# Decay amplifudes to three hadrons from lattice QCD

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In collaboration with: M. T. Hansen and S. R. Sharpe

## Decay amplitudes to three hadrons from finite-volume matrix elements

#### Maxwell T. Hansen $^1$ , Fernando Romero-López $^2$ , and Stephen R. Sharpe $^3$

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#### DESCLAIMET

This talk is based on the RFT three-particle formalism [Hansen, Sharpe]
Other approaches are FVU [Döring, Mai] and NREFT [Hammer, Pang, Rusetsky]

Analytic Expansions of Two- and Three-Particle Excited-State Energies	Dr Dorota Maria Grabowska	
	19:15 - 19:30	
Three-particle quantization condition for nondegenerate particles	Stephen Sharpe	
	19:30 - 19:45	
Three-hadron s- and d-wave interactions from lattice QCD	Dr Andrew Hanlon	
	19:45 - 20:00	
Parameters of the a1(1260) resonance from lattice QCD	Maxim Mai	
	20:00 - 20:15	
Three pion interactions from the lattice	Ruairí Brett	
	20:15 - 20:30	
Three-particle finite-volume formalism for \$\pi^+ \pi^+ K^+\$ and related systems	Tyler Blanton	
	20:30 - 20:45	
Infinite volume, three-body scattering formalisms in the presence of bound states	Sebastian Dawid	
	20:45 - 21:00	

- We have a (relativistic) three-particle formalism for three-particle scattering [Hansen, Sharpe]
  - Recent extensions: nonidentical pions, nondegenerate particles. [Hansen, FRL, Sharpe] + [Blanton, Sharpe]
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$A_g^C$	$A_g^N$	Ref.	
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- O In the two-particle formalism this is the well-known Lellouch-Lüscher formalism
  - lacksquare It has been applied successfully in  $K o\pi\pi$  processes [RBC/UKQCD, see plenary by C. Kelly]

#### O Finite-volume correlator:

$$C_L^{\mathsf{M}}(E, \mathbf{P}) = \int_{-\infty}^{\infty} dx_0 \int_L d^3x \, e^{i(Ex^0 - \mathbf{P} \cdot \mathbf{x})} \langle 0 | \mathrm{T}\sigma(x)\sigma^{\dagger}(0) | 0 \rangle \,,$$

#### **O** Finite-volume correlator:

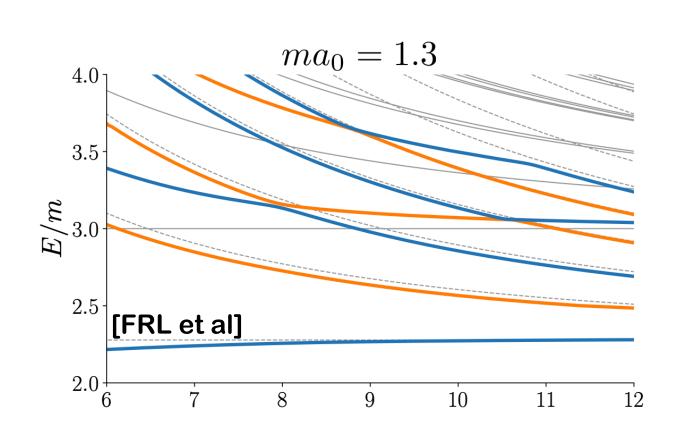
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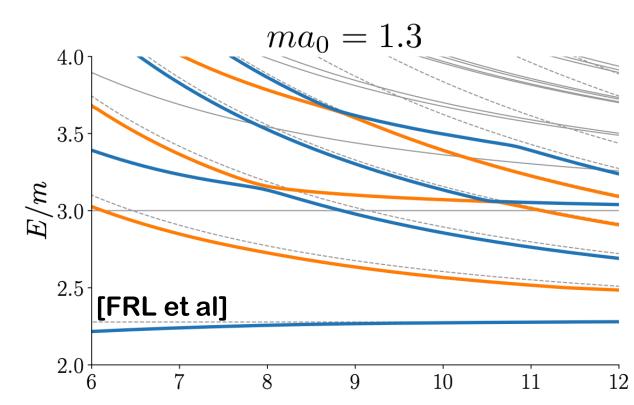
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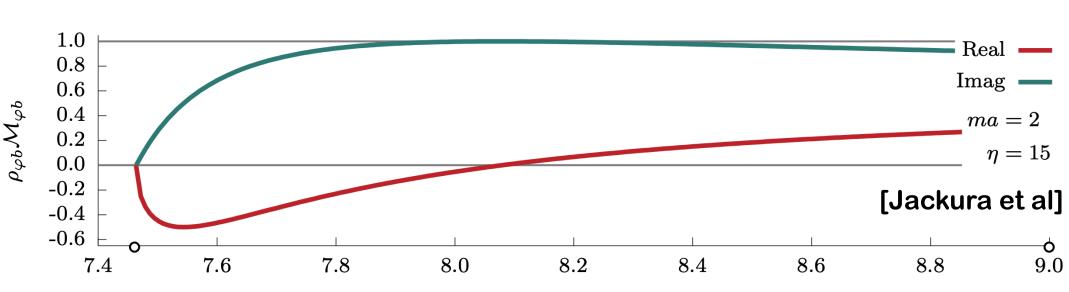
O Step 1: three-particle quantization condition:

