

# Perturbative predictions for color superconductivity on the lattice

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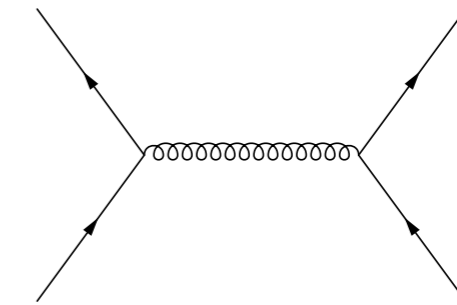
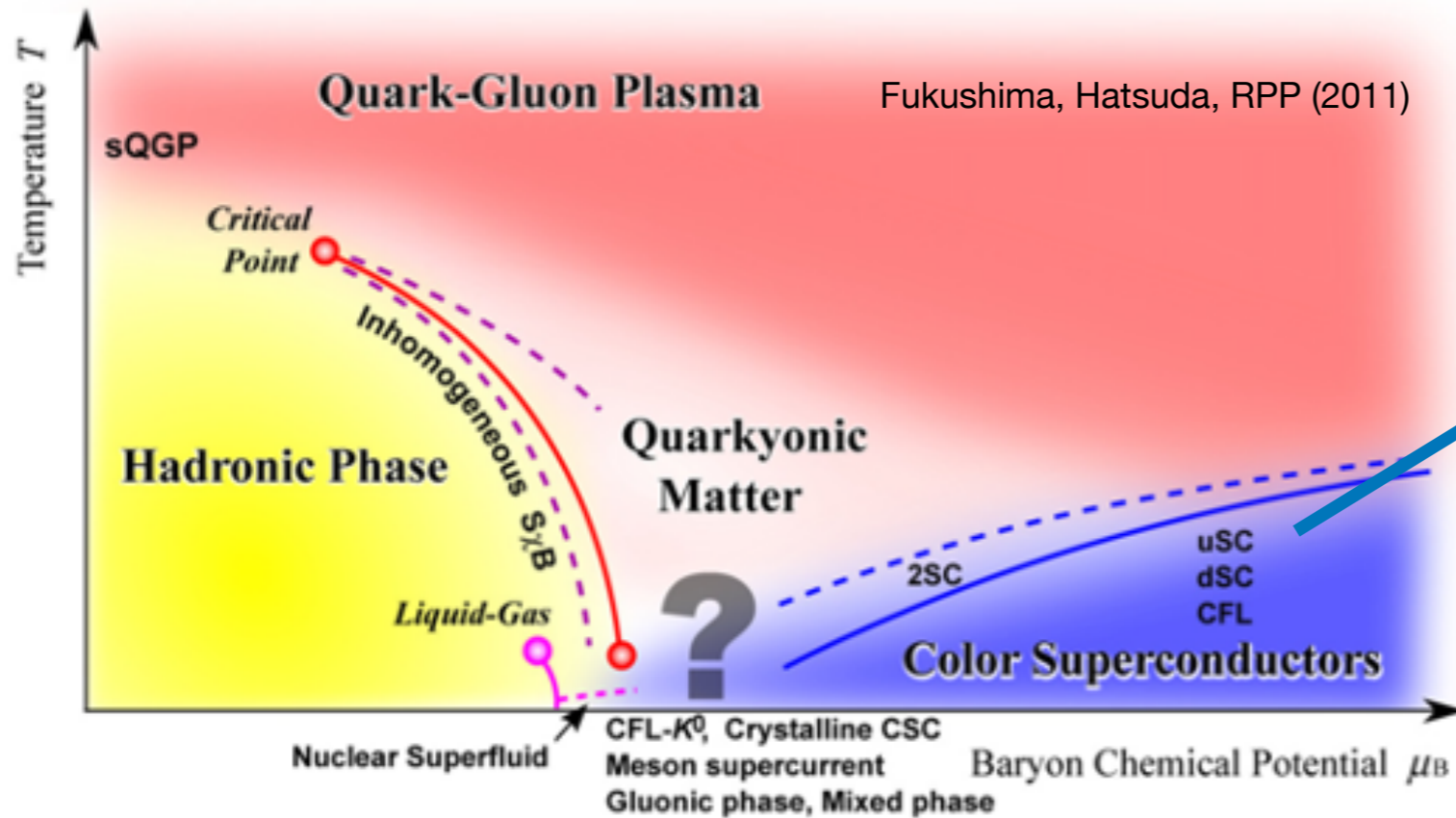
In collaboration with

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Y. Namekawa (YITP), J. Nishimura (KEK), A. Tsuchiya (Shizuoka), S. Tsutsui (RIKEN)

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# Color superconductivity (CSC)

