

# Lee-Yang edge singularities in 2+1 flavor QCD with imaginary chemical potential.

Guido Nicotra



27th July Tuesday

**Bielefeld-Parma Collaboration:**

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G. Nicotra, C. Schmidt, S. Singh, K. Zambello, F. Ziesché



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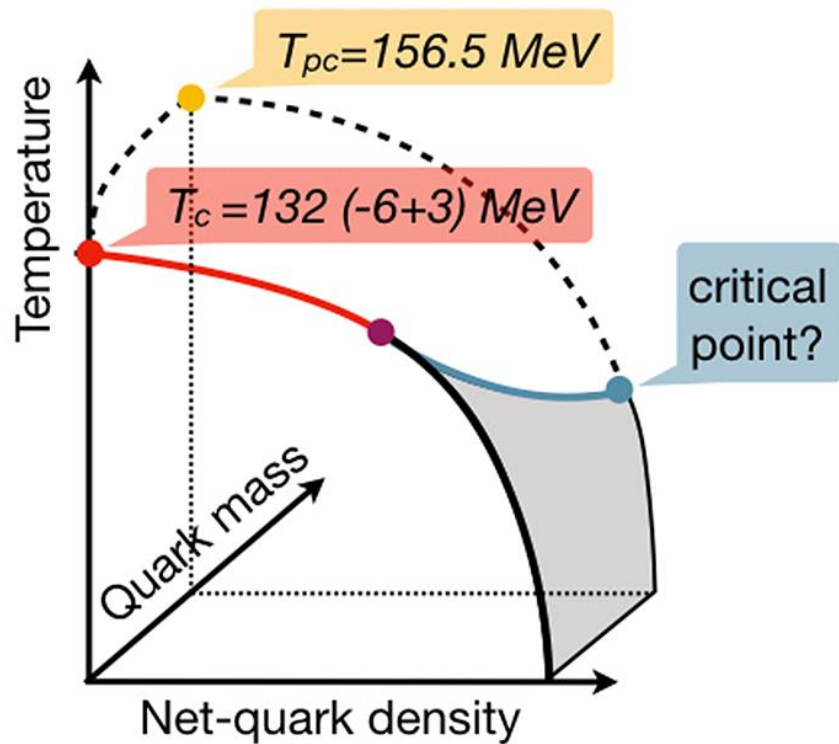


UNIVERSITÄT  
BIELEFELD

Fakultät für Physik



# Introduction



In the search of the critical point.

Sign problem:

Introduce a quark chemical potential into the action;  
Determinant of the Dirac's Matrix become complex;  
No probabilistic interpretation for MC simulations.

How we manage it:

Simulations with imaginary baryon chemical potential  
in the time direction.

What interests us:

Universal critical behavior of phase transitions in the  
complex plane.

# Lattice settings and observables

2+1 flavors, HISQ, Rational Hybrid Monte Carlo (RHMC);  
Bielefeld GPU code;  
 $N_\sigma^3 \times N_\tau : 24^3 \times 4, 36^3 \times 6.$

