

# Hadronic vacuum polarization of the muon on 2+1+1-flavor HISQ ensembles: an update.

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on behalf of the Fermilab Lattice, HPQCD, and MILC collaborations.

7/26/2021

# Muon's magnetic moment

❖ Fermilab experiment announced first measurement, new exp. average (BNL+FNAL).

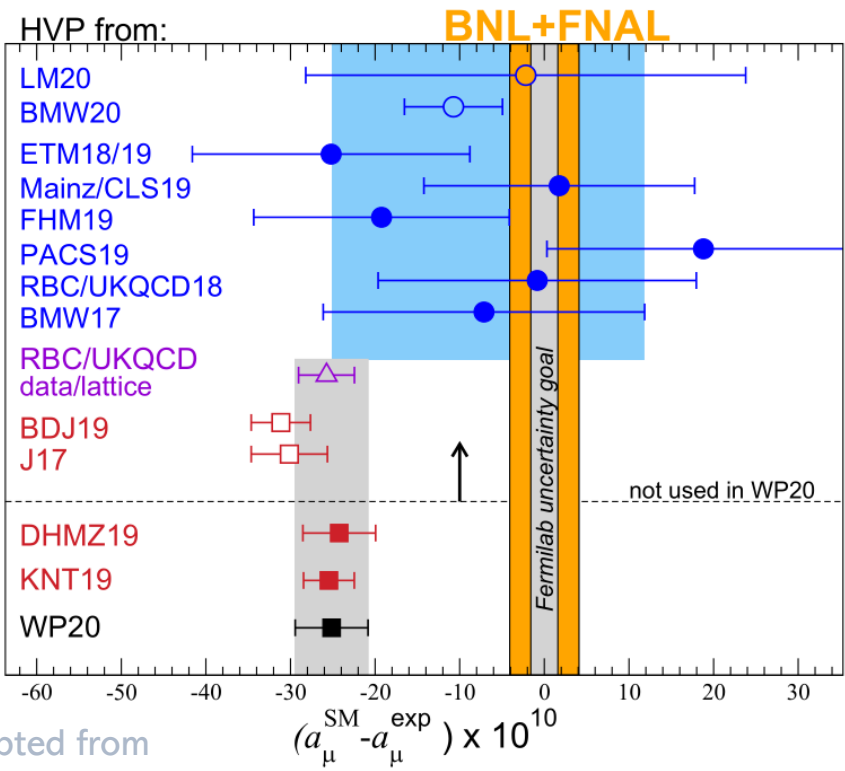
$$a_\mu(\text{Exp}) = 116\,592\,061(41) \times 10^{-11}$$

❖ 4.2 $\sigma$  tension with white paper result (WP20).

## Anomalous magnetic moment

$$\mu = g \frac{e}{2m_\mu} \mathbf{S} \Rightarrow a_\mu \equiv (g - 2)/2$$

$g = 2$  at tree level

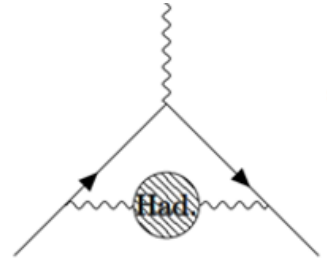


Adapted from 2006.04822

❖ HVP: leading order hadronic vacuum polarization contribution.


❖ Dominates the Standard Model uncertainty.

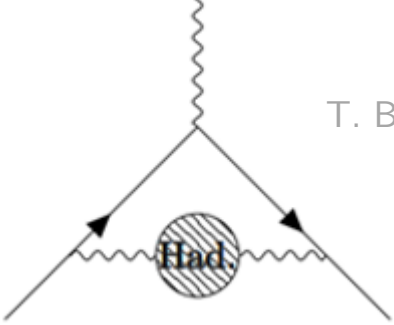
$a_\mu^{\text{HVP, LO}}$  :



# Leading Order Hadronic Vacuum Polarization

❖ Energy-momentum representation.

FT:   $\sim \hat{\Pi}_{\text{Had.}}(q^2),$



T. Blum, Phys. Rev. Lett. 91, 052001



$$a_{\mu}^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dq^2 K(q^2) \hat{\Pi}_{\text{Had.}}(q^2)$$

❖ Time-momentum representation.

