

Study of the EoS of dense QCD in an external magnetic field

Victor V. Braguta, Natalia V. Kolomojets, Andrey Yu. Kotov, Artem A. Roenko

Joint Institute for Nuclear Research, Russia

Nikita Yu. Astrakhantsev

Universität Zürich, Switzerland

Alexander A. Nikolaev

Department of Physics, College of Science, Swansea University, Great Britain

LATTICE21

July 28, 2021

- 1 Introduction
- 2 Simulation setup
- 3 Results (preliminary)
- 4 Conclusions

Introduction

Extreme conditions:

- high T
- $\mu_B \neq 0$
- large eB .

There is success in study of EoS

- at finite density

D. Bollweg et al., Nucl. Phys. A 1005, 121835 (2021)
[arXiv:2002.01837 [hep-lat]]; others

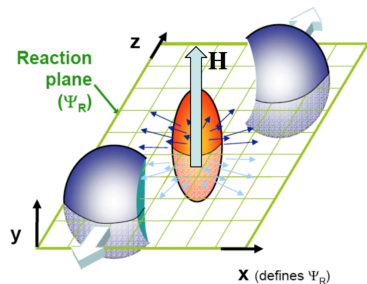
- at nonzero magnetic field

M. D'Elia et al., Phys. Rev. D 98, no.5, 054509 (2018)
[arXiv:1808.07008 [hep-lat]]; others

There are much less investigations at both $\mu_B \neq 0$ and $eB \neq 0$

H. T. Ding et al., Eur. Phys. J. A 57, no.6, 202 (2021)
[arXiv:2104.06843 [hep-lat]]

We want to obtain EoS from 2+1 LQCD accounting for both $\mu_B \neq 0$ and $eB \neq 0$



$$eB \sim 0.1 \text{ GeV}^2$$

⋈

Strongly influences QCD properties,
in particular, phase diagram

Thermodynamics on the lattice

$$p = -\frac{\Omega}{V} = \frac{T}{V} \ln \mathcal{Z} \quad \leftarrow \text{cannot be measured directly}$$

$$\int \star \mathcal{Z} = \int \mathcal{D}[U] \mathcal{D}[\psi, \bar{\psi}] e^{-S(\psi, \bar{\psi}, U)} \quad \star /$$

Derivatives of p can be measured!

$$n_q = \frac{N_q}{V} = \frac{\partial p}{\partial \mu_q} \quad - \text{quark number density}$$

$\mu_B = \mu_u + 2\mu_d$	$n_B = (n_u + n_d + n_s)/3$
$\mu_Q = \mu_u - \mu_d$	$n_Q = (2n_u - n_d - n_s)/3$
$\mu_S = \mu_d - \mu_s$	$n_S = -n_s$

J. N. Günther et al., Nucl. Phys. A **967**, 720 (2017) [arXiv:1607.02493 [hep-lat]]

$$\frac{p}{T^4} = c_0(T) + c_2(T) \left(\frac{\mu_B}{T}\right)^2 + c_4(T) \left(\frac{\mu_B}{T}\right)^4 + c_6(T) \left(\frac{\mu_B}{T}\right)^6 + \mathcal{O}(\mu_B^8)$$

$$\frac{n}{\mu_B T^2} = \frac{T}{\mu_B} \cdot \frac{d(p/T^4)}{d(\mu_B/T)} = 2c_2 + 4c_4 \left(\frac{\mu_B}{T}\right)^2 + 6c_6 \left(\frac{\mu_B}{T}\right)^4 \quad \leftarrow \text{coefficients can be found from fit}$$

$(\mu_u, \mu_d, \mu_s \text{ are such that } \langle n_S \rangle = 0, \langle n_Q \rangle = 0.4 \langle n_B \rangle)$

Calculation of c_0 :

S. Borsanyi et al., Phys. Lett. B **730**, 99 (2014) [arXiv:1309.5258 [hep-lat]].

G. S. Bali et al., JHEP **08**, 177 (2014) [arXiv:1406.0269 [hep-lat]].

Our choice of chemical potentials:

$$\mu_u = \mu_d = \mu_q; \quad \mu_s = 0$$

\Downarrow

$$\mu_B = 3\mu_q; \quad \mu_Q = 0; \quad \mu_S = \mu_q$$

Lattice setup

- Tree level improved Symanzik gauge action.
- Staggered 2 + 1 fermionic action.
- Stout smearing improvement.
- Imaginary chemical potential: $\mu_q = i\mu_I$.
- External magnetic field:

$$\vec{B} = B\vec{e}_z; \quad B = \text{const}$$

$$A_y^{\text{ext}} = Bx/2, \quad A_x^{\text{ext}} = -By/2, \quad A_\mu^{\text{ext}} = 0, \quad \mu = z, t$$

- Splitting of the rooted determinant:

$$\mathcal{Z} = \int \mathcal{D}U e^{-S_G} [\det D(B, m_u, q_u)]^{\frac{1}{4}} [\det D(B, m_d, q_d)]^{\frac{1}{4}} [\det D(B, m_s, q_s)]^{\frac{1}{4}}$$

$$D(n|f) = \frac{1}{2a} \sum_{\mu} \eta_{\mu}(n) \left[u_{\mu}(B, q, n) e^{ia\mu_I \times \delta_{\mu 4}} U_{\mu}(n) \delta_{f, n+\hat{\mu}} - u_{\mu}^*(B, q, f) e^{-ia\mu_I \times \delta_{\mu 4}} U_{\mu}^{\dagger}(f) \delta_{f, n-\hat{\mu}} \right] + m \delta_{f, n}$$

$$u_x(B, q, n_x, n_y, n_z, n_t) = e^{-ia^2 q B n_y/2}, \quad n_x \neq N_x - 1,$$

$$u_y(B, q, n_x, n_y, n_z, n_t) = e^{ia^2 q B n_x/2}, \quad n_y \neq N_y - 1,$$

$$u_x(B, q, N_x - 1, n_y, n_z, n_t) = e^{-ia^2 q B (N_x + 1) n_y/2},$$

$$u_y(B, q, n_x, N_y - 1, n_z, n_t) = e^{ia^2 q B (N_y + 1) n_x/2}.$$

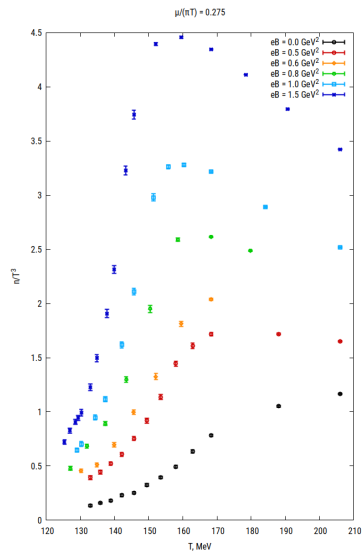
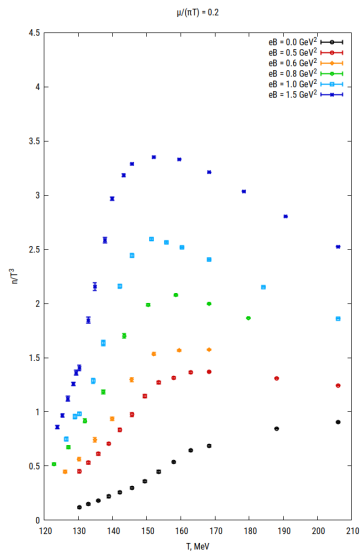
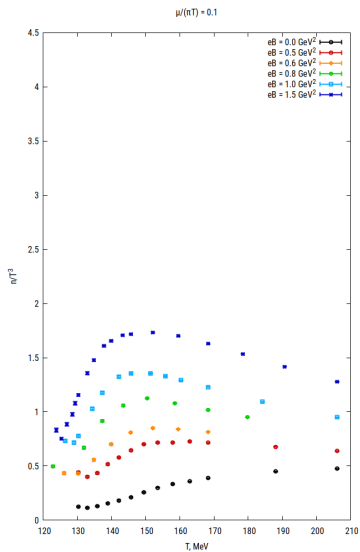
$$\text{Periodic boundary conditions} \quad \Rightarrow \quad eB = \frac{6\pi k}{N_x N_y a^2}, \quad k \in \mathbb{Z}$$

Simulation parameters:

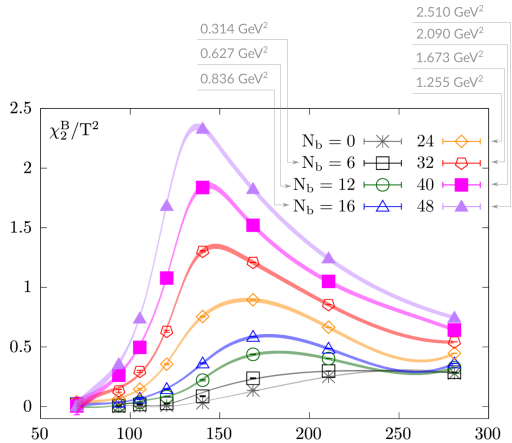
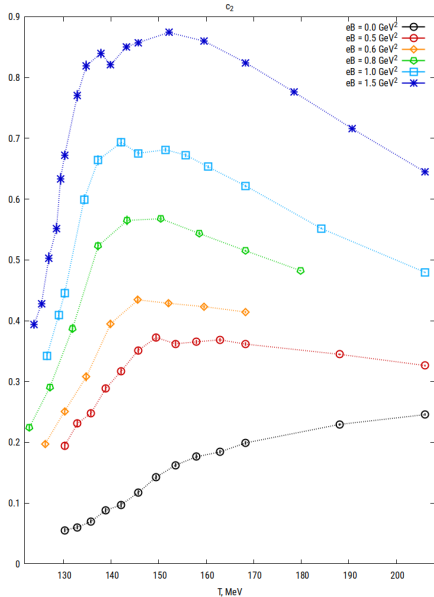
Lattice	6×24^3	8×32^3	10×40^3
$eB, \text{ GeV}^2$	0.0, 0.5, 0.6, 0.8, 1.0, 1.5	0.6	0.6
$T, \text{ MeV}$	123 - 206	126 - 198	126 - 198

physical quark masses.

Results: densities at 6×24^3 lattice



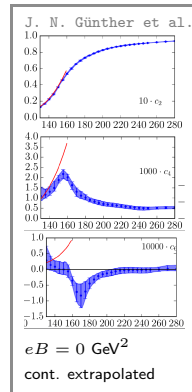
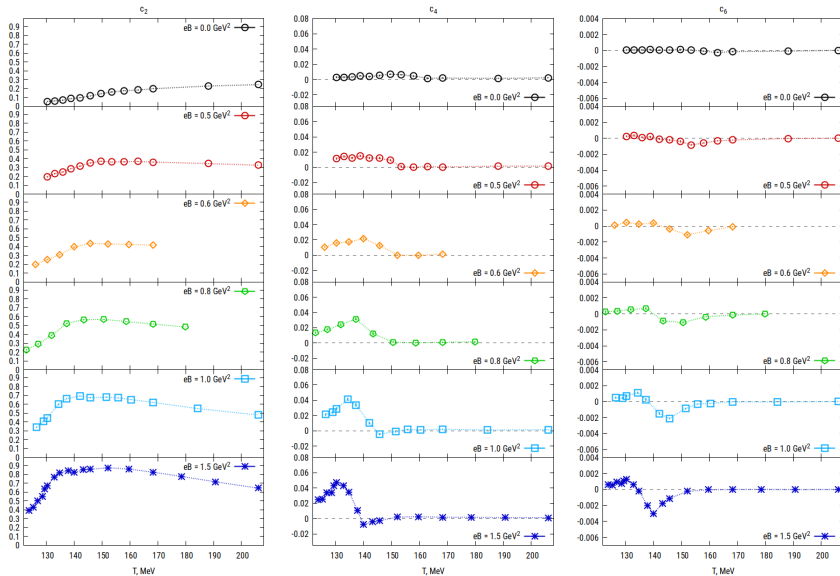
Results: expansion coefficient c_2 at 6×24^3 lattice