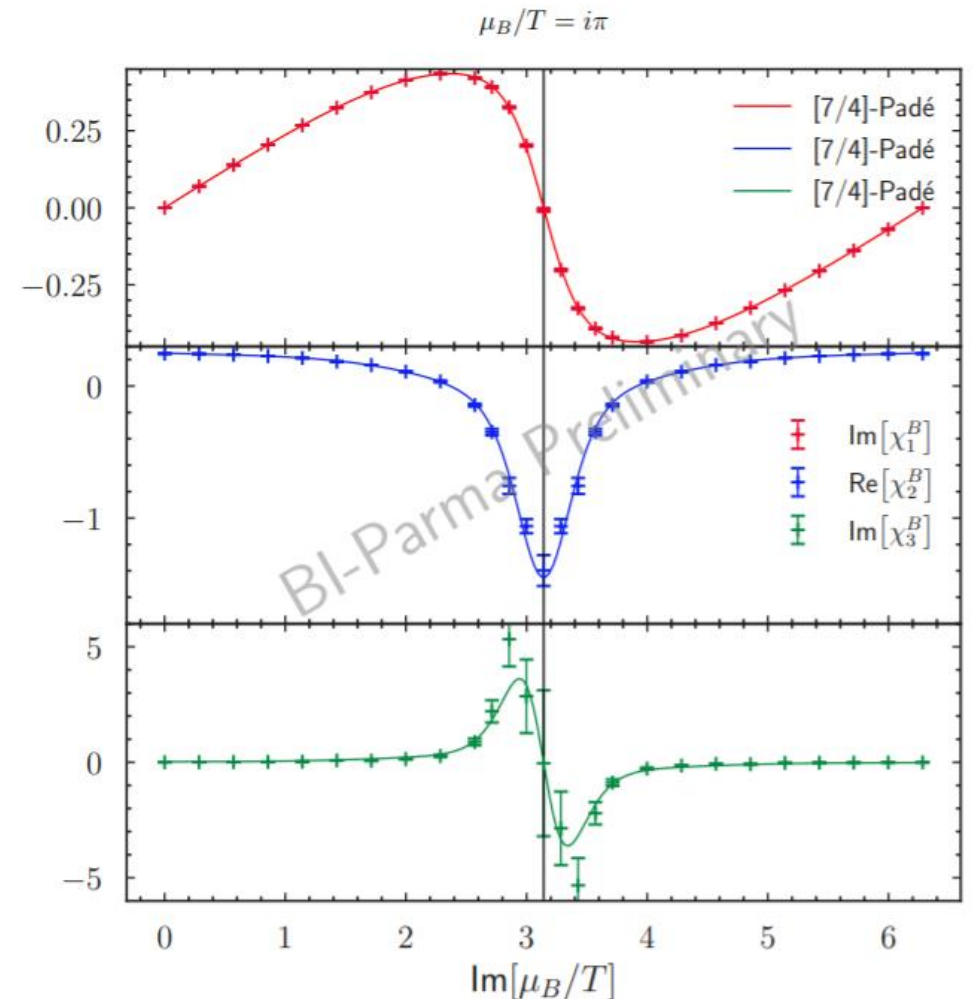
$$Z(T, V, \mu) = \int DU \cdot \det(M_l(i\mu))^{1/2} \cdot \det(M_s(i\mu))^{1/4} \cdot e^{-S_G}$$

$$\chi_n^B = \frac{1}{n! V T^3} \frac{\partial^n \ln Z}{\partial (\mu_B/T)}$$

From Taylor series at different purely imaginary baryon chemical potential we have built multi-point Padé approximants:

$$R_n^m(x) = \frac{P_m(x)}{1 + Q_n(x)} = \frac{\sum_{i=0}^m a_i x^i}{1 + \sum_{j=1}^n b_j x^j}$$

[Simran Singh, Lattice 2021, Tuesday at 6.45 EST]



# Magnetic Equation of State (EoS) and Lee-Yang Edge singularity

[Michael E. Fisher, Phys. Rev. Lett. 40, 1610 (1978)]

We can write the order parameter as:

$$M(t, h) = h^{\frac{1}{\delta}} f_G(z) + M_{reg}(t, h) \quad \text{Magnetic EoS}$$

$t$  = field responsible for the deviation from criticality without explicit symmetry breaking;

$h$  = symmetry breaking field;

$f_G(z)$  = scaling function;

$z$  = scaling variable  $t/h^{1/\beta\delta}$ ;

$\beta, \delta$  = universal critical exponent.

