# Production and signatures of multiflavour dark matter scenarios

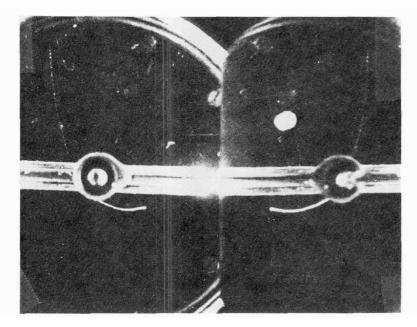
Alejandro Ibarra



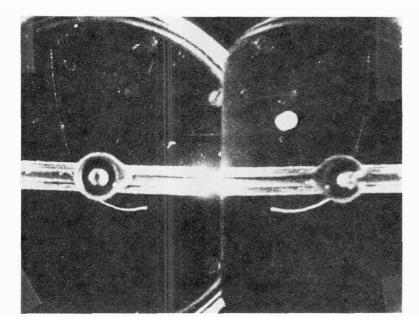


Based on 2103.10392, 1912.09458, in collaboration with Johannes Herms.

# <u>Introduction</u>



Anderson, Neddermeyer '36

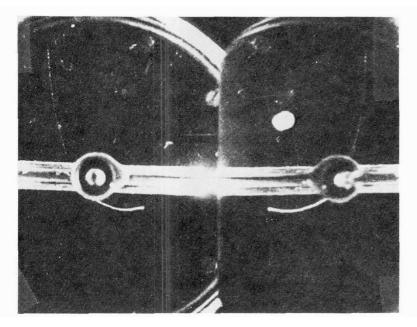


Anderson, Neddermeyer '36

Same charge as the electron.

Mass between the electron and the proton mass.

Mesotron

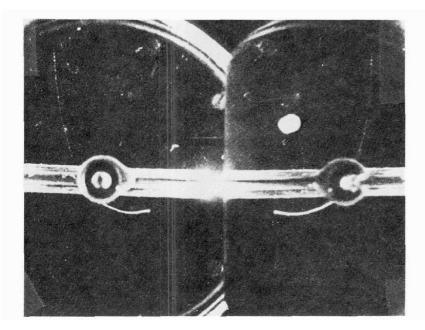


Anderson, Neddermeyer '36

Same charge as the electron.

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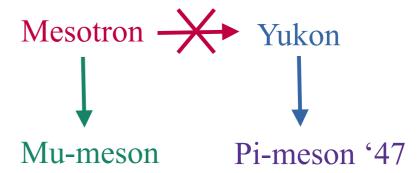
Mesotron — Yukon

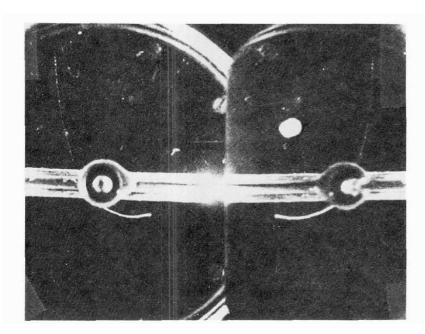


Anderson, Neddermeyer '36

Same charge as the electron.

Mass between the electron and the proton mass.

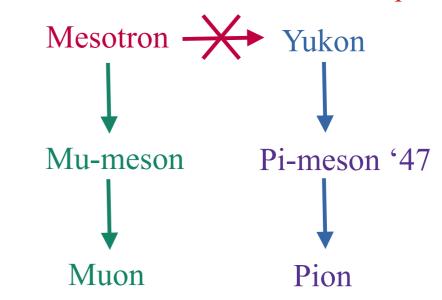


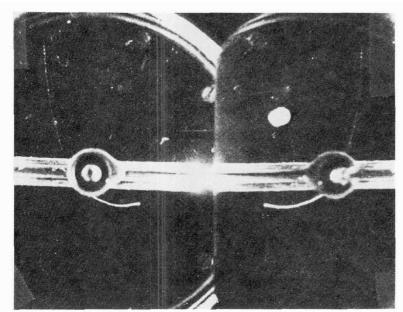


Anderson, Neddermeyer '36

Same charge as the electron.

Mass between the electron and the proton mass.

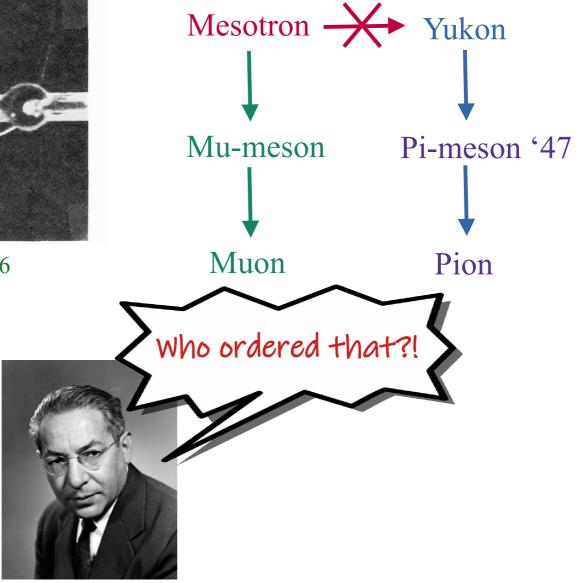




Anderson, Neddermeyer '36

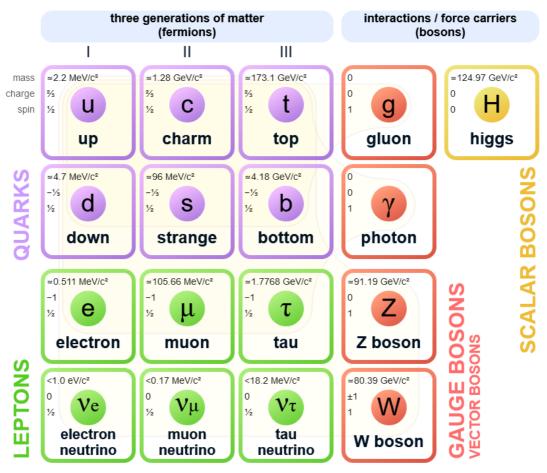
Same charge as the electron.

Mass between the electron and the proton mass.



Isidor Rabi

#### Standard Model of Elementary Particles



As of today, we still don't understand why all known fermions are replicated three times, with different masses (Yukawa couplings?)

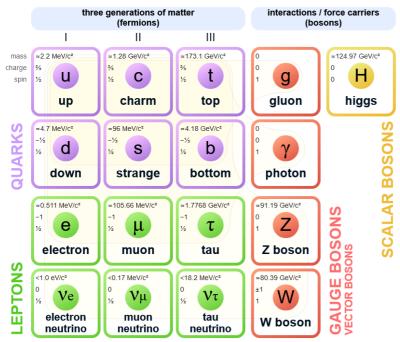
# <u>Introduction</u>

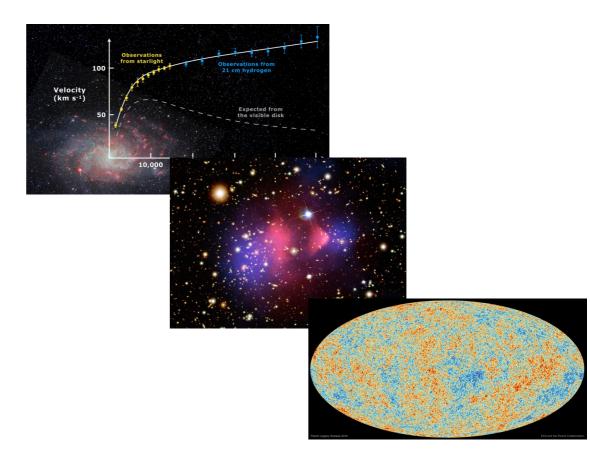
• Who ordered the heavier generations? Our Universe today is mostly made of first generation fermions: u,d,e (plus three massive neutrinos).

• However, the heavier generations could have played a crucial role in shaping our Universe (CP violation in baryogenesis)

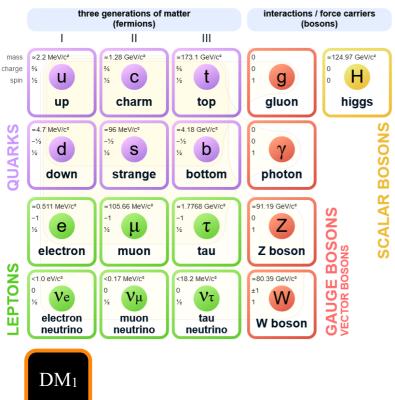
•Also, the heavier generations lead to many phenomena in experiments, which helped to elucidate the structure of the Standard Model.

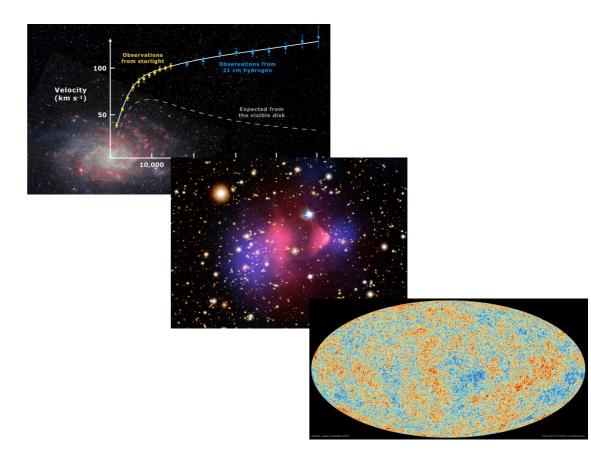
#### **Standard Model of Elementary Particles**



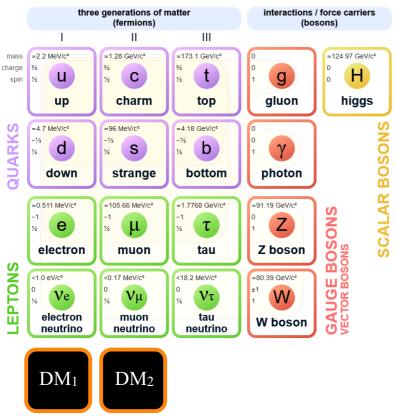


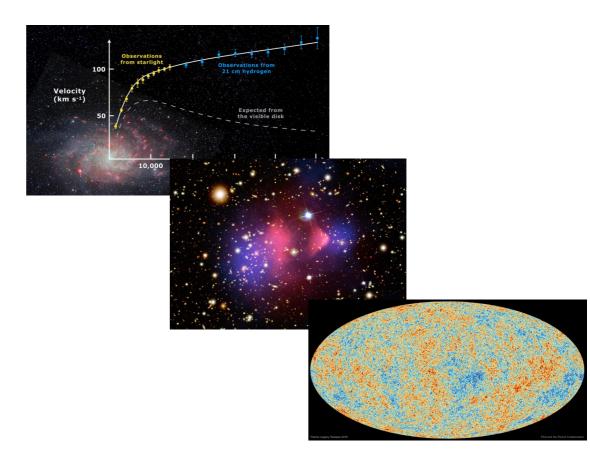
#### **Standard Model of Elementary Particles**



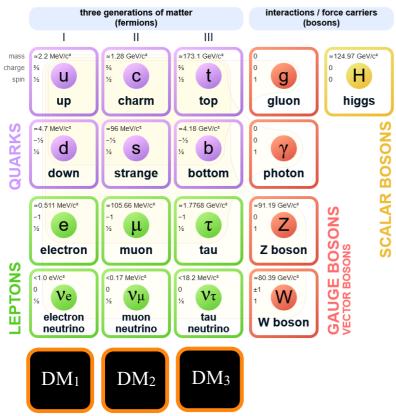


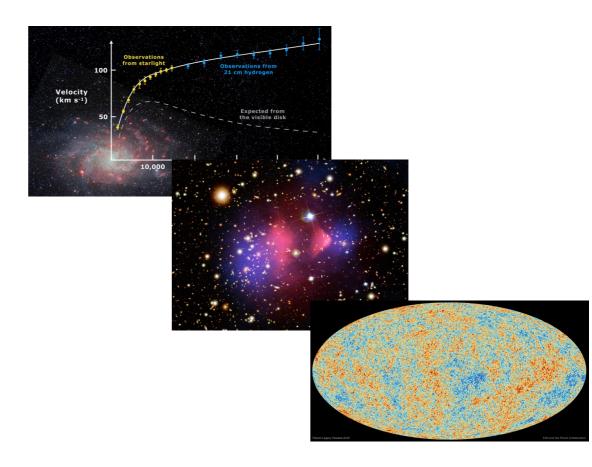
#### **Standard Model of Elementary Particles**



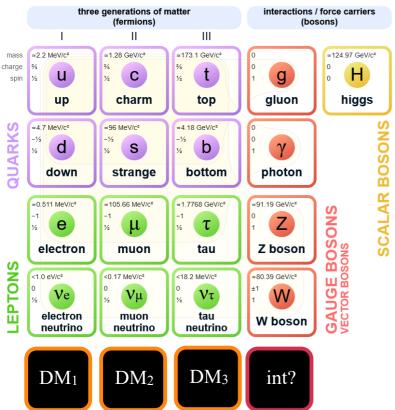


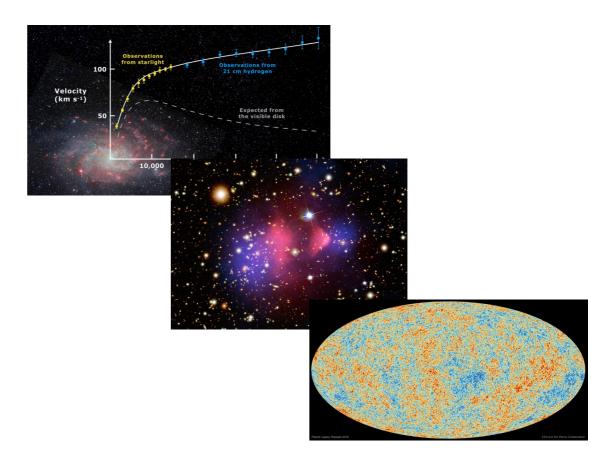
#### **Standard Model of Elementary Particles**





#### **Standard Model of Elementary Particles**





# <u>Introduction</u>

• Who ordered the heavier generations? Our Universe today is mostly made of first generation fermions: u,d,e (plus three massive neutrinos).

Could the DM be made by two long-lived components?

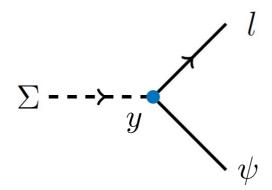
• However, the heavier generations could have played a crucial role in shaping our Universe (CP violation in baryogenesis)

Could the heavy DM flavors also play a role in shaping our Universe?

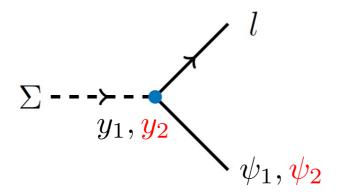
•Also, the heavier generations lead to many phenomena in experiments, which helped to elucidate the structure of the Standard Model.

Could the heavy DM flavors lead to new signals, that will help us elucidate the structure of the dark sector?

 $\bullet$  Consider a Majorana singlet fermion as DM candidate. Assume that it couples to the right-handed leptons through a scalar mediator  $\Sigma$ .



• Consider a Majorana singlet fermion as DM candidate. Assume that it couples to the right-handed leptons through a scalar mediator  $\Sigma$ . Add a second generation of the Majorana singlet fermion.