H. T. Ding et al., Eur. Phys. J. A57, no.6, 202 (2021)
[arXiv:2104.06843 [hep-lat]]

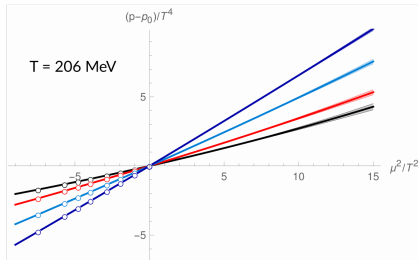
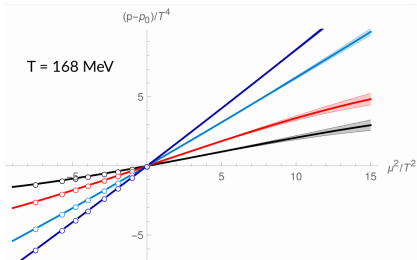
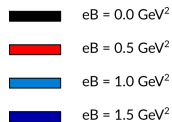
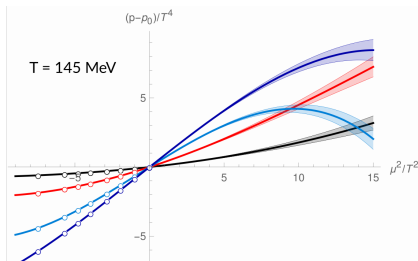
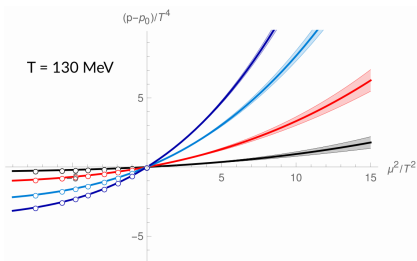
PRELIMINARY

Results: expansion coefficients at 6×24^3 lattice



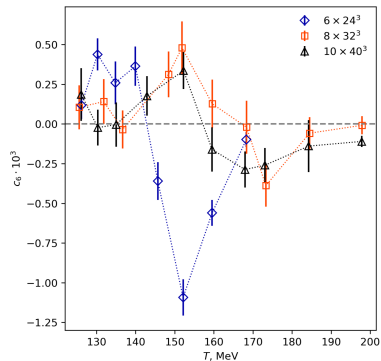
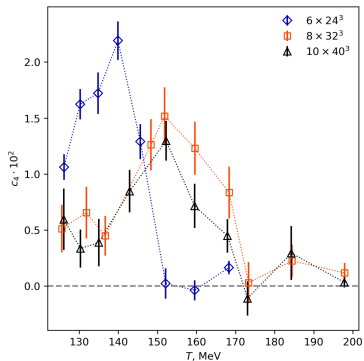
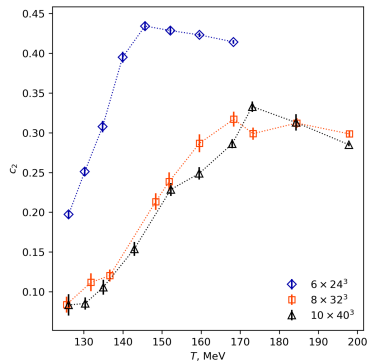
PRELIMINARY

Results: pressure



Results: comparison of coefficients for different lattices

$$eB = 0.6 \text{ GeV}^2$$



PRELIMINARY

Conclusions

- Simulations at non-zero chemical potential and with external magnetic field have been carried out on 3 lattices.
- First results on expansion coefficients c_2 , c_4 , c_6 in external magnetic field have been obtained.
- Strong dependence of the EoS on magnetic field.