

2+1 flavor fine lattice simulation at finite temperature with domain wall fermions

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1: YITP, 2: R-CCS, 3: Osaka, 4: KEK

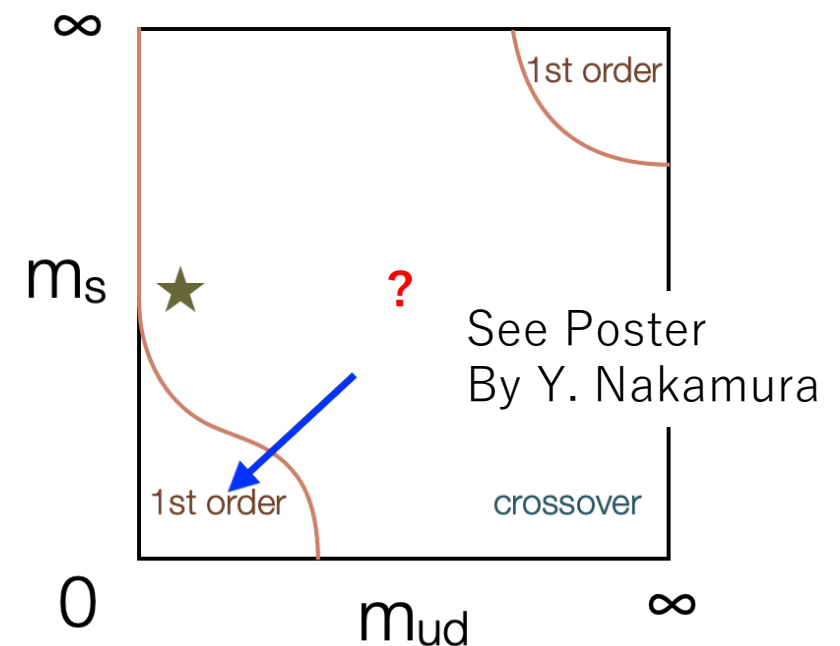
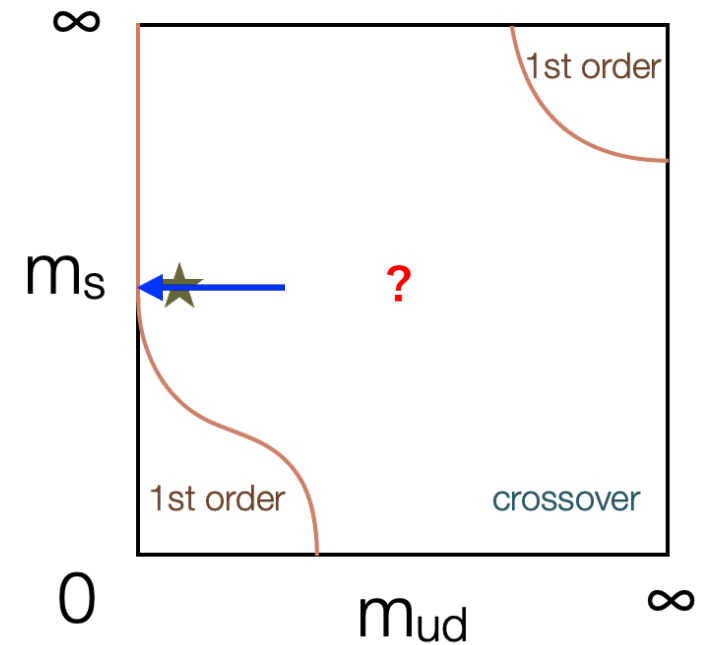
Lattice 2021
June 28, 2021

acknowledgements

- Codes used:
 - HMC
 - Grid / Regensburg → poster by N. Meyer
 - Measurements:
 - BQCD
 - Bridge++ (Wilson multigrid on Fugaku → talk by I.Kanamori (Wed))
- MEXT program
成果創出加速プログラム
Program for Promoting Researches on the Supercomputer Fugaku
 - Simulation for basic science: from fundamental laws of particles to creation of nuclei
- Computers
 - Oakforest-PACS
 - Polaire and Grand Chariot at Hokkaido University
 - supercomputer Fugaku provided by the RIKEN Center for Computational Science
- Fugaku: software / performance → plenary talk by Y. Nakamura (Fri)

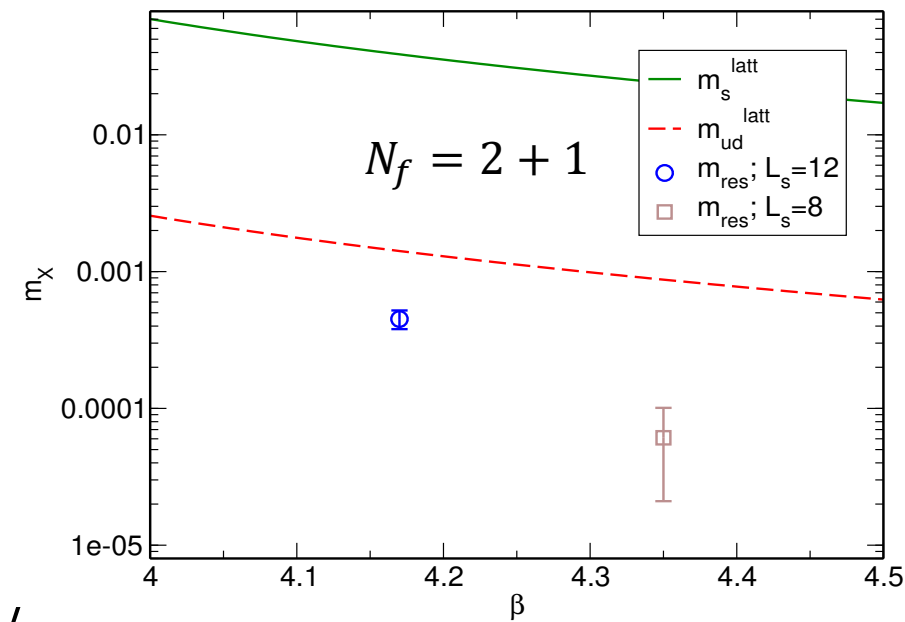
Intro

- $N_f=2+1$ thermodynamic property
 - through chiral symmetric formulation
 - Order of the transition
 - (pseudo) critical temperature
 - Location of the phase boundary
 - Near the physical point
- Chiral symmetric formulation
 - Ideal to treat flavor $SU(2)$ and $U(1)_A$ properly
 - Domain wall fermion (DWF) : practical choice
- DWF and chirality
 - Fine lattice needed
 - Aiming for $a < 0.08$ fm (eventually)
 - Current search domain: $0.08 < a < 0.12$ fm