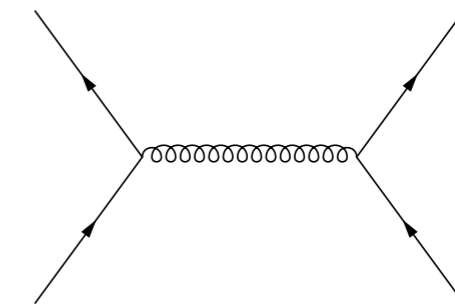
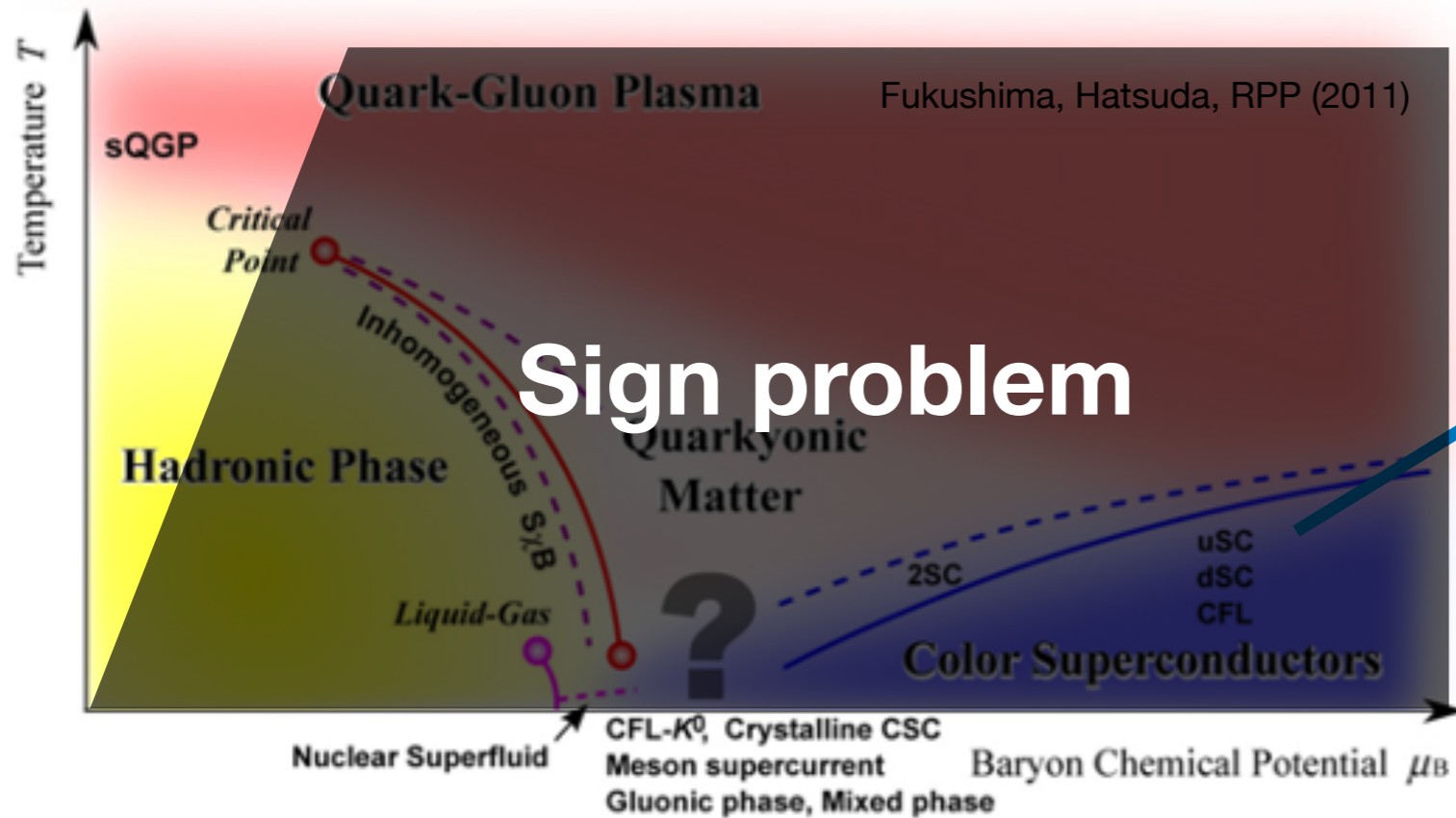
Barrois, NPB (1977); PhD thesis (1979), Frautschi (1978)  
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**Cooper instability in color-antitriplet channel**  
 (attractive channel)

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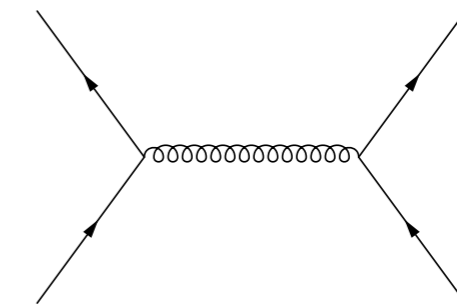
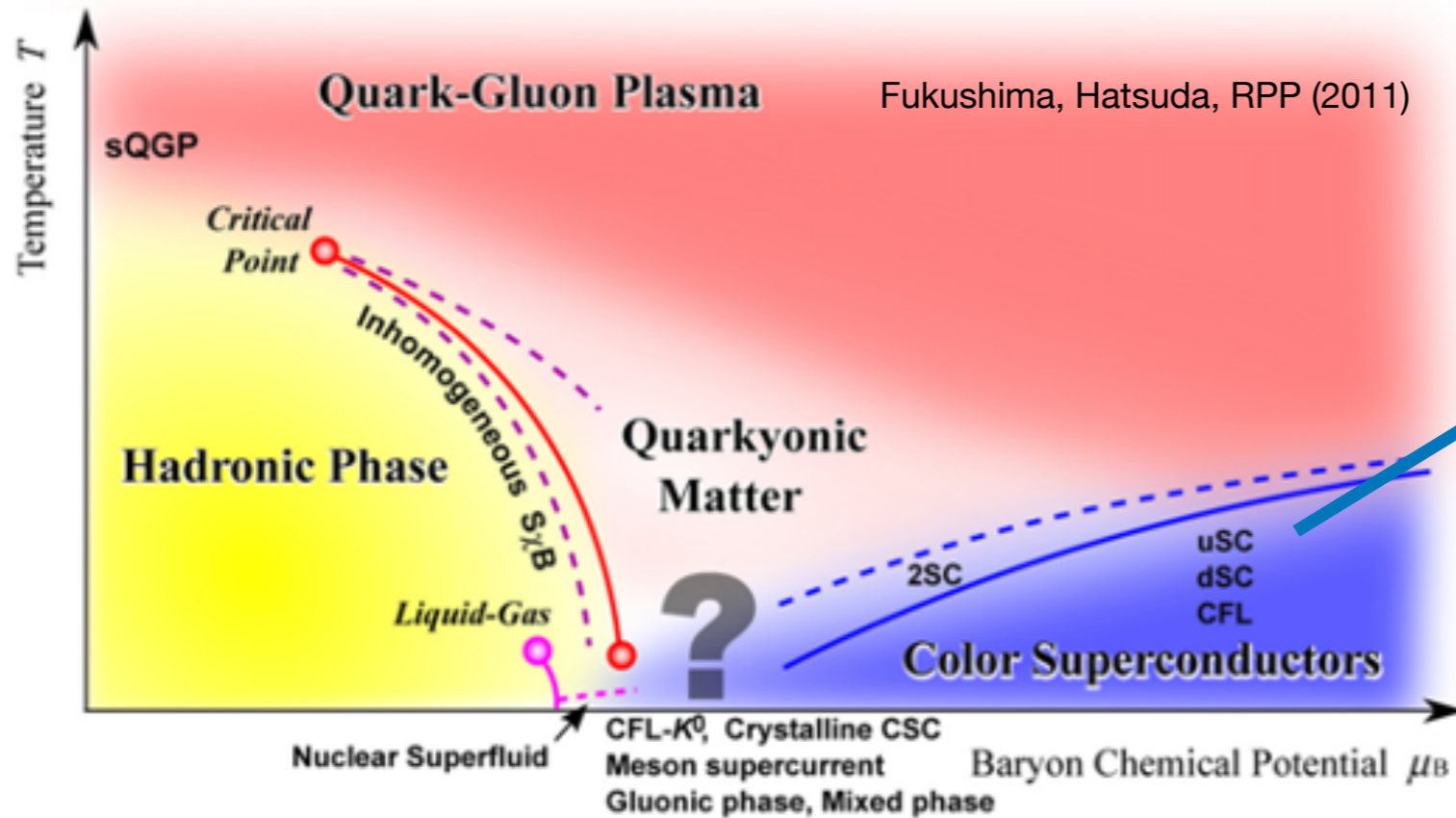
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# Color superconductivity (CSC)

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**Cooper instability in color-antitriplet channel**  
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## Exploration of CSC by first-principles lattice-QCD methods?

