

Chromo-electric screening length in 2+1 flavor QCD

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What is the chromo-electric screening length in QCD ?

Use smeared Polyakov loop correlators to enhance the signal at large distances

PP, Sebastian Steinbeißer, Johannes Weber, work in progress



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Deconfinement and color screening

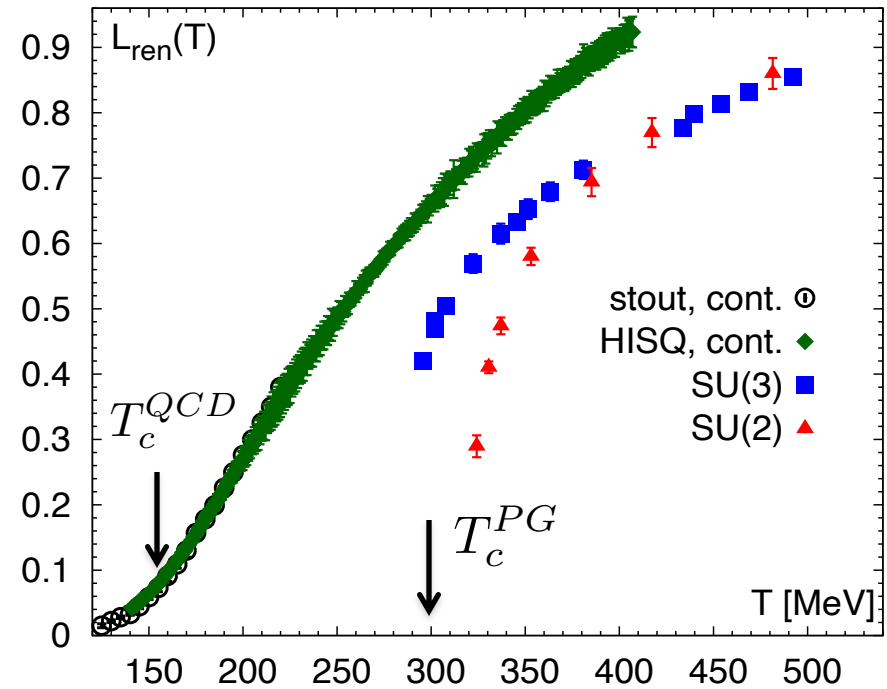
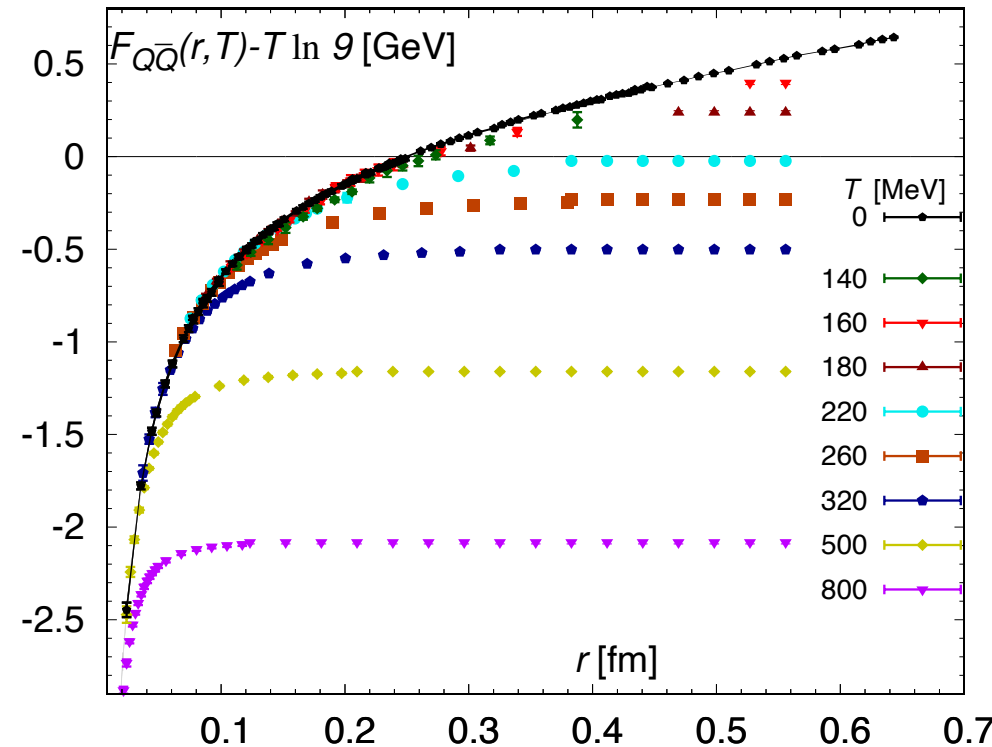
Onset of color screening is described in terms of the free energy of static quark anti-quark pair