$$\det(F_3^{-1} + \mathcal{K}_{df,3}) = 0.$$

O Step 2: three-particle scattering amplitude:

$$\mathcal{M}_{3,L}^{(u,u)} = \mathcal{D}^{(u,u)} + \mathcal{L}_{L}^{(u)} \frac{1}{1 + \mathcal{K}_{df,3} F_3} \mathcal{K}_{df,3} \mathcal{R}_{L}^{(u)},$$





O Consider real scalars, "pion" and "kaon"

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$$\mathcal{H}_W(x) = c_W \frac{K(x)\phi(x)^3}{3!}$$

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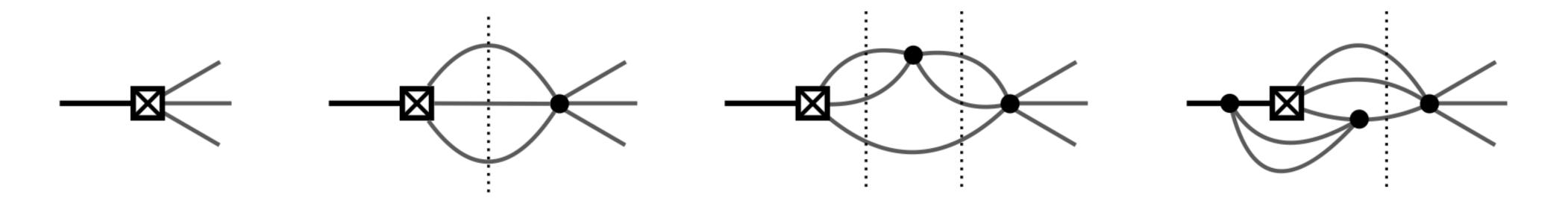
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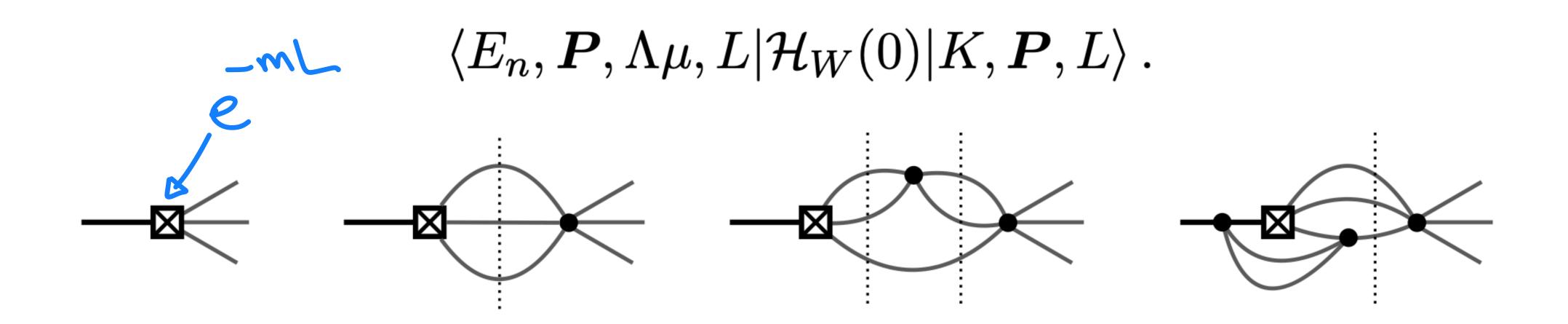
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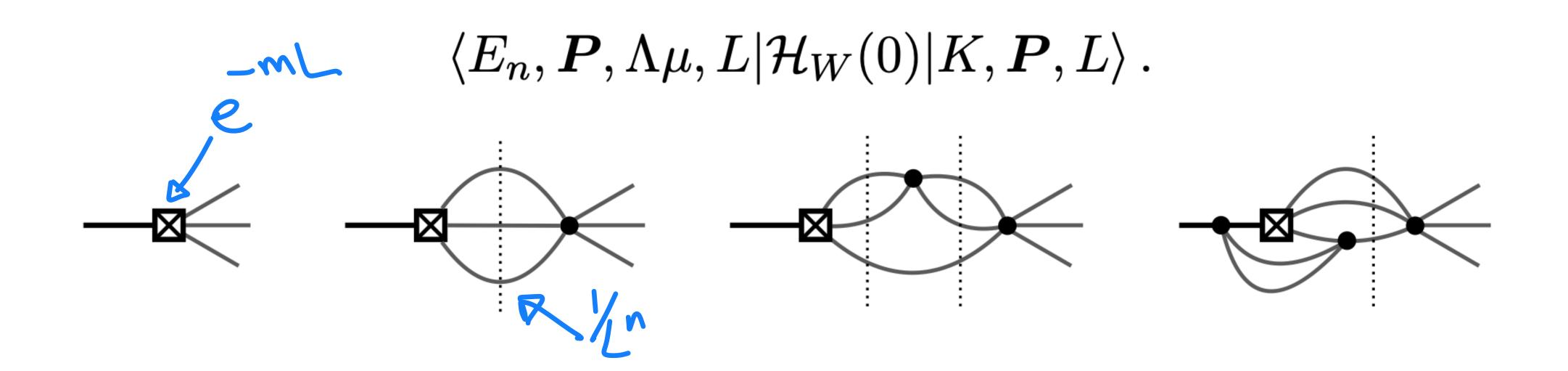
O Tune the box size such that:  $E_{3\phi}(L) \simeq M_K$ 

$$\langle E_n, \mathbf{P}, \Lambda \mu, L | \mathcal{H}_W(0) | K, \mathbf{P}, L \rangle$$
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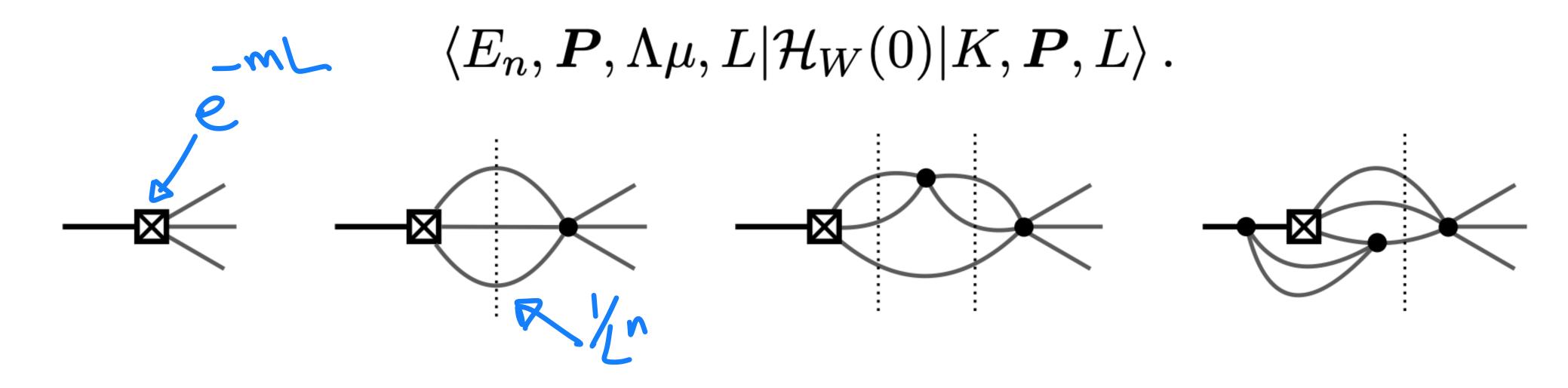
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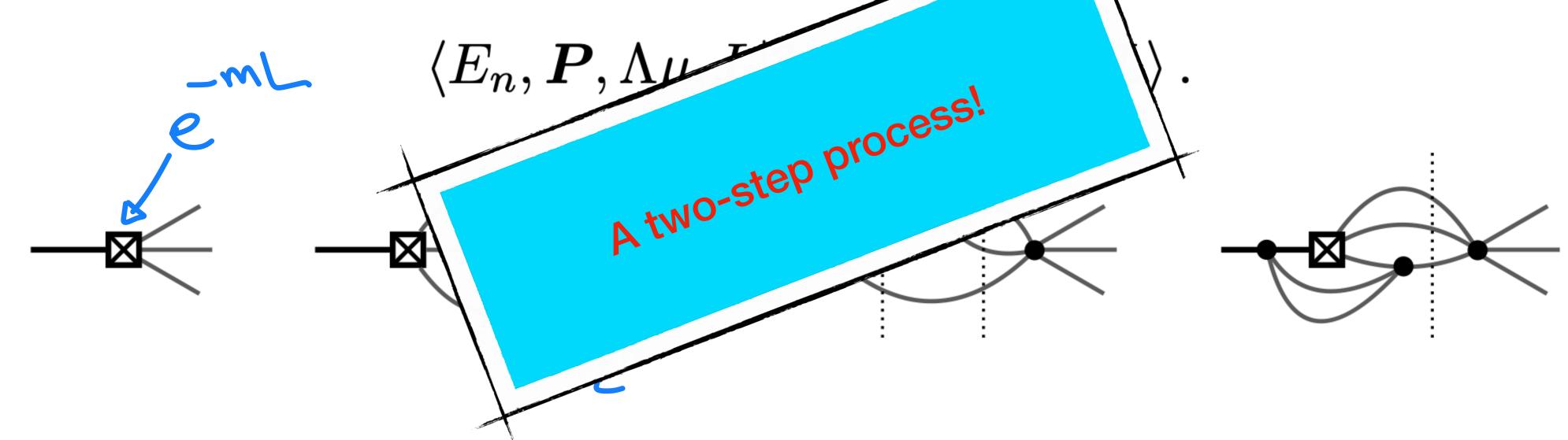
Trom the lattice, one can get the one-to-three finite-volume matrix element:



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O Calculate the residue of the quantization condition at the solution with  $E=E_{3\phi}(L)\simeq M_K$ 

$$\mathcal{R}_{\Lambda\mu}(E_n^{\Lambda}, \mathbf{P}, L) = \lim_{P_4 \to iE_n^{\Lambda}} -(E_n^{\Lambda} + iP_4) \, \mathbb{P}_{\Lambda\mu} \cdot \frac{1}{F_3^{-1} + \mathcal{K}_{df,3}} \cdot \mathbb{P}_{\Lambda\mu}$$

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O Use the measured matrix element to compute an intermediate quantity

$$\sqrt{2E_K({m P})}L^3\langle E_n,{m P},\Lambda\mu,L|{\cal H}_W(0)|K,{m P},L
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O  $A_{K3\pi}^{\mathrm{PV}}$  is however scheme dependent! —— Meed step 2!

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$$A^{\mathsf{iso}} = \sum_{n=0}^{\infty} \Delta^n A^{\mathsf{iso,n}} \,, \qquad \qquad \Delta = rac{m_K^2 - 9 m_\pi^2}{9 m_\pi^2} \,. \qquad \qquad \Delta_i = rac{s_i - 4 m_\pi^2}{9 m_\pi^2} \,,$$

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O Need as many matrix elements as parameters!