Candidates of methods to overcome sign problem:

- **Complex Langevin method**
- Lefschetz thimble
- Tensor renormalization group
- Path optimization

...



**Tsutsui's talk**  
 (application to CSC)

# Perturbative study of CSC on the lattice

## Perturbative study

- **Valid for small coupling  $g$  (high density, small volume, ...)**
- Useful for **cross check** of results by first-principles lattice-QCD methods

Study in infinite continuum case  
based on perturbative QCD

(Review) Alford, Schmitt, Rajagopal, Schäfer RMP (2008)

Gap eq.

$$\Delta_{k,r} = \frac{g^2}{4} \int \frac{d^3q}{(2\pi)^3} \sum_s Z(\epsilon_{q,s}) \frac{\Delta_{q,s}}{\epsilon_{q,s}} \tanh\left(\frac{\epsilon_{q,s}}{2T}\right) \\ \times \left[ D_\ell(p) \mathcal{T}_{00}^s(\mathbf{k}, \mathbf{q}) + D_t(p, \epsilon_{q,s}, \epsilon_{k,r}) \mathcal{T}_t^s(\mathbf{k}, \mathbf{q}) \right]$$

- Functional equation for gap  $\Delta_{q,s}$   
**Some ansatz on  $\Delta_{q,s}$  is necessary.**
- Gauge-symmetry-breaking cutoff

## Analysis on the finite lattice

- **Gap eq. is reduced to coupled equations for finite number of elements**
  - Suitable for numerical study
- **Lattice regularization does not break gauge symmetry**

# Our main results

## A new method to analyze CSC quantitatively based on the lattice QCD perturbation

- Application of Nambu formalism to the lattice system
- Analysis around critical point
  - **First application of (generalized) Thouless criterion to QCD**
  - **Numerical method without any ansatz on the form of Cooper pairs**

## Quantitative prediction of parameter region with which CSC occurs

- $N_f = 4$  staggered fermions (with corrections of results reported in APLAT 2020)
- Wilson fermions ( $N_f \geq 2$ )
- Critical coupling constant

## Prediction for the structure of Cooper pairs

- Symmetry in flavor space
- Momentum modes composing Cooper pairs

# Method

Dyson eq.  
in Nambu basis

$$\Sigma = \mathbf{S}^{-1} - \mathbf{S}_{\text{free}}^{-1}$$

Emerged by  $\Sigma_{12(21)}$

$$\mathbf{S} = \begin{pmatrix} \langle \psi \bar{\psi} \rangle & \langle \psi \psi \rangle \\ \langle \bar{\psi} \bar{\psi} \rangle & \langle \bar{\psi} \psi \rangle \end{pmatrix}$$

# Method

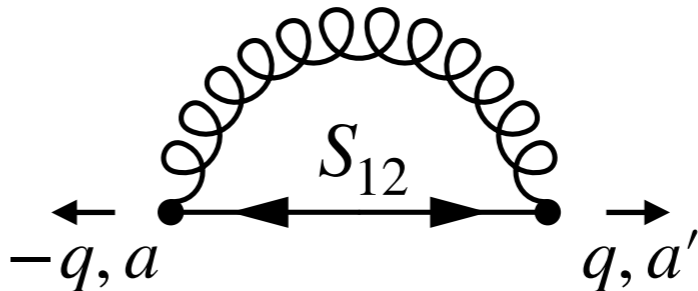
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Gap eq.  
in lowest order

$$\Sigma_{12}(qaa') =$$


$q$ : momentum  
 $a, a'$ : internal d.o.f.

- Complicated non-linear eq. for  $\Sigma_{12} \dots$



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$$\Sigma_{12}(qaa') = \text{Diagram}$$

$q$ : momentum  
 $a, a'$ : internal d.o.f.

- Complicated non-linear eq. for  $\Sigma_{12} \dots$
- **However, it is reduced to linear eq. near critical point, where  $\Sigma_{12}$  becomes small but nonzero value.**

Critical point  
(Thouless criterion)

$$\frac{1}{V} \sum_{q'bb'} M_{(qaa')(q'bb')} \Sigma_{12}(q'bb') = \beta \Sigma_{12}(qaa')$$

