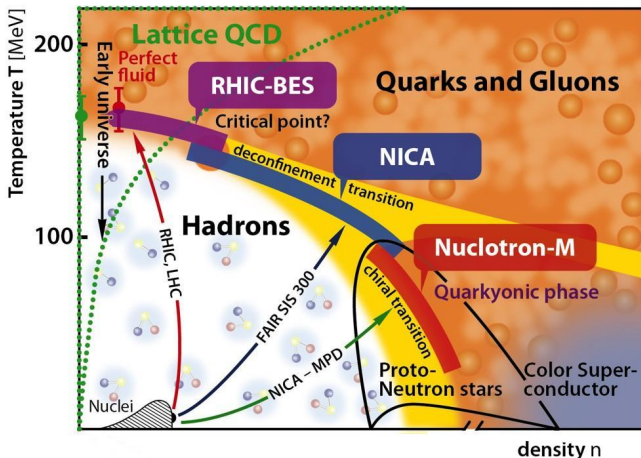


## Electromagnetic conductivity of quark-gluon plasma at non-zero baryon density

N. Astrakhantsev, V. Braguta, M. Cardinali, M. D'Elia,  
L. Maio, F. Sanfilippo, [A. Trunin](#), A. Vasiliev

# QCD phase diagram



Electromagnetic conductivity:

$$\mathbf{j} = \sigma \mathbf{E}$$

is one of the **transport coefficients** of QGP

# How to measure conductivity?

- perturbative QCD and kinetic theory at weak couplings
- holography, effective models
- Lattice QCD

On the lattice, quenched and dynamical fermion results available for  $\mu = 0$   
(see, e.g. recent review [EPJA **57**, 118 (2021); 2008.12326])

$\mu \neq 0$  – sign problem

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- QCD-like theories  
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- simulation at imaginary chemical potential  $i\mu$   
this work

# How to measure conductivity on the lattice?

Kubo relation:

$$\sigma = \frac{1}{6} \lim_{\omega \rightarrow 0} \frac{\rho^{ii}(\omega)}{\omega},$$

$\rho^{\mu\nu}(\omega) = \int d^4x e^{i\omega t} \langle [j^\mu(t, \mathbf{x}), j^\nu(0)] \rangle$  – current-current spectral function

$j^\mu(x) = \sum_f (eq_f) \bar{\psi}_f(x) \gamma^\mu \psi_f(x)$  – electromagnetic current

On the lattice we measure in Euclidean time

$$G_E(\tau) = \frac{1}{V} \sum_{\mathbf{x}} \langle j_i(\tau, \mathbf{x}) j_i(0) \rangle$$

$$G_E(\tau) = \int_0^\infty \frac{d\omega}{2\pi} K(\omega, \tau) \rho(\omega), \quad K(\omega, \tau) = \frac{\cosh[\omega(\tau - 1/2T)]}{\sinh(\omega/2T)}$$

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We use Backus-Gilbert method with Tikhonov regularisation

# Lattice setup

- $N_f = 2 + 1$  staggered fermions
- physical pion and strange masses
- $T \simeq 200, 250$  MeV
- imaginary chemical potential:

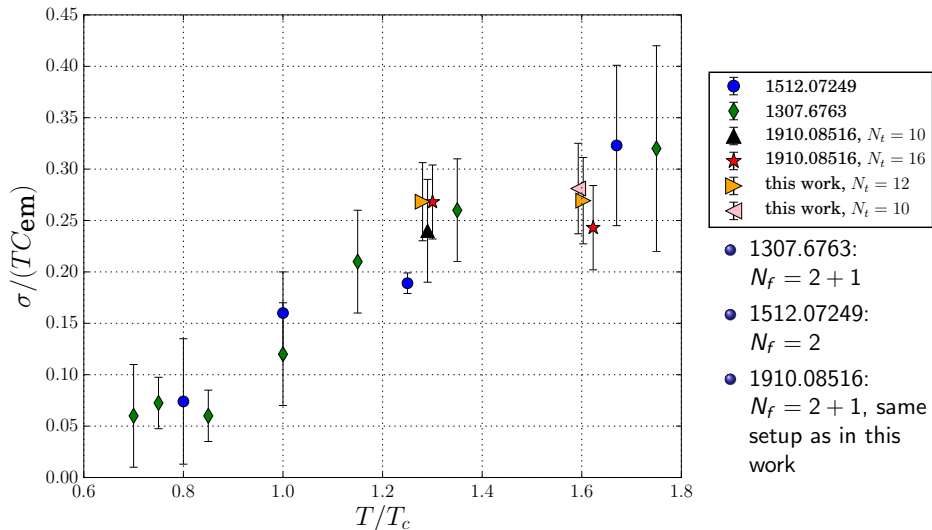
$$\mu_s = 0$$

$$a\mu_u = a\mu_d = 0.140, 0.200, 0.245, 0.285$$

$N_s$	$N_t$	$a$ [fm]	$T$ [MeV]
48	10	0.099	200
48	10	0.079	250
48	12	0.082	200
48	12	0.066	250



