattering of  $tt \rightarrow tt$  through the exchange of (heavy) bosons
- Rather different kinematics as compared to QCD 4-top production:
  - Scale uncertainties are different — LO scale dependence is not a good estimate of the actual uncertainties
  - No reason to expect (theoretical) correlations among contribution (and similar behaviour when computing NLO corrections)

Leading order	X-section LHC-13	scale dependence
QCD <sup>2</sup>	6.8 fb	+73% -40%
EW <sup>2</sup>	2.6 fb	+45% -30%
QCD*EW	-1.8 fb	+62% -36%

# Four-top production and Top Yukawa coupling

[Cao, Chen & Liu, 2016]



$$\begin{aligned} \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g+Z/\gamma} &\propto |\mathcal{M}_g + \mathcal{M}_{Z/\gamma}|^2, \\ \sigma^{\text{SM}}(t\bar{t}t\bar{t})_H &\propto |\mathcal{M}_H|^2, \\ \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{\text{int}} &\propto \mathcal{M}_{g+Z/\gamma} \mathcal{M}_H^\dagger + \mathcal{M}_{g+Z/\gamma}^\dagger \mathcal{M}_H \end{aligned}$$

$$\sigma(t\bar{t}t\bar{t}) = \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g+Z/\gamma} + \kappa_t^2 \sigma_{\text{int}}^{\text{SM}} + \kappa_t^4 \sigma^{\text{SM}}(t\bar{t}t\bar{t})_H$$

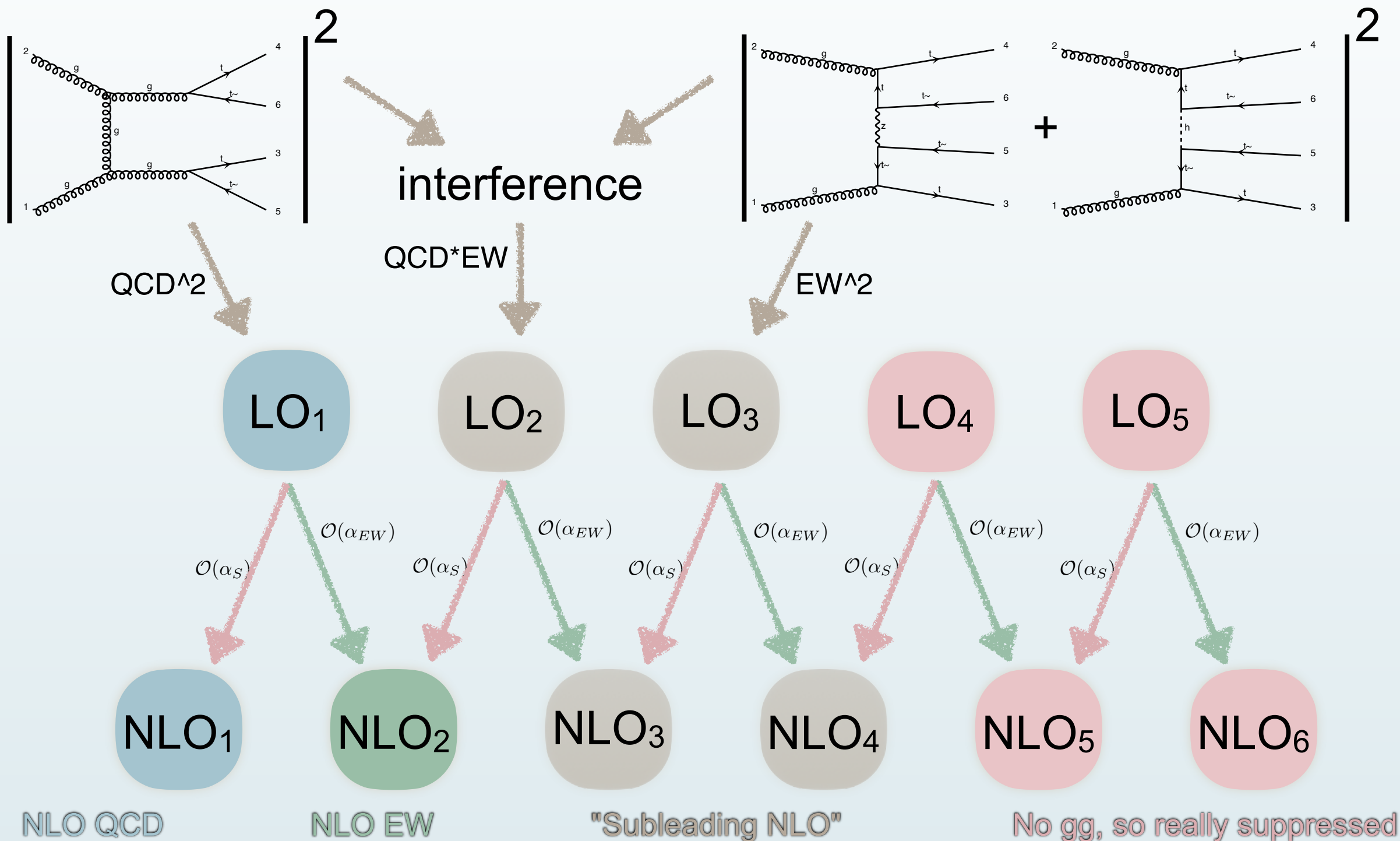
- Four-top production can be used to constrain/measure a **anomalous top Yukawa coupling** independently from the Higgs width
  - kappa-framework
- Large contributions from subleading **LO<sub>i</sub>**, with large cancelations
  - **How do NLO corrections affect these?**