$$\beta = 2N_c/g^2$$

$$\frac{1}{\beta} M_{(qaa')(q'bb')} = \text{Diagram}$$

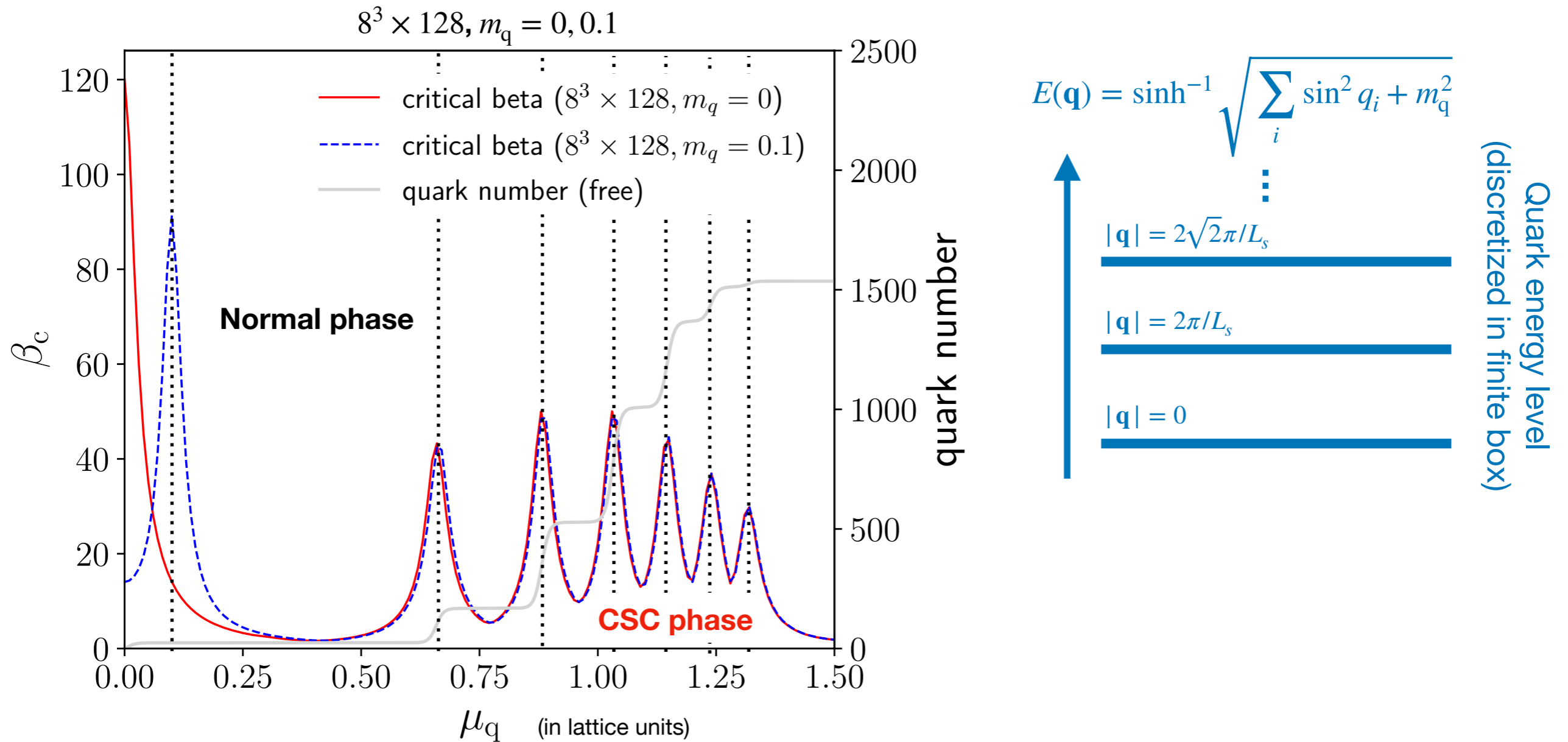
Max. eigenvalue of  $M$  = Critical  $\beta$  for CSC

Eigenvector = Shape of gap  $\Sigma_{12}$

Analysis near critical point without any ansatz!

Numerical method: power iteration

# Critical $\beta$ with $N_f = 4$ staggered fermion

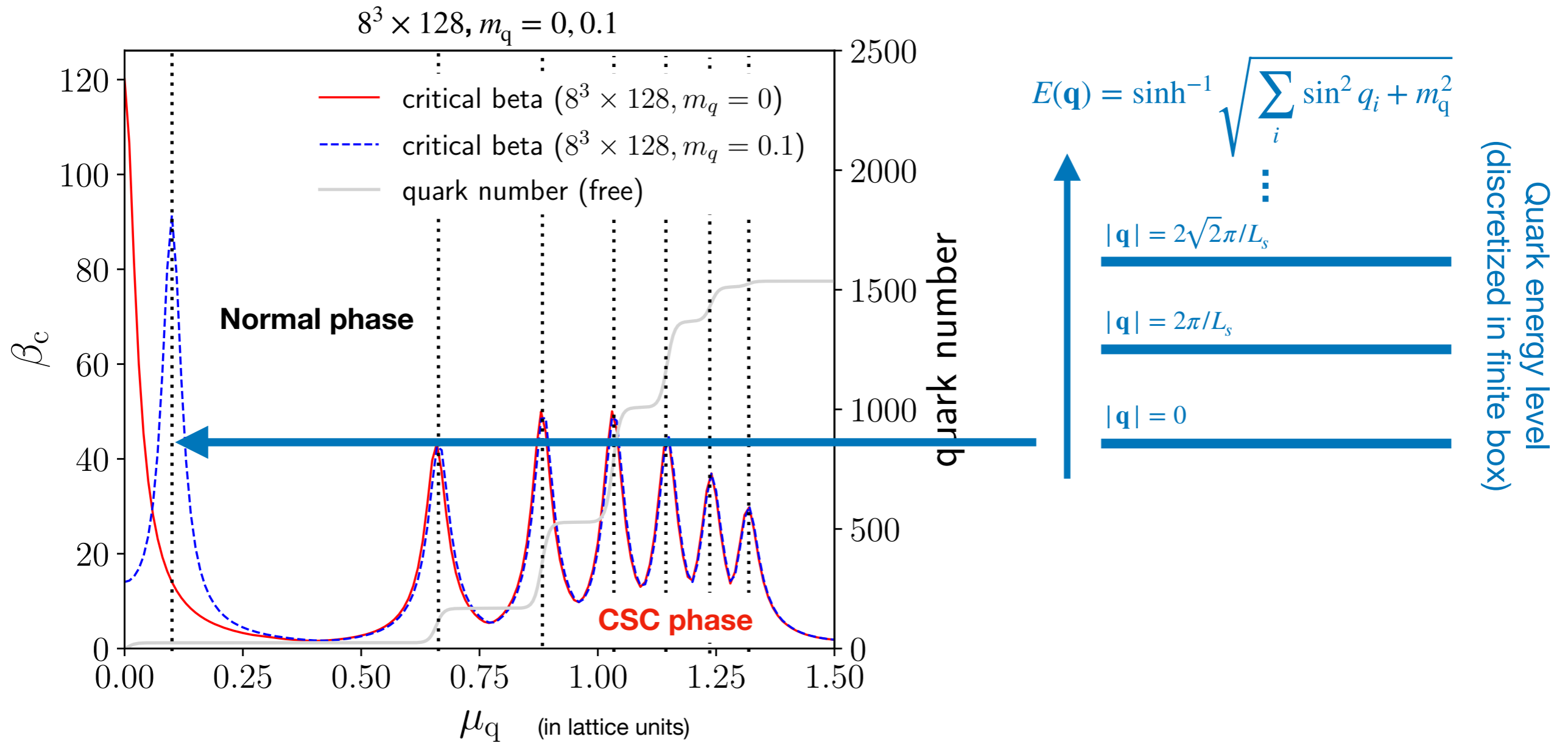


- **CSC phase extends to larger  $\beta$  when  $\mu_q = E(\mathbf{q})$ .**

Cooper pairs are easy to form when modes exist around  $\mu_q$

- **Finite mass shifts the energy levels and thus the peak positions.**
- Peak near origin may be the finite-size effect.