$$L(x) = \mathcal{P} \exp \left(-ig \int_0^{1/T} d\tau A_0(x, \tau) \right) \quad \exp(-F_{Q\bar{Q}}(r, T)/T) = \frac{1}{9} \langle \text{tr} L(r) \text{tr} L^\dagger(0) \rangle$$

$$L_{ren} = \exp(-F_Q(T)/T)$$

$$F_{Q\bar{Q}}(r \rightarrow \infty, T) = 2F_Q(T)$$

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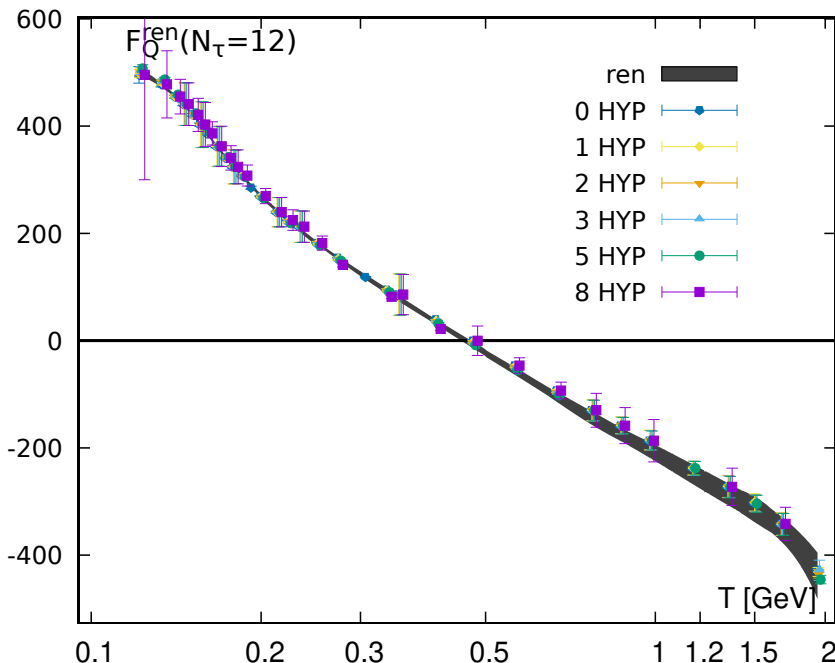
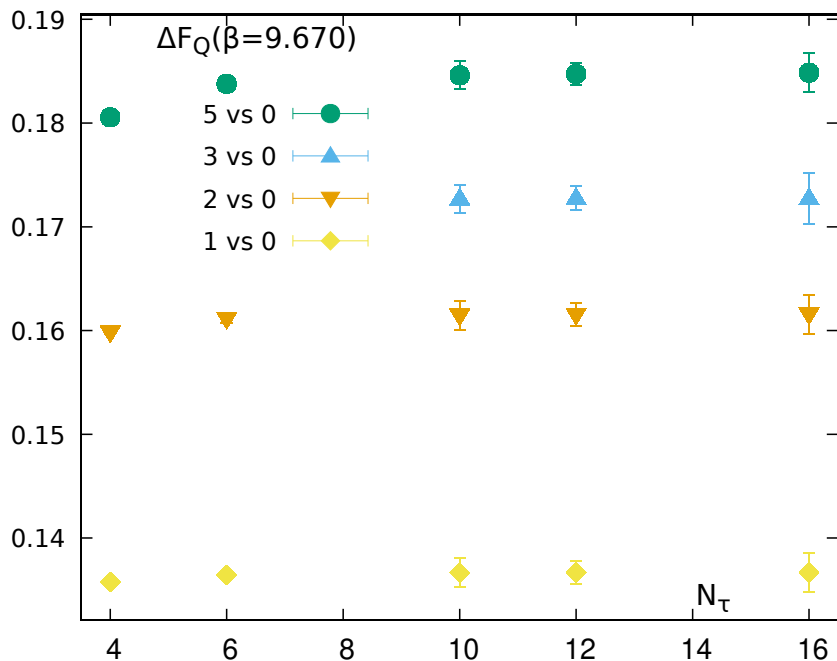
What is the screening mass (length) ? Challenge : poor signal for large separations

