

# Recent Progress on Nucleon Form Factors

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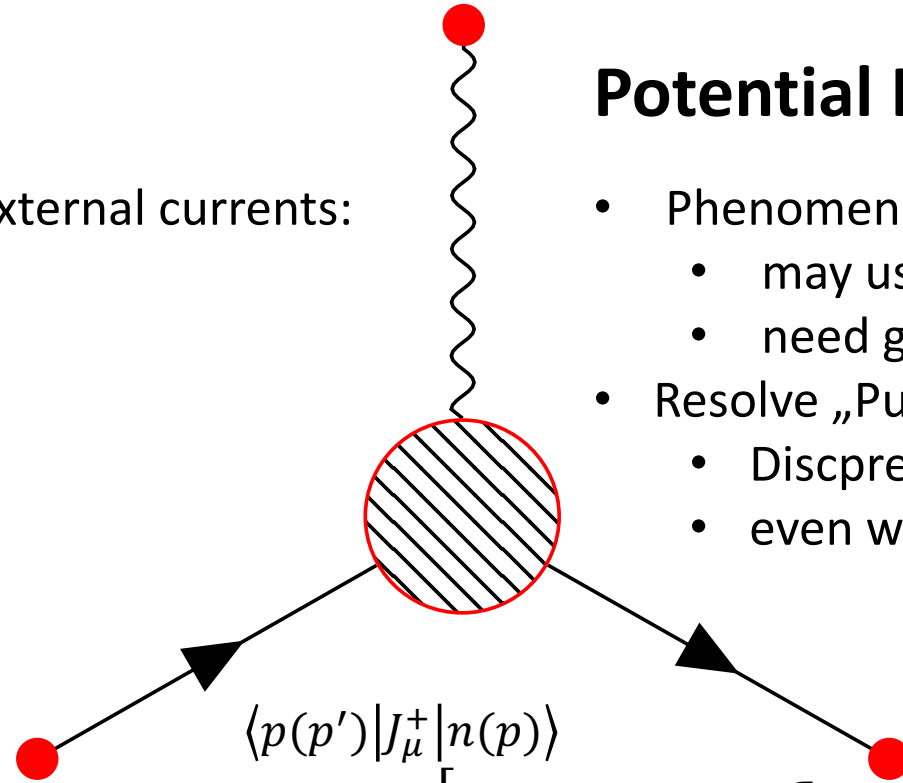
# Form Factors

## What are they?

- Parameterize response to external currents: Form Factors
- Form Factor universal: Enter many processes
- In principle measurable

## Potential Impact of Lattice

- Phenomenology :
  - may use Lattice results as input
  - need good accuracy
- Resolve „Puzzles“:
  - Discrepancies between experiments
  - even with less accuracy



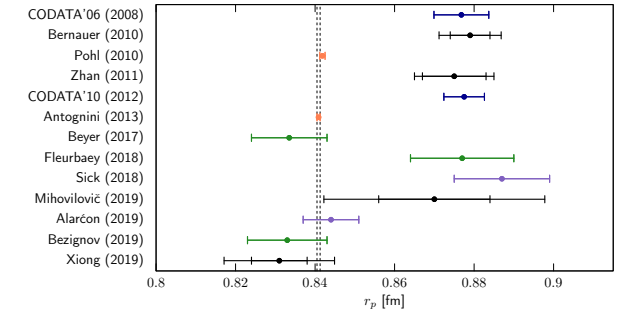
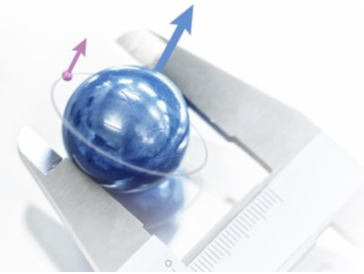
Example V-A current:

$$\begin{aligned} & \langle p(p') | J_\mu^+ | n(p) \rangle \\ &= \bar{u}^p(p') \left[ \gamma_\mu F_1^{CC}(q^2) + i \frac{\sigma_{\mu\nu}}{2m_N} q^\nu F_2^{CC}(q^2) + \frac{q_\mu}{m_N} F_S^{CC}(q^2) \right. \\ &+ \gamma_\mu \gamma_5 F_A^{CC}(q^2) + \frac{\gamma_5 q_\mu}{m_N} F_P^{CC}(q^2) \\ &+ \left. \frac{\gamma_5 (p' + p)_\mu}{m_N} F_T^{CC}(q^2) \right] u^n(p) \end{aligned}$$

# Impact of Form Factors

- Proton Radius Puzzle

- » Provide ab-initio calculation



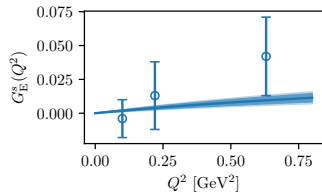
Taken from Bernauer, EPJ Web Conf. 234 (2020)

- Precision Tests of SM

- » Via strangeness FF → Parity Violation Experiments

- Lattice determinations of strange FF very precise

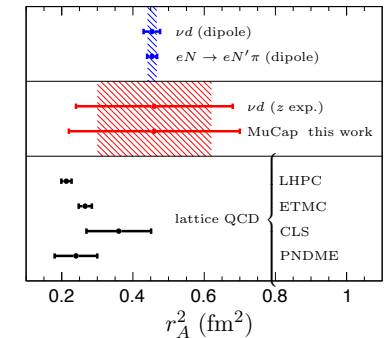
Taken from D.D. et al. *Phys.Rev.Lett.* 123 (2019) 21, 212001