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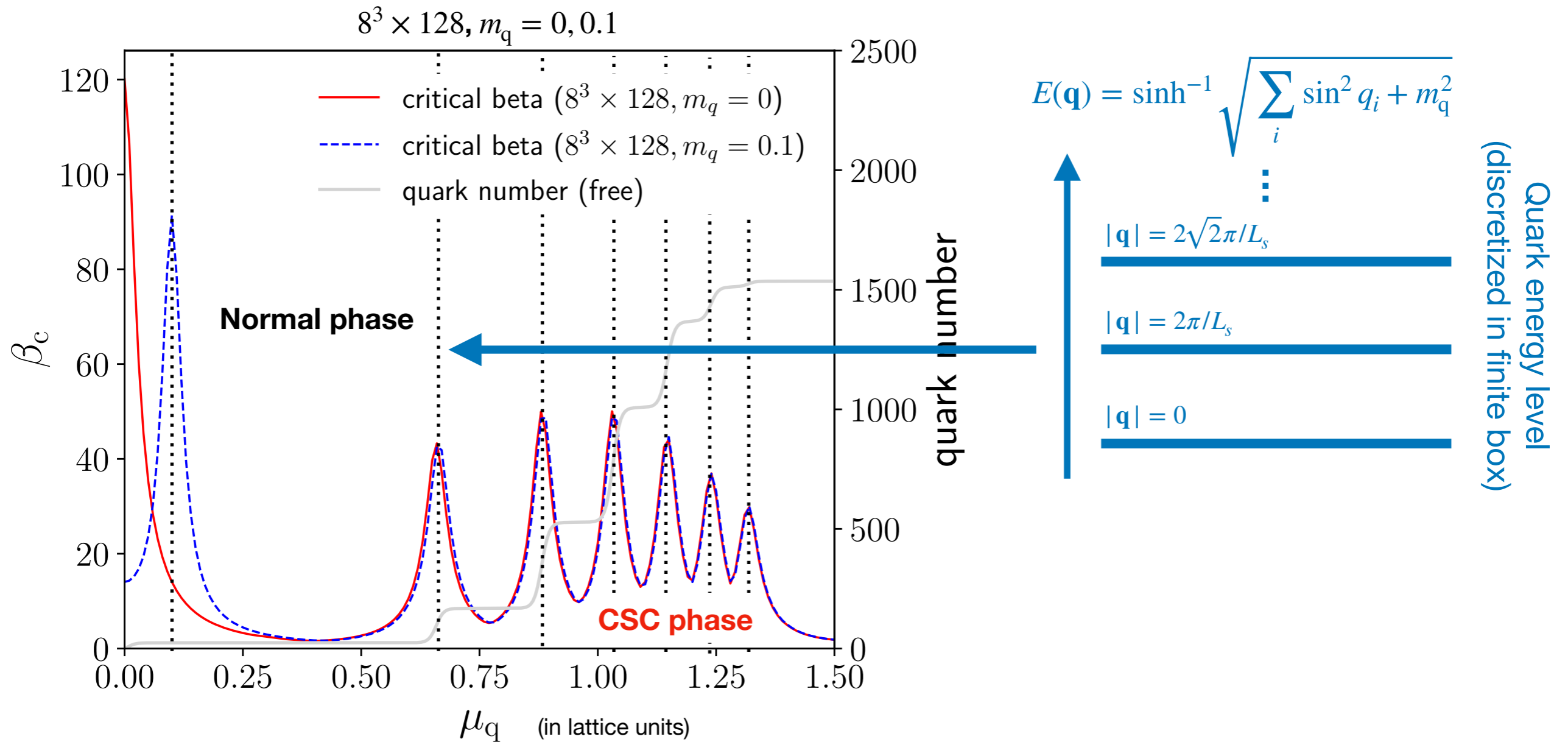


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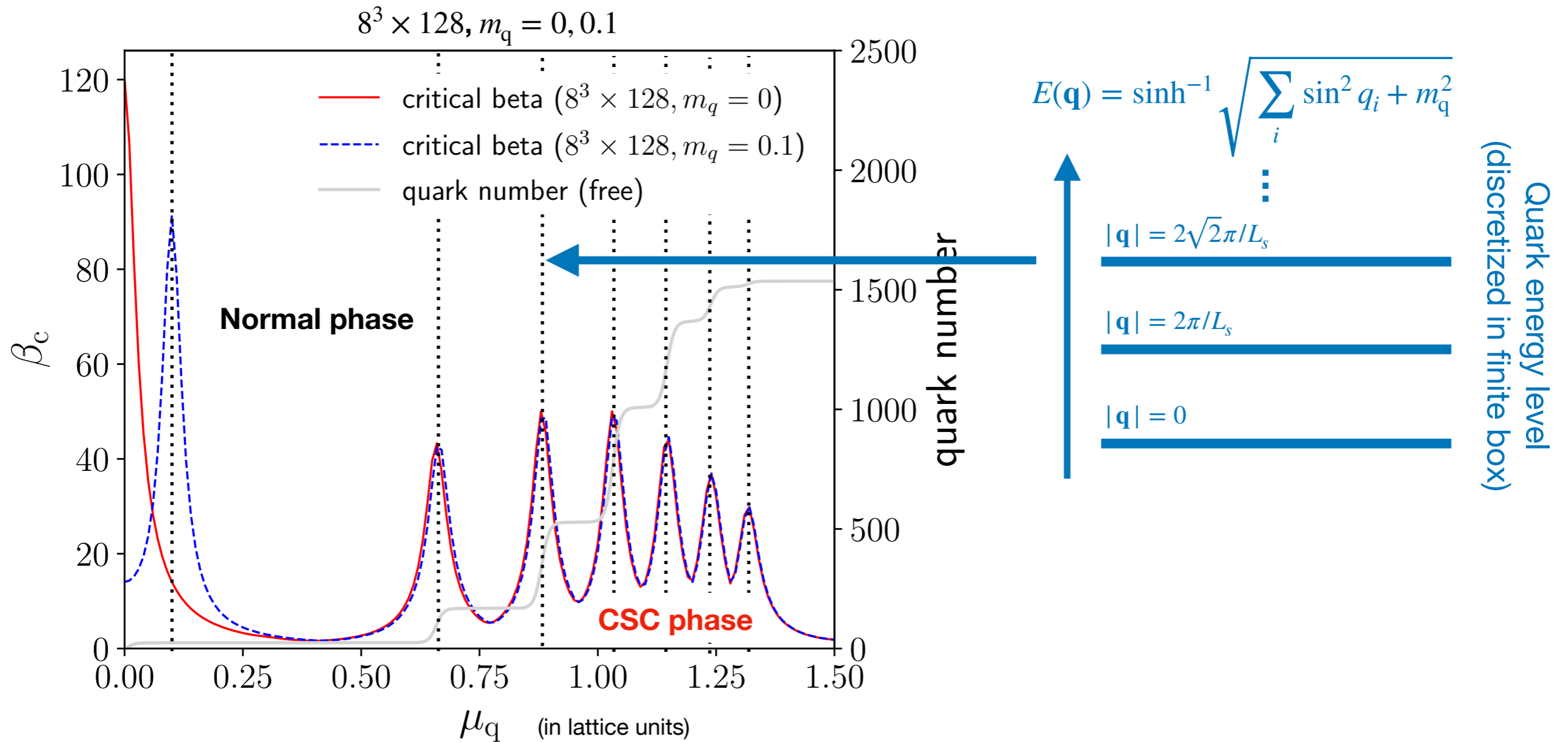


