

# Review of muon g-2



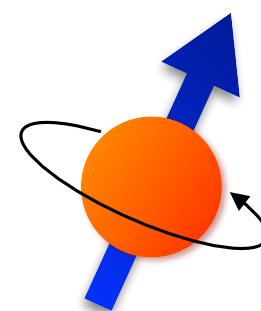
Aida X. El-Khadra  
University of Illinois



# Outline

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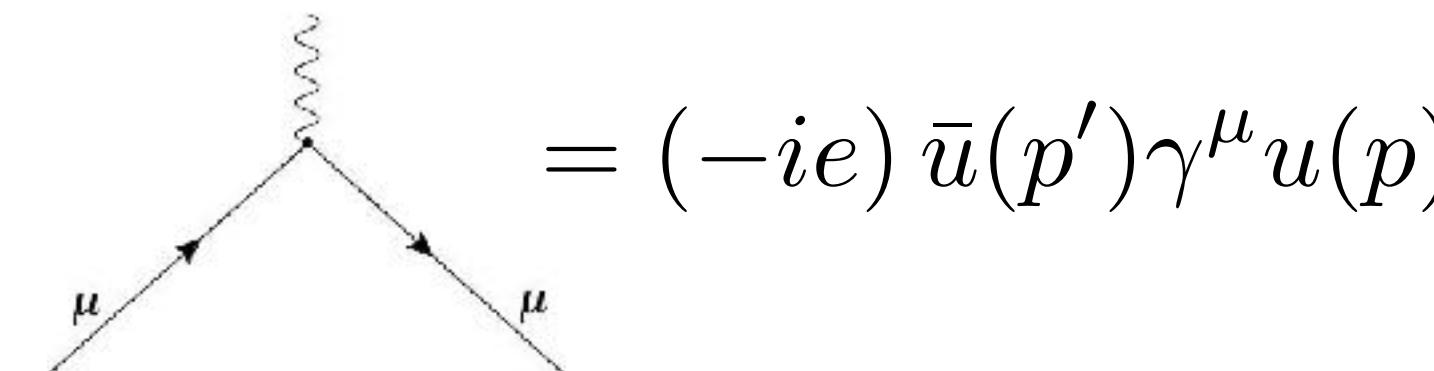
- Introduction
- Theory vs experiment
- Muon g-2 Theory Initiative
- g-2 SM contributions
- Dispersive, data driven methods for
  - Hadronic Vacuum Polarization (HVP)
  - Hadronic Light-by-Light scattering (HLbL)
- Lattice HVP
- Lattice HLbL
- Summary and Outlook



# Anomalous magnetic moment

The magnetic moment of charged leptons ( $e, \mu, \tau$ ):  $\vec{\mu} = g \frac{e}{2m} \vec{S}$

Dirac (leading order):  $g = 2$



Quantum effects (loops):

All SM particles contribute

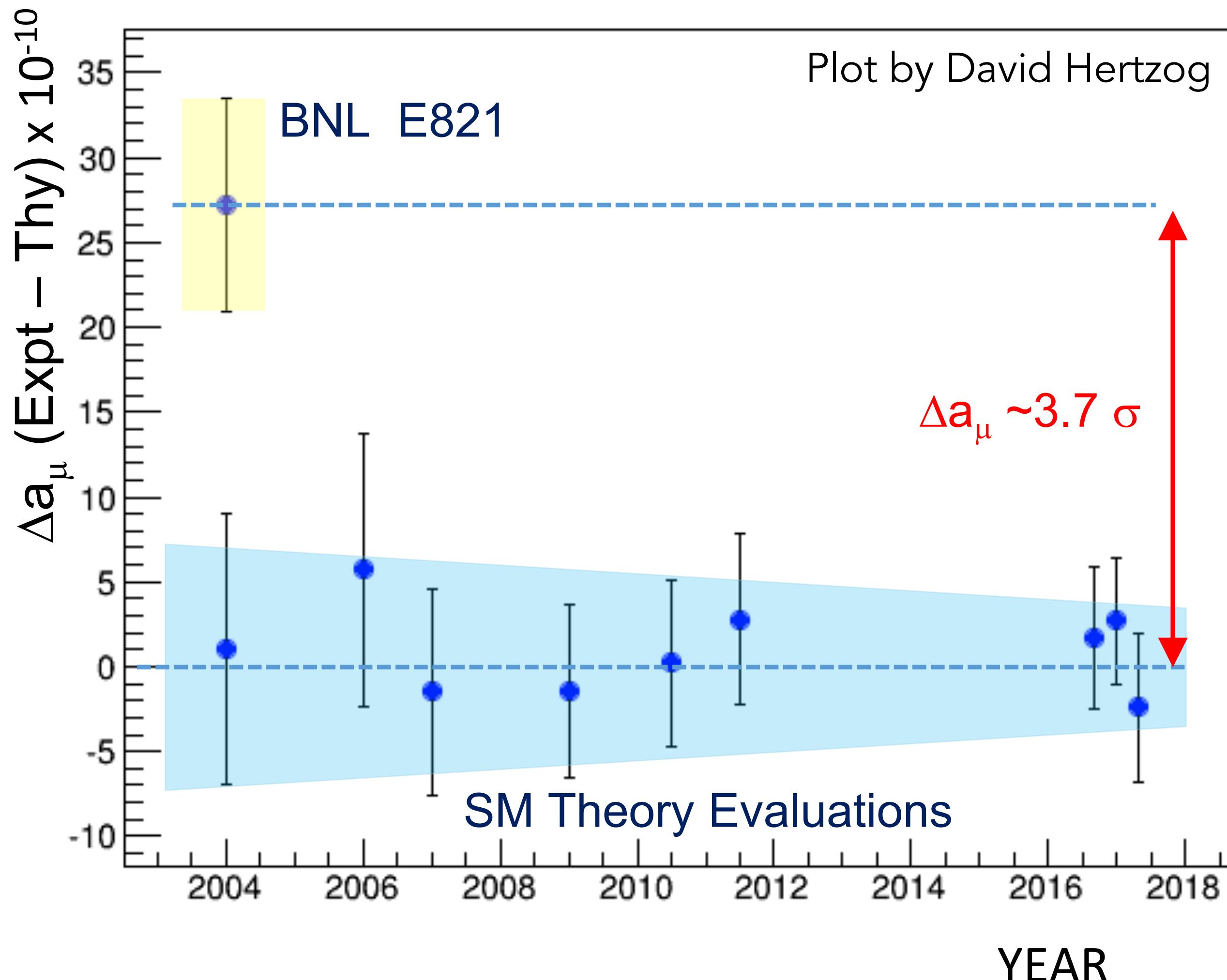
$$= (-i e) \bar{u}(p') \left[ \gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2m} F_2(q^2) \right] u(p)$$

Note:  $F_1(0) = 1$  and  $g = 2 + 2 F_2(0)$

Anomalous magnetic moment:

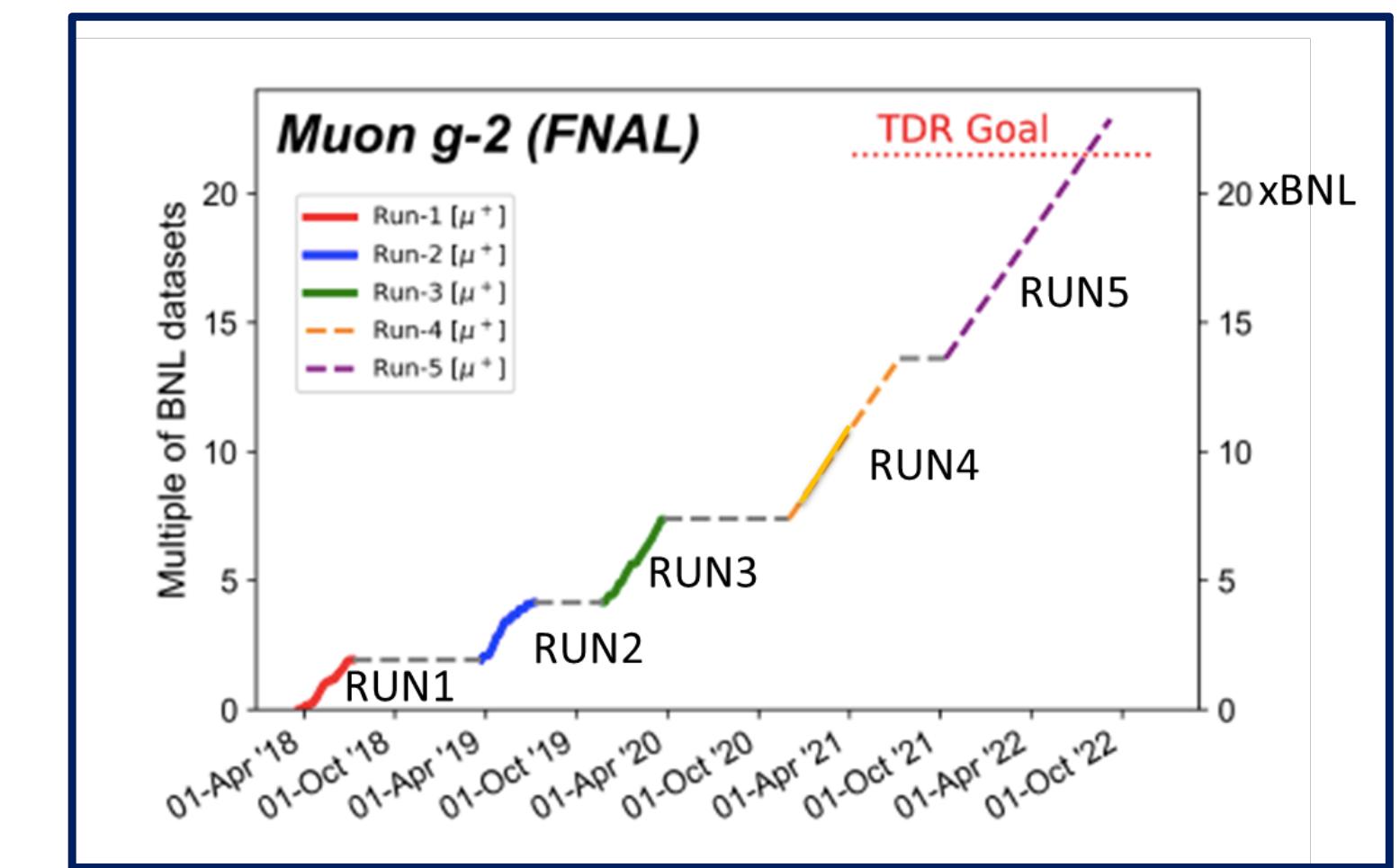
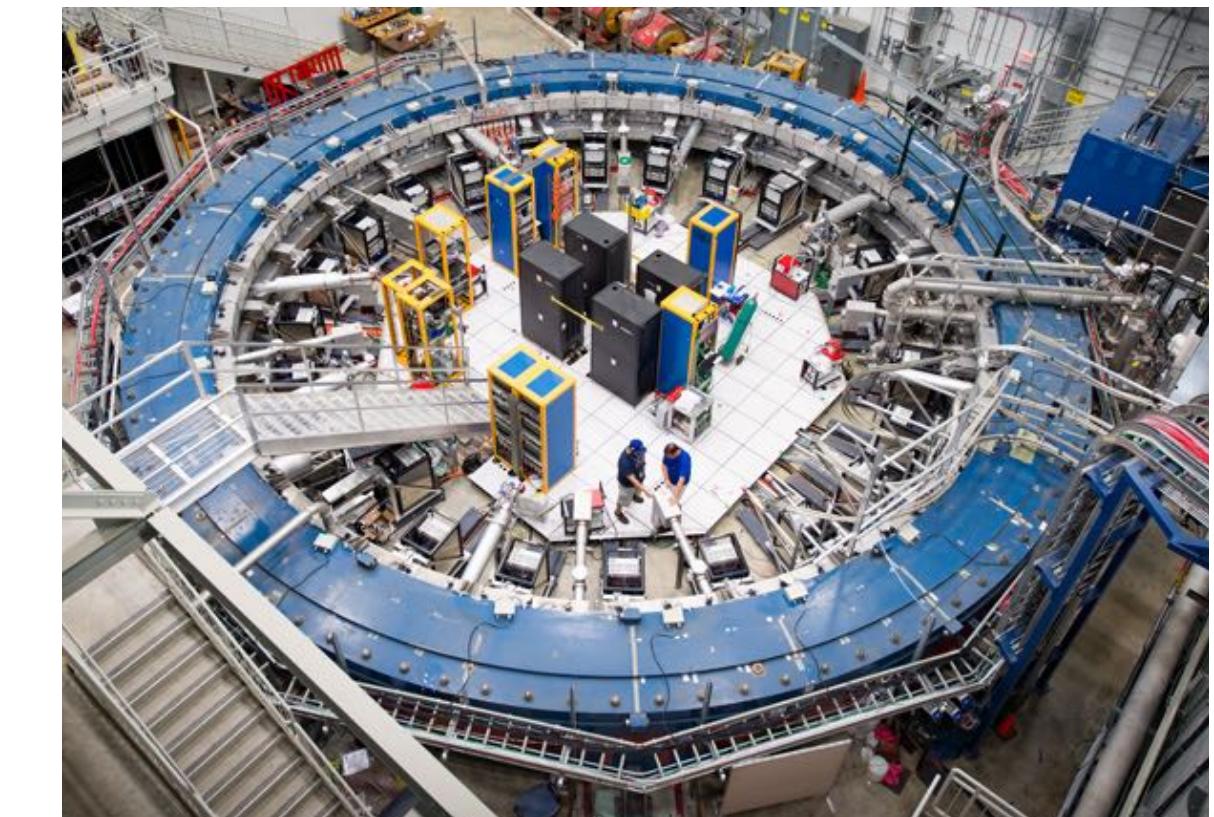
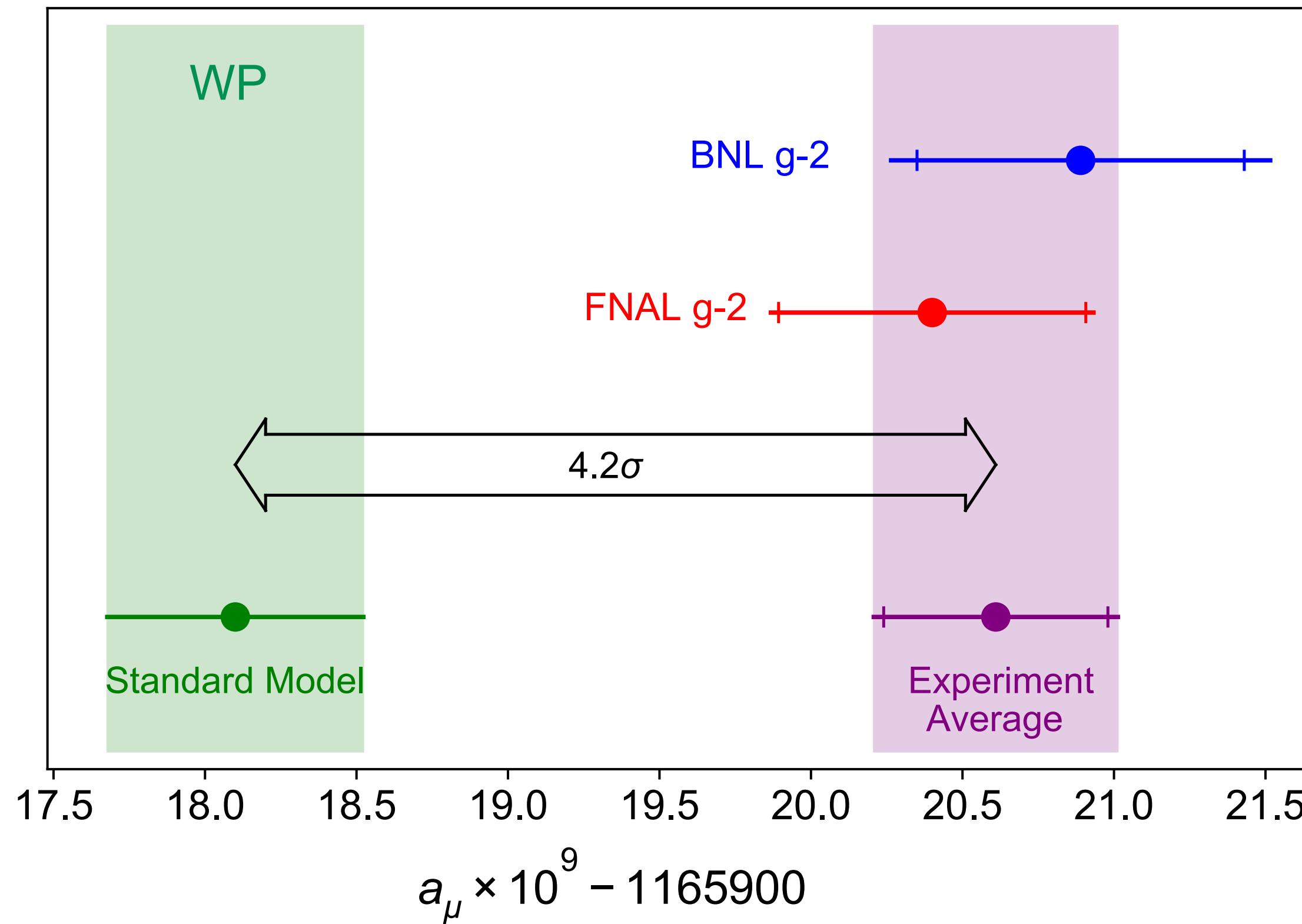
$$a \equiv \frac{g - 2}{2} = F_2(0)$$

# Muon g-2: history of experiment vs theory



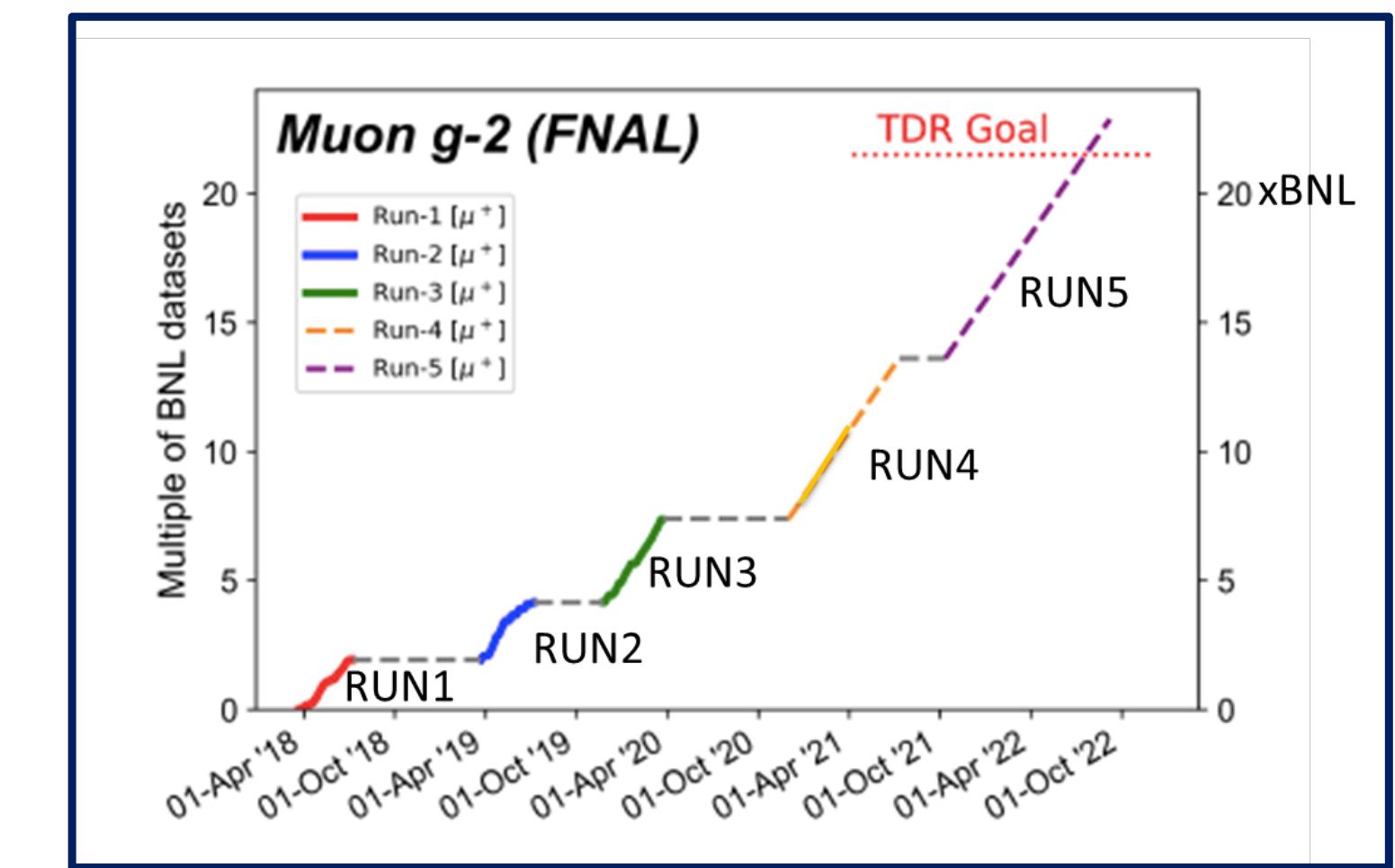
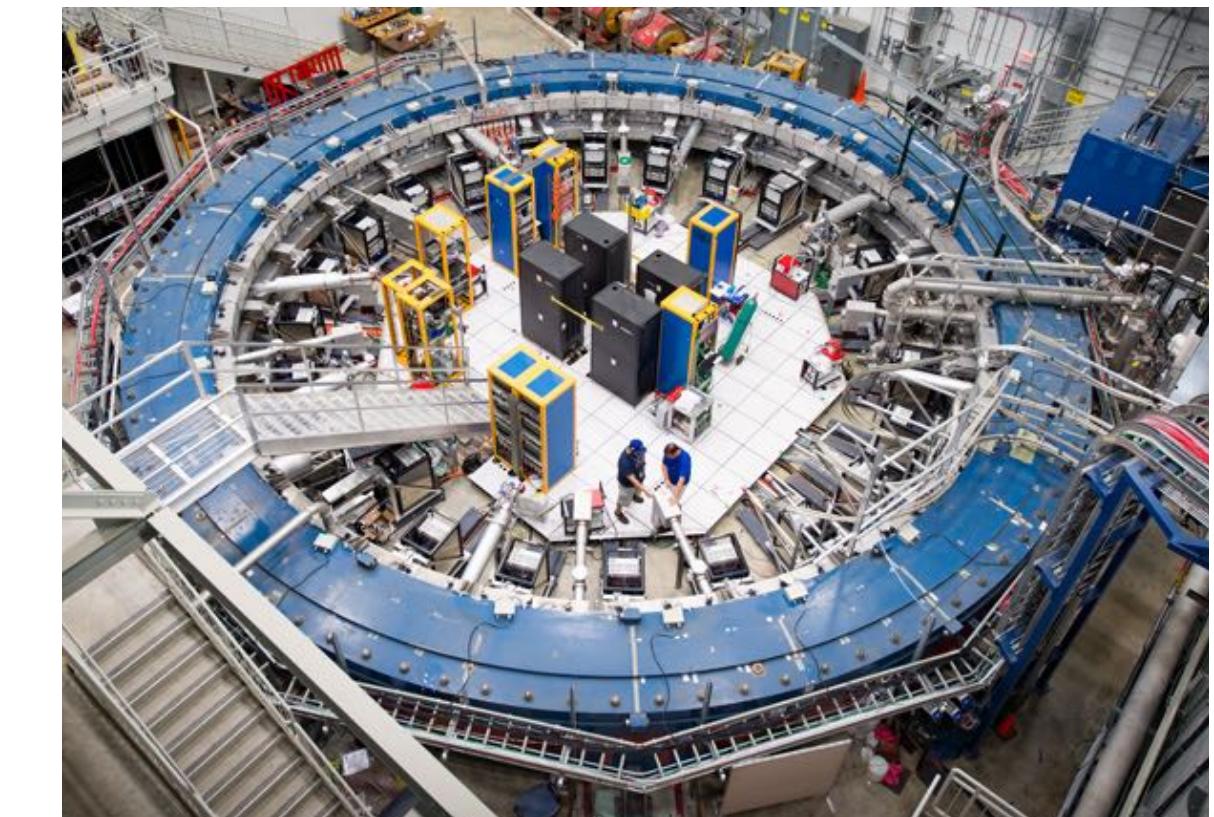
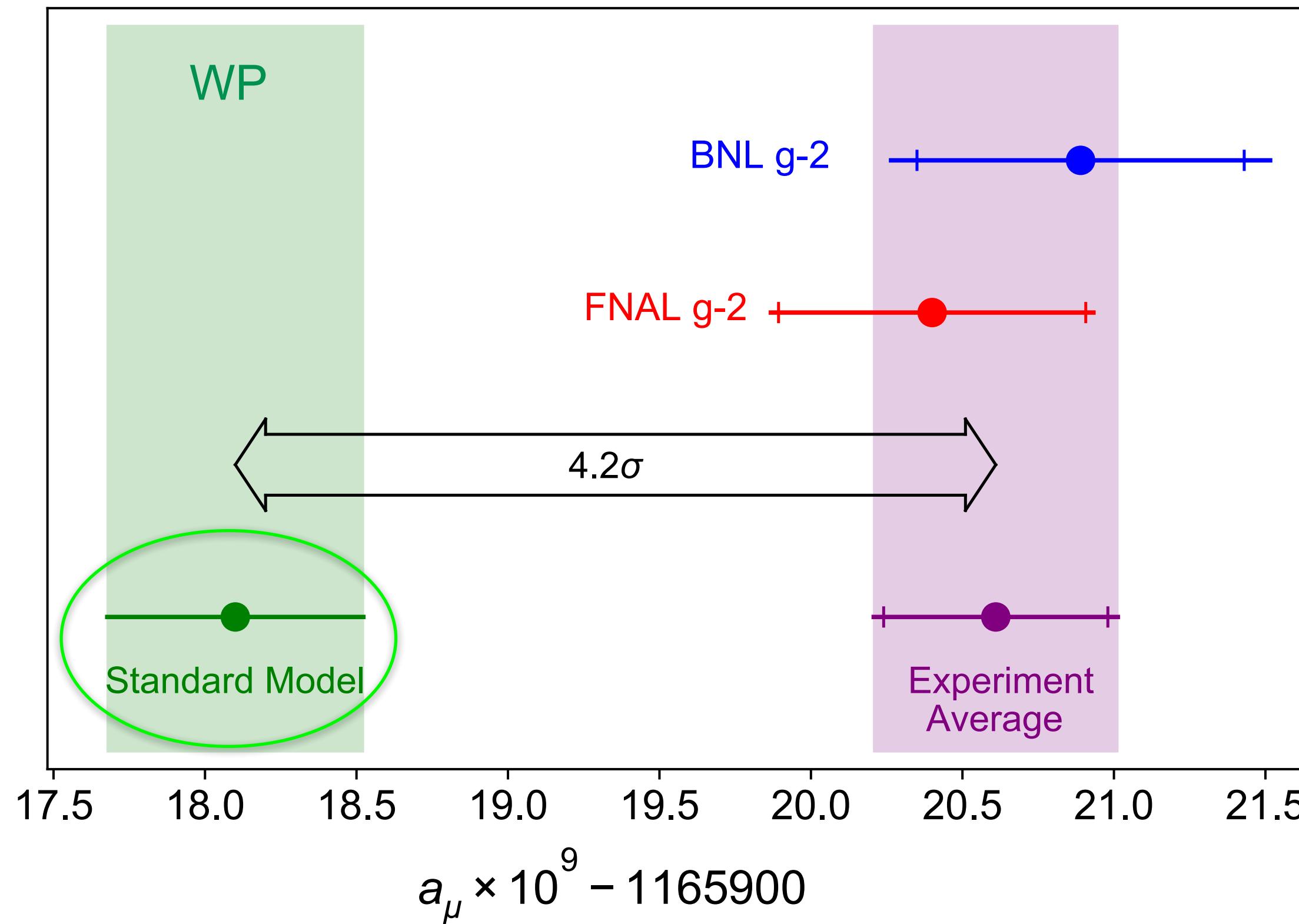
# Muon g-2: experiment

- ⌚ The Fermilab experiment released the measurement result from their run 1 data on 7 April 2021.  
[B. Abi et al, [Phys. Rev. Lett. 124, 141801 \(2021\)](#)]
- ⌚ Analysis of runs 2 and 3 is now underway.



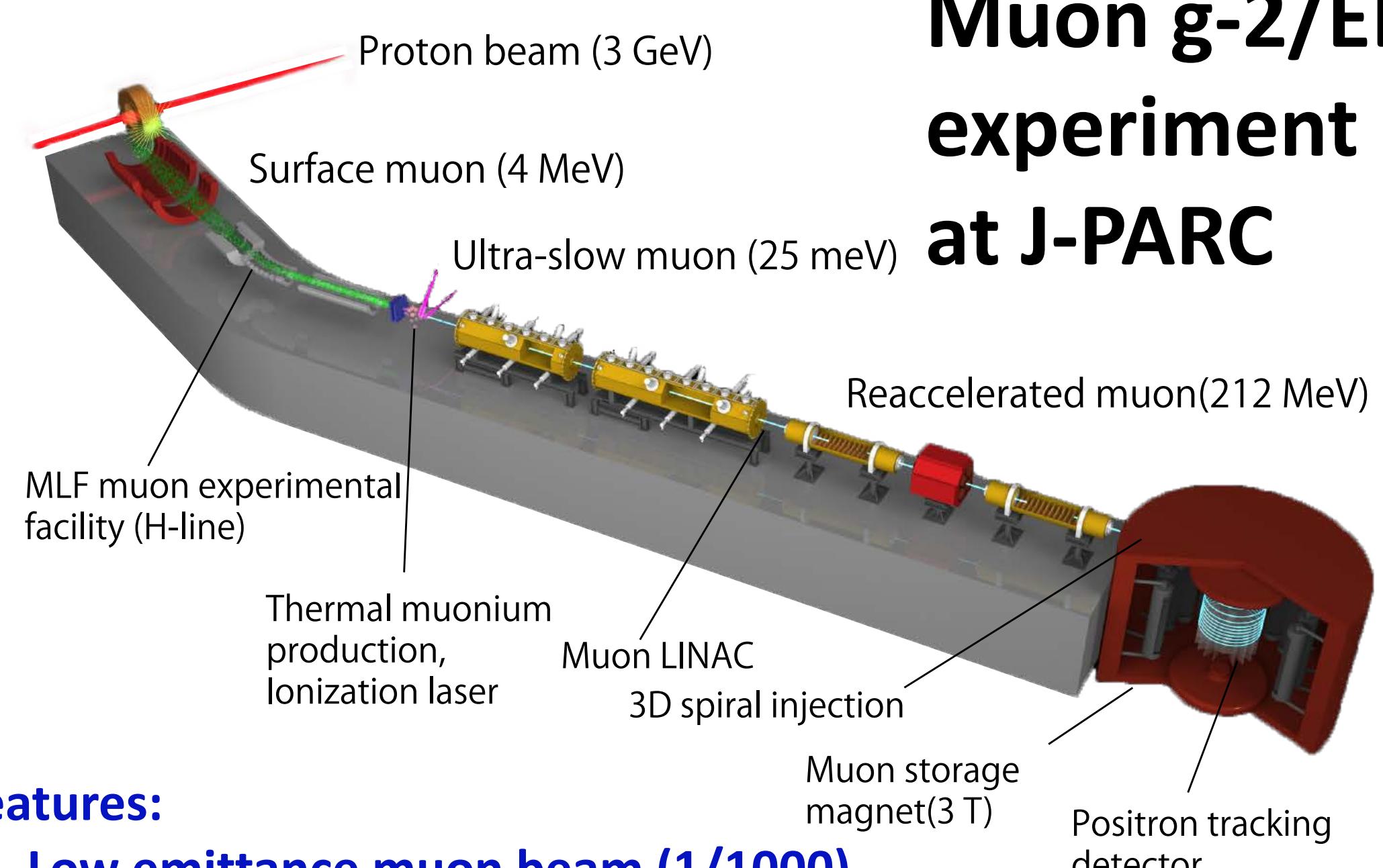
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[B. Abi et al, [Phys. Rev. Lett. 124, 141801 \(2021\)](#)]
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# Muon g-2: experiment

T. Mibe for E34 @ INT g-2 workshop



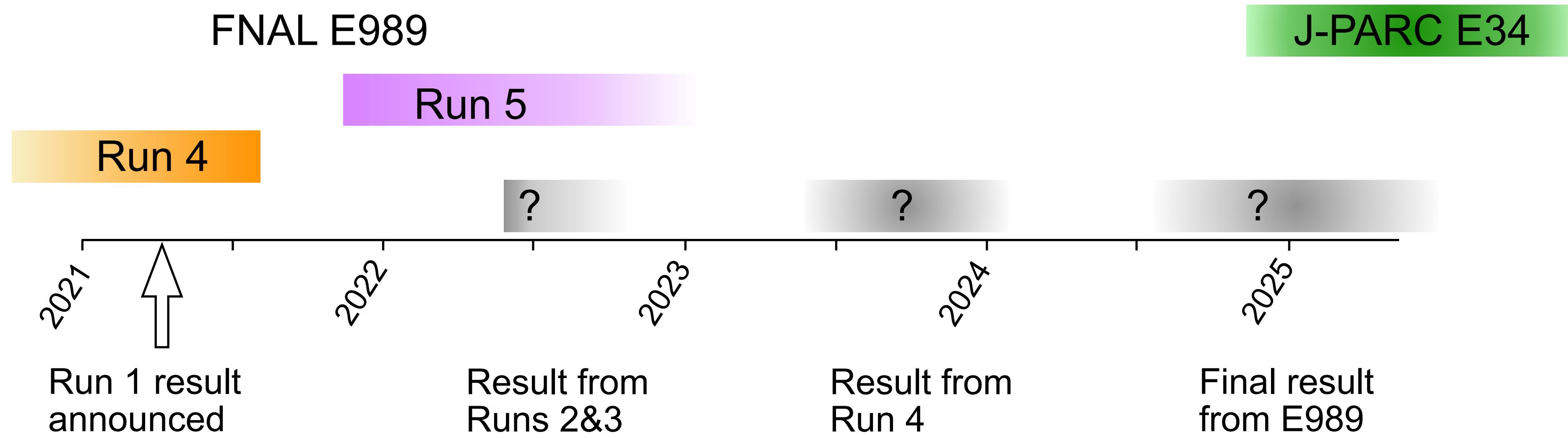
## Muon g-2/EDM experiment at J-PARC

### Features:

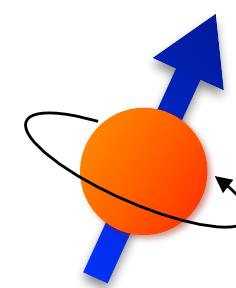
- **Low emittance muon beam (1/1000)**
- **No strong focusing (1/1000) & good injection eff. (x10)**
- **Compact storage ring (1/20)**
- **Tracking detector with large acceptance**
- **Completely different from BNL/FNAL method**

- 2018:  
Stage II approval by IPNS and IMSS directors.
- March 2019:  
Endorsed by KEK-SAC as a near-term priority
- 2020:  
Funding request
- 2024+:  
data taking runs

# Timeline of g-2 experiments

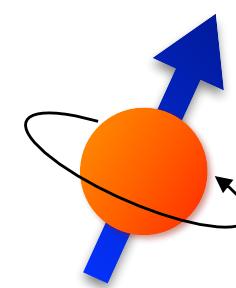


- The Fermilab and J-PARC g-2 experiments will yield increasingly precise measurements of  $a_\mu$  through the rest of this decade.
- In addition, future efforts may include negative muon runs at Fermilab and new muon EDM experiments.
  - ➡ Continue efforts to improve on the precision of the SM predictions.
  - ➡ Theory Initiative: plan to continue efforts and update WP with new SM predictions



# Muon g-2 Theory Initiative

- ➊ Maximize the impact of the Fermilab and J-PARC experiments  
➡ **quantify and reduce the theoretical uncertainties on the hadronic corrections**
- ➋ summarize the theory status and assess reliability of uncertainty estimates
- ➌ organize workshops to bring the different communities together:  
[First plenary workshop @ Fermilab: 3-6 June 2017](#)  
[HVP workshop @ KEK: 12-14 February 2018](#)  
[HLbL workshop @ U Connecticut: 12-14 March 2018](#)  
[Second plenary workshop @ HIM \(Mainz\): 18-22 June 2018](#)  
[Third plenary workshop @ INT \(Seattle\): 9-13 September 2019](#)  
[Lattice HVP at high precision workshop \(virtual\): 16-20 November 2020](#)  
[Fourth plenary workshop @ KEK \(virtual\): 28 June - 02 July 2021](#)  
Fifth plenary workshop @ Higgs Centre (Edinburgh): early September 2022 (tentative)
- ➍ 1<sup>st</sup> White Paper published in 2020 (132 authors, 82 institutions, 21 countries)  
[T. Aoyama et al, [arXiv:2006.04822](#), Phys. Repts. 887 (2020) 1-166.]
- ➎ 2<sup>nd</sup> White Paper: First discussions @ KEK meeting