Bernecker & Meyer, 2011

$$a_{\mu}^{\text{HVP,LO}} = 4\alpha^2 \int_0^{\infty} dt \tilde{K}(t) C(t), \quad C(t) = \sum_i^3 \int d^3x \langle J_i(x) J_i(0) \rangle$$

$$J_i = \sum_f Q_f \bar{q}_f \gamma_i q_f$$

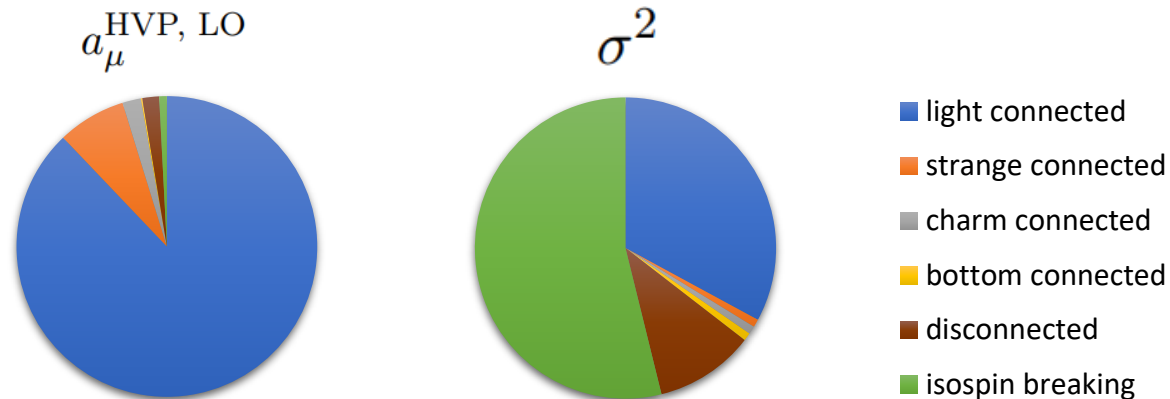
❖ Contributions separated by

- ❖ connected & disconnected wick contractions
- ❖ quark flavor
- ❖ Isospin symmetric & breaking contributions



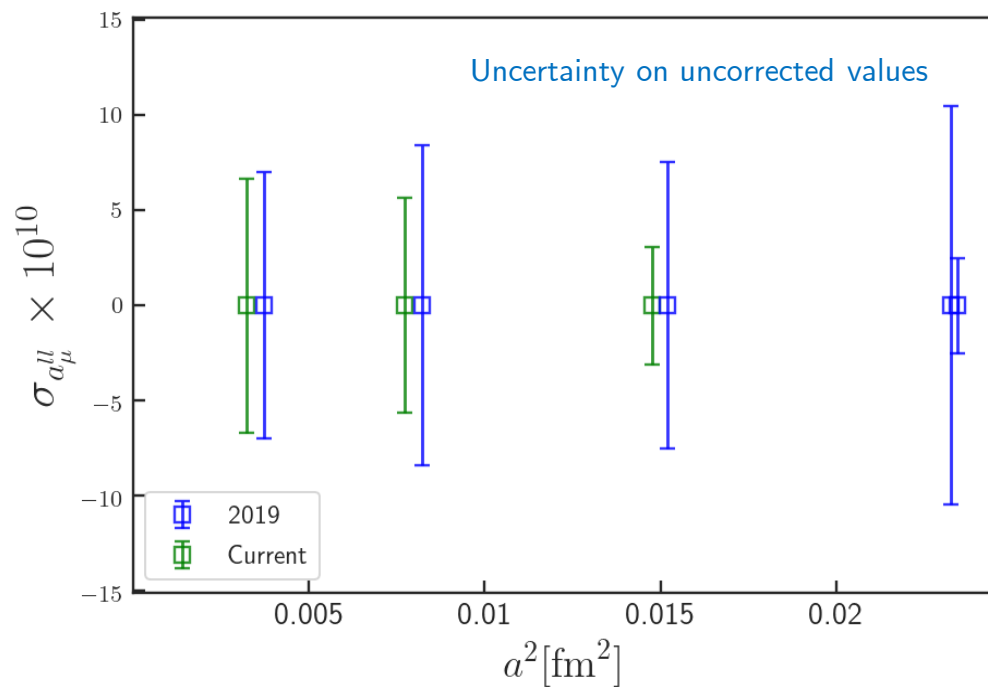
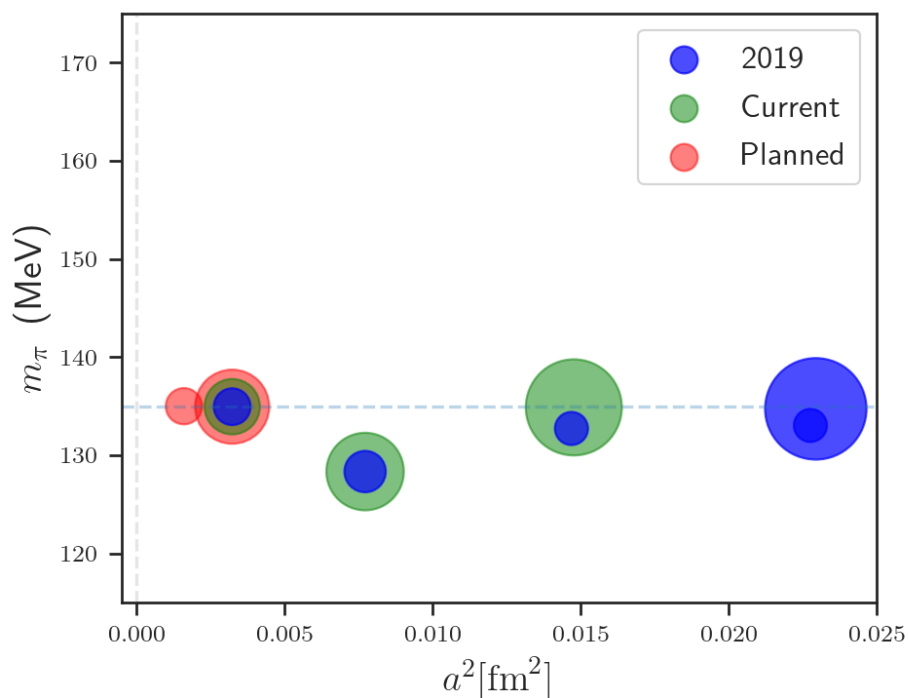
$$10^{10} a_{\mu}^{\text{HVP, LO}} = 699(15)_{u,d}(1)_{s,c,b}$$