• Free parameters:  $m_1, m_2, m_{\Sigma}, y_1, y_2$  (+ quartic couplings)

initial state		final state		$y_i$ scaling
$\Sigma^+$	$\Sigma^-$	$\begin{array}{c c} \gamma \\ \gamma \\ Z \\ W^+ \\ H \\ H \\ q \\ l^+ \end{array}$	$\begin{array}{c} \gamma \\ Z \\ \overline{Z} \\ W^{-} \\ \overline{H} \\ \overline{Z} \\ \overline{q} \\ \overline{l}^{-} \end{array}$	1
$\Sigma^+$	$\Sigma^+$	$l^+$	$\frac{l}{l^+}$	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2 y_2^2  y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$\frac{y_1,y_1,y_2,y_2}{y_i^2}$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \hline \Sigma^{\pm} \end{array}$	$\frac{\psi_j}{l^{\pm}}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_j^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	$y_i^2$
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$y_i^2 \\ y_i^2$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

initial state		final state		$y_i$ scaling
		$\frac{\gamma}{\gamma}$	$\frac{\gamma}{Z}$	
$\Sigma^+$	$\Sigma^-$	$Z$ $W^+$	$\frac{Z}{W^{-}}$	1
	<b>_</b>	H	$\frac{H}{Z}$	
			$\frac{Z}{\bar{q}}$	
		$\frac{q}{l^+}$	$l^-$	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2 y_2^2$
$\Sigma^+$	$\Sigma^+$	$l^+$	$l^+$	$y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$y_i^2$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \hline \Sigma^{\pm} \end{array}$	$l^{\pm}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_j^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	0
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$\frac{y_i^2}{y_i^2}$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

Annihilation

initial state		final	state	$y_i$ scaling
		$\gamma$ $\gamma$	$\frac{\gamma}{Z}$	
$\Sigma^+$	$\Sigma^-$	$X = W^+$	$\frac{Z}{W^-}$	1
		H	$\frac{H}{Z}$	
		q	$\frac{Z}{\bar{q}}$	
		$l^+$	$l^-$	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2 y_2^2  y_1^4, y_1^2 y_2^2, y_2^4$
$\Sigma^+$	$\Sigma^+$	$l^+$	$l^+$	$y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$y_i^2$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \Sigma^{\pm} \end{array}$	$l^{\pm}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_j^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	$y_i^2$
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$y_i^2$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

Co-annihilation

initial state		final state		$y_i$ scaling
		$\begin{array}{c c} \gamma \\ \gamma \\ Z \end{array}$	$\frac{\gamma}{Z}$	
$\Sigma^+$	$\Sigma^-$	$W^+$	$W^-$	1
		H	$\frac{H}{Z}$	
			$\frac{Z}{\bar{q}}$	
		$l^+$	<i>l</i> -	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2 y_2^2$
$\Sigma^+$	$\Sigma^+$	<i>l</i> +	$l^+$	$y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$y_i^2$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 \\ y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \Sigma^{\pm} \end{array}$	$l^{\pm}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_i^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	$y_i^2$
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$y_i^2$ $y_i^2$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2 \\ y_1^2 y_2^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

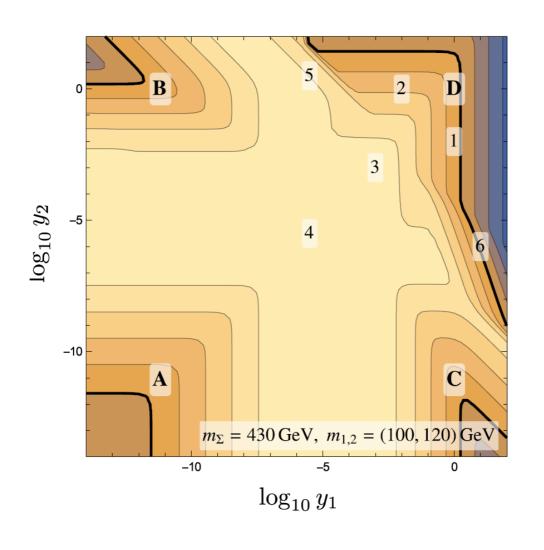
 $\Sigma\Sigma$  pair annihilation

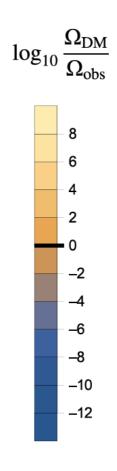
initial state		final state		$y_i$ scaling
$\Sigma^+$	$\Sigma^-$	$\gamma$ $\gamma$	$\frac{\gamma}{Z}$	
		$X = W^+$	$\frac{Z}{W^{-}}$	1
		H	H	
		$\frac{H}{q}$	$\frac{Z}{\bar{q}}$	
		$l^+$	$l^-$	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2 y_2^2$
$\Sigma^+$	$\Sigma^+$	$l^+$	$l^+$	$y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$y_i^2$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \Sigma^{\pm} \end{array}$	$l^{\pm}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_j^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	$y_i^2$
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$y_i^2$ $y_i^2$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

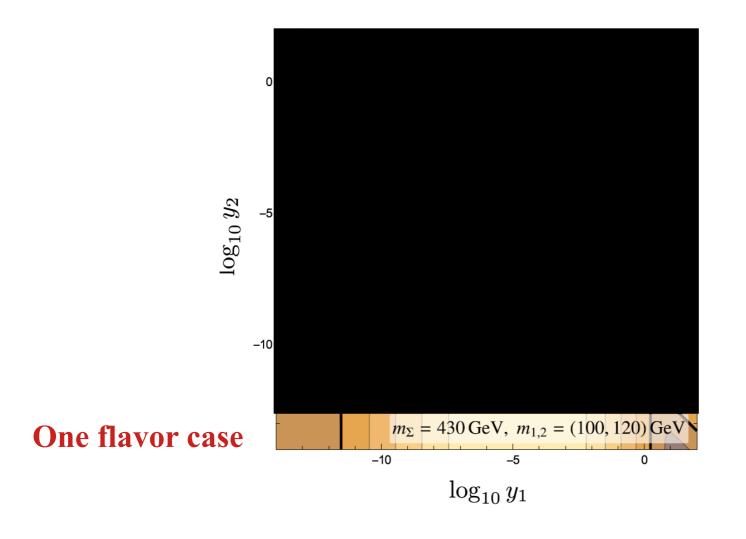
Conversion

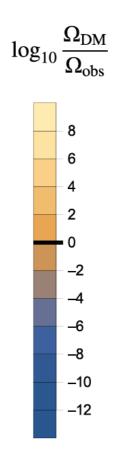
initial state		final	state	$y_i$ scaling
		$\begin{array}{c c} \gamma \\ \gamma \\ Z \end{array}$	$\frac{\gamma}{Z}$	
$\Sigma^+$	$\Sigma^-$	$W^+$	$W^ H$	1
		H	Z	
		$l^+$	$\frac{\bar{q}}{l^-}$	$1, y_1^2, y_1^4, y_2^2, y_2^4, y_1^2y_2^2$
$\Sigma^+$	$\Sigma^+$	$l^+$	$l^+$	$y_1^4, y_1^2 y_2^2, y_2^4$
$\psi_i$	$\Sigma^+$	$l^+$	$\gamma, Z$	$y_i^2$
$\psi_i$	$\psi_j$	$l^+$	$l^-$	$y_i^2 y_j^2$
$\begin{array}{c c} \psi_i \\ \Sigma^{\pm} \\ \hline \Sigma^{\pm} \end{array}$	$l^{\pm}$	$\psi_j$	$l^{\pm}$	$y_i^2 y_j^2$
$\Sigma^{\pm}$	$\gamma, Z$	$\psi_i$	$l^{\pm}$	$y_i^2$
$\Sigma^{\pm}$	$l^{\pm}$	$\psi_i$	$\gamma$	$y_i^2 $ $y_i^2$
$\Sigma^{\pm}$		$\psi_i$	$l^{\pm}$	$y_i^2$
$\psi_2$		$\psi_1$	$l^+l^-$	$y_1^2 y_2^2$

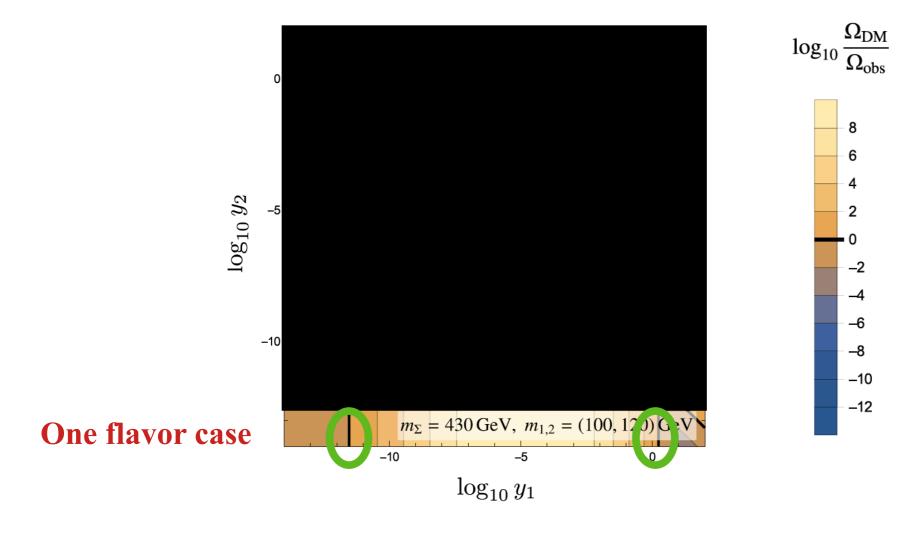
Decays (freeze-in)

