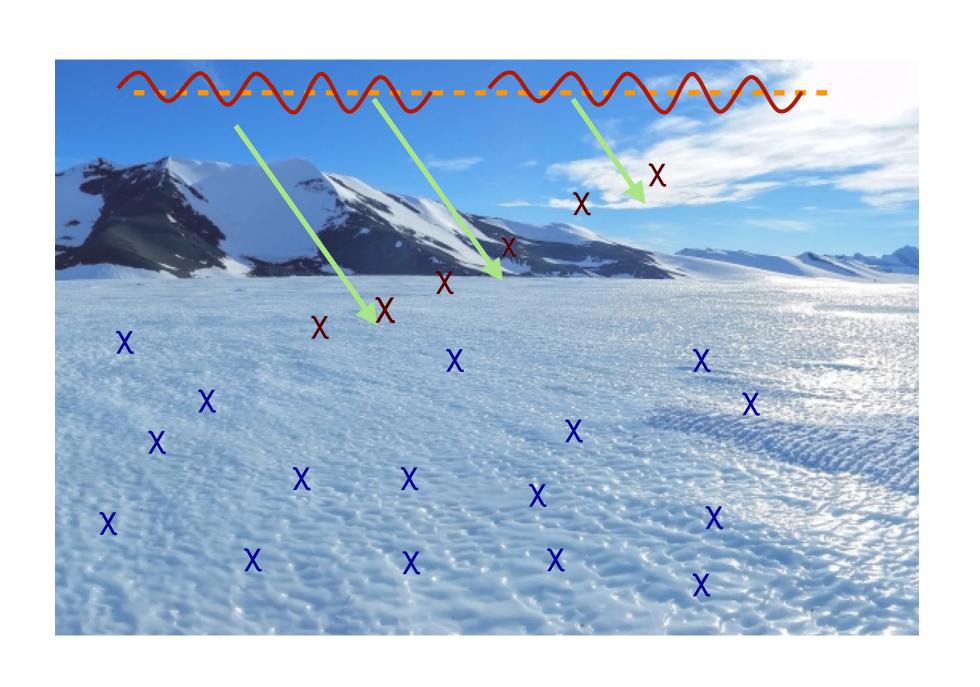
Freeze-in versus Glaciation: Freezing into a thermalized hidden sector

Nicolas Fernandez



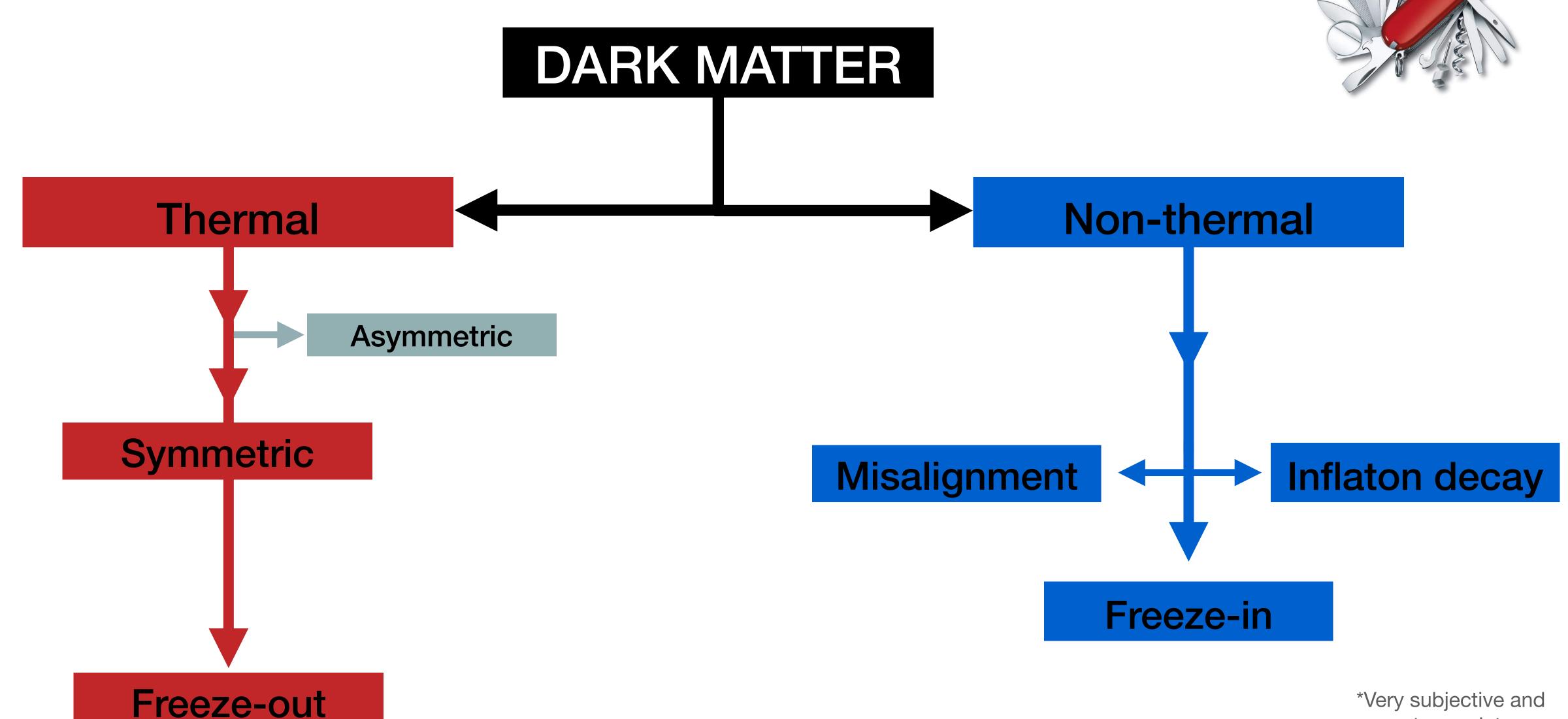
MOCa 2021

June 2021



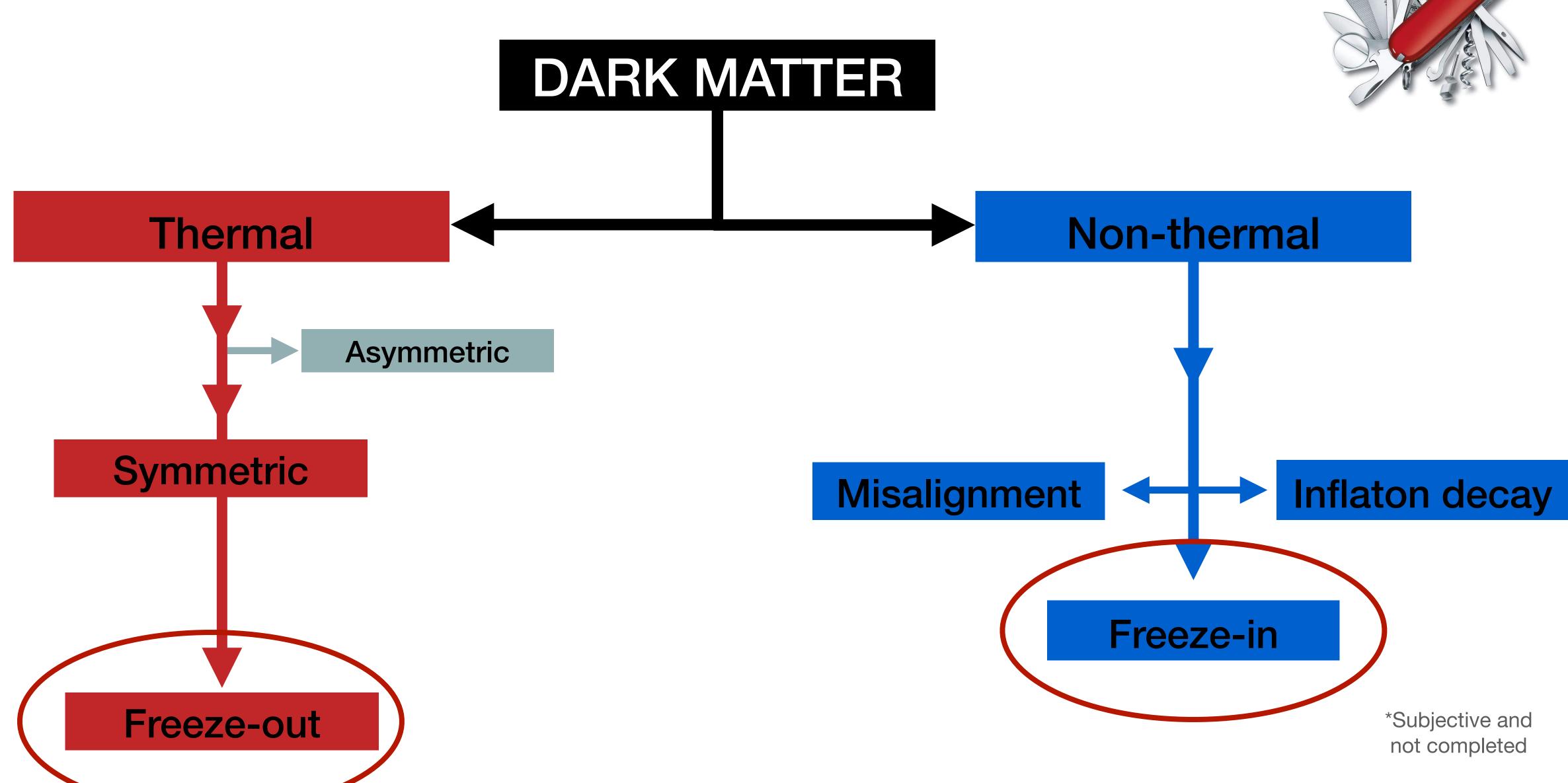
Based on NF, Kahn and Shelton in prep.