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One needs to symmetrize and take the infinite-volume limit:

$$T_{K3\pi}(\mathbf{k}, \widehat{\mathbf{a}}^*) \equiv S \{ T_{K3\pi}(\mathbf{k})_{\ell m} \} ,$$
  
=  $T_{K3\pi}^{(u)}(\mathbf{k}, \widehat{\mathbf{a}}^*) + T_{K3\pi}^{(u)}(\mathbf{a}, \widehat{\mathbf{b}}^*) + T_{K3\pi}^{(u)}(\mathbf{b}, \widehat{\mathbf{k}}^*) ,$ 

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Similar to integral equations in three-to-three scattering! [Jackura et al, Hansen et al.]

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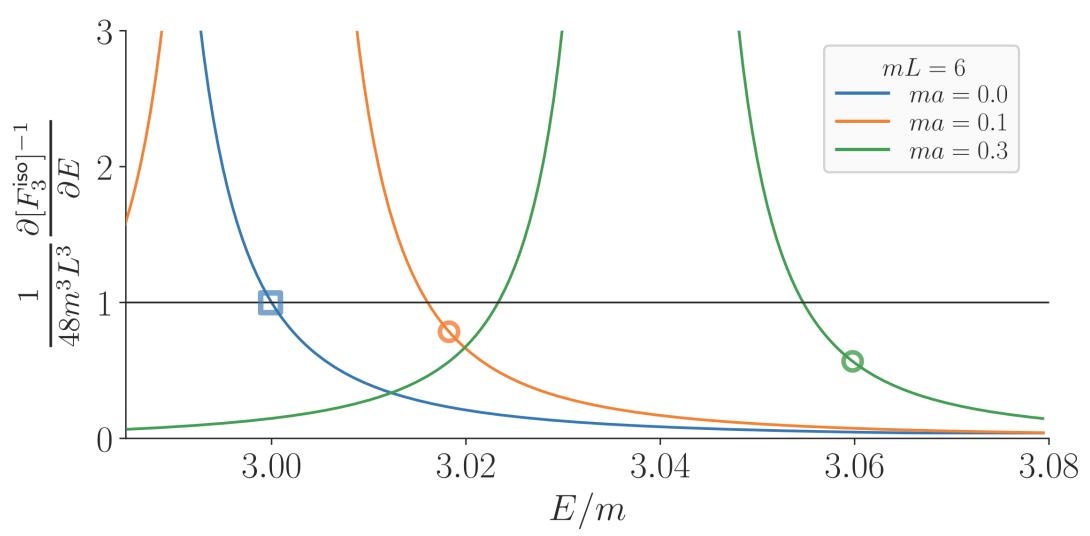
"Lettouch-Lüscher" Ehree-particle factor

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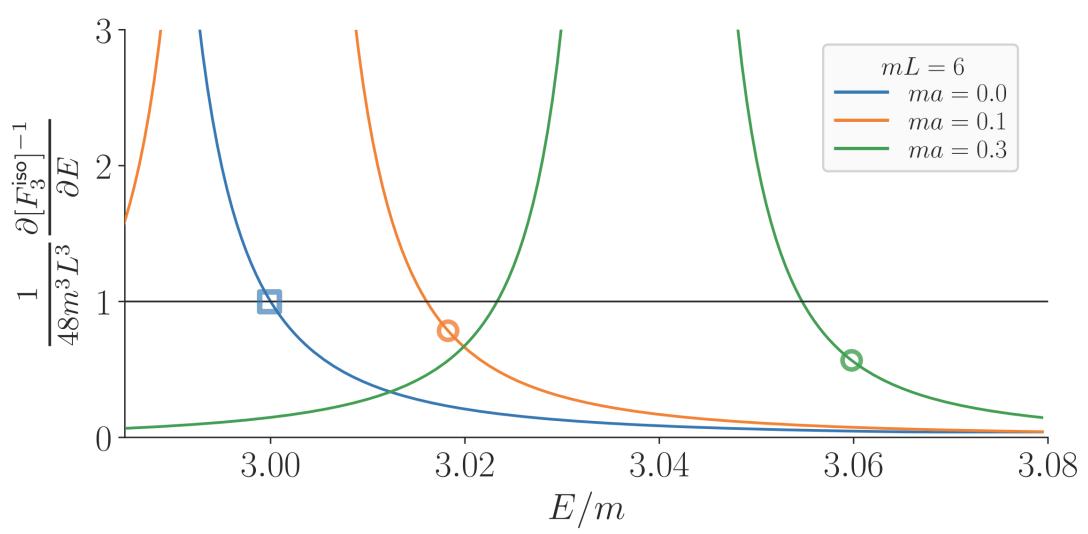