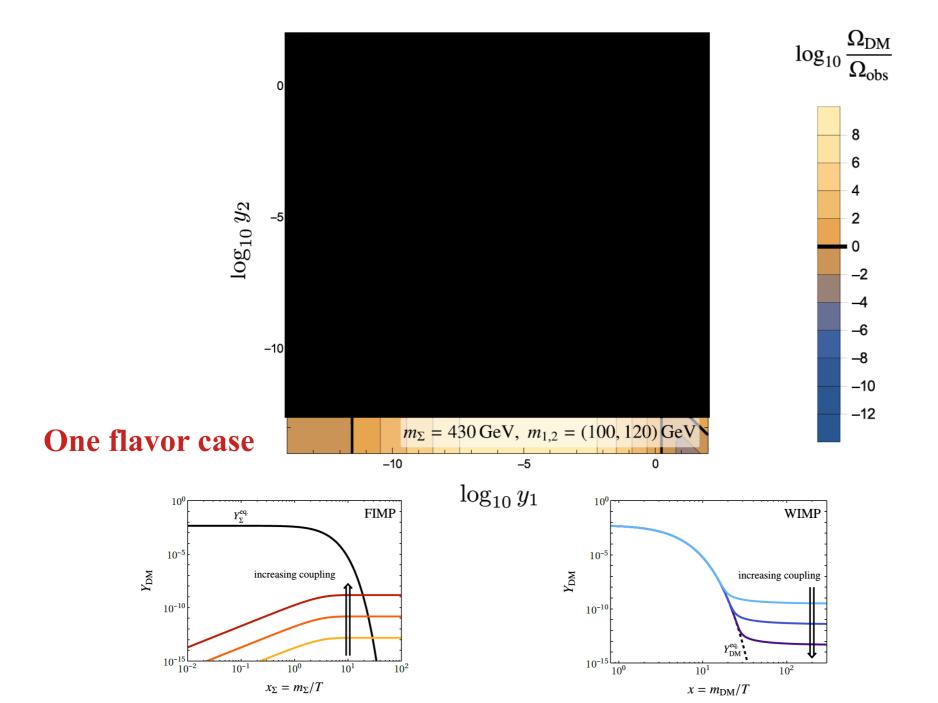


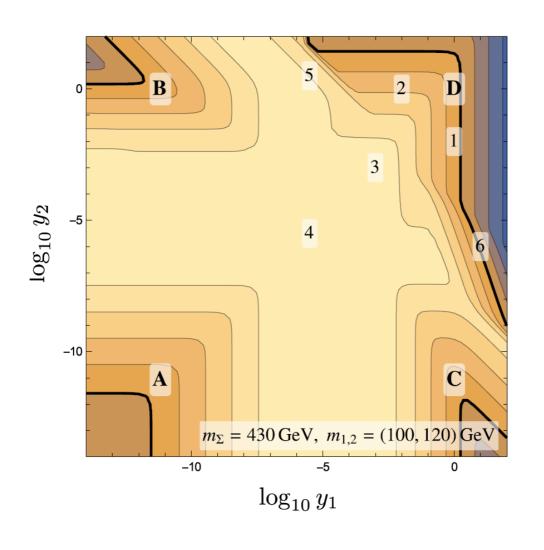


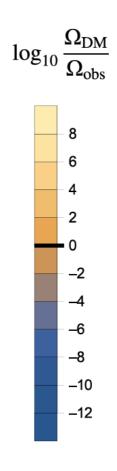


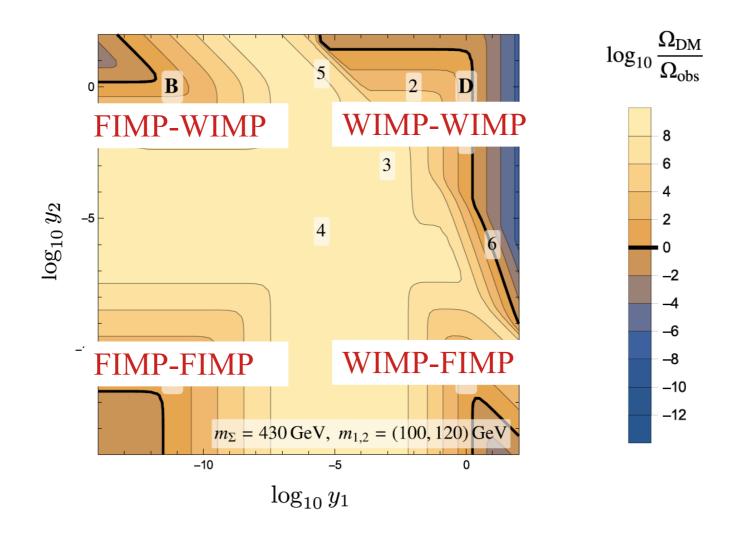


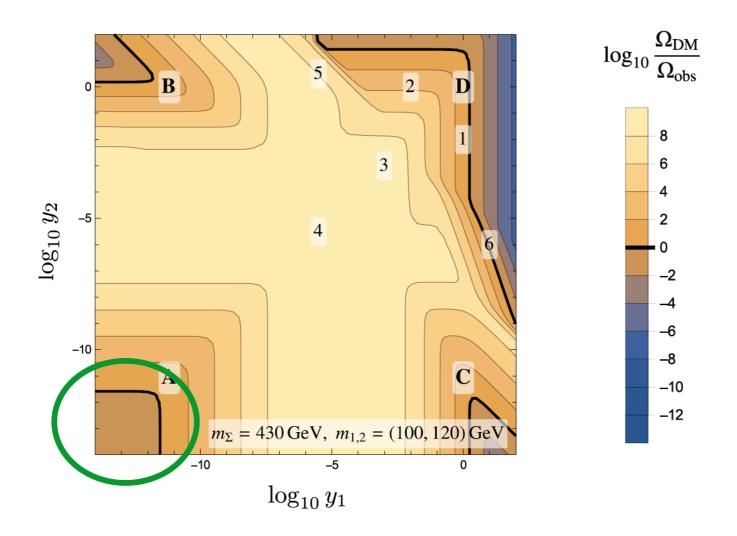




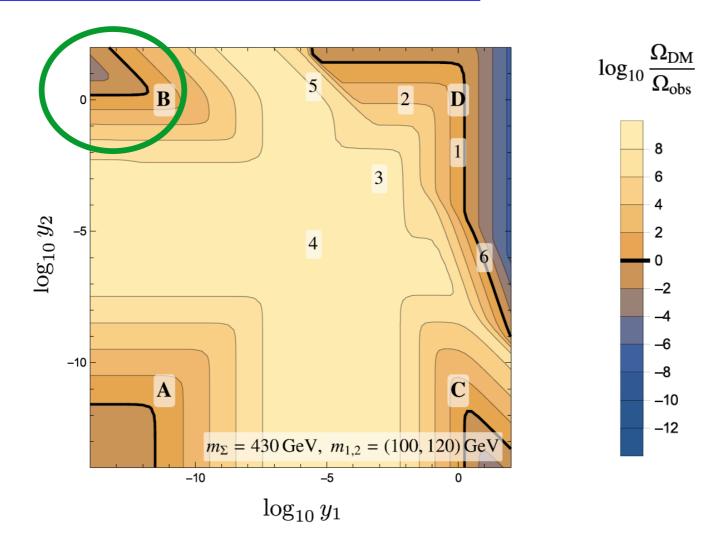




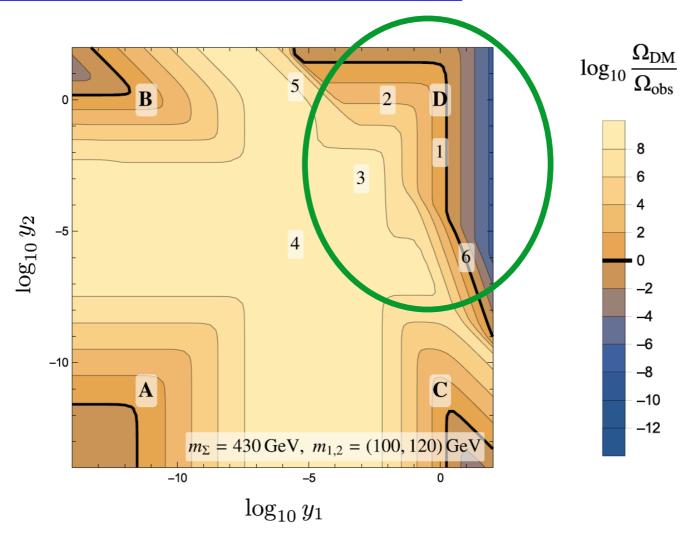




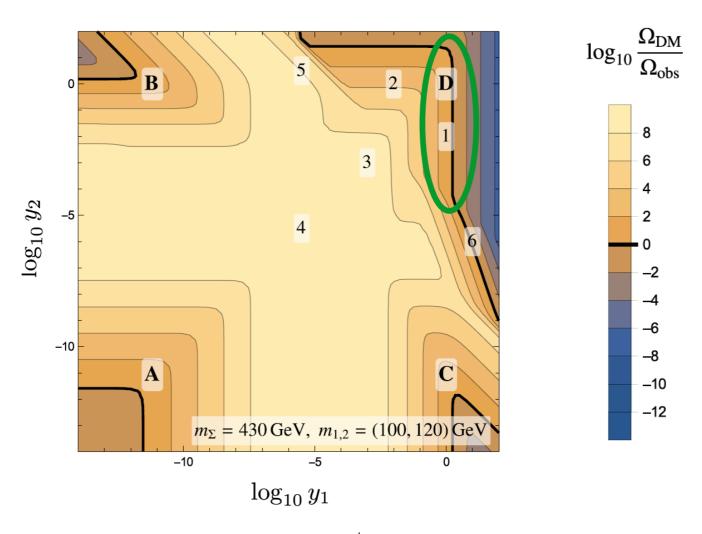
 $\psi_1$  and  $\psi_2$  not thermalized. Abundances evolve independently. Abundance set by  $\Sigma \to \psi_i l$ ,  $\Sigma l \to \psi_i \gamma$ , etc.



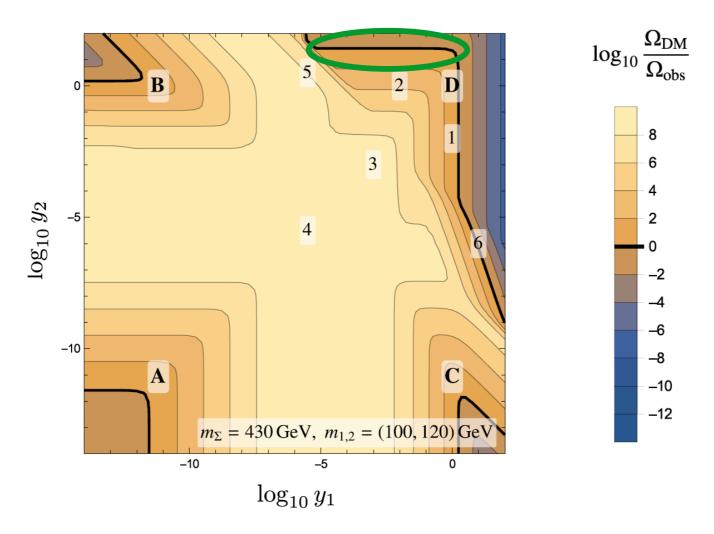