Dark Matter Production Flow*



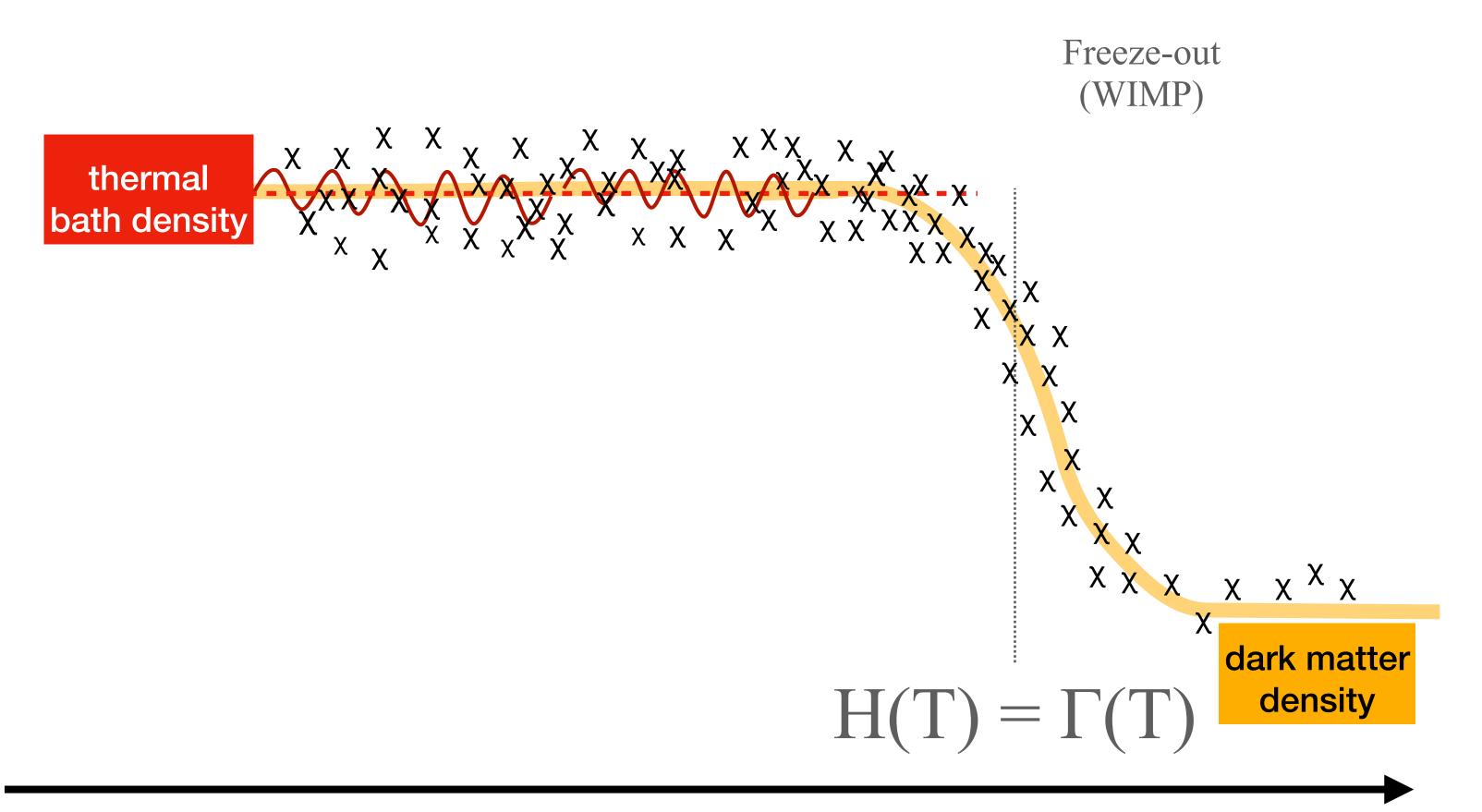
*Very subjective and not complete

Dark Matter Production Flow*

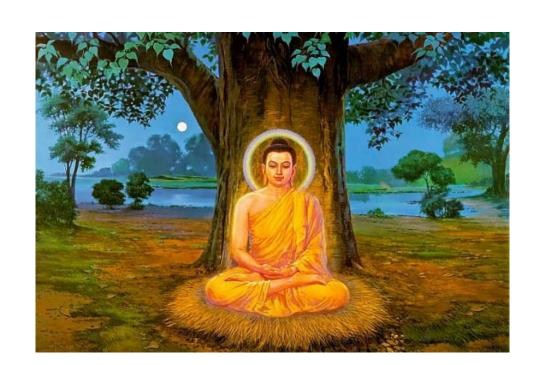


Freeze-out

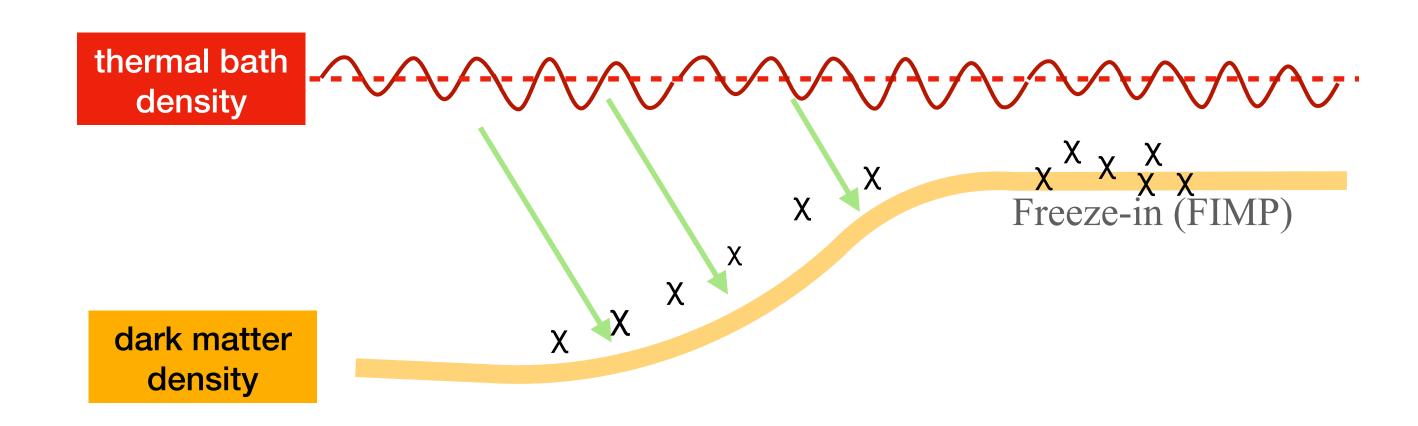
Or the art of getting rid of stuff



- Relic abundance is independent of initial conditions
- Fine with BBN (masses > few MeV)
- Experimentally testable.
 Past ~15 years

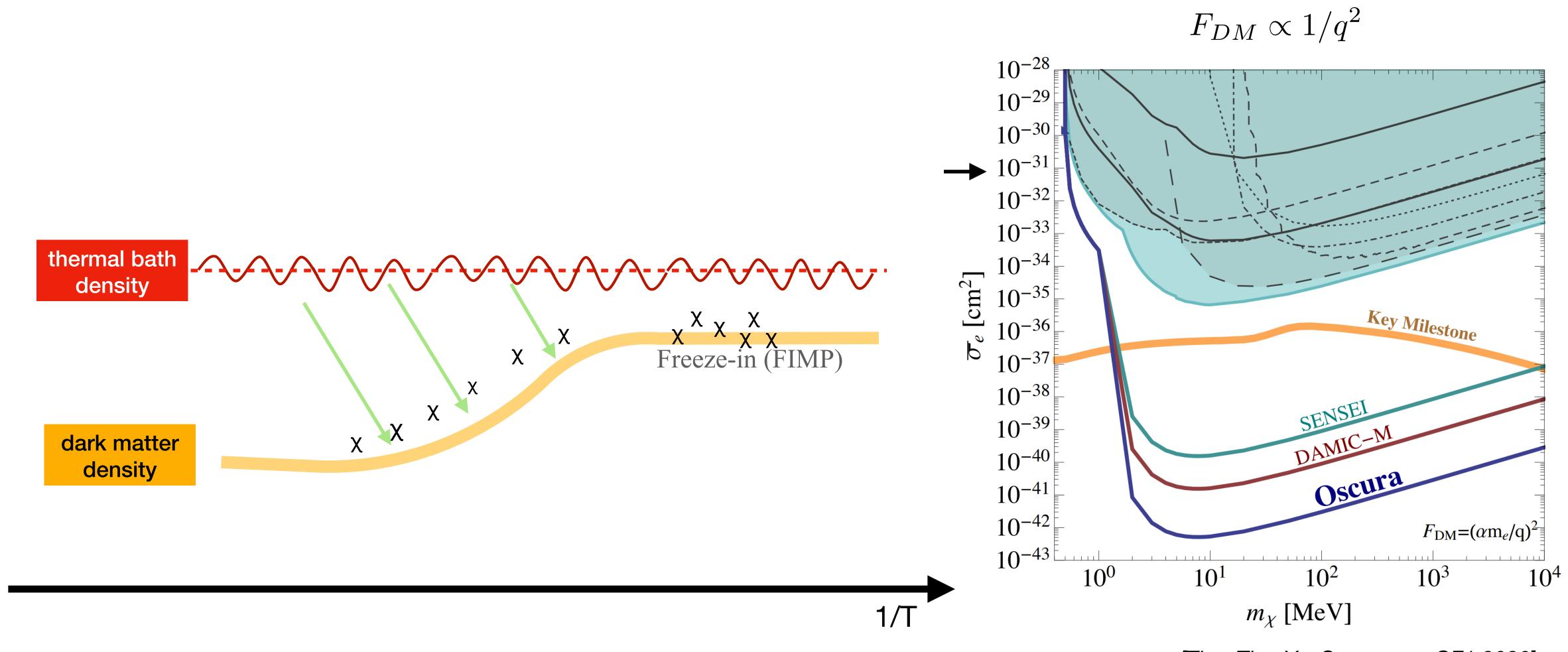


Or the art of getting enough with less

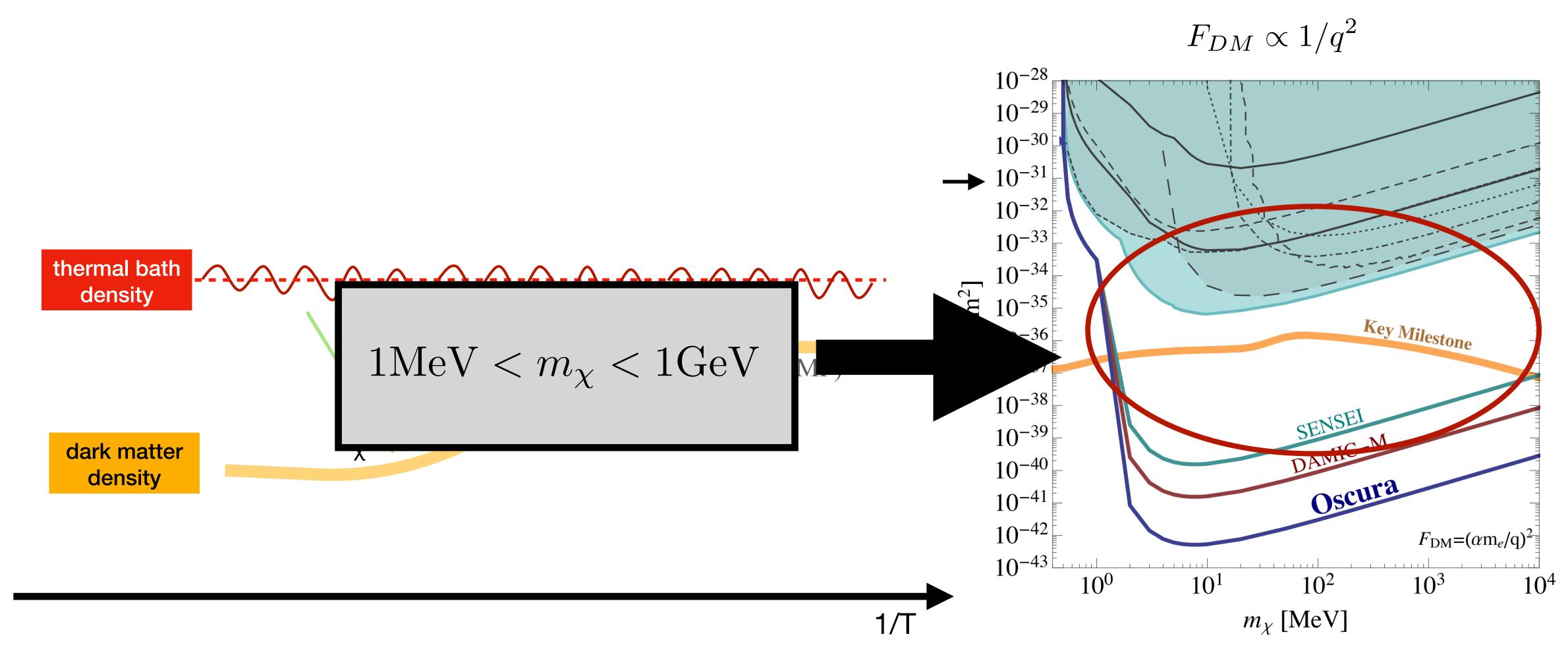


- Relic abundance is independent of initial conditions*
- Fine with BBN and Neff (masses > O(10) keV)
- Experimentally testable soon! Very exciting!

Looking in to the future

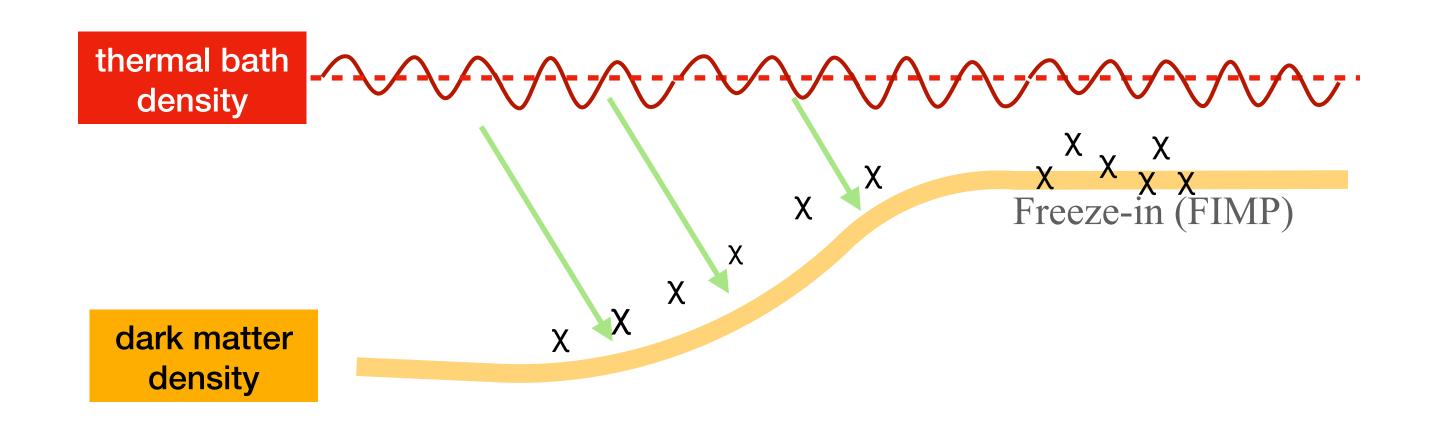


Looking in to the future



[Tien-Tien Yu: Snowmass CF1 2020]