$N_f=2$ Möbius DWF

- Lessons learned
 - Chiral symmetry important for discussing
 - chiral, $U(1)_A$ problems
 - Reweighting to overlap essential
 - For reweighting to be successful for DW - OV
 - Fine lattice needed (efficiency of reweighting): $a \approx 0.1$ fm
 - Smoothness of configuration & smallness of m_{res}
 - For reweighting to be successful in general
 - Large volume is problematic
 - It may not work for further finer lattices
- Expectation
 - Finer the lattice, smaller m_{res}
 - DWF itself eventually becomes good enough
 - Aiming fine lattice DWF simulation would help in any sense



$N_f=2$ Möbius DWF

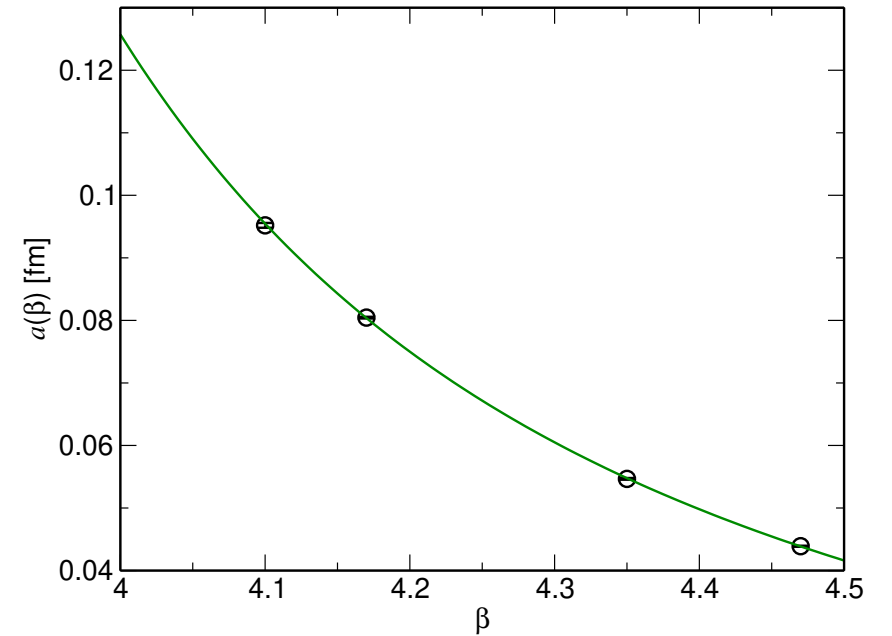
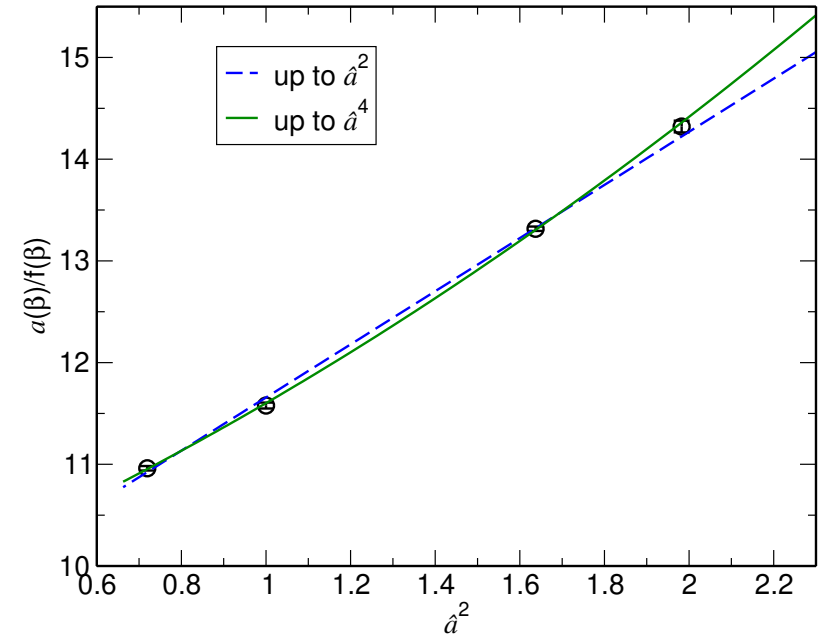
- Action
 - Tree-level improved Symanzik gauge
 - stout-smear, scale-factor 2 Shamir
- So far studied
 - $U(1)_A$ and chiral symmetry
 - Topological charge
 - Chiral susceptibility \rightarrow Fukaya (poster)
- Simulation setup
 - Fix β
 - Fix N_t
 - Vary m

$$N_f = 2 + 1$$

- Action: same as $N_f = 2$
- Simulation setup (we follow most of the simulations by now)
 - Fix β
 - Fix N_t
 - Fix m_s^{latt} near physical
 - Vary m_l^{latt}
 - Aiming to understand the role of chiral symmetry, $U(1)_A$, topology
 - See next talk by K. Suzuki
- \leftrightarrow fix physics and vary T in this study
 - Line of Constant Physics
 - Aiming to study the (pseudo) criticality w/ fixed physics

