



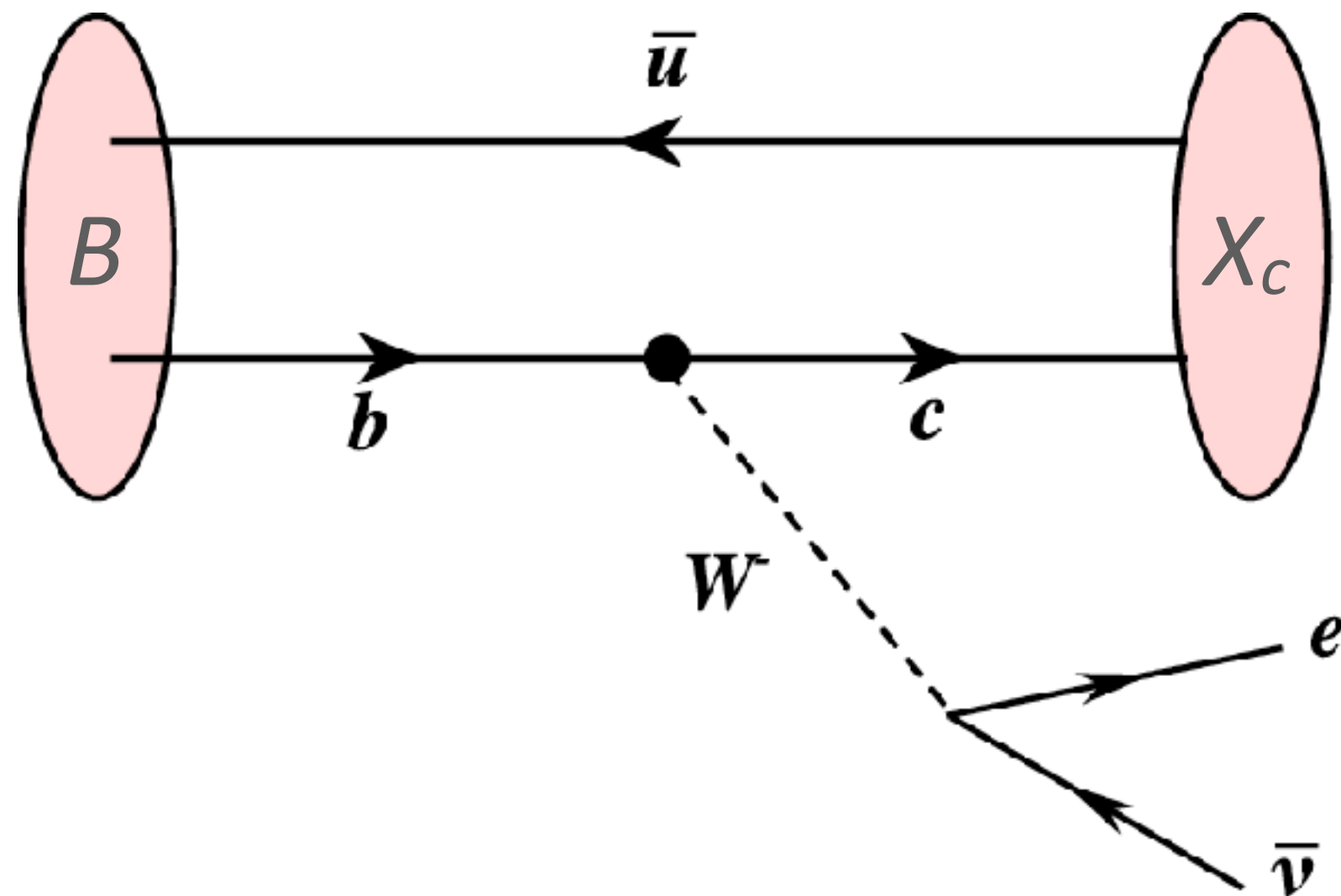
# Composition of the inclusive semi-leptonic decay of B meson

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Jul 28, 2021

# Semi-leptonic B decay



**Exclusive:** specific final states ( $D, D^*, \dots$ )

**Inclusive:** sum over all final states,  
can be computed using PT (or OPE)

A new method to compute the “sum” in LQCD

Gambino and SH, arXiv:2005.13730

from the forward-Compton amplitude.

$$\langle B(\mathbf{0}) | \tilde{J}_\mu^\dagger(-\mathbf{q}; t) \left| \tilde{J}_\nu(\mathbf{q}; 0) | B(\mathbf{0}) \right\rangle$$

all possible states contribute

# Composition of inclusive decays (exp't):

Bernlochner @ ICHEP 2012

Charm state $X_c$	$\mathcal{B}(B^+ \rightarrow X_c \ell^+ \nu)$	
$D$	$(2.31 \pm 0.09) \%$	
$D^*$	$(5.63 \pm 0.18) \%$	
$\sum D^{(*)}$	$(7.94 \pm 0.20) \%$	
$D_0^* \rightarrow D \pi$	$(0.41 \pm 0.08) \%$	} broad states $(0.86 \pm 0.12) \%$
$D_1^* \rightarrow D^* \pi$	$(0.45 \pm 0.09) \%$	
$D_1 \rightarrow D^* \pi$	$(0.43 \pm 0.03) \%$	} narrow states $(0.84 \pm 0.04) \%$
$D_2^* \rightarrow D^{(*)} \pi$	$(0.41 \pm 0.03) \%$	
$\sum D^{**} \rightarrow D^* \pi$	$(1.70 \pm 0.12) \%$	
$D \pi$	$(0.66 \pm 0.08) \%$	
$D^* \pi$	$(0.87 \pm 0.10) \%$	
$\sum D^* \pi$	$(1.53 \pm 0.13) \%$	
$\sum D^{(*)} + \sum D^* \pi$	$(9.47 \pm 0.24) \%$	
$\sum D^{(*)} + \sum D^{**} \rightarrow D^{(*)} \pi$	$(9.64 \pm 0.23) \%$	
Inclusive $X_c$	$(10.92 \pm 0.16) \%$	

Sum of all identified states is less than **ALL**.