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# Flavor structure of Cooper pairs

## Color-antisymmetric Lorentz-scalar condensate

obtained from  $\Sigma_{12}$  (up to overall factor)

$$S_{12,a_1 a_2}^{(-)}(N) = \epsilon_{\alpha\beta 3} \left\langle \psi_{a_1}^{\alpha t}(N) C \psi_{a_2}^{\beta}(0) \right\rangle$$

color indices      charge-conjugation operator      flavor indices

# Flavor structure of Cooper pairs

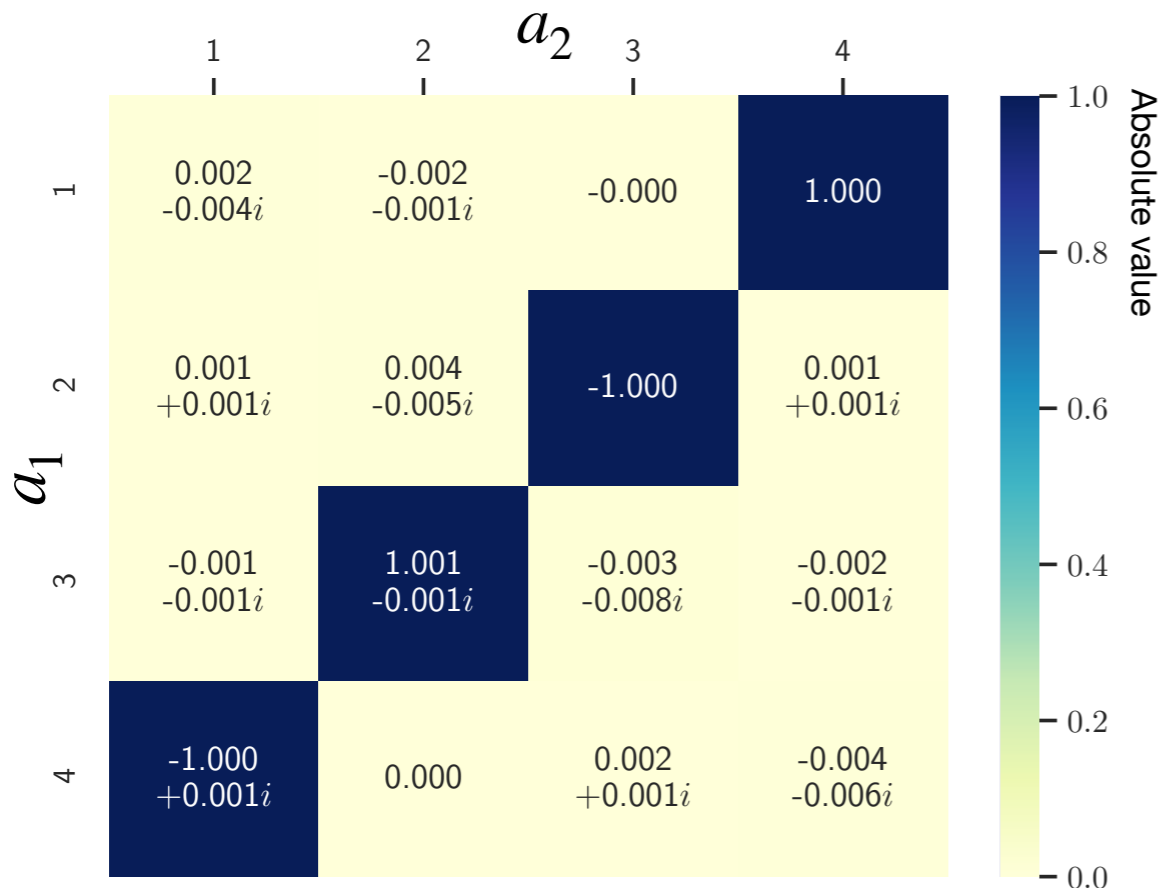
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Flavor dependence of  $S_{12,a_1a_2}^{(-)}(0)$  (normalized by  $S_{12,14}^{(-)}(0)$ )



$8^3 \times 128$ ,  $m_q = 0$ ,  $\mu_q = 0.66$ , at the critical  $\beta$

Standard representation for Euclidian  $\gamma$  matrices is used:

$$\gamma_i = \begin{pmatrix} 0 & -i\sigma^i \\ i\sigma^i & 0 \end{pmatrix} \quad \gamma_4 = \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix}$$

## Flavor structure of the condensate:

$$S_{12,a_1a_2}^{(-)}(N) = \kappa(N) \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}_{a_1a_2} = \kappa(N)(t_2 t_4)_{a_1a_2}$$

( $t_\mu = \gamma_\mu^t$  acts on the flavor space)