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Similar approximation to [Müller, Rusetsky]

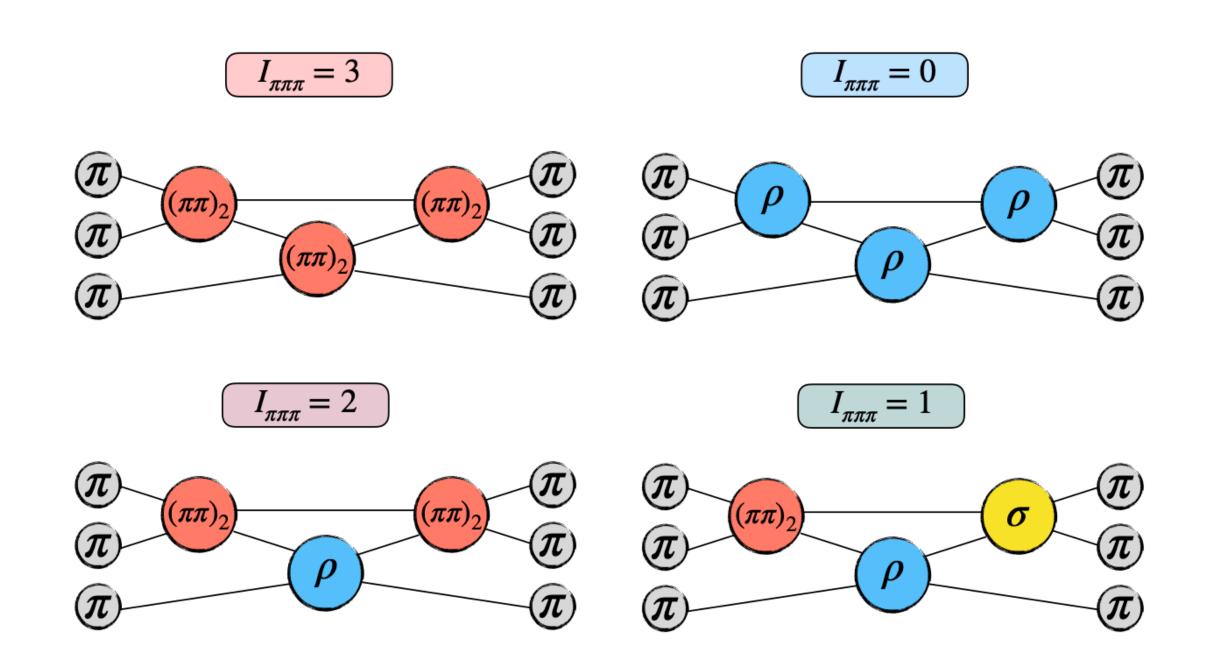
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#### Formalism for three pions