All values from [HFAG 2010]. For the values of  $B \rightarrow D \pi \ell \bar{\nu}_\ell$  and  $B \rightarrow D^* \pi \ell \bar{\nu}_\ell$  an uncertainty weighted average of both isospin modes was calculated assuming a 100% correlation between both values.

$\Rightarrow$  'Gap' of  $(1.45 \pm 0.29) \%$  emerges which is not accounted for

Can we understand why using lattice data? Need an access to both incl and excl.

# Inclusive rate

Differential decay rate:

$$d\Gamma \sim |V_{cb}|^2 l^{\mu\nu} W_{\mu\nu}$$

Structure function:

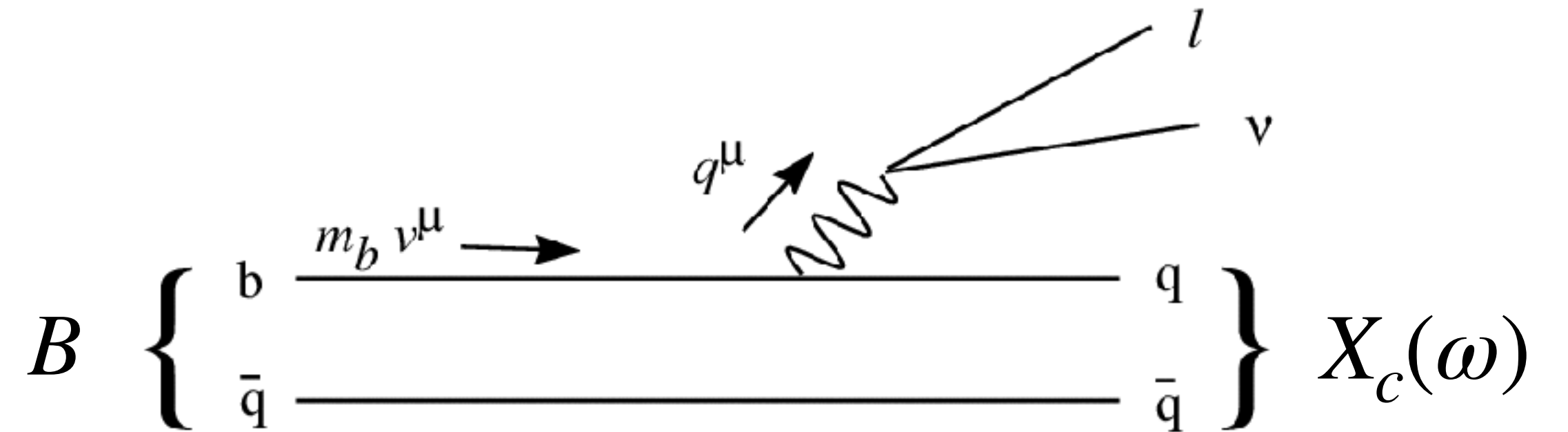
$$W_{\mu\nu} = \sum_X (2\pi)^2 \delta^4(p_B - q - p_X) \frac{1}{2M_B} \langle B(p_B) | J_\mu^\dagger(0) | X \rangle \langle X | J_\nu(0) | B(p_B) \rangle$$

$$\rightarrow \langle B(\mathbf{0}) | \tilde{J}_\mu^\dagger(-\mathbf{q}; t) \delta(\omega - \hat{H}) \tilde{J}_\nu(\mathbf{q}; 0) | B(\mathbf{0}) \rangle$$

Decay rate:

$$\Gamma \propto \int_0^{\mathbf{q}_{\max}^2} d\mathbf{q} \int_{\sqrt{m_D^2 + \mathbf{q}^2}}^{m_B - \sqrt{\mathbf{q}^2}} d\omega K(\omega; \mathbf{q}^2) \langle B(\mathbf{0}) | \tilde{J}^\dagger(-\mathbf{q}) \delta(\omega - \hat{H}) \tilde{J}(\mathbf{q}) | B(\mathbf{0}) \rangle$$

