

A novel approach to lattice QCD at finite baryon density

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in Collaboration with:

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partially based on JHEP 05 (2020) 088; Giordano, Kapas, Katz, Negradi, Pásztor

Approaches to finite density lattice QCD

In addition to the sign problem, known approaches to finite density QCD suffer from additional serious problems. E.g.

- Taylor and imaginary μ : **analytic continuation problem**
- Reweighting: **overlap problem**
- Complex Langevin: **convergence issues**
- ...

This talk:

→ a method where the only problem is the sign problem

If the sign problem is dealt with by sufficient statistics, the results are reliable, and errors (on a fixed lattice) are statistical only.

Reweighting in general

Target theory: Z_w Simulated theory: Z_r

$$Z_w = \int \mathcal{D}U w(U) \quad w(U) = \det M[U, \mu] e^{-S_g[U]} \in \mathbb{C}$$

$$Z_r = \int \mathcal{D}U r(U) \quad r(U) > 0$$

$$\langle O \rangle_w = \frac{\int \mathcal{D}U w(U) O(U)}{\int \mathcal{D}U w(U)} = \frac{\int \mathcal{D}U r(U) \frac{w(U)}{r(U)} O(U)}{\int \mathcal{D}U r(U) \frac{w(U)}{r(U)}} = \frac{\langle \frac{w}{r} O \rangle_r}{\langle \frac{w}{r} \rangle_r}$$

Two problems that are exponentially hard in the volume:

- $\frac{w}{r} \in \mathbb{C} \rightarrow$ **sign problem**
- Tails of $\rho(\frac{w}{r})$ long \rightarrow **overlap problem** \leftarrow The first bottleneck when reweighting from $\mu = 0$ (see talk by Kornél Kapás)

Sign reweighting

$$Z = \int \mathcal{D}U e^{-S_g} \det M = \int \mathcal{D}U e^{-S_g} \text{Re det } M$$

- Beware: the substitution $\det M \rightarrow \text{Re det } M$ can be done in Z but not in generic expectation values.
- Can calculate e.g. $\frac{\partial^n \log Z}{\partial \mu_{ud}^n}$, $\frac{\partial^n \log Z}{\partial m_{ud}^n}$ and $\frac{\partial^n \log Z}{\partial \beta^n}$

A new choice of a theory to reweight to and from:

$$\begin{aligned} w &= e^{-S_g} \text{Re det } M \\ r &= e^{-S_g} |\text{Re det } M| \end{aligned} \Rightarrow \frac{w}{r} \equiv \epsilon = \pm 1$$

- The weights are $\epsilon = \pm 1 \rightarrow$ No tail, **no overlap problem**
- $\langle \pm \rangle_r$ measures the strength of the **sign problem**

Early mentions of the idea:

de Forcrand, Kim, Takaishi: hep-lat/0209126; Nucl.Ph.B Proc.S. 119 (2003)

Alexandru, Faber, Horvath, Liu: hep-lat/0507020; PRD72 114513

Simulation algorithm

- We modify a standard RHMC at $\mu = 0$ by including an extra term in the accept-reject step.
- At the end of the trajectory $\Delta H_{\text{total}} = \Delta H_{\mu} + \Delta H_0$ is calculated, where

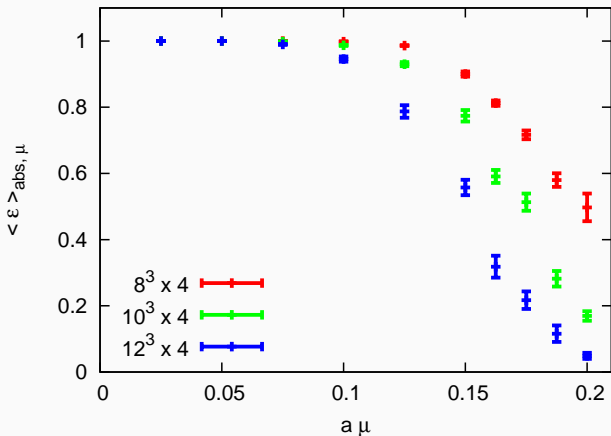
$$\Delta H_{\mu} = \Delta \left(\log \left| \text{Re det}^{1/2} M_{ud}(\mu_{ud}) \right| - \log \text{det}^{1/2} M_{ud}(0) \right)$$

calculated with the reduced matrix [Hasenfratz, Toussaint; '92].

- Frequent tunnelings between the \pm sectors
- Acceptance $> 50\%$

First test of the new method - unimproved staggered $N_\tau = 4$

Strength of the sign problem at $T_c(\mu)$

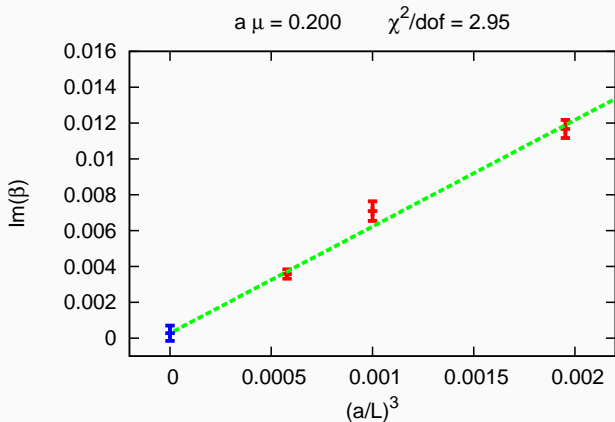


Giordano, Kapas, Katz, Nogradi, Pasztor; JHEP 05 (2020) 088

For simplicity we take $\mu_s = 0$ and $\mu_u = \mu_d = \mu_B/3$

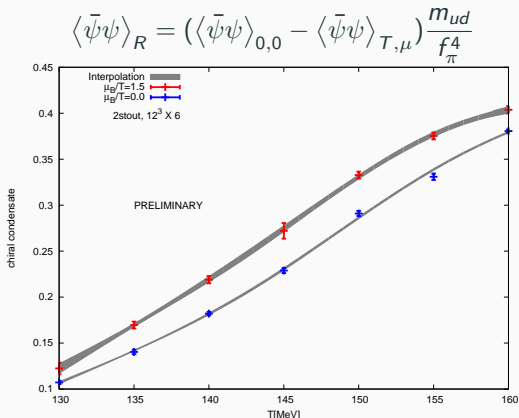
First test of the new method - unimproved staggered $N_\tau = 4$

Finite volume scaling at $\mu_B/T = 2.4$



Consistent with Fodor, Katz; JHEP 04 (2004) 050. BUT: to start being relevant for phenomenology, a much better lattice action has to be used

Second test - 2stout $N_\tau = 6$ (PRELIMINARY)



At $T = 130\text{MeV}$ and $\mu_B/T = 1.5$ the strength of the sign problem is:

$$12^3 \times 6 \rightarrow \langle \epsilon \rangle = 0.95 \pm 0.01 \quad 16^3 \times 6 \rightarrow \langle \epsilon \rangle = 0.70 \pm 0.03$$

\rightarrow can go further both in μ_B and in lattice size

Summary

- Current methods to study finite density QCD are typically not bottlenecked by the sign problem itself
- In particular reweighting from $\mu = 0$ is bottlenecked by the overlap problem
- We proposed a new reweighting method that is free from the overlap problem in the weights and is therefore only bottlenecked by the sign problem itself
- First test: CEP for unimproved staggered at $N_\tau = 4$, expected to be dominated by cut-off effects
- Second test: 2stout at $N_\tau = 6$; preliminary; width of transition at $\mu_B/T = 1.5 \approx$ width at 0