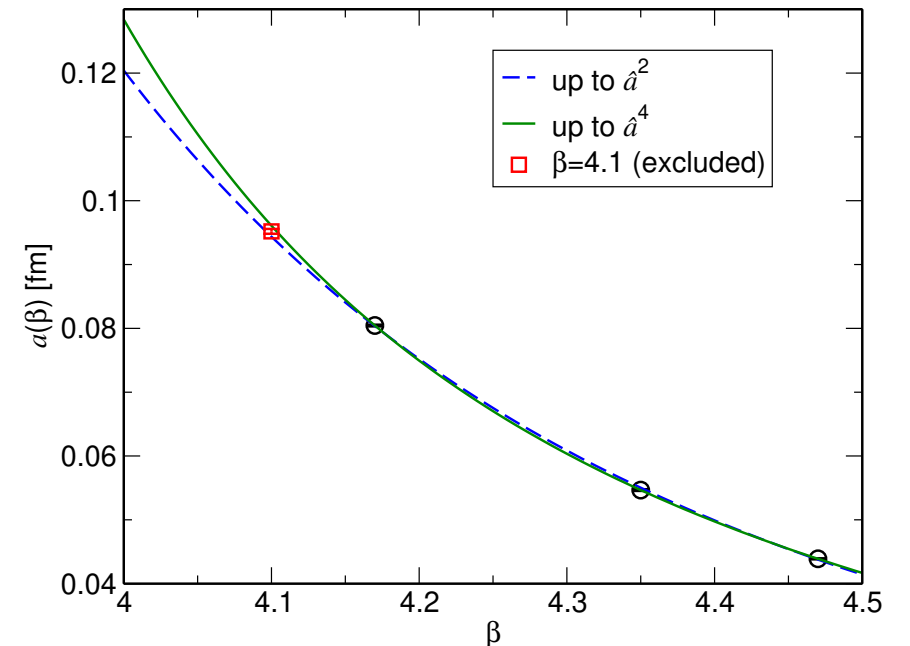
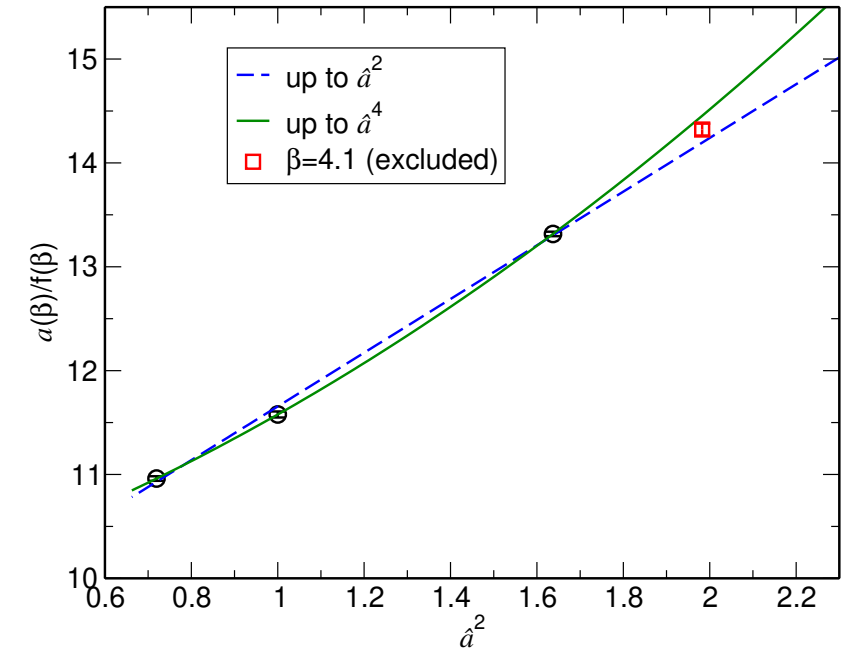
$N_f=2+1$ Möbius DWF

- $a(\beta)$
- Using
 - JLQCD T=0 lattices with t_0 meas.
 - $a=0.080, 0.055, 0.044$ fm (published)
 - $a=0.095$ fm (pilot study)
 - Parameterization of Edwards et al (1998)
 - $a = c_0 f(g^2) (1 + c_2 \hat{a}(g)^2 + c_4 \hat{a}(g)^4)$.
 - $\hat{a}(g)^2 \equiv [f(g^2)/f(g_0^2)]^2$,
 - $f(g^2) \equiv (b_0 g^2)^{-b_1/2b_0^2} \exp\left(-\frac{1}{2b_0 g^2}\right)$,
 - $b_0 = \frac{1}{(4\pi)^2} \left(11 - \frac{2}{3}N_f\right)$, $b_1 = \frac{1}{(4\pi)^4} \left(102 - \frac{38N_f}{3}\right)$,
 - Fit to \hat{a}^4 works well



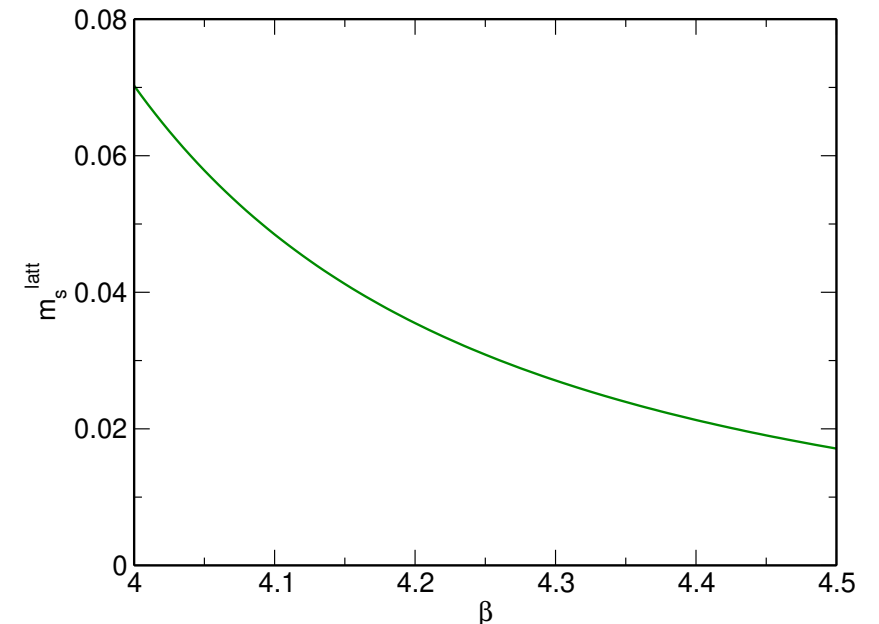
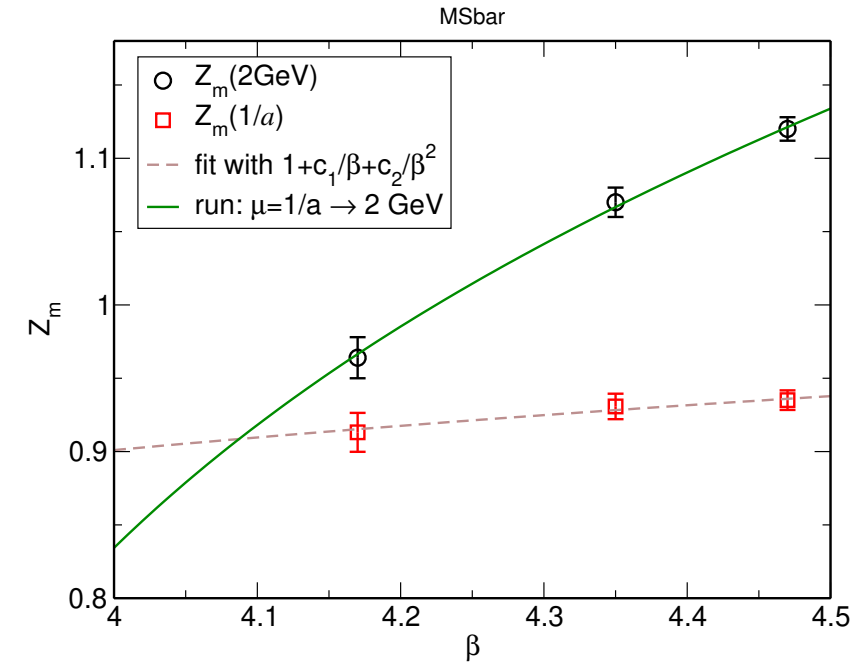
$N_f=2+1$ Möbius DWF

- $a(\beta)$ precision over the range
- Test excluding coarsest one
 - 1 % diff : fit \leftrightarrow measurement @ $\beta=4.1$
 - Difference $O(\hat{a}^4) - O(\hat{a}^2)$ fits: good measure of error (maybe overestimating)
- Full range fit
 - @ $\beta = 4.0$ error is \sim few %
 - The fit may be regarded as renormalized trajectory
 - Continuum limit will absorb the error