known kinematical factor



# Energy integral

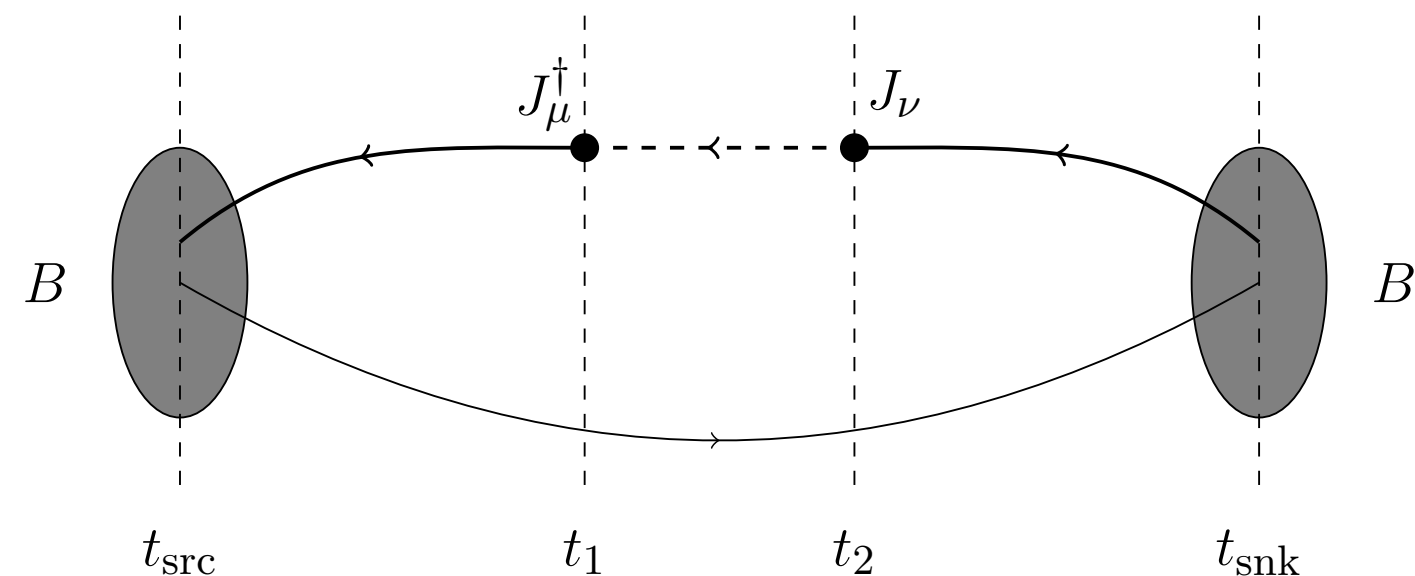
$$\Gamma \propto \int_0^{q_{\max}^2} d\mathbf{q} \int_{\sqrt{m_D^2 + \mathbf{q}^2}}^{m_B - \sqrt{\mathbf{q}^2}} d\omega K(\omega; \mathbf{q}^2) \langle B(\mathbf{0}) | \tilde{J}^\dagger(-\mathbf{q}) \delta(\omega - \hat{H}) \tilde{J}(\mathbf{q}) | B(\mathbf{0}) \rangle$$

$$= \langle B(\mathbf{0}) | \tilde{J}^\dagger(-\mathbf{q}) K(\hat{H}; \mathbf{q}^2) \tilde{J}(\mathbf{q}) | B(\mathbf{0}) \rangle$$



Lattice Compton amplitude:

$$\langle B(\mathbf{0}) | \tilde{J}_\mu^\dagger(-\mathbf{q}; t) \tilde{J}_\nu(\mathbf{q}; 0) | B(\mathbf{0}) \rangle \longrightarrow \langle B(\mathbf{0}) | \tilde{J}^\dagger(-\mathbf{q}) e^{-\hat{H}t} \tilde{J}(\mathbf{q}) | B(\mathbf{0}) \rangle$$



Approximation:

$$K(\hat{H}) = k_0 + k_1 e^{-\hat{H}} + k_2 e^{-2\hat{H}} + \dots + k_N e^{-k_N \hat{H}}$$

# Chebyshev approx:

Bailas, Ishikawa, SH, arXiv:2001.11779

(shifted) Chebyshev polynomials

$$T_0^*(x) = 1$$

$$T_1^*(x) = 2x - 1$$

$$T_2^*(x) = 8x^2 - 8x + 1$$

⋮

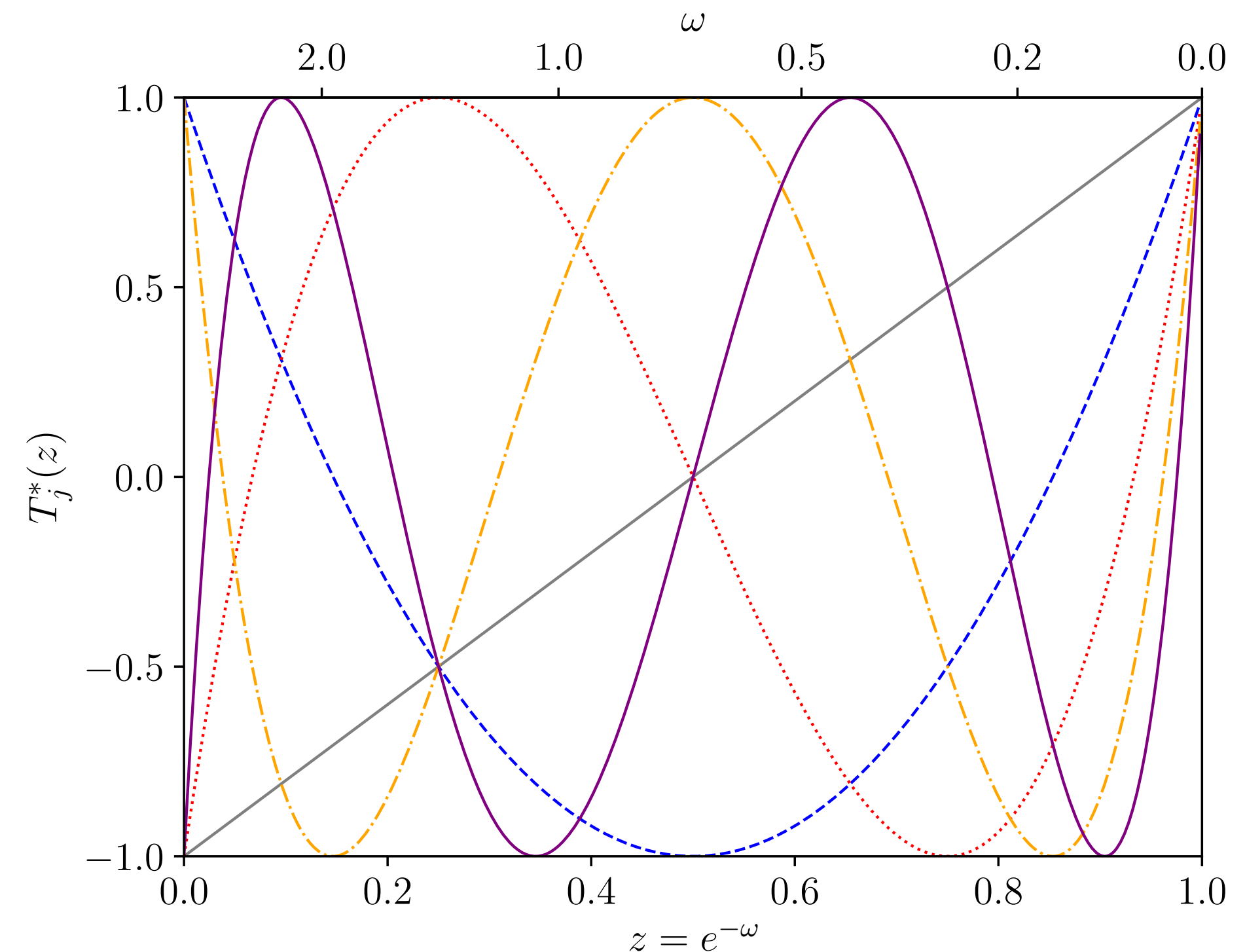
$$T_{j+1}^*(x) = 2(2x - 1)T_j^*(x) - T_{j-1}^*(x)$$

each term corresponds to the correlator, because  $x = e^{-\omega}$

“Best” approximation can be obtained with

$$c_j^* = \frac{2}{\pi} \int_0^\pi d\theta S \left( -\ln \frac{1 + \cos \theta}{2} \right) \cos(j\theta)$$

$$K(\omega) \simeq \frac{c_0}{2} + \sum_{j=1}^N c_j^* T_j^*(e^{-\omega})$$



“best” = maximal deviation is minimal



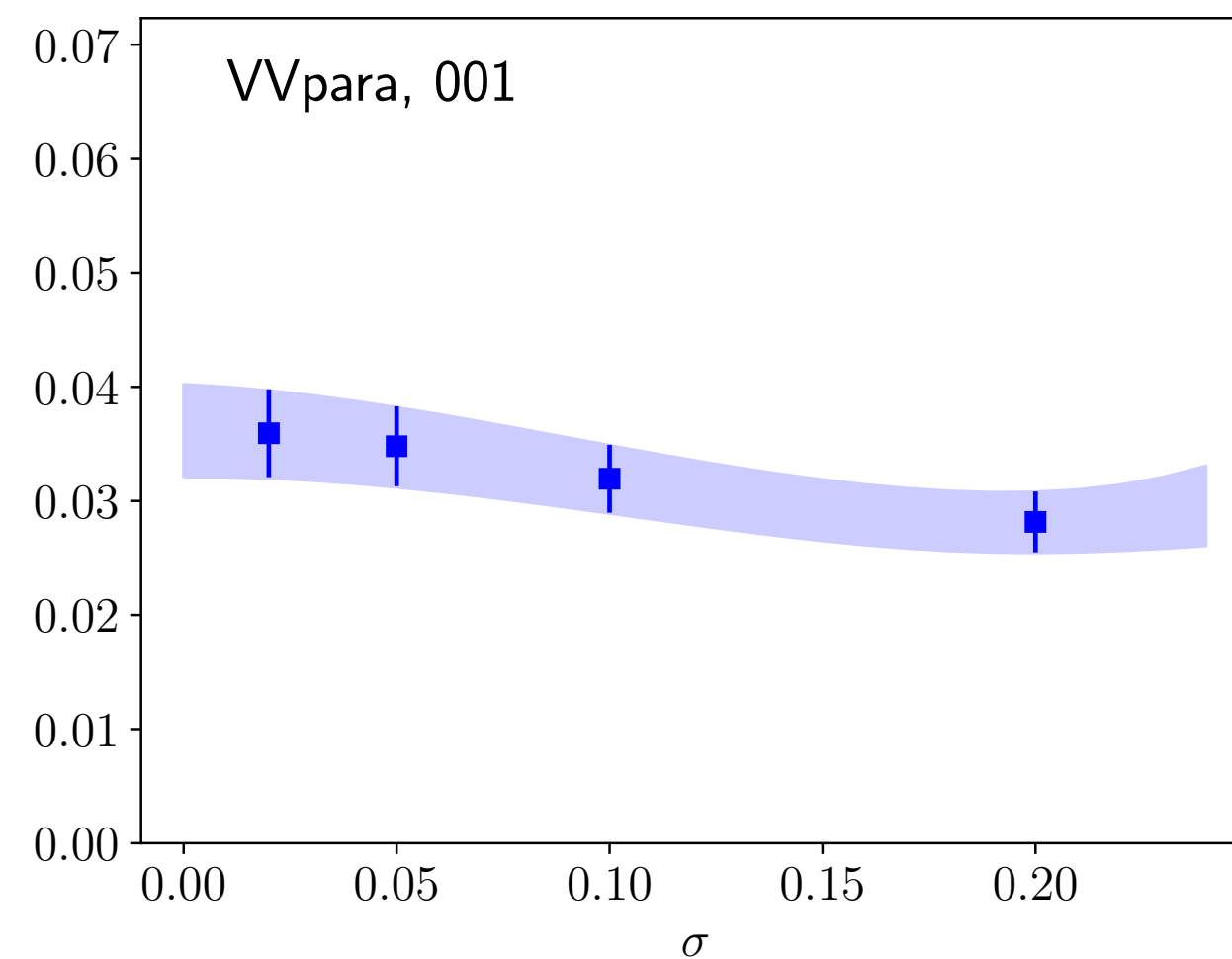
# Kernel to approximate

To implement the upper limit of integ

$$K(\omega) \sim e^{2\omega t_0} \underbrace{(m_B - \omega)^l}_{\text{kinematical factor}} \theta(m_B - |\mathbf{q}| - \omega)$$