	8 TeV	14 TeV
$\sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g+Z/\gamma}$	1.193 fb,	12.390 fb,
$\sigma^{\text{SM}}(t\bar{t}t\bar{t})_H$	0.166 fb,	1.477 fb,
$\sigma^{\text{SM}}(t\bar{t}t\bar{t})_{\text{int}}$	-0.229 fb,	-2.060 fb.

QCD<sup>2</sup> (points to first row)  
EW<sup>2</sup> (points to second and third rows)  
QCD\*EW (points to second and third rows)



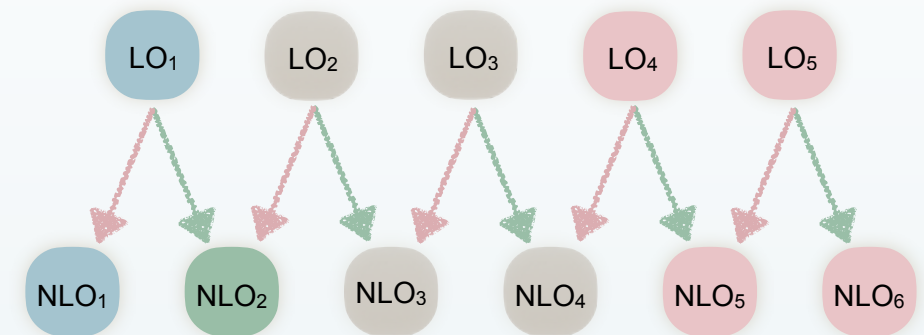
# Perturbative expansion at NLO



# EW & Subleading NLO corrections

- **Additive vs Multiplicative EW corrections**

- Differ at NNLO; multiplicative is preferred if QCD and EW corrections factorise. Not the case here.



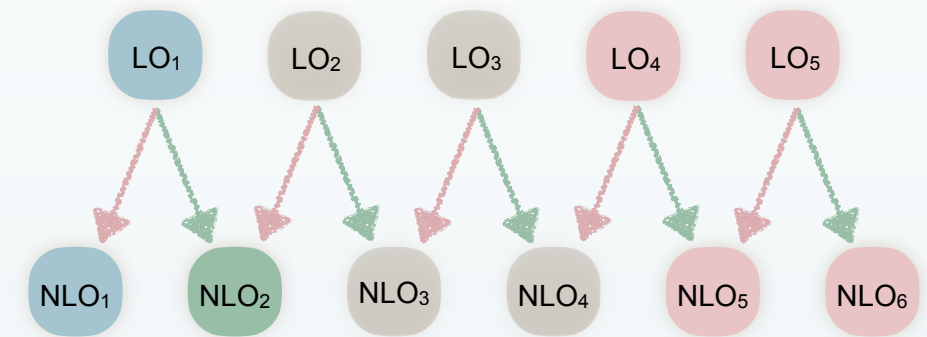
- Not possible to single out the Higgs Yukawa coupling and compute QCD corrections to those separately

- When computing EW corrections, the Yukawa is linked to the top mass and other EW couplings through renormalisation: they are not independent
- Need complete SMEFT to be able to compute NLO corrections and study in the impact of NLO corrections: currently beyond our capabilities
- At NLO, also contributions that have single, triple, quintuple powers of the Yukawa coupling (and not only double or quadruple)

- **Compute NLO corrections in the Standard Model**

- Central scale values in our studies follow the recipe by Maltoni et al. [arXiv:1507.05640]

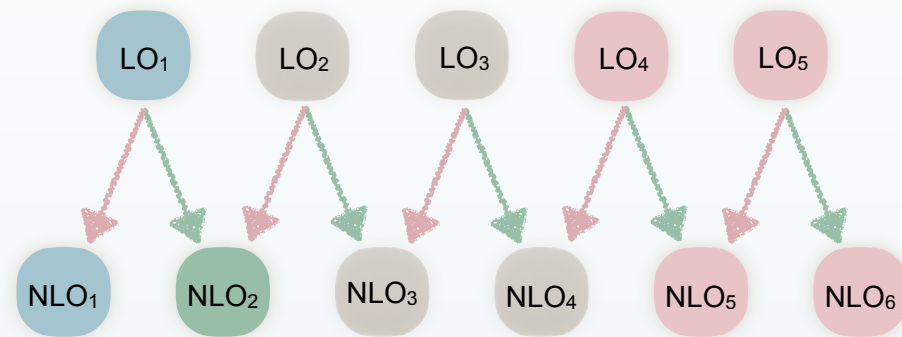
# 4-top: total cross section



$\sigma[\text{fb}]$	$\text{LO}_{\text{QCD}}$	$\text{LO}_{\text{QCD}} + \text{NLO}_{\text{QCD}}$	LO	LO + NLO	$\frac{\text{LO}(+\text{NLO})}{\text{LO}_{\text{QCD}}(+\text{NLO}_{\text{QCD}})}$
$\mu = H_T/4$	$6.83^{+70\%}_{-38\%}$	$11.12^{+19\%}_{-23\%}$	$7.59^{+64\%}_{-36\%}$	$11.97^{+18\%}_{-21\%}$	1.11 (1.08)

- NLO corrections are large but within the LO scale uncertainty band
- "Subleading" NLO corrections increase the X-section by an additional 8%
- Remaining scale uncertainty of the order of 20% at NLO

# NLO 4-top production in the SM



- $LO_2$  and  $LO_3$  have large cancelations
- $NLO_2$  and  $NLO_3$  mainly given by QCD corrections on top of them
  - large and strongly dependent on the scale choice
- However, the sum of  $NLO_2 + NLO_3$  very stable and small
- Different scale choices have even more extreme cancelations between  $NLO_2$  and  $NLO_3$

[RF, Pagani, Zaro, 2017]