PhysRevD (2019)  
101.034512



- ❖ Light-quark, connected contribution (90% of total). Uncertainties dominated by statistics, continuum extrapolation, scale-setting & finite volume effects.
- ❖ This talk will focus on the calculation of this contribution including an exclusive long-distance tail calculation.
- ❖ Isospin breaking and disconnected contributions: [Poster by Craig McNeile](#)

# Simulation Details

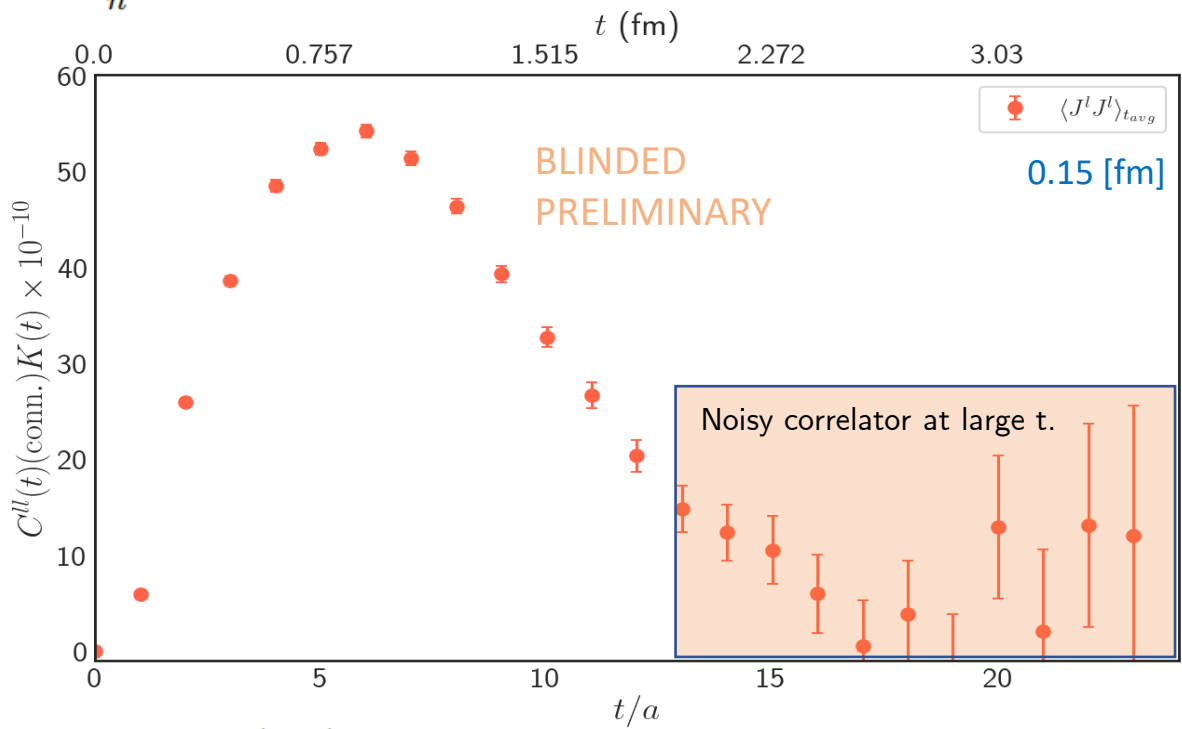


- ❖ Nf=2+1+1 HISQ ensembles with physical light-quark masses.
- ❖ 4 lattice spacings [0.06-0.15] fm, planning to add a 5<sup>th</sup> lattice spacing at 0.042 fm.
- ❖ All ensembles (except largest spacing) with  $m_\pi L > 3.7$ .

# Exclusive two-pion study

$$a_{\mu}^{ll}(\text{conn.}) = 4\alpha^2 \int_0^{\infty} dt \tilde{K}(t) C(t)_{\text{conn.}}^{ll} \quad , \quad C(t)_{\text{conn.}}^{ll} = \sum_i^3 \int d^3x \langle J_i^l(x) J_i^l(0) \rangle_{\text{conn.}}$$

$$C(t)_{\text{conn.}}^{ll} = \sum_n \left| \langