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# Constraints on DM interactions using WDs

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With Nicole Bell, Sandra Robles, Giorgio Busoni and Michael Virgato

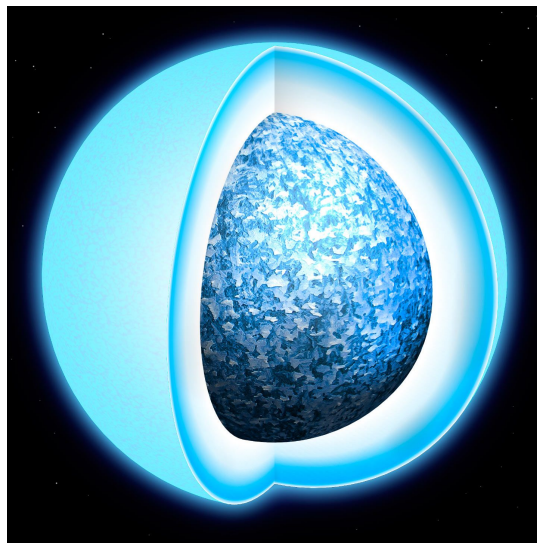
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[arXiv:2104.14367](https://arxiv.org/abs/2104.14367)

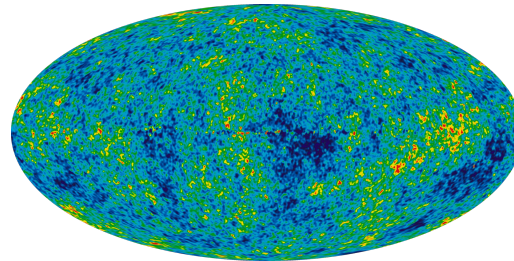
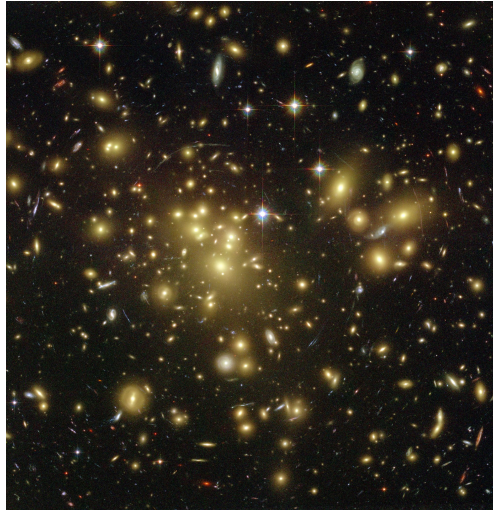
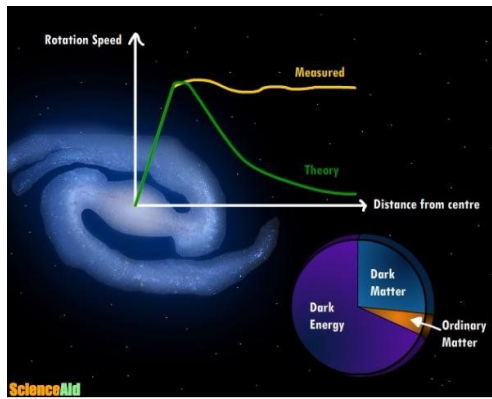
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# Outline



- Introduction
  - Dark matter capture
- Capture in WDs
  - Capture by scattering on ions
    - Finite temperature effects and DM evaporation
- Results
- Summary

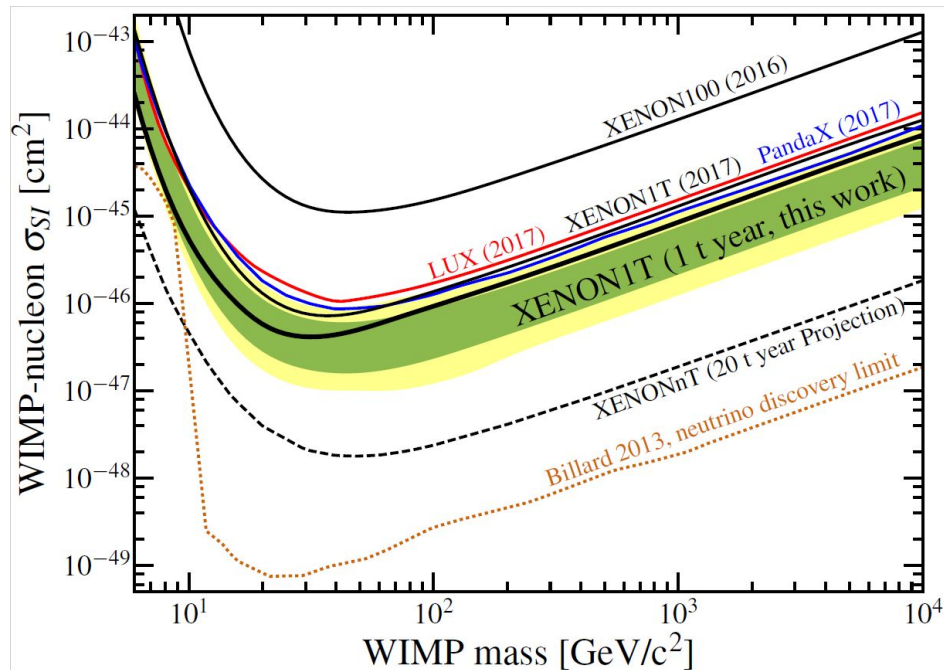


# DM direct detection searches

## Limitations

- Background is still a challenging
- sensitivity is limited by their threshold energy and target mass
- Non-relativistic regime of DM interactions

XENON collaboration (PRL) 2018



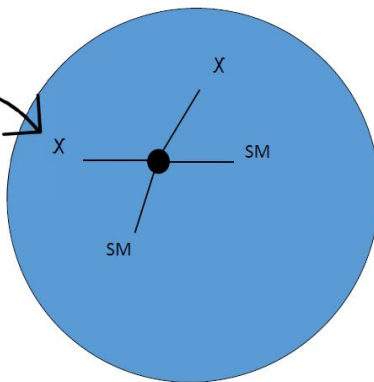
# An alternative approach

## Capture in compact objects

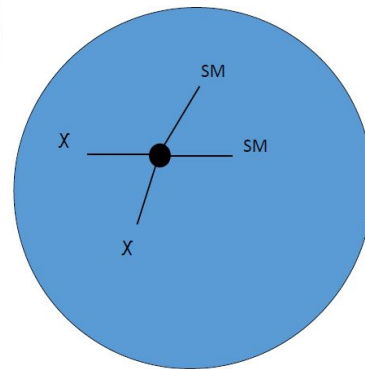
DM can accumulate in the core of these objects in considerable amounts

Complementary to direct detection experiments.

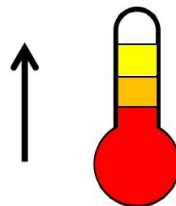
- Velocity (momentum) dependent interactions.
- Sub-Gev regime



1. Capture



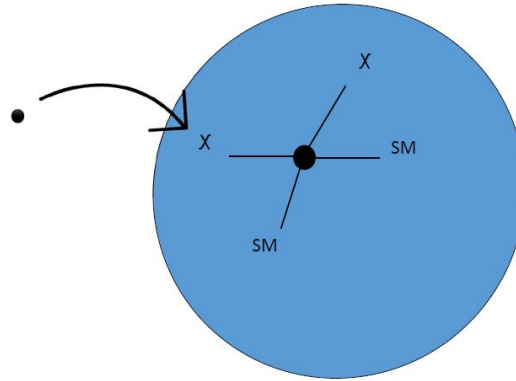
2. Annihilation



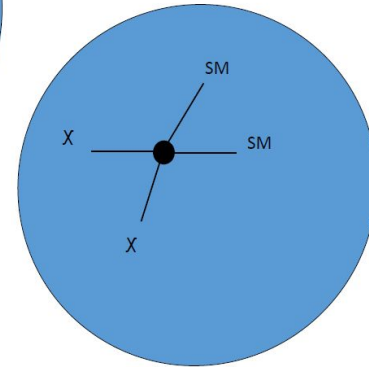
# An alternative approach

## DM scatters off stellar material

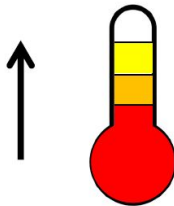
- If loses enough energy it becomes gravitationally bound to the star
- Accumulates and annihilates in the centre of the star.



1. Capture



2. Annihilation



# White Dwarfs

Here we are using M4 WDs

We assume a **Maxwell Boltzmann** (MB) distribution,

$$\rho_\chi = 798 \text{ GeV/cm}^3$$

[arXiv:1001.2737](https://arxiv.org/abs/1001.2737)

$$f_{MB}(u_\chi)$$

DM capture rate due to scattering on ions: assuming WDs are made of only **one element**: Carbon

→ **Inner structure**



# DM interactions

We calculate bounds on the cutoff scale of the dimension 6 EFT SI operators that describe DM interactions with WD targets: Including the response function of carbon

$$\text{cut off scale} \rightarrow \Lambda_q$$

- We decompose the EFT operators in the basis of non-relativistic operators (NR)
- We then use these coefficients, and the nuclear response functions to obtain the DM-ion scattering amplitude

$$|\overline{\mathcal{M}}_T|^2 = \frac{m_T^2}{m_N^2} \sum_{N,N',i,j} C_i^N C_j^{N'} F_{ij}^{NN'}(q_{tr}^2)$$



# DM interactions

We calculate bounds on the cutoff scale of the dimension 6 EFT SI operators that describe DM interactions with WD targets: Including the response function of carbon

Name	Operator	Coupling	$M_N^{\text{NR}}$
D1	$\bar{\chi}\chi \bar{N}N$	$ic_N^S/\Lambda_q^2$	$4\frac{ic_N^S}{\Lambda_q^2}m_\chi m_N \mathcal{O}_1^{\text{NR}}$
D2	$\bar{\chi}\gamma^5\chi \bar{N}N$	$ic_N^S/\Lambda_q^2$	$-4\frac{ic_N^S}{\Lambda_q^2}m_N \mathcal{O}_{11}^{\text{NR}}$
D5	$\bar{\chi}\gamma_\mu\chi \bar{N}\gamma^\mu N$	$c_N^V/\Lambda_q^2$	$4\frac{c_N^V}{\Lambda_q^2}m_\chi m_N \mathcal{O}_1^{\text{NR}}$
D6	$\bar{\chi}\gamma_\mu\gamma^5\chi \bar{N}\gamma^\mu N$	$c_N^V/\Lambda_q^2$	$8\frac{c_N^V}{\Lambda_q^2}(m_\chi m_N \mathcal{O}_8^{\text{NR}} + m_\chi \mathcal{O}_9^{\text{NR}})$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi \bar{N}\sigma^{\mu\nu} N$	$ic_N^T/\Lambda_q^2$	$8\frac{ic_N^T}{\Lambda_q^2}(m_\chi \mathcal{O}_{11}^{\text{NR}} - m_N \mathcal{O}_{10}^{\text{NR}} - 4m_\chi m_N \mathcal{O}_{12}^{\text{NR}})$

1501.03729

# Capture rate

- Assuming that the WD is **optically thin** to DM scattering , the capture rate is given by

$$C_{opt,thin} = \frac{\rho_\chi}{m_\chi} \int_0^{R_\star} 4\pi r^2 \int_0^\infty du_\chi \frac{f_{MB}(u_\chi)}{u_\chi} w(r) \Omega^-(w)$$

- An absolute upper limit on the capture rate arises when we assume the maximum capture probability: **Geometric limit**

$$C_{geom} \propto \frac{1}{m_\chi}$$

- We introduce the **Star Opacity** by including the optical factor in the capture rate computation

$$\eta(\tau_\chi) = e^{-\tau_\chi}$$

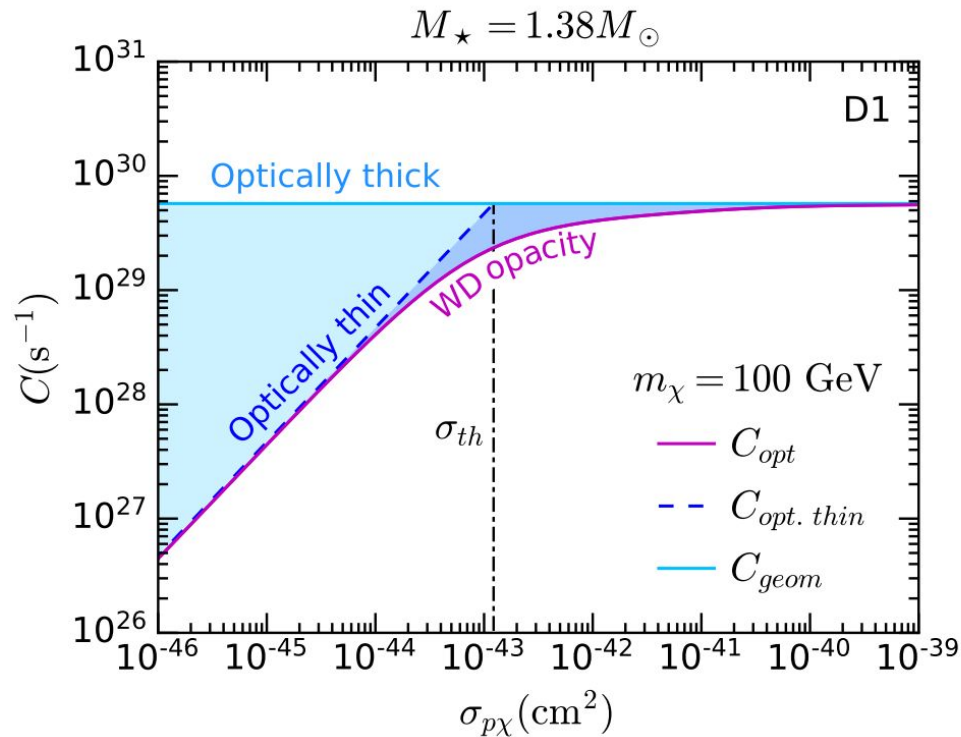
# Capture rate

- Geometric limit:

$$\Omega^-(w) \rightarrow 1$$

- Optically thin limit:

$$\eta(\tau_\chi) \rightarrow 1$$

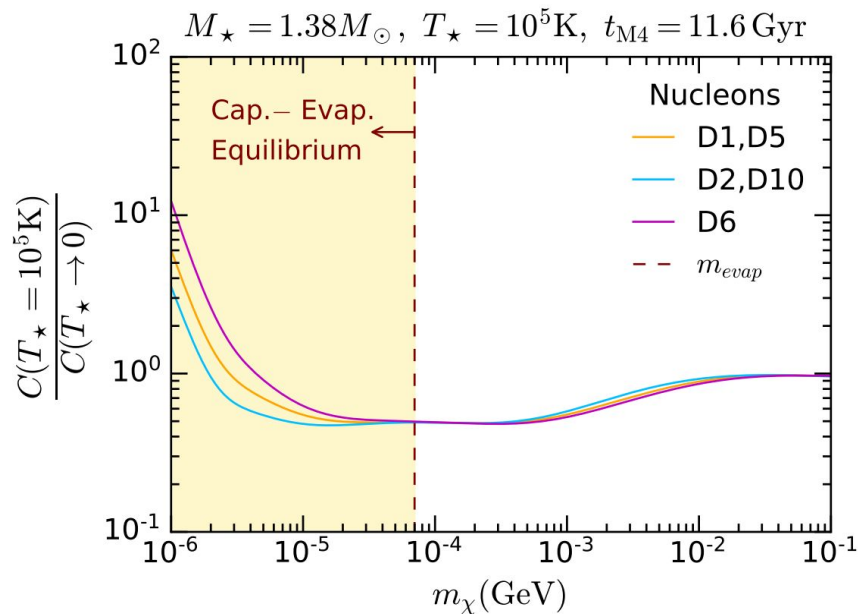


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# Other considerations: Finite temperature