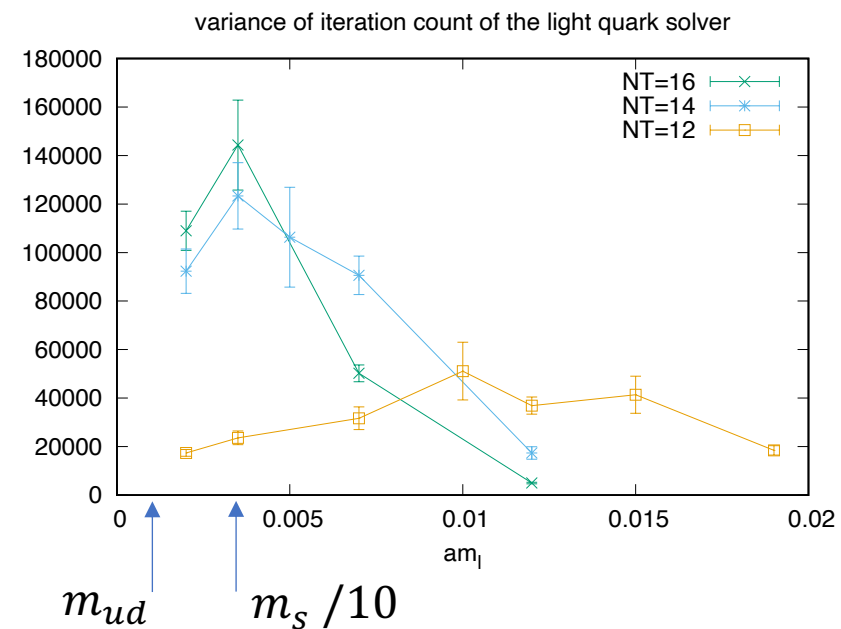
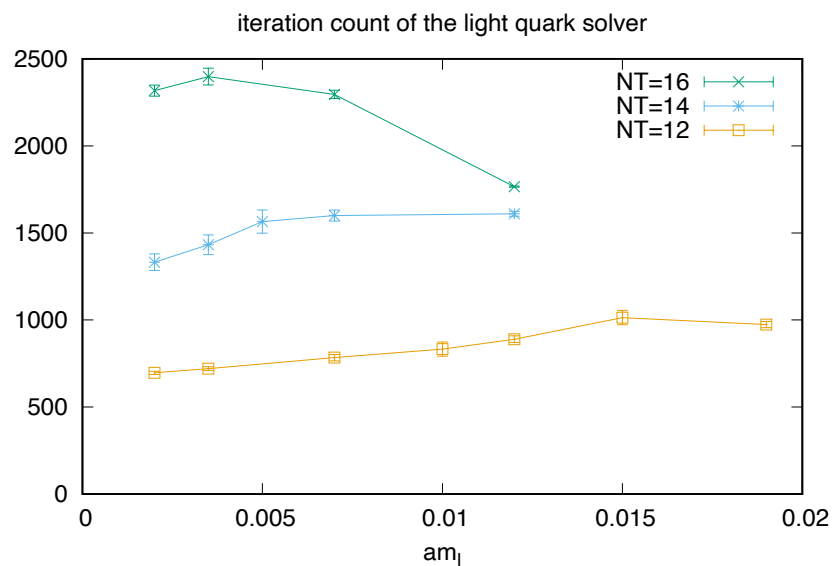
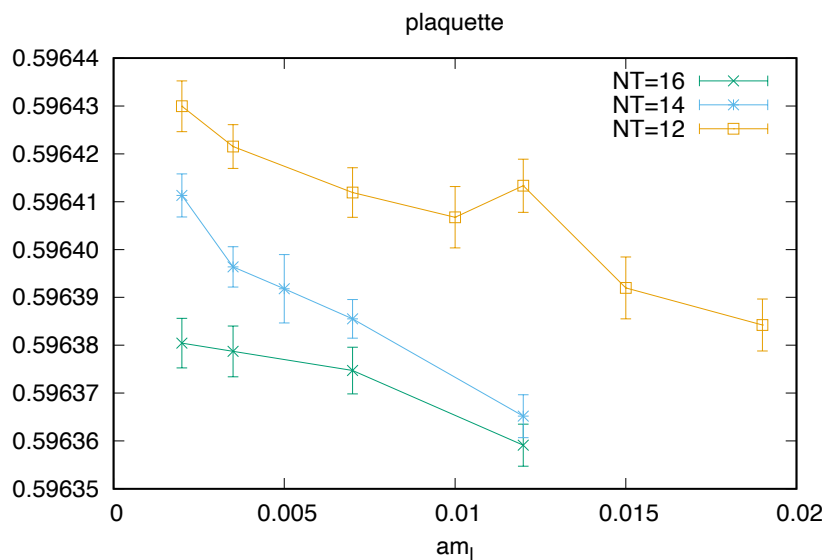
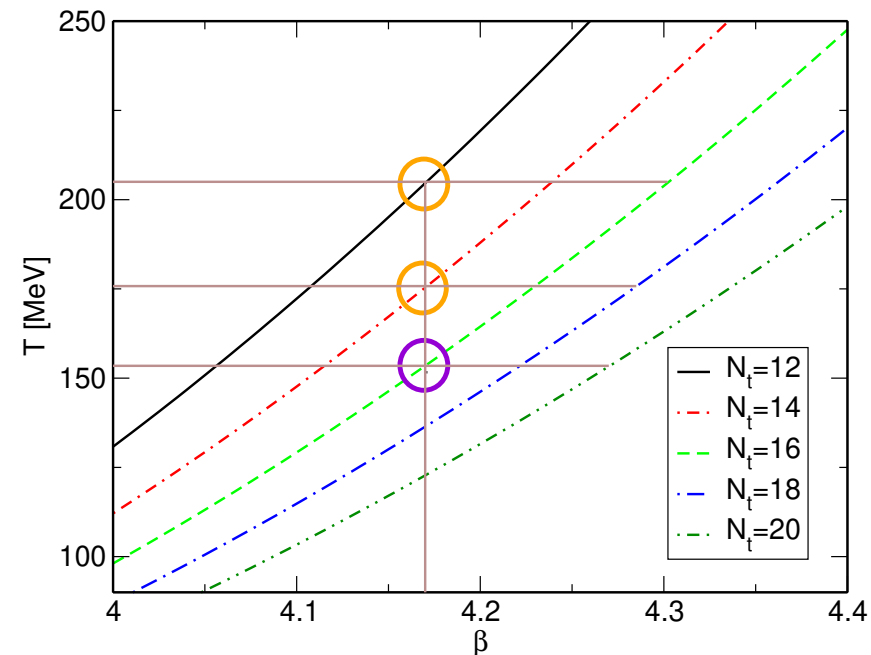
$N_f=2+1$ Möbius DWF LCP

- Quark mass as function of β [fixed physics]
- We use quark mass input
 - $m_s = 92 \text{ MeV}$ (MSb 2GeV)
 - $\frac{m_s}{m_{ud}} = 27.4$ (See for example FLAG 2019)
 - $m_q^R = Z_m \cdot (am_q^{\text{latt}}) \cdot a^{-1}(\beta)$
- Parameterizing $Z_m(\beta)$
 - Take $Z_m(2\text{GeV})$ w/ NPR Tomii et al 2016
 - $Z_m(2\text{GeV}) \rightarrow Z_m(a^{-1})$ NNNLO pert.
 - No (large) $\log(a\mu)$
 - Should behave like $1 + d_1 g^2 + d_2 g^4 + \dots$
 - Fit $Z_m(a^{-1})$ with $1 + c_1 \beta^{-1} + c_2 \beta^{-2}$
 - $Z_m(a^{-1}) \rightarrow Z_m(2\text{GeV})$ NNNLO pert.



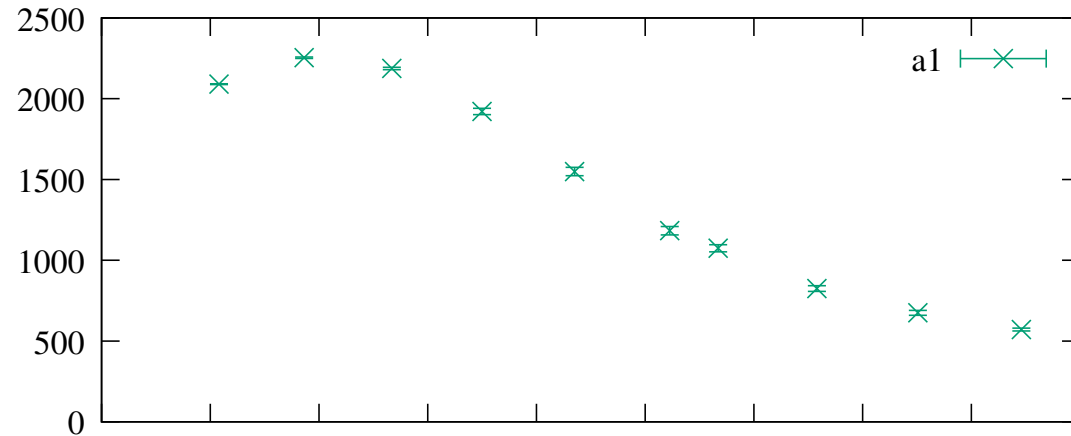
Simulation range

- $T - \beta$ relation $T = 1/(aN_t)$
- Information from fixed β simulation
 - $N_t = 12, 14$
 - $m \rightarrow 0$ study (next talk by K. Suzuki)
 - $N_t = 16 : T \sim 150 \text{ MeV}$
 - $N_s = 32, L_s = 12$
 - $m_{ud}^{latt} = 0.0014, m_s^{latt} = 0.0388$

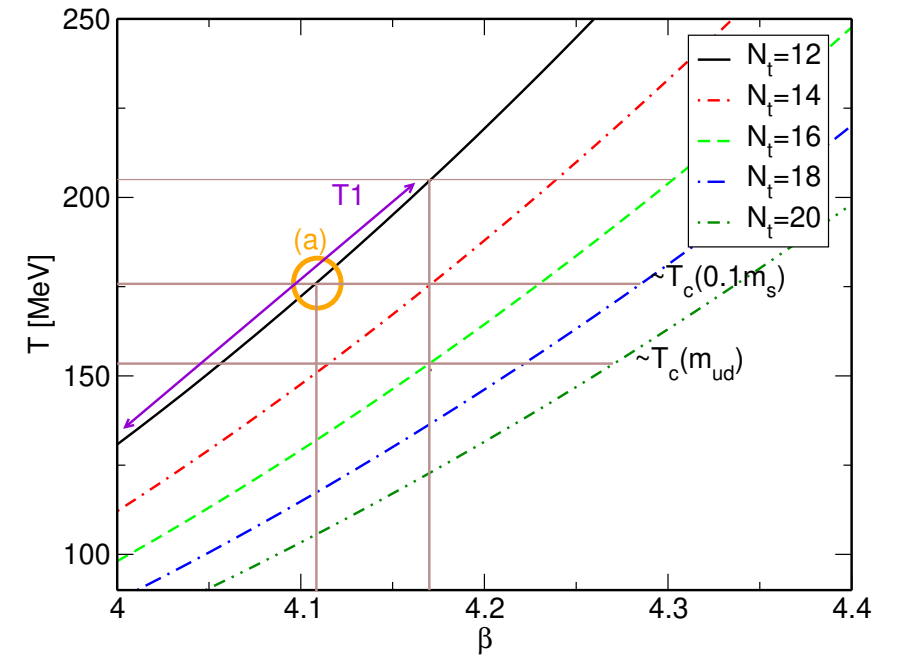
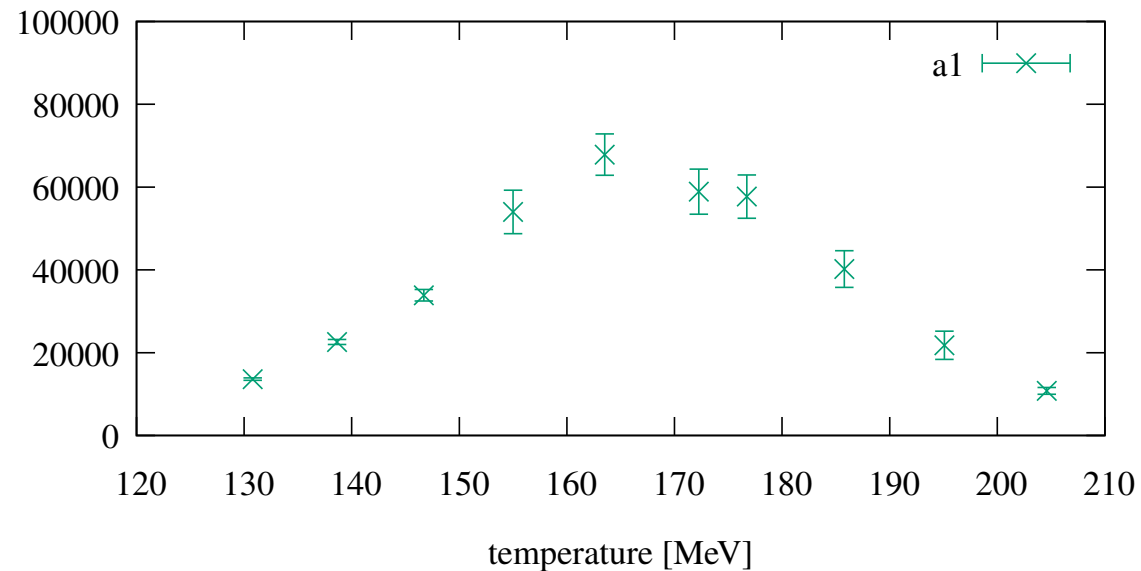