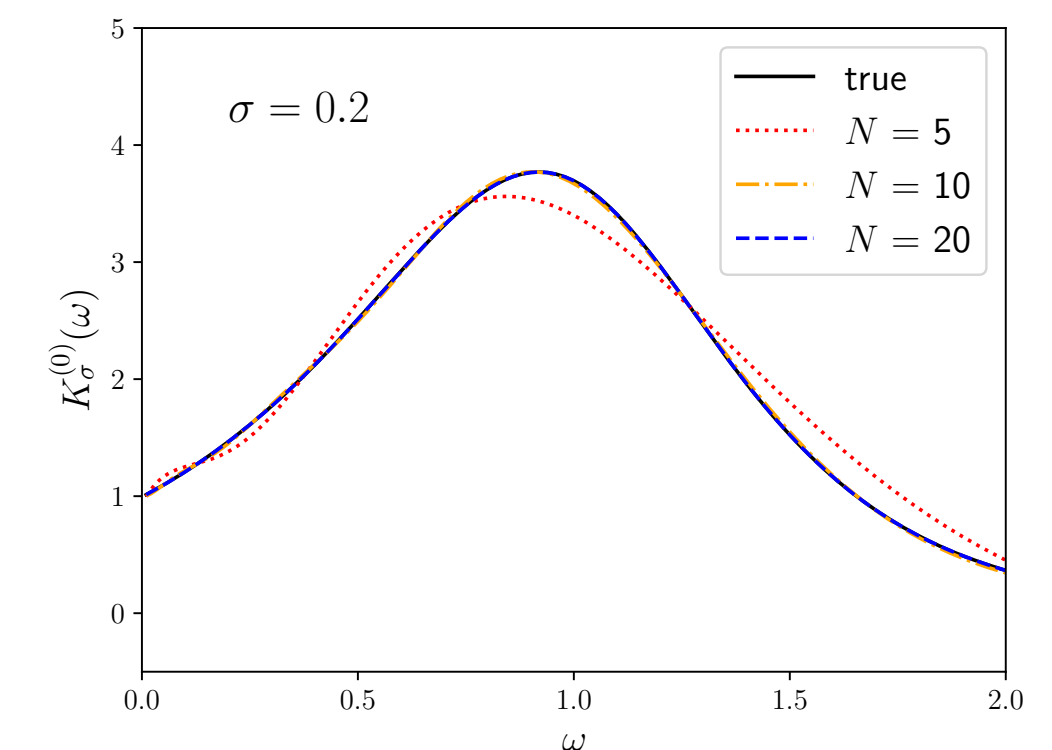
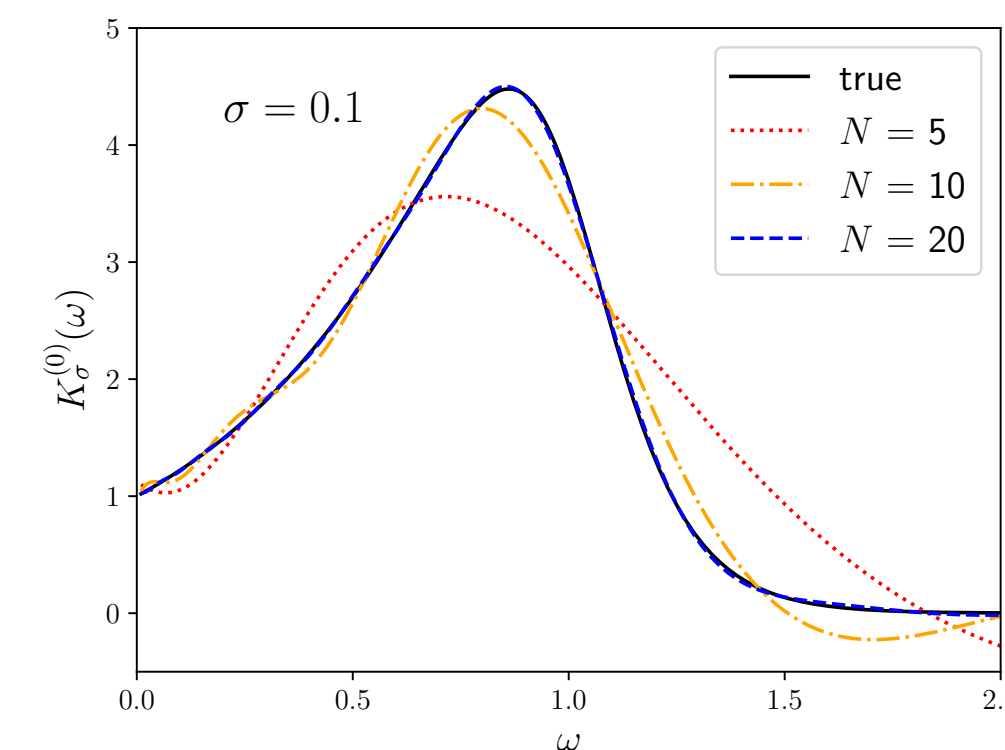
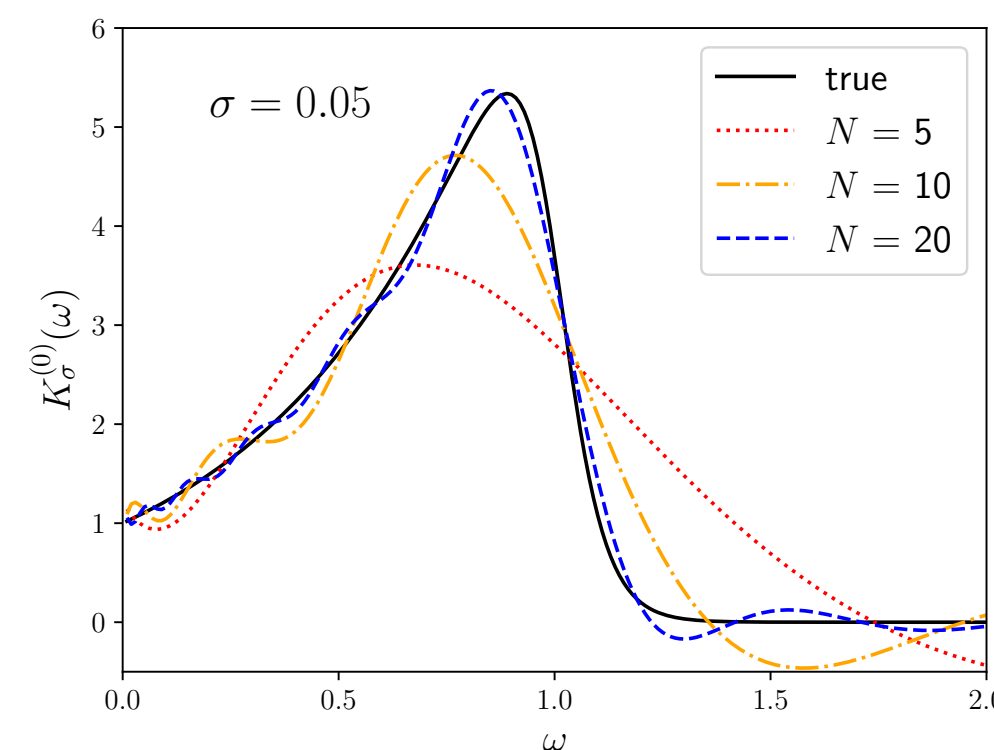
kinematical factor