# Muon g-2 Theory Initiative

## Steering Committee

Simon Eidelman



[photo by Hartmut Wittig]

(1948-2021)

- Gilberto Colangelo (Bern)
- Michel Davier (Orsay)
- Aida El-Khadra (UIUC & Fermilab)
- Martin Hoferichter (Bern)
- Christoph Lehner (Regensburg University & BNL)
- Laurent Lellouch (Marseille)
- Tsutomu Mibe (KEK) J-PARC Muon g-2/EDM experiment
- Lee Roberts (Boston) Fermilab Muon g-2 experiment
- Thomas Teubner (Liverpool)
- Hartmut Wittig (Mainz)

website: <https://muon-gm2-theory.illinois.edu>

# Muon g-2: SM contributions

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$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

QED

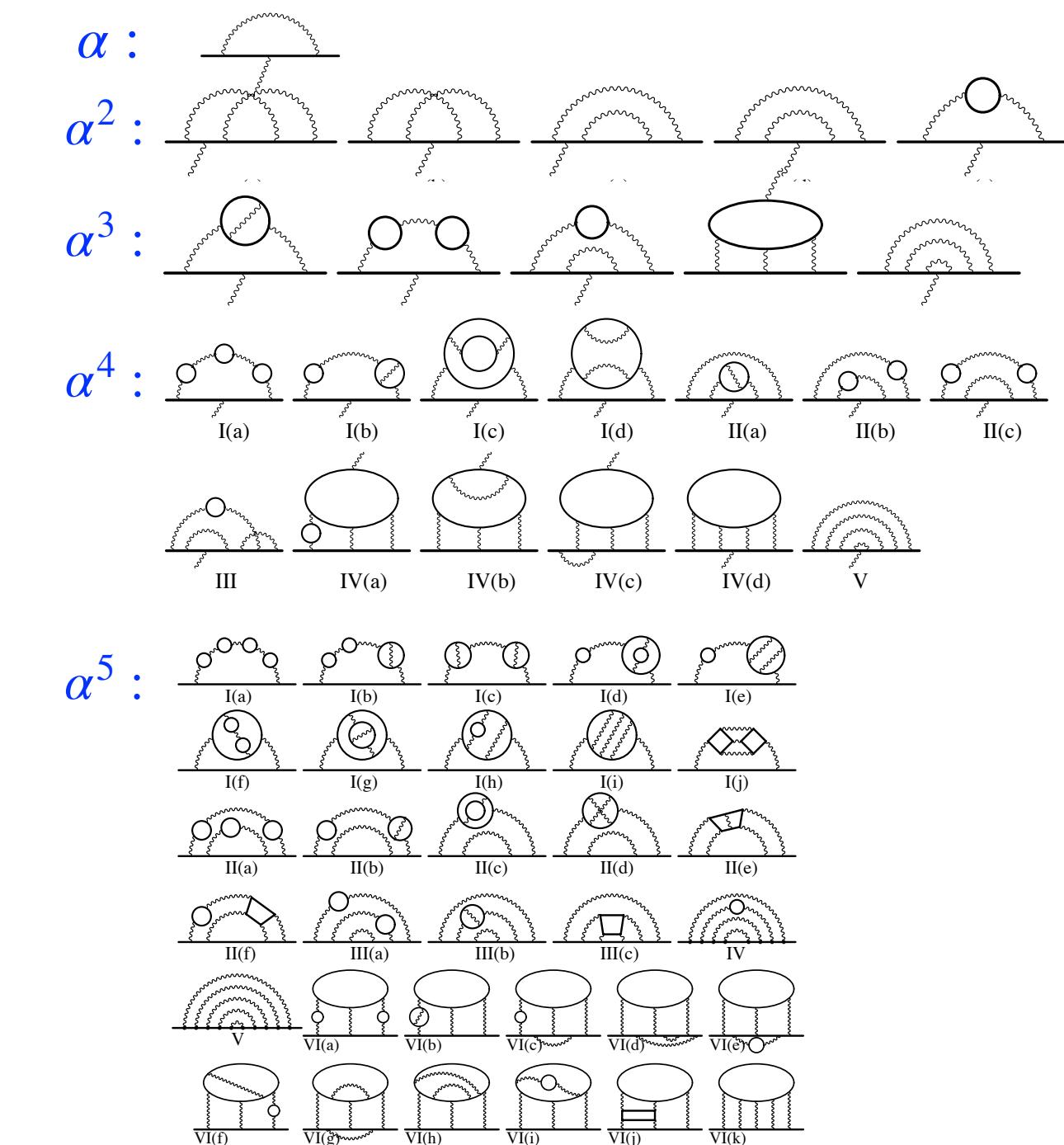
$$a_\mu(\text{QED}) = A_1 + A_2 \left( \frac{m_\mu}{m_e} \right) + A_2 \left( \frac{m_\mu}{m_\tau} \right) + A_3 \left( \frac{m_\mu}{m_e}, \frac{m_\mu}{m_\tau} \right)$$

$$A_i = \sum_{n=0} \left( \frac{\alpha}{\pi} \right)^n A_i^{2n}$$

$n$	# of diagrams	Contribution $\times 10^{11}$
1	1	116140973.32
2	7	413 217.63
3	71	30141.90
4	891	381.00
5	12672	5.08

$$a_\mu(\text{QED}) = 116\,584\,718.9(1) \times 10^{-11}$$

[T. Aoyama et al, arXiv:1205.5370, PRL;  
 T. Aoyama, T. Kinoshita, M. Nio, Atoms 7 (1) (2019) 28]

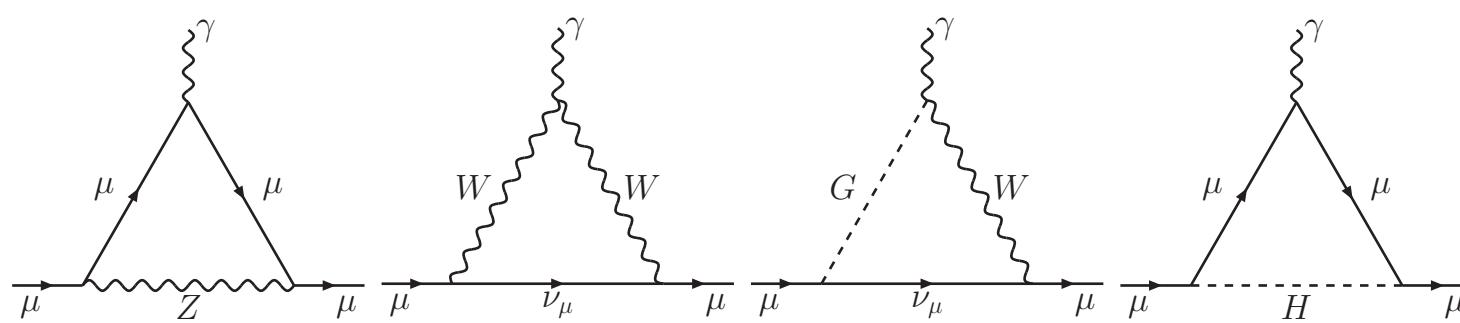


# Muon g-2: SM contributions

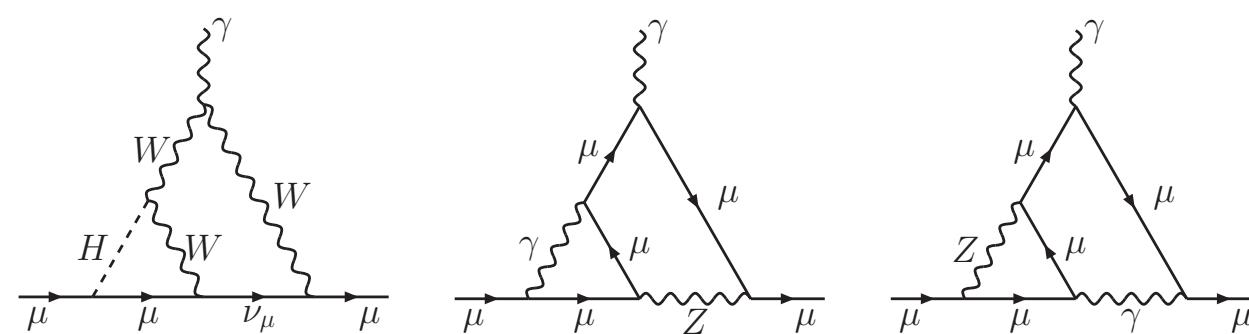
$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

Electroweak  
(contributions from W,Z,H,.. bosons)

1-loop

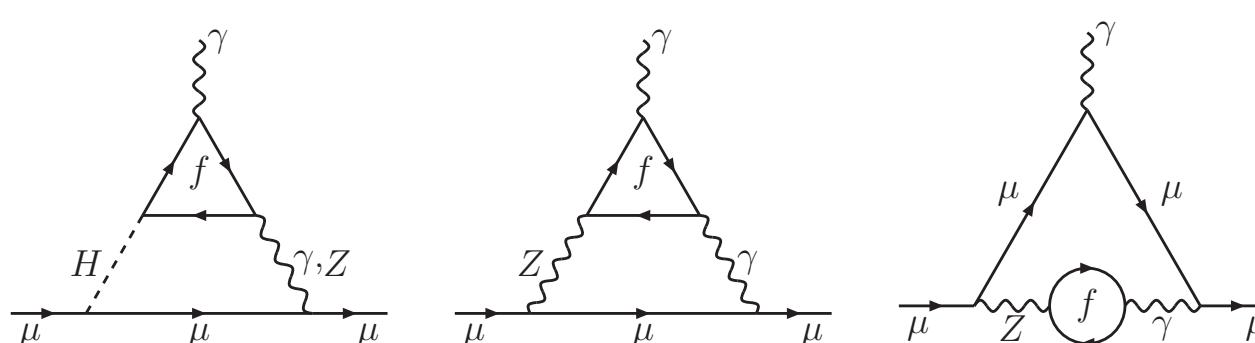


2-loop



$$a_\mu(\text{EW}) = 153.6(1.0) \times 10^{-11}$$

[A. Czarnecki et al, hep-ph/0212229, PRD;  
C. Gnendinger et al, arXiv:1306.5546, PRD]



# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

leading hadronic



- ◆ The hadronic contributions are written as:

$$\begin{aligned} a_\ell(\text{hadronic}) = & \quad a_\ell^{\text{HVP, LO}} + a_\ell^{\text{HVP, NLO}} + a_\ell^{\text{HVP, NNLO}} + \dots \\ & + a_\ell^{\text{HLbL}} \quad + a_\ell^{\text{HLbL, NLO}} + \dots \end{aligned}$$

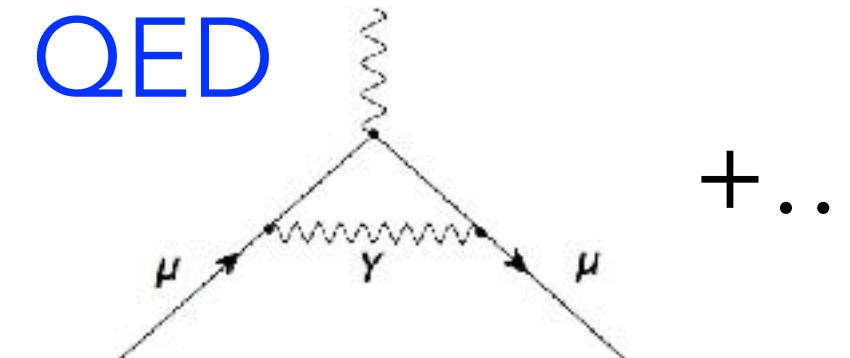
$$\alpha^2 \qquad \alpha^3 \qquad \alpha^4$$

$$\sim 10^{-7}$$

# Muon g-2: SM contributions

$$a_\mu = a_\mu(\text{QED}) + a_\mu(\text{EW}) + a_\mu(\text{hadronic})$$

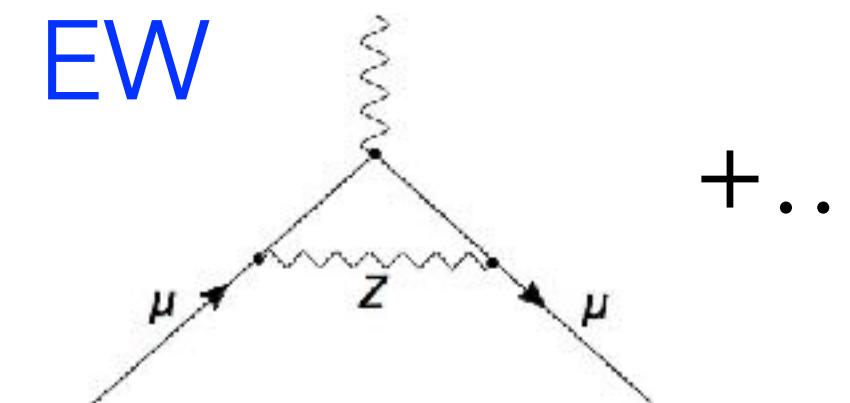
QED



$$116\,584\,718.9(1) \times 10^{-11}$$

0.001 ppm

EW



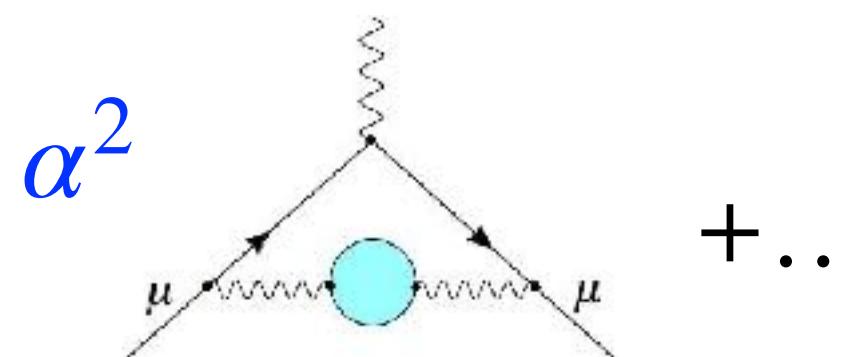
$$153.6(1.0) \times 10^{-11}$$

0.01 ppm

Hadronic...

...Vacuum Polarization (HVP)

$\alpha^2$



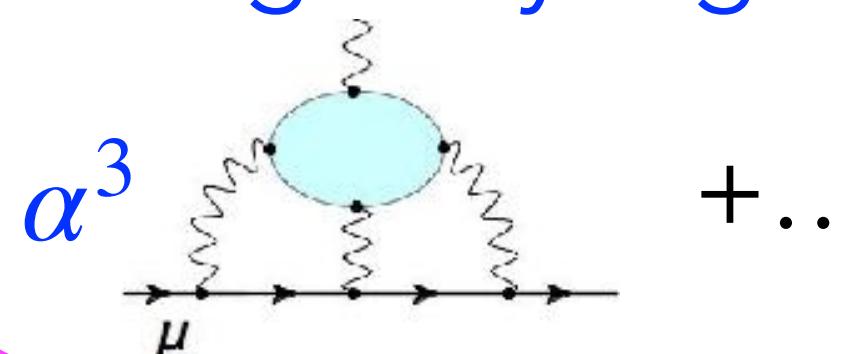
$$6845(40) \times 10^{-11}$$

[0.6%]

0.34 ppm

...Light-by-Light (HLbL)

$\alpha^3$



$$92(18) \times 10^{-11}$$

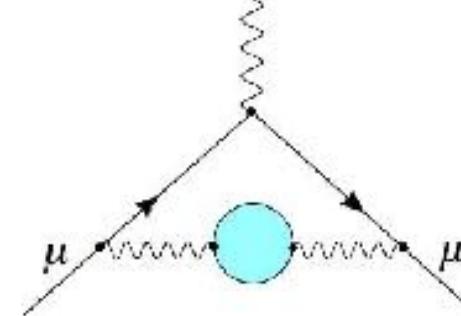
[20%]

0.15 ppm

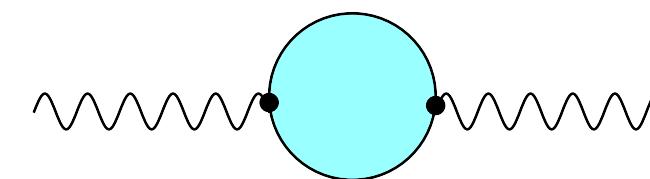
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# Hadronic vacuum polarization



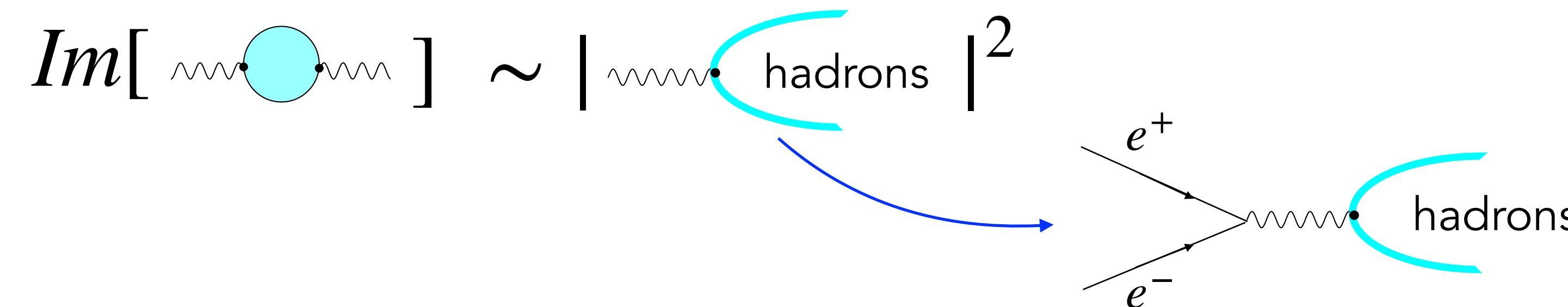
$$\hat{\Pi}(q^2) = \Pi(q^2) - \Pi(0)$$

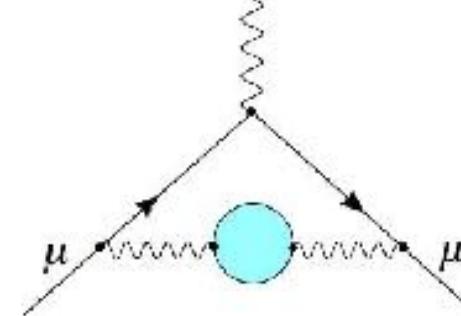
$$\Pi_{\mu\nu} = \int d^4x e^{iqx} \langle j_\mu(x) j_\nu(0) \rangle = (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(q^2)$$

Leading order HVP correction:

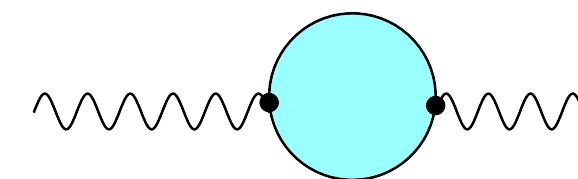
$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int dq^2 \omega(q^2) \hat{\Pi}(q^2)$$

- Use optical theorem and dispersion relation to rewrite the integral in terms of the hadronic  $e^+e^-$  cross section:





# Hadronic vacuum polarization



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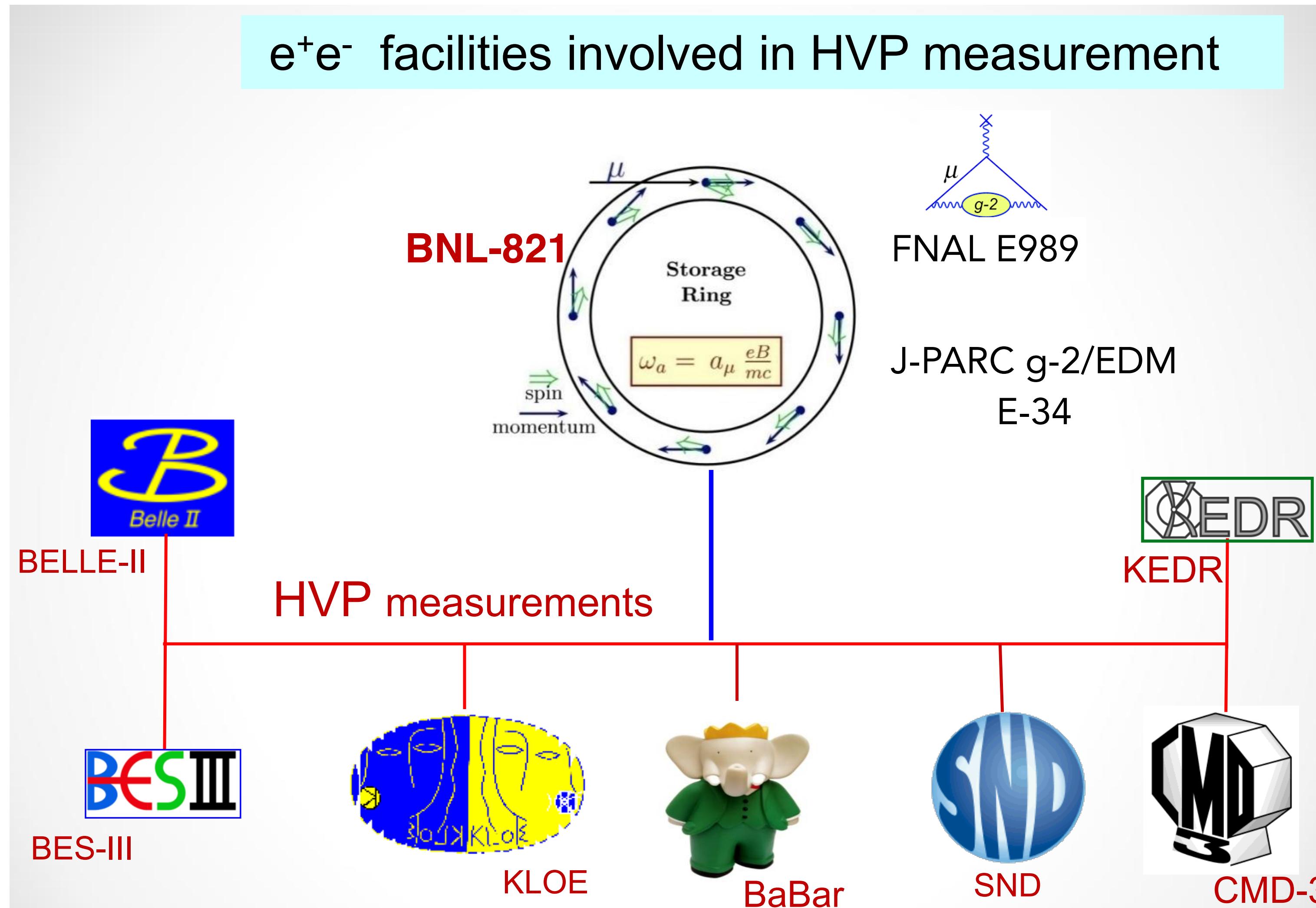
$$a_\mu^{\text{HVP,LO}} = \frac{m_\mu^2}{12\pi^3} \int ds \frac{\hat{K}(s)}{s} \sigma_{\text{exp}}(s)$$

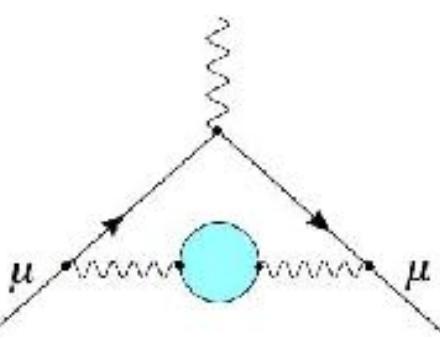
Dominant contributions from low energies  
 $\pi^+\pi^-$  channel: 73% of total

- Use direct integration method, summing up cross sections for all possible hadronic channels up to  $\sim 2$  GeV

# Experimental Inputs to HVP

S. Serednyakov (for SND) @ HVP KEK workshop



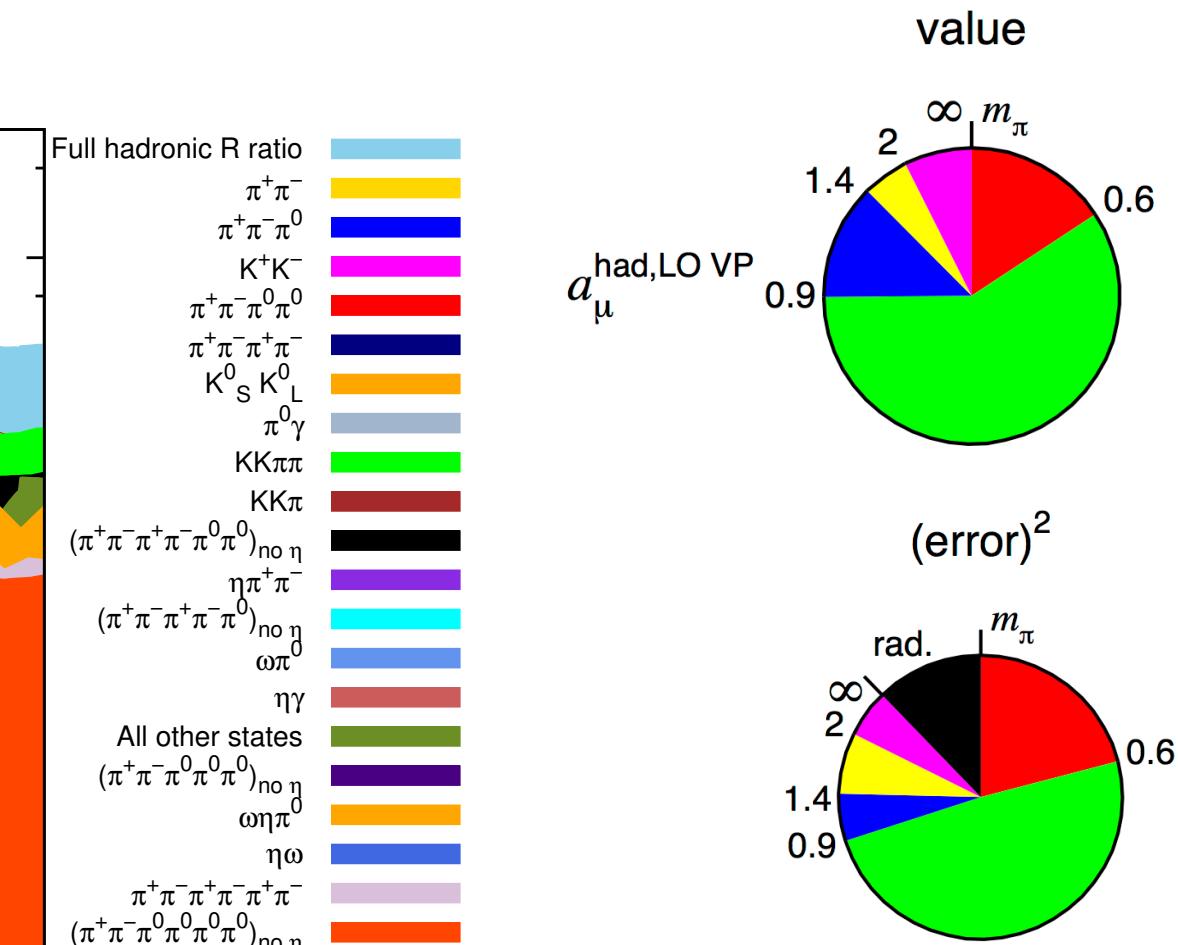
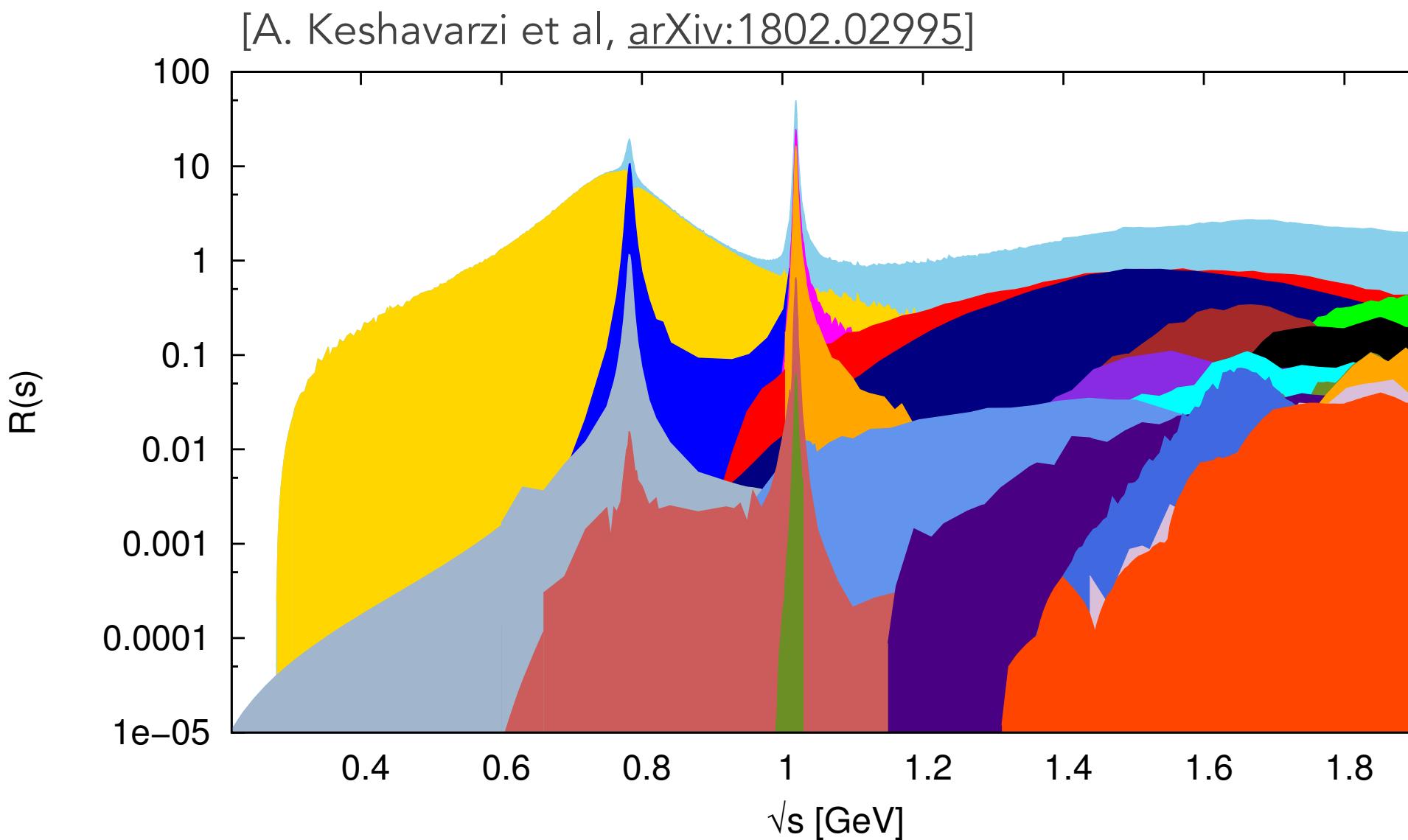


# HVP: data-driven

Z. Zhang for DHMZ @ INT g-2 workshop

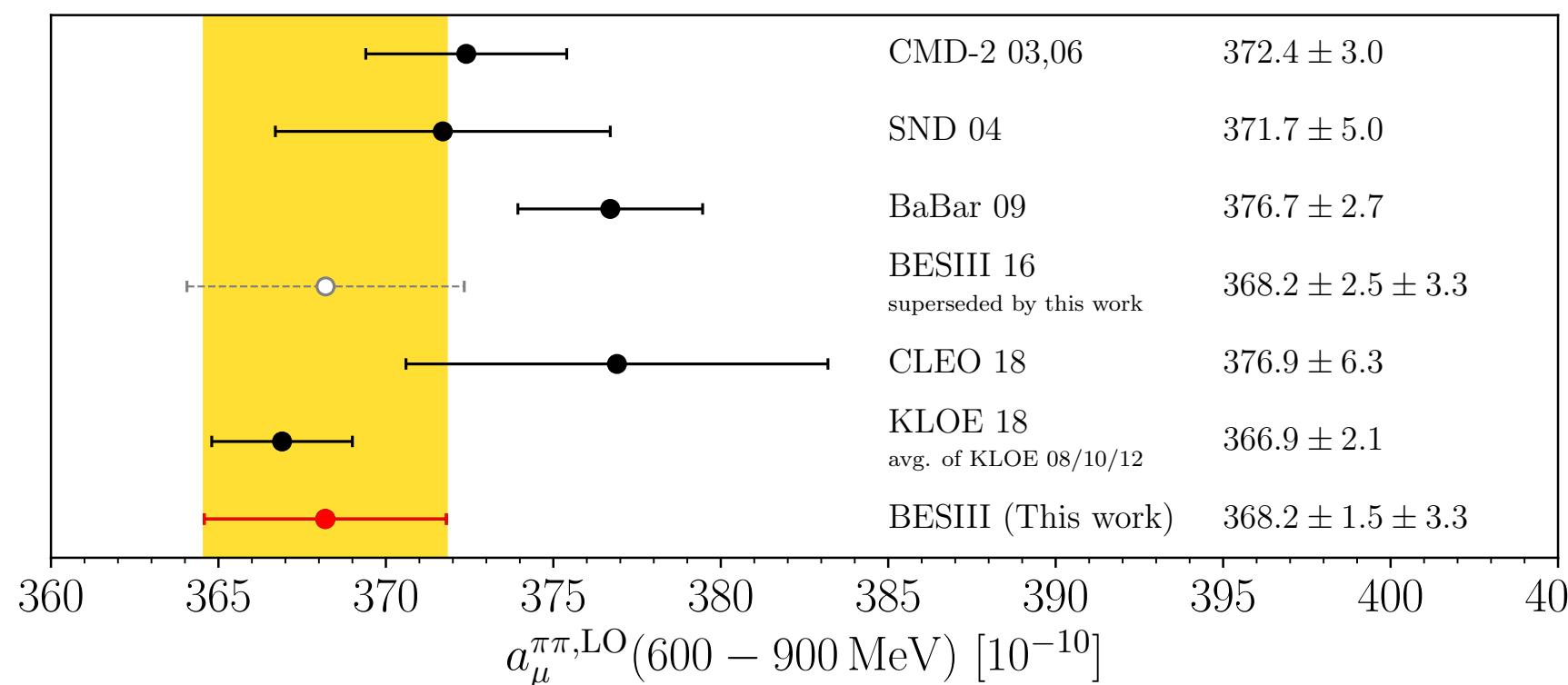
[M. Davier et al, arXiv:1908.00921]

Channel	$a_\mu^{\text{had}, \text{LO}} [10^{-10}]$
$\pi^0\gamma$	$4.29 \pm 0.06 \pm 0.04 \pm 0.07$
$\eta\gamma$	$0.65 \pm 0.02 \pm 0.01 \pm 0.01$
$\pi^+\pi^-$	$507.80 \pm 0.83 \pm 3.19 \pm 0.60$
$\pi^+\pi^-\pi^0$	$46.20 \pm 0.40 \pm 1.10 \pm 0.86$
$2\pi^+2\pi^-$	$13.68 \pm 0.03 \pm 0.27 \pm 0.14$
$\pi^+\pi^-2\pi^0$	$18.03 \pm 0.06 \pm 0.48 \pm 0.26$
$2\pi^+2\pi^-\pi^0$ ( $\eta$ excl.)	$0.69 \pm 0.04 \pm 0.06 \pm 0.03$
$\pi^+\pi^-3\pi^0$ ( $\eta$ excl.)	$0.49 \pm 0.03 \pm 0.09 \pm 0.00$
$3\pi^+3\pi^-$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$
$2\pi^+2\pi^-2\pi^0$ ( $\eta$ excl.)	$0.71 \pm 0.06 \pm 0.07 \pm 0.14$
$\pi^+\pi^-4\pi^0$ ( $\eta$ excl., isospin)	$0.08 \pm 0.01 \pm 0.08 \pm 0.00$
$\eta\pi^+\pi^-$	$1.19 \pm 0.02 \pm 0.04 \pm 0.02$
$\eta\omega$	$0.35 \pm 0.01 \pm 0.02 \pm 0.01$
$\eta\pi^+\pi^-\pi^0$ (non- $\omega, \phi$ )	$0.34 \pm 0.03 \pm 0.03 \pm 0.04$
$\eta 2\pi^+2\pi^-$	$0.02 \pm 0.01 \pm 0.00 \pm 0.00$
$\omega\eta\pi^0$	$0.06 \pm 0.01 \pm 0.01 \pm 0.00$
$\omega\pi^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.94 \pm 0.01 \pm 0.03 \pm 0.00$
$\omega(\pi\pi)^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.07 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega$ (non- $3\pi, \pi\gamma, \eta\gamma$ )	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$
$K^+K^-$	$23.08 \pm 0.20 \pm 0.33 \pm 0.21$
$K_SK_L$	$12.82 \pm 0.06 \pm 0.18 \pm 0.15$
$\phi$ (non- $K\bar{K}$ , $3\pi, \pi\gamma, \eta\gamma$ )	$0.05 \pm 0.00 \pm 0.00 \pm 0.00$
$K\bar{K}\pi$	$2.45 \pm 0.05 \pm 0.10 \pm 0.06$
$K\bar{K}2\pi$	$0.85 \pm 0.02 \pm 0.05 \pm 0.01$
$K\bar{K}3\pi$ (estimate)	$-0.02 \pm 0.01 \pm 0.01 \pm 0.00$
$\eta\phi$	$0.33 \pm 0.01 \pm 0.01 \pm 0.00$
$\eta K\bar{K}$ (non- $\phi$ )	$0.01 \pm 0.01 \pm 0.01 \pm 0.00$
$\omega K\bar{K}$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega 3\pi$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.06 \pm 0.01 \pm 0.01 \pm 0.01$
$7\pi$ ( $3\pi^+3\pi^-\pi^0$ + estimate)	$0.02 \pm 0.00 \pm 0.01 \pm 0.00$
$J/\psi$ (BW integral)	$6.28 \pm 0.07$
$\psi(2S)$ (BW integral)	$1.57 \pm 0.03$
$R$ data [3.7 – 5.0] GeV	$7.29 \pm 0.05 \pm 0.30 \pm 0.00$
$R_{\text{QCD}}$ [1.8 – 3.7 GeV] <sub>uds</sub>	$33.45 \pm 0.28 \pm 0.65$ <sub>dual</sub>
$R_{\text{QCD}}$ [5.0 – 9.3 GeV] <sub>udsc</sub>	$6.86 \pm 0.04$
$R_{\text{QCD}}$ [9.3 – 12.0 GeV] <sub>udscb</sub>	$1.21 \pm 0.01$
$R_{\text{QCD}}$ [12.0 – 40.0 GeV] <sub>udscb</sub>	$1.64 \pm 0.00$
$R_{\text{QCD}}$ [> 40.0 GeV] <sub>udscb</sub>	$0.16 \pm 0.00$
$R_{\text{QCD}}$ [> 40.0 GeV] <sub>t</sub>	$0.00 \pm 0.00$
<b>Sum</b>	$693.9 \pm 1.0 \pm 3.4 \pm 1.6 \pm 0.1_\psi \pm 0.7_{\text{QCD}}$