 $\psi_2$  thermalized. Abundance set by  $\psi_2\psi_2 \to l^+l^ \psi_1$  not thermalized. Abundance set by  $\Sigma \to \psi_i l, \ \Sigma l \to \psi_i \gamma, \psi_2 l \to \psi_1 l$ 



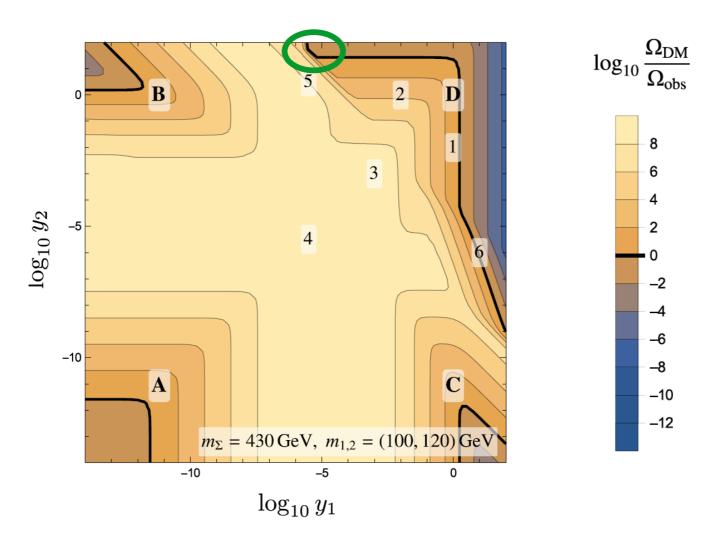
 $\psi_1$  and  $\psi_2$  thermalized.



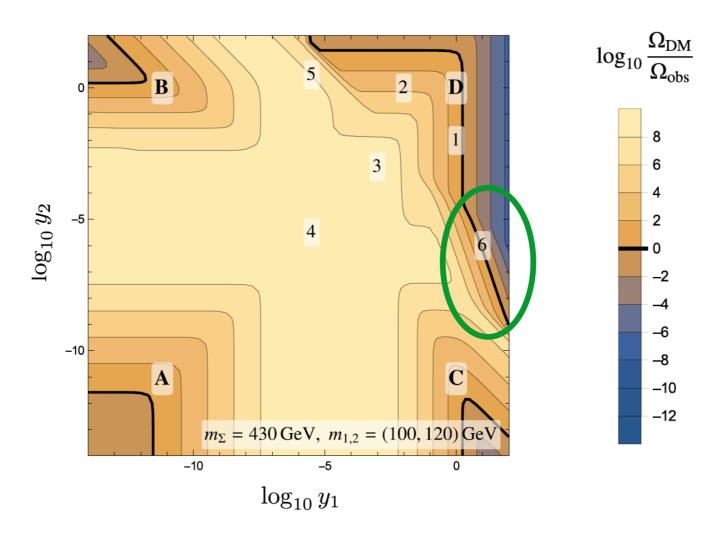
 $\psi_1$  abundance set by  $\psi_1\psi_1 \to l^+l^ \psi_2$  irrelevant



DM abundance set by  $\psi_2$  coannihilation  $\psi_2\psi_2 \to l^+l^-$ 

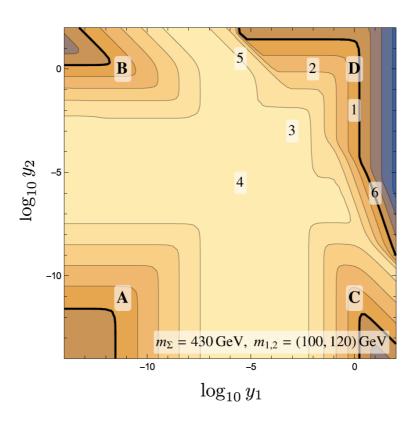


 $\psi_1 \rightarrow \psi_2$  conversion driven freeze-out  $\psi_1$  drops out of chemical equilibrium before  $\psi_2$  coannihilation becomes effective