## Consistent with the following symmetries:

- Antisymmetry for flavor indices
- $U(1)$  symmetry of the staggered fermions:

$$\psi(N) \rightarrow e^{i\theta\gamma_5 \otimes t_5} \psi(N)$$

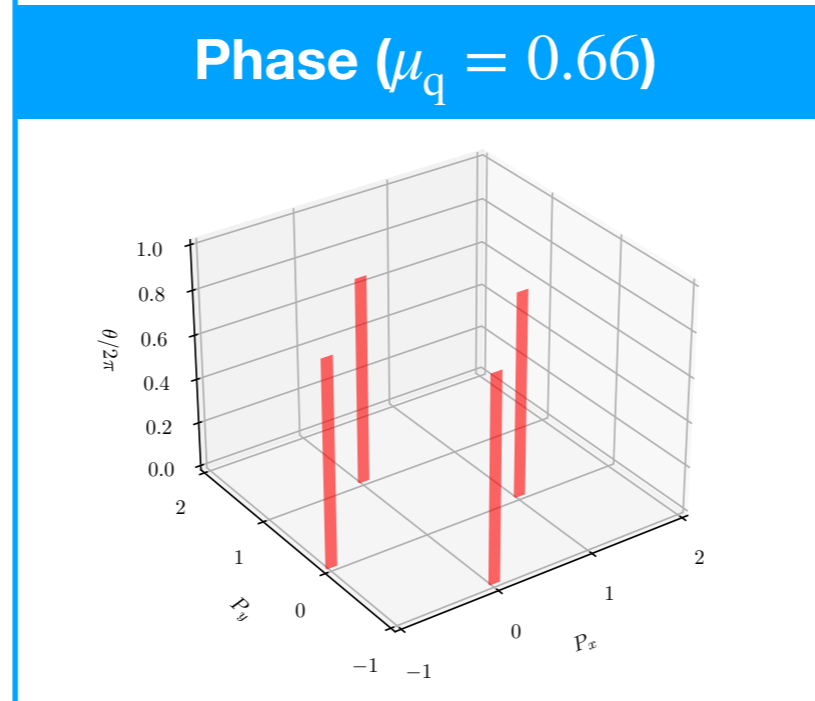
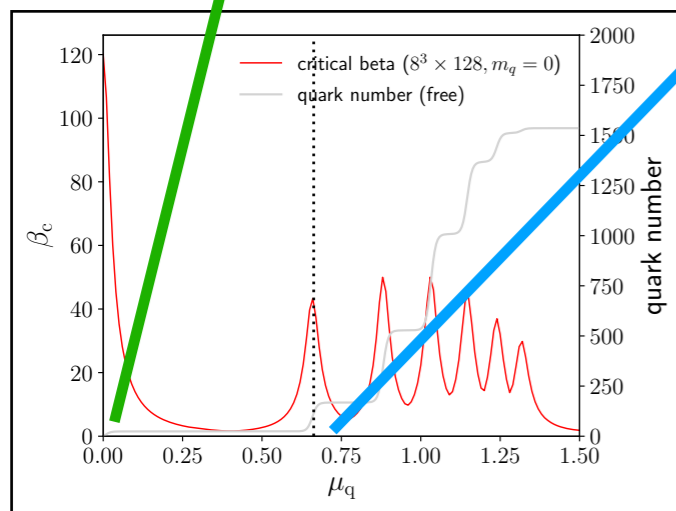
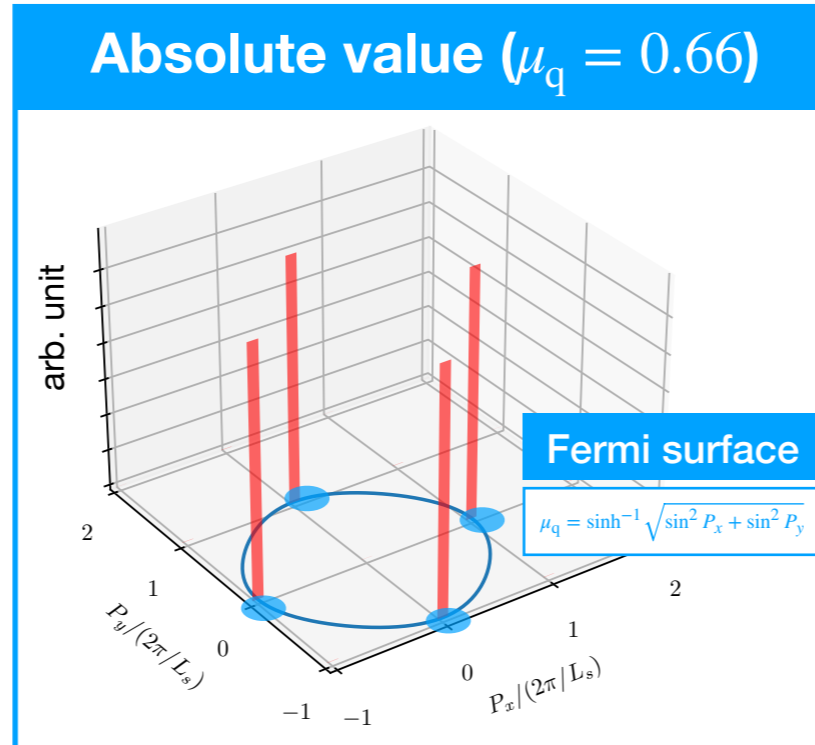
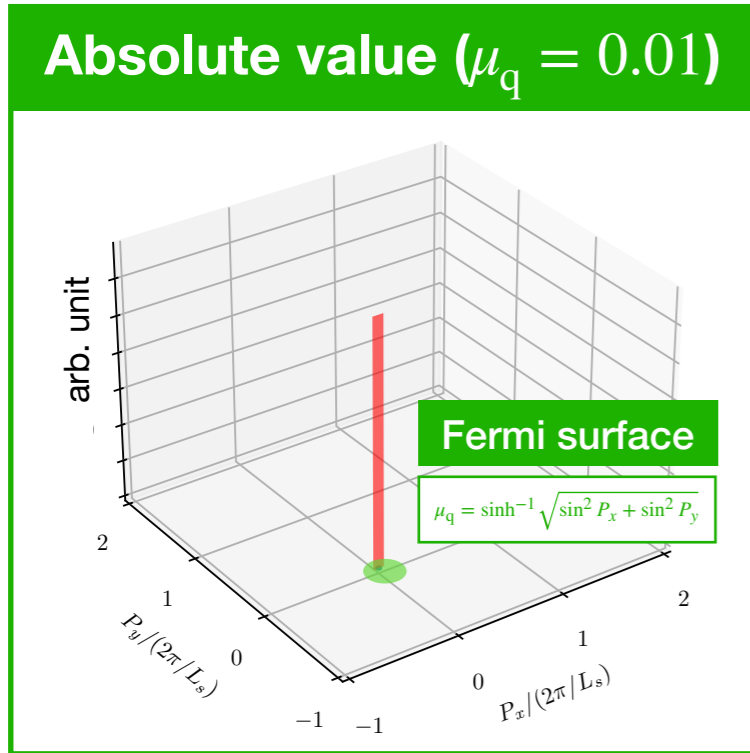
- Result does not change with different initial cond. for power iteration  
→ No degeneracy.
- Small but nonzero values of other components may be caused by the numerical error of the Fourier transform.