Tensions between BaBar and KLOE data sets:

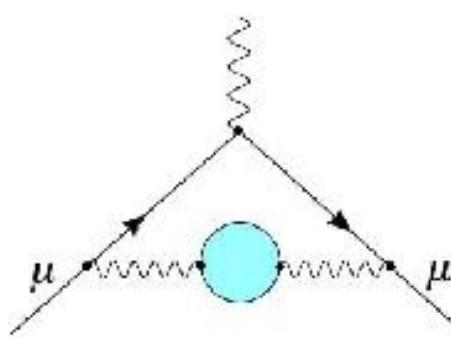
[M. Ablikim et al (BES III), arXiv:2009.05011]



- Cross checks using analyticity and unitarity relating pion form factor to  $\pi\pi$  scattering

- Combinations of data sets affected by tensions

➡ conservative merging procedure



# HVP: data-driven

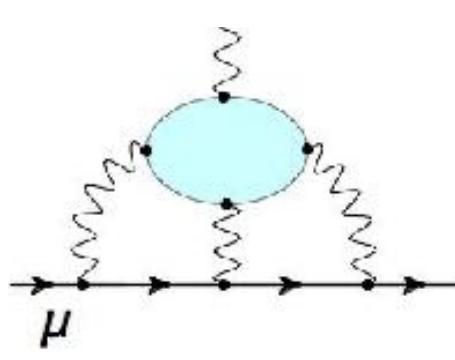
## Conservative merging procedure

[B. Malaescu @ INT g-2 workshop]

to obtain a realistic assessment of the underlying uncertainties:

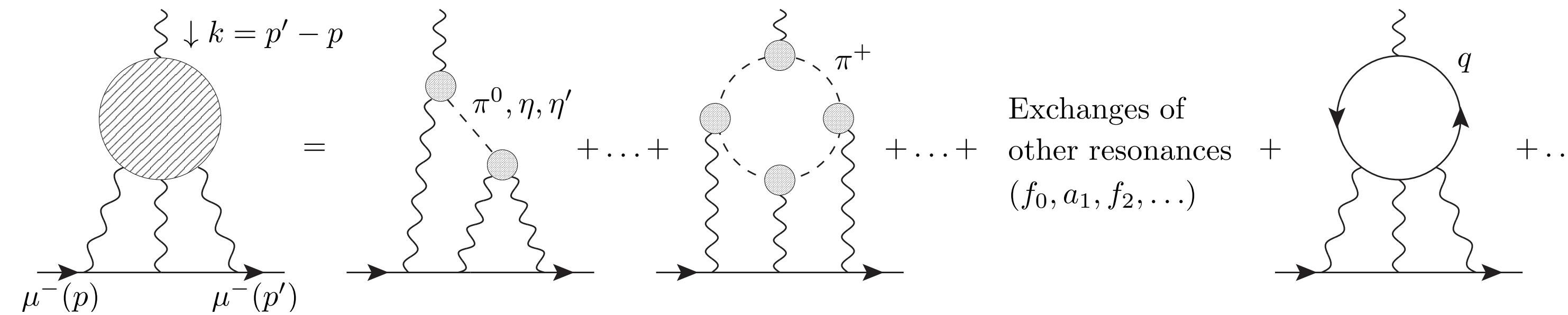
- account for tensions between data sets
- account for differences in methodologies for compilation of experimental inputs
- include correlations between systematic errors
- constraints and cross checks from unitarity & analyticity constraints  
[Colangelo et al, 2018; Anantharayan et al, 2018; Davier et al, 2019; Hoferichter et al, 2019]
- Full NLO radiative corrections [Campanario et al, 2019]

$$\Rightarrow a_\mu^{\text{HVP,LO}} = 693.1 (2.8)_{\text{exp}} (2.8)_{\text{sys}} (0.7)_{\text{DV+pQCD}} \times 10^{-10} = 693.1 (4.0) \times 10^{-10}$$

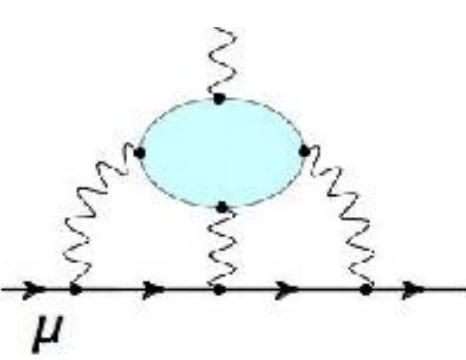


# Hadronic Light-by-light

Hadronic light-by-light: Target:  $\lesssim 10\%$  total error

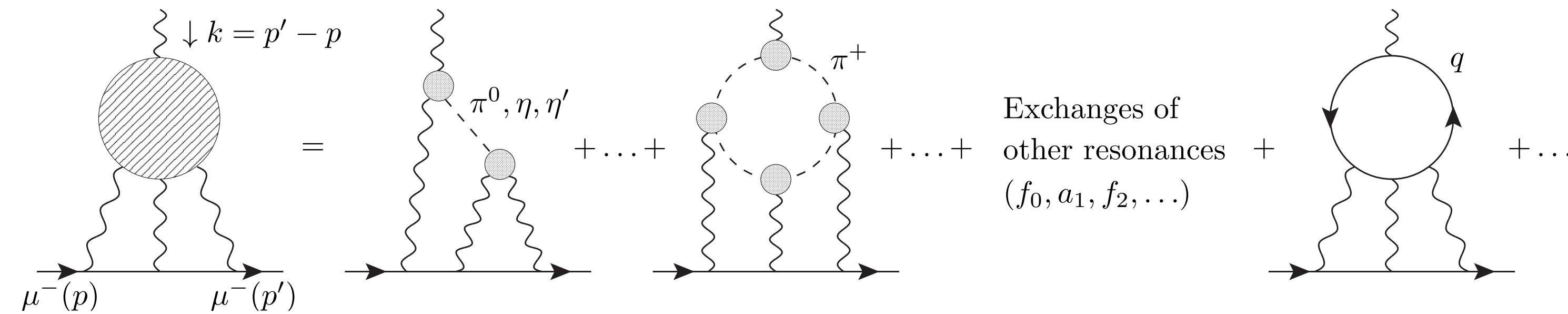


- ◆ previous estimates “Glasgow consensus” use models of QCD
- ◆ used to evaluate individual contributions to HLbL scattering tensor
- ◆ theory error not well determined and not improvable



# Hadronic Light-by-light

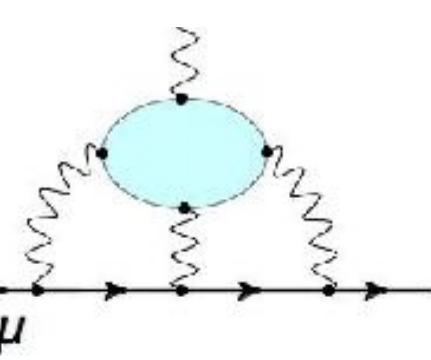
Hadronic light-by-light: Target:  $\lesssim 10\%$  total error



## Dispersive approach:

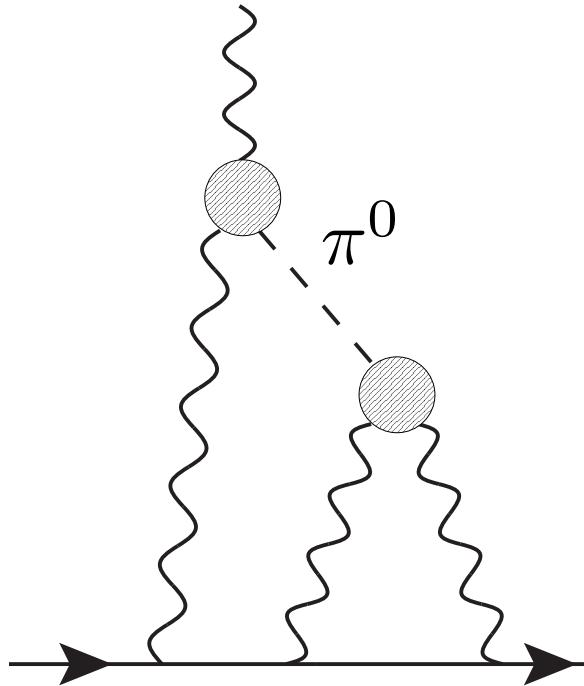
[Colangelo et al, 2014; Pauk & Vanderhaegen 2014; ...]

- ◆ model independent
- ◆ significantly more complicated than for HVP
- ◆ provides a framework for data-driven evaluations
- ➡ ◆ can also use lattice results as inputs



# HLbL: dispersive

Three independent results for the pion pole contribution:  
[G. Colangelo @ INT g-2 workshop]



- ▶ Dispersive calculation of the pion TFF

Hoferichter et al. (18)

$$a_\mu^{\pi^0} = 63.0_{-2.1}^{+2.7} \times 10^{-11}$$

- ▶ Padé-Canterbury approximants

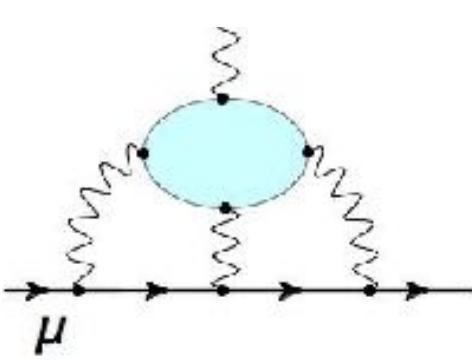
Masjuan & Sanchez-Puertas (17)

$$a_\mu^{\pi^0} = 63.6(2.7) \times 10^{-11}$$

- ▶ Lattice

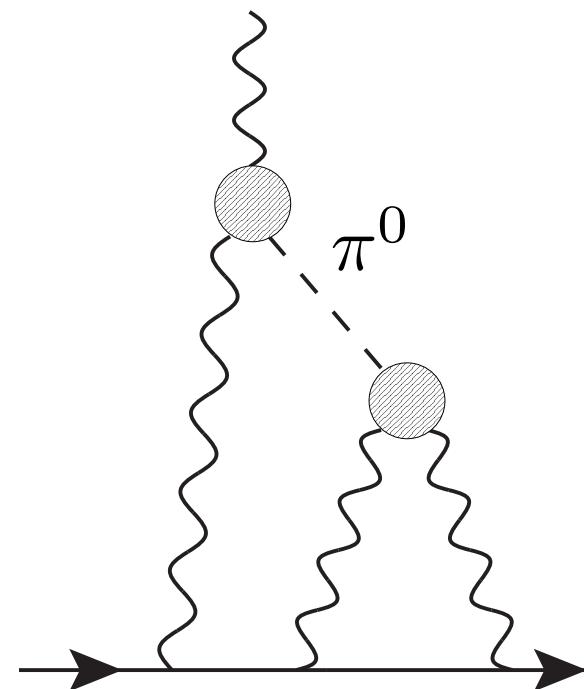
Gérardin, Meyer, Nyffeler (19)

$$a_\mu^{\pi^0} = 62.3(2.3) \times 10^{-11}$$



# HLbL: dispersive

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- Lattice

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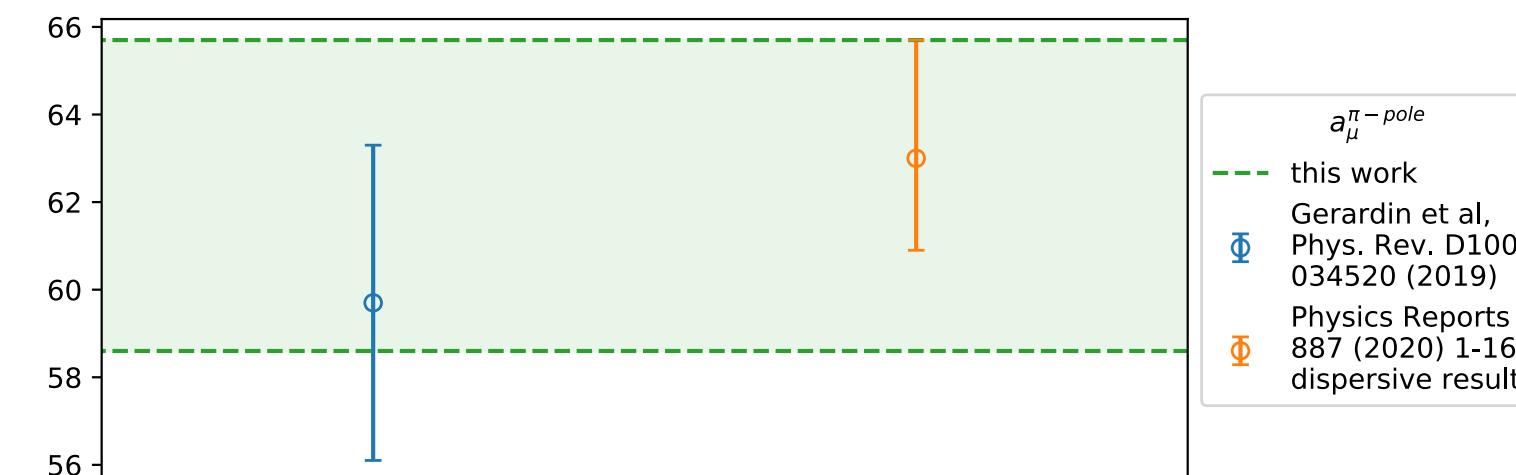
$$a_\mu^{\pi^0} = 62.3(2.3) \times 10^{-11}$$

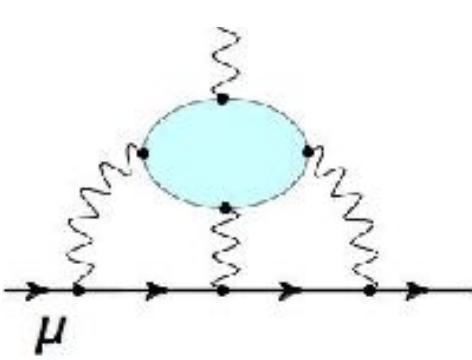
- From the data shown on the previous slide with  $t_{cut} \in [0.88, 2.08]$  fm we get without considering the systematic error a **preliminary**

$a_\mu^{\pi-pole} \in [58.6, 65.7] \times 10^{-11}$  with a 2%-6% statistical error.

⌚ Tuesday, 5:00-8:00 US EDT  
 🎤 Sebastian Burri (ETMc)

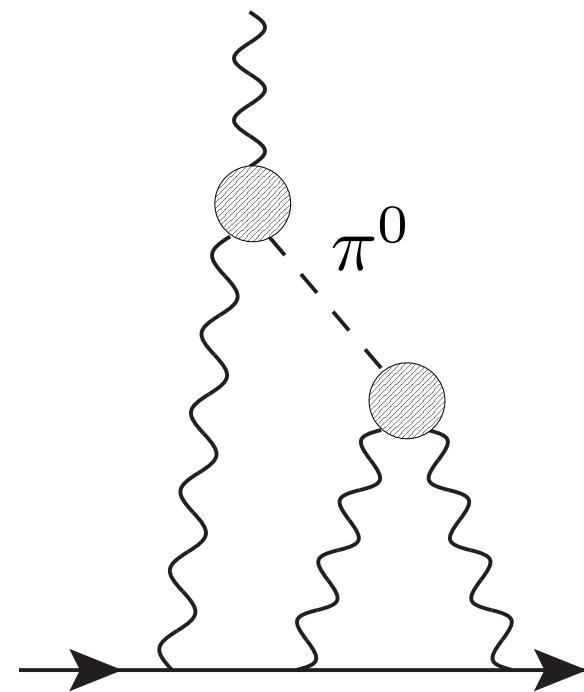
New calculation of pion transition form factor on ETMc twisted-mass ensemble with phys. mass and  $a \sim 0.08$  fm





# HLbL: dispersive

Three independent results for the pion pole contribution:  
 [G. Colangelo @ INT g-2 workshop]



- ▶ Dispersive calculation of the pion TFF

Hoferichter et al. (18)

$$a_\mu^{\pi^0} = 63.0^{+2.7}_{-2.1} \times 10^{-11}$$

- ▶ Padé-Canterbury approximants

Masjuan & Sanchez-Puertas (17)

$$a_\mu^{\pi^0} = 63.6(2.7) \times 10^{-11}$$

- ▶ Lattice

Gérardin, Meyer, Nyffeler (19)

$$a_\mu^{\pi^0} = 62.3(2.3) \times 10^{-11}$$

⌚ Tuesday, 5:00-8:00 US EDT  
 🎤 Willem Verplanke (BMWc)

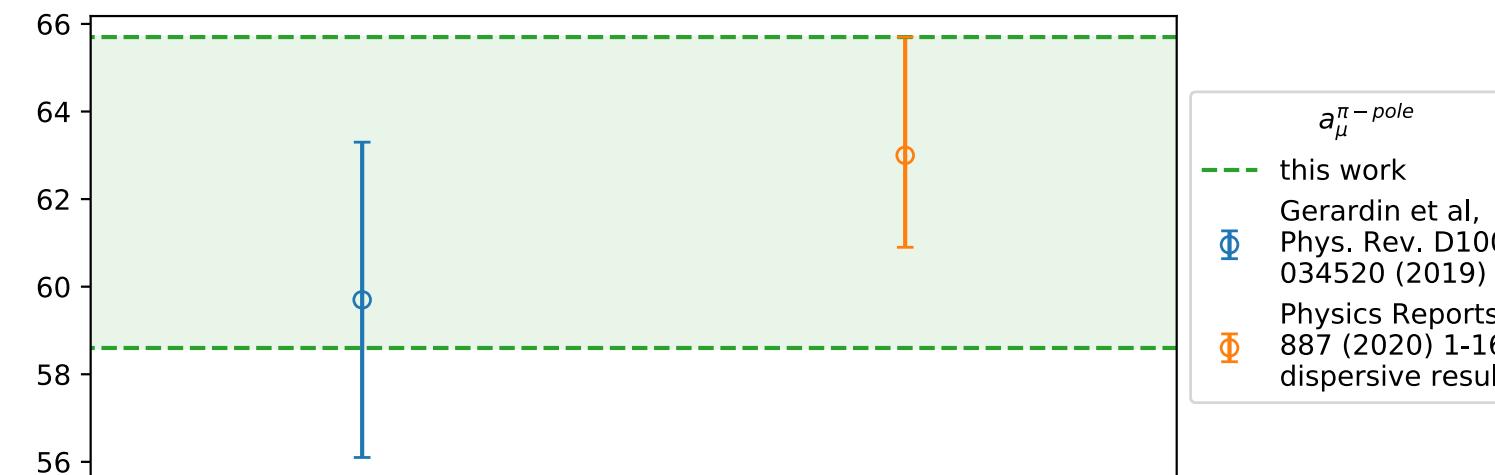
calculation of  $\eta, \eta'$  transition form factors  
 on staggered ensembles (BMWc)

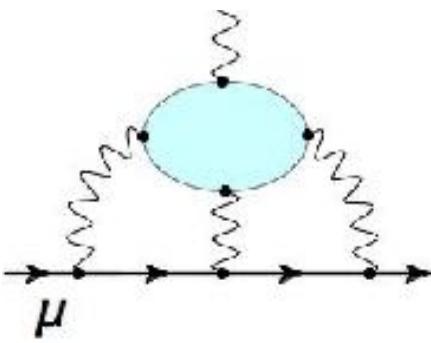
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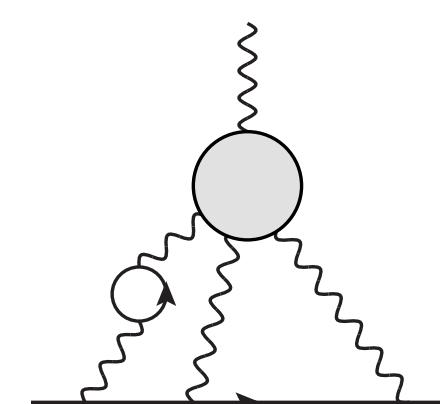




# HLbL: dispersive

Comparison:

Contribution	PdRV(09) [471]	N/JN(09) [472, 573]	J(17) [27]	Our estimate
$\pi^0, \eta, \eta'$ -poles	114(13)	99(16)	95.45(12.40)	93.8(4.0)
$\pi, K$ -loops/boxes	-19(19)	-19(13)	-20(5)	-16.4(2)
$S$ -wave $\pi\pi$ rescattering	-7(7)	-7(2)	-5.98(1.20)	-8(1)
subtotal	88(24)	73(21)	69.5(13.4)	69.4(4.1)
scalars	-	-	-	- 1(3)
tensors	-	-	1.1(1)	
axial vectors	15(10)	22(5)	7.55(2.71)	6(6)
$u, d, s$ -loops / short-distance	-	21(3)	20(4)	15(10)
$c$ -loop	2.3	-	2.3(2)	3(1)
total	105(26)	116(39)	100.4(28.2)	92(19)



NLO HLbL contribution:

$$a_\mu^{\text{HLbL,NLO}} = 2(1) \times 10^{-11}$$

# Outline

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- Introduction
- Theory vs experiment
- Muon g-2 Theory Initiative
- g-2 SM contributions
- Dispersive, data driven methods for
  - Hadronic Vacuum Polarization (HVP)
  - Hadronic Light-by-Light scattering (HLbL)
- Lattice HLbL
- [Lattice HVP](#)
- Summary and Outlook

# g-2 & related talks/posters @ Lattice 2021

## QCD in searches for New Physics

- (light-quark) connected HVP, windows
  - Monday, 13:00-15:00 US EDT
    - Finn Stokes (BMWc) FV effects
    - Kalman Szabo (BMWc) cont. limit
    - Shaun Lahert (Fermilab-HPQCD-MILC)
    - Chris Aubin (Aubin et al)
  - Tuesday, 5:00-8:00 US EDT
    - Hartmut Wittig (Mainz)
    - Christoph Lehner (RBC/UKQCD)
    - Davide Giusti (ETMc)

- IB corrections, disc. HVP
  - Tuesday, 5:00-8:00 US EDT
    - Andreas Risch (Mainz)
    - Letizia Parato (BMWc)
  - Poster, Wednesday, 8:00-9:00 US EDT
    - C. McNeile (Fermilab-HPQCD-MILC)

- $\Delta\alpha$  and  $\Delta \sin^2 \theta_W$ 
  - Tuesday, 5:00-8:00 US EDT
    - Teseo San Jose (Mainz)
    - Kohtaro Miura (Mainz)

- HLbL contributions, PS transition form factors
  - Tuesday, 5:00-8:00 US EDT
    - Willem Verplanke (BMWc)
    - Sebastian Burri (ETMc)
    - En-Hung Chao (Mainz) complete HLbL

- Scale Setting
  - Monday, 13:00-15:00 US EDT
    - Lukas Varnhorst (BMWc)
  - Thursday 5:00-8:00 US EDT Friday, 5:00–8:00 US EDT
    - Alexander Segner (Mainz) Ben Strassberger (Mainz)

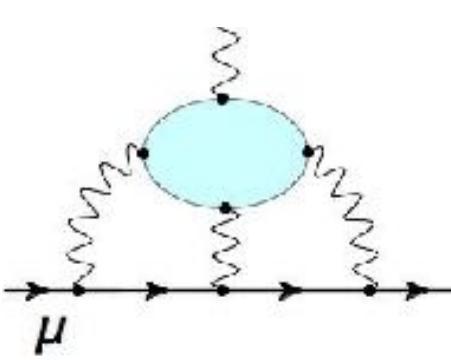
- Cut-off effects
  - Tuesday, 5:00-8:00 US EDT
    - Tim Harris (NEPhEU QCD) thermal observables
  - Friday, 5:00-8:00 US EDT
    - Nicolai Husung (DESY) log corrections

- Variance reduction
  - Tuesday, 5:00-8:00 US EDT
    - Leonardo Giusti (Milan) Multi-level integration
  - Tuesday, 13:00-15:00 US EDT
    - Tej Kanwar (MIT) contour deformation

Hadron  
spect.

SM params

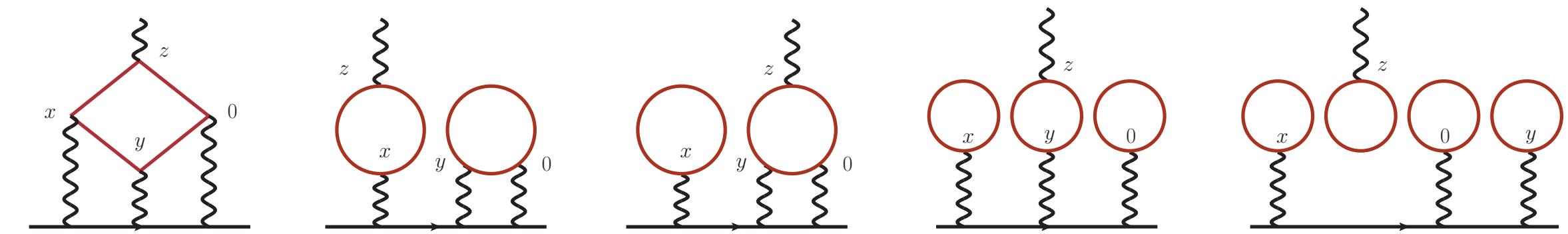
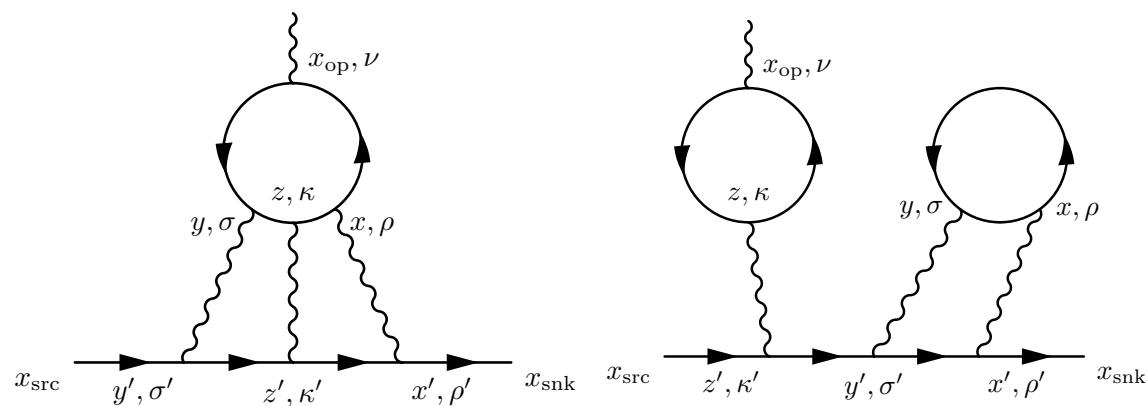
Algorithms



# HLbL: lattice

Hadronic light-by-light: Target:  $\lesssim 10\%$  total error

Two independent and complete direct lattice calculations of  $a_\mu^{\text{HLbL}}$



## ◆ RBC/UKQCD

[T. Blum et al, arXiv:1610.04603, 2016 PRL; arXiv:1911.08123, 2020 PRL]

## ◆ QCD + QED<sub>L</sub> (finite volume)

⇒  $1/L^2$  FV effects

stochastic evaluation of position space sums

Feynman gauge photon propagators

DWF ensembles at/near phys mass,

$a \approx 0.08 - 0.2$  fm,  $L \sim 4.5 - 9.3$  fm

## ◆ Mainz group

[E. Chao et al, arXiv:2104.02632]

## ◆ QCD + QED (infinite volume & continuum)

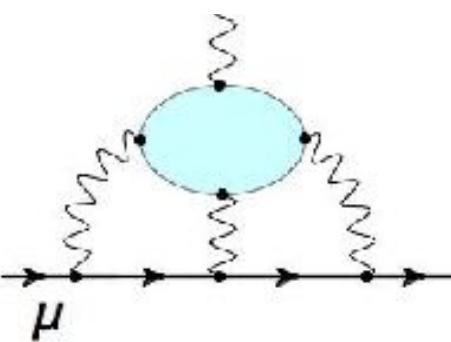
⇒  $e^{-m_\pi L}$  FV effects

semi-analytic QED kernel function

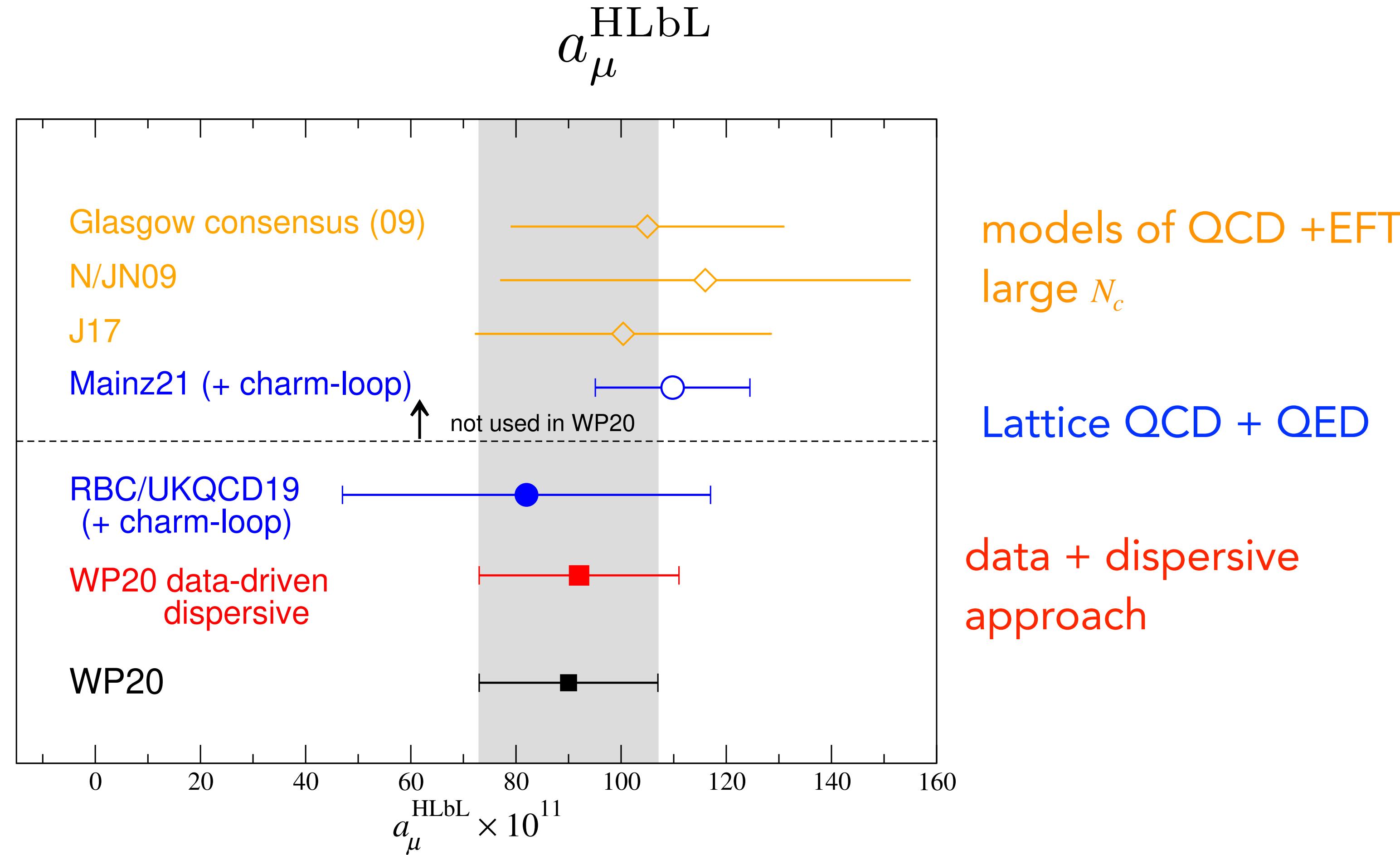
CLS (2+1 Wilson-clover) ensembles

$m_\pi \sim 200 - 430$  MeV,  $a \approx 0.05 - 0.1$  fm,  $m_\pi L > 4$

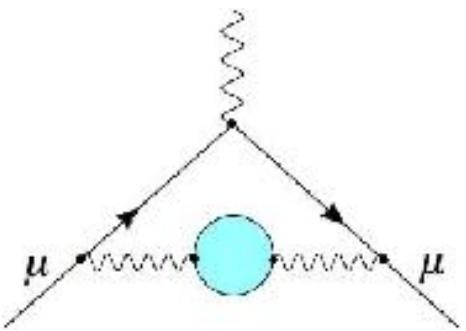
## ◆ Cross checks between RBC/UKQCD & Mainz approaches in White Paper at unphysical pion mass



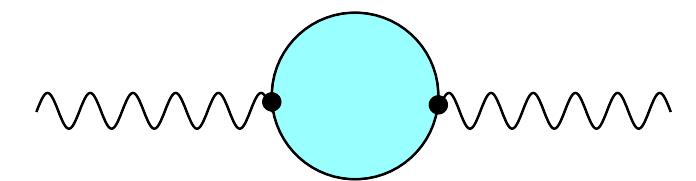
# HLbL: Comparison



Now well-determined in two independent approaches, systematically improvable



# Lattice HVP: Introduction



$$\hat{\Pi}(q^2)$$

Leading order HVP correction:

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int dq^2 \omega(q^2) \hat{\Pi}(q^2)$$

[B. Lautrup, A. Peterman, E. de Rafael, Phys. Rep. 1972;  
E. de Rafael, Phys. Let. B 1994; T. Blum, PRL 2002]

- Calculate  $a_\mu^{\text{HVP,LO}}$  in Lattice QCD

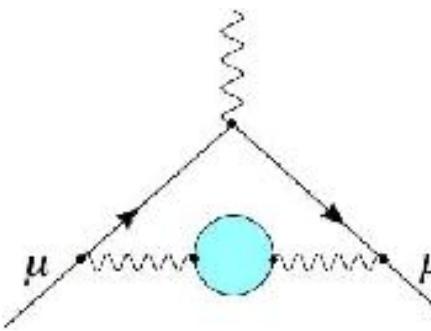
Compute correlation function:  $C(t) = \frac{1}{3} \sum_{i,x} \langle j_i(x, t) j_i(0, 0) \rangle$

and  $\hat{\Pi}(Q^2) = 4\pi^2 \int_0^\infty dt C(t) \left[ t^2 - \frac{4}{Q^2} \sin^2 \left( \frac{Qt}{2} \right) \right]$

[D. Bernecker and H. Meyer, arXiv:1107.4388,  
EPJA 2011]

Obtain  $a_\mu^{\text{HVP,LO}}$  from an integral over Euclidean time:

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \tilde{w}(t) C(t)$$

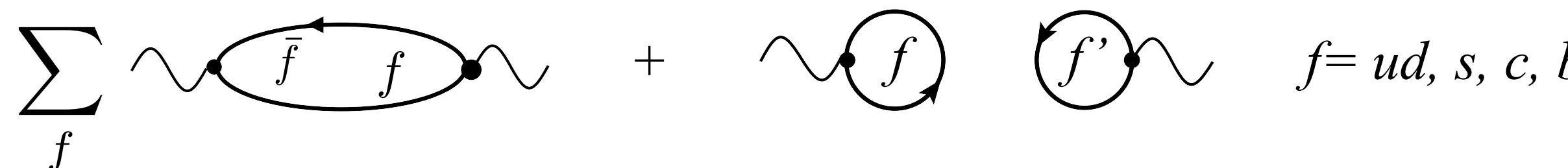


# Lattice HVP: Introduction

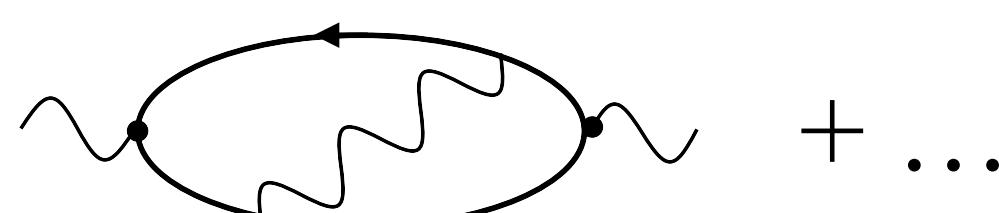
Calculate  $a_\mu^{\text{HVP}}$  in Lattice QCD:

$$a_\mu^{\text{HVP,LO}} = \sum_f a_{\mu,f}^{\text{HVP,LO}} + a_{\mu,\text{disc}}^{\text{HVP,LO}}$$

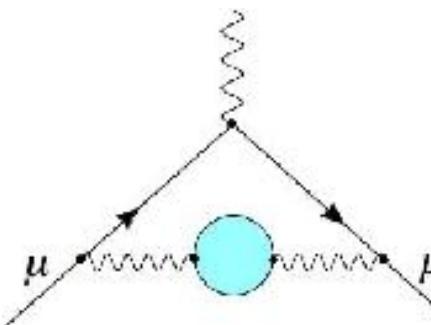
- Separate into connected for each quark flavor + disconnected contributions  
(gluon and sea-quark background not shown in diagrams)
- Note: almost always  $m_u = m_d$



- need to add QED and strong isospin breaking ( $\sim m_u - m_d$ ) corrections:

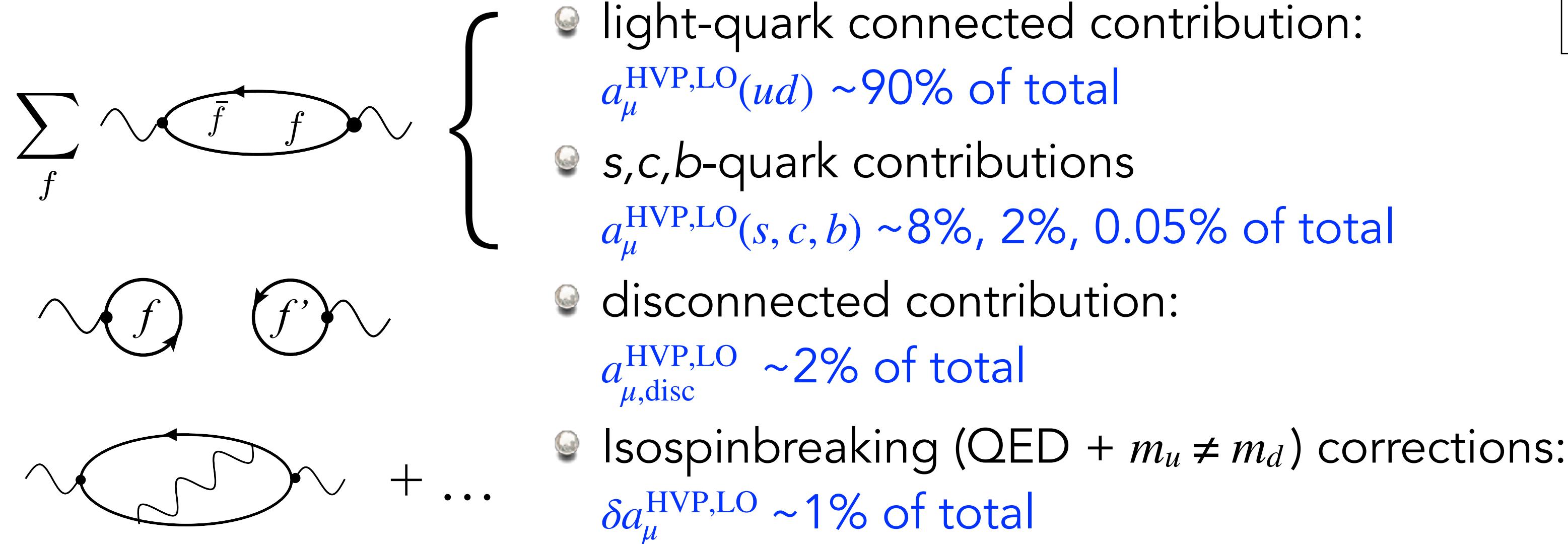


- either perturbatively on isospin symmetric QCD background
- or by using QCD + QED ensembles with  $m_u \neq m_d$



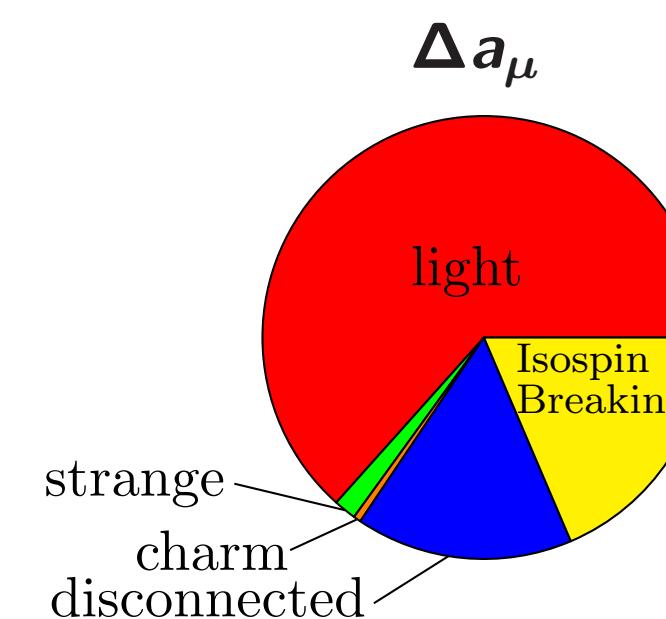
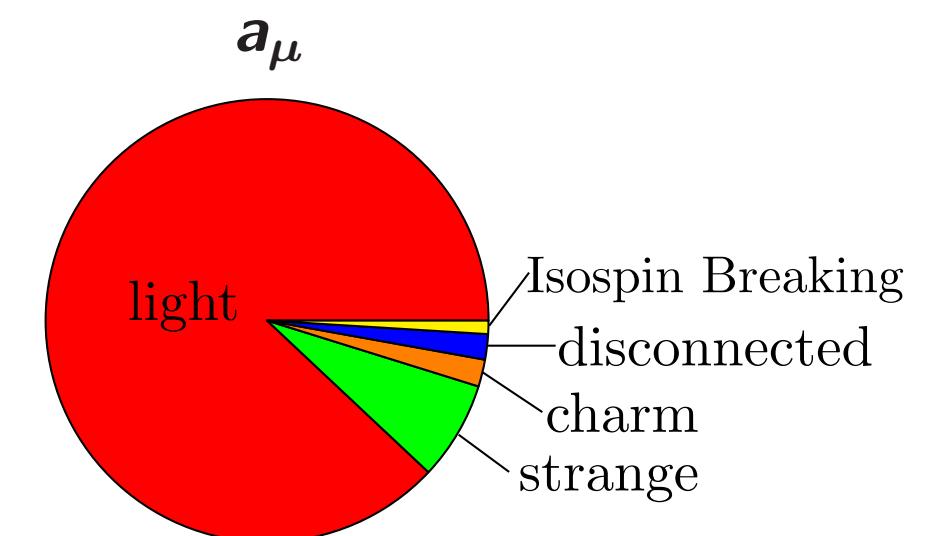
# Lattice HVP: Introduction

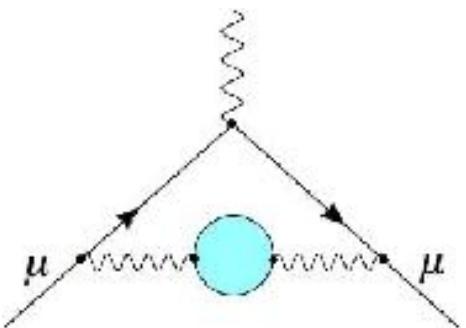
- Target: ~ 0.2% total error



$$a_\mu^{\text{HVP,LO}} = a_\mu^{\text{HVP,LO}}(ud) + a_\mu^{\text{HVP,LO}}(s) + a_\mu^{\text{HVP,LO}}(c) + a_{\mu, \text{disc}}^{\text{HVP,LO}} + \delta a_\mu^{\text{HVP,LO}}$$

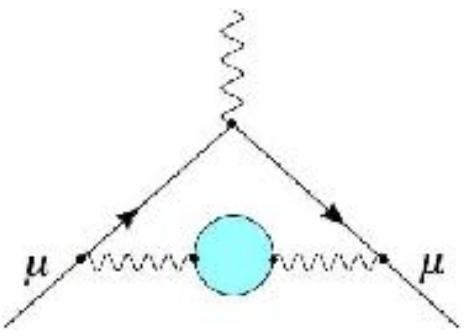
V. GÜLPER @ Lattice HVP workshop



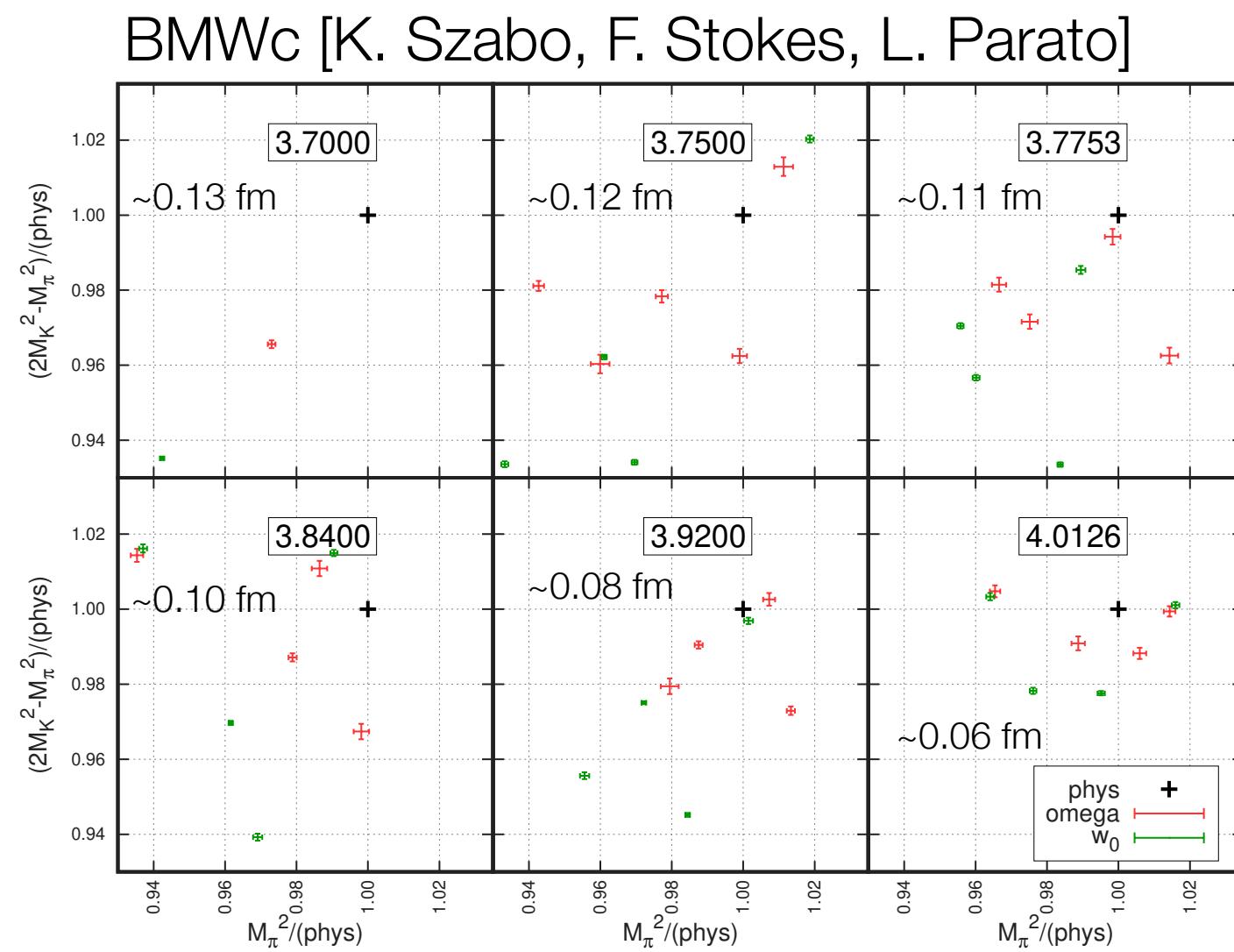


# Lattice HVP: Introduction

- ➊ Target: ~ 0.2% total error
- ➋ Challenges:
  - ✓ needs ensembles with (light sea) quark masses at their physical values
  - ✓ finite volume corrections
    - growth of statistical errors at long-distances
    - Continuum extrapolation
    - scale setting
    - disconnected contribution
    - QED and strong isospin breaking corrections ( $m_u \neq m_d$ )
- ➌ Focus on windows in Euclidean times [T. Blum et al, arXiv:1801.07224, 2018 PRL]
  - disentangle systematics/statistics from long distance/FV and discretization effects
    - ➡ valuable cross checks
  - intermediate window easy to compute & compare with disperse methods



# Lattice HVP: Ensemble parameters

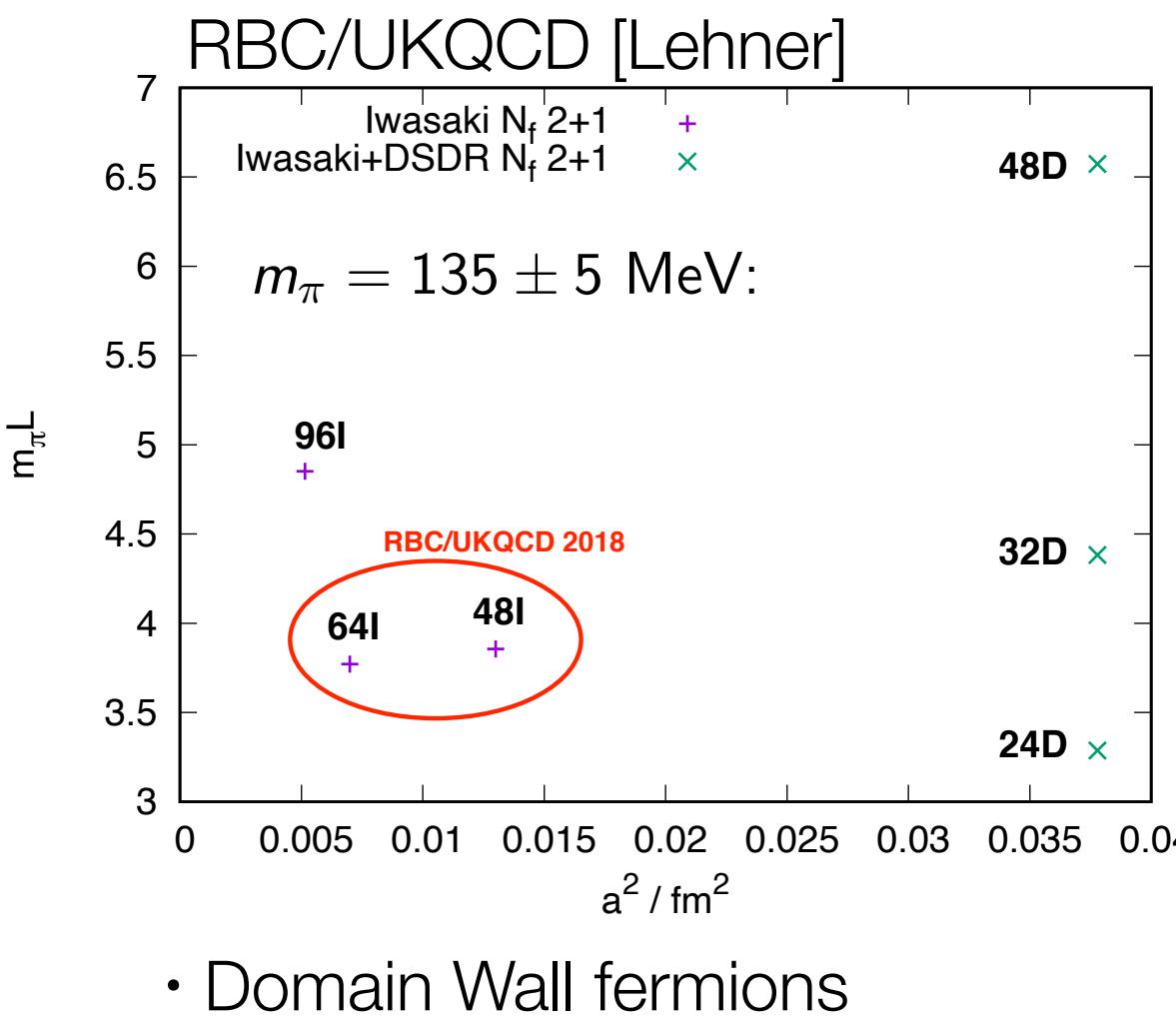


- Stout-smeared staggered fermions 2+1+1
- $L \sim 6 - 11 \text{ fm}$

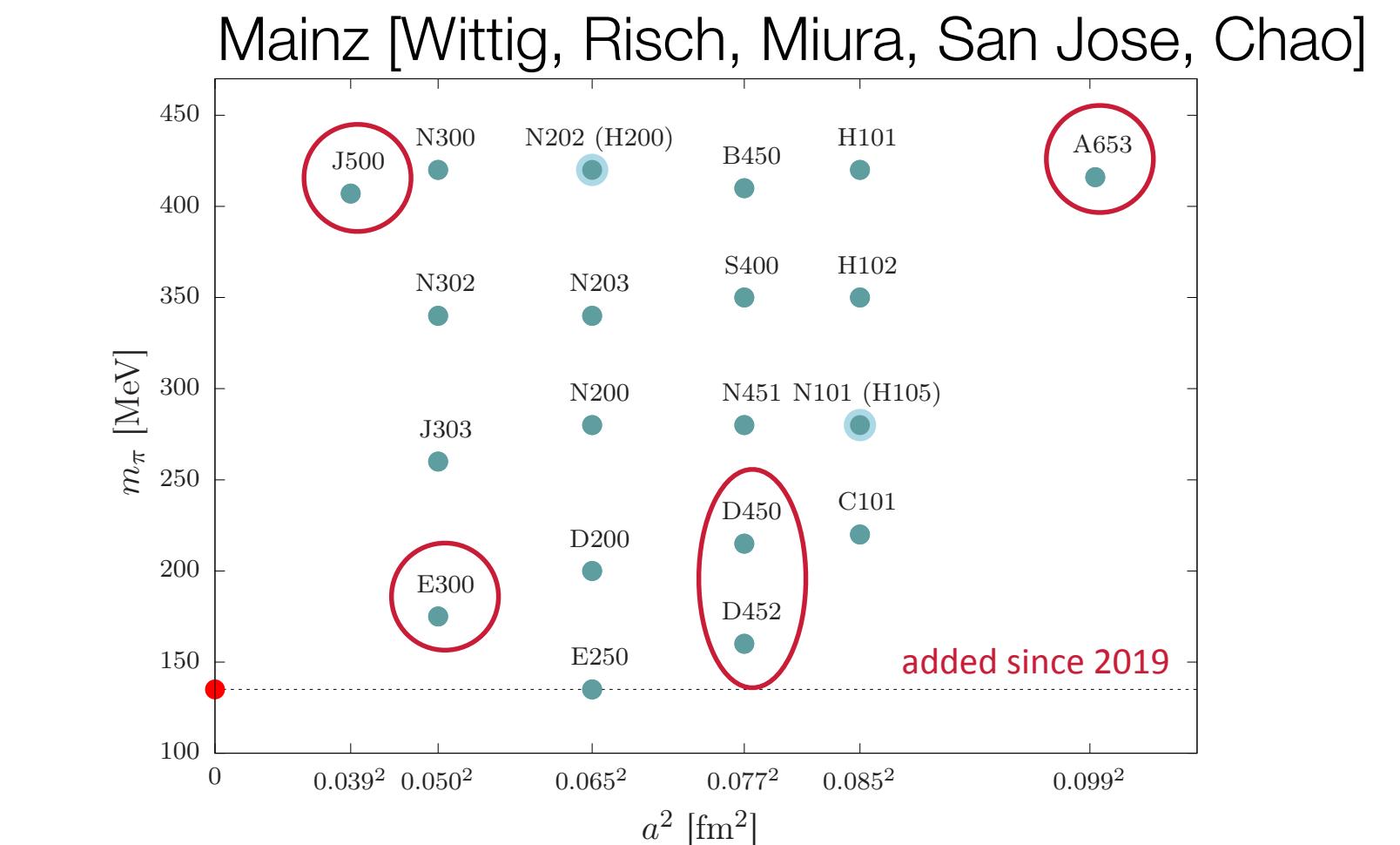
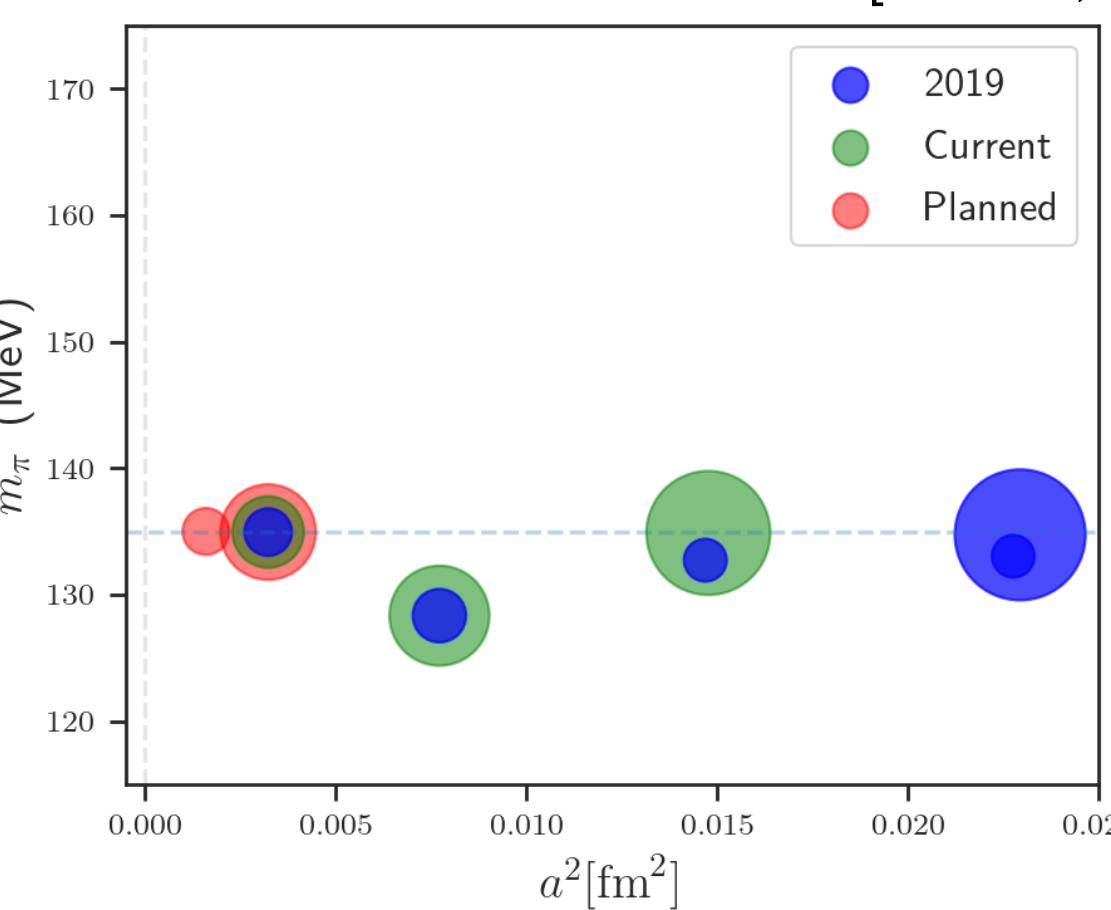
**ETMc [D. Giusti]**

Pion masses in the range 220 - 490 MeV  
4 volumes @  $M_\pi \approx 320 \text{ MeV}$  and  $a \approx 0.09 \text{ fm}$   
 $M_\pi L \approx 3.0 \div 5.8$

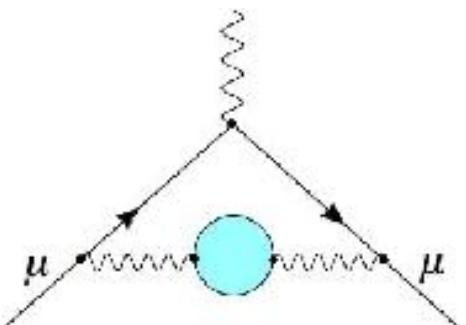
- Twisted-mass Wilson fermions, 2+1+1
- Plan to include phys. mass ensemble in future



Fermilab-HPQCD-MILC [Lahert, McNeile]



- Highly Improved Staggered Quarks (HISQ) 2+1+1+1
- $L \sim 5 - 6 \text{ fm}$
- Subset also used by Aubin et al [Aubin]



# Lattice HVP: Finite Volume

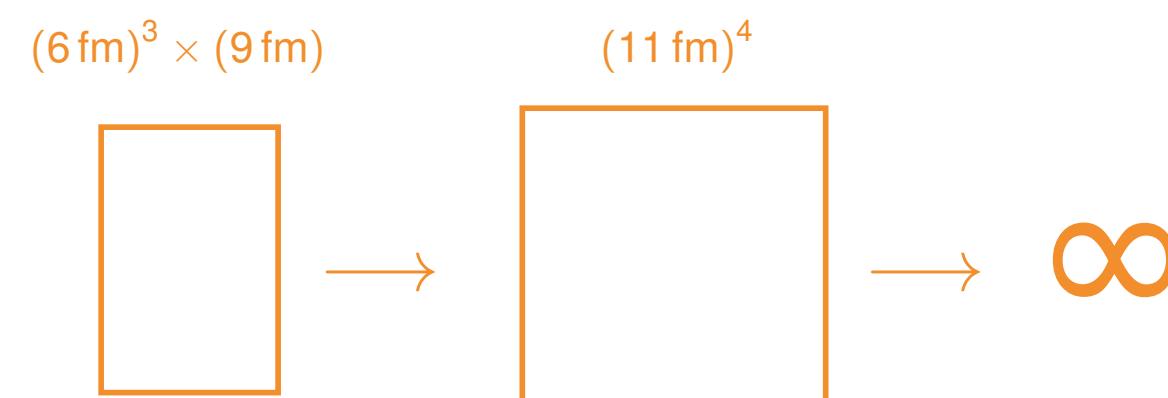
Monday, 13:00-15:00 US EDT  
Finn Stokes (BMWc)

Tuesday, 5:00-8:00 US EDT  
Hartmut Wittig (Mainz)

## Model comparison

- Two more models for finite L (but not T)
    - Generic field-theory approach [Hansen & Patella '19, '20] (HP) relates the finite-size effect to  $F_\pi(k)$
    - Rho-pion-gamma model [Chakraborty et al '17] (RHO) incorporates the  $\rho(770)$  resonance directly into a  $\chi$ PT-like framework
  - Compare finite L corrections for reference volume in infinite-T limit
  - All four models agree within  $\sim 2.5 \times 10^{-10}$
- |                |      |                   |
|----------------|------|-------------------|
| NNLO $\chi$ PT | 16.7 | $\times 10^{-10}$ |
| MLLGS          | 18.8 | $\times 10^{-10}$ |
| HP             | 17.7 | $\times 10^{-10}$ |
| RHO            | 16.2 | $\times 10^{-10}$ |

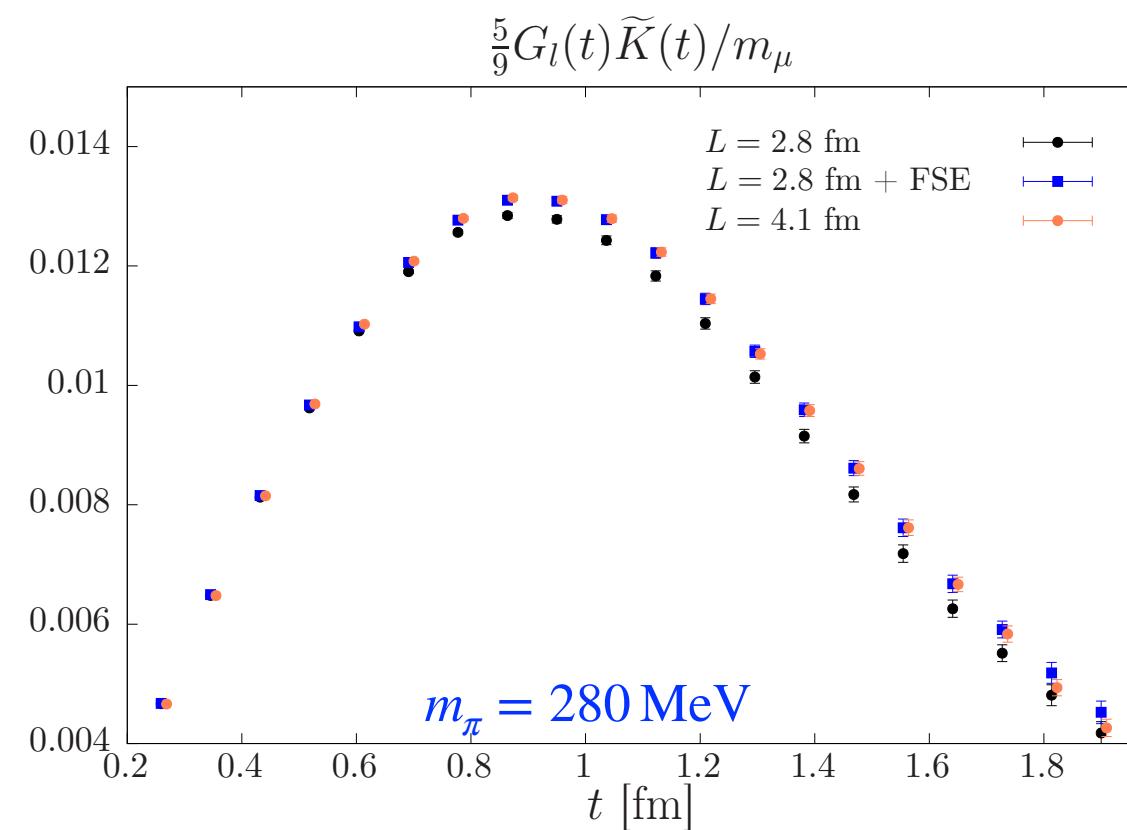
## Residual correction



4HEX	$18.1 \pm 2.4$	$\times 10^{-10}$	
NNLO $\chi$ PT	$15.7$	$0.6 \pm 0.3$	$\times 10^{-10}$
MLLGS	$17.8$		$\times 10^{-10}$

Mainz method (aka MLL):  $G(t, L) \xrightarrow{t \rightarrow \infty} \sum_n |A_n|^2 e^{-\omega_n t}$        $G(t, \infty) = \int_0^\infty d\omega \omega^2 \rho(\omega^2) e^{-\omega|t|}$

Both  $|A_n|$  and  $\rho(\omega^2)$  can be related to the pion form factor  $F_\pi(\omega) \Rightarrow G(t, \infty) - G(t, L)$   
[Meyer 2011, Francis et al. 2013, Della Morte et al. 2017; Lellouch & Lüscher 2001]



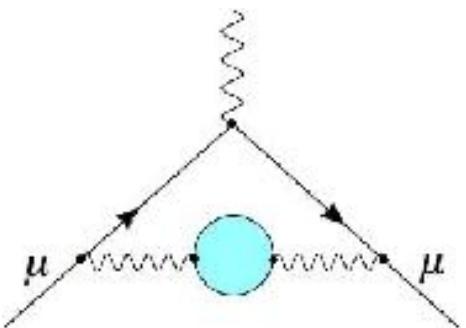
## Full HVP contribution:

Explicit verification at  $m_\pi = 280\text{ MeV}$ :  
→ method works remarkably well  
 $m_\pi L = 4$  ( $L = 6.2\text{ fm}$ )  
 $\Rightarrow \Delta_{\text{FV}} a_\mu^{\text{hyp}} = 22.6 \cdot 10^{-10}$  (3%)

In agreement with direct lattice calculations:  
PACS, BMWc

[Shintani & Kuramashi 2019, BMWc 2020]

- Other direct calculations by RBC/UKQCD [Lehner @ 2019 INT workshop, Shintani & Kuramashi, PRD 2019] are consistent (but with larger errors).

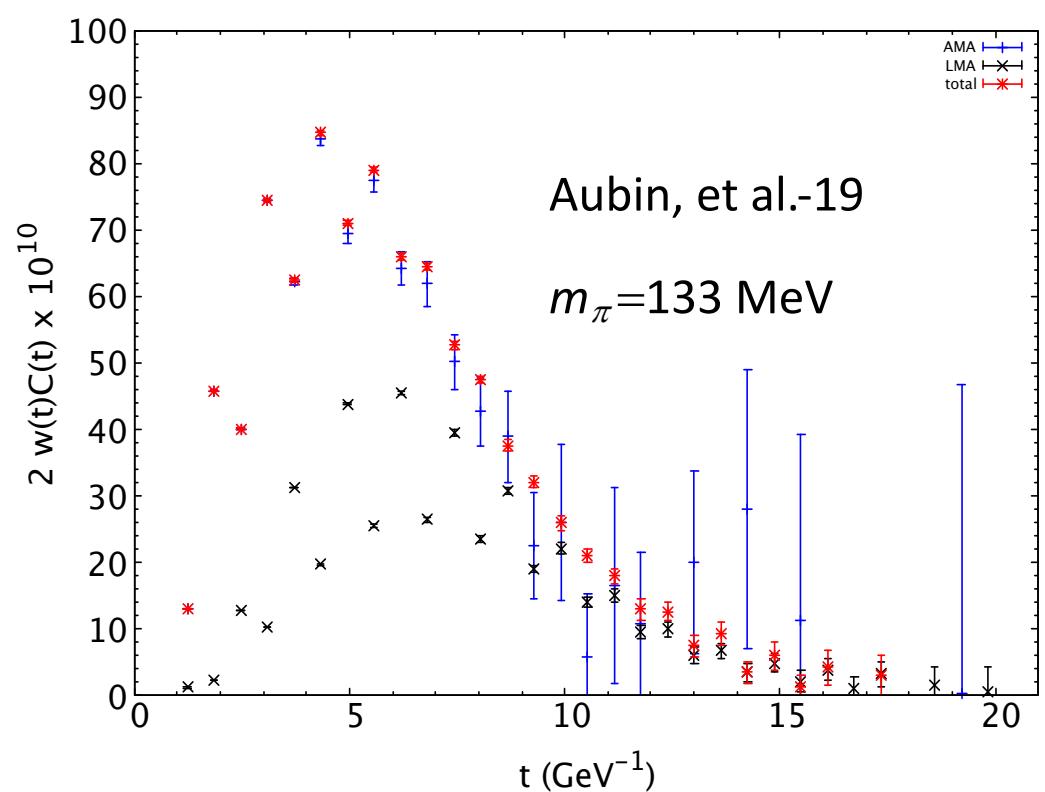


# Lattice HVP: long-distance tail

$$G(t) = \frac{1}{3} \sum_{i,x} \langle j_i(x,t) j_i(0,0) \rangle$$

- Use noise reduction methods (AMA, LMA,...):

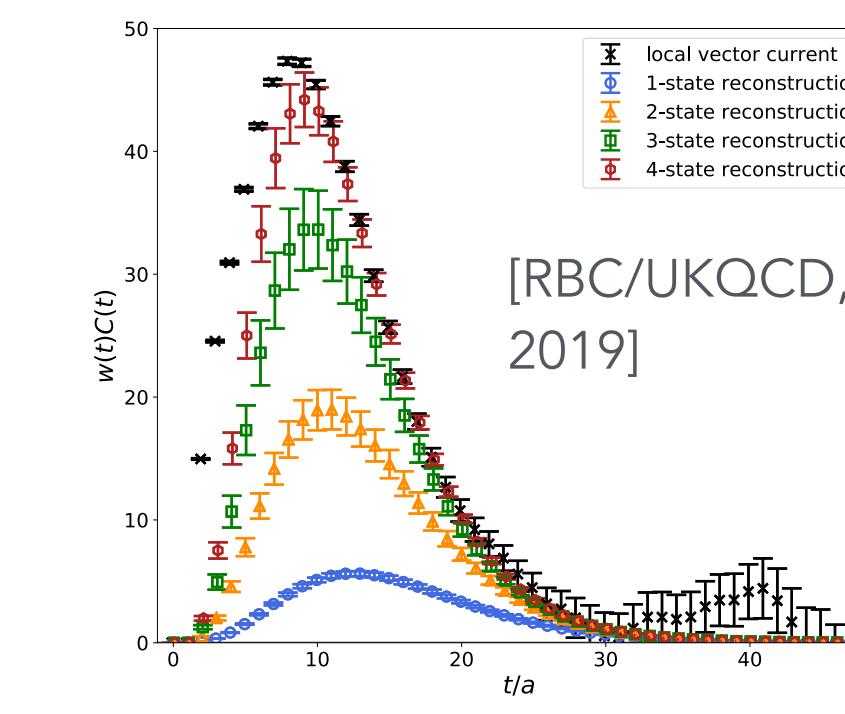
Aubin et al, RBC/UKQCD, BMWc, Mainz, ...