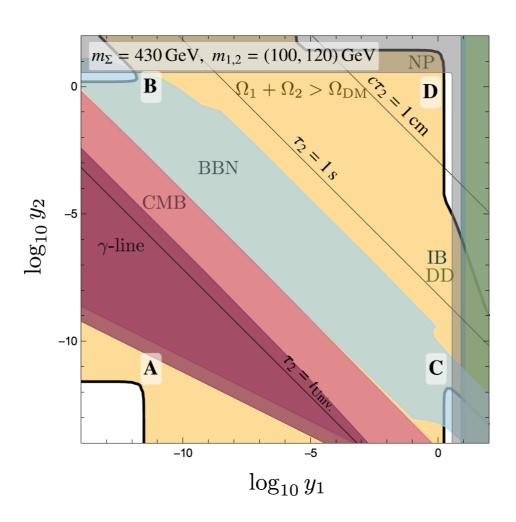


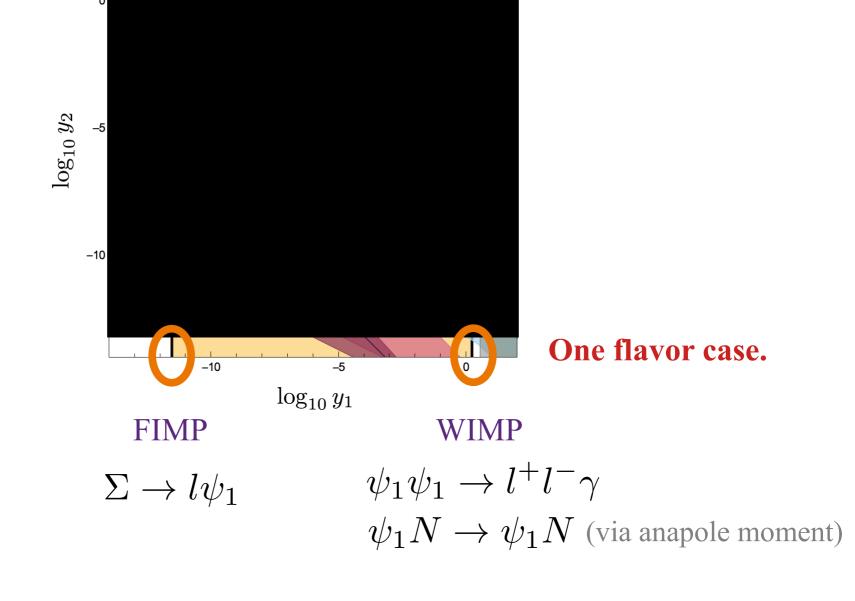
 $\psi_2$  long-lived. DM abundance set by  $\psi_1\psi_1 \to l^+l^-, \psi_2 l \to \psi_1 l,$   $\psi_2 \to \psi_1 l^+l^-$ 

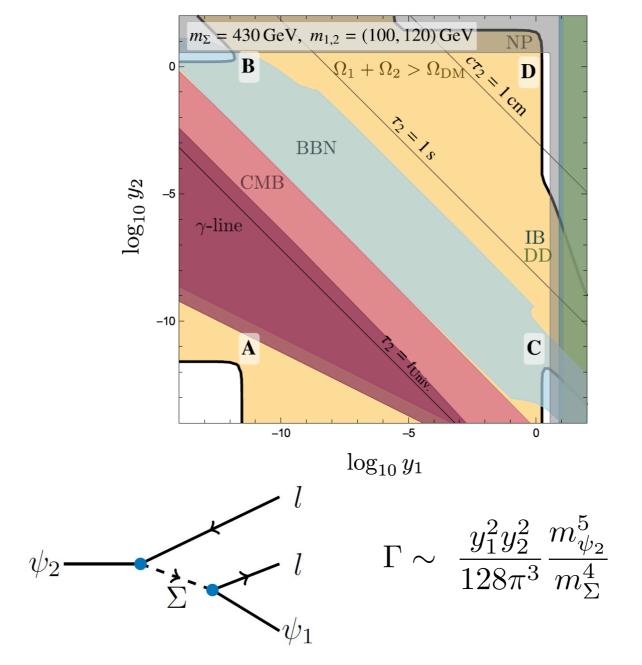
#### A toy model of multiflavor DM



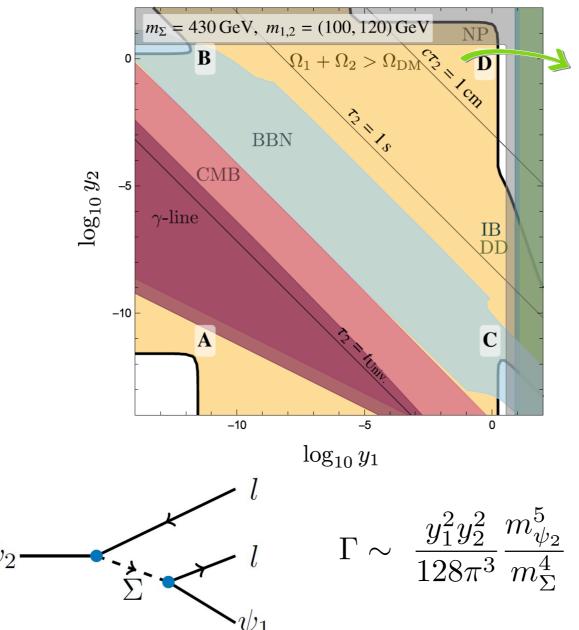
- With one DM flavor, the value of the Yukawa coupling is fixed.
- With two DM flavors, many more possibilities. Larger couplings are compatible with the observed DM abundance → enhanced signals
- Bonus: new signals from  $\psi_2 \to \psi_1 \gamma$ ,  $\psi_2 \to \psi_1 e^+ e^-$ , etc.