Initial simulations on LCP

iteration count of the light quark solver

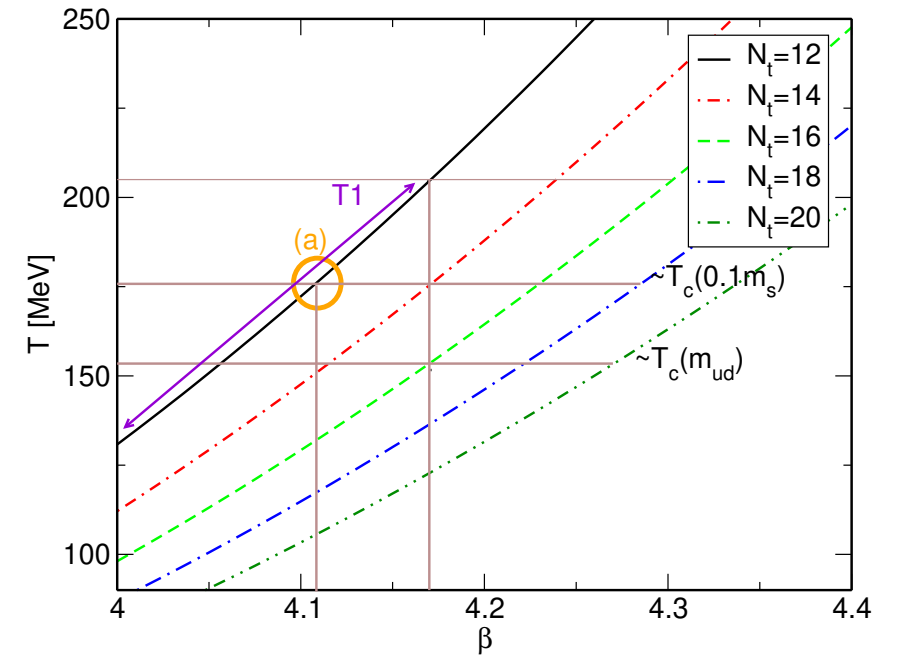
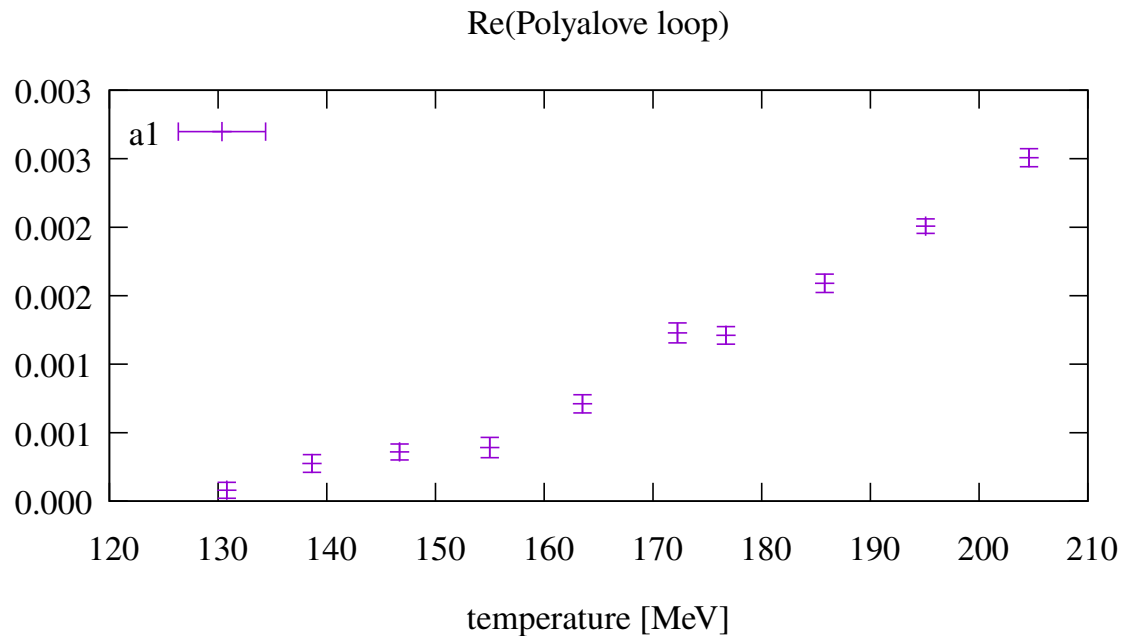
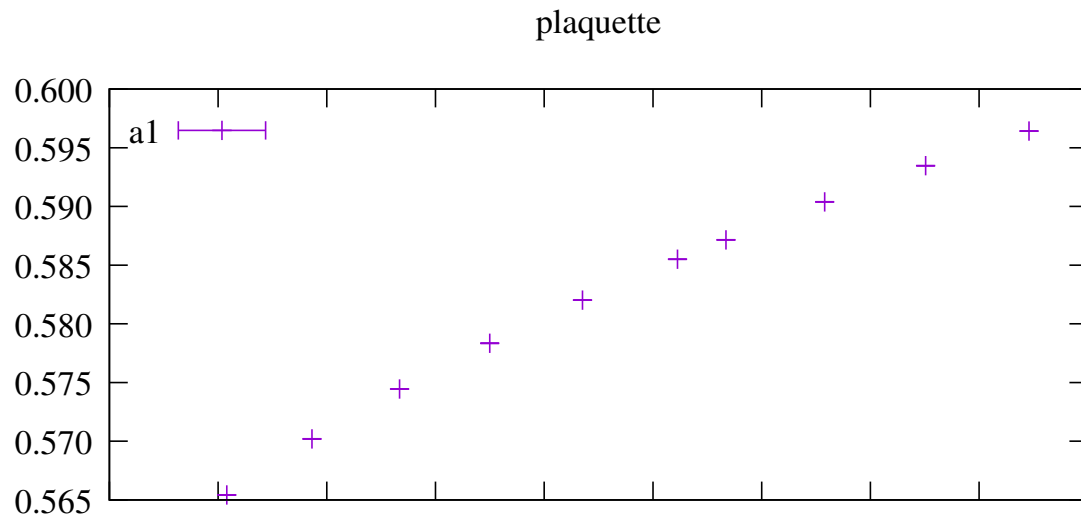