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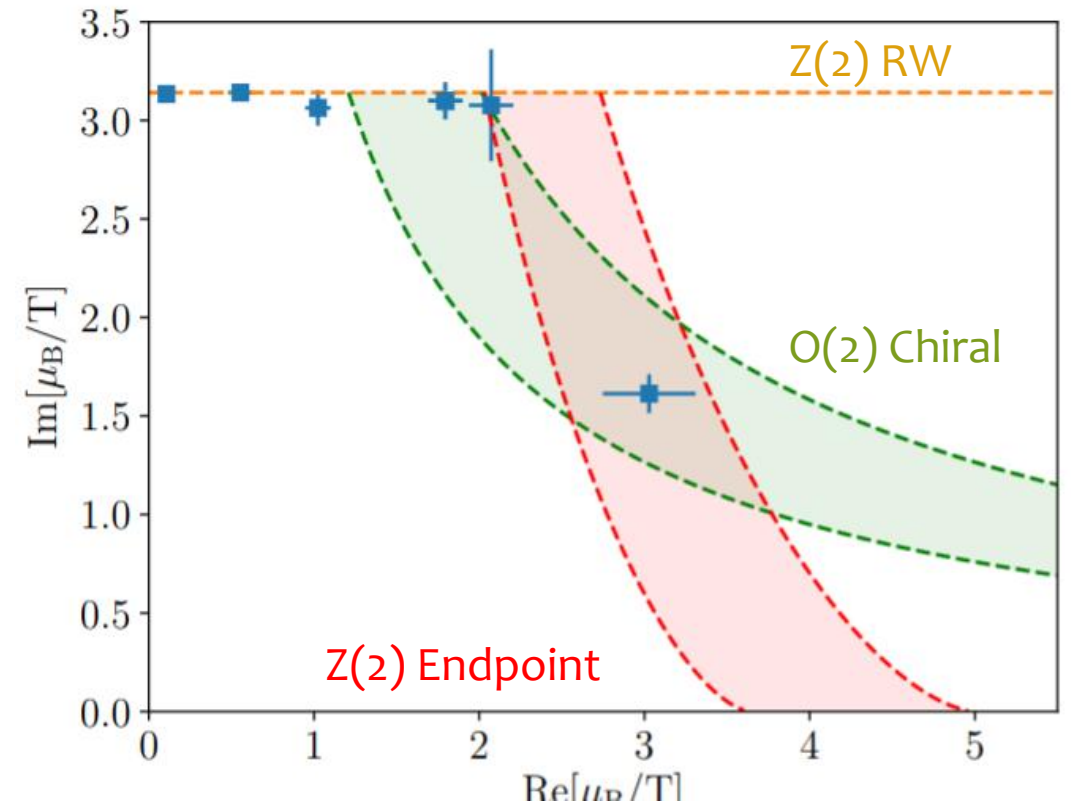
$\beta, \delta$  = universal critical exponent.

Approaching the critical point,  $M_{reg}(t, h)$  does not have singularities in this region.

$h^{\frac{1}{\delta}} f_G(z)$  governs the non analytic behavior.

The singularity of the scaling function in  $z = z_c$  is called **Lee Yang edge singularity (LYE)**.

Now we have to **choose the scaling fields**  $t, h$  in agreement with the transition.