- » Via axial FF → Vital input to neutrino-nucleus scattering

- Lattice competitive to z-exp extractions of experiments

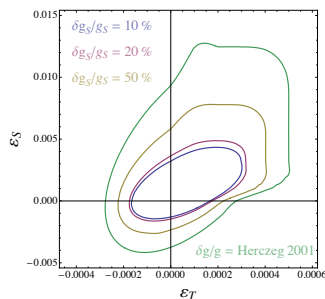
A. Meyer, Tue ID: 291



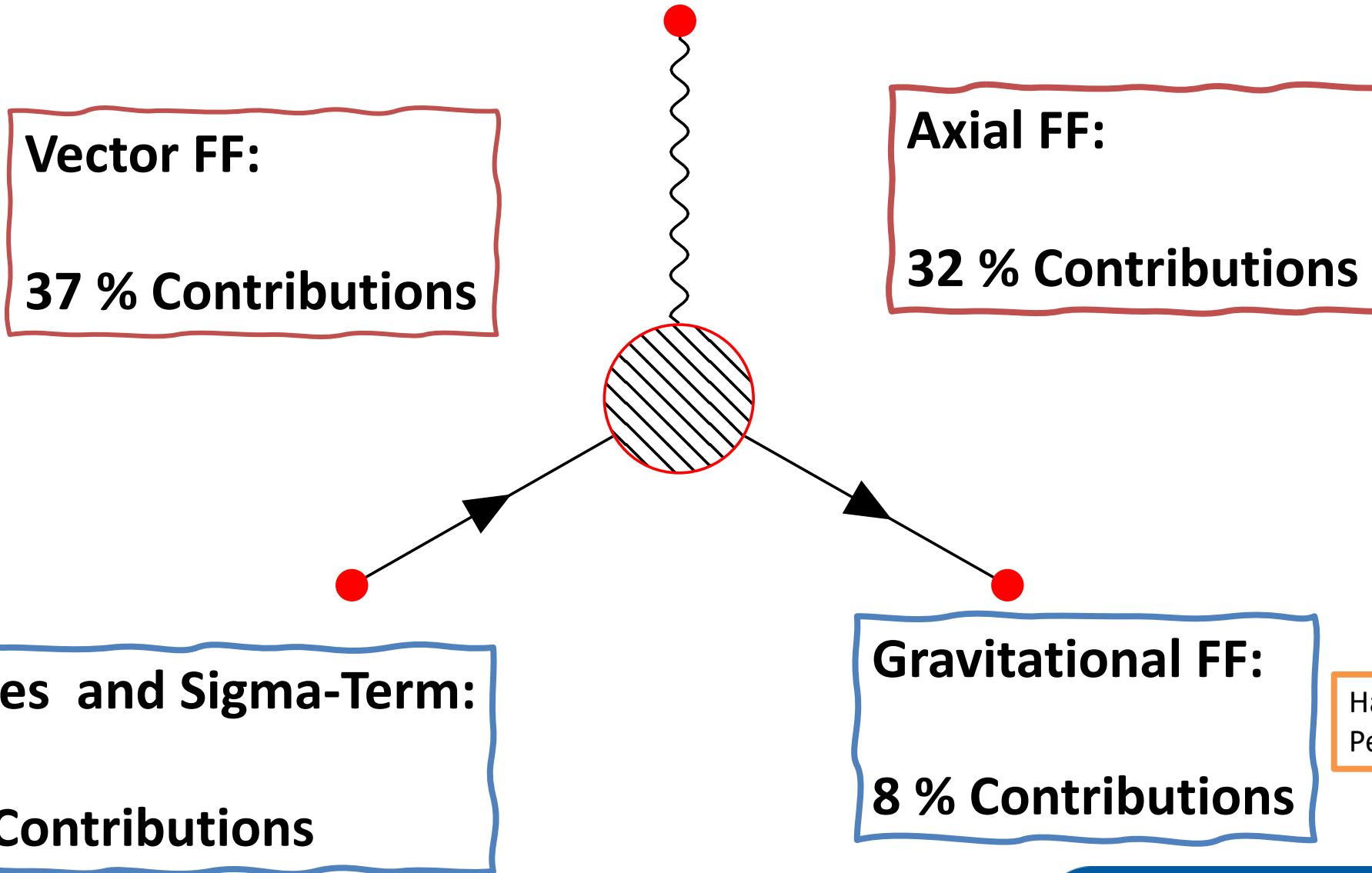
Taken from A. Kronfeld, et al. *Eur. Phys. J. A* 55, 196 (2019)

- » Via Charges → Constraining BSM EFT couplings

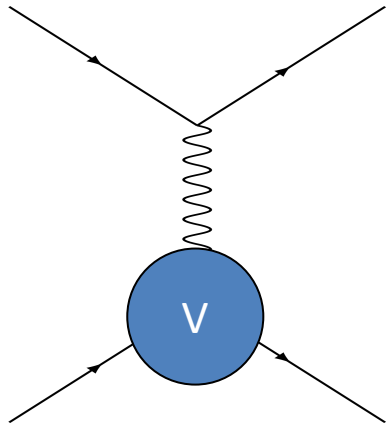
Taken from T. Bhattacharya et al., *Phys. Rev. D* 85, 054512 (2012)



# My apologies for omissions ...



# Nucleon Form Factors



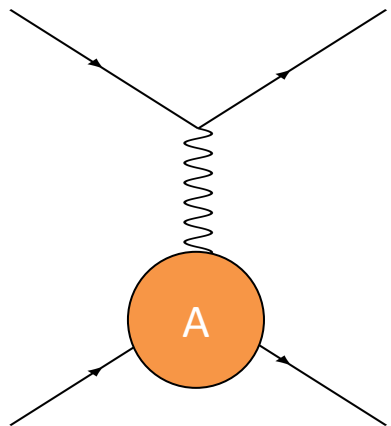
$$\langle N(p') | \bar{q} \gamma_\mu q | N(p) \rangle = \bar{u}(p') \left[ \gamma_\mu F_1^q(Q^2 = -q^2) + i \frac{\sigma_{\mu\nu}}{2m_N} q^\nu F_2^q(Q^2) \right] u(p)$$

$$G_E(Q^2) = F_1^q(Q^2) - \frac{Q^2}{2m_N} F_2^q(Q^2) = \left( 1 - \frac{1}{6} \langle r_E^2 \rangle Q^2 + \dots \right)$$

$$G_M(Q^2) = F_1^q(Q^2) + F_2^q(Q^2) = \mu \left( 1 - \frac{1}{6} \langle r_M^2 \rangle Q^2 + \dots \right)$$

Dirac

Pauli

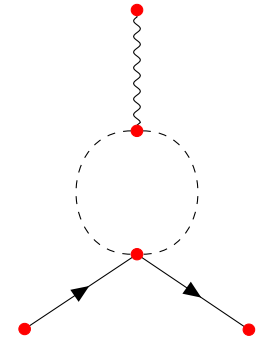


$$\langle N(p') | \bar{q} \gamma_\mu \gamma_5 q | N(p) \rangle = \bar{u}(p') \left[ \gamma_\mu G_A(Q^2) + i \frac{q^\mu}{2m_N} G_P(Q^2) \right] \gamma_5 \frac{\tau^3}{2} u(p)$$

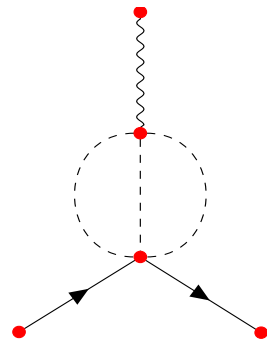
$$G_A(Q^2) = g_A \left( 1 - \frac{1}{6} \langle r_A^2 \rangle Q^2 + \dots \right)$$

$$2m_N G_A(Q^2) - \frac{Q^2}{2m_N} G_P(Q^2) = \frac{2M_\pi^2 F_\pi}{M_\pi^2 + Q^2} G_{\pi N}(Q^2) \quad (\text{PCAC})$$