Smear by “sigmoid” with a width  $\sigma$   
Need to take a limit of  $\sigma \rightarrow 0$



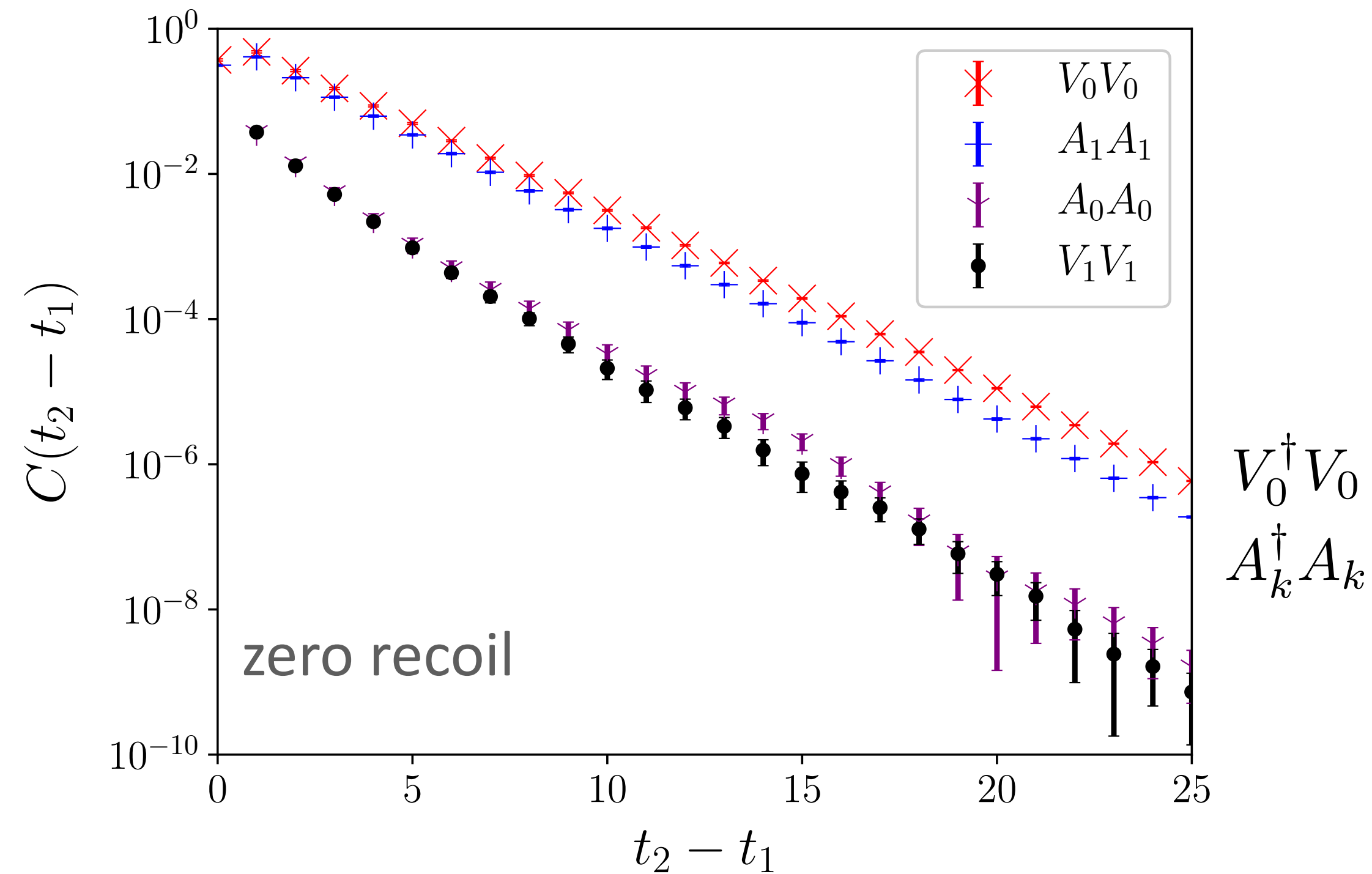
narrow

wide



# Compton amplitude (S-wave)

$$\langle B(\mathbf{0}) | \tilde{J}_\mu^\dagger(-\mathbf{q}; t) \tilde{J}_\nu(\mathbf{q}; 0) | B(\mathbf{0}) \rangle$$



Pilot lattice computation [JLQCD setup]

- On a lattice of  $48^3 \times 96$  at  $1/a = 3.6$  GeV
- Strange spectator quark
- physical charm quark mass
- (unphysically) light b quark  $\sim 2.7$  GeV
- 100 configs x 4 src

S-wave (D and D\*)

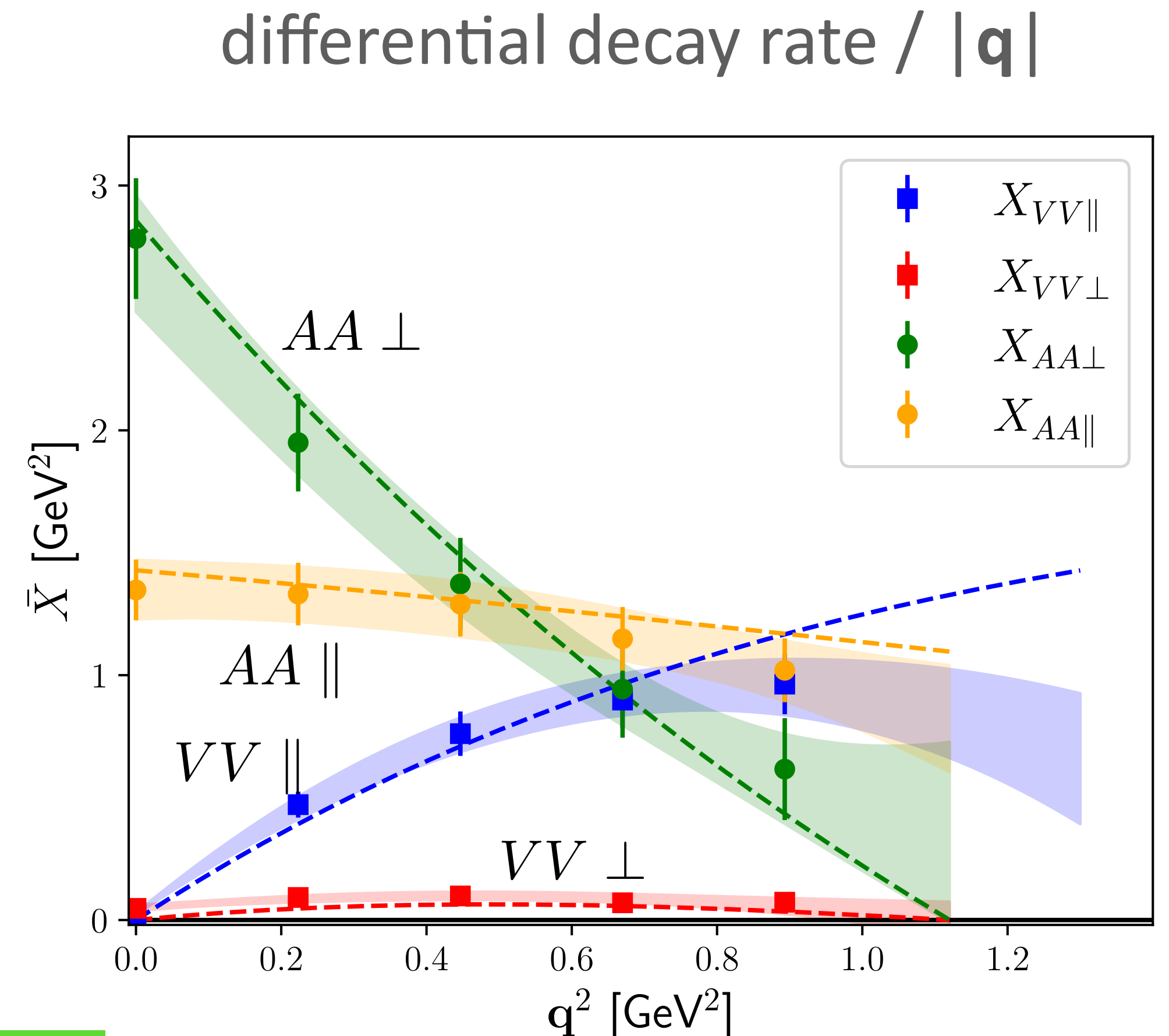
- Very well approximated by a single-exp  
= no sign of excited state contrib.
- No wave function overlap of excited states when  $m_b = m_c$  and zero recoil ( $V_0 V_0$ )



# Inclusive decay rate

- Breakdown to individual channels: VV and AA; parallel and perp with respect to the recoil momentum
- Compared to exclusive contributions estimated from  $B \rightarrow D^{(*)}$  form factors (dashed line), that are separately calculated.
- $VV_{||}$  dominated by  $B \rightarrow D$
- All others by  $B \rightarrow D^*$

Inclusive rate is dominated by the ground states, naturally because of the unphysical b quark mass.