For WDs in M4 finite temperature effects are expected to be important when

$$v_T > v_d$$



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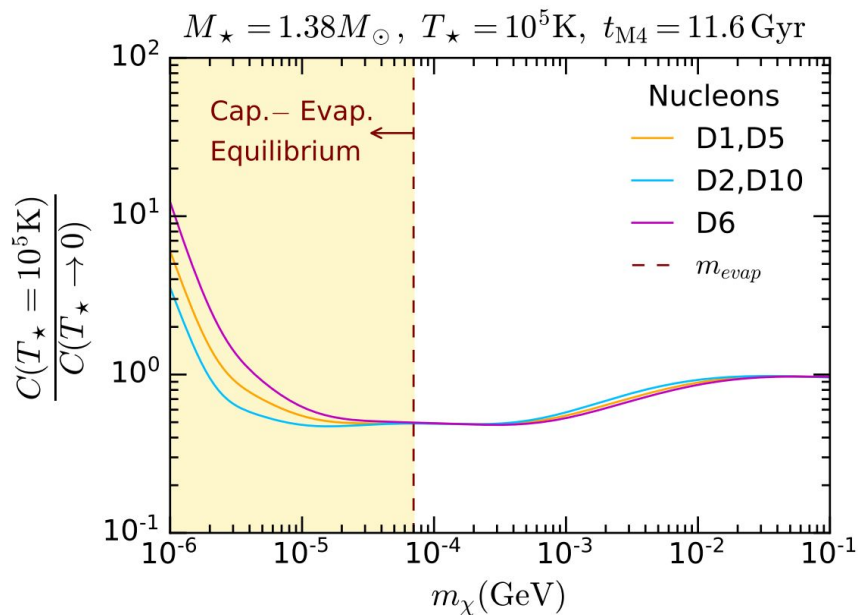
# Other considerations: Evaporation

Accreted DM accumulated in the WD core could also escape from the star: this is the so called **evaporation**

- Final DM velocity greater than the local escape velocity
- we can estimate the evaporation mass  $m_{evap}$  as the DM mass for which

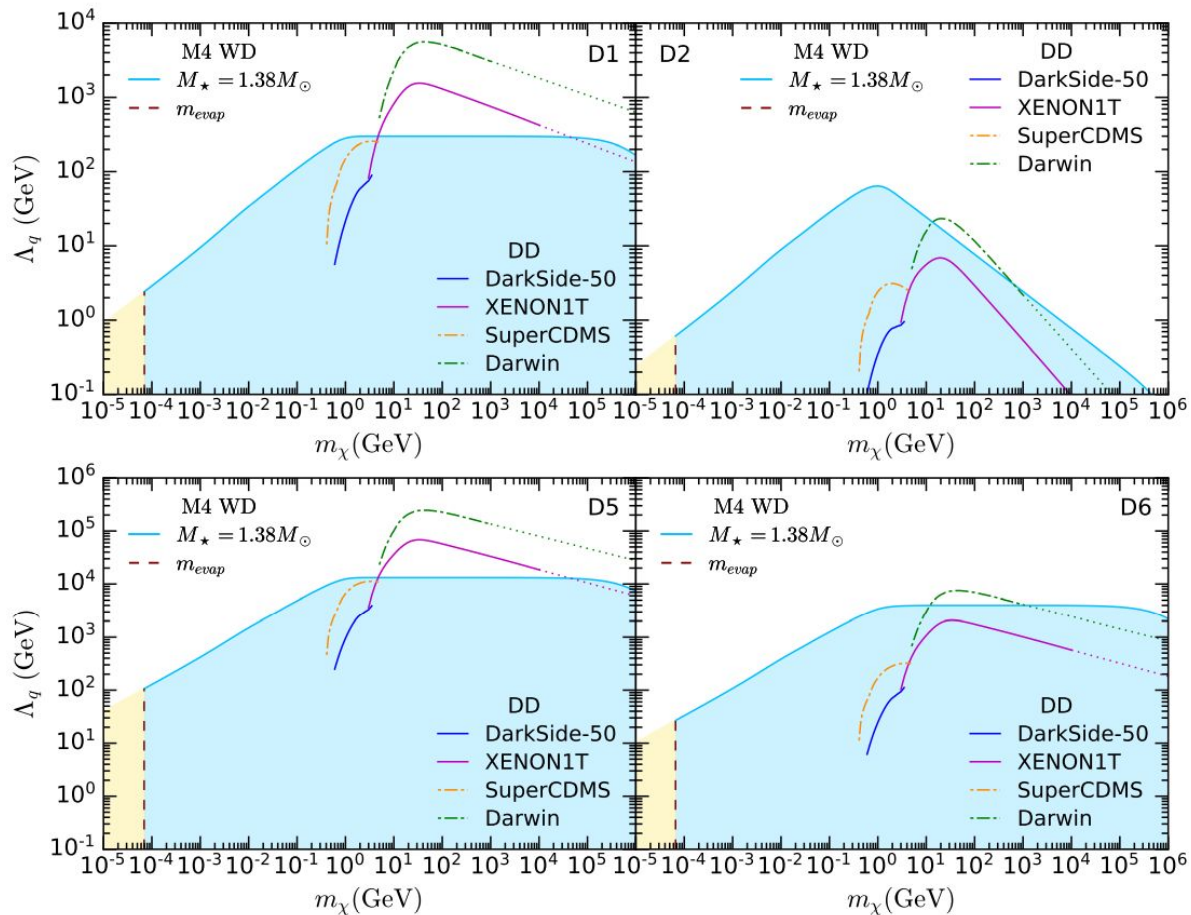
$$E(m_\chi)t_\star \sim 1$$

# Other considerations: Evaporation



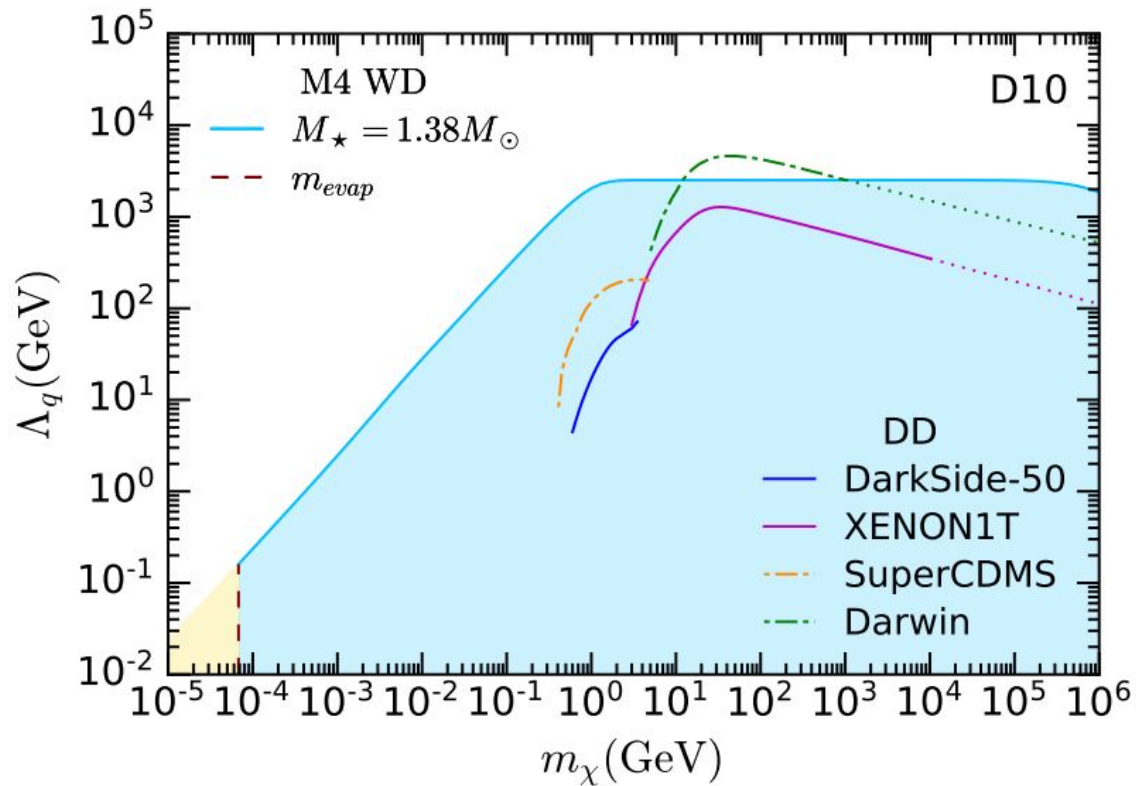
$$m_{evap} \sim 70 \text{ keV}$$

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# Summary

- We have discussed DM capture in compact objects as complementary to DD experiments.
- We have introduced M4 WDs to set constraints on DM interactions
- We have discussed the SI dimension 6 EFT operators.
- The capture rate in
  - Optically thin limit
  - Geometric limit
  - Complete treatment (Star opacity)

# Summary

- Discussed some other considerations in this calculation:
  - Temperature effects
  - Dark matter evaporation
- Limits on the cut-off scale: WDs improve constraints in the **Sub-GeV** mass region

# Thank you