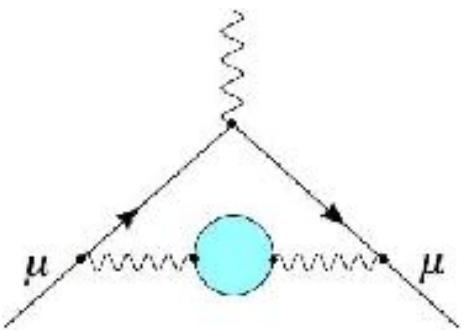


- Spectral reconstruction (RBC/UKQCD, Mainz):

- ◆ obtain low-lying finite-volume spectrum ( $E_n, A_n$ ) in dedicated study using additional operators that couple to two-pion states
- ◆ use to reconstruct  $G(t > t_c)$
- ◆ can be used to improve bounding method:

$$G(t) \rightarrow G(t) - \sum_{n=0}^N A_n^2 e^{-E_n t}$$



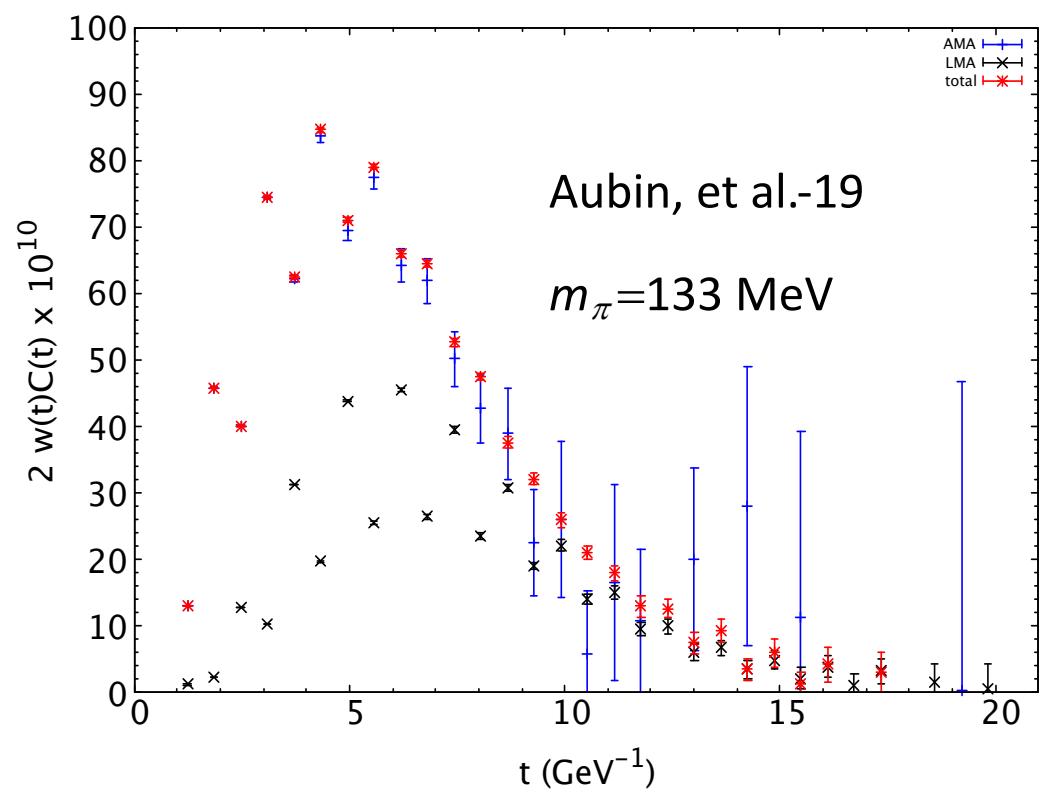


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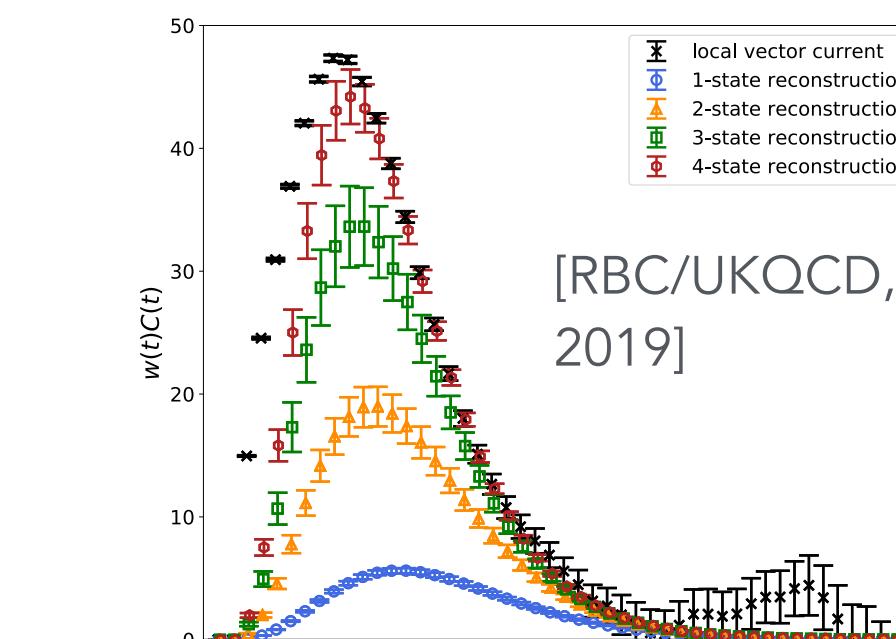
Aubin et al, RBC/UKQCD, BMWc, Mainz, ...



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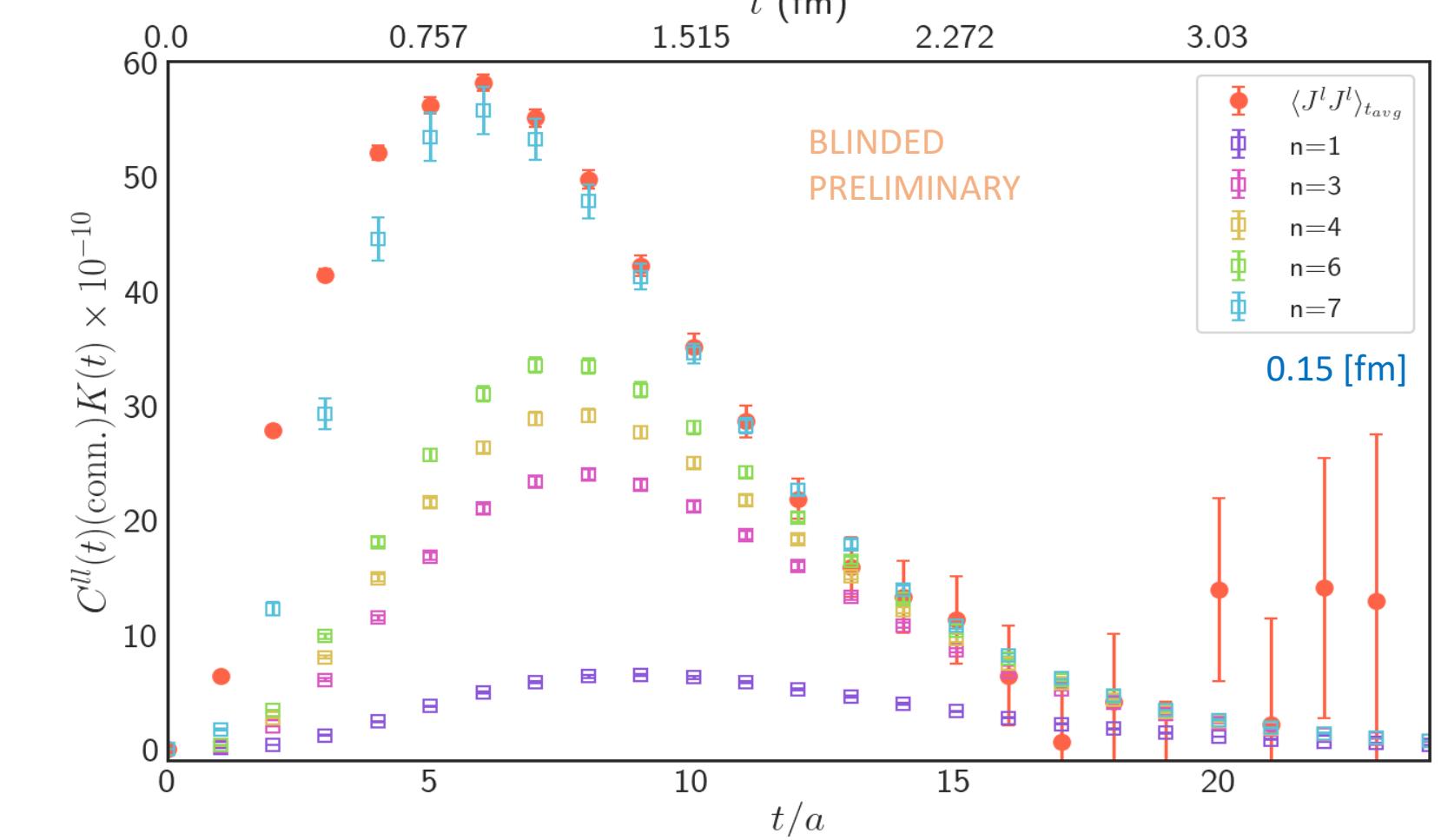
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- ◆ use to reconstruct  $G(t > t_c)$
- ◆ can be used to improve bounding method:

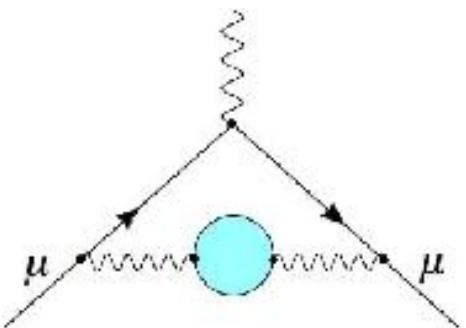
$$G(t) \rightarrow G(t) - \sum_{n=0}^N A_n^2 e^{-E_n t}$$



⌚ Monday, 13:00-15:00 US EDT  
🎤 Shaun Lahert  
 (Fermilab-HPQCD-MILC)

- First calculation with staggered multi-pion operators



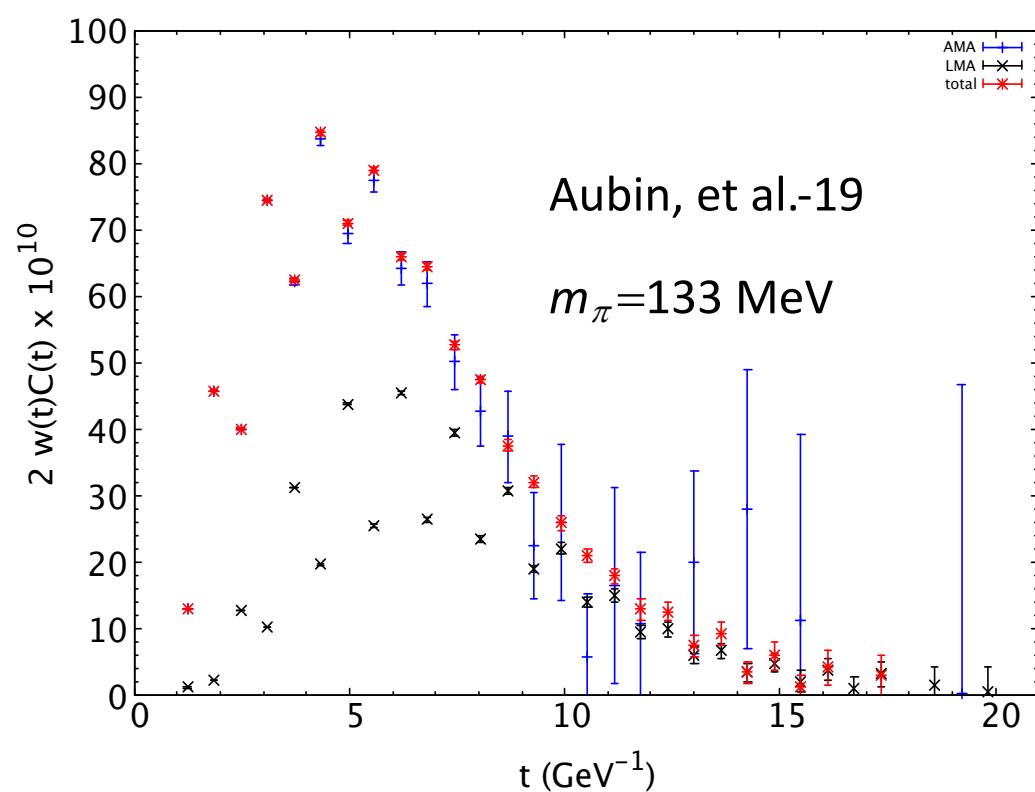


# Lattice HVP: long-distance tail

$$G(t) = \frac{1}{3} \sum_{i,x} \langle j_i(x,t) j_i(0,0) \rangle$$

- Use noise reduction methods (AMA, LMA,...):

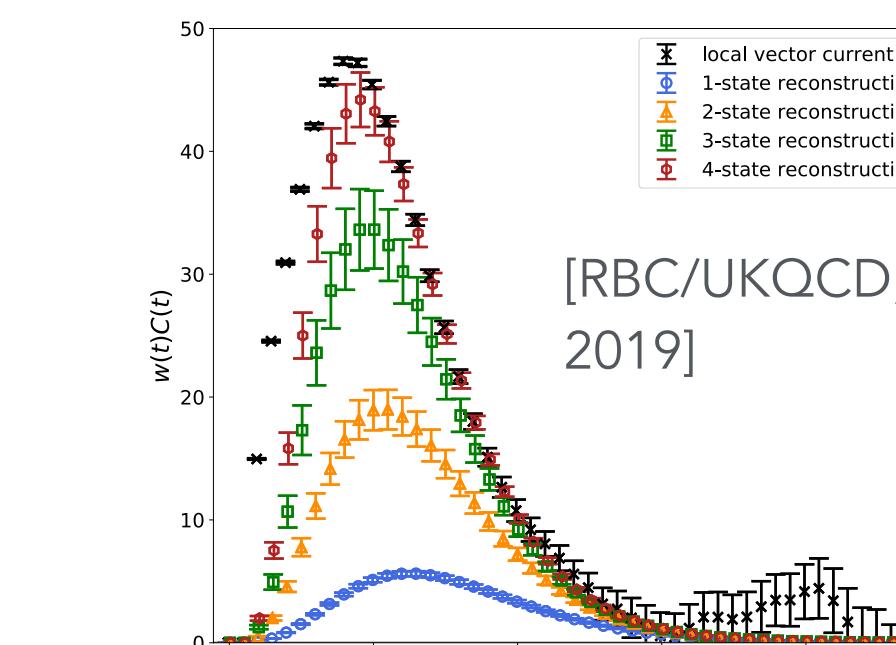
Aubin et al, RBC/UKQCD, BMWc, Mainz, ...



- Spectral reconstruction (RBC/UKQCD, Mainz):

- ◆ obtain low-lying finite-volume spectrum ( $E_n, A_n$ ) in dedicated study using additional operators that couple to two-pion states
- ◆ use to reconstruct  $G(t > t_c)$
- ◆ can be used to improve bounding method:

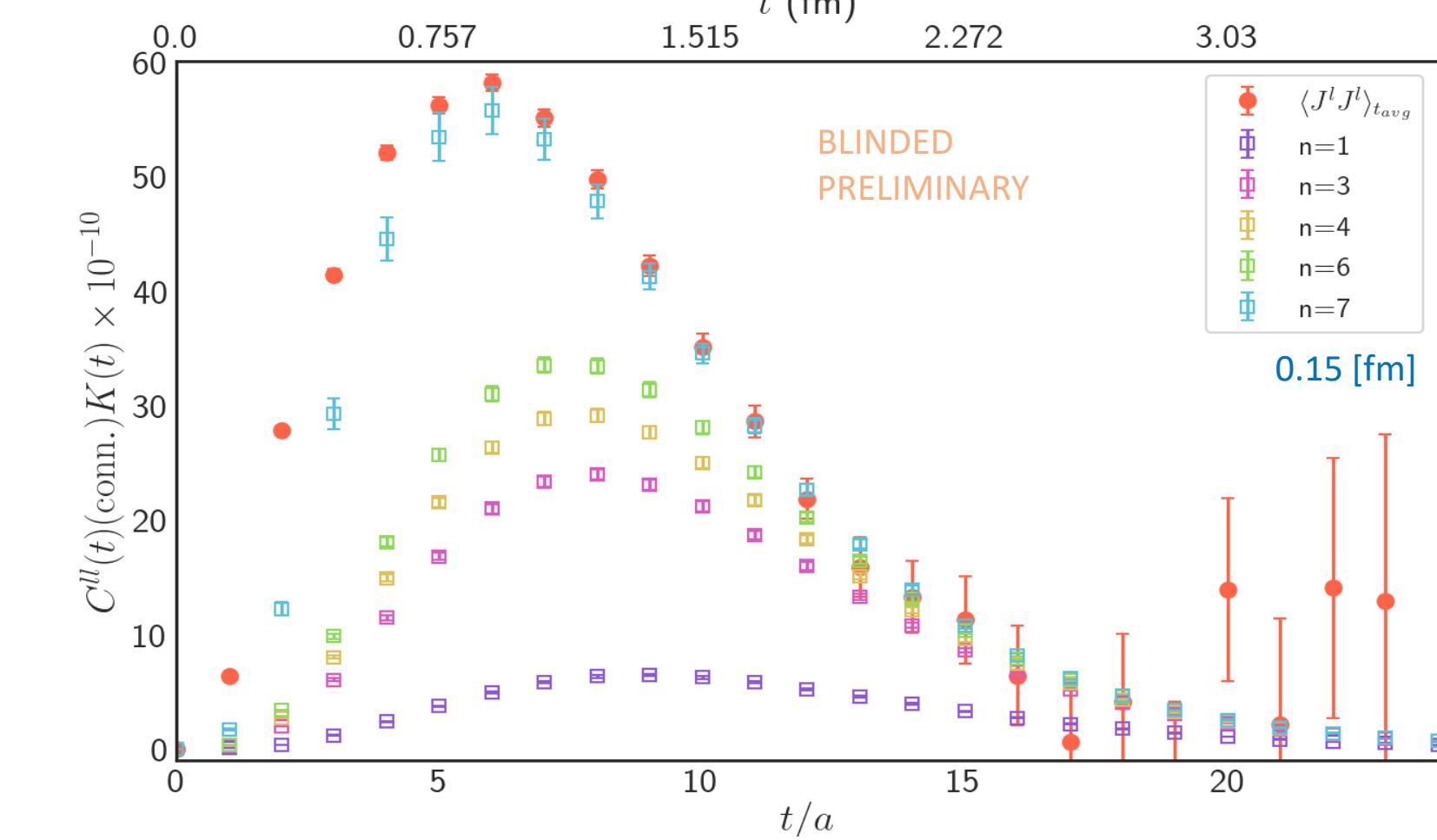
$$G(t) \rightarrow G(t) - \sum_{n=0}^N A_n^2 e^{-E_n t}$$

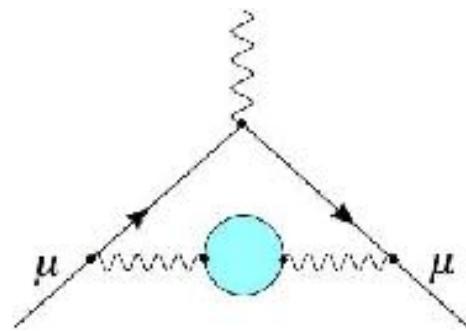


⌚ Tuesday, 5:00-8:00 US EDT	⌚ Leonardo Giusti (Milan) Multi-level integration
⌚ Tuesday, 13:00-15:00 US EDT	⌚ Tej Kanwar (MIT) contour deformation

⌚ Monday, 13:00-15:00 US EDT  
\_MICROPHONE Shaun Lahert  
 (Fermilab-HPQCD-MILC)

- First calculation with staggered multi-pion operators



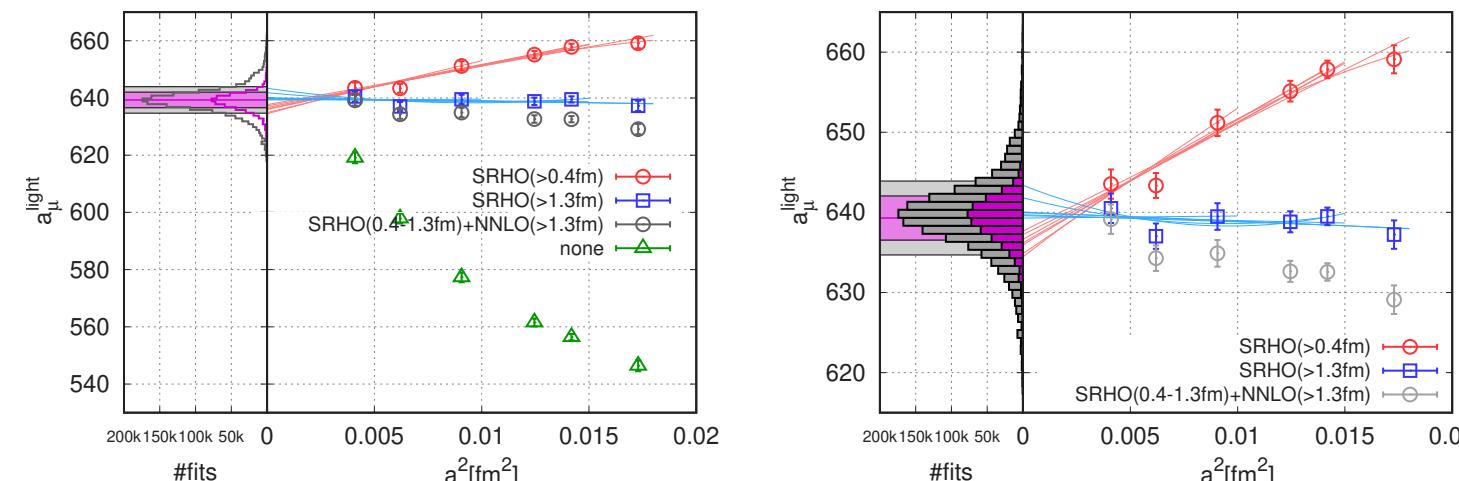


# Lattice HVP: continuum extrapolation

Monday, 13:00-15:00 US EDT  
Kalman Szabo (BMWc)

## Taste improvement II

- $a_\mu(a) \rightarrow a_\mu(a) - a_\mu^{\text{SRHO}}(a) + a_\mu^{\text{RHO}}$
- reduces lattice artefact, also makes  $a^2$  dependence linear



SRHO improvement gives central value. Systematic errors by:

- change starting point of improvement  $t = 0.4 \rightarrow 1.3$  fm
- skip coarse lattices
- change  $\Gamma = 0$  and  $\Gamma = 3$
- replace SRHO by NNLO SXPT above 1.3 fm

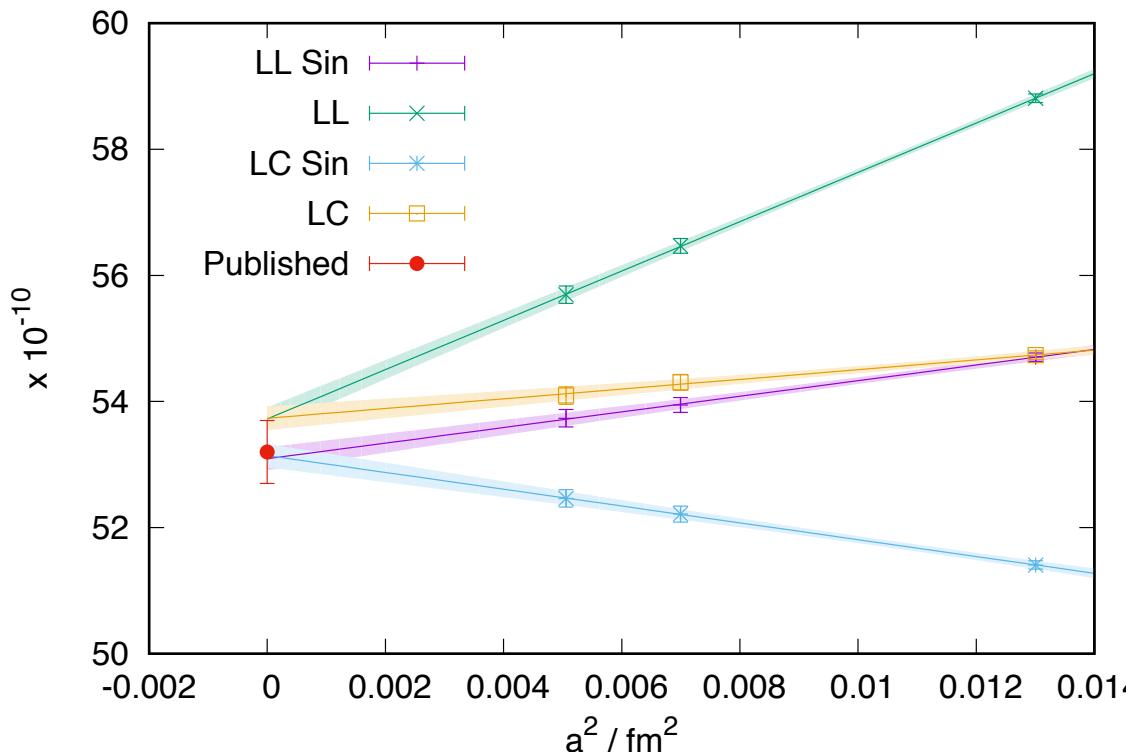
- Large taste-breaking effects with BMW set-up
  - uncorrected data not easily fit to power series, i.e.

$$1 \quad A_0 + A_1 [a^2] + A_2 [a^2]^2$$

$$2 \quad A_0 + A_1 [a^2 \alpha_s^3(\frac{1}{a})] + A_2 [a^2 \alpha_s^3(\frac{1}{a})]^2$$

Tuesday, 5:00-8:00 US EDT  
Christoph Lehner (RBC/UKQCD)

- Third lattice spacing for strange data ( $a^{-1} = 2.77$  GeV with  $m_\pi = 234$  MeV with sea light-quark mass corrected from global fit):

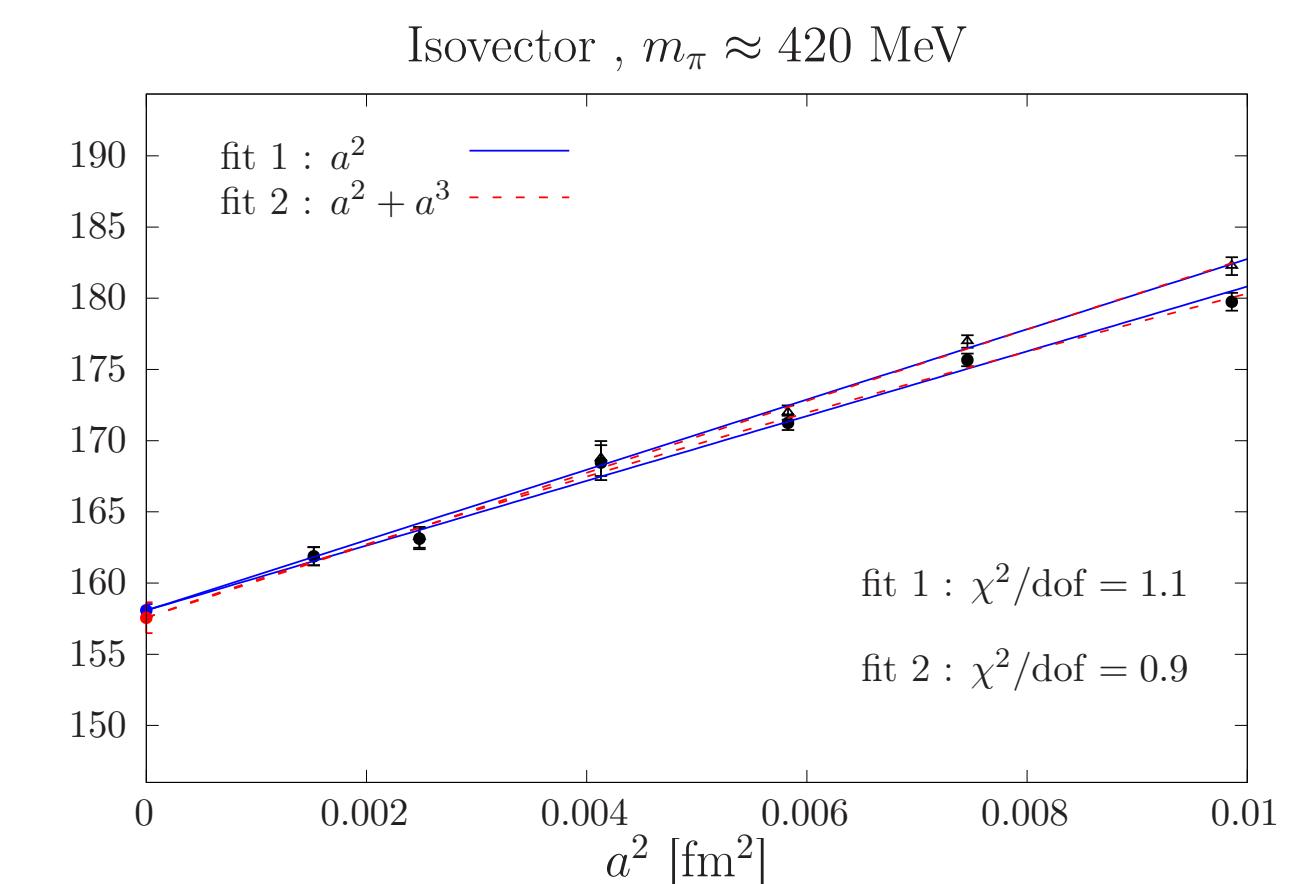


- For light quark use new 96I ensemble at physical pion mass. Data still being generated on Summit in USA and Booster in Germany ( $a^{-1} = 2.77$  GeV with  $m_\pi = 139$  MeV)

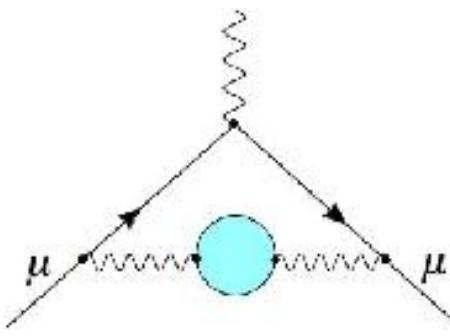
- RBC:  
Currently adding add a third lattice spacing

- Fermilab-HPQCD-MILC:  
planning to add a 5th lattice spacing (0.042 fm).

Tuesday, 5:00-8:00 US EDT  
Hartmut Wittig (Mainz)



- Mainz and ETMc perform combined chair and continuum extrapolation



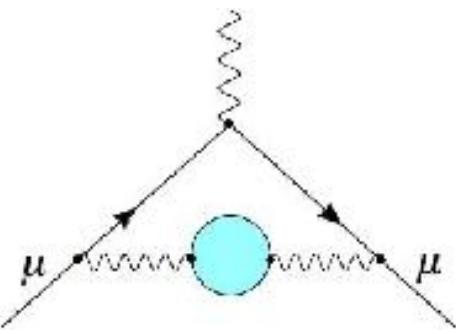
# Lattice HVP: Scale Setting

- $a_\mu$  is dimensionless, but depends on the lattice indirectly, through masses in lattice units in the Kernel. In particular,  $am_\mu$ :

$$\frac{\delta a_\mu^{\text{hvp}}}{a_\mu^{\text{hvp}}} = \underbrace{\frac{1}{a_\mu^{\text{hvp}}} \left| a \frac{da_\mu^{\text{hvp}}}{da} \right|}_{\approx 1.8} \frac{\delta a}{a}$$

[H. Wittig @ 1st Muon g-2 Theory Initiative workshop;  
Della Morte et al, Lattice 2017]

- need a good physical quantity to determine lattice spacing to high precision (< 0.2%).  
Currently in use:
  - $f_\pi$  — depends on  $V_{ud}$  and requires radiative QED corrections
  - $\Omega$  baryon mass (RBC/UKQCD, BMW)  
also being adopted by Mainz, Fermilab-HPQCD-MILC



# Lattice HVP: Scale Setting

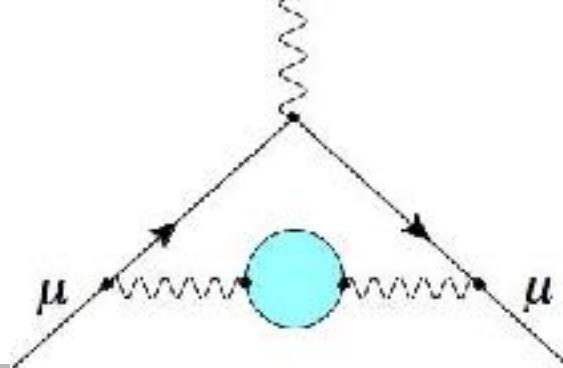
- $a_\mu$  is dimensionless, but depends on the lattice indirectly, through masses in lattice units in the Kernel. In particular,  $am_\mu$ :

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also being adopted by Mainz, Fermilab-HPQCD-MILC

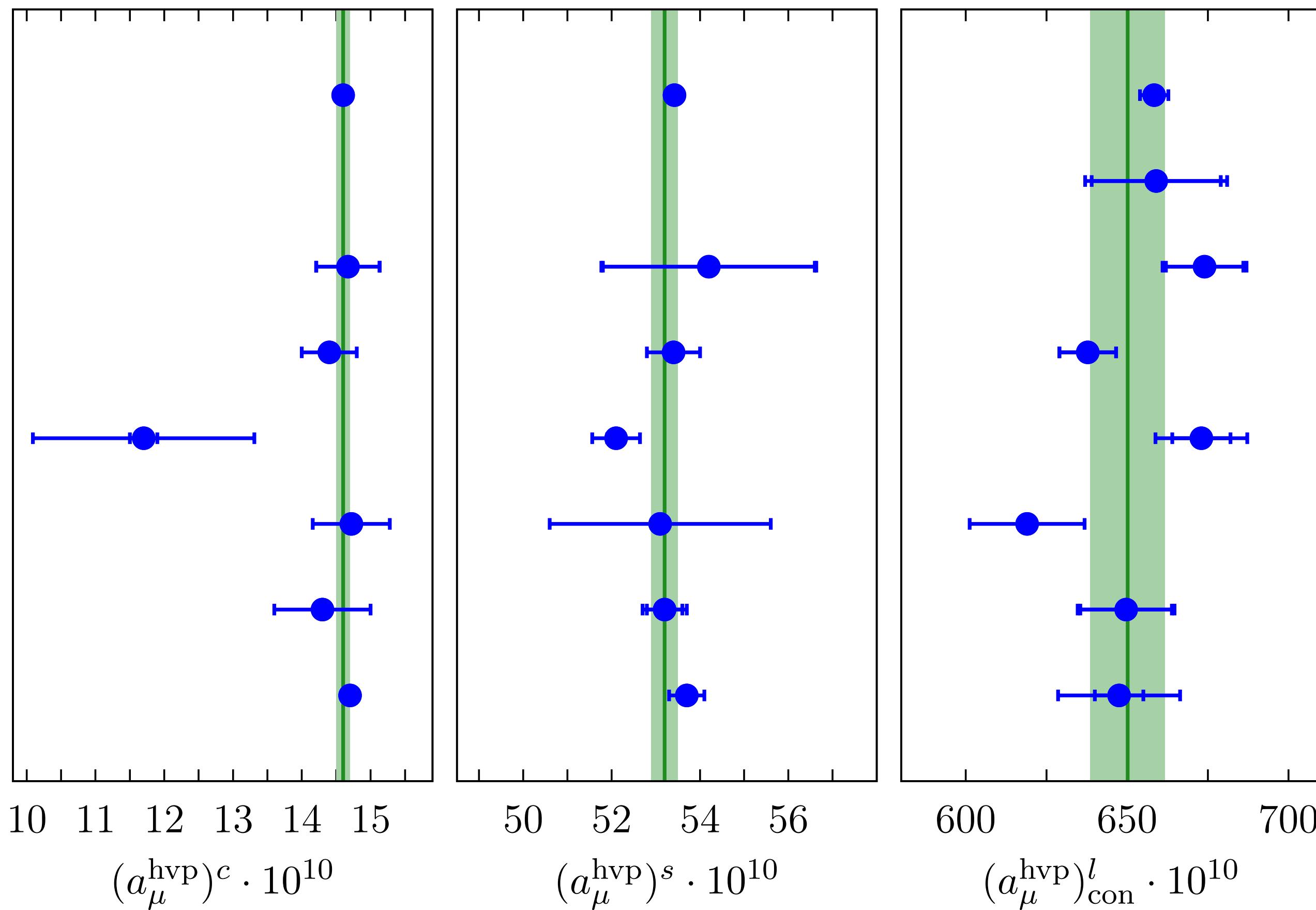
⌚ Monday, 13:00-15:00 US EDT  
🗣 Lukas Varnhorst (BMWc)  
⌚ Thursday 5:00-8:00 US EDT Friday, 5:00–8:00 US EDT  
🗣 Alexander Segner (Mainz) Ben Strassberger (Mainz)