$$\delta_{(N)LO_i}(\mu) = \frac{\Sigma_{(N)LO_i}(\mu)}{\Sigma_{LO_{QCD}}(\mu)}$$

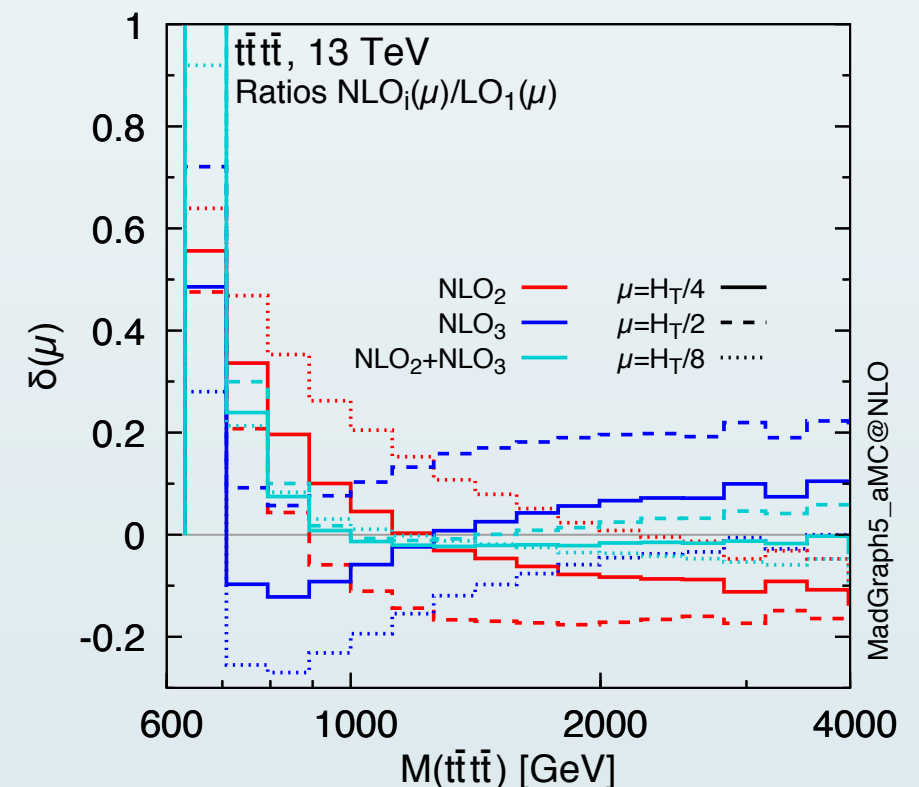
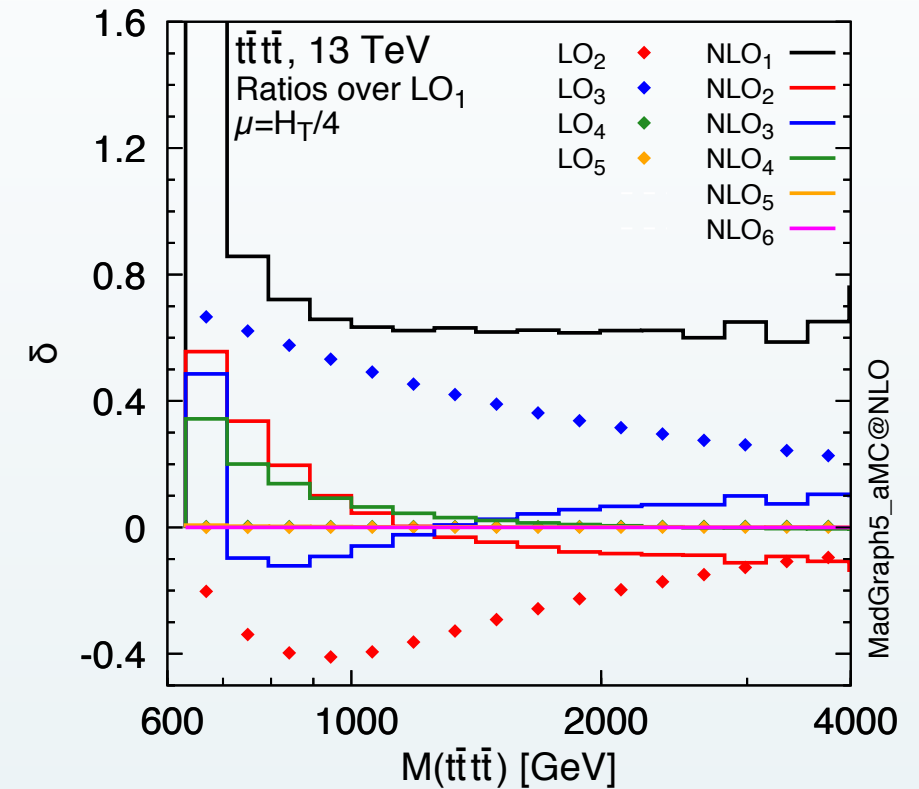
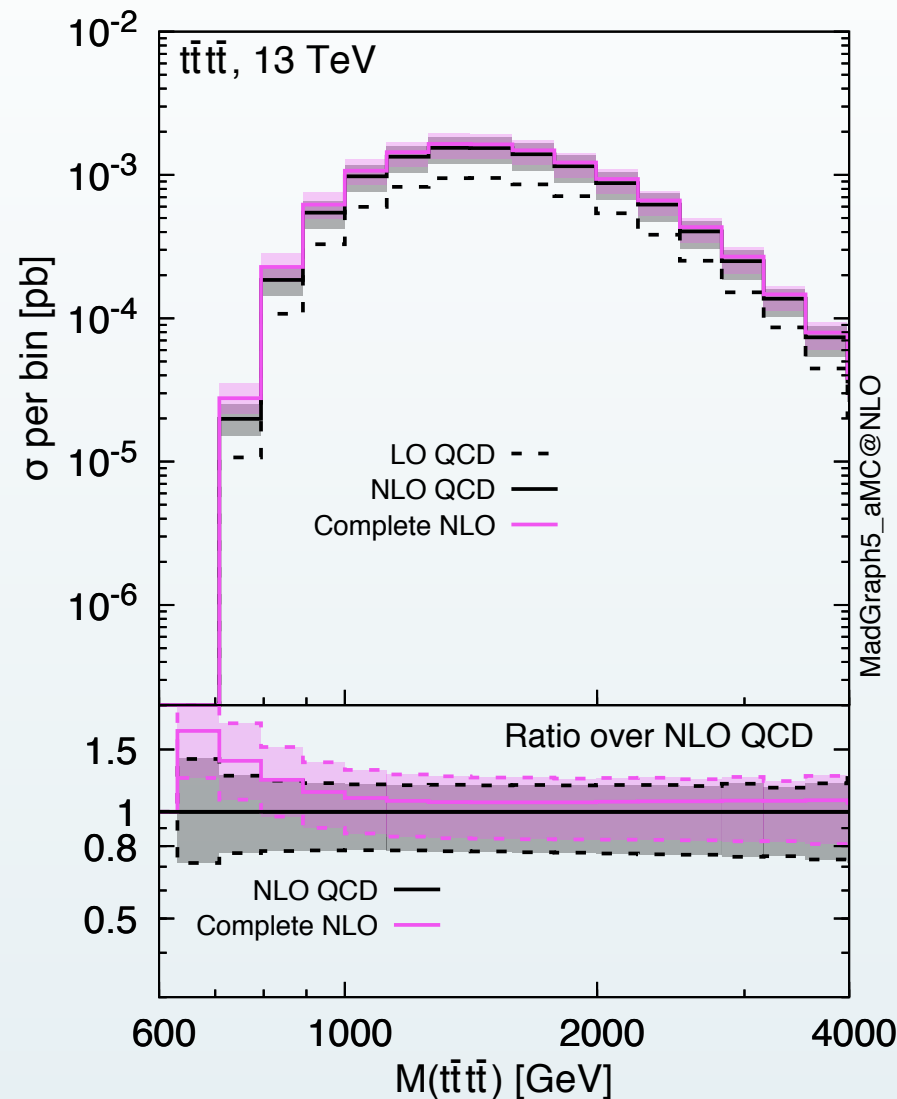
$\sigma[\text{fb}]$	$LO_{QCD}$			Naive expectation
$\mu = H_T/4$	$\mu = H_T/8$	$\mu = H_T/4$	$\mu = H_T/2$	
	$6.83^{+70\%}_{-38\%}$			
$\delta[\%]$	$\mu = H_T/8$	$\mu = H_T/4$	$\mu = H_T/2$	
$LO_2$	-26.0	-28.3	-30.5	10%
$LO_3$	32.6	39.0	45.9	1%
$LO_4$	0.2	0.3	0.4	0.1%
$LO_5$	0.02	0.03	0.05	0.01%
$NLO_1$	14.0	62.7	103.5	10%
$NLO_2$	8.6	-3.3	-15.1	1%
$NLO_3$	-10.3	1.8	16.1	0.1%
$NLO_4$	2.3	2.8	3.6	0.01%
$NLO_5$	0.12	0.16	0.19	0.001%
$NLO_6$	< 0.01	< 0.01	< 0.01	0.0001%
$NLO_2 + NLO_3$	-1.7	-1.6	0.9	