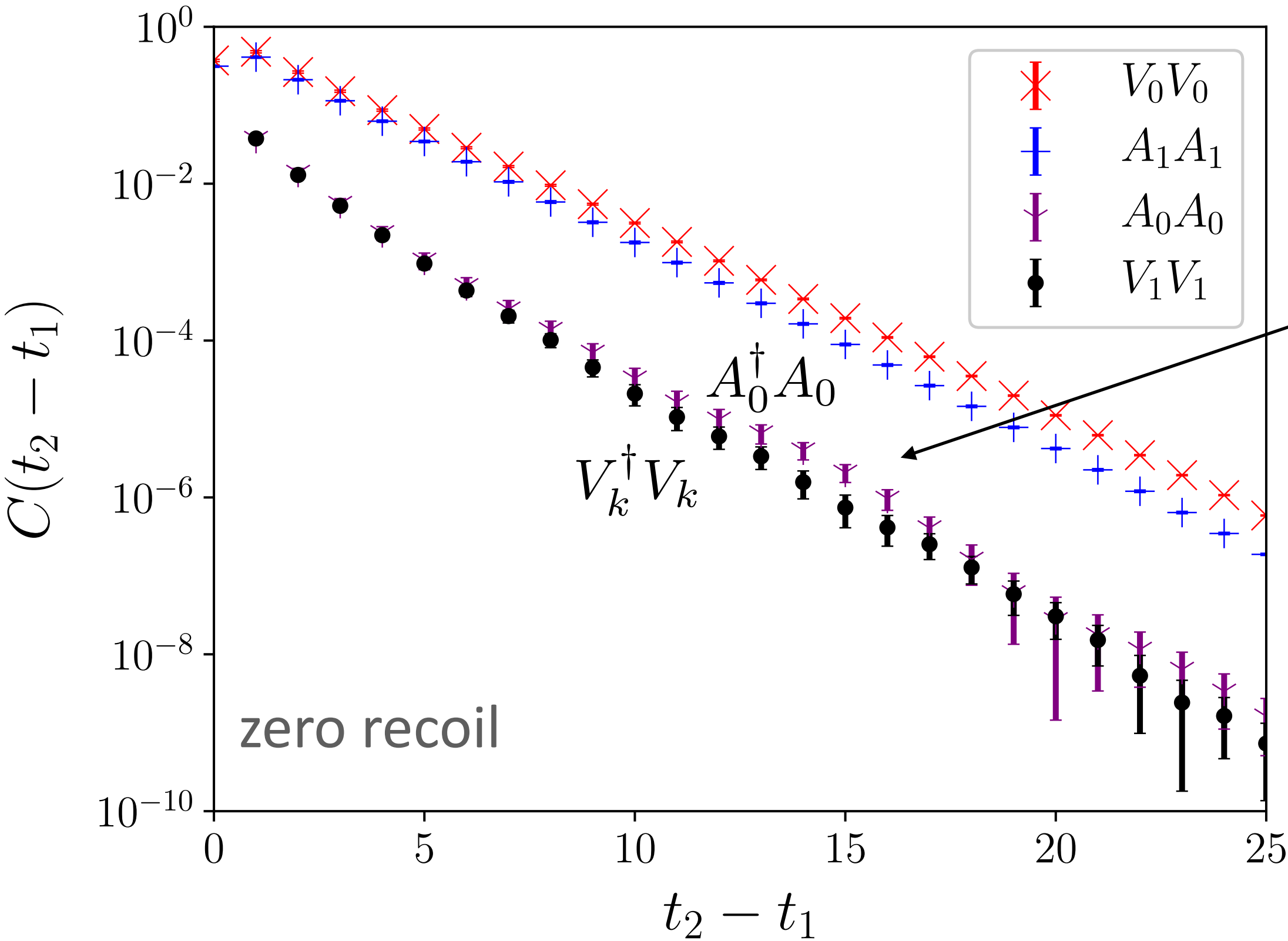


# Compton amplitude (P-wave)

$$\langle B(\mathbf{0}) | \tilde{J}_\mu^\dagger(-\mathbf{q}; t) \tilde{J}_\nu(\mathbf{q}; 0) | B(\mathbf{0}) \rangle$$

P-wave

- Opposite parity; can't be S-wave.
- Slightly heavier; Amplitude is  $\sim 50\times$  smaller.



$$|g_+(1)|^2 e^{-m_{D_0^*} t}$$

HQET Leibovich, Ligeti, Stewart, Wise (1998)

$$g_+(1) = -\frac{3}{2} \left( \frac{1}{2m_c} + \frac{1}{2m_b} \right) (\bar{\Lambda}^* - \bar{\Lambda}) \zeta(1)$$

0.35 GeV     $\sim 0.5$

Isgur-Wise func for P-wave (0<sup>+</sup>)

# Summary

- Framework to compute inclusive decay rate is now available. The energy integral can be reconstructed from Euclidean lattice correlators.
- Compton amplitudes contain the excited state contributions. They are suppressed.
  - S-wave: small wave function overlap.
  - P-wave: starts from  $1/m$  or non-zero recoil
- Decomposition into different channels. Will be compared with pQCD/OPE. See,
  - Sandro Maechler “Comparison of lattice QCD results for inclusive semi-leptonic decays of B meson with OPE”, Thu 6:30am
- See also the next talk (Jun-Sik Yoo), about an application to inelastic  $\ell N$  scattering.