Smearred Polyakov loops

Use several steps of HYP smearing, $h=0,1,2,3,5, 8$ on gauge configurations to improve the noise

2+1 flavor QCD, HISQ action, $m_l = m_s/20 \leftrightarrow m_\pi = 160$ MeV, $N_\sigma \times N_\tau$ lattices with $N_\tau = 4, 6, 8, 10, 12$ and 16 , $N_\sigma/N_\tau = 4$ or 6 $T = 155 - 6000$ MeV

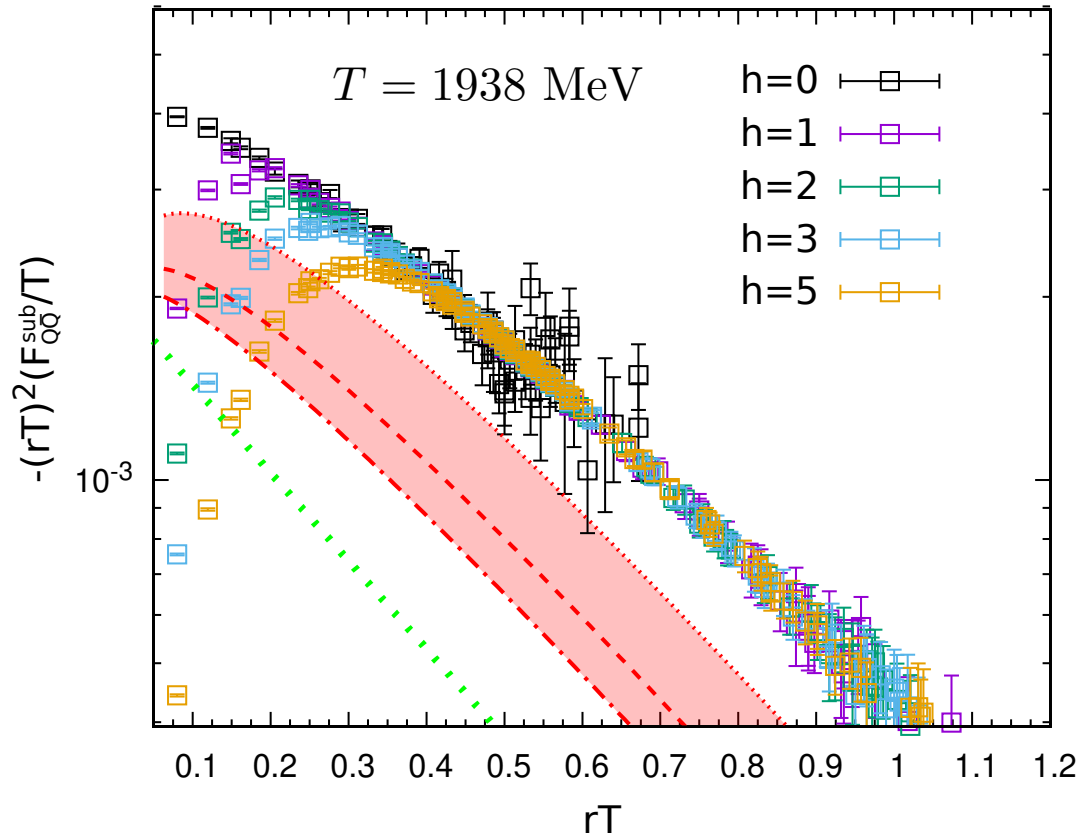
$$\Delta F_Q = F_Q^h - F_Q^{h'}$$



HYP smearing only causes a T -independent shift of F_Q , the T -dependence is not affected by smearing for $N_\tau > 4$

Free energy of static quark anti-quark with smearing

$$F_{Q\bar{Q}}^{sub}(r, T) = F_{Q\bar{Q}}(r, T) - 2F_Q(T)$$



- HYP smearing largely improves the signal at large r for $F_{Q\bar{Q}}(r, T)$
- There is an overlap region for different smearing levels
- There is no distortion for $r > 0.4/T$ where screening sets in
- The lattice results are significantly larger than the LO and NLO perturbative result for $0.4 < rT < 0.8$