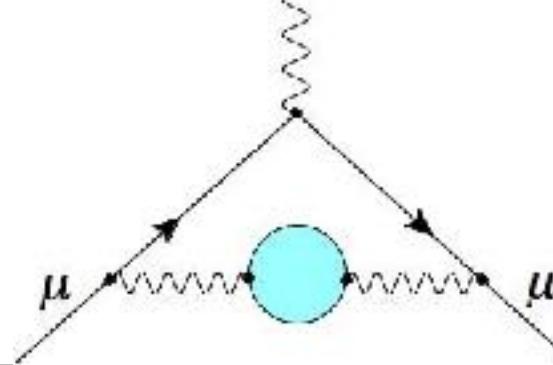


# Lattice HVP: $c, s, \ell$ connected

H. Wittig @ Lattice HVP workshop

$$\sum_f \sim \text{---} \bullet \bar{f} \text{---} f \text{---} \sim \text{---}$$

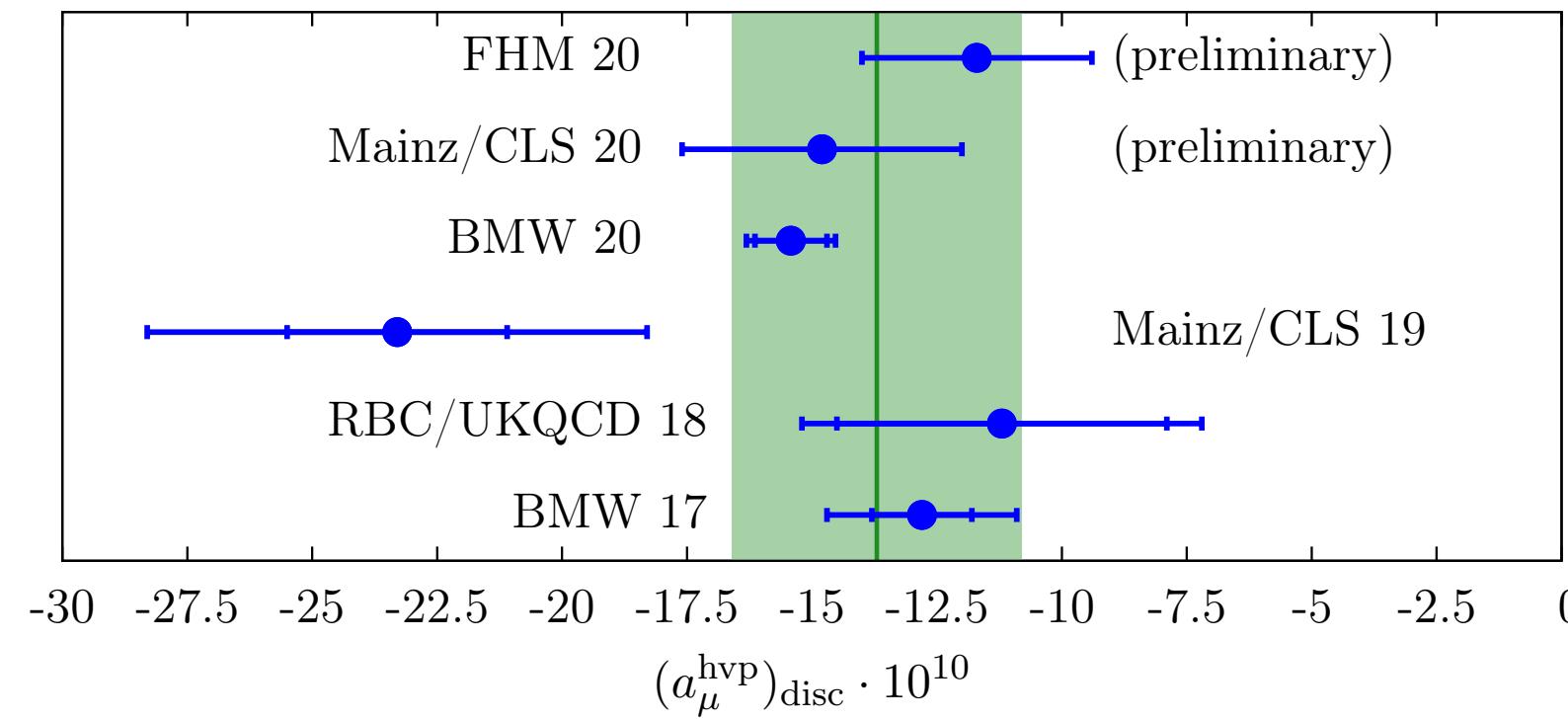




# Lattice HVP: disconnected corrections

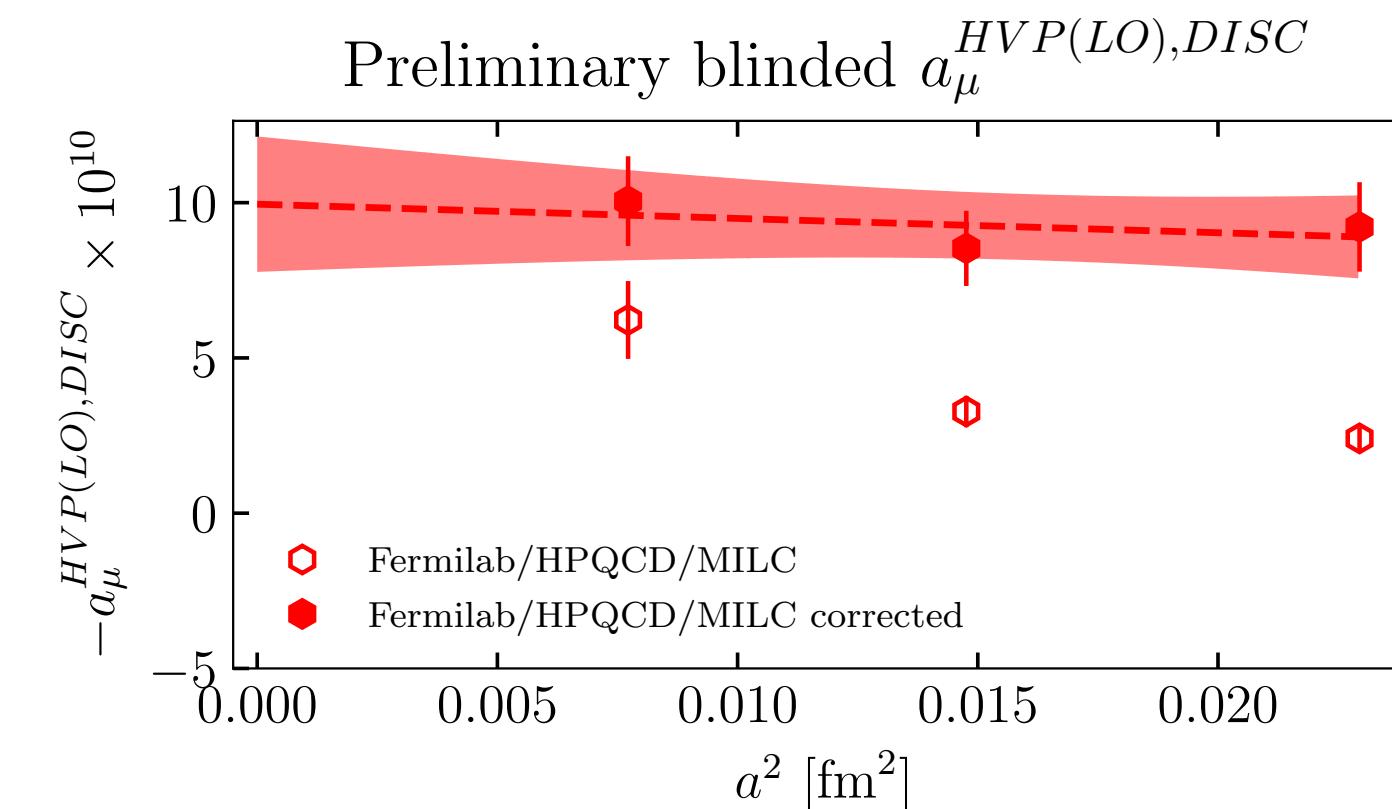


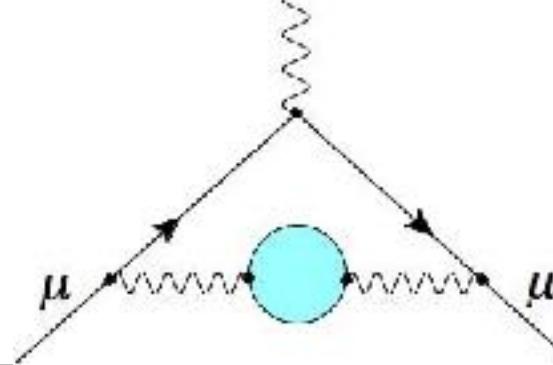
H. Wittig @ Lattice HVP workshop



Convergence towards precise & consistent results

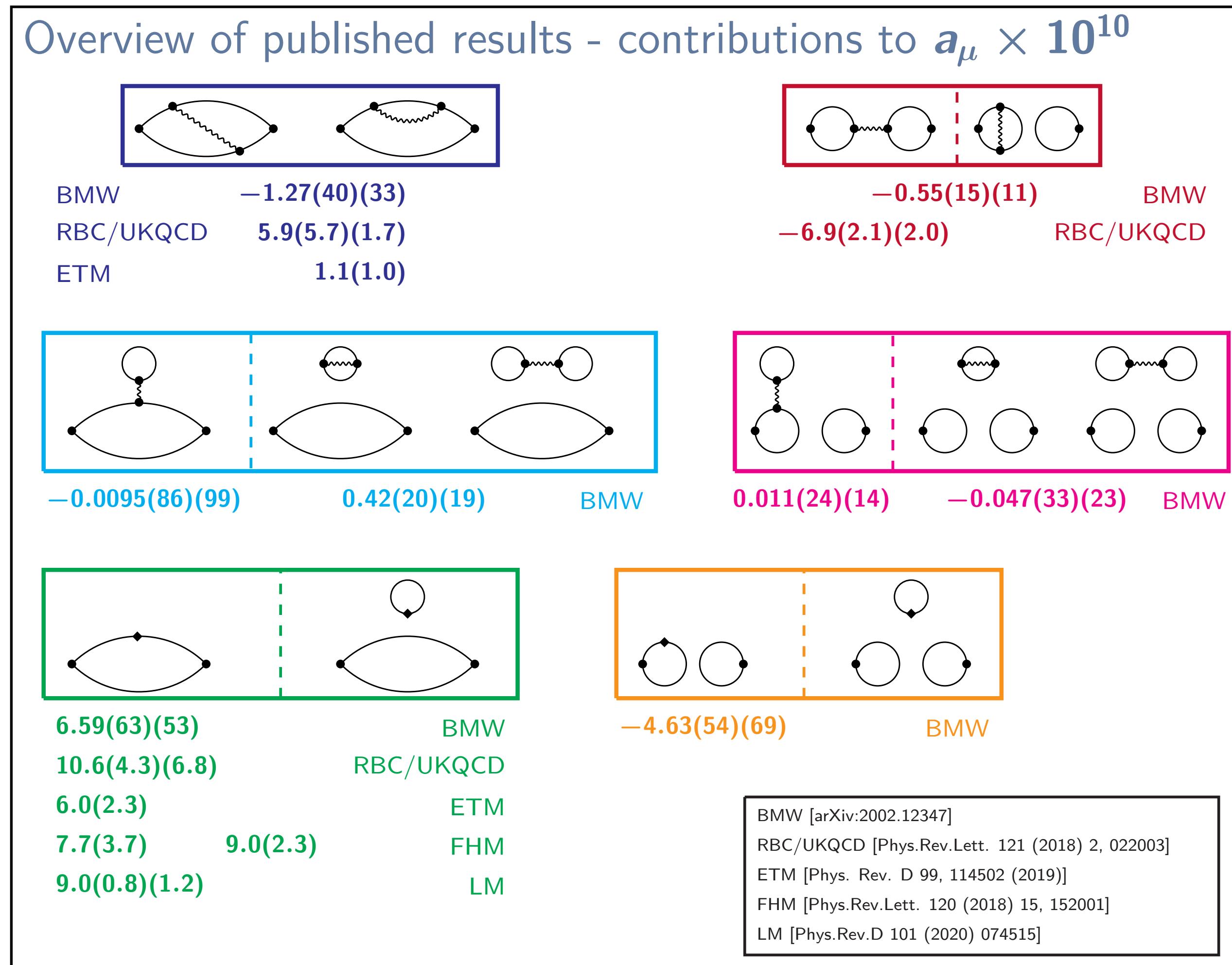
Poster, Wednesday, 8:00-9:00 US EDT  
C. McNeile (Fermilab-HPQCD-MILC)



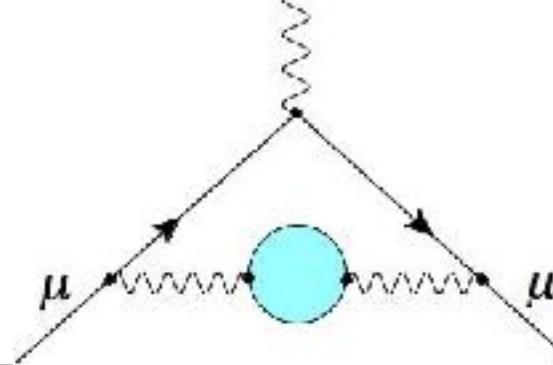


# Lattice HVP: Isospin corrections

V. Gülpers @ Lattice HVP workshop

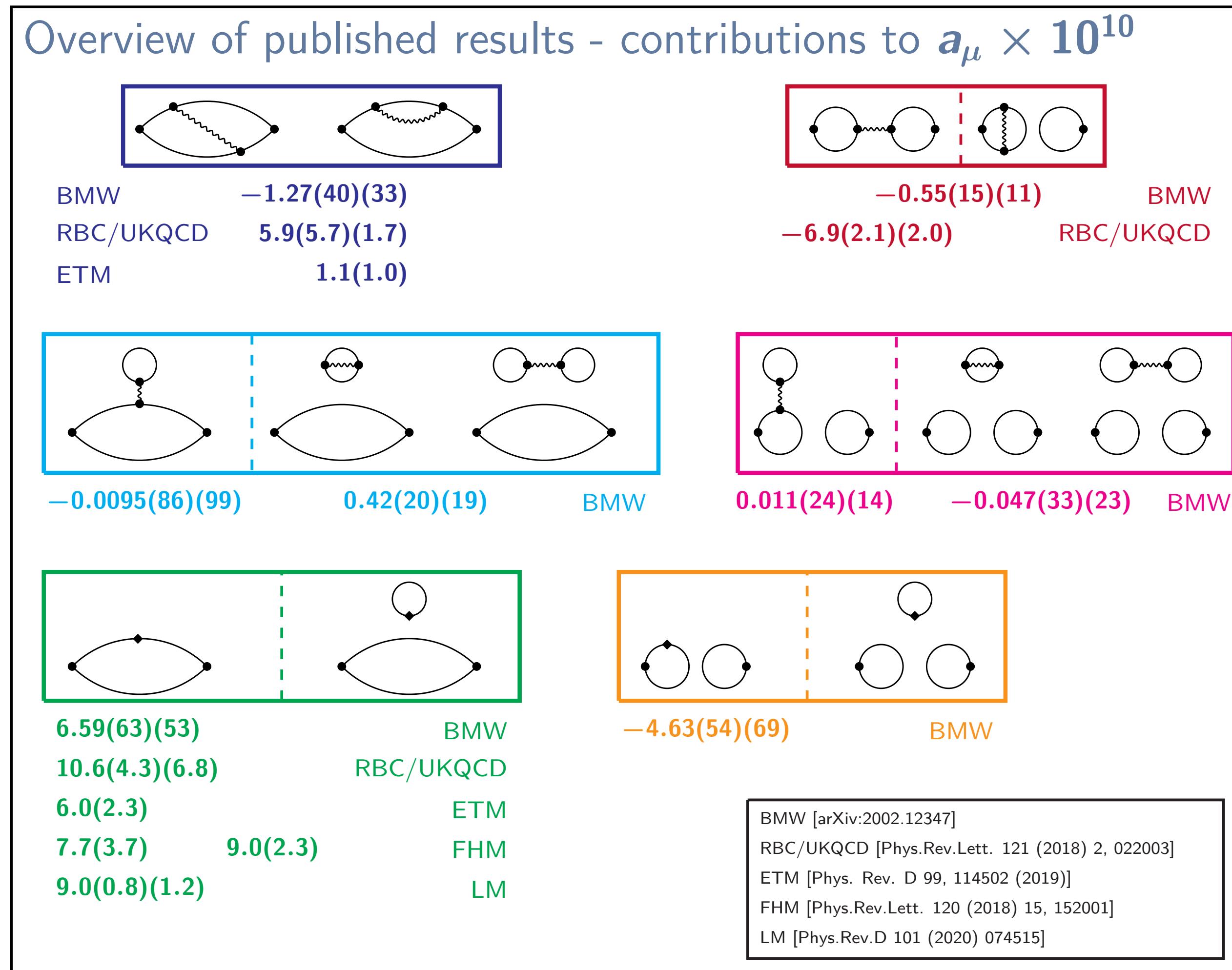


- Some tensions between lattice results for individual contributions.
- Large cancellations between individual contributions:  
 $\delta a_\mu^{\text{IB}} \lesssim 1\%$
- Ongoing efforts presented by Mainz and Fermilab-HPQCD-MILC.



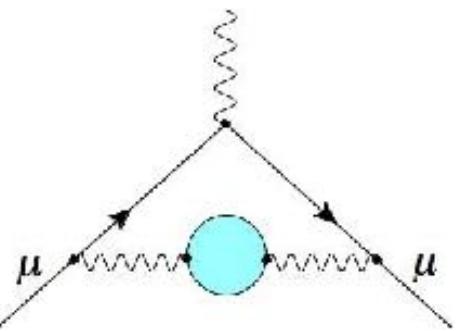
# Lattice HVP: Isospin corrections

V. Gülpers @ Lattice HVP workshop



- Some tensions between lattice results for individual contributions.
- Large cancellations between individual contributions:  
 $\delta a_\mu^{\text{IB}} \lesssim 1\%$
- Ongoing efforts presented by Mainz and Fermilab-HPQCD-MILC.

- ⌚ Tuesday, 5:00-8:00 US EDT
  - ⌚ Andreas Risch (Mainz)
  - ⌚ Letizia Parato (BMWc)
- ⌚ Poster, Wednesday, 8:00-9:00 US EDT
  - ⌚ C. McNeile (Fermilab-HPQCD-MILC)



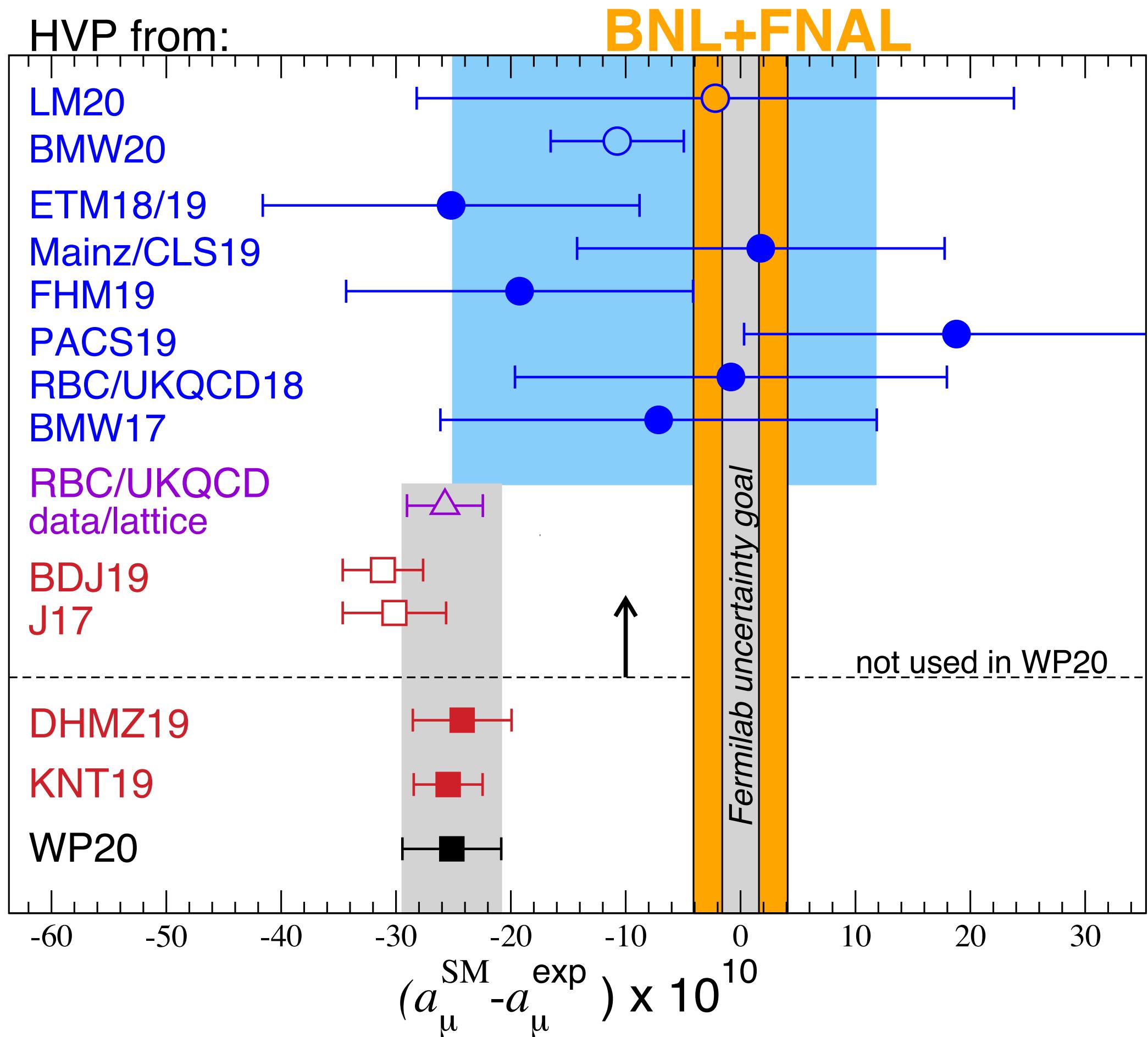
# HVP: Comparison

Lattice QCD + QED

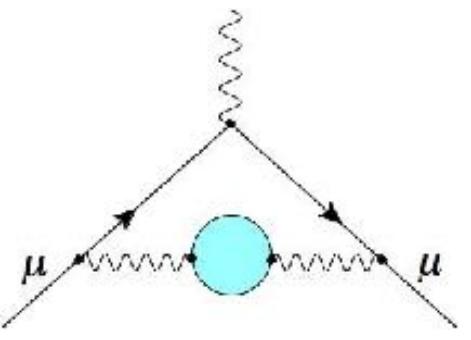
hybrid: combine data & lattice

data driven  
+ unitarity/analyticity  
constraints

$$a_\mu^{\text{SM}} \leftarrow a_\mu^{\text{HVP}} + [a_\mu^{\text{QED}} + a_\mu^{\text{Weak}} + a_\mu^{\text{HLbL}}]$$



Where do we go  
from here?



# Lattice HVP: windows in Euclidean time

$$a_\mu^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \tilde{w}(t) C(t)$$

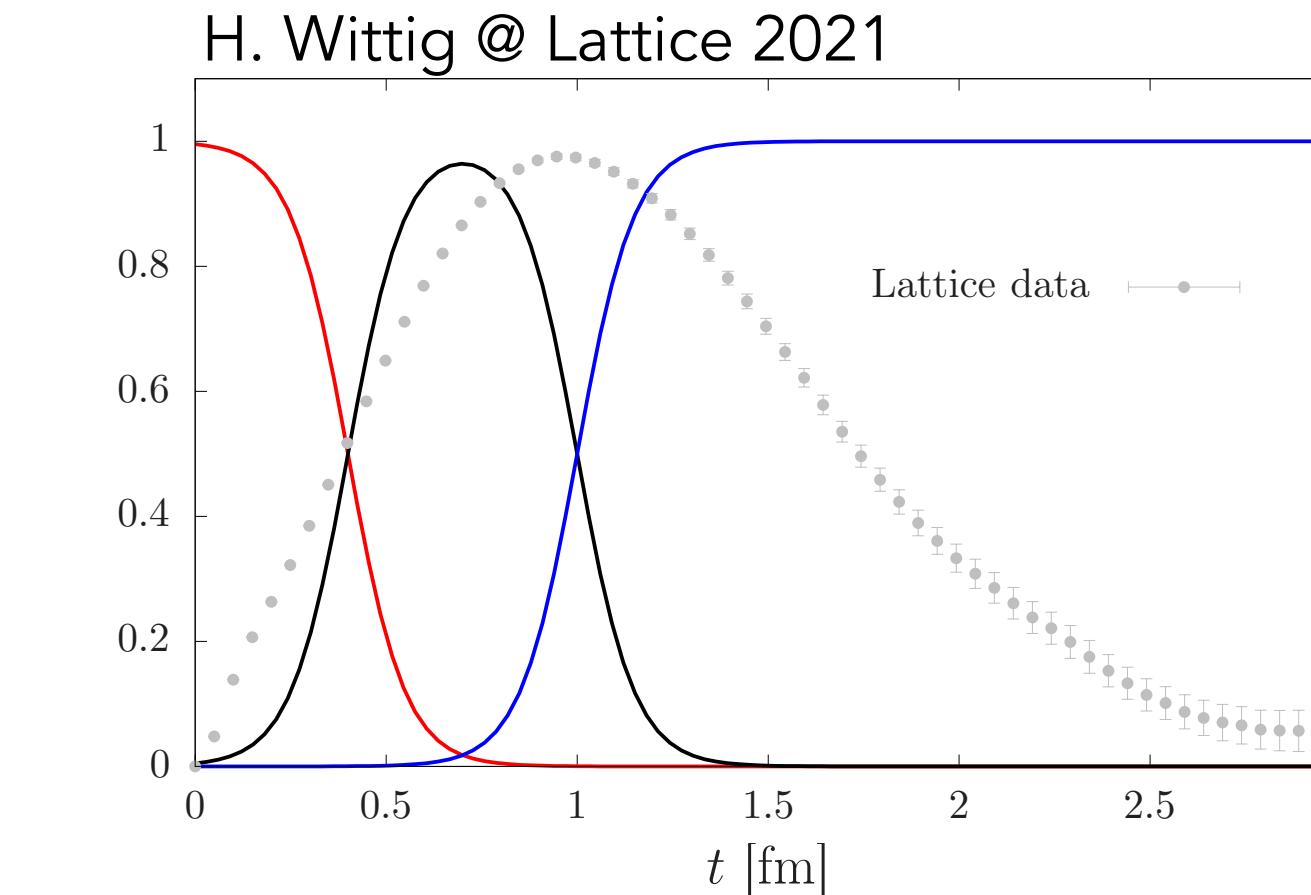
- Use windows in Euclidean time to consider the different time regions separately.

Short Distance (SD)       $t : 0 \rightarrow t_0$

Intermediate (W)       $t : t_0 \rightarrow t_1$

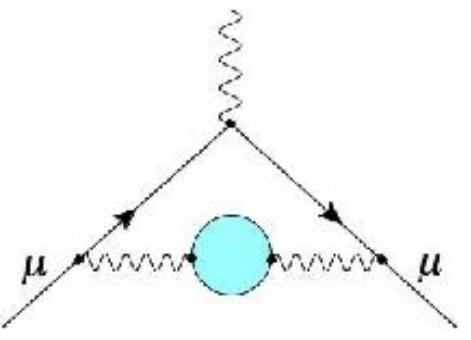
Long Distance (LD)       $t : t_1 \rightarrow \infty$

$$t_0 = 0.4 \text{ fm}, t_1 = 1.0 \text{ fm}$$



- Compute each window separately (in continuum, infinite volume limits,...) and combine

$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

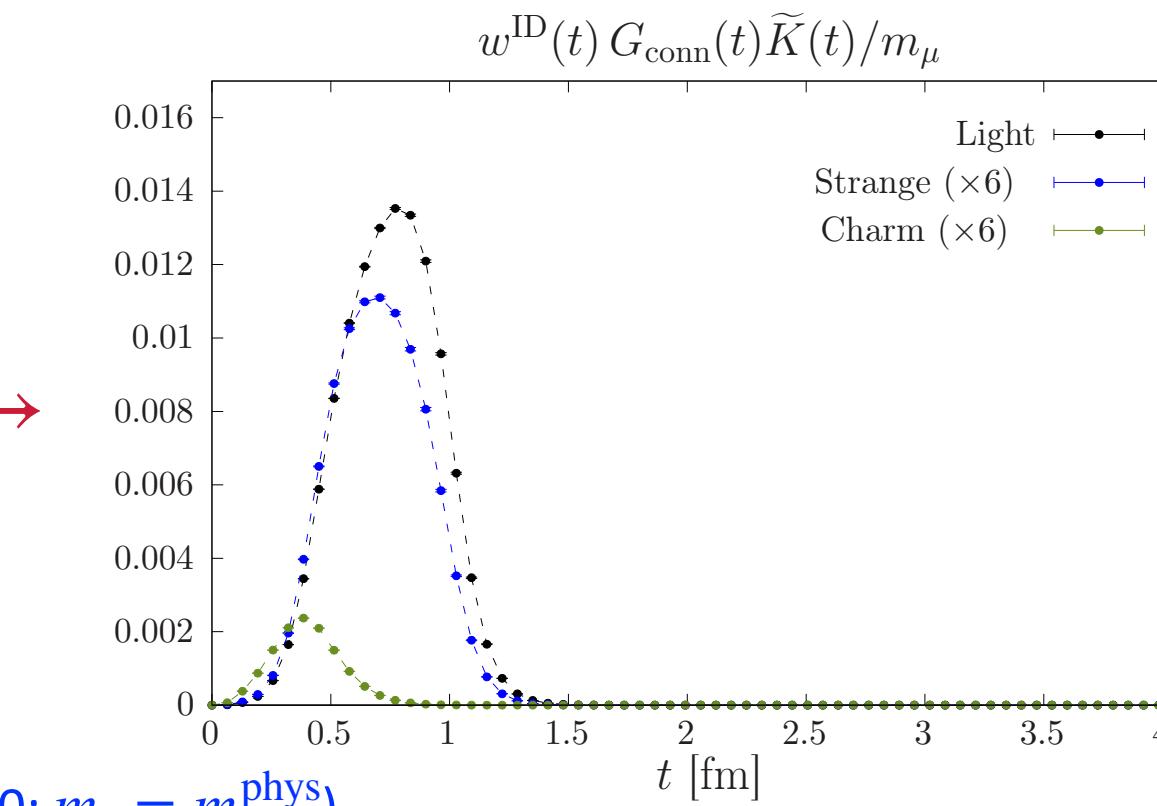
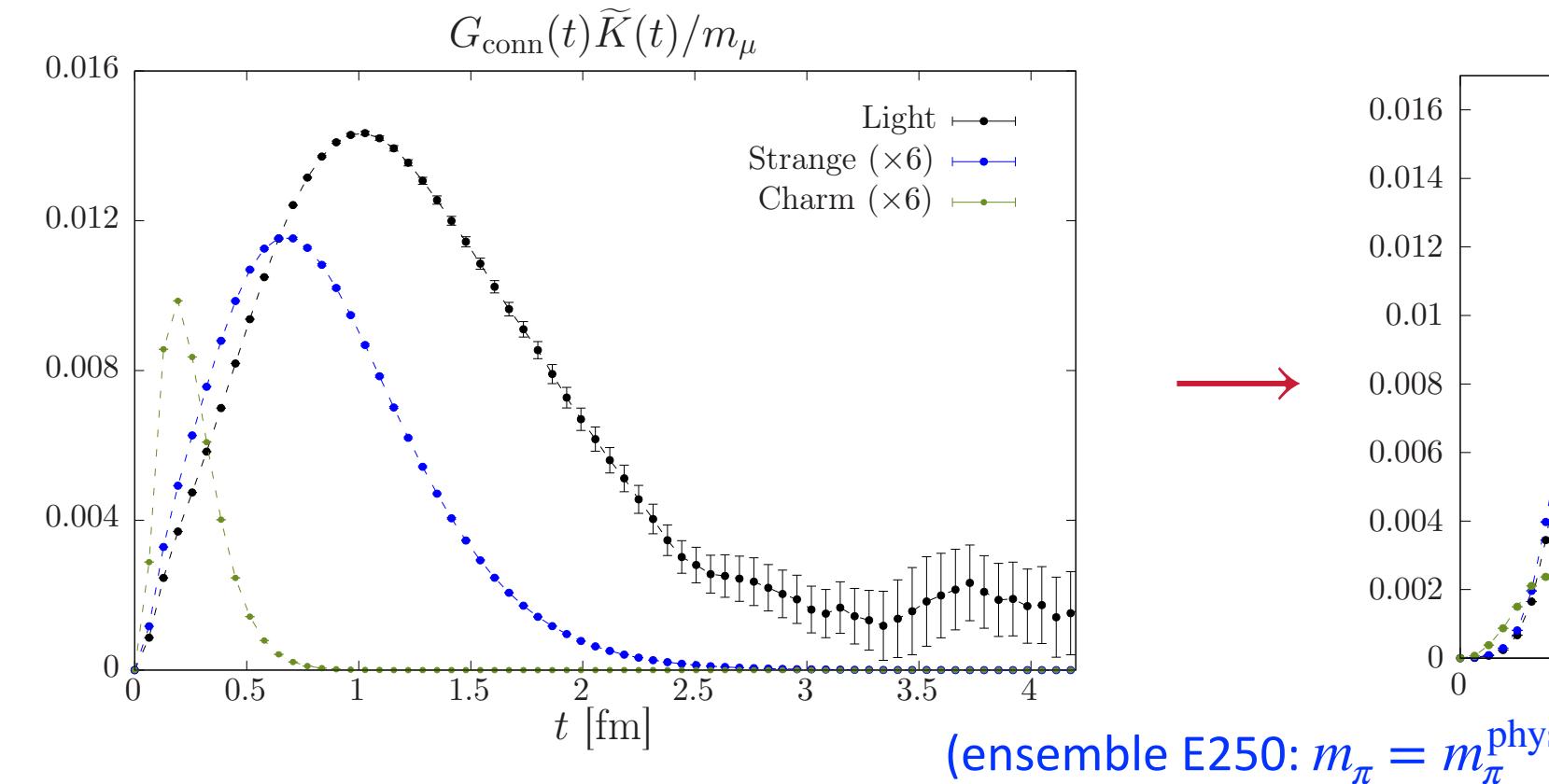


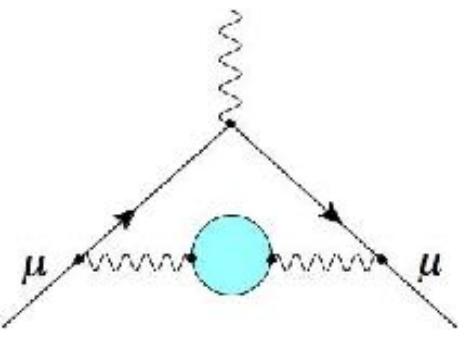
# Lattice HVP: windows in Euclidean time

## Intermediate window

⌚ Tuesday, 5:00-8:00 US EDT  
👤 Hartmut Wittig (Mainz)

Long-distance tail of the integrand is suppressed:



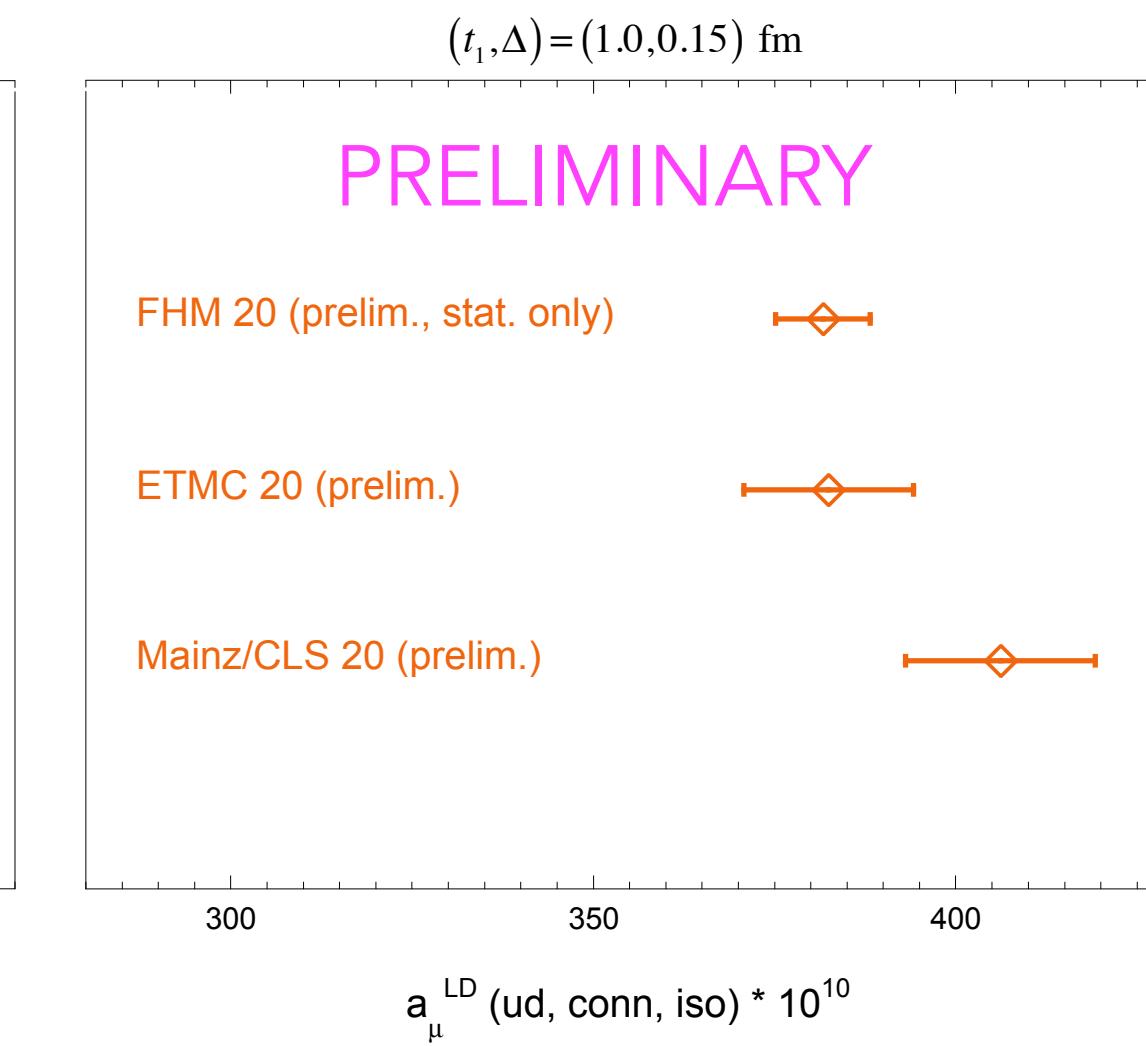
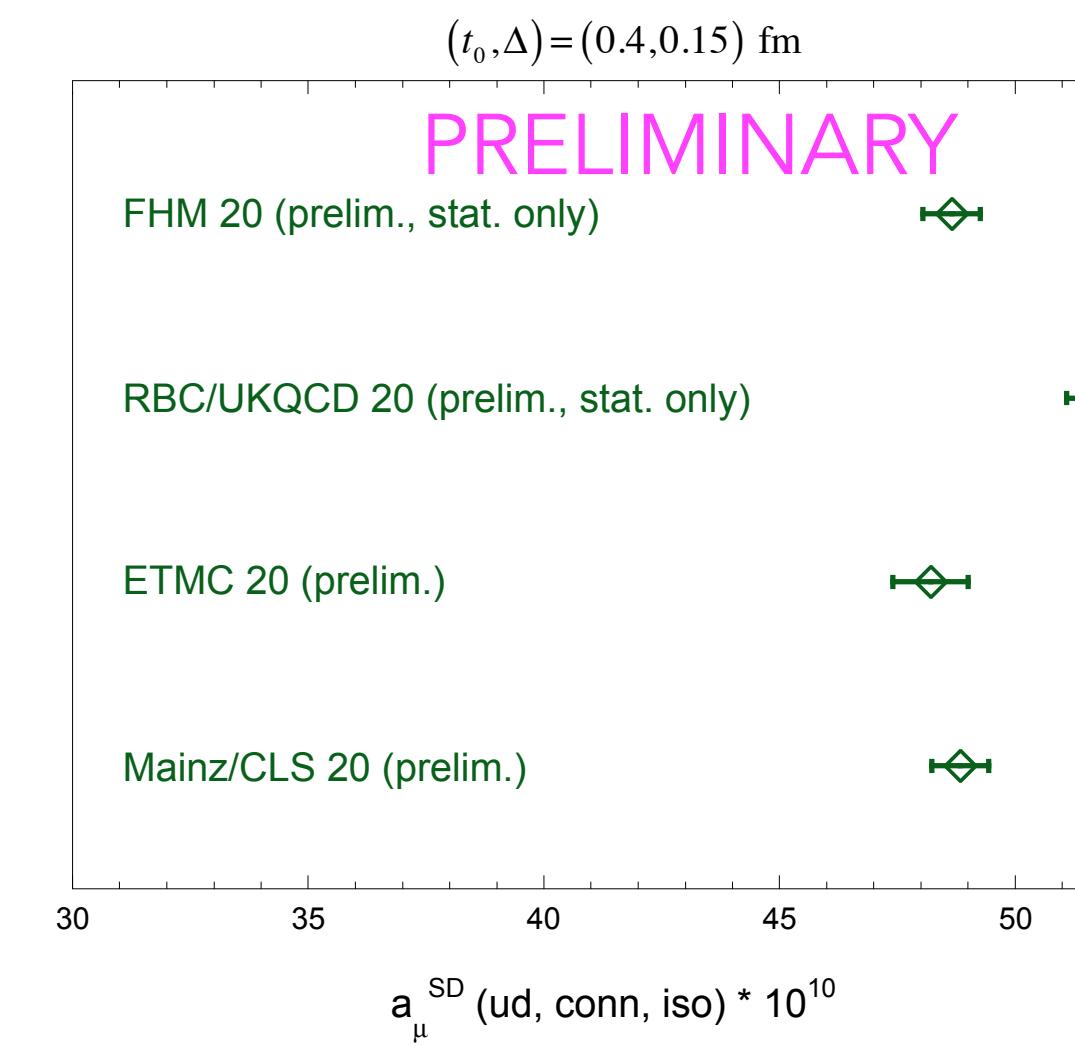
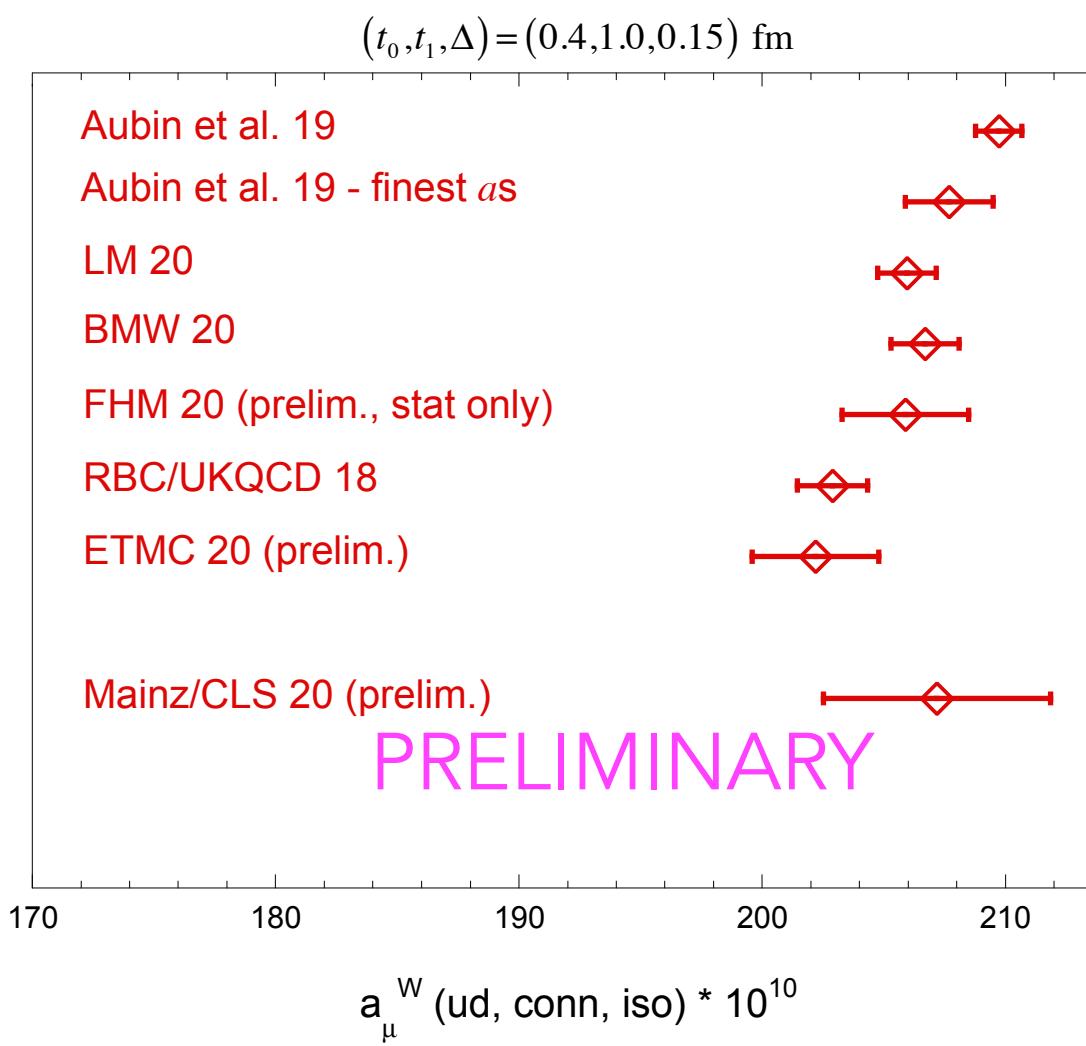


# Lattice HVP: windows in Euclidean time

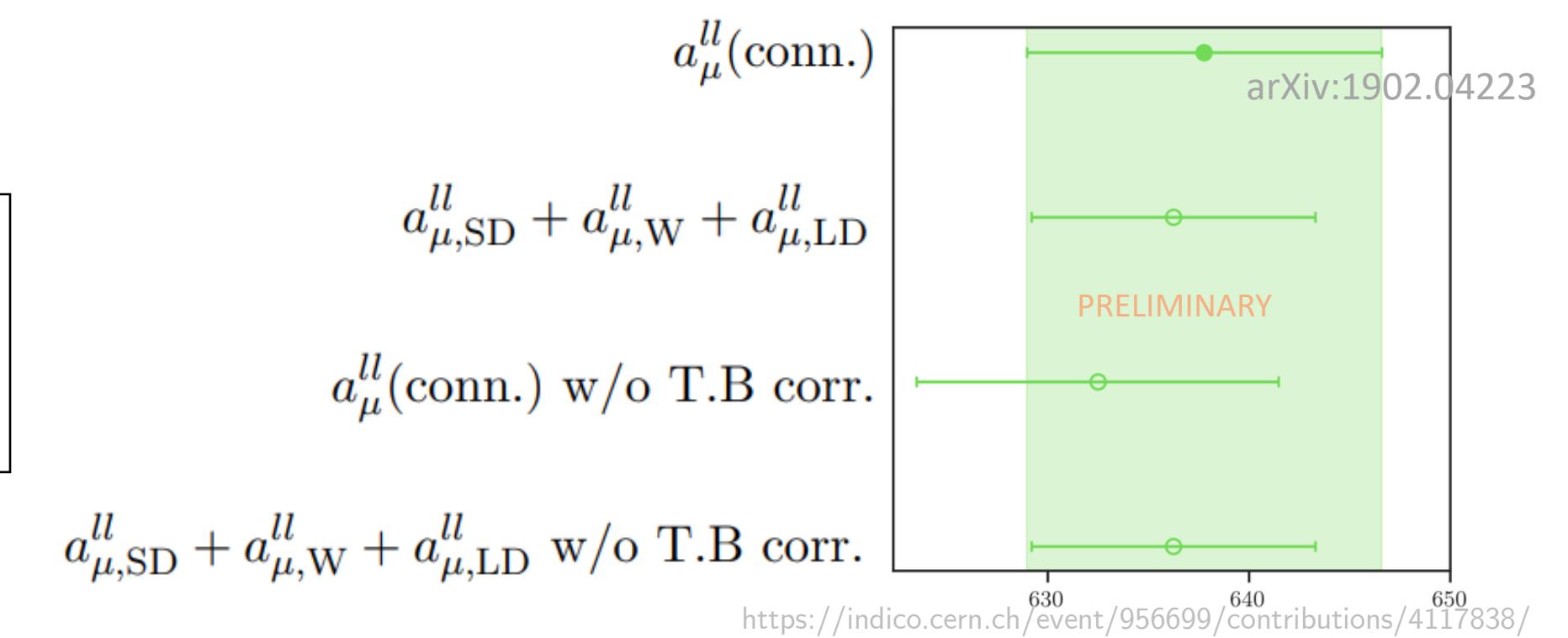
H. Wittig @ Lattice HVP workshop

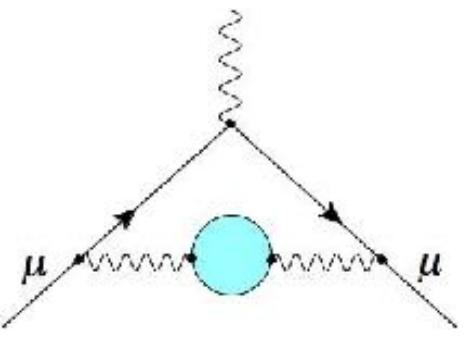
$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

## "Window" quantities



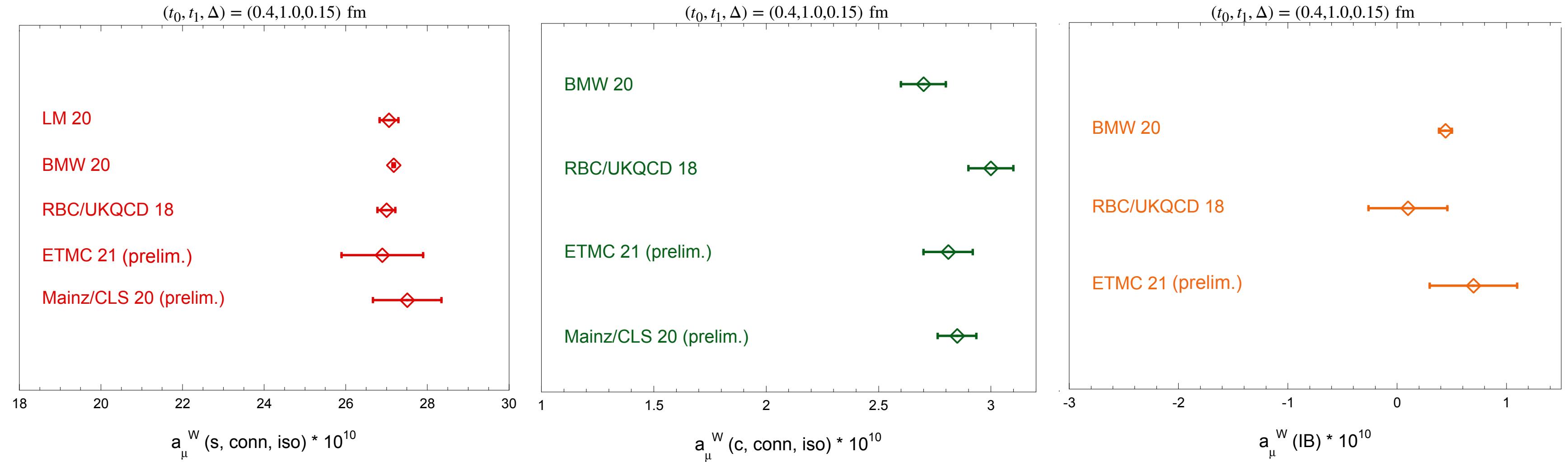
Monday, 13:00-15:00 US EDT  
Shaun Lahert  
(Fermilab-HPQCD-MILC)



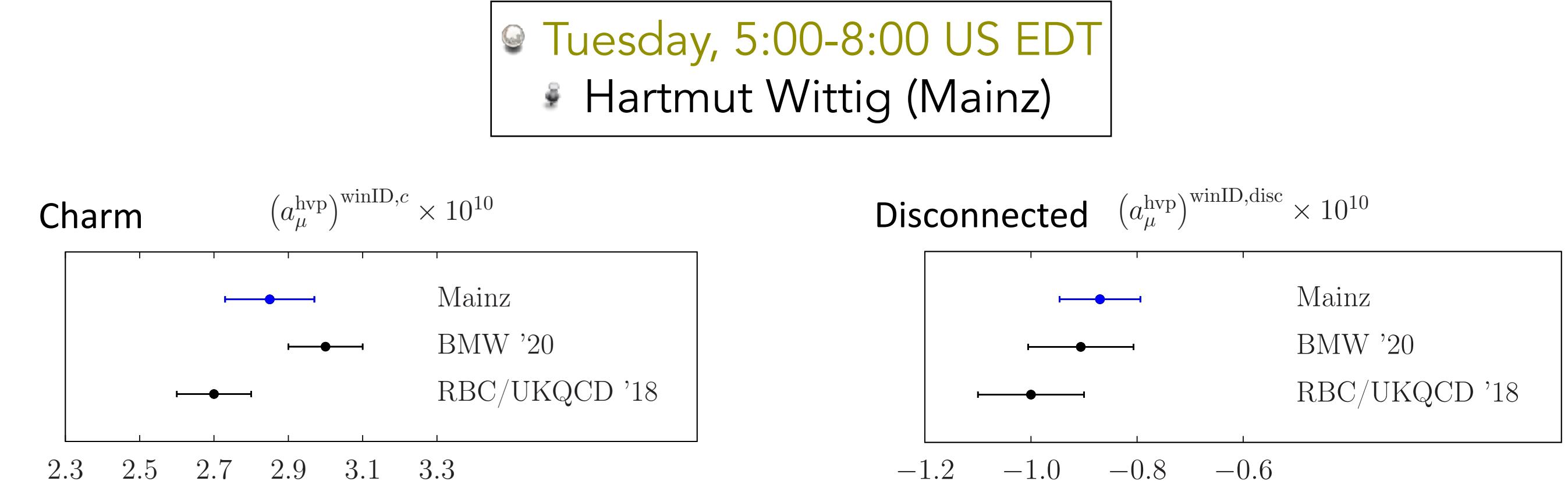
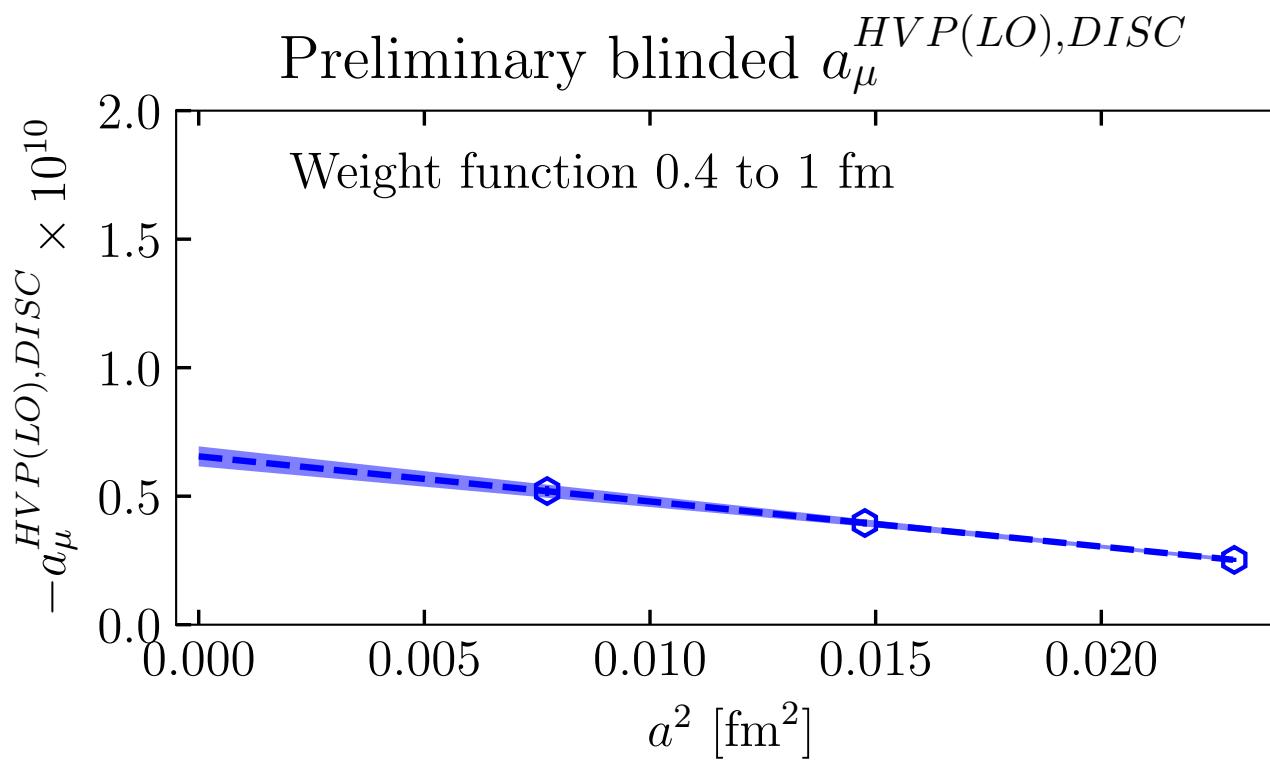


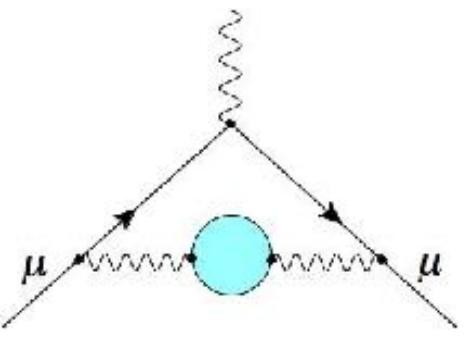
# Lattice HVP: windows in Euclidean time

⌚ Tuesday, 5:00-8:00 US EDT  
🗣️ Davide Giusti (ETMc)



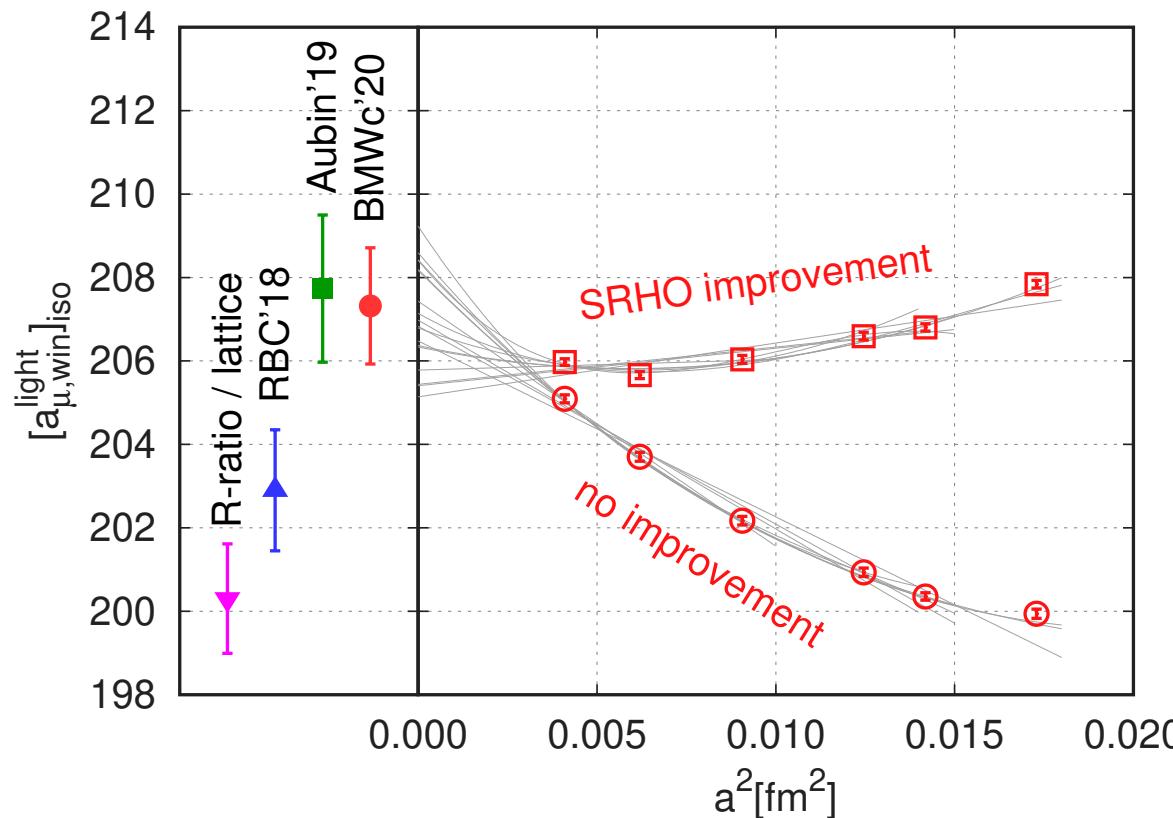
⌚ Poster, Wednesday, 8:00-9:00 US EDT  
🗣️ C. McNeile (Fermilab-HPQCD-MILC)





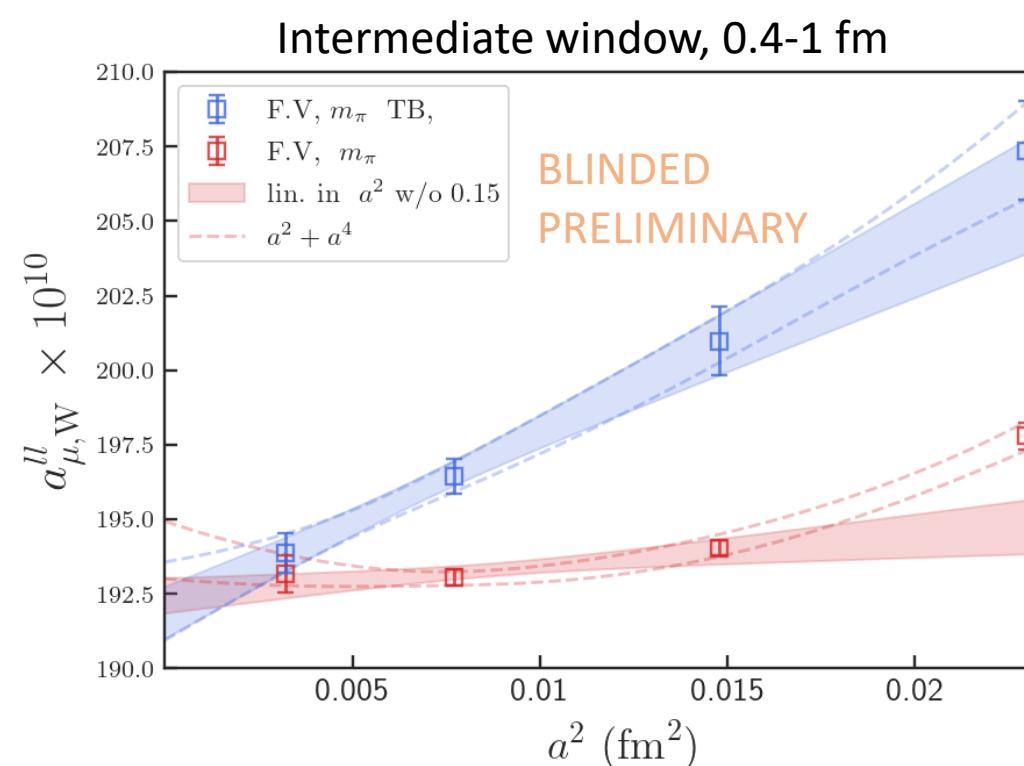
# Lattice HVP: intermediate window (ud)

⌚ Monday, 13:00-15:00 US EDT  
🎤 Kalman Szabo (BMWc)



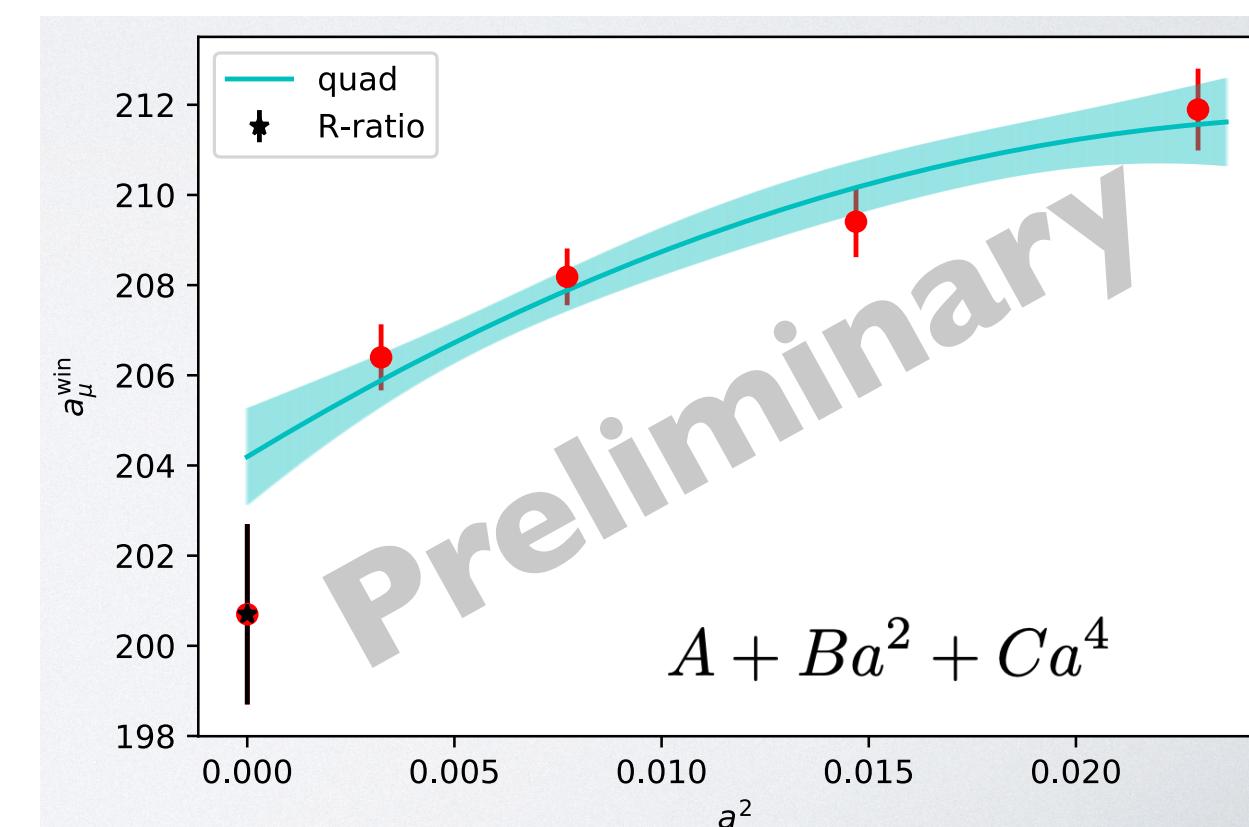
-3.7  $\sigma$  tension with data-driven evaluation  
-2.2  $\sigma$  tension with RBC/UKQCD18

⌚ Monday, 13:00-15:00 US EDT  
🎤 Shaun Lahert  
(Fermilab-HPQCD-MILC)



- ❖ Corrections from  $\rho - \gamma - \pi\pi$  model (leading order).
- ❖ Good consistency between extrapolations of data with(out) discretization effect corrections.