Trying to poke holes



- Relic abundance is independent of initial conditions
- Fine with BBN and Neff (masses > O(10) keV)
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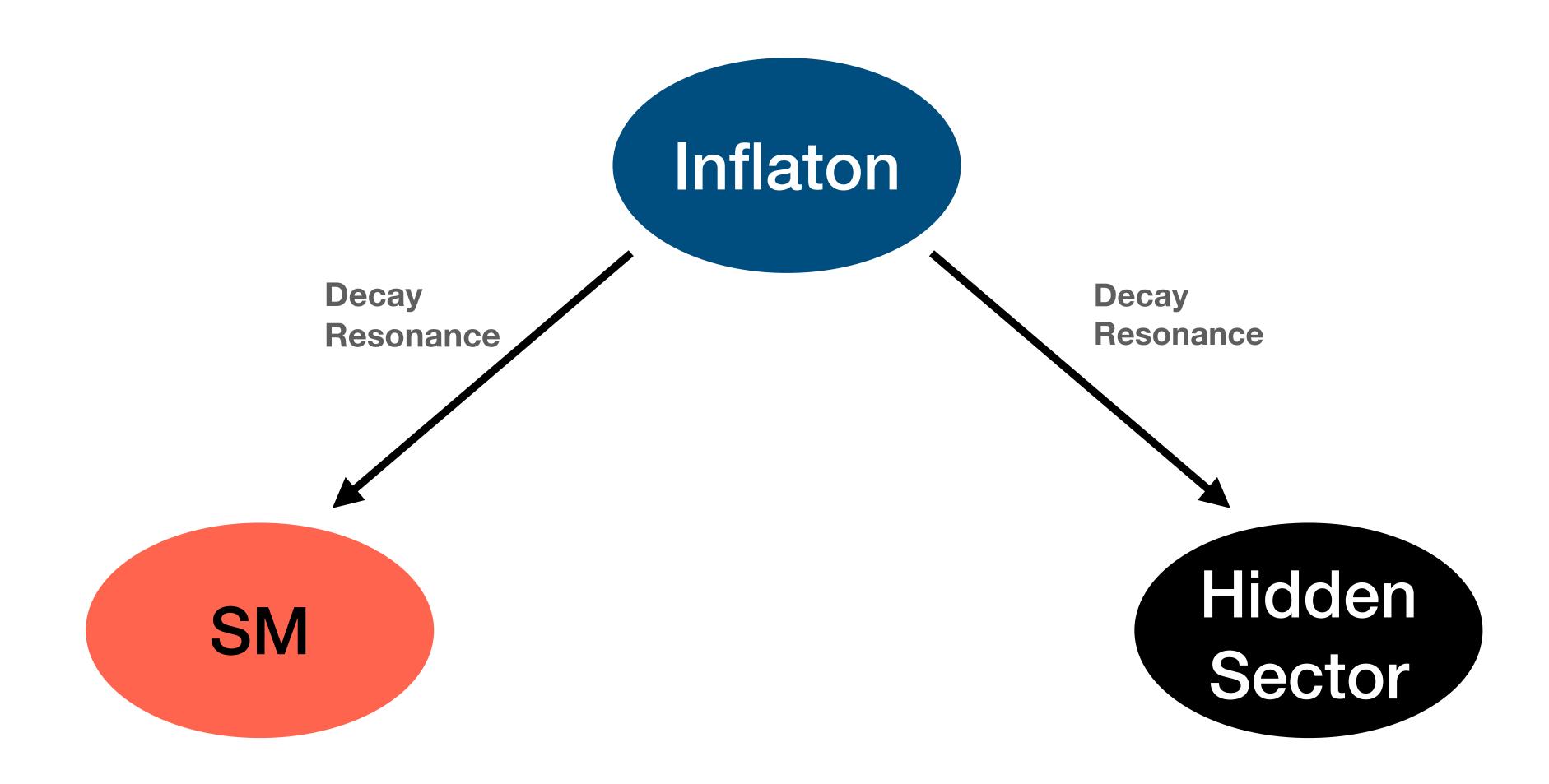
Trying to poke holes

The standard freeze-in paradigm could have a hidden UV sensitivity in that the initial DM population is assumed to be exactly zero.

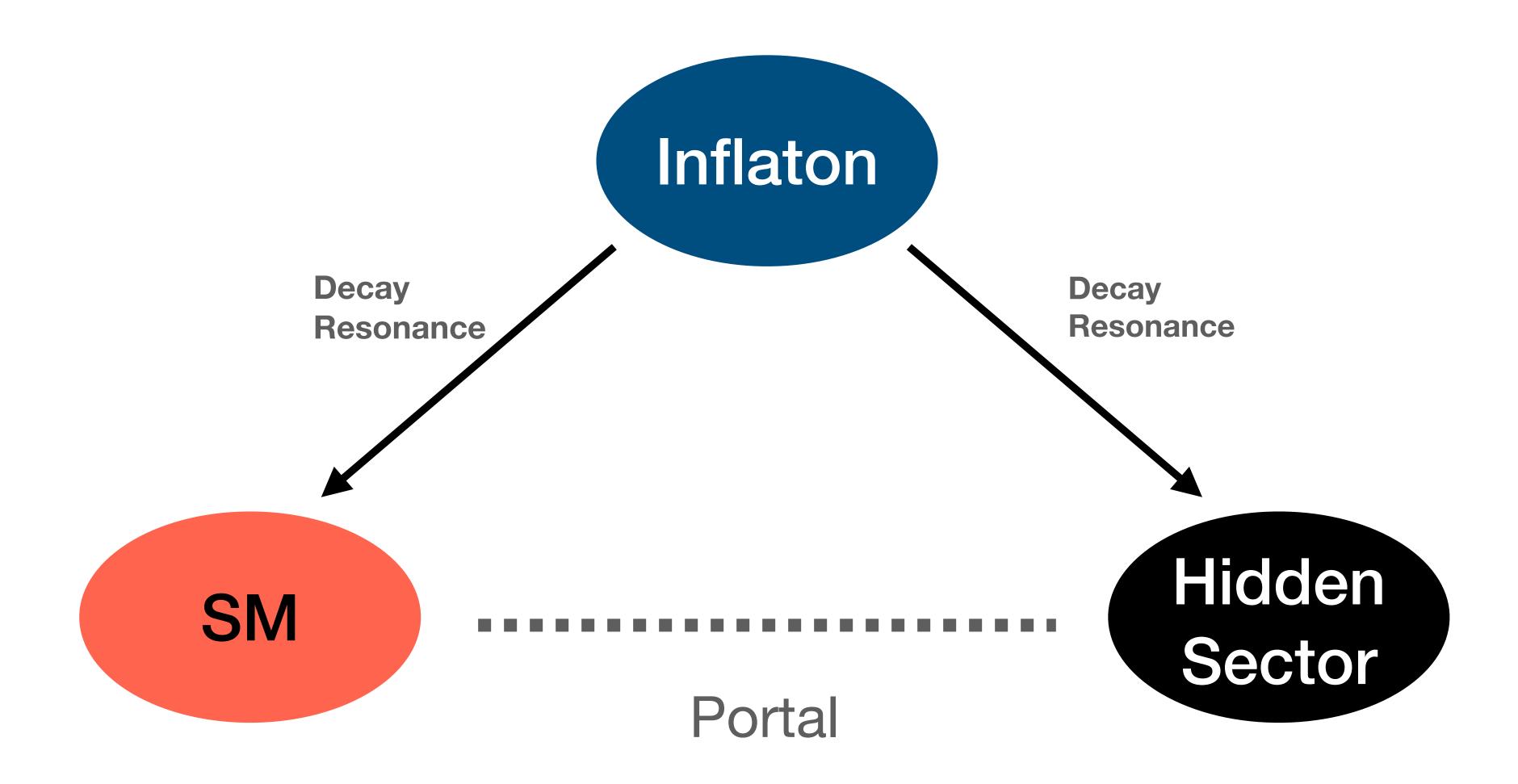
thermal l

dark m

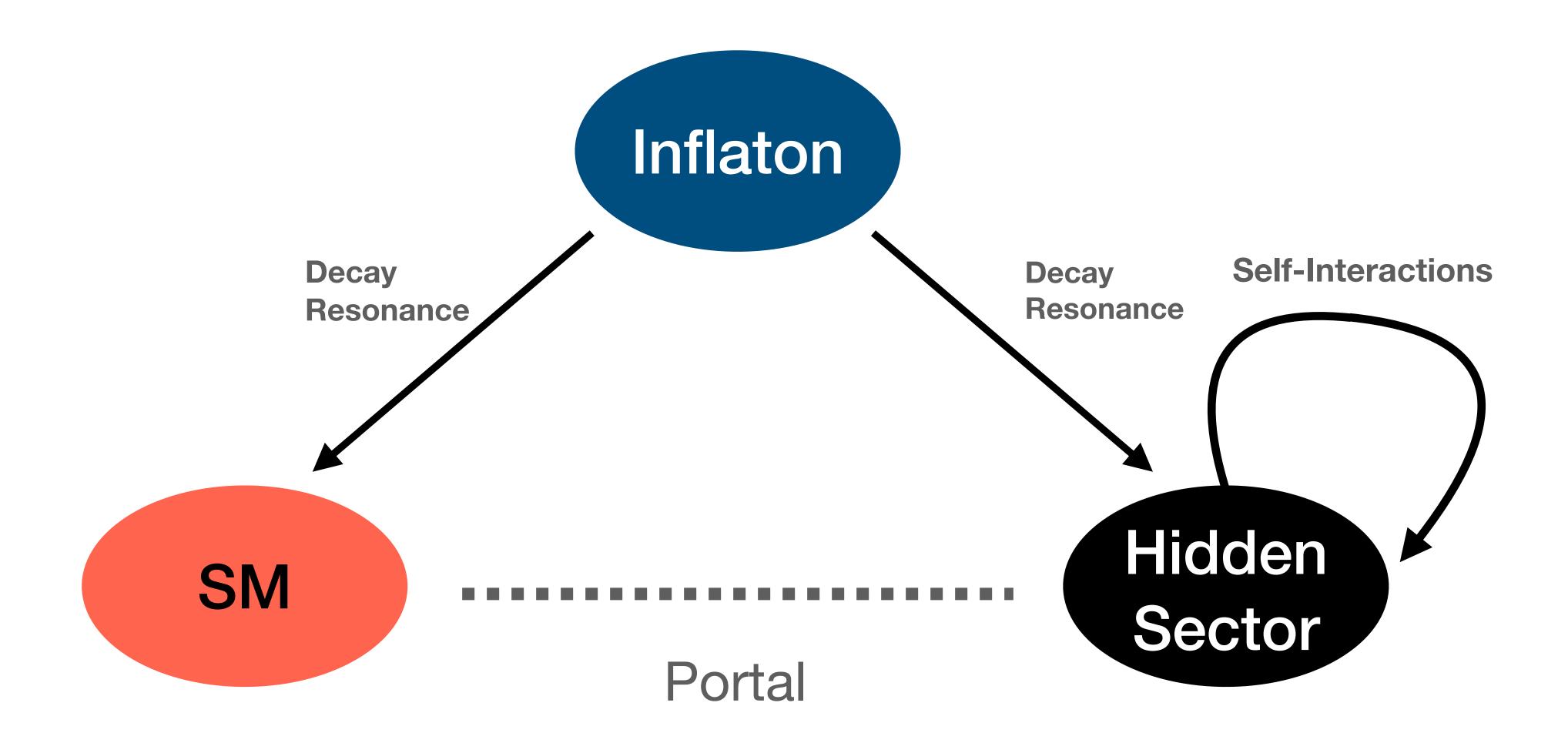
What if...



What if...



What if...



Let us explore an explicit model: Kinetic mixing portal Dark photon

Explicit Model: Kinetic mixing portal Dark photon

$$\mathcal{L} = -\frac{1}{4}\tilde{Z}'_{\mu\nu}\tilde{Z}'^{\mu\nu} \left(-\frac{\epsilon}{2\cos\theta_W} \tilde{Z}'_{\mu\nu}\tilde{B}^{\mu\nu} \right) - \frac{1}{2}m_{Z_D}^2 \tilde{Z}_{D\mu}\tilde{Z}^{\mu}_D + g_{\chi}J^{\mu}_D \tilde{Z}_{D\mu} + \bar{\chi}\left(i\gamma^{\mu}\partial_{\mu} - m_{\chi}\right)\chi\,,$$

Rotating away the mixing term

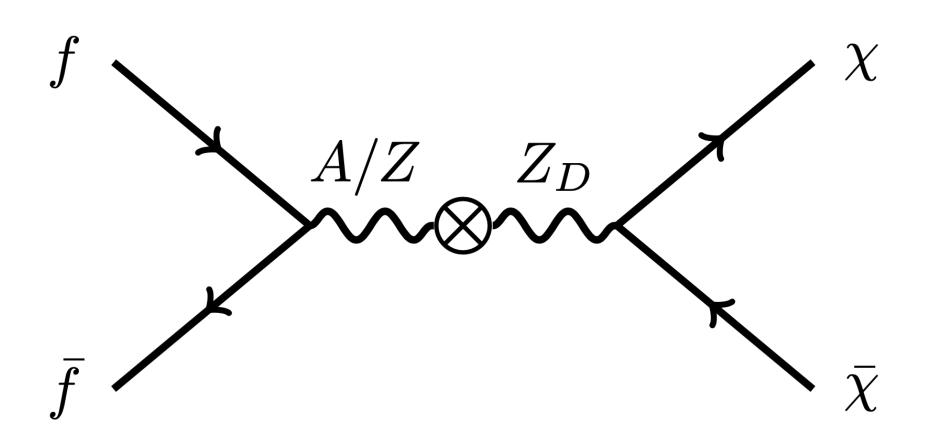
Gauge boson mass eigenstates

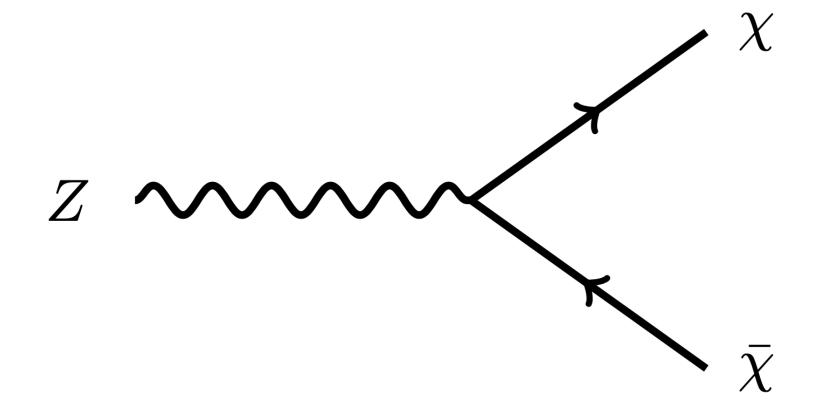
$$m_{Z_D} \ll m_Z$$
 (ultra light mediator)

$$\mathcal{L} \supset -\epsilon e J_{\rm EM}^{\mu} Z_{D\mu} + \epsilon g_{\chi} \tan \theta_W J_D^{\mu} Z_{\mu} + g_{\chi} J_D^{\mu} Z_{D\mu} ,$$