variance of iteration count of the light quark solver



- $N_t = 12$ (T1)
- $m = 0.1m_s$ (a)
- $N_s = 24, L_s = 12$

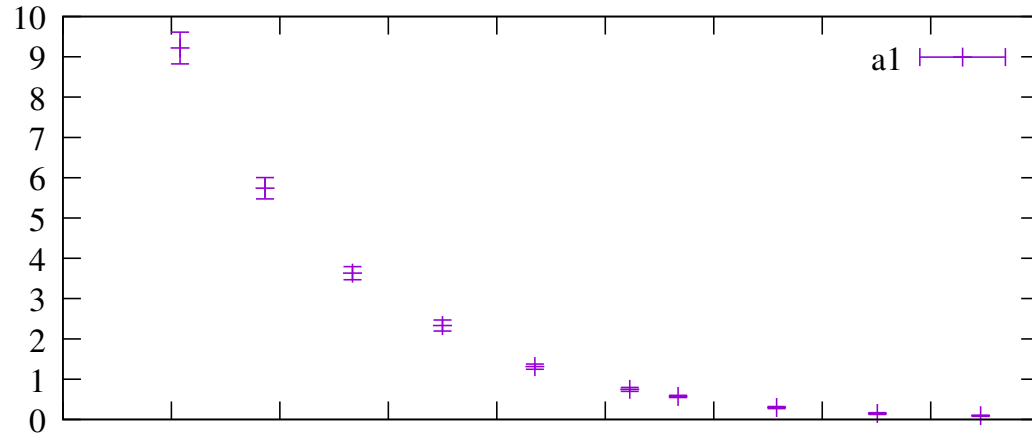
Initial simulations on LCP



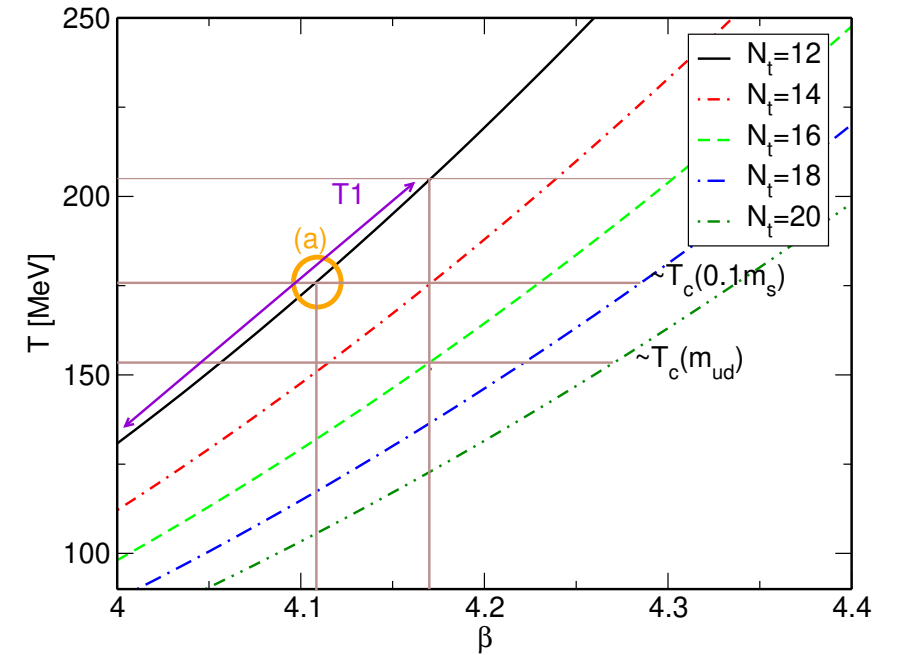
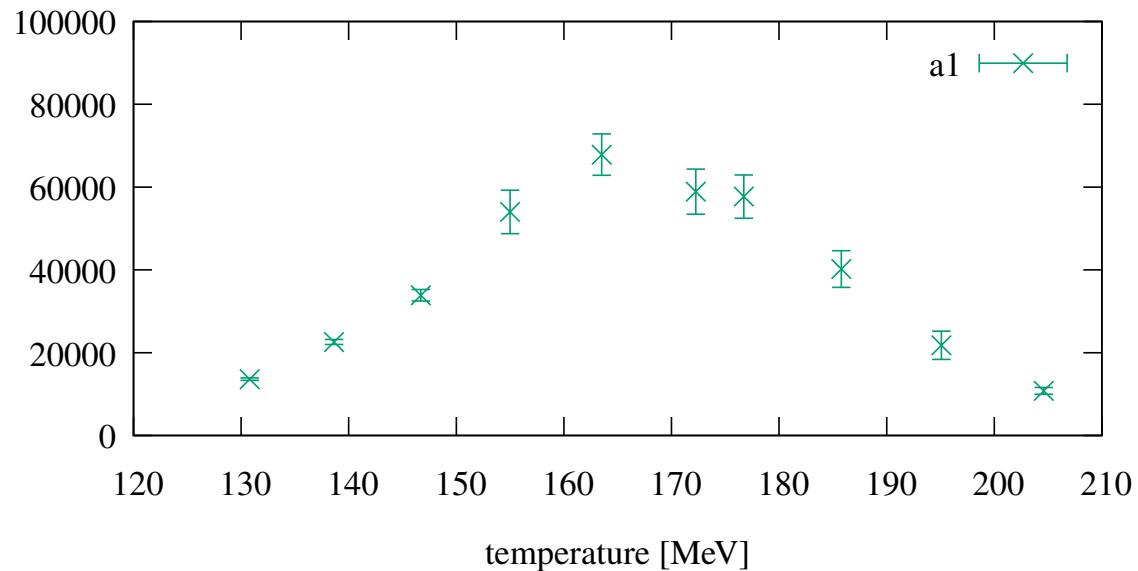
- $N_t = 12$ (T1)
- $m = 0.1m_s$ (a)
- $N_s = 24, L_s = 12$

Initial simulations on LCP

topological charge squared



variance of iteration count of the light quark solver



- $N_t = 12$ (T1)
- $m = 0.1m_s$ (a)
- $N_s = 24, L_s = 12$

Summary and outlook

- Summary
 - Möbius DWF simulation for $T > 0$ with $N_t \geq 12$
 - $\leftrightarrow N_t = 8$ by HotQCD (2012)
 - Along the Line of Constant Physics
 - First simulations with $m = 0.1 m_s$, $N_s/N_t = 2$
 - Underway using Fugaku
- Outlook
 - Statistics is increasing
 - Measurements esp, fermionic
 - Closer to physical mud
 - Another lattice spacing
 - Larger volume

Simulation plan

- T1-(a)
 - $N_t = 12$
 - $m = 0.1m_s$
 - $N_s = 24, L_s = 12$
 - Now underway
- T1-(b)
 - $N_t = 12$
 - $m \simeq m_{ud}$
 - $N_s = 24, L_s = 12$
 - Mass tuning is necessary
 - $m_{res} \simeq m_{ud}$
- T2-(c)
 - $N_t = 16$
 - $m = 0.1m_s$
 - $N_s = 32, L_s = 12$
 - This is straight forward