⌚ Monday, 13:00-15:00 US EDT  
🎤 Chris Aubin (Aubin et al)

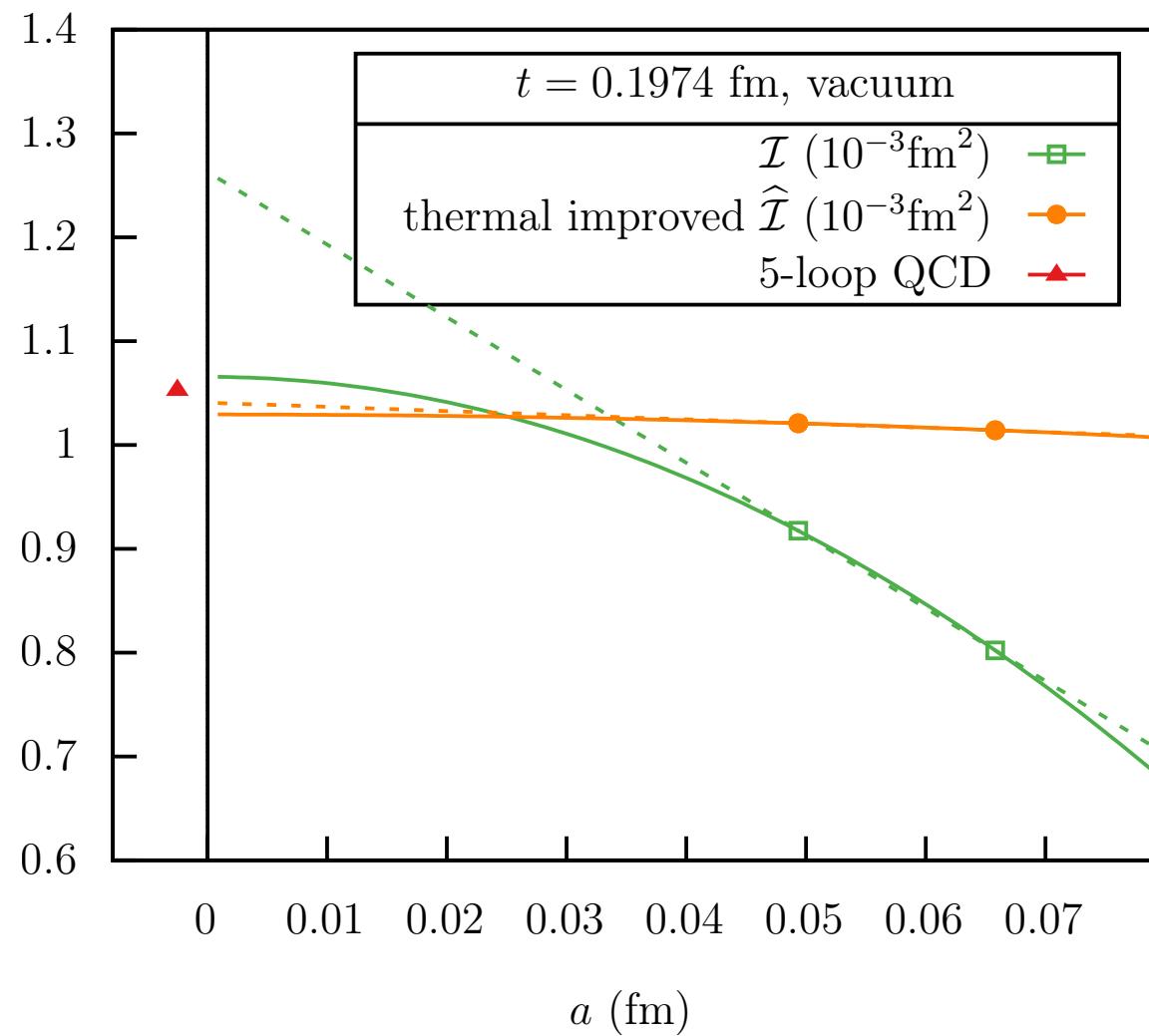


Ongoing work:

- vary functional form of extrapolation, e.g., include  $a_s^m a^{2n}$  terms in expansion
- RBC/UKQCD adding a 3rd lattice spacing
- Blind analyses by Fermilab-HPQCD-MILC and RBC/UKQCD

# Short-distance corrections

⌚ Tuesday, 5:00-8:00 US EDT  
🎙 Tim Harris (NEPhEU QCD)



- Use thermal observables and gauge ensembles at finite temperature and very fine lattice spacings to resolve discretization effects in calculations of short-distance quantities
- Construct improved observables with better smaller discretization effects
- Can be applied to SD window and  $\Delta\alpha$
- New step-scaling method for computing  $\Delta\alpha$  at large  $Q^2$

⌚ Friday, 5:00-8:00 US EDT  
🎙 Nicolai Husung (DESY)

- Compute anomalous dimensions of higher dimensional operators in Symanzik EFT
- Use to guide continuum extrapolation
- In most cases studied:  $\hat{\Gamma} \gtrsim 0$

# Summary

---

- ★ The QED and EW contributions are known very precisely  
Hadronic contributions determine the uncertainty in the SM prediction.
- ★ dispersive HVP: ~0.6% error [0.34ppm]  
based on well-tested experimental data, will be improved with new measurements (coming soon).
- ★ lattice HVP: first LQCD calculation with sub-percent uncertainty by BMWc  
but in tension with data-driven approach  
a lot of activity and progress, expect more sub-precision results soon.
- ★ dispersive HLBL: ~20% error [0.15ppm]  
newly developed dispersive approach with almost fully quantified errors  
systematically improvable
- ★ lattice HLBL: two complete lattice calculations  
consistent with each other and with data-driven result  
systematically improvable

# Outlook

- ★ Theory Initiative:

- a “WP process” for Lattice HVP and HLbL results

- ★ Starting to developing guidelines for assessing lattice HVP calculations
  - panel discussion at the KEK workshop

- ★ Quantities to calculate:

- windows for lattice-only cross checks for all components

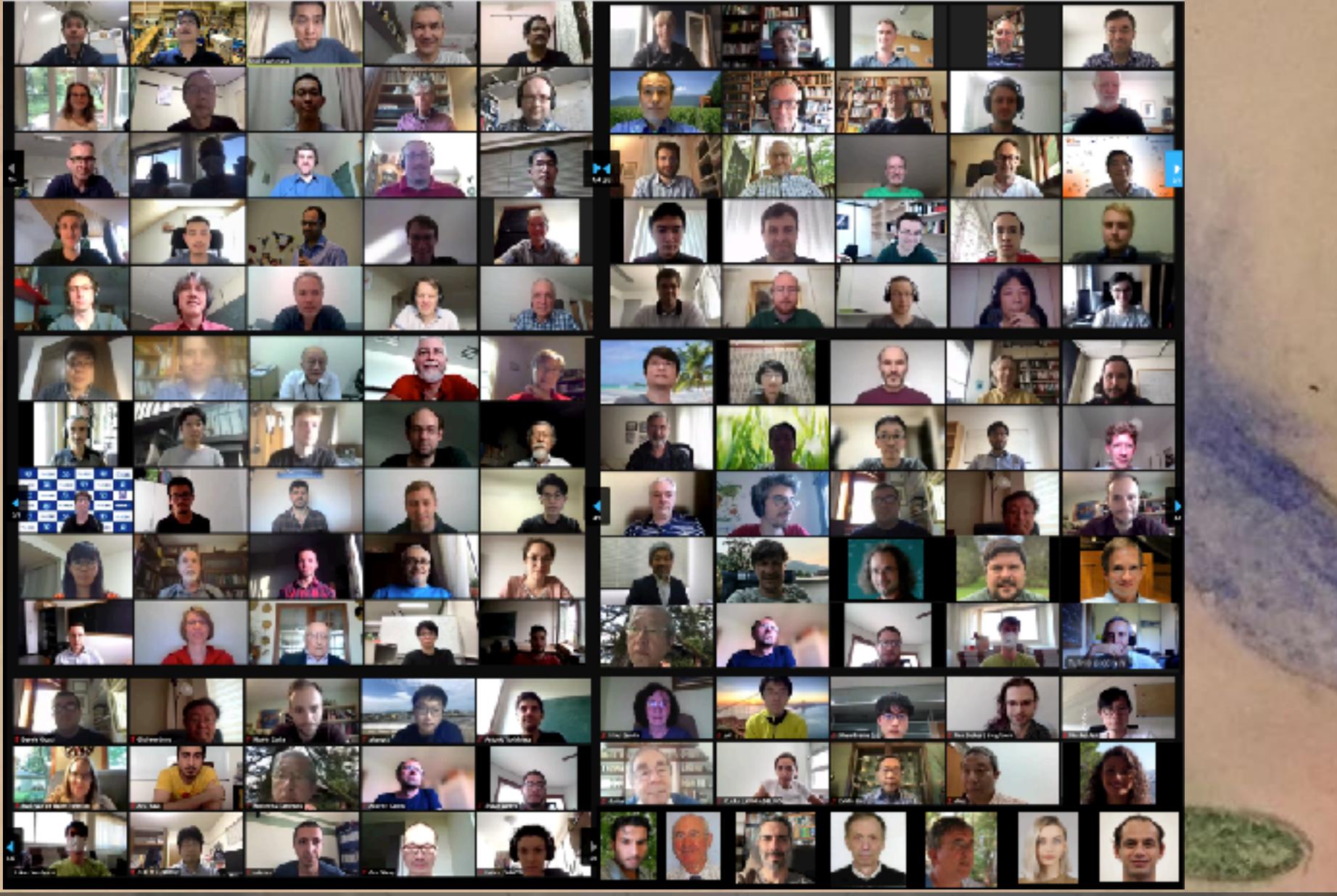
$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

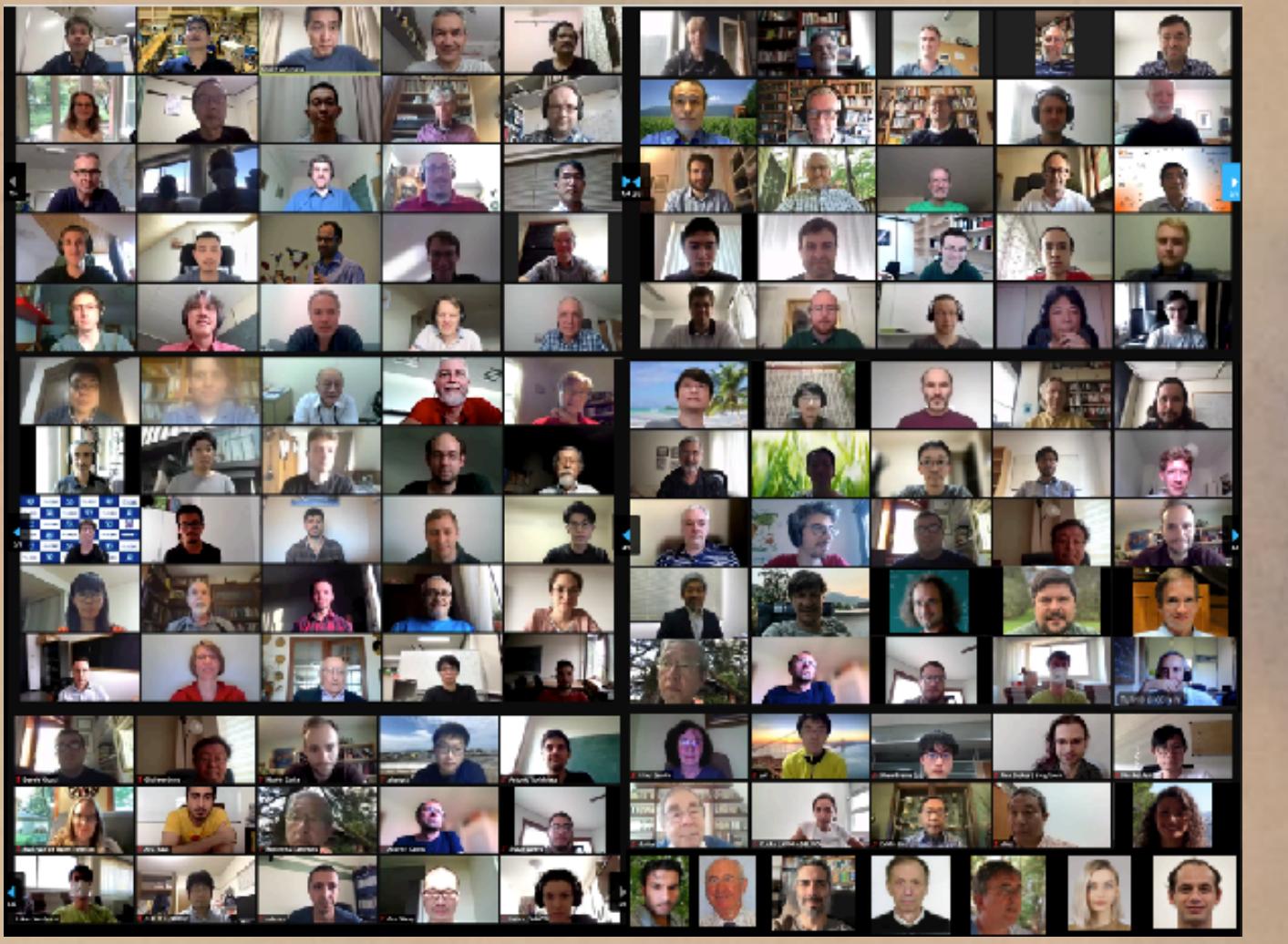
- $l=1,0$  (light quark)

- derivatives of  $a_\mu$  (or windows) w.r.t. parameters

- ★ Prescription for defining isospin limit and separating QED & SIB contributions

- ★ Blinding lattice calculations is good practice to avoid unintended bias





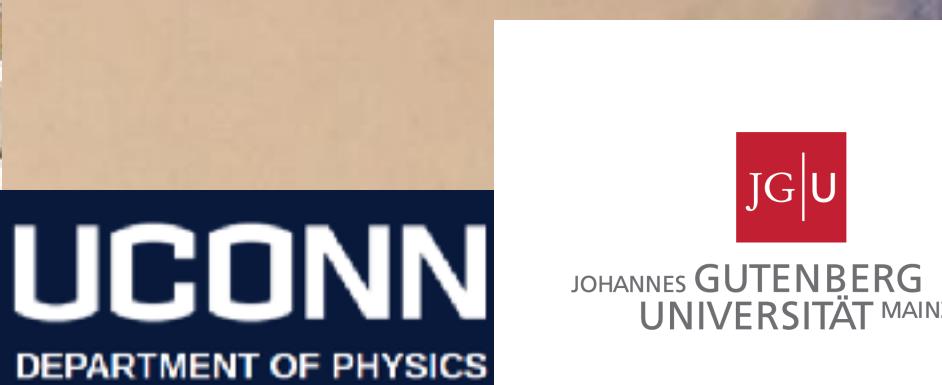
UNIVERSITY of  
WASHINGTON



Office of  
Science



Thank you!



# Appendix

# Muon g-2 Theory Initiative

## WP section authors:

- [Section 2: Data-driven evaluations of HVP](#)

M. Benayoun, C. M. Carloni Calame, H. Czyz, M. Davier, S. I. Eidelman, M. Hoferichter, F. Jegerlehner, A. Keshavarzi, B. Malaescu, D. Nomura, M. Passera, T. Teubner, G. Venanzoni, Z. Zhang

- [Section 3: Lattice QCD calculations of HVP](#)

T. Blum, M. Bruno, M. Ce, C. T. H. Davies, M. Della Morte, A. X. El-Khadra, D. Giusti, Steven Gottlieb, V. Guelpers, G. Herdoiza, T. Izubuchi, C. Lehner, L. Lellouch, M. K. Marinkovic, A. S. Meyer, K. Miura, A. Portelli, S. Simula, R. Van de Water, G. von Hippel, H. Wittig

- [Section 4: Data-driven and dispersive approach to HLbL](#)

J. Bijnens, G. Colangelo, F. Curciarello, H. Czyz, I. Danilkin, F. Hagelstein, M. Hoferichter, B. Kubis, A. Kupsc, A. Nyffeler, V. Pascalutsa, E. Perez del Rio, M. Procura, C. F. Redmer, P. Sanchez-Puertas, P. Stoffer, M. Vanderhaeghen

- [Section 5: Lattice approaches to HLbL](#)

N. Asmussen, T. Blum, A. Gerardin, M. Hayakawa, R. J. Hudspith, T. Izubuchi, L. Jin, C. Lehner, H. B. Meyer, A. Nyffeler

- [Section 6: The QED contributions to  \$a\_\mu\$](#)

T. Aoyama, T. Kinoshita, M. Nio

- [Section 7: The electroweak contributions to  \$a\_\mu\$](#)

D. Stoeckinger, H. Stoeckinger-Kim

# Updated WP Summary Table

Contribution	Value $\times 10^{11}$	References
Experimental average (E989+E821)	116592061(41)	<a href="#">Phys.Rev.Lett. 124, 141801</a>
HVP LO ( $e^+e^-$ )	6931(40)	Refs. [2–7]
HVP NLO ( $e^+e^-$ )	-98.3(7)	Ref. [7]
HVP NNLO ( $e^+e^-$ )	12.4(1)	Ref. [8]
HVP LO (lattice, $uds\bar{c}$ )	7116(184)	Refs. [9–17]
HLbL (phenomenology)	92(19)	Refs. [18–30]
HLbL NLO (phenomenology)	2(1)	Ref. [31]
HLbL (lattice, $uds$ )	79(35)	Ref. [32]
HLbL (phenomenology + lattice)	90(17)	Refs. [18–30, 32]
QED	116 584 718.931(104)	Refs. [33, 34]
Electroweak	153.6(1.0)	Refs. [35, 36]
HVP ( $e^+e^-$ , LO + NLO + NNLO)	6845(40)	Refs. [2–8]
HLbL (phenomenology + lattice + NLO)	92(18)	Refs. [18–32]
Total SM Value	116 591 810(43)	Refs. [2–8, 18–24, 31–36]
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	251(59)	

website: <https://muon-gm2-theory.illinois.edu>

# Connections

$$\sigma(e^+e^- \rightarrow \text{hadrons}) \Leftrightarrow a_\mu^{\text{HVP}} \Leftrightarrow \Delta\alpha_{\text{had}}(M_Z^2)$$

- $\Delta\alpha_{\text{had}}(M_Z^2)$  also depends on the hadronic vacuum polarization function, and can be written as an integral over  $\sigma(e^+e^- \rightarrow \text{hadrons})$ , but weighted towards higher energies.
- a shift in  $a_\mu^{\text{HVP}}$  also changes  $\Delta\alpha_{\text{had}}(M_Z^2)$ :  $\Rightarrow$  EW fits  
[Crivellin et al 2020, Keshavarsi et al 2020, Malaescu & Scott 2020]  
If the shift is due to differences in the low ( $\lesssim 2 \text{ GeV}$ ) energy region, the impact on  $\Delta\alpha_{\text{had}}(M_Z^2)$  and EW fits is small.
- A shift in  $a_\mu^{\text{HVP}}$  from low ( $\lesssim 2 \text{ GeV}$ ) energies  $\Rightarrow \sigma(e^+e^- \rightarrow \pi\pi)$  must satisfy unitarity & analyticity constraints  $\Rightarrow F_\pi^V(s)$  can be tested with lattice calculations  
[Colangelo, Hoferichter, Stoffer 2021]

# Connections

Martin Hoferichter @ Lattice HVP workshop

## Hadronic running of $\alpha$ and global EW fit

$e^+ e^-$	KNT, DHMZ	EW fit <a href="#">HEPFit</a>	EW fit <a href="#">GFitter</a>	guess based on <a href="#">BMWc</a>
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \times 10^4$	276.1(1.1)	270.2(3.0)	271.6(3.9)	277.8(1.3)
difference to $e^+ e^-$		$-1.8\sigma$	$-1.1\sigma$	$+1.0\sigma$

- **Time-like formulation:**

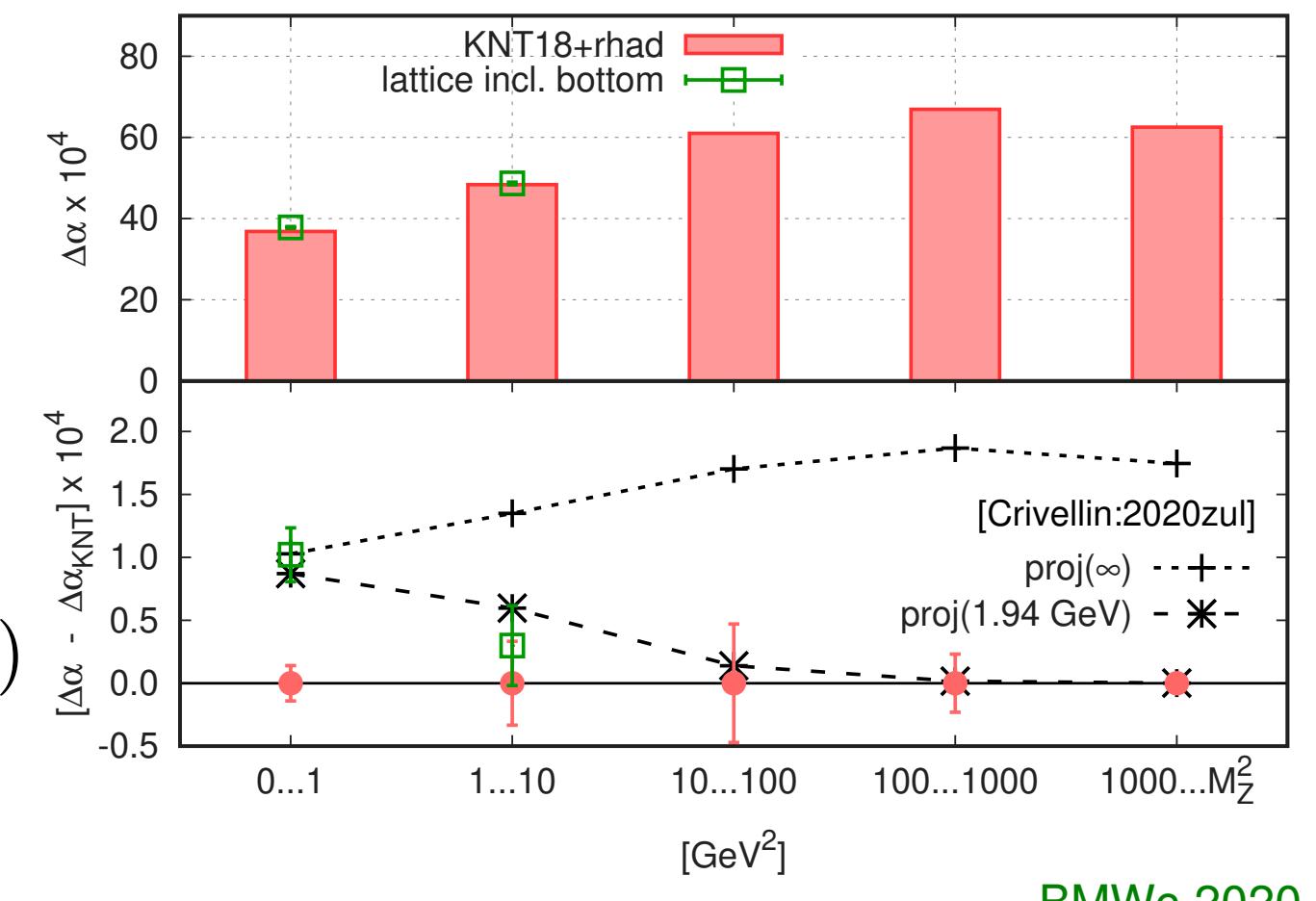
$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = \frac{\alpha M_Z^2}{3\pi} P \int_{s_{\text{thr}}}^{\infty} ds \frac{R_{\text{had}}(s)}{s(M_Z^2 - s)}$$

- **Space-like formulation:**

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = \frac{\alpha}{\pi} \hat{\Pi}(-M_Z^2) + \frac{\alpha}{\pi} (\hat{\Pi}(M_Z^2) - \hat{\Pi}(-M_Z^2))$$

- Global EW fit

- Difference between [HEPFit](#) and [GFitter](#)  
implementation mainly treatment of  $M_W$
- Pull goes into **opposite direction**



BMWc 2020

More in talks by M. Passera, B. Malaescu (phenomenology)  
and K. Miura, T. San José (lattice)

# Connections

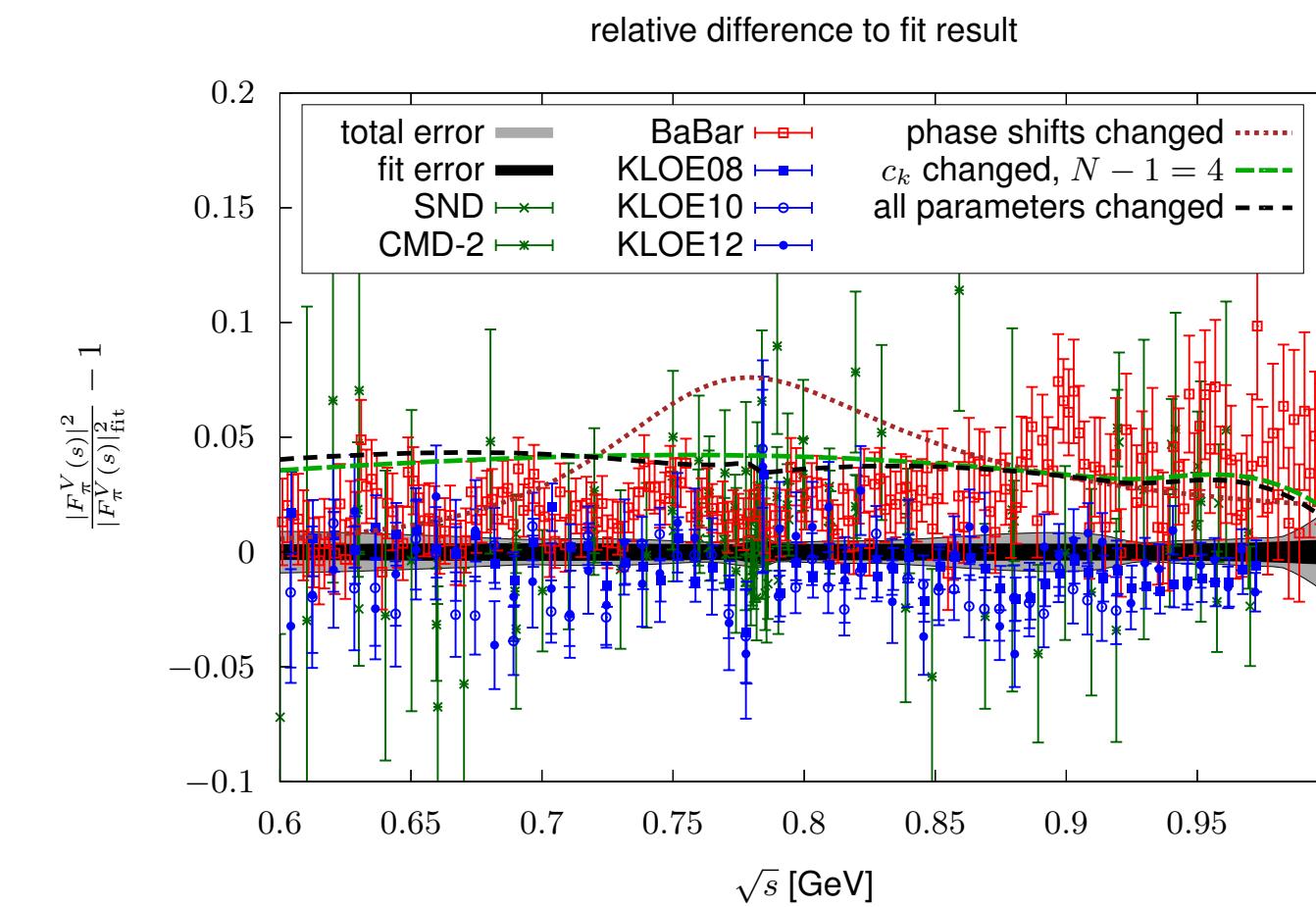
Peter Stoffer @ Lattice HVP workshop

Constraints on the two-pion contribution to HVP

arXiv:2010.07943 [hep-ph]

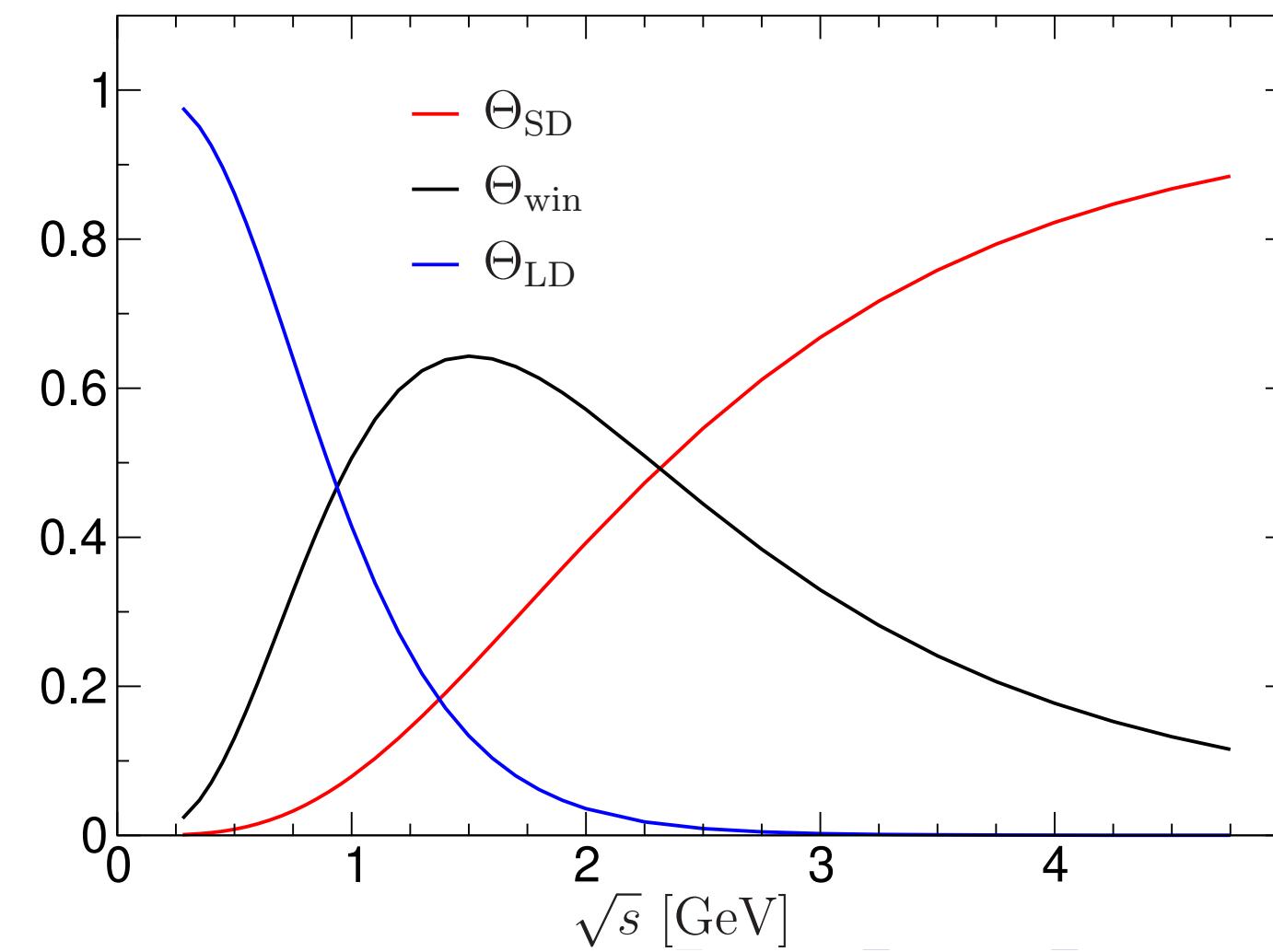
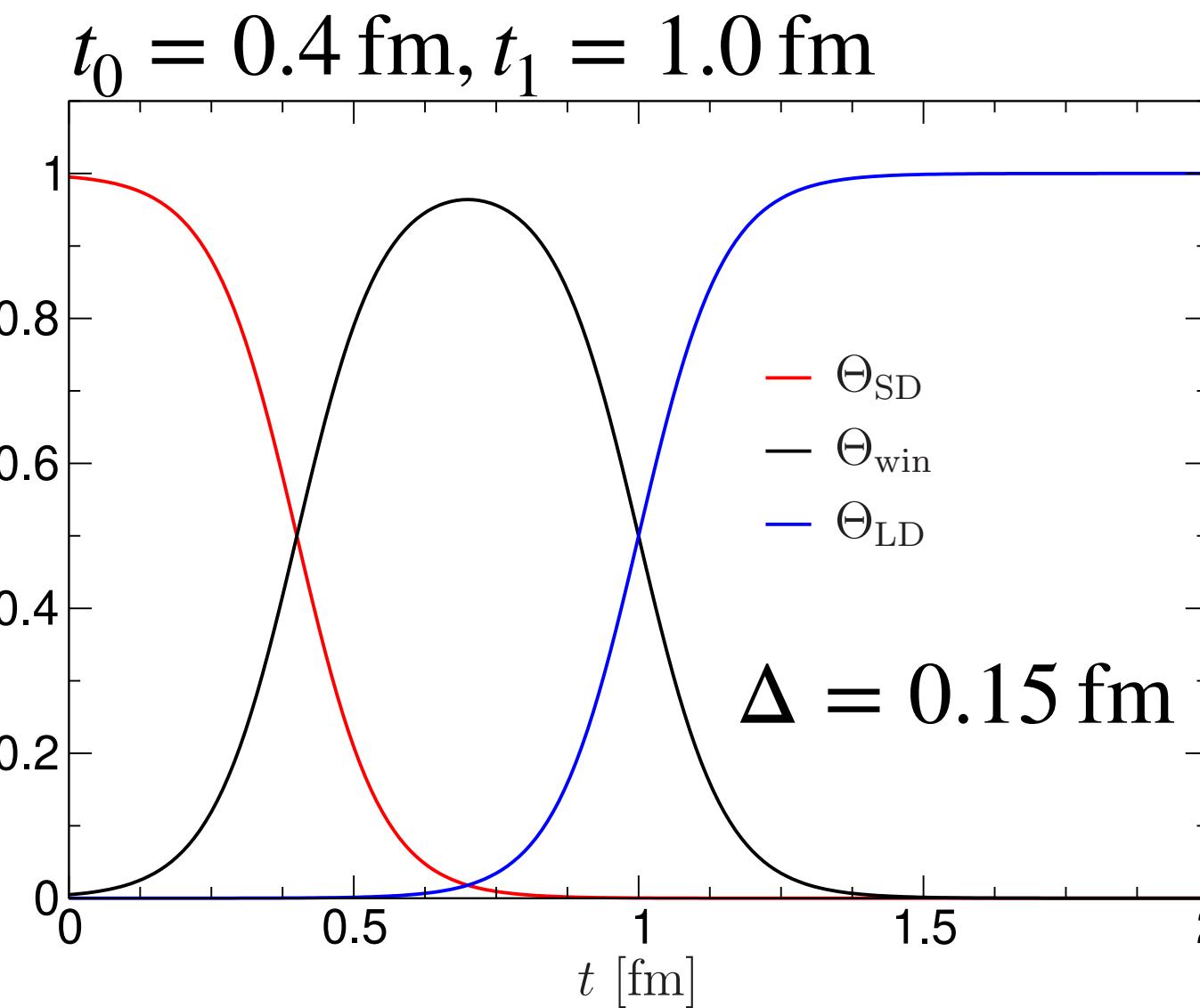
Modifying  $a_\mu^{\pi\pi}|_{\leq 1 \text{ GeV}}$

- “low-energy” scenario: local changes in cross section of  $\sim 8\%$  **around  $\rho$**
- “high-energy” scenario: impact on **pion charge radius** and space-like VFF  $\Rightarrow$  chance for **independent lattice-QCD checks**
- requires **factor  $\sim 3$  improvement** over  $\chi$ QCD result:  
 $\langle r_\pi^2 \rangle = 0.433(9)(13) \text{ fm}^2$   
 $\rightarrow$  arXiv:2006.05431 [hep-ph]



# Windows: Euclidean time vs $\sqrt{s}$

Martin Hoferichter @ Lattice HVP workshop



[ $t_0, t_1]$ intermediate window	percentage captured of $\pi\pi$ channel $\leq 1 \text{ GeV}$		
	SD	intermediate	LD
[0.4, 1.0] fm	3	28	69
[1.0, 2.0] fm	31	51	18
[1.0, 2.5] fm	31	61	9
[1.0, 3.0] fm	31	65	4

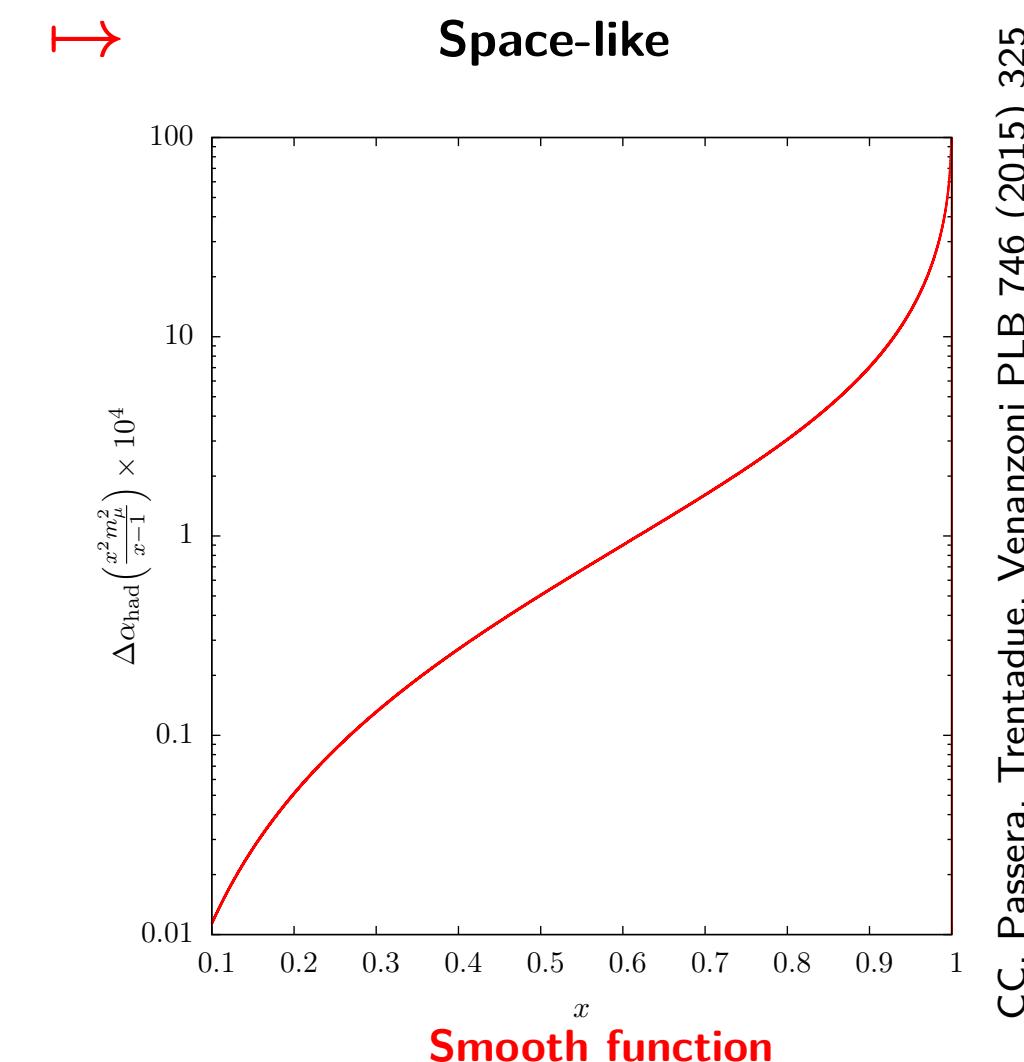
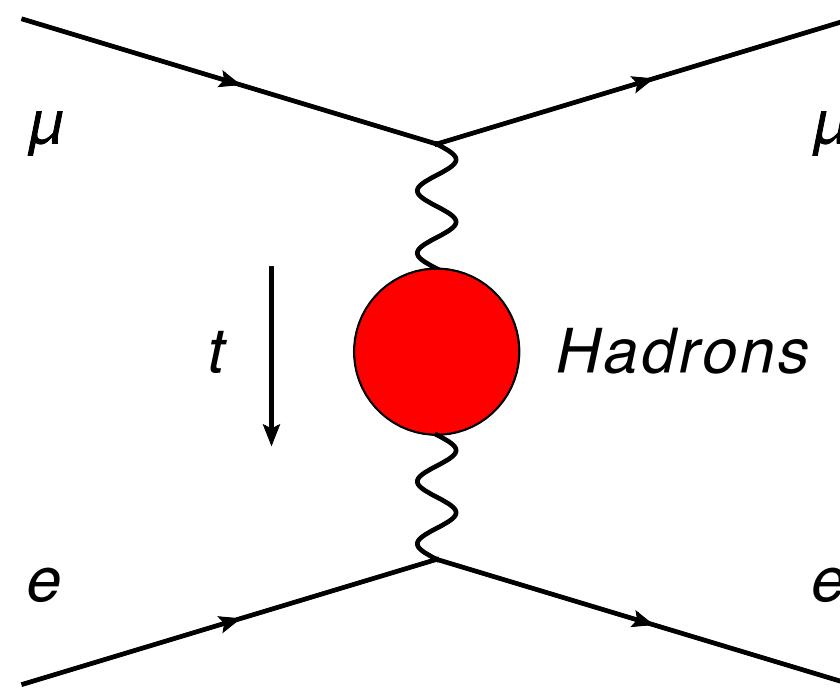
SD:  $[0, t_0]$   
LD:  $[t_1, \infty]$   
intermediate:  $[t_0, t_1]$

For intermediate window:  
 $\sim 30\%$  from  $\sigma(\pi\pi) \lesssim 1 \text{ GeV}$

# Hadronic vacuum polarization

$\mu$ -e elastic scattering to measure  $a_\mu^{\text{HVP}}$

LOI June 2019 [P. Banerjee et al, [arXiv:2004.13663](https://arxiv.org/abs/2004.13663), Eur.Phys.J.C 80 (2020)]



- use CERN M2 muon beam (150 GeV)
- Physics beyond colliders program @ CERN
- LOI June 2019
- pilot run in 2021
- full apparatus in 2023-2024