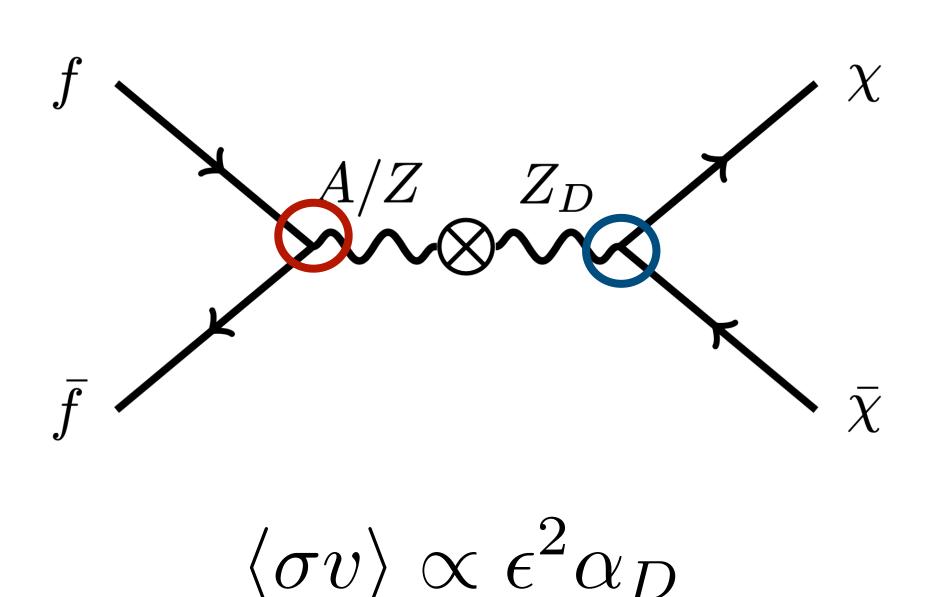
$$\mathcal{L} \supset -\epsilon e J_{\rm EM}^{\mu} Z_{D\mu} + \epsilon g_{\chi} \tan \theta_W J_D^{\mu} Z_{\mu} + g_{\chi} J_D^{\mu} Z_{D\mu} ,$$

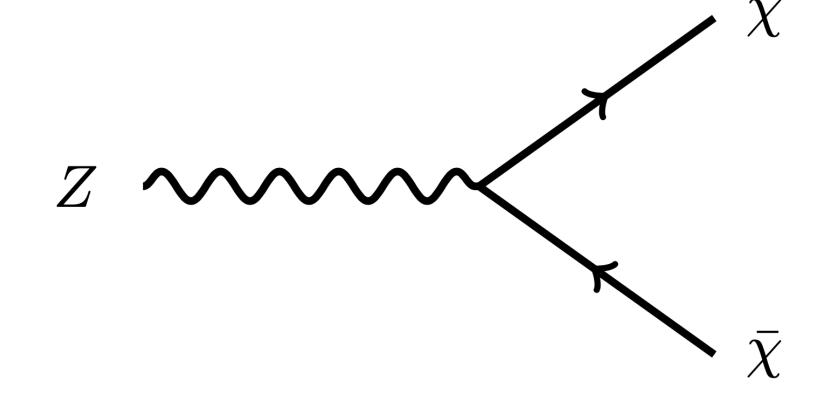
 $1 \mathrm{MeV} < m_{\chi} < 1 \mathrm{GeV}$



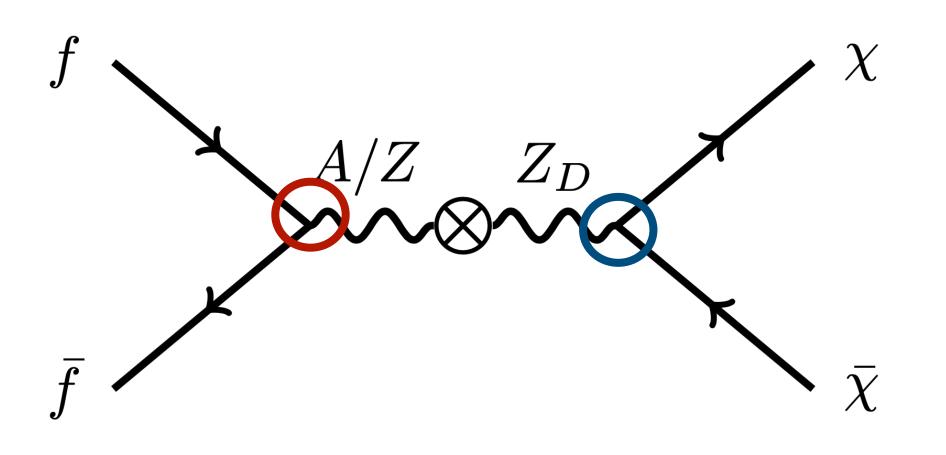


$$\mathcal{L} \supset -\epsilon e J_{\mathrm{EM}}^{\mu} Z_{D\mu} + \epsilon g_{\chi} \tan \theta_W J_D^{\mu} Z_{\mu} + g_{\chi} J_D^{\mu} Z_{D\mu}, \qquad \qquad \alpha_D = \frac{g_{\chi}^2}{4\pi}$$

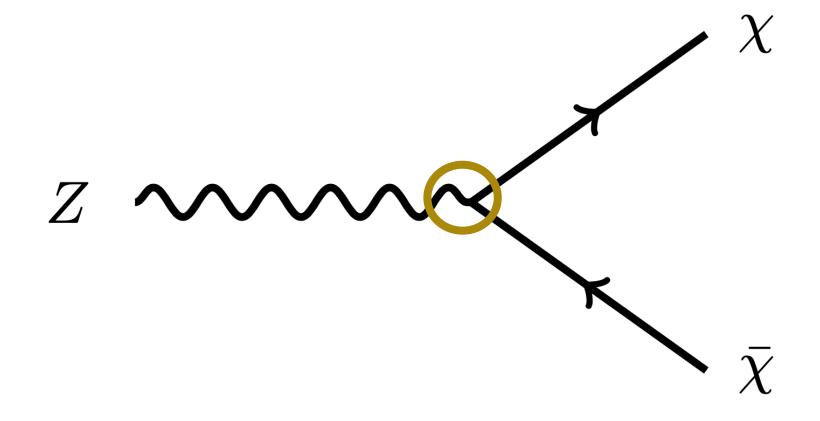




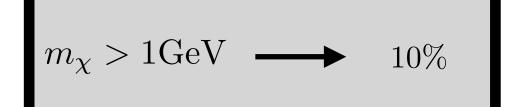
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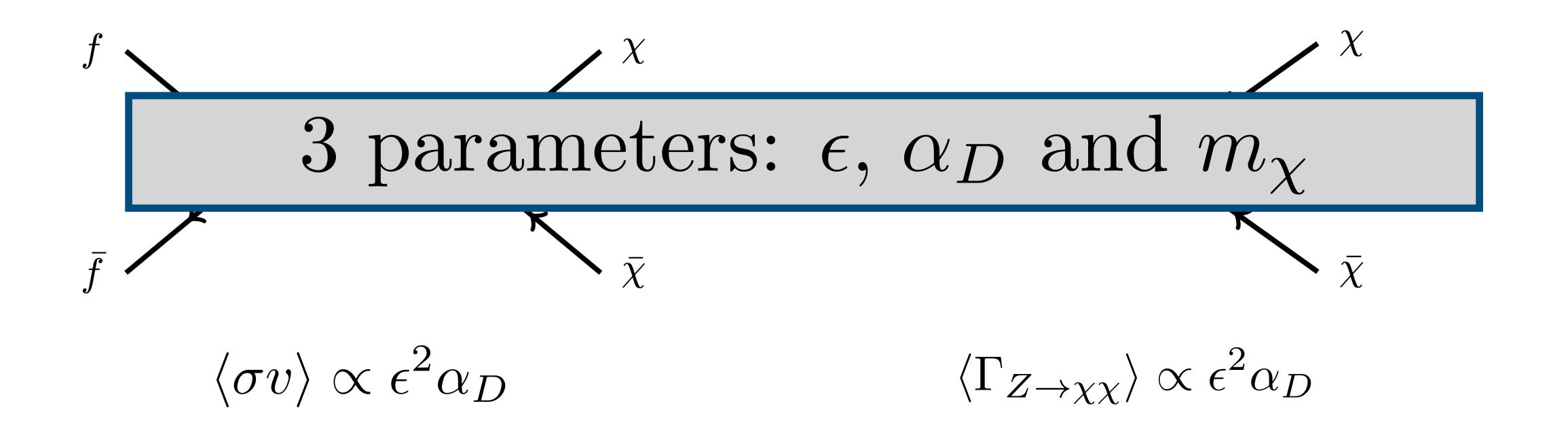
$$\langle \sigma v \rangle \propto \epsilon^2 \alpha_D$$



$$\langle \Gamma_{Z \to \chi \chi} \rangle \propto \epsilon^2 \alpha_D$$



$$\mathcal{L} \supset -\epsilon e J_{\mathrm{EM}}^{\mu} Z_{D\mu} + \epsilon g_{\chi} \tan \theta_W J_D^{\mu} Z_{\mu} + g_{\chi} J_D^{\mu} Z_{D\mu} , \qquad \qquad \alpha_D = \frac{g_{\chi}^2}{4\pi}$$



$$m_{\chi} > 1 \text{GeV} \longrightarrow 10\%$$

$$\mathcal{L} \supset -\epsilon e J_{\mathrm{EM}}^{\mu} Z_{D\mu} + \epsilon g_{\chi} an heta_W J_D^{\mu} Z_{\mu} + g_{\chi} J_D^{\mu} Z_{D\mu} , \qquad \qquad \alpha_D = rac{g_{\chi}^2}{4\pi}$$

f

The relic abundance is given by the combination

1

$$\langle \sigma v \rangle \propto \epsilon^2 \alpha_D$$
 $\langle \Gamma_{Z \to \chi \chi} \rangle \propto \epsilon^2 \alpha_D$

Boltzmann Equation

Number density of DM:

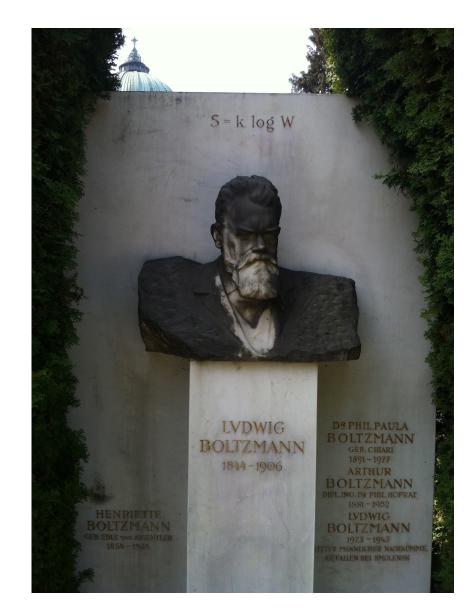
$$\dot{n}_{\chi} + 3Hn_{\chi} = -\langle \sigma v \rangle_{fo}^{\tilde{T}}(n_{\chi}^2 - n_{eq}^2(\tilde{T})) + \langle \sigma v \rangle_{fi}^{T}n_{eq}^2$$

Energy density of the HS:

$$\dot{\rho}_{HS} + 3H(\rho_{HS} + P_{HS}) = \langle \sigma v E \rangle_{fi}^T n_{eq}^2(T)$$



Boltzmann Equation



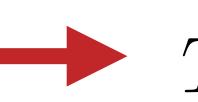
Number density of DM:

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Energy density of the HS:



$$\dot{\rho}_{HS} + 3H(\rho_{HS} + P_{HS}) = \langle \sigma v E \rangle_{fi}^T n_{eq}^2(T)$$



 Instantaneous kinetic equilibration assumed

Boltzmann Equation

$$\chi\chi\leftrightarrow Z_DZ_D$$

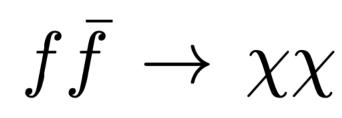
Number density of DM:



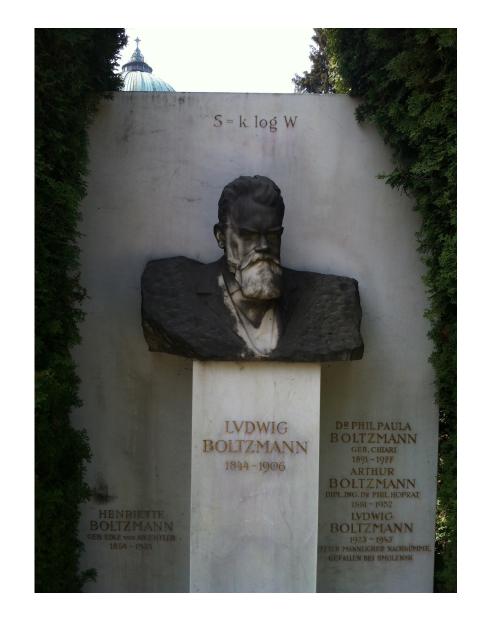
$$\dot{n}_\chi + 3Hn_\chi = -\langle \sigma v \rangle_{fo}^{\tilde{\scriptscriptstyle T}}(n_\chi^2 - n_{eq}^2(\tilde{\scriptscriptstyle T})) + \langle \sigma v \rangle_{fi}^{\scriptscriptstyle T} n_{eq}^2$$



Energy density of the HS:



$$\dot{\rho}_{HS} + 3H(\rho_{HS} + P_{HS}) = \langle \sigma v E \rangle_{fi}^T n_{eq}^2(T)$$