13 TeV

◆  $LO_4$ ,  $(N)LO_5$  and  $NLO_6$  only qqbar initial state. Hence, very small

# four-top invariant mass

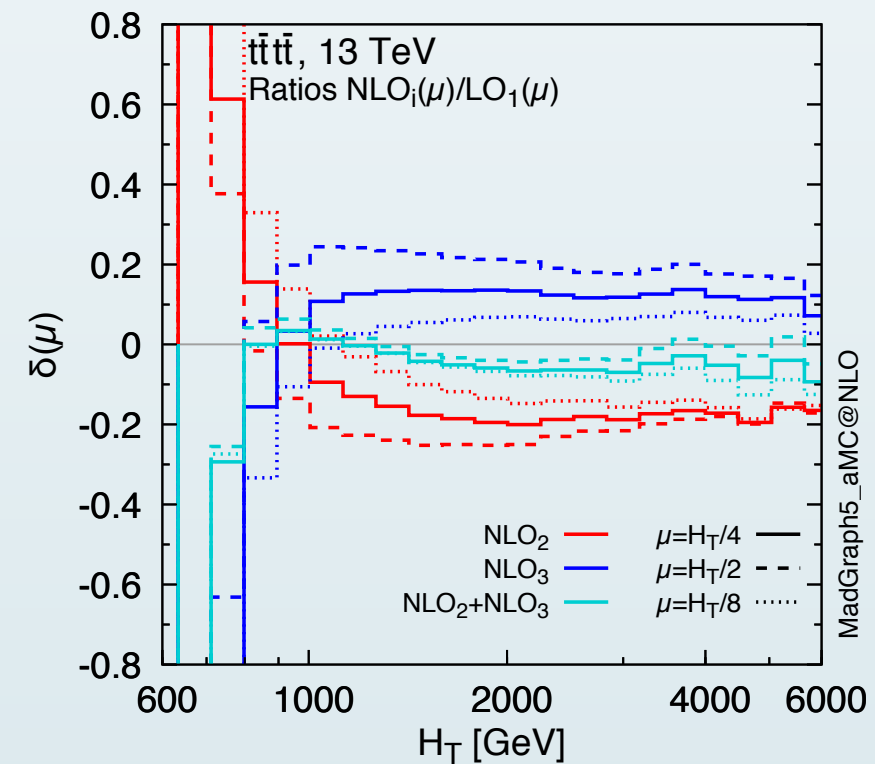
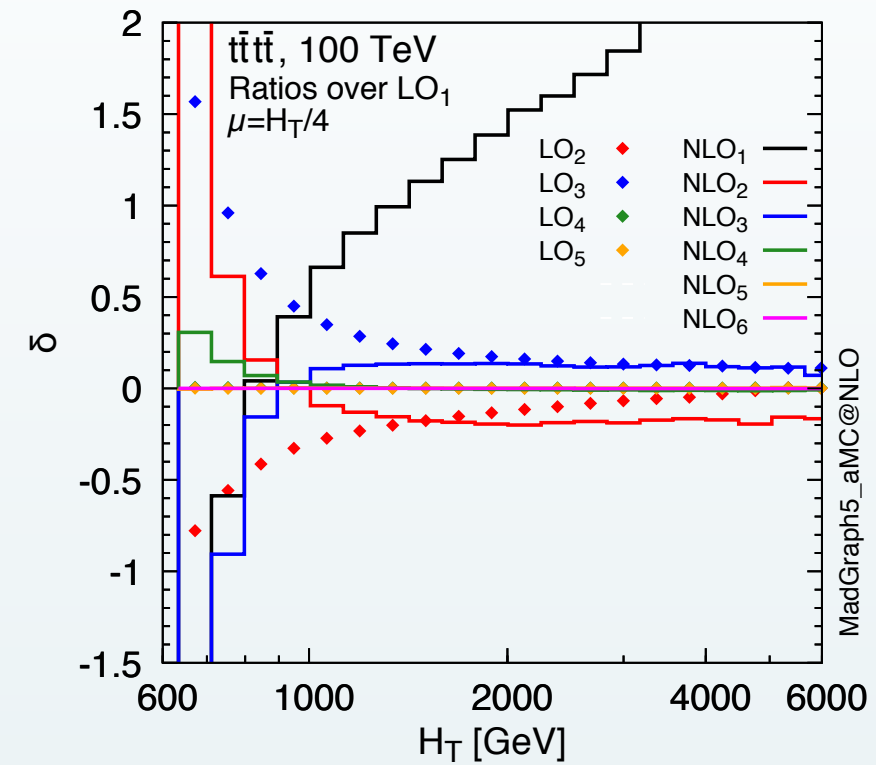
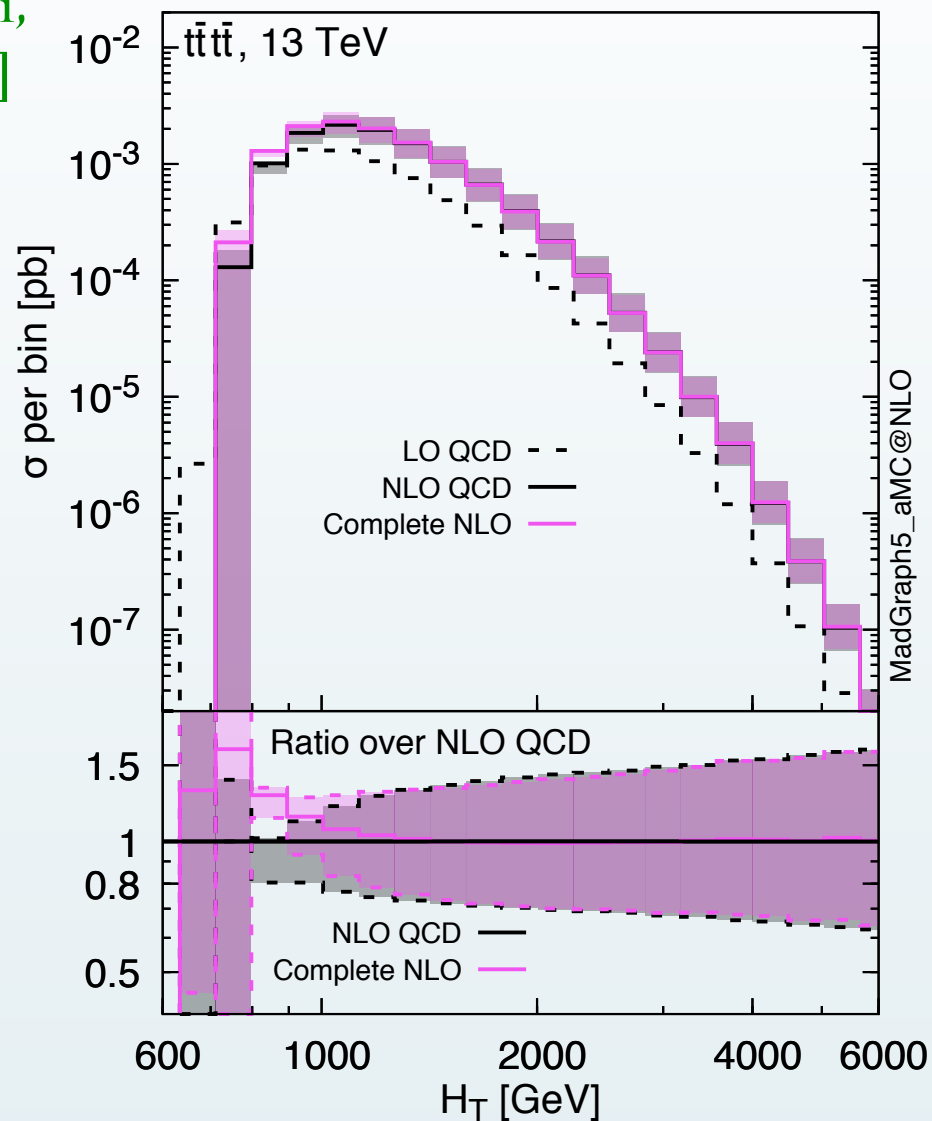
[RF, Pagani,  
Zaro, 2017]