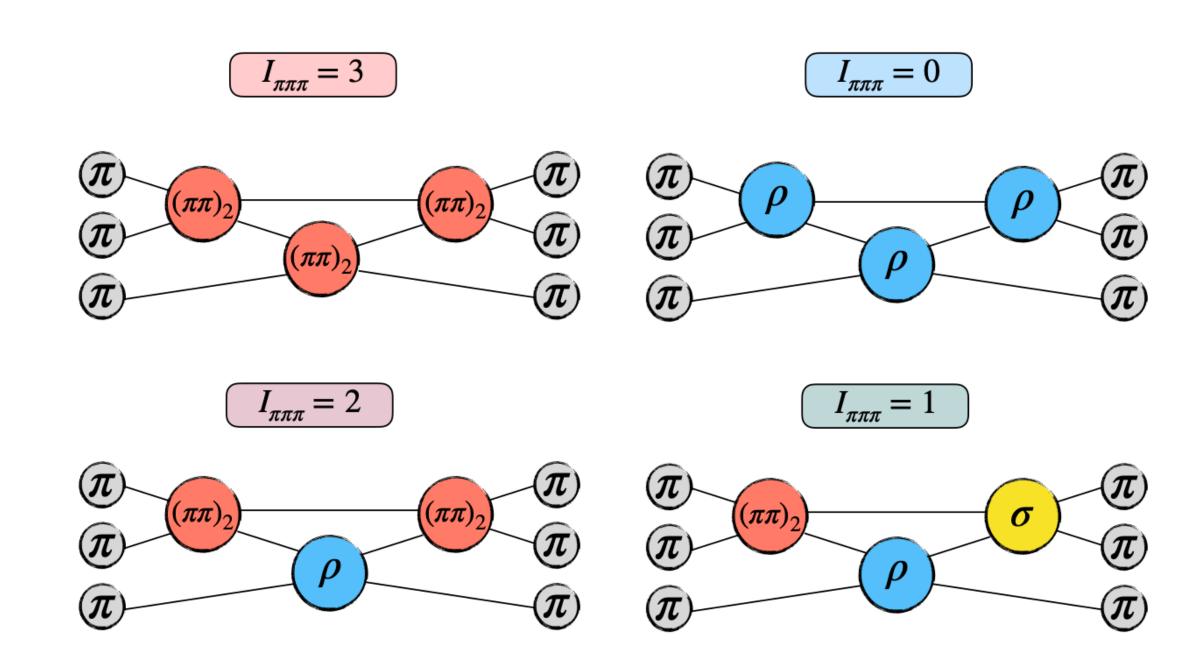
# Formalism for three pions

Use the three-pion finite volume formalism from [Hansen, FRL, Sharpe, arXiv:2003.10974]



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$$\det[1 - \mathbf{K}_{\mathrm{df},3}^{[I]}(E^{\star}) \, \mathbf{F}_{3}^{[I]}(E, \mathbf{P}, L)] = 0$$

Expressions for decays are formally very similar with additional flavor index:

$$\sqrt{2E_K(\boldsymbol{P})}L^3\langle E_n^{\Lambda,[I]},\boldsymbol{P},I,I_3,\Lambda\mu,L|\mathcal{H}_W(0)|K,\boldsymbol{P},L\rangle=\mathbf{v}^{\dagger}\mathbf{A}_{K3\pi}^{\mathrm{PV},[I]}.$$

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This can be used to treat phenomenologically relevant processes:

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$$\begin{split} \sqrt{2E_K(\boldsymbol{P})}L^3\langle E_n^{\Lambda,[I]},\boldsymbol{P},I,I_3,\Lambda\mu,L|\mathcal{H}_W(0)|K,\boldsymbol{P},L\rangle &= \mathbf{v}^\dagger\mathbf{A}_{K3\pi}^{\mathrm{PV},[I]}\,.\\ \mathbf{R}_{\Lambda\mu}^{[I,I_3]}(E_n^{\Lambda,[I]},\boldsymbol{P},L) &= \lim_{P_4\to iE_n^{\Lambda,[I]}} -(E_n^{\Lambda,[I]}+iP_4)\mathbb{P}_{\Lambda\mu}^{[I,I_3]} \frac{(-i)}{1/\mathbf{F}_3^{[I]}-\mathbf{K}_{\mathrm{df},3}^{[I]}}\mathbb{P}_{\Lambda\mu}^{[I,I_3]} &\equiv \mathbf{v}\,\mathbf{v}^\dagger\,. \end{split}$$

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isospin-breaking effect

### Summary and Outlook

- We have derived a formalism for three-particle decays.
- O Formalism for identical particles similar to [Müller, Rusetsky]
- C Extension to generic three-pion systems based on previous three-pion quantization condition
- O Explicit expressions to treat  $K o \pi\pi\pi$ ,  $\gamma^* o \pi\pi\pi$ ,  $\eta o \pi\pi\pi$ .
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