Instantaneous kinetic equilibration assumed

Competing terms

Energy injection:

$$Z o \chi \bar{\chi}$$

$$f\bar{f} \to \chi \bar{\chi}$$

$$\epsilon^2 \alpha_D$$

Kinetic thermalization:

$$\chi\chi \to \chi\chi \qquad \chi\bar{\chi} \to \chi\bar{\chi}$$

$$\chi \bar{\chi} \to \chi \bar{\chi}$$

$$\chi Z_D \to \chi Z_D \quad \chi \bar{\chi} \to Z_D Z_D$$

$$\chi \bar{\chi} \to Z_D Z_D$$

$$\alpha_D^2$$

$$\epsilon$$
 α_D

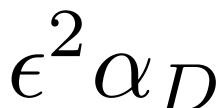
VS

Competing terms

Energy injection:

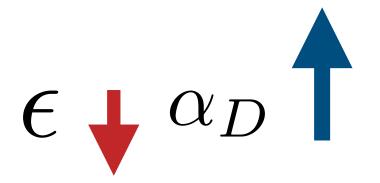
$$Z o \chi \bar{\chi}$$

$$f \bar{f} \to \chi \bar{\chi}$$





Maira Dutra's talk for Z'



VS

Kinetic thermalization:

$$\chi\chi \to \chi\chi$$

$$\chi\chi \to \chi\chi \qquad \chi\bar{\chi} \to \chi\bar{\chi}$$

$$\chi Z_D \to \chi Z_D$$

$$\chi Z_D \to \chi Z_D \quad \chi \bar{\chi} \to Z_D Z_D$$

$$\alpha_D^2$$

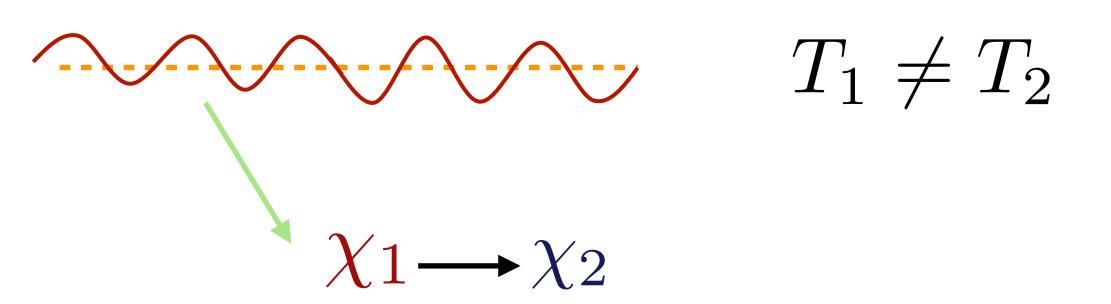


Self-interactions could be important

Let us check two things:

- Instantaneous kinetic equilibration assumption
- Initial conditions

Instantaneous kinetic equilibration of DM



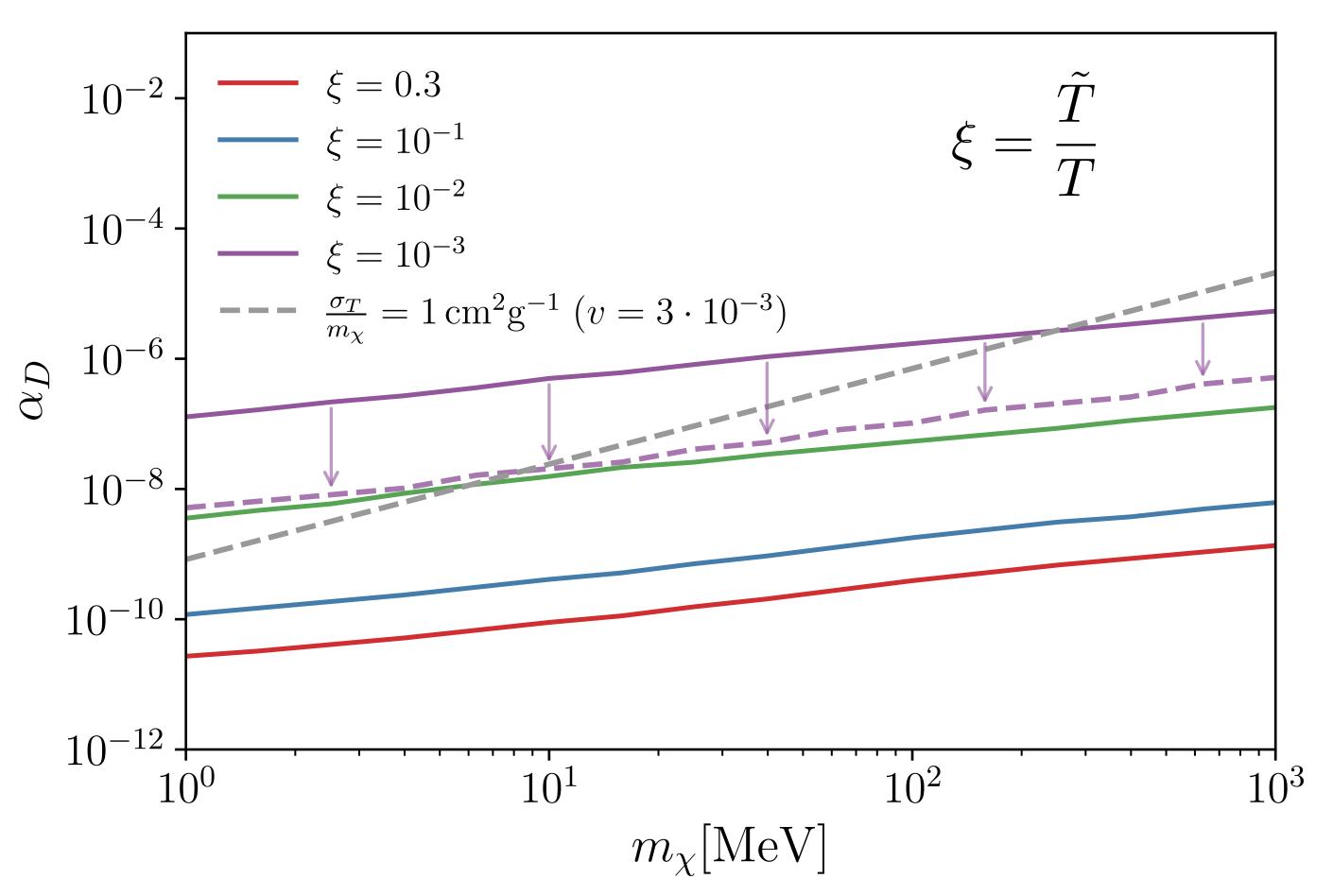
$$T_1 \neq T_2$$

- Two different temperatures in the initial state
- Relativistic

Thermally averaged momentum loss:

$$\Gamma_{p \, \text{loss}} pprox \langle \frac{dp}{dt} \rangle \frac{1}{\langle p \rangle} = \frac{n_{2 \text{eq}}(\tilde{T}) \langle \sigma_T vp \rangle}{\langle p \rangle} \gtrsim H$$

Instantaneous kinetic equilibration of DM Vs DM self-interaction constraints

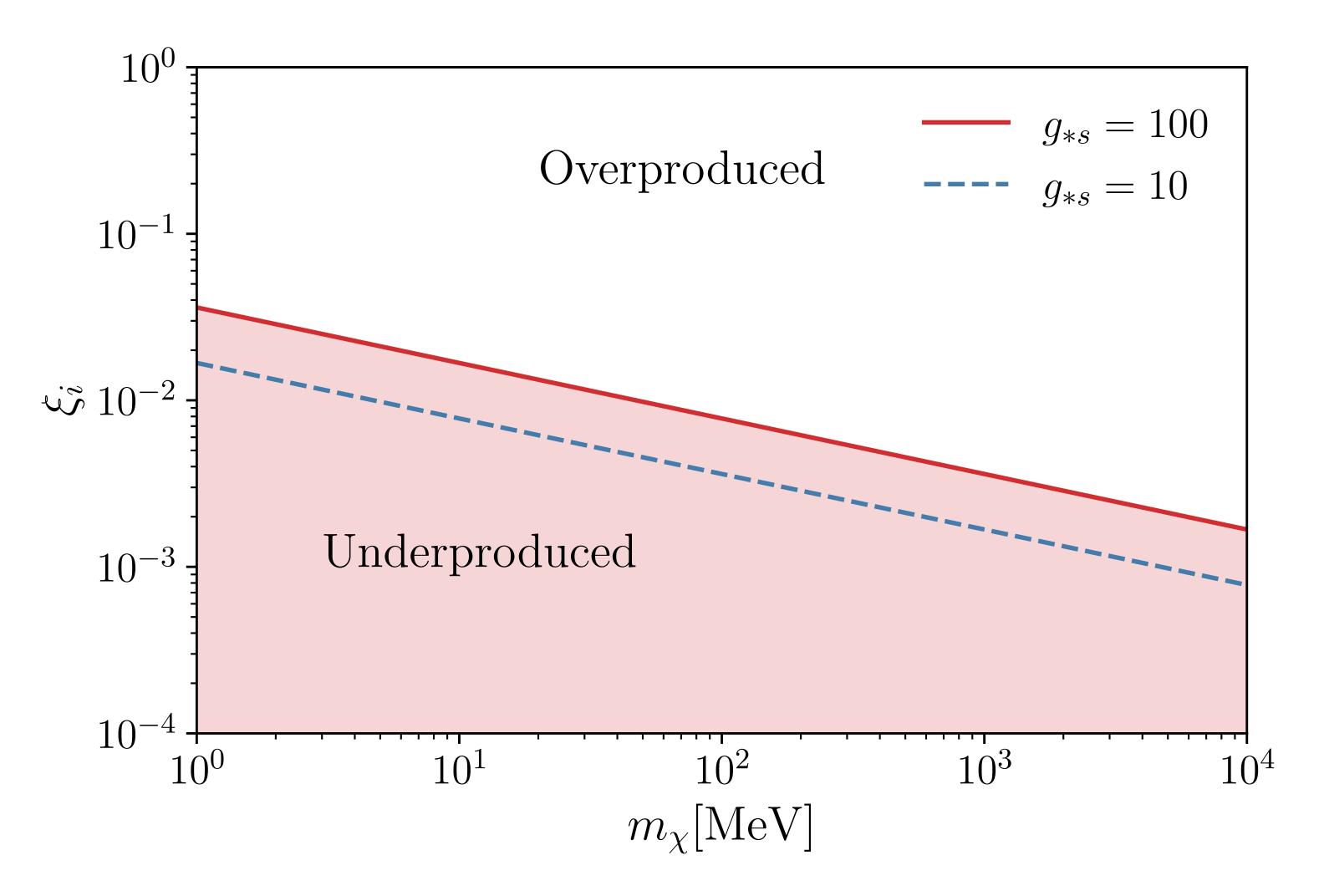


$$\chi_1\chi_2 \rightarrow \chi_1\chi_2$$

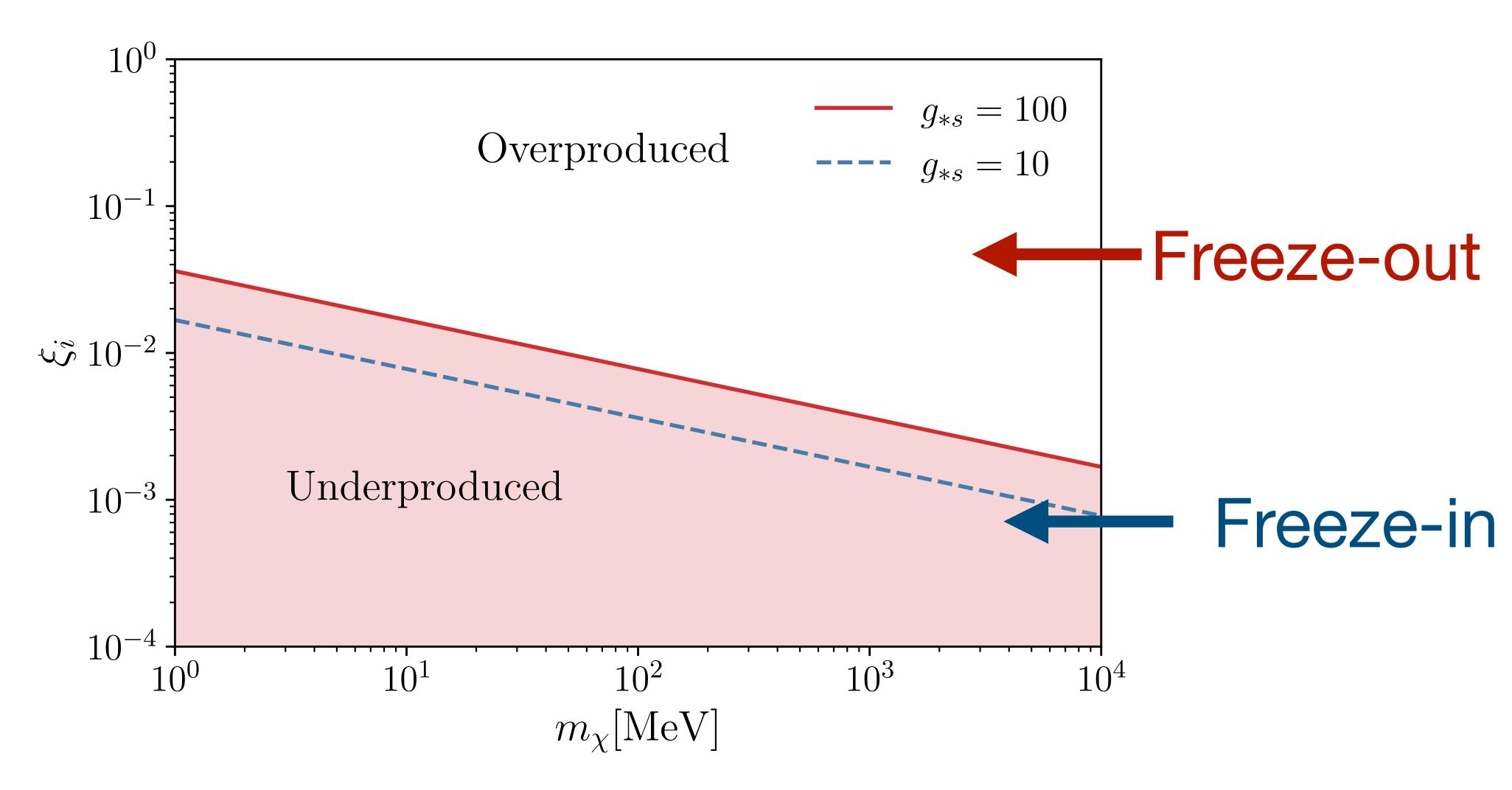
$$\chi_1 \bar{\chi}_2 \rightarrow \chi_1 \bar{\chi}_2$$