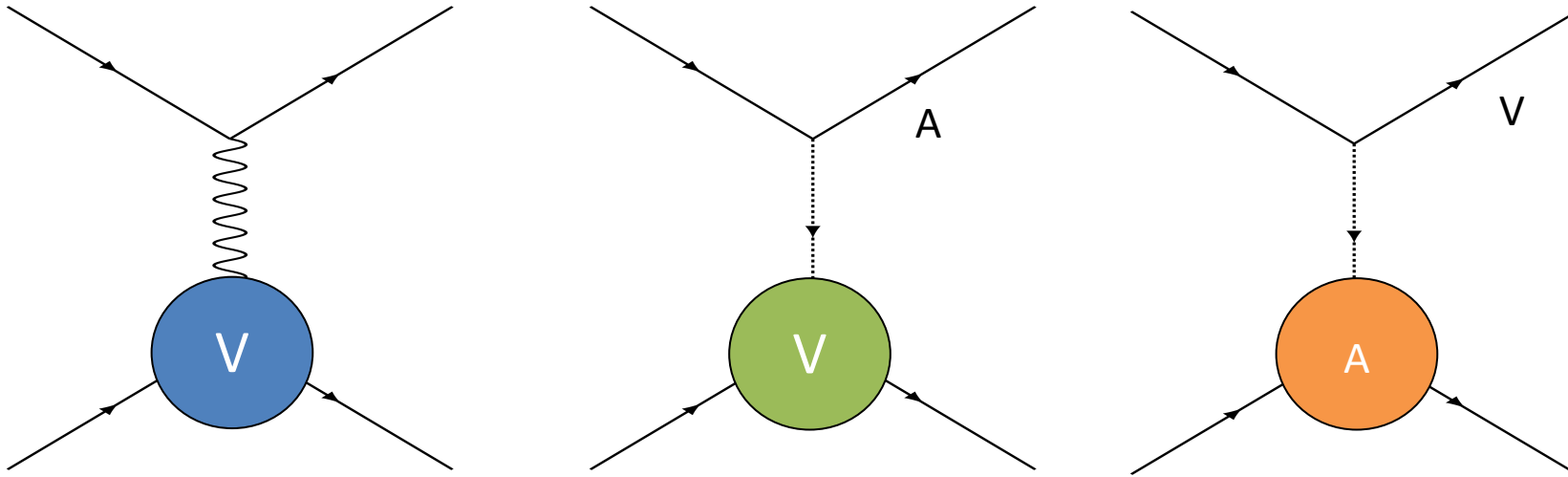
$$\langle r_{E/M}^2 \rangle \xrightarrow{M_\pi \rightarrow 0} \left\{ \ln M_\pi, \frac{1}{M_\pi} \right\}$$



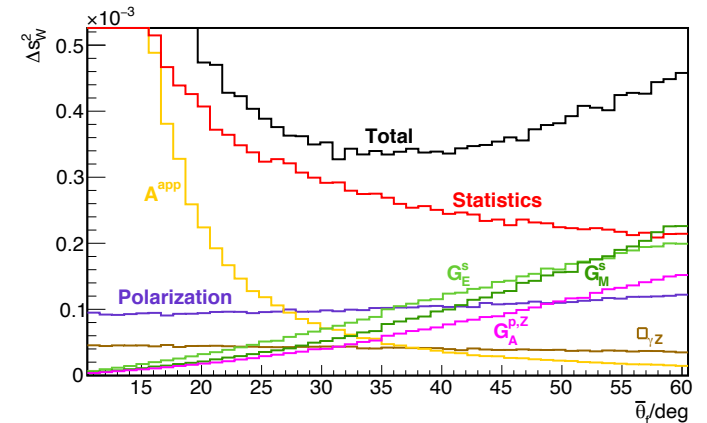
$$\langle r_A^2 \rangle \xrightarrow{M_\pi \rightarrow 0} \text{const.}$$



# Parity Violation



Projected Error Budget for P2:



Taken from arxiv:1802.04759

$$A = \frac{d\sigma_R - d\sigma_L}{d\sigma_R + d\sigma_L} = -\frac{G_F Q^2}{4\sqrt{2}\pi\alpha} \frac{\epsilon G_E^Y G_E^Z + \tau G_M^Y G_M^Z - (1 - 4\sin^2\theta_W) \epsilon' G_M^Y G_A}{\epsilon (G_E^Y)^2 + \tau (G_M^Y)^2}$$

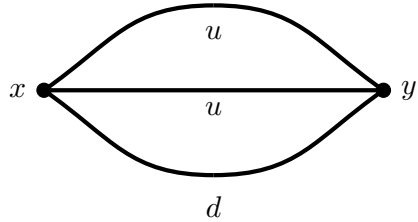
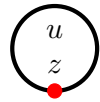
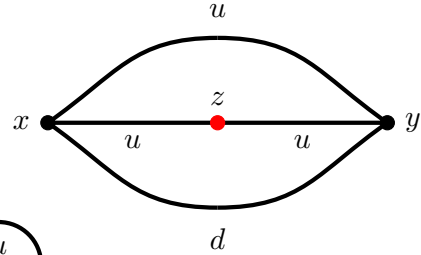
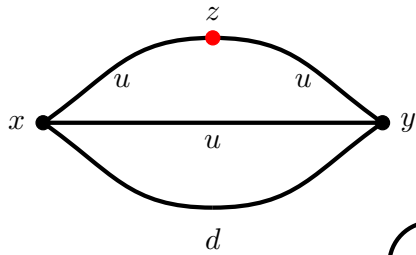
$$\tau = \frac{Q^2}{4m_N^2}, \epsilon = \left(1 + 2(1 + \tau) \tan^2 \frac{\theta}{2}\right)^{-1}, \epsilon' = \sqrt{\tau(1 + \tau)(1 - \epsilon^2)}$$

$$G_{E/M}^{Z,p} = (1 - 4\sin^2\theta_W) G_{E/M}^{Y,p} - G_{E/M}^{Y,n} - G_{E/M}^S$$

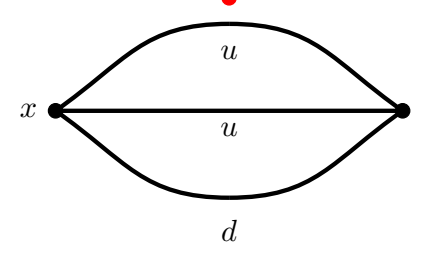
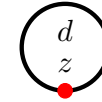
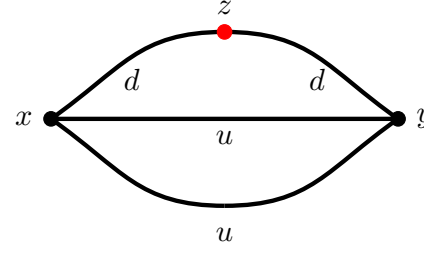
- Sensitive to the Weak Charge
  - Test of SM at low energies
  - Need e/m FF (strange)
  - Need axial FF (strange)
- (Decomposition assuming Isospin Symmetry)

# Computational Setup Lattice

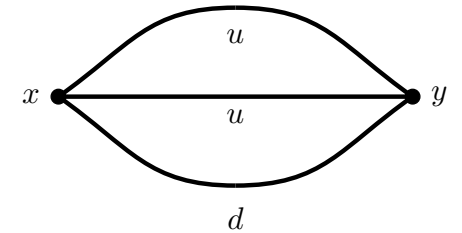
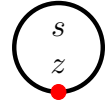
*u*-quark:



*d*-quark:



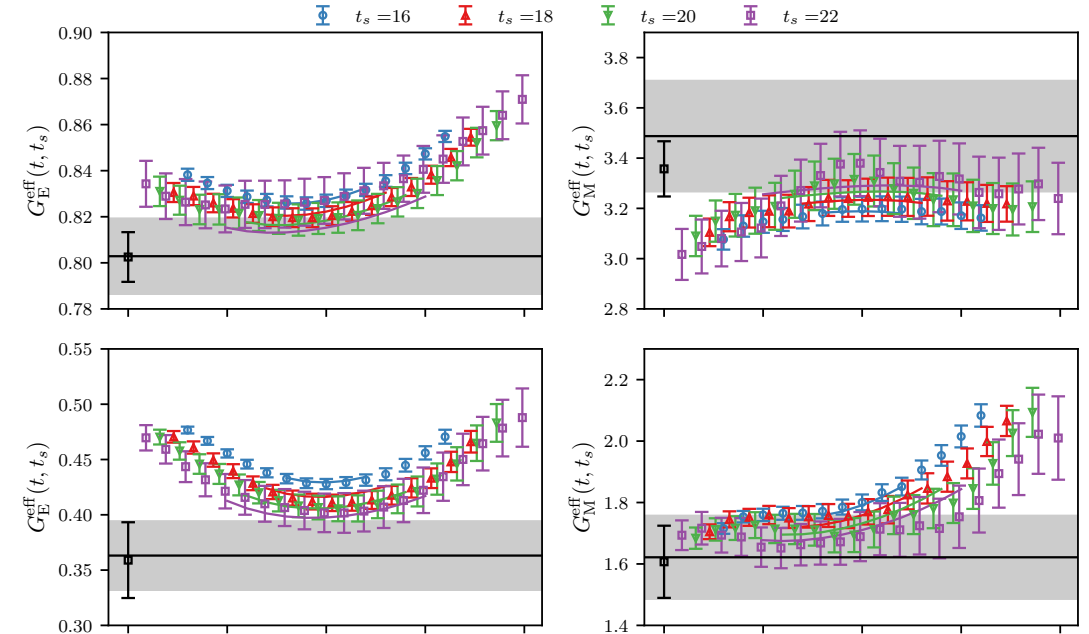
*s*-quark:



Ratio gives access to Form Factors:

$$R_{V_\mu}^s(z_0, \mathbf{q}; y_0, \mathbf{p}'; \Gamma_\nu) = \frac{C_{3,V_\mu}^s(\mathbf{q}, z_0; \mathbf{p}', y_0; \Gamma_\nu)}{C_2(\mathbf{p}', y_0)}$$

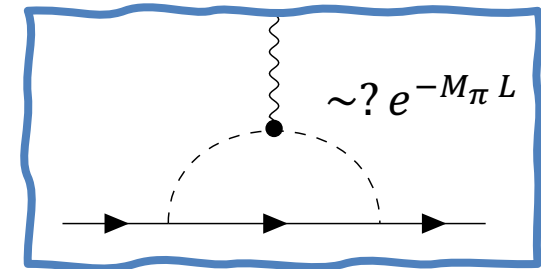
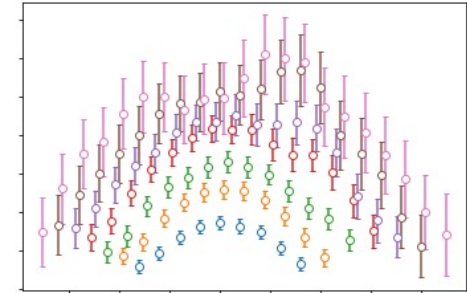
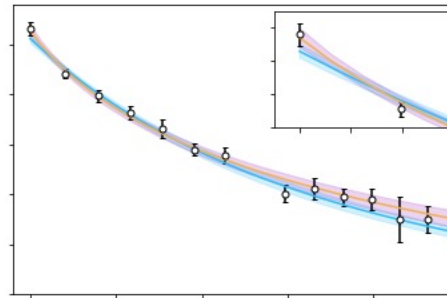
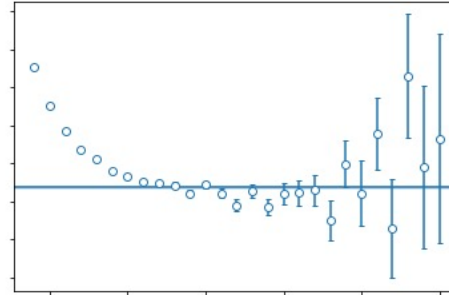
$$\times \sqrt{\frac{C_2(\mathbf{p}', y_0)C_2(\mathbf{p}', z_0)C_2(\mathbf{p}'-\mathbf{q}, y_0-z_0)}{C_2(\mathbf{p}'-\mathbf{q}, y_0)C_2(\mathbf{p}'-\mathbf{q}, z_0)C_2(\mathbf{p}', y_0-z_0)}}.$$



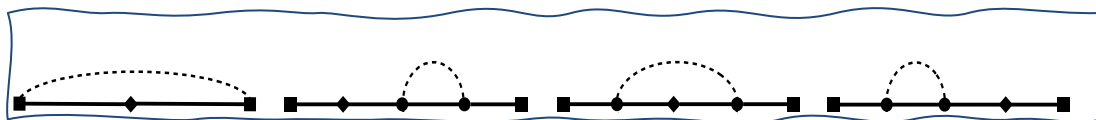
Example of these Ratios

# Sources of Uncertainty

- Statistical Accuracy
- Excited State Contamination
- Model Dependence
- Extrapolations
  - Chiral
  - Continuum
  - Finite Size



Aside: Can calculate Excited States directly in ChPT even for FF (reliable for large  $t_{sep}$ )



O. Bär, Tue ID: 70

Taken from O. Bär, H.Colic, *Phys.Rev.D* 103 (2021) 11, 114514



# Excited States

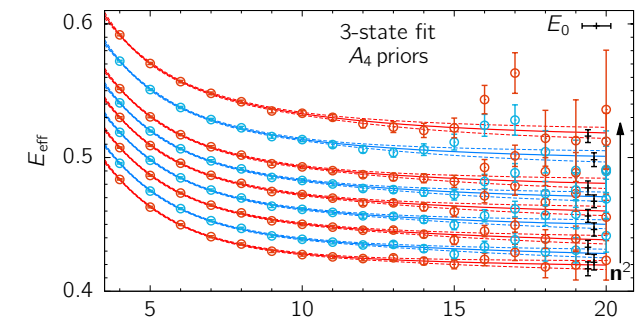
- Write any correlation function generically as

$$C(t, t_{sep}) = a + b e^{-t \Delta E_1} + c e^{-(t_{sep}-t)\Delta E_1} + \dots$$

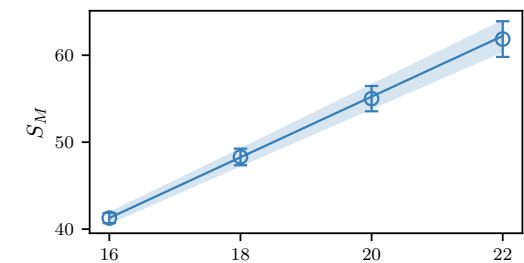
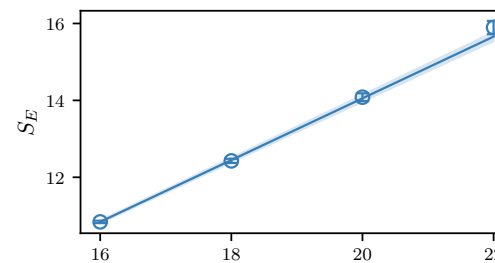
- „Just“ do the fit (Multistate method)
- Summation method

$$\sum_{t=1}^{t_{sep}-1} C(t, t_{sep}) = a(t_{sep} - 1) + \tilde{b} e^{-t_{sep}\Delta E_1} + \tilde{c} e^{-t_{sep}\Delta E_1} + \dots$$

- Note that:  $t_{sep} \gg t, (t_{sep} - t)$
- $a$  = ground state ME
- $b, \tilde{b}$ - and  $c, \tilde{c}$ -terms ES