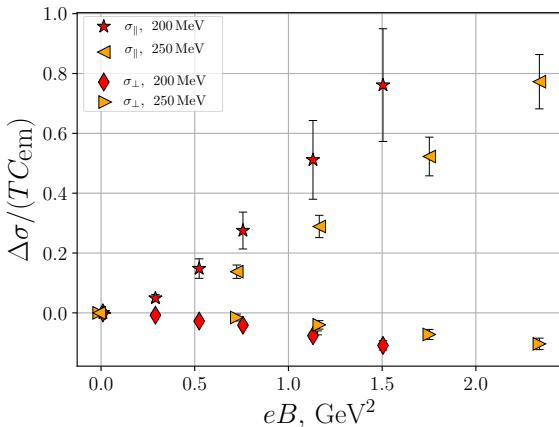
# Electromagnetic conductivity: $\mu = 0$



# Electromagnetic conductivity: $\mu = 0$

Conductivity at external magnetic field  $\mathbf{B}$ :

$$\Delta\sigma(B) \equiv \sigma(B) - \sigma(B = 0)$$



Conductivity rises in direction  $\parallel \mathbf{B} \implies$  chiral magnetic effect

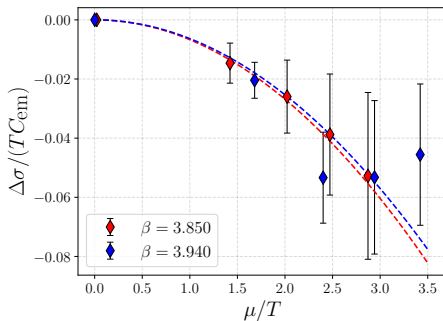
[PRD 102, 054516 (2020); 1910.08516]

# Electromagnetic conductivity: imaginary $\mu$

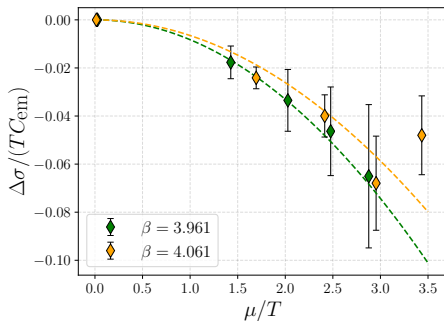
$$\Delta\sigma \equiv \sigma(\mu) - \sigma(\mu = 0)$$

Preliminary:

$T = 200$  MeV



$T = 250$  MeV



$$\frac{\sigma(\mu, T)}{\sigma(\mu = 0, T)} \simeq 1 + A_2(T) \left(\frac{\mu}{T}\right)^2$$

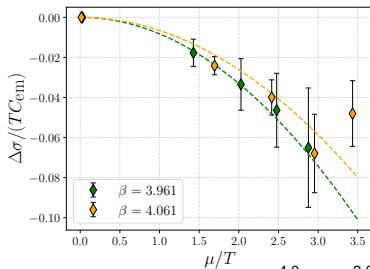
We are at imaginary  $\mu$ :

$$\mu \rightarrow i\mu_{\text{Re}} \quad \Rightarrow$$

Conductivity rises with baryon density

# Electromagnetic conductivity: imaginary $\mu$

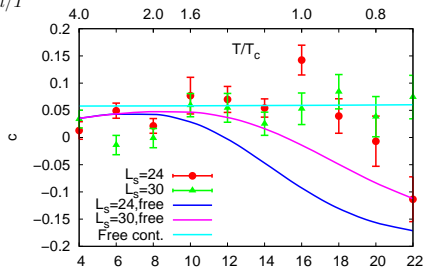
Preliminary:



$$\frac{\sigma(\mu, T)}{\sigma(\mu = 0, T)} \simeq 1 + A_2(T) \left(\frac{\mu}{T}\right)^2$$

$$A_2 \simeq 0.02 - 0.08$$

- large uncertainty from reconstruction
- agreement with SU(2) within errors



SU(2) [PRD **102**, 094510 (2020);  
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# Summary

- The first calculation of the conductivity in QCD at physical pion mass and finite baryon density
- Conductivity increases with chemical potential
- Small  $A_2 \implies$  very weak dependence on chemical potential
- The obtained value of  $A_2$  agrees with SU(2) results within errors

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