$$\chi_1 Z_D \to \chi_1 Z_D$$

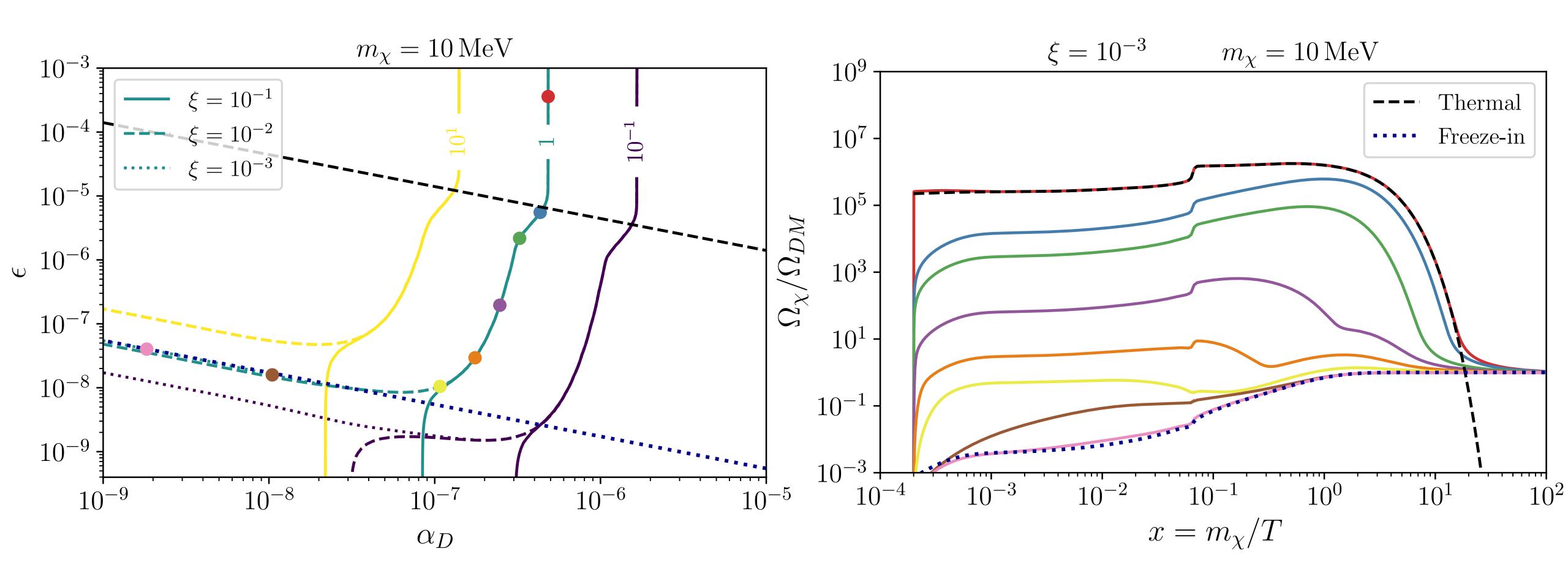
Initial condition DM

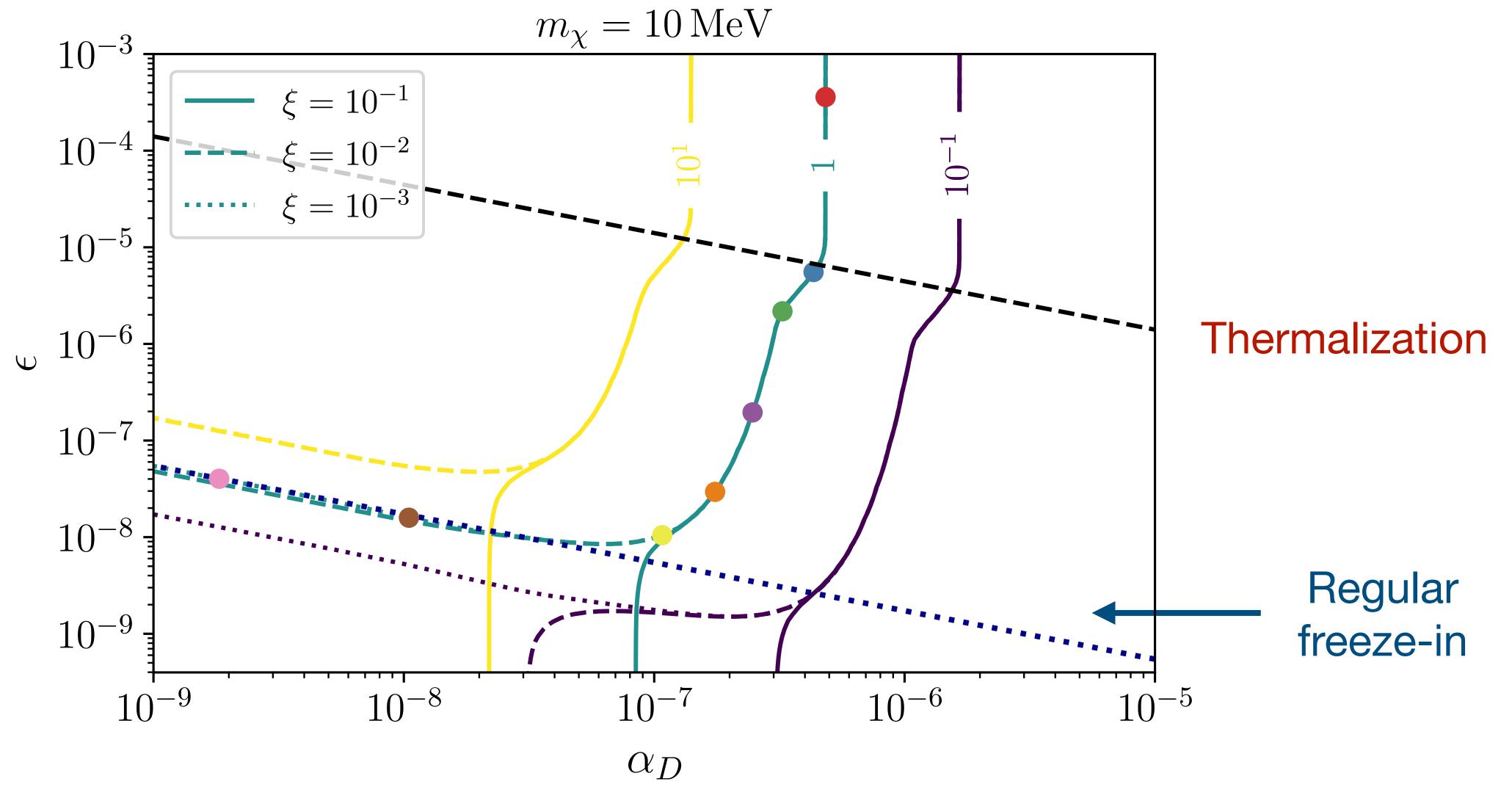


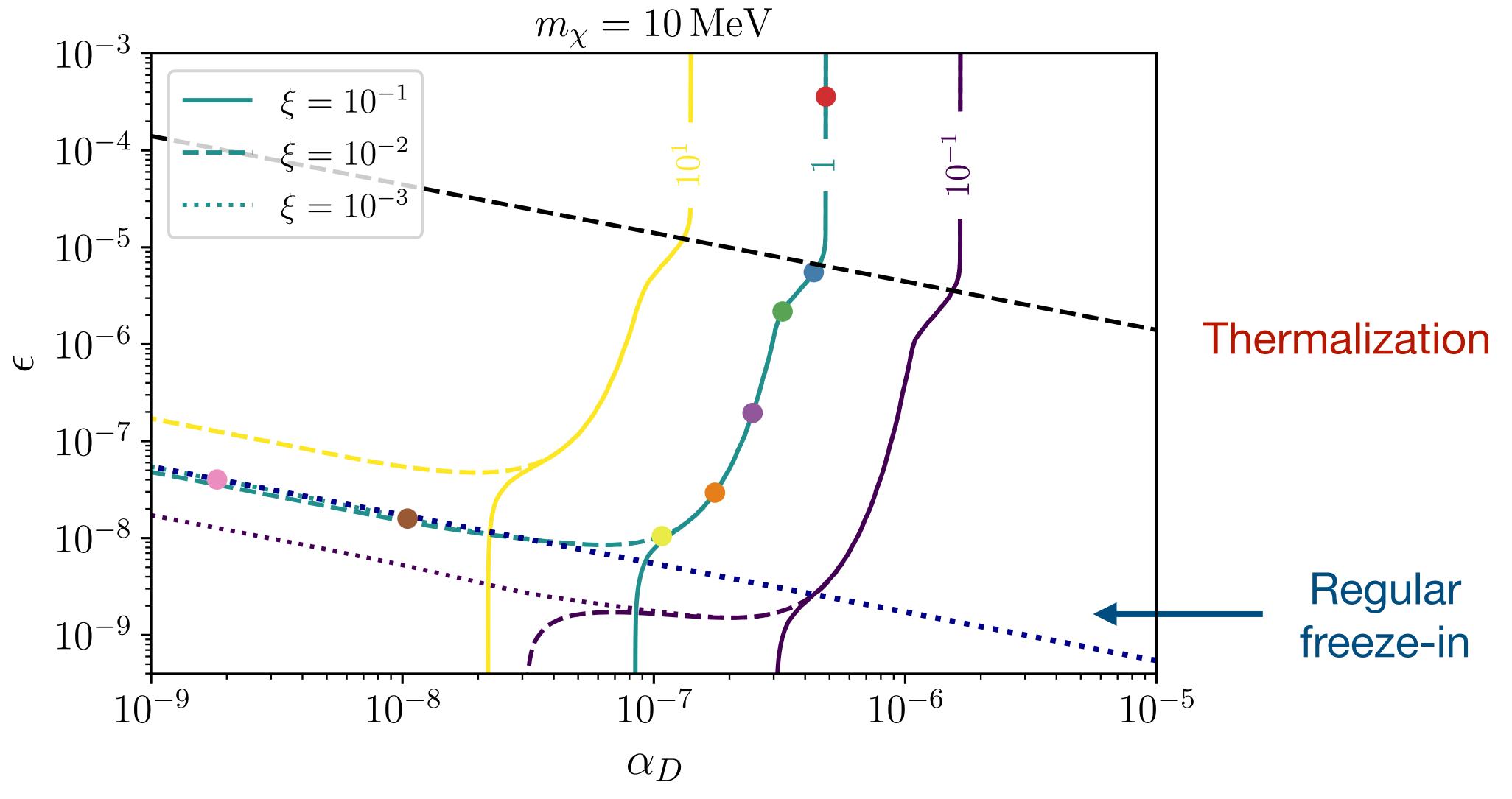
Freeze-in or freeze-out?

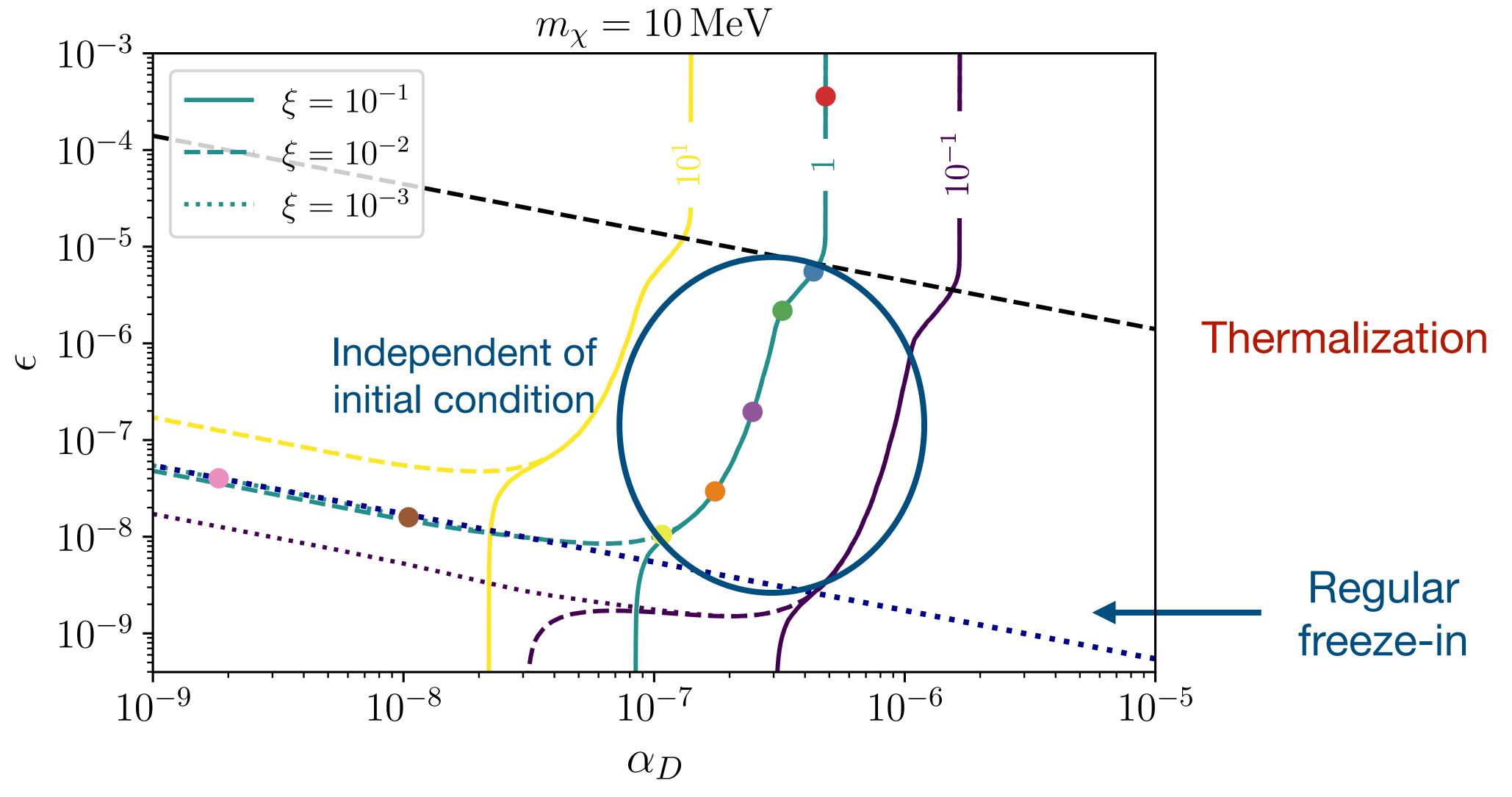


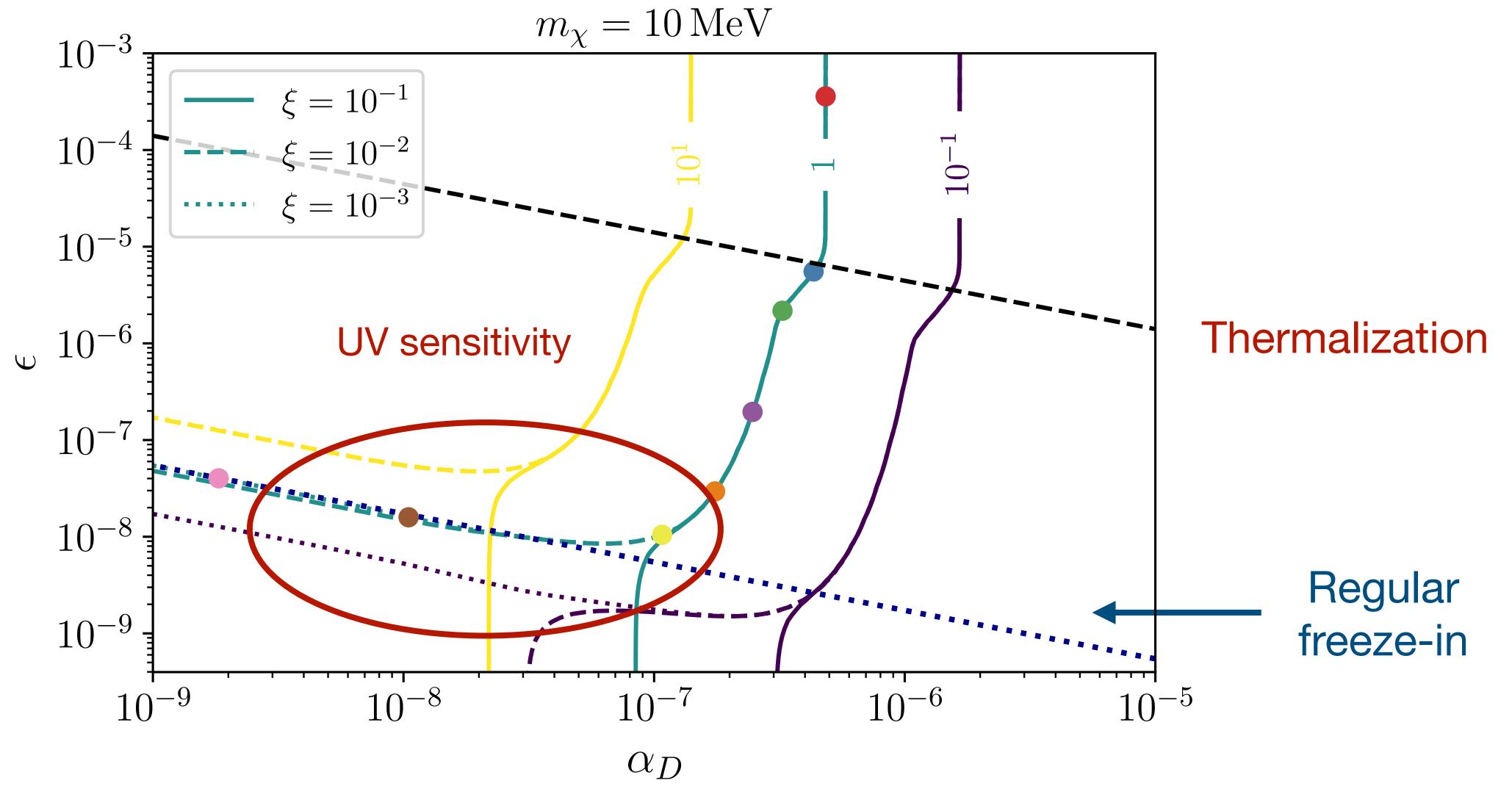
Parameter space and DM relic abundance



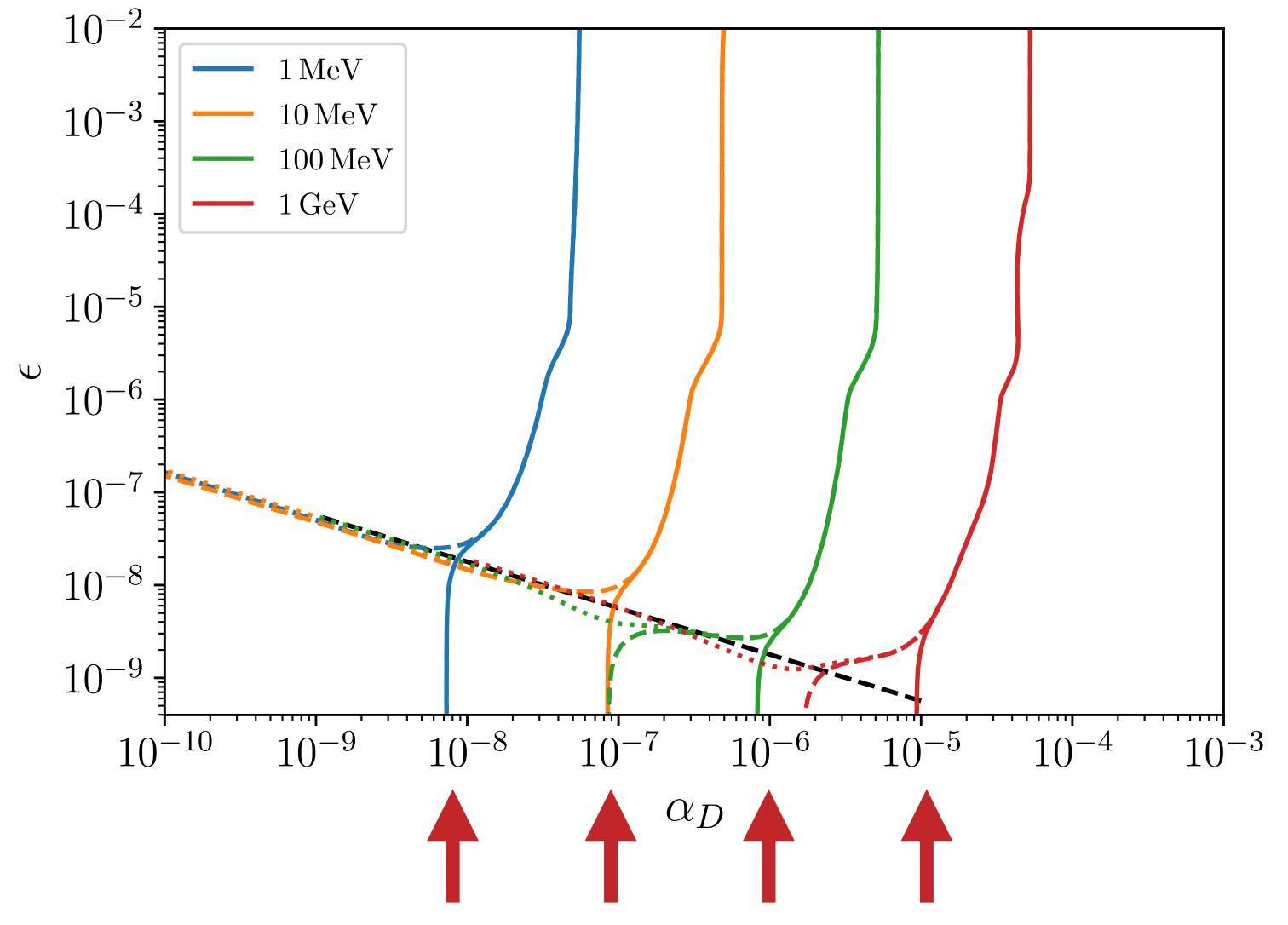








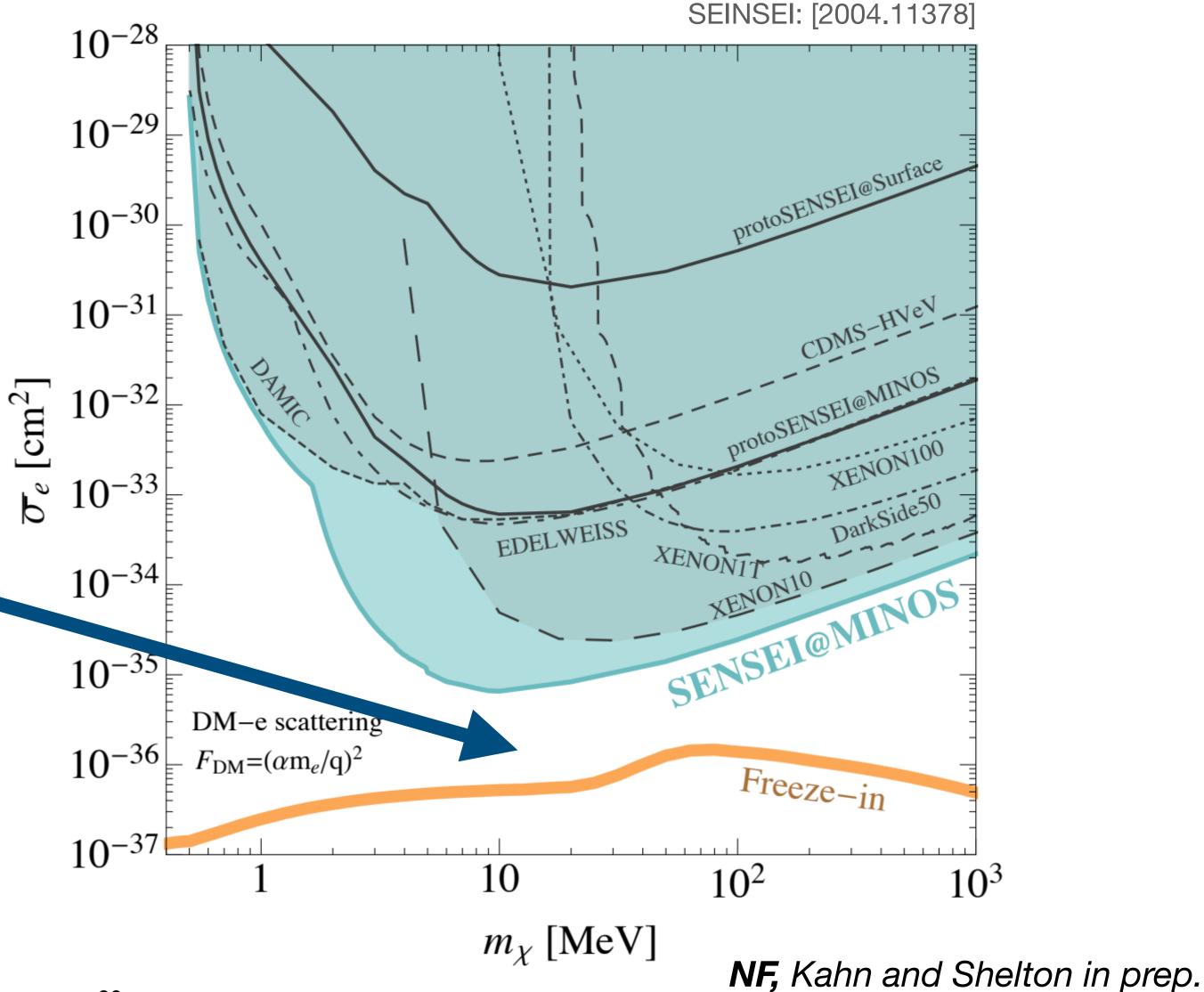
DM relic abundance and parameter space



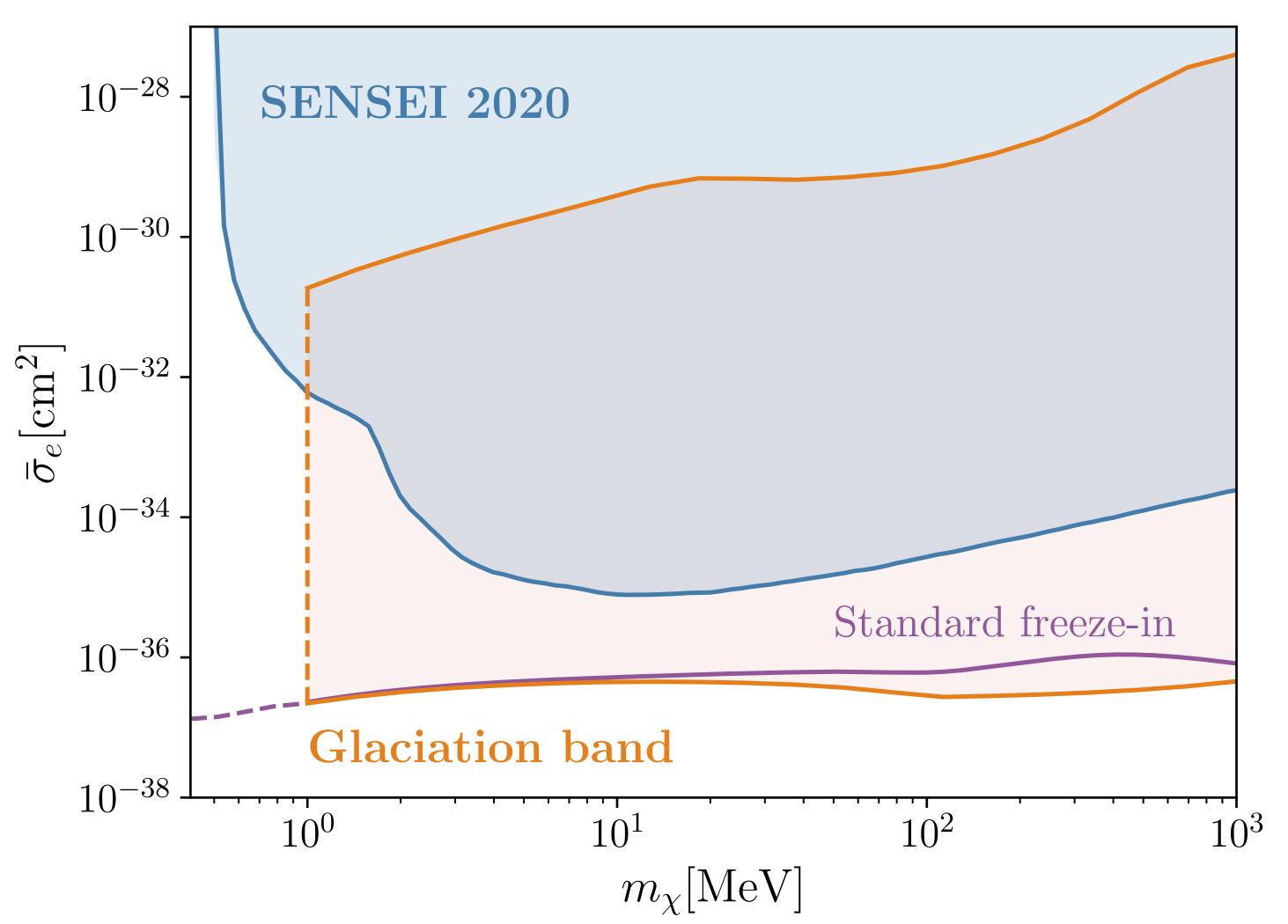
Self-interactions are important

SENSEI 2020

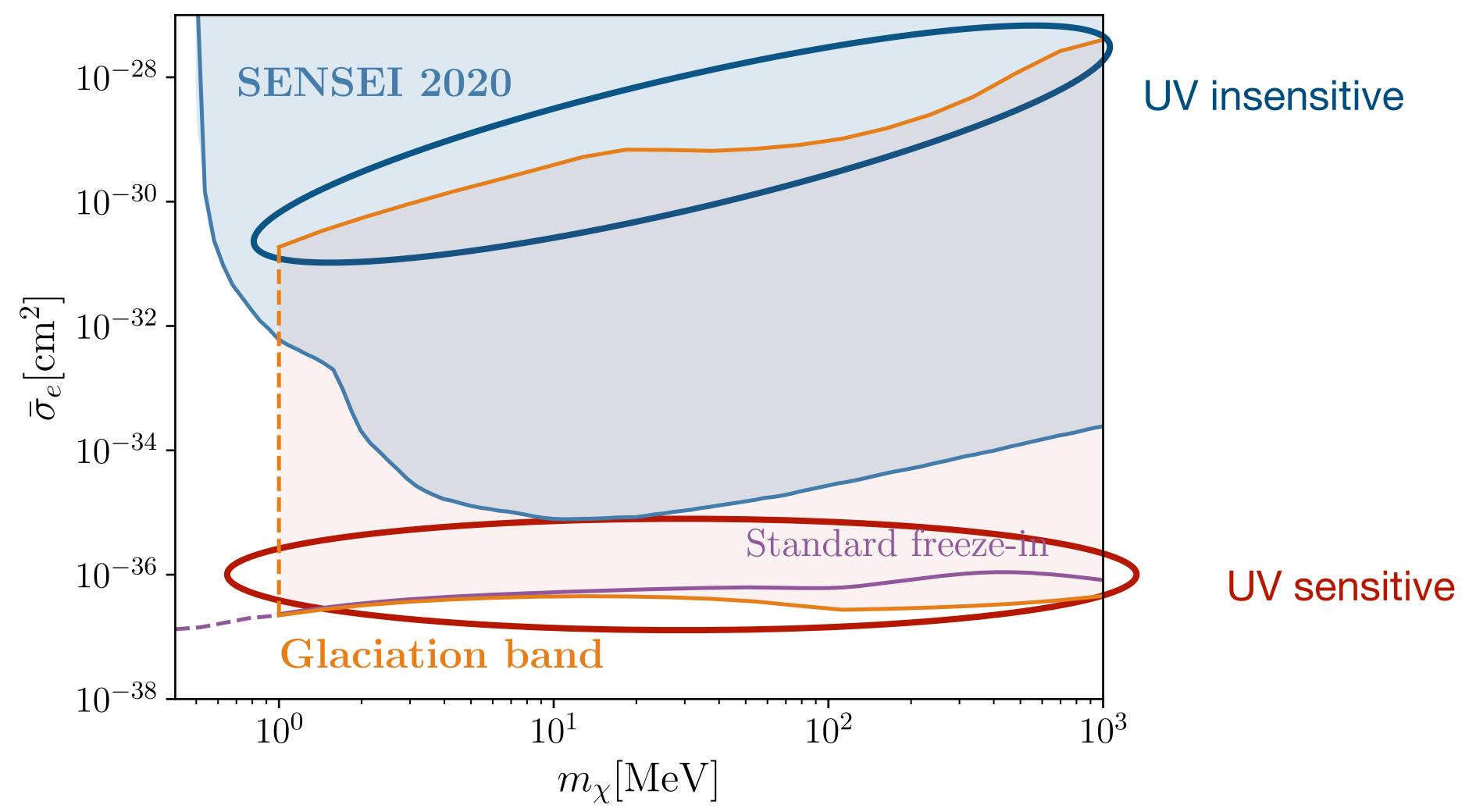
 How the the freeze-in line is affected?



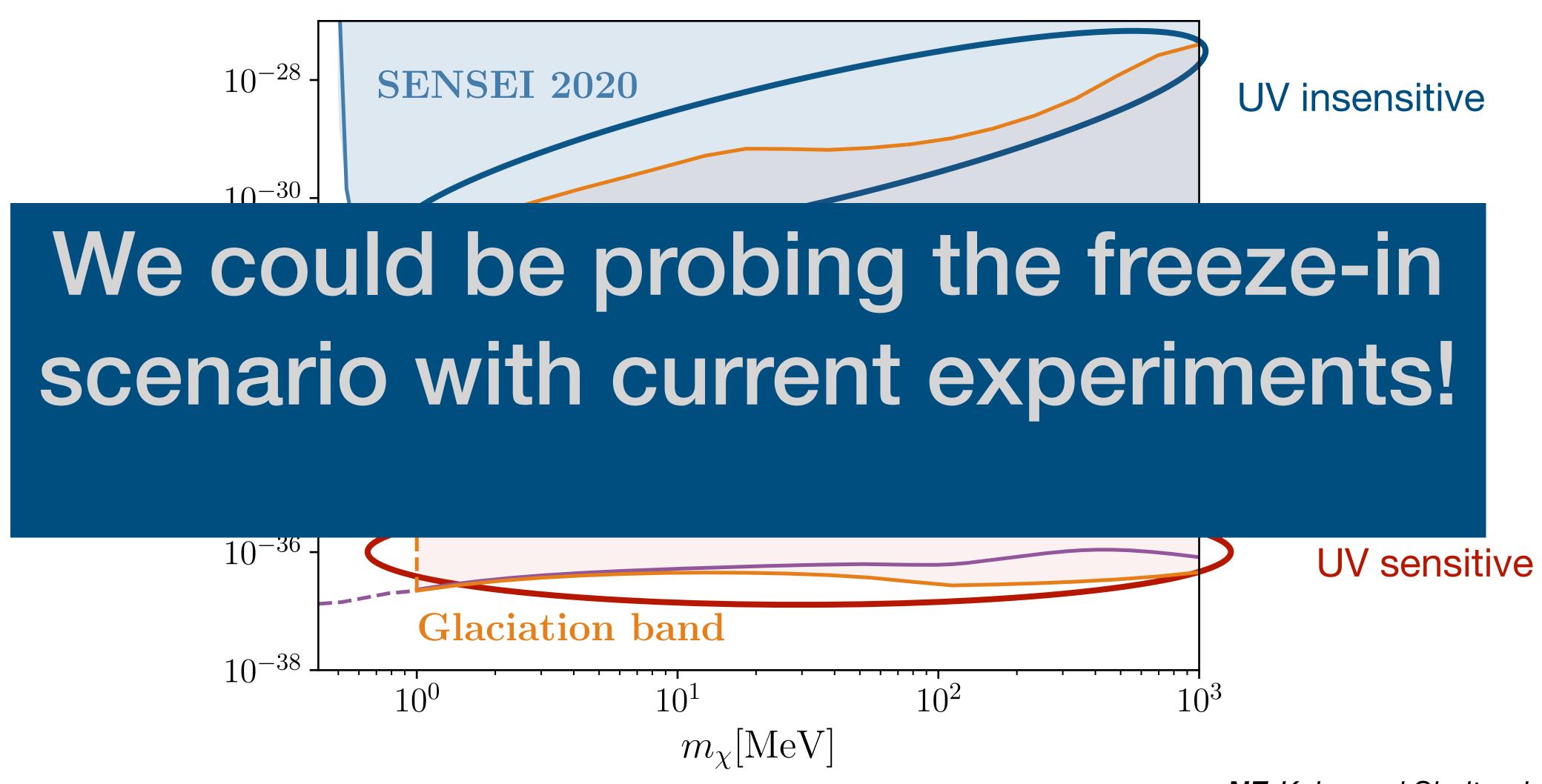
Current Experiments



Current Experiments are testing this parameter space



Current Experiments are testing this parameter space



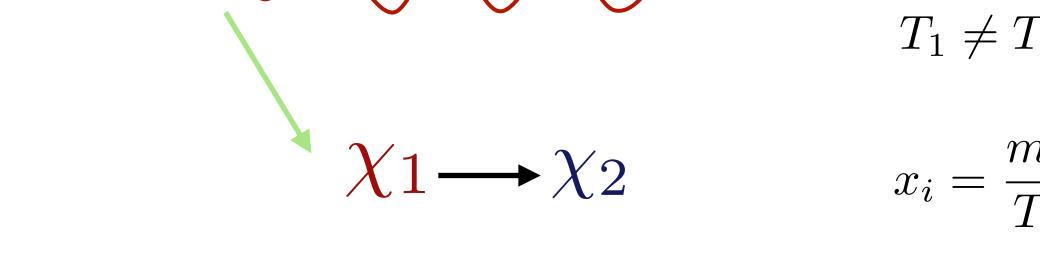
Conclusion

The standard freeze-in paradigm, this same combination of couplings appears in the annihilation cross section, leading to a 1-to-1 relation between thermal history parameter space and direct detection parameter space. As soon as one allows for an initial thermalized population in the dark sector or sizable self-interactions, this "freeze-in line" expands to a "glaciation band" because there are multiple points in the $\epsilon - \alpha_D$ plane which achieve the correct relic abundance.

Gracias

Instantaneous kinetic equilibration of DM

Momentum transferred:



$$C_{1 \, 2 \to 3 \, 4}^{p}(T, \tilde{T}) = n_{1}^{\text{eq}}(T) n_{2}^{\text{eq}}(\tilde{T}) \langle \sigma v p \rangle$$

$$= -\frac{g_{1} g_{2} T^{4} \tilde{T}^{3}}{32 \pi^{4}} \int_{\tilde{s}}^{\infty} d\tilde{s} \, \frac{\lambda^{\frac{1}{2}}(\tilde{s}^{2}, x_{1}, x_{2})}{\tilde{s}} \sigma(s) \left(\lambda(\tilde{s}^{2}, x_{1}, x_{2}) K_{2}(\tilde{s}) + 4\tilde{s} x_{1}^{2} K_{1}(\tilde{s})\right)$$

$$s = \tilde{s}^2 T \tilde{T} + (T - \tilde{T})(T x_1^2 - \tilde{T} x_2^2), \quad \tilde{s}_{\min} = x_1 + x_2$$

- Two different temperatures in the initial state
- Relativistic
- One Integration variable left