# Cooper pairs in momentum space (Brillouin zone)

$P_x$  and  $P_y$ -dependence of  $\tilde{\kappa}(P_x, P_y, 0, \pi/L_t)$

$$\tilde{S}_{12,a_1a_2}^{(-)}(P) = \tilde{\kappa}(P)(t_2 t_4)_{a_1 a_2}$$

$8^3 \times 128 (= L_s^3 \times L_t), m_q = 0$



- Cooper pairs are actually formed by the quarks on the Fermi surface.

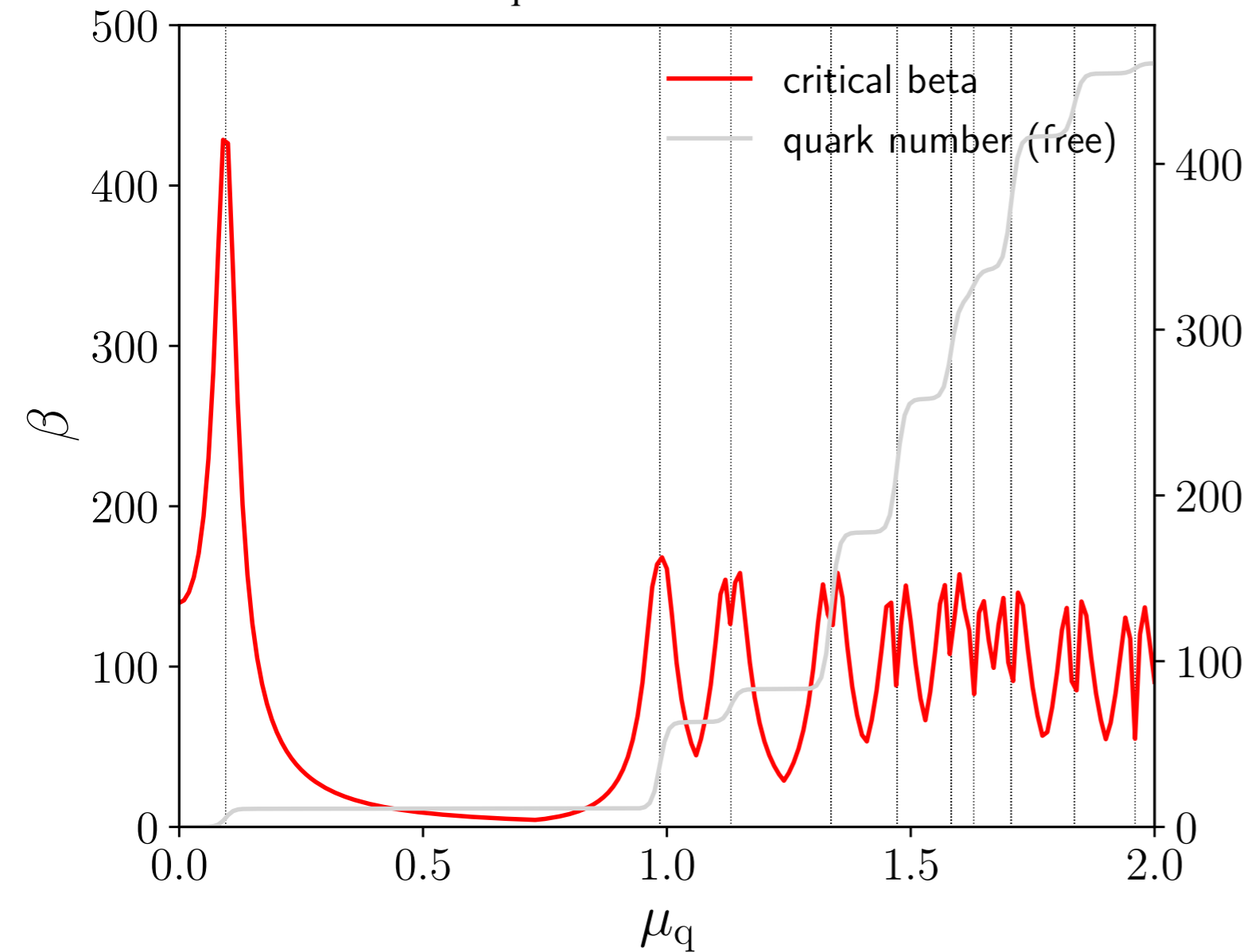
- All modes have the same phase.

These are obtained from calc. without any ansatz.



# Critical $\beta$ with Wilson fermion ( $N_f \geq 2$ )

$4^3 \times 128, m_q = 0.1, \kappa = 0.12195 (r = 1)$



Vertical dotted lines:  $\mu_q = E(\mathbf{q})$

Energy levels of free Wilson fermions

$$E(\mathbf{q}) = 2 \sinh^{-1} \sqrt{\frac{\sum_i \sin^2 q_i + \left(m_q + 2 \sum_i \sin^2 \frac{q_i}{2}\right)^2}{4 \left(1 + m_q + 2 \sum_i \sin^2 \frac{q_i}{2}\right)}}$$

- **Peaks corresponding to  $E(\mathbf{q})$  appear as in the case of staggered fermions.**
- **Splitting of peaks is observed at high  $\mu_q$ .**

Related to structure of Cooper pairs?  
(under investigation)

# Conclusion

## Study of color superconductivity (CSC) based on lattice QCD perturbation theory

A new method to analyze CSC quantitatively

- First application of (generalized) Thouless criterion to QCD
- Numerical method without any ansatz on the form of Cooper pairs

Quantitative prediction for critical lattice  $\beta$

- $N_f = 4$  staggered and Wilson fermions ( $N_f \geq 2$ )
- Peak structure in  $(\mu_q, \beta)$  plane (effect of discretized energy levels of quarks)

Prediction of the structure of Cooper pairs

- Symmetry in flavor space
- Momentum modes composing Cooper pairs

### Outlook

- Analysis of Cooper pair with Wilson fermions
- Continuum limit, scaling
- Complex Langevin analysis of CSC in the parameter region given by our perturbative analysis (talk by Tsutsui)