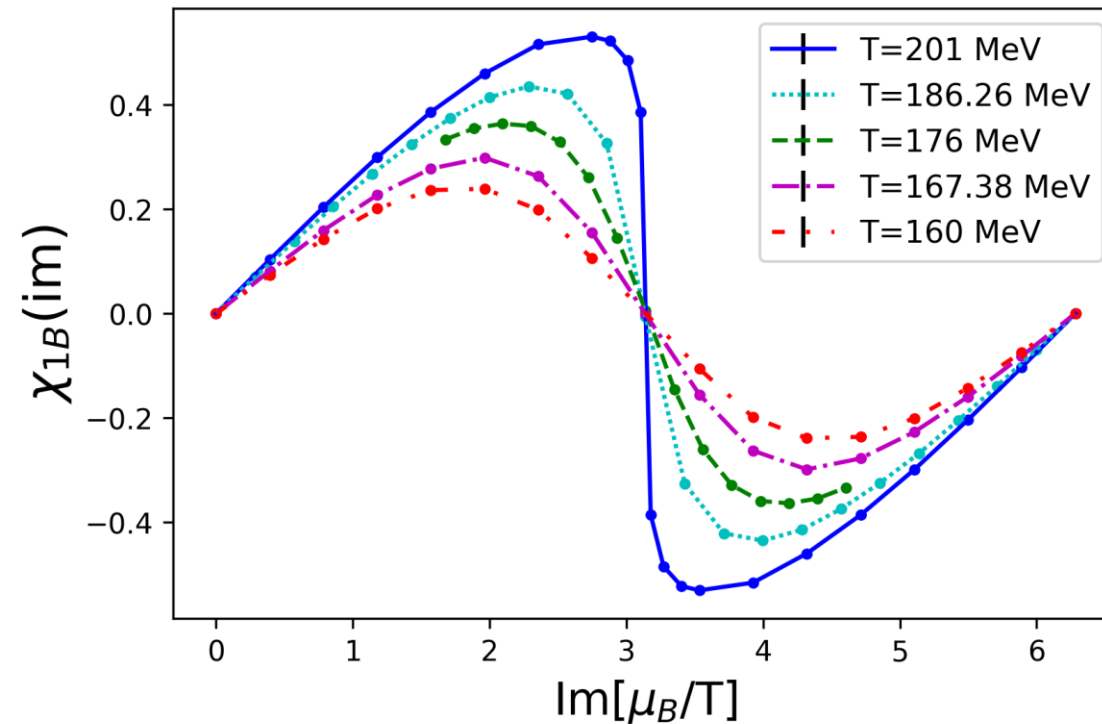
# Roberge-Weiss transition scaling

Lattice size:  $24^3 \times 4$

$T_{RW}=201$  MeV

[Jishnu Goswami et al, PoS CORFU2018 (2019) 162]

Cumulant measurements for  $Nt = 4$



[Simran Singh, Lattice 2021, Tuesday at 6.45 EST]

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Symmetry of the transition:  $Z(2)$

We set:

$$h = \frac{1}{h_0} \frac{\text{Im}(\frac{\mu_B}{T}) - \pi}{\pi} \quad t = \frac{1}{t_0} \frac{T - T_{RW}}{T_{RW}}$$

Temperature scaling:

$$\text{Re}[\mu_B/T] = \pi \left( \frac{z_0}{|z_c|} \right)^{\beta\delta} \left| \frac{T_{RW} - T}{T_{RW}} \right|^{\beta\delta}$$

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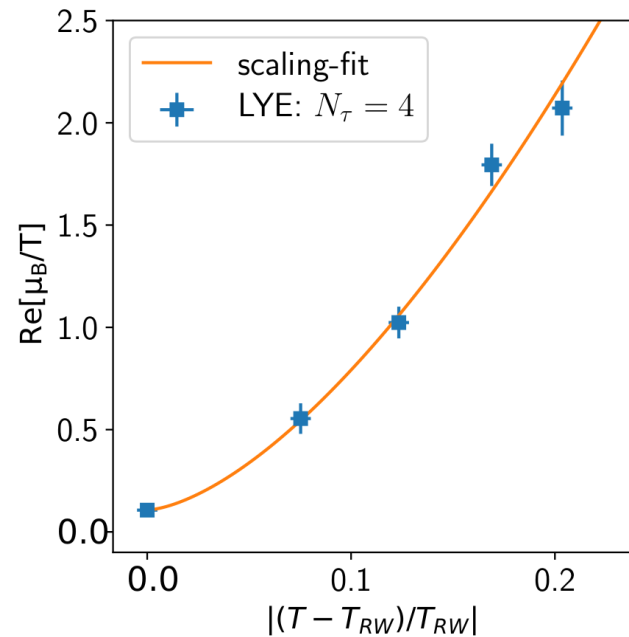
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Fit  $f(x) = a \cdot t^{\beta\delta} + b$

Known 3D Ising critical exp:

$$\beta = 0.3264$$

$$\delta = 4.789$$

$$z_c = 2.452 \pm 0.025$$

[Andrew Connelly et al, arXiv:2006.12541v2]