Correlators of real and imaginary part of the Polyakov loop

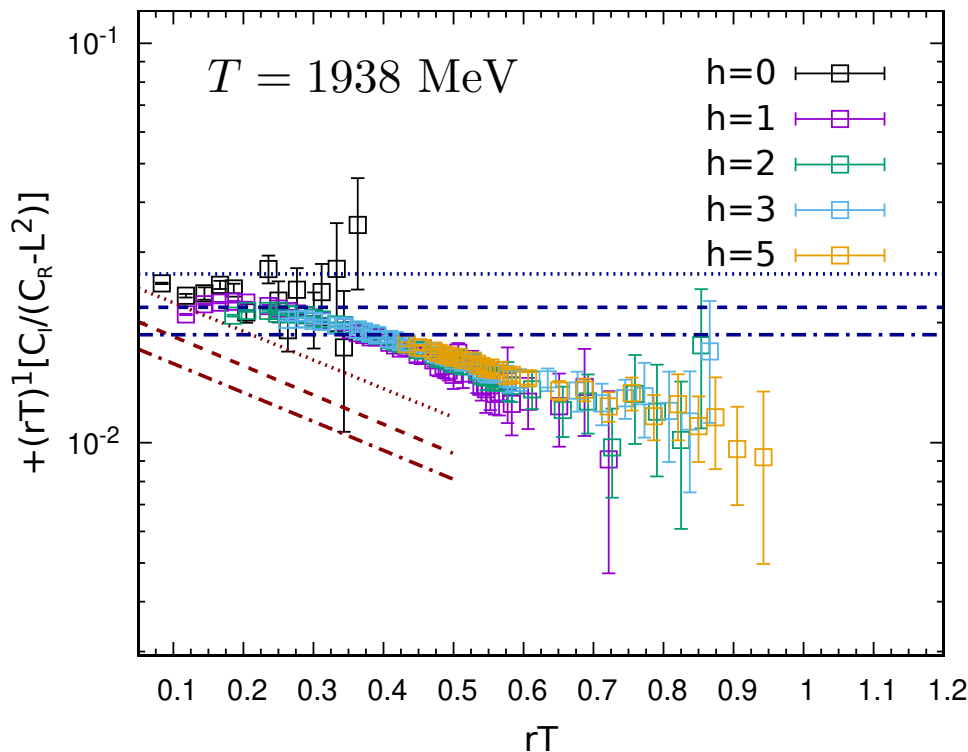
$$C_{PL}(r, T) = C_{PL}^R(r, T) + C_{PL}^I(r, T),$$

$$C_{PL}^R(r, T) = \langle \text{Re}L(r) \cdot \text{Re}L(0) \rangle, C_{PL}^I(r, T) = \langle \text{Im}L(r) \cdot \text{Im}L(0) \rangle.$$

LO perturbative result:

$$C_{PL}^R(r, T) = 1 + 2F_Q^{(0)} + g^4 \frac{N^2 - 1}{8N^2} \left(\frac{e^{-m_{DR}}}{4\pi r T} \right)^2, \quad C_{PL}^I(r, T) = g^6 \frac{N^2 - 1}{8N^2} \frac{N^2 - 4}{12N} \left(\frac{e^{-m_{DR}}}{4\pi r T} \right)^3$$

↑
↑
2-gluon exchange
3-gluon exchange



data with $r/a < \#HYP$ are cut away

- HYP smearing largely improves the signal at large r for $F_{Q\bar{Q}}(r, T)$
- There is an overlap region for different smearing levels
- The lattice results agree with LO result for result for $rT < 0.4$ but are above the LO result at large distances