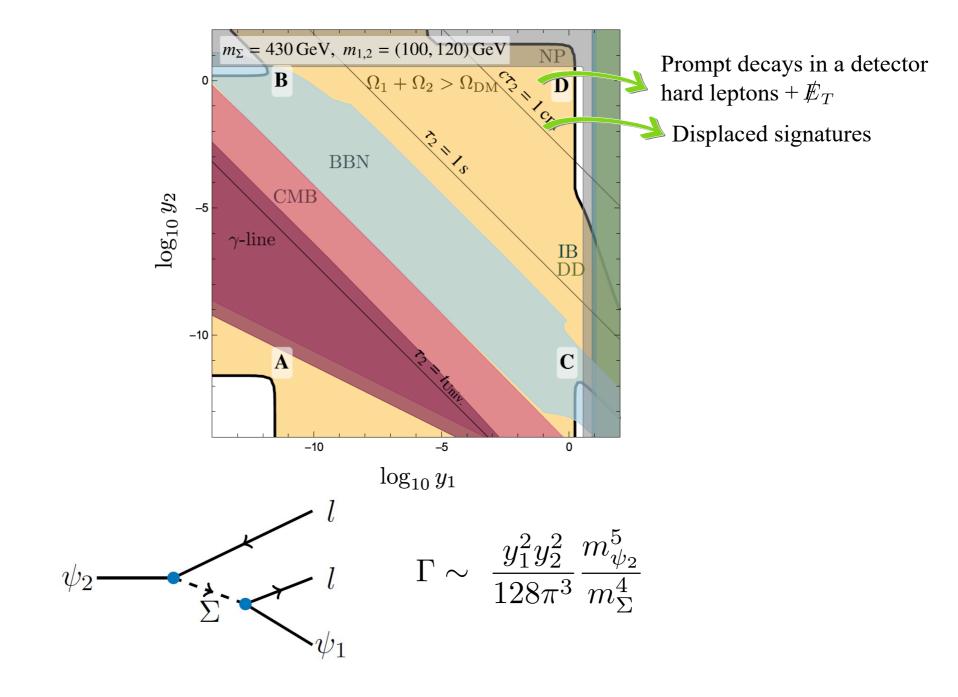


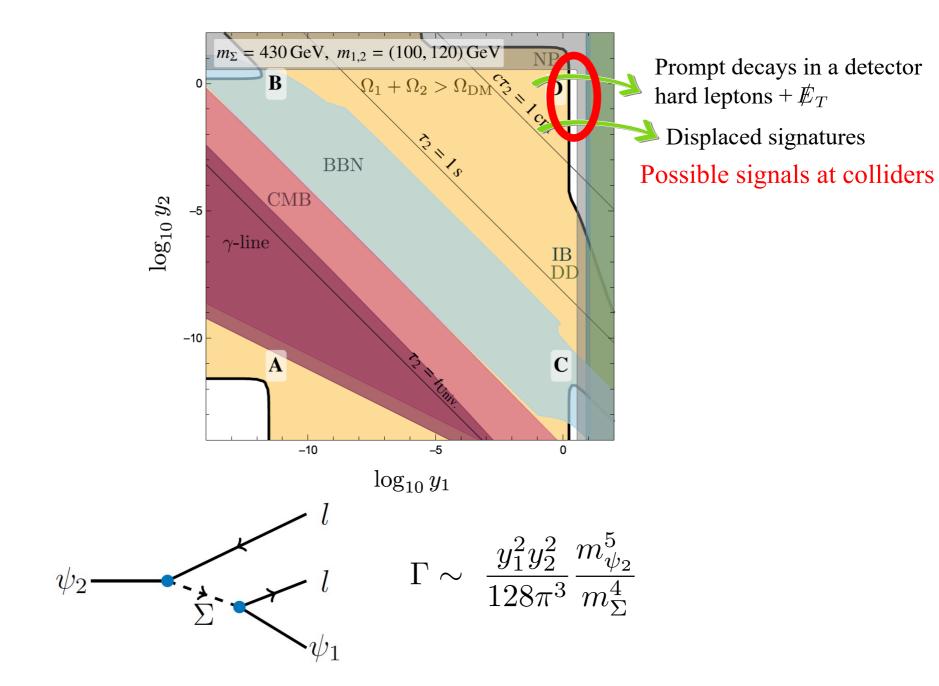


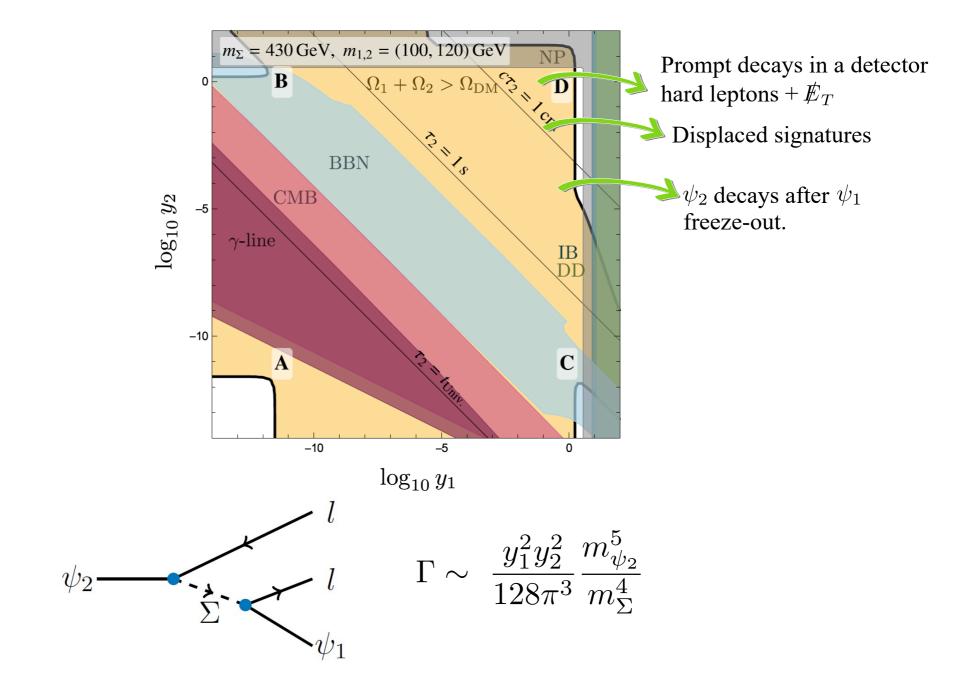
Isocontours of  $\psi_2$  lifetime

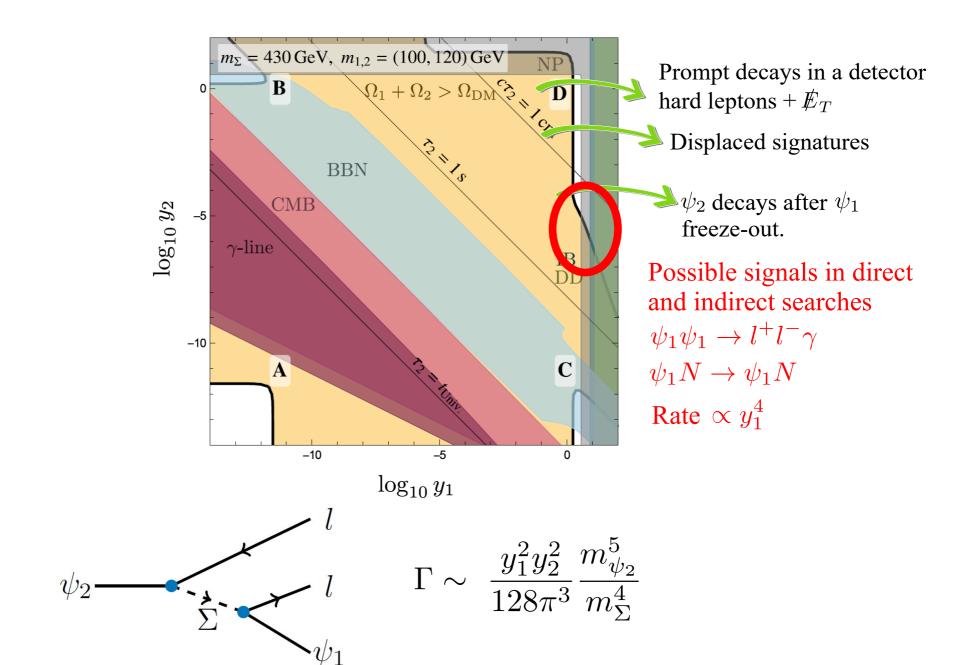


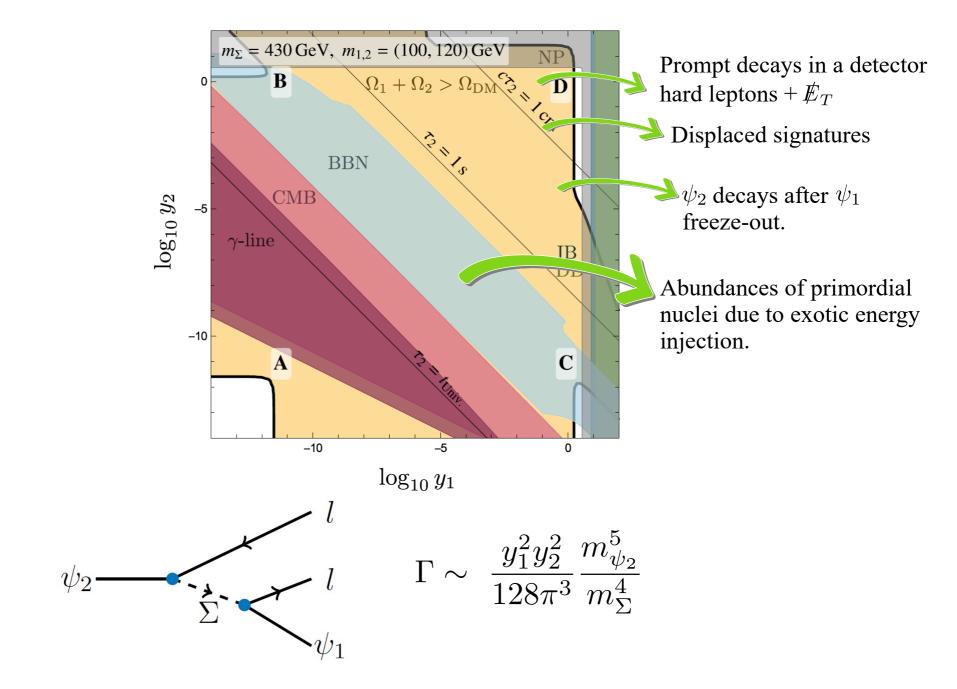
Prompt decays in a detector hard leptons  $+ E_T$ 