**Preliminary result:**

$$\chi_{test}^2 \approx 0.8$$

$$z_0 = 9.28 \pm 0.26$$



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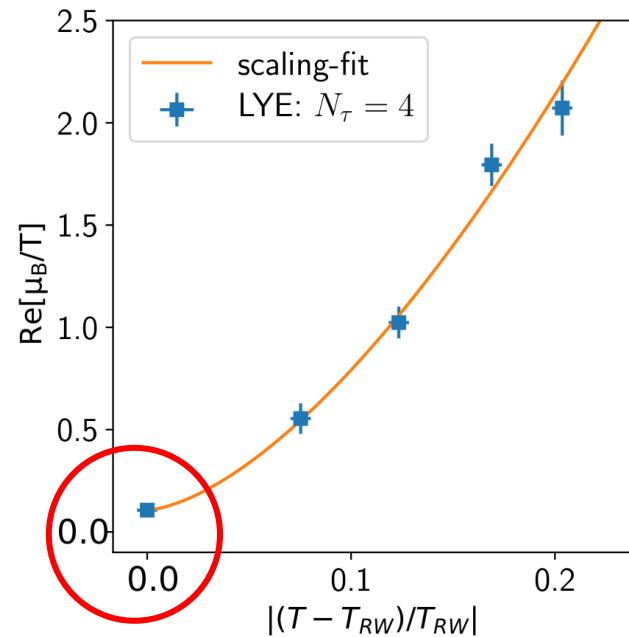
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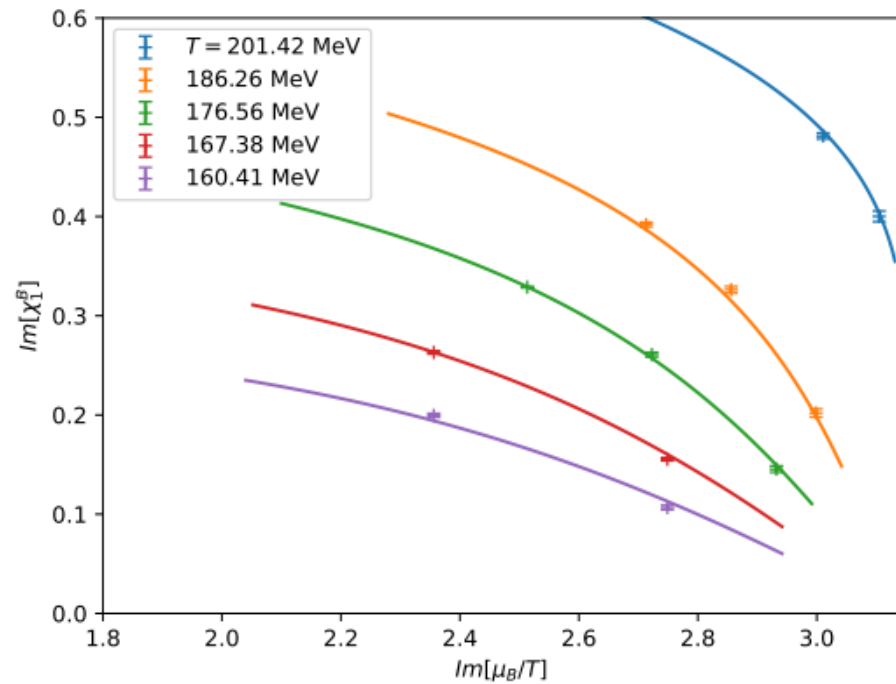
$$\chi_{test}^2 \approx 0.8$$

$$z_0 = 9.28 \pm 0.26$$

Fit parameters **slightly** correlated;  
Non zero singularity in  $t=0$ .

# Magnetic EoS as alternative method to obtain $z_0$

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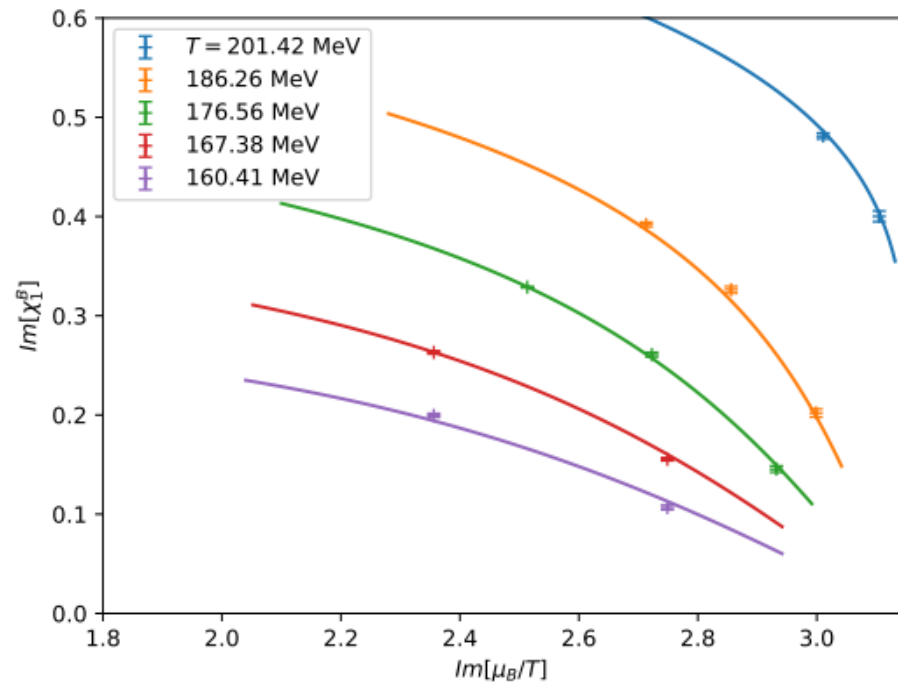


