





Transverse momentum broadening in real-time lattice simulations of the glasma

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Transverse momentum broadening in real-time lattice simulations of the glasma for LATTICE 2021

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based on:

[1] A. Ipp D. I. Müller and D. Schuh, Phys. Rev. D **102** (2020) 074001, [hep-ph/2001.10001]

[2] A. Ipp D. I. Müller and D. Schuh, Phys. Let. B **810** (2020) 135810, [hep-ph/2009.14206]

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Introduction

high energy nuclei collisions \rightarrow jets hit detector jet quenching, elliptical shape of the jet \rightarrow momentum broadening seeds of jets: highly energetic partons (created by hard scatterings during collision)

ightarrow all stages of the medium that is created by the collision may contribute to momentum broadening

consider first stage: glasma (pre-equilibrium precursor state of the quark-gluon plasma)

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high energy nuclei collisions \rightarrow Color Glass Condensate framework classical effective theory for high energy QCD; main idea: separation of scales:

- hard partons: high momentum (valence quarks, high energy gluons)
 - ightarrow described by thin sheets of classical color charge (Lorentz contracted)
 - \rightarrow color configuration frozen with respect to QCD time scales (time dilation); specified by MV model
- soft partons: low momentum (low energy gluons)
 - ightarrow described by a highly occupied color field

effective degrees of freedom: color currents and classical color fields equations of motion: classical Yang-Mills equations

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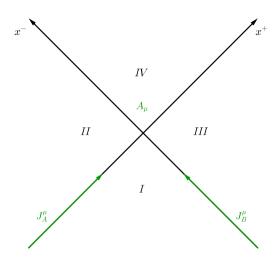
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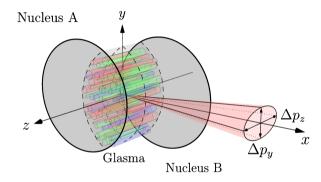
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consider limit of infinitely thin nuclei ightarrow 2+1D system, independent of rapidity

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color fields exert Lorentz forces on the parton no deflection in this limit, but the parton accumulates momentum

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components of jet broadening parameter \hat{q} :

$$\hat{q}_i(\tau) = \frac{\mathrm{d}}{\mathrm{d}\tau} \langle p_i^2(\tau) \rangle_q$$

Wong equations:

$$egin{aligned} rac{\mathrm{d} oldsymbol{p}_{\mu}}{\mathrm{d} au} &= g Q^a(au) rac{\mathrm{d} x^{
u}}{\mathrm{d} au} F^a_{\mu
u}(au) \ rac{\mathrm{d} Q^a}{\mathrm{d} au} &= g rac{\mathrm{d} x^{\mu}}{\mathrm{d} au} f^{abc} A^b_{\mu}(au) Q^c(au) \end{aligned}$$

solution for a quark:

$$\langle p_i^2(au)
angle_q = rac{g^2}{N_c} \int\limits_0^ au \mathrm{d} au' \int\limits_0^ au \mathrm{d} au'' \langle \mathrm{Tr}\left[f^i(au')f^i(au'')
ight]
angle$$

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$$\langle \rho_i^2(\tau) \rangle_q = \frac{g^2}{N_c} \int_0^{\tau} d\tau' \int_0^{\tau} d\tau'' \langle \text{Tr} \left[f^i(\tau') f^i(\tau'') \right] \rangle$$

 f^i is a function representing the color rotated Lorentz force:

$$f^{y}(\tau) = U(\tau) (E_{y}(\tau) - B_{z}(\tau)) U^{\dagger}(\tau)$$

$$f^{z}(\tau) = U(\tau) (E_{z}(\tau) + B_{y}(\tau)) U^{\dagger}(\tau)$$

lightlike Wilson line in the fundamental representation along particle trajectory:

$$U(au,0) = \mathcal{P} \exp \left(-ig\int\limits_0^ au d au' A_x(au')
ight)$$

result for gluon can be found by Casimir scaling:

$$\langle p_i^2 \rangle_g = \frac{C_A}{C_F} \langle p_i^2 \rangle_q$$

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finite difference method to solve Yang-Mills equations numerically transverse plane $\to N_T \times N_T$ lattice with periodic boundary conditions degrees of freedom:

- $U_{x,\hat{i}}(\tau_n)$ gauge links
- $A_{x,n}(\tau_n)$ vector field
- $P_x^i(\tau_{n+\frac{1}{2}})$ conjugate momenta
- $P_{x}^{\eta}(\tau_{n+\frac{1}{2}})$ conjugate momentum

Yang-Mills equations \rightarrow leapfrog scheme

initial conditions \rightarrow discretized

lattice expressions for E_i and $B_i \rightarrow f^i$, replace integrals by sums:

$$\langle p_{(y,z)}^2(t_n) \rangle_q pprox rac{g^2 a_T^2}{N_c} \langle \operatorname{Tr} \left[\left(\sum_{i=0}^n f^{(y,z)}(t_n) \right)^2
ight]
angle$$

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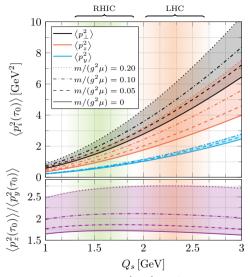


Figure: Accumulated transverse momentum (top) and momentum broadening anisotropy (bottom) at $\tau_0=0.6\,{\rm fm/c}$

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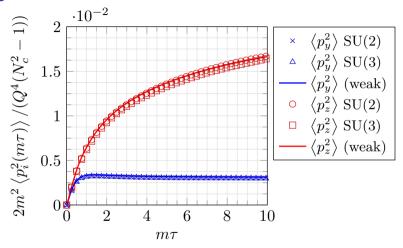


Figure: Accumulated transverse momentum in the dilute glasma: weak field approximation and lattice approximation

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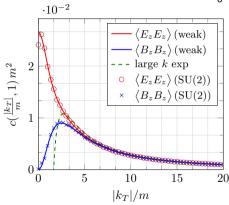
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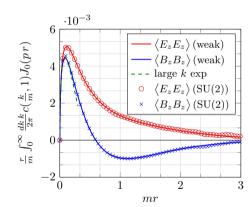
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$$\langle p_{(y,z)}^2(\tau)\rangle = \int\limits_0^\infty dk \, g(\tau,k) \, c_{(B,E)}(k)$$





(a) Initial correlators in momentum space

(b) Initial correlators in coordinate space multiplied by the dimensionless distance *mr*

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this talk:

- momentum broadening in the very early stages of heavy ion collisions
- amount, anisotropy and its origin
- weak field approximation and lattice approximation

future work:

- relax approximations (nuclei moving at the speed of light, test parton moving at the speed of light, no backreaction of the parton in the glasma)
- energy loss
- improve understanding of physics behind the form of the initial correlators

open source code: gitlab.com/openpixi/curraun

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