- Large cancellations between **NLO<sub>2</sub>** and **NLO<sub>3</sub>** also at the differential level
- **NLO<sub>4</sub>** large at threshold

# $H_T$ distribution

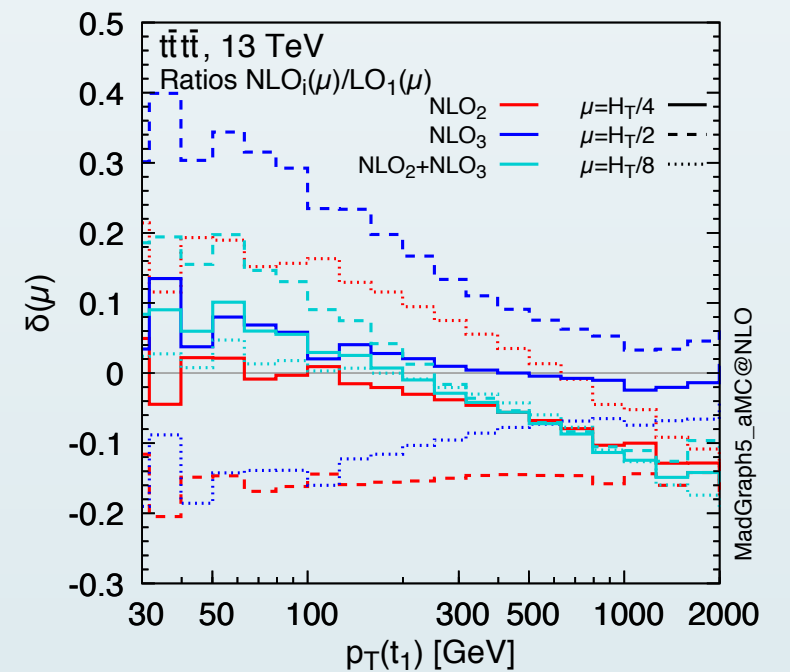
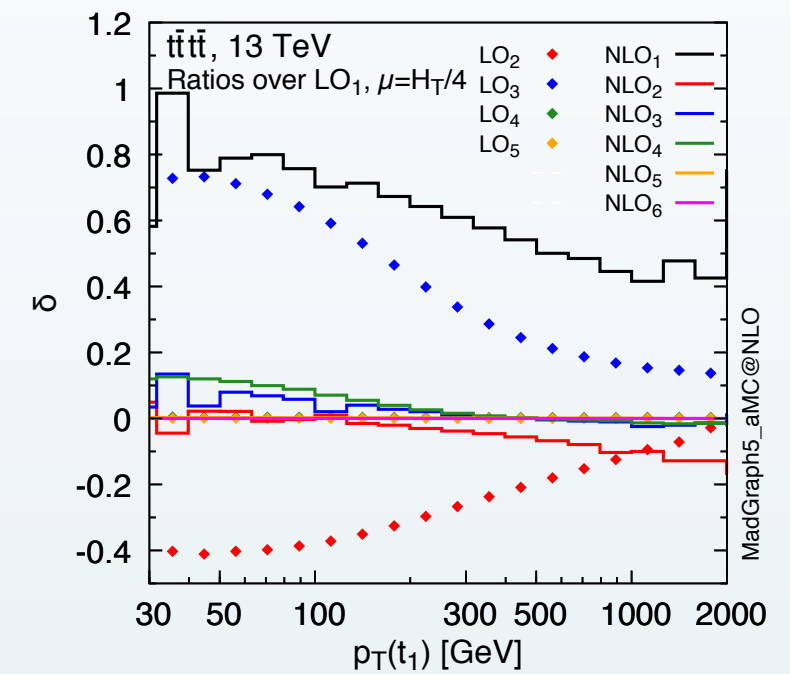
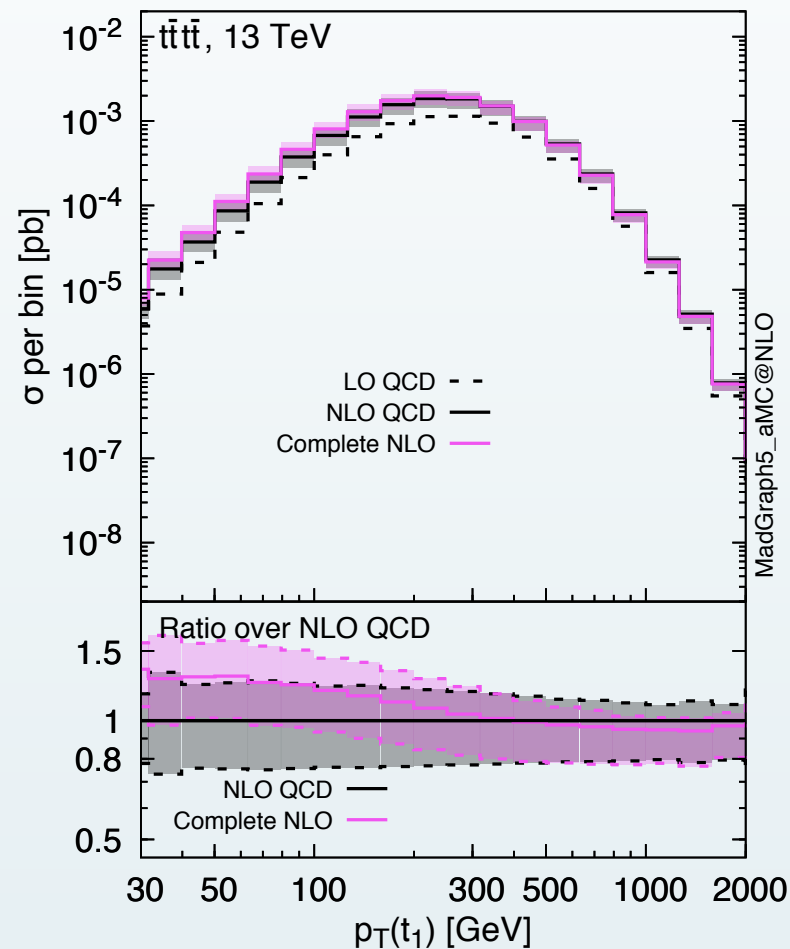
[RF, Pagani,  
Zaro, 2017]



- Similar to 4-top invariant mass

[RF, Pagani,  
Zaro, 2017]

# $p_T$ of hardest top



- Large "subleading" corrections at small transverse momentum



# Conclusions

- NLO corrections to 4-top production are special:
  - For no other process  $LO_2$  is as large as here.
  - For no other process  $NLO_2$  is dominated by QCD corrections (and not EW corrections)
- **Large cancelations among  $NLO_i$  corrections**
  - Surprising!
  - Not possible to estimate how this affects BSM physics (e.g. modified Yukawa coupling)
    - Is there still a cancelation?
    - Would require computing the complete-NLO corrections in the SMEFT, which is currently out-of-reach