$$f_G(z) = \frac{M(t, h) - M_{reg}(t, h)}{h^{\frac{1}{\delta}}}$$

with  $M_{reg}(t, h) = a_2 \cdot t \cdot h$

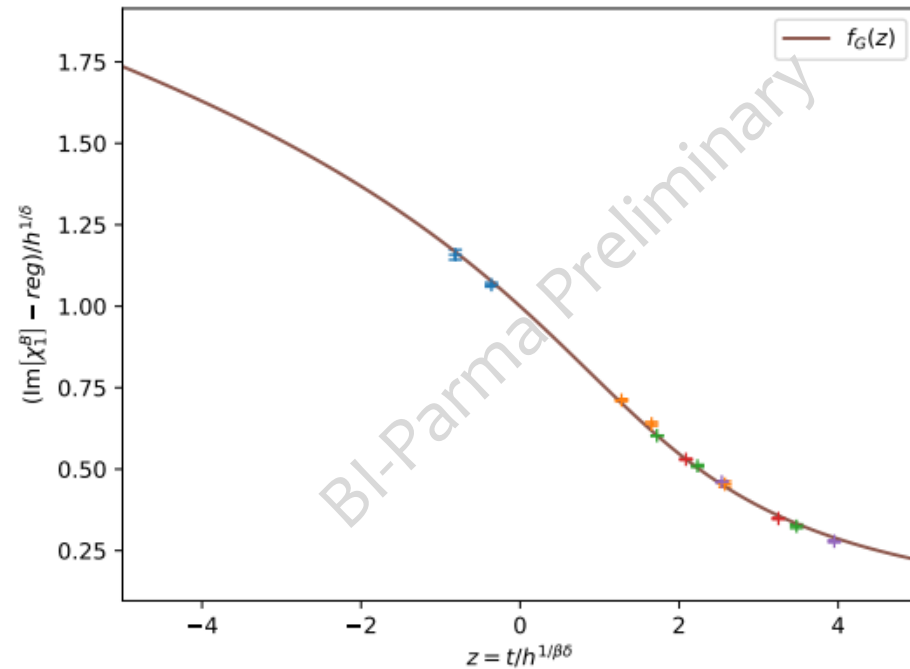
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Preliminary

$$z_0 = 5.316 \pm 0.017$$

$$\chi_{test}^2 \approx 6$$

LYE singularity method:

$$(z_0 = 9.28 \pm 0.26)$$

Highly correlated fit parameters!