Taken from Y-C. Jang Talk

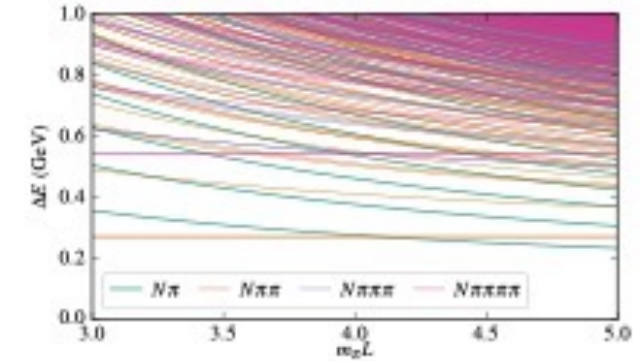


# Multistate

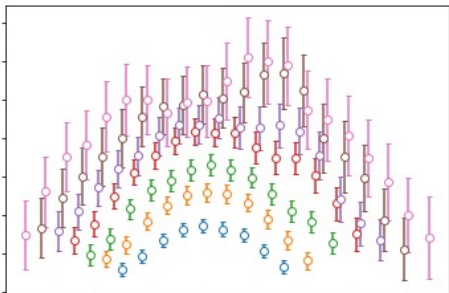
- Multistate Fits

- How many states to include? (states become dense)
- Use results from Spectrum?
- Gaps universal (?)  
⇒ Do Simultaneous Fits

- » Possibly (very) large covariance matrices
  - Make this less demanding by taking 2pt-functions spectrum
    - as is or
    - as priors for 3pt-corr.
- » Data might NOT constrain all parameters  
⇒ Stabilize via priors
- » Assumptions about universality between 2pt- and 3-pt might not be justified



Taken from J. Green arxiv:1812.10574

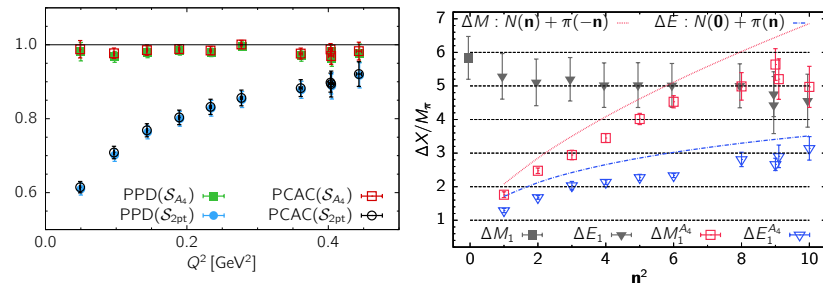


See e.g. R. Gupta, Mon ID: 527

# Multistate

- How do we know if asymptotics is reached?  
In general not easy to see

- Axial Form Factors use PCAC as consistency check

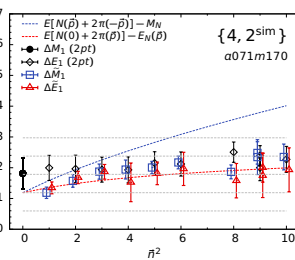
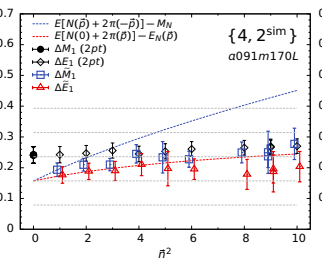
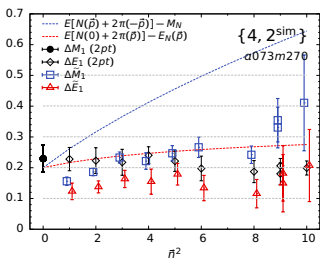
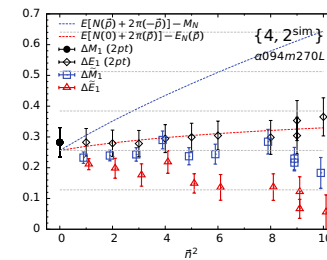
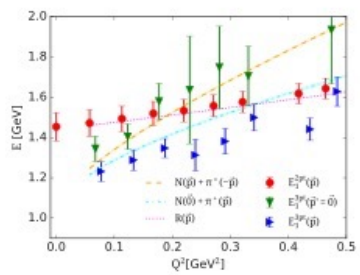
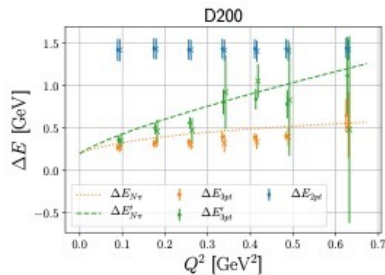


Y-C. Jang, R. Gupta, B. Yoon, T. Bhattacharya, Phys. Rev. Lett. 124, 072002 (2020)

$$2\hat{m}G_P(Q^2) = 2MG_A(Q^2) - \frac{Q^2}{2M} \tilde{G}_P(Q^2)$$

Y-C. Jang, Tue ID:519

- Gaps might not be universal between 2-pt and 3-pt (Less severe for vector?)



Bali et. al (RQCD) JHEP05 (2020) 126

Alexandrou, et al. (ETMC) Phys.Rev.D 103 (2021) 3

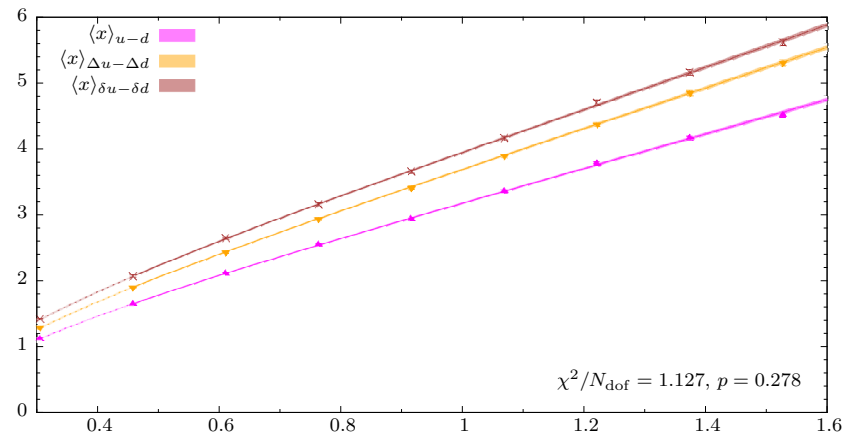
S. Park, et. al (NME) arXiv:2103.05599

G. Koutsou, Tue ID: 401

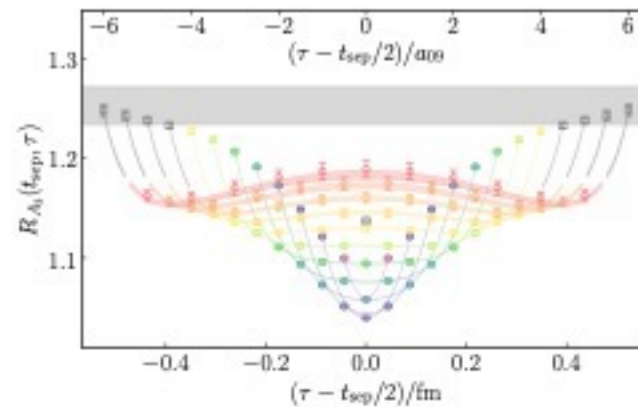
# Summation

- Easy to apply but errors usually larger
- Suppression of **ES** paramterically larger
  - » Use more  $t_{seps}$  in Fits (Deviation from linearity)