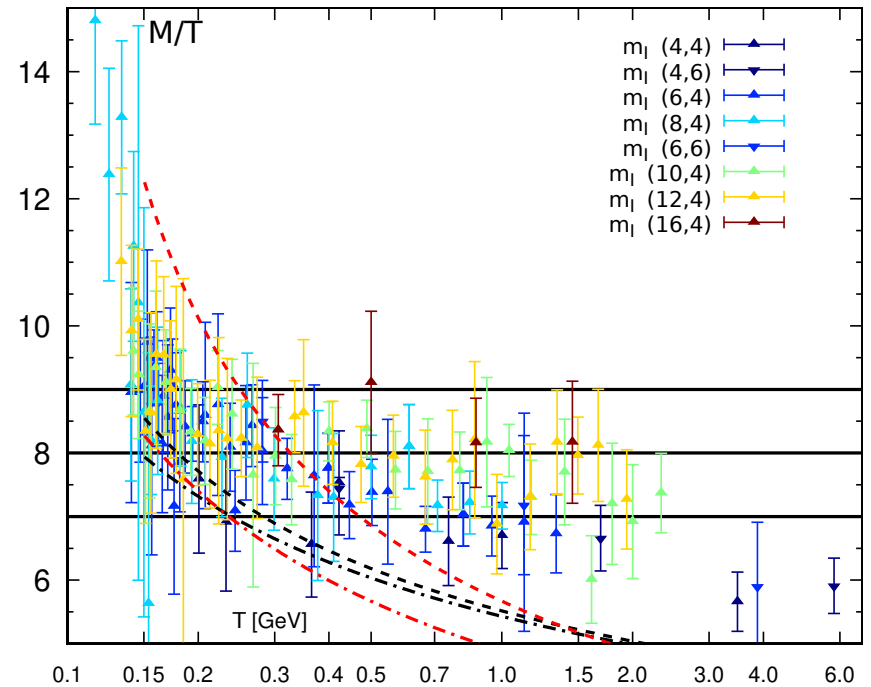
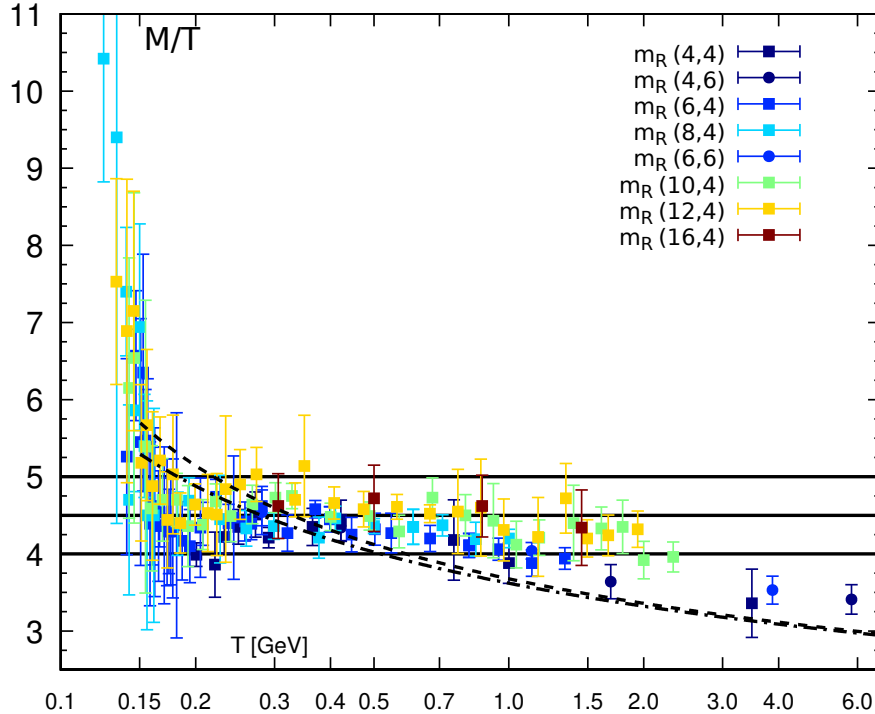
Electric screening masses from the Polyakov loop correlators

The large distance behavior of the correlation functions is governed by screening masses:

$$C_R(r, T) \sim \exp(-m_R(T)r)/r, \quad C_I \sim \exp(-m_I(T)r)/r$$

In the weak coupling picture

$$m_I/m_R \simeq 3/2 \text{ and } m_R \simeq 2m_D, m_I \simeq 3m_D$$



No quantitative agreement between the lattice and the perturbative results but the lattice Results are consistent with the picture of chromo-electric gluon exchange picture:

$$m_I/m_R \simeq 1.75$$

Conclusions

- HYP smearing largely improves the signal in the Polyakov loop and Polyakov loop correlators and the distortion due to smearing can be controlled
- The free energy of static $Q\bar{Q}$ obtained on the lattice for $0.4 < rT < 0.8$ is significantly larger than the NLO perturbative prediction
- The correlator of the imaginary part of the Polyakov loop is well described by LO perturbative result for $rT < 0.4$, but is larger than the LO result for $rT > 0.4$
- We determined the screening masses for the correlator of the real and imaginary part of the Polyakov loop and found $m_R/T \simeq 4.5$ and $m_I/T \simeq 8$ corresponding to screening length of $(0.38 - 0.44)T^{-1}$; Similar results in [Borsanyi et al, JHEP 04 \(2015\) 138](#)

[Datta, Gupta PRD67 \(2003\) 054503](#) (for the ratio of the screening masses in quenched QCD)

Back-up