# Chiral transition scaling

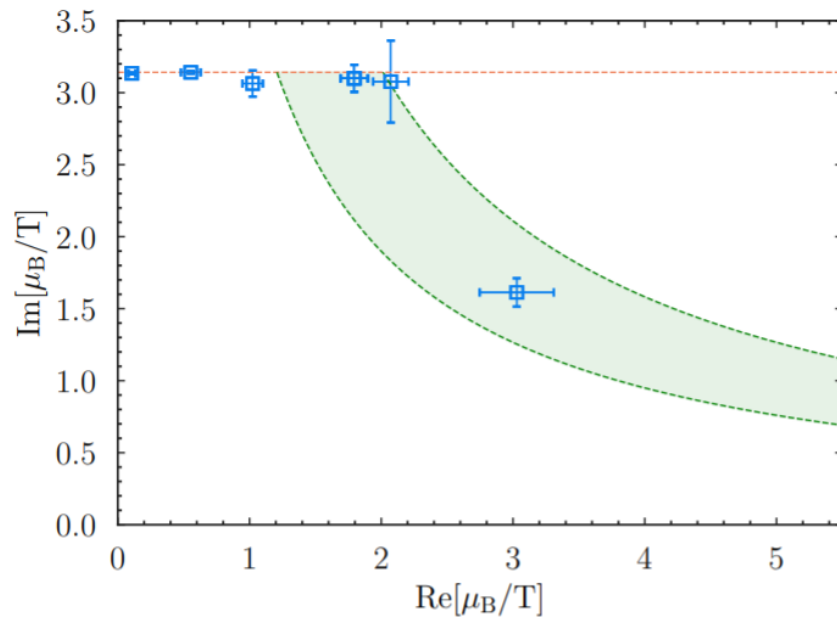
Lattice  $36^3 \times 6$ ,  $T_c = 147$  MeV

Symmetry of the transition:  $O(2)$

$$t = \frac{1}{t_0} \left[ \frac{T - T_c}{T_c} - \kappa_2^B \left( \frac{\mu_B}{T} \right)^2 \right] \quad h = \frac{1}{h_0} \left( \frac{m_l}{m_s} \right)$$

We find the scaling:  $\frac{\mu_B}{T} = \left[ \frac{1}{\kappa_2^B} \left( \frac{T - T_c}{T_c} - \frac{z_c}{z_0} \left( \frac{m_l}{m_s} \right)^{1/\beta\delta} \right) \right]^{1/2}$

[S. Mukherjee & V. Skokov, arxiv 1909.04639v2]



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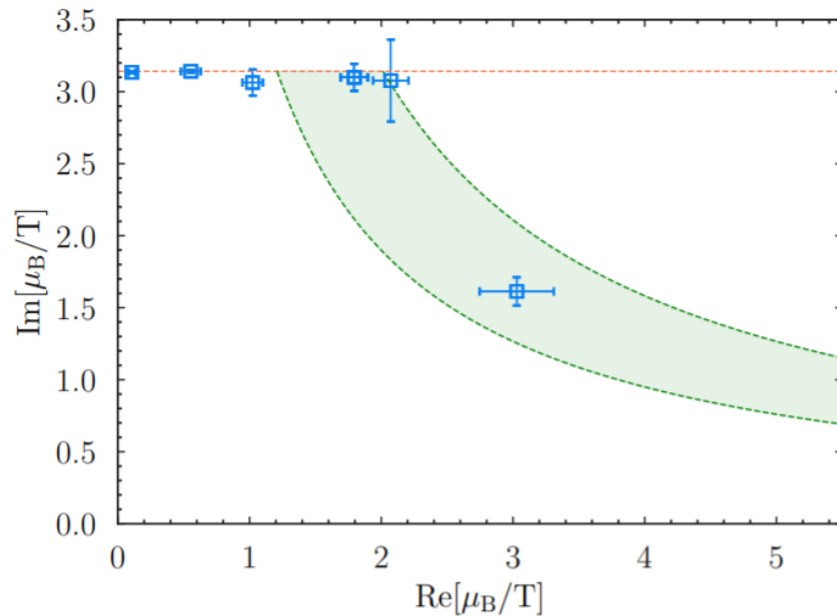
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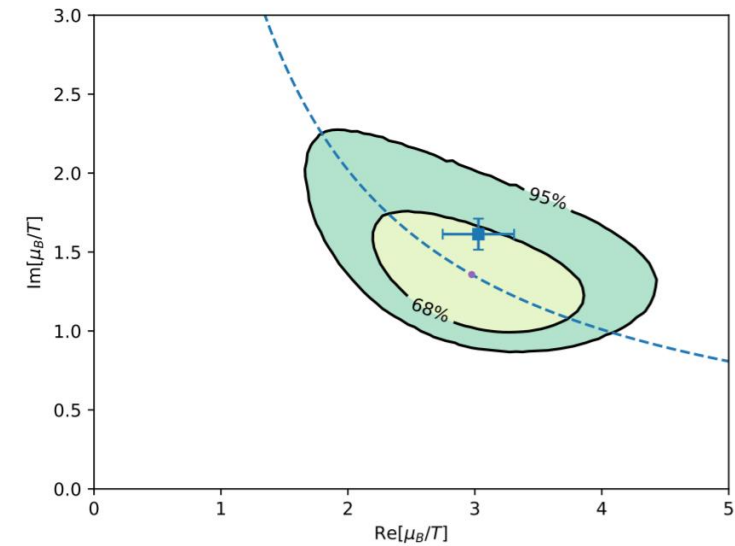
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How do we know that is a chiral LYE?