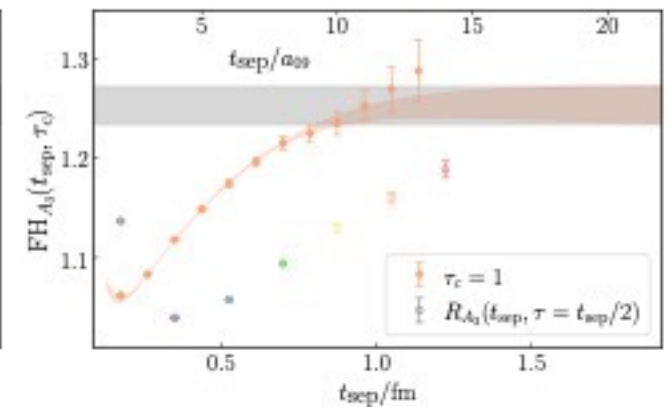
$$\sum_{t=1}^{t_f-1} C(t, t_f) = a(t_{sep} - 1) - \tilde{b}e^{-t_{sep}\Delta E_1} - \tilde{c}e^{-t_{sep}\Delta E_1} + \dots$$



Taken from K. Ottnad, Mon ID: 229



Taken from A. Walker-Loud, Mon ID: 612



Callat: arXiv:2104.05226

# Model Dependence

- Derived Quantities depend on model for  $Q^2$  dependence

$$C^q \left( 1 - \frac{1}{6} \langle r_E^2 \rangle Q^2 + \dots \right)$$

- » Dipole is not very flexible!

J. Bernauer et al., *Phys.Rev.C* 90 (2014) 1, 015206

....

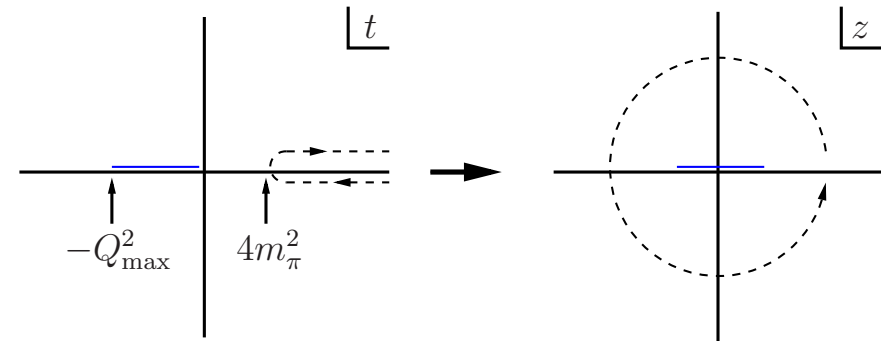
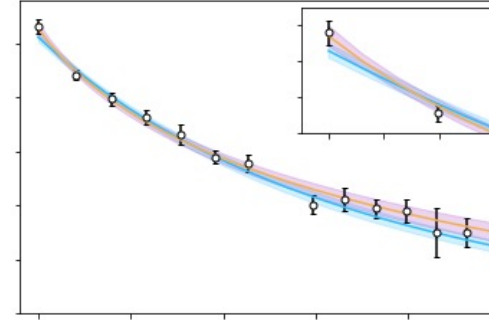
- » Use z-expansion, e.g.

$$G_{E/M}(Q^2) = \sum_{k=1/0}^5 a_k^{E/M} z(Q^2)^k,$$
$$z(Q^2) = \frac{\sqrt{t_{\text{cut}} + Q^2} - \sqrt{t_{\text{cut}}}}{\sqrt{t_{\text{cut}} + Q^2} + \sqrt{t_{\text{cut}}}}.$$

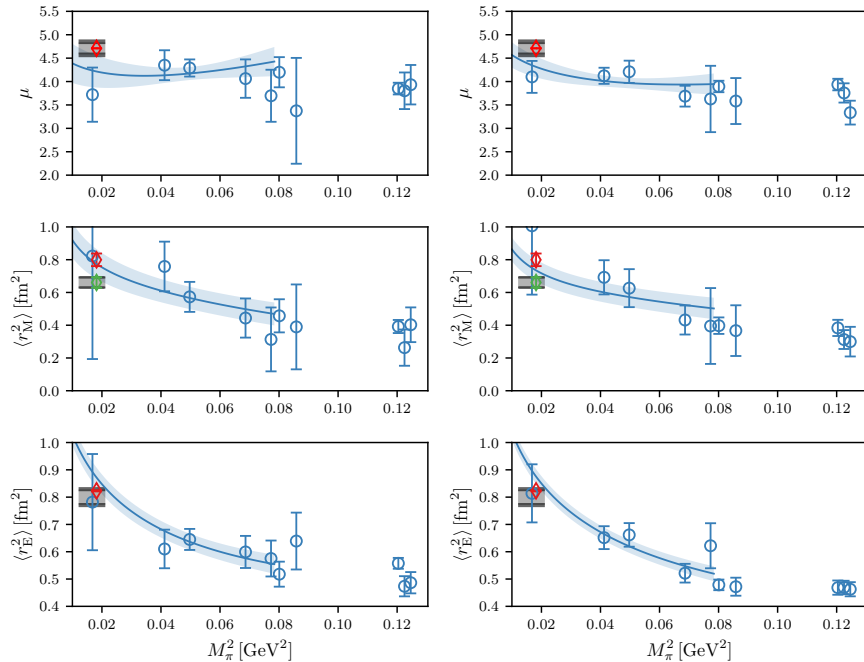
R. Hill, G. Paz, *Phys. Rev. D* 82 (2010), 113005

- » Direct methods avoid this altogether!

- Most recently K-I. Ishikawa (PACS) arxiv:2107.07085



# Model Dependence & CCF - Combine



D.D. et. al, Phys. Rev. D 103, 094522 (2021)

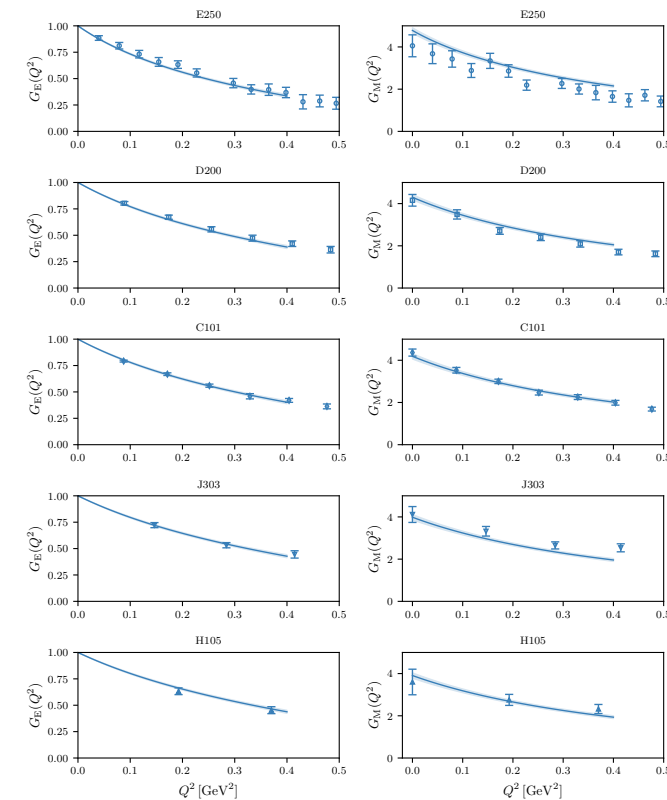
## Z-Expansion Fits

- Do the z-exp for each ensemble
- Perform CCF Fits
- Number of  $Q^2$  points lost at this stage
- Combine z-exp and CCF (larger errors)