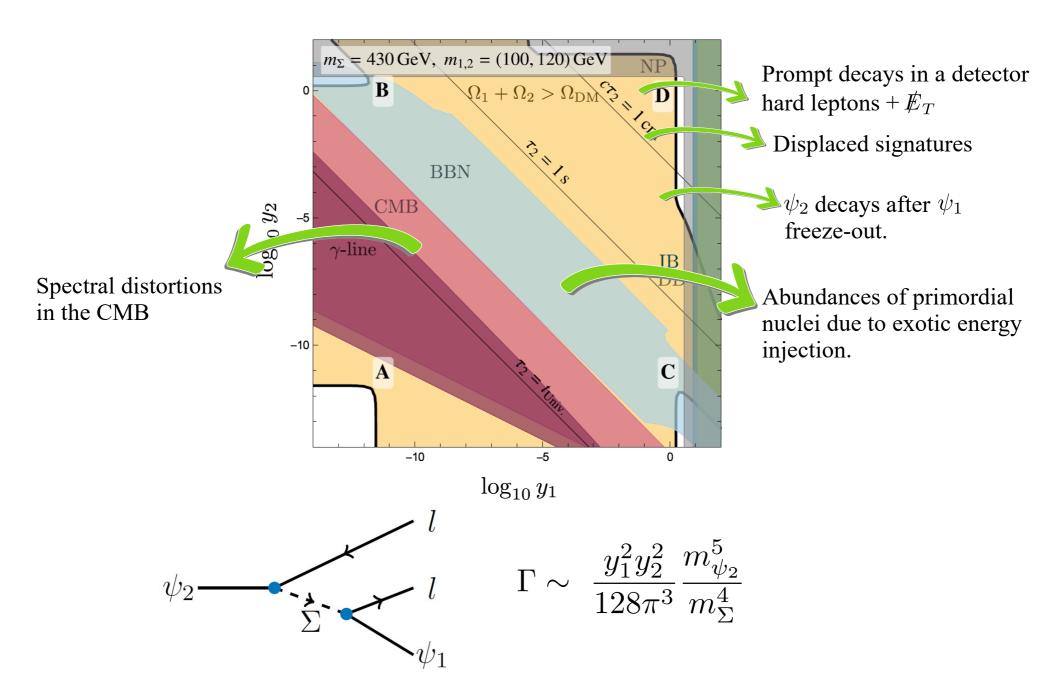


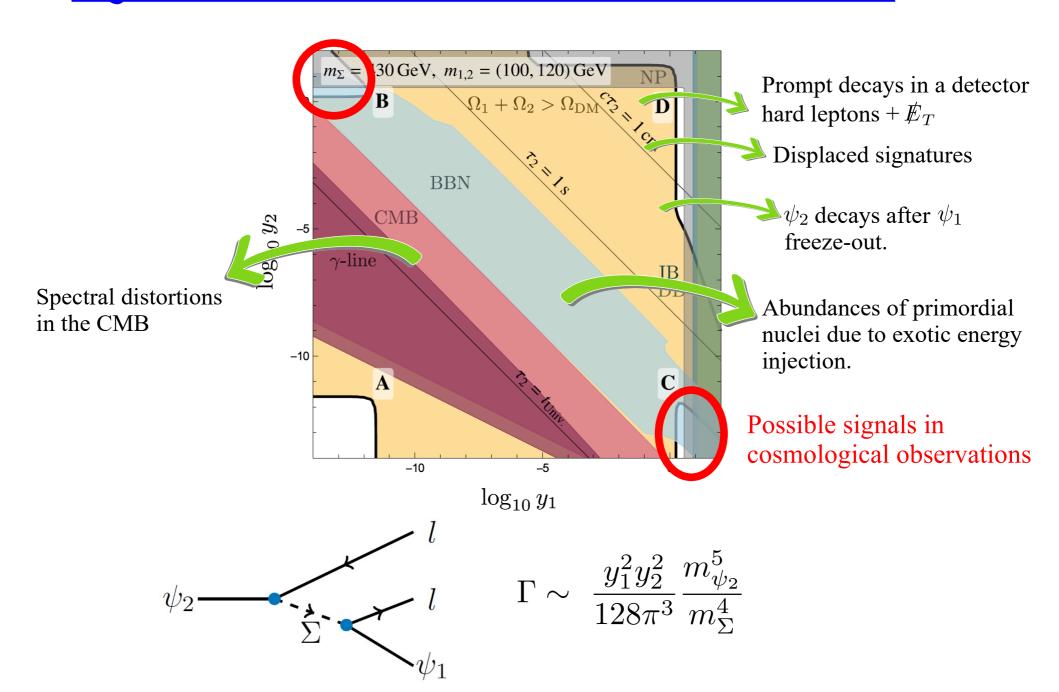


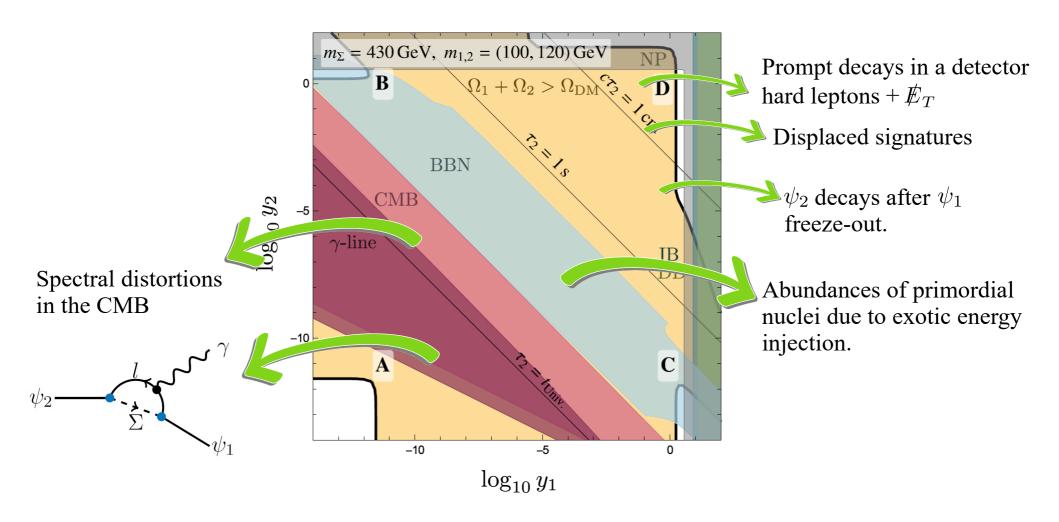


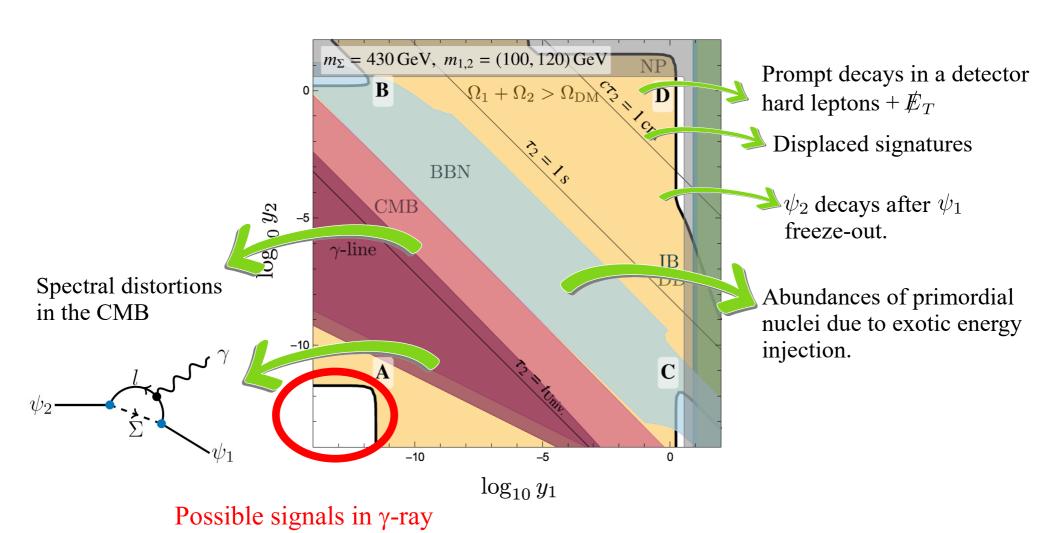




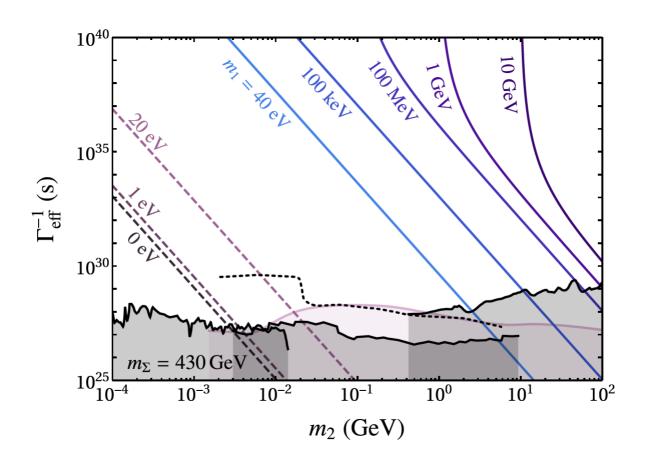








or cosmic ray experiments



### **Conclusions**

- It is not unlikely that the dark matter particle could come in several flavors/generations.
- Even if the heavier generations are not present today in our Universe, they could have played a role in the dark matter production (may be also in generating DM asymmetries??)
- The existence of several flavors could enhance the detection rate at experiments, or lead to new signatures.