Parameter free prediction from HotQCD collaboration

[HotQCD private communications]



They are in good agreement

# Critical endpoint scaling: an outlook

Hypothetical Scenario

Symmetry of the transition:  $Z(2)$   
Unknown map to universal theory

Linear map:

$$t = \alpha_t(T - T_{cep}) + \beta_t(\mu_B - \mu_B^{cep})$$

$$h = \alpha_h(T - T_{cep}) + \beta_h(\mu_B - \mu_B^{cep})$$

[Gokce Basar, arXiv:2105.08080]

Scaling law:

$$\mu_{LY}(T) \approx c_1(T - T_{cep}) + ic_2|z_c|^{-\beta\delta}(T - T_{cep})$$

Many parameters are unknown!

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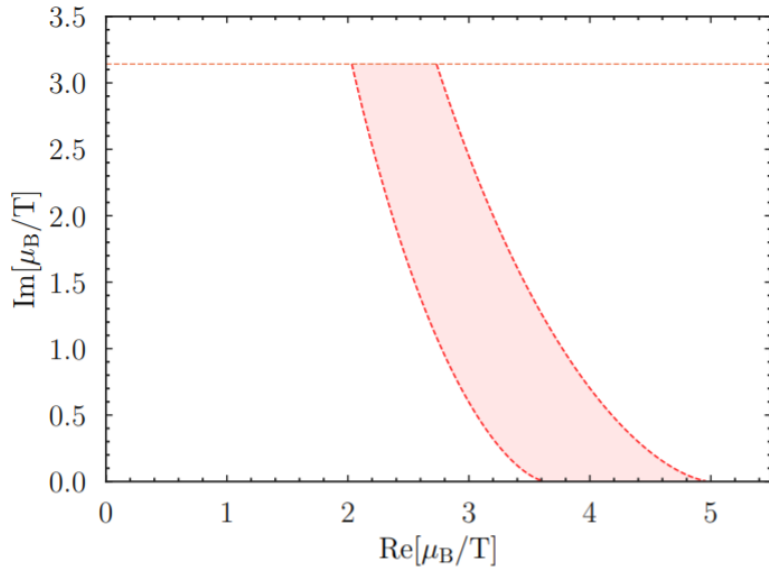
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Scaling law:

$$\mu_{LY}(T) \approx c_1(T - T_{cep}) + i c_2 |z_c|^{-\beta\delta} (T - T_{cep})$$

Many parameters are unknown!

Possible scenario:

$$T_{cep} = T_{pc} (1 - \kappa_2^B (\mu_{cep}/T_{pc})^2)$$

$$T_{pc} = 156.5 \text{ MeV}$$

$$c_1 = 0.024 \quad (\text{slope of the transition line at the critical point})$$

$$c_2 |z_c|^{-\beta\delta} = 0.5 \quad (\text{depends on the relative angle between the } h \text{ and } t \text{ axes})$$

G. Basar used this scaling to identify the **critical point in the Gross-Neveu model**.

In principle it should work for QCD as well, if one is able to get enough **LYE in the red region**.

# Summary & conclusions

Imaginary chemical potential and multi-points Padè approximants;  
LYE and scaling functions introduced;  
Correct scaling of our LYE in the Roberge-Weiss region;  
One LYE in the chiral region in good agreement with the prediction;  
Critical endpoint region was introduced.

## Future perspectives

Get more points on the chiral region;  
Possibly start to study the critical end point region.



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Thanks for the attention

# Backup