Y-C. Yang (PNDME) Phys. Rev. D 101, 014507 (2020)

## Direct Fits Using Chiral EFT

- Use Chiral EFT T. Bauer, et al, Phys. Rev. C86, 065206 (2012).  
Less freedom at small  $Q^2$  vs z-exp (smaller errors)
- Can be more aggressive with cuts in  $Q^2$
- Results from usual z-expansion consistent



# Lots of Variations

- In the end still have lots of variations ( $Q^2, M_\pi^2, \mathcal{O}(a^2), \mathcal{O}(e^{-M_\pi L})$ )
- No clear winner
- Perform averages based on AIC weights

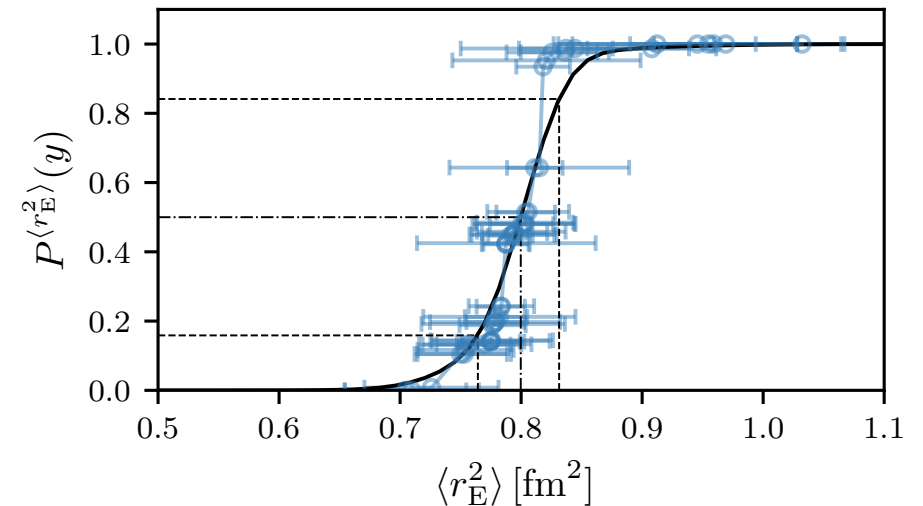
E. Neil, Tue ID: 316

$$w_i^{\text{AIC}} = \frac{e^{-\frac{1}{2}\text{AIC}_i}}{\sum_j e^{-\frac{1}{2}\text{AIC}_j}}$$

- Treat estimates as random variable

$$P^x(y) = \int_{-\infty}^y \sum_i^n w_i \mathcal{N}(y'; x_i, \sigma_i^2) dy'$$

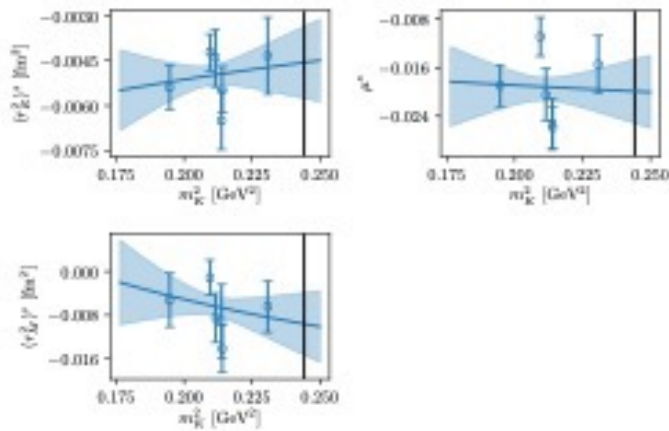
Strategy from S. Borsanyi et al. (2020), *Nature* **593**, 51–55 (2021)



Weighting applied to isovector electric radius.

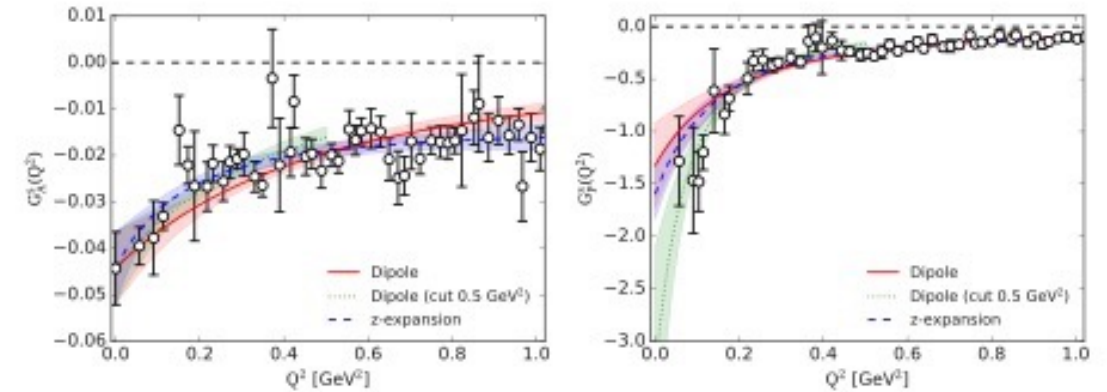
# Results for Isovector Strange FF

- Disconnected Diagrams only
- Very precise estimates
- CCF mild

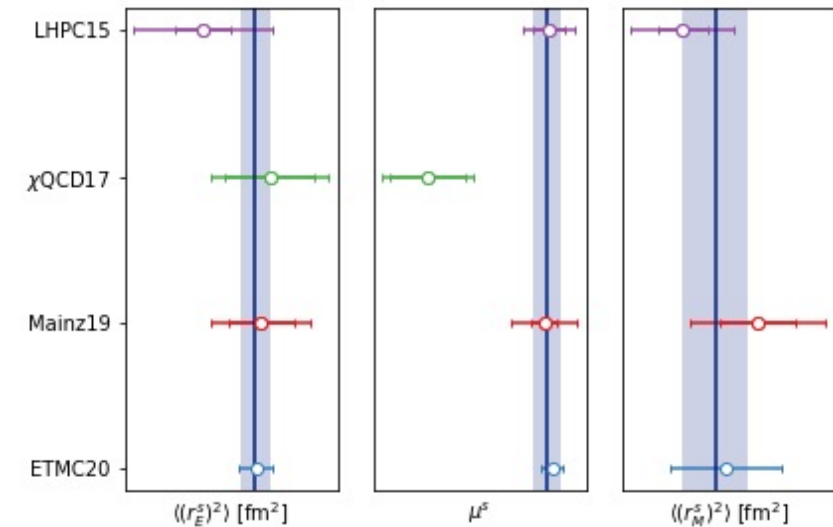


Taken from *PRL* 123 (2019) 21, 212001

- Consistent picture from Lattice non-zero radii  
(Blue Band PDG-style average)



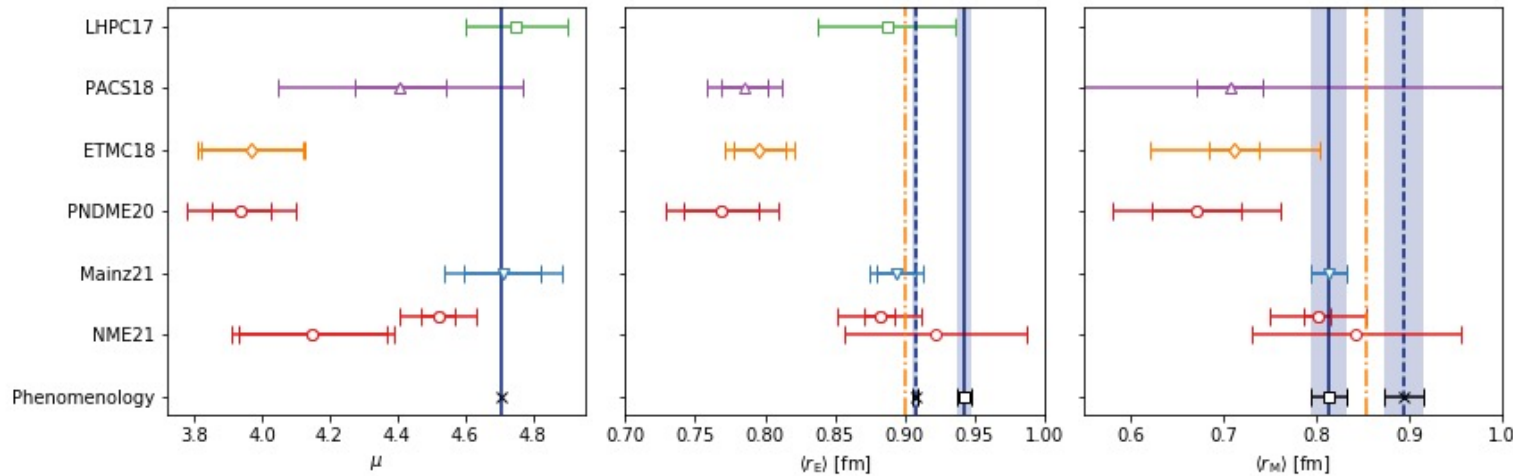
Taken from C. Alexandrou, et al. arxiv:2106.13468



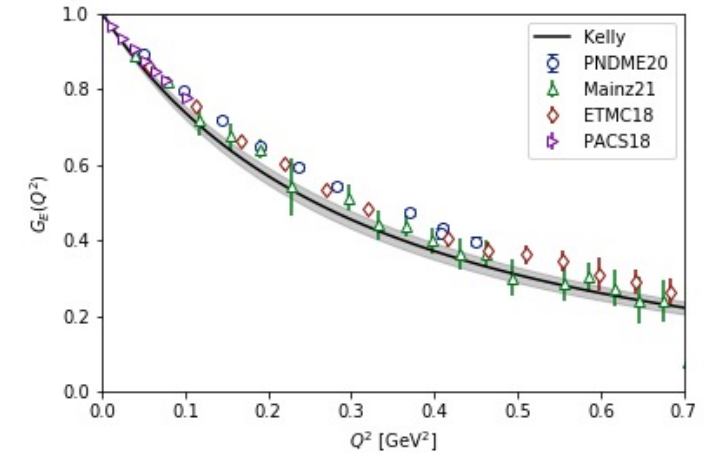


# Results Isovector Vector FF

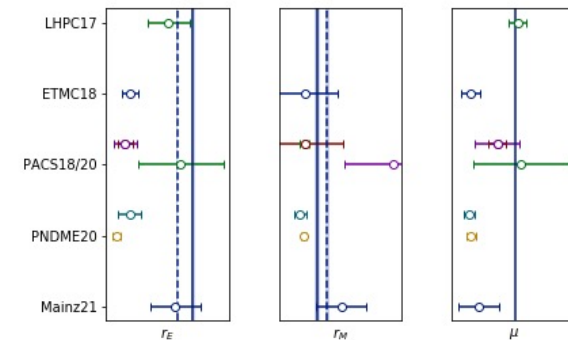
- Excited State Contamination: Summation or Multistate
- $Q^2$  -dependence: via z-Expansion/Dipole/EFT/Pade
- CCF extrapolations performed



Current status after CCF



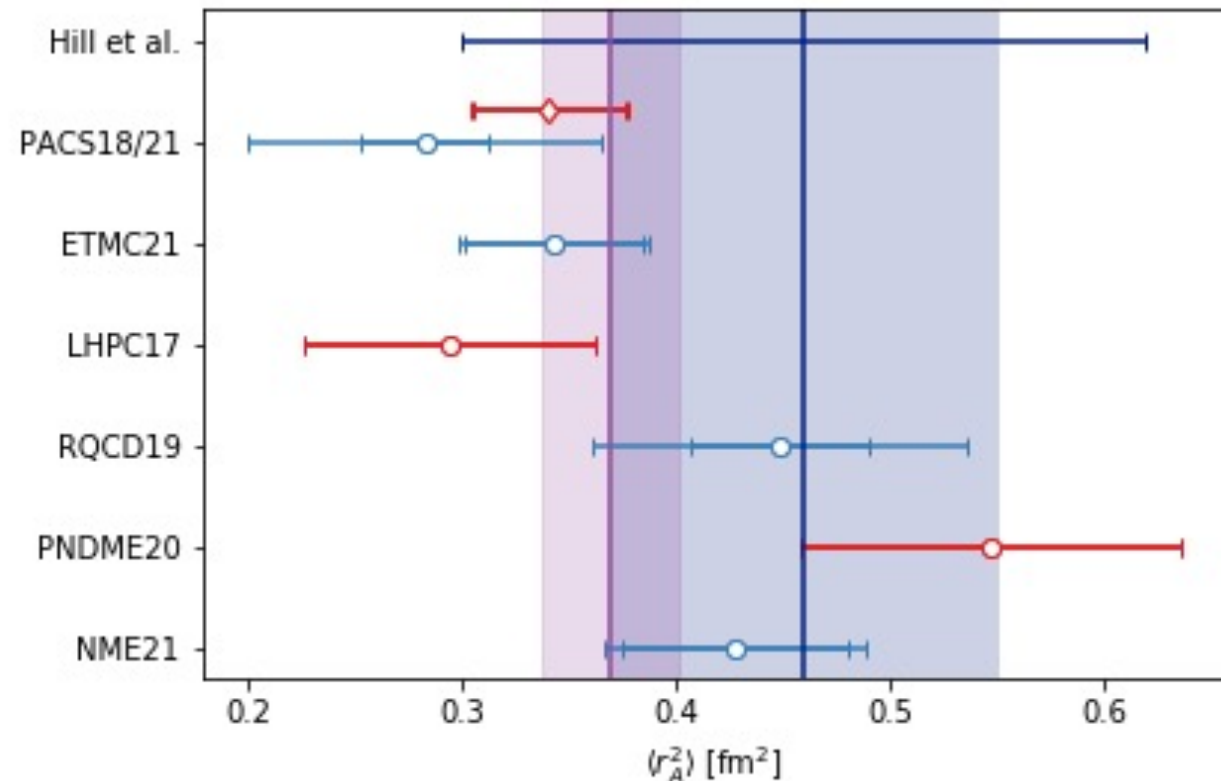
Recent ensembles @ Physical pion mass



Isoscalar FF e.g. M. Salg, Tue ID: 406

# Results Isovector Axial FF

- Excited State Contamination: Summation/Multistate ( $\pi N$ )
- $Q^2$  -dependence param. via z-Expansion
- CCF extrapolations are performed
- Lattice:  $\sim 10\%$  statistical error



(Determination with 20 % acc blue shaded area)  
(Purple Band PDG-Style Average)

New analyses e.g. T. Schulz, Tue ID: 86

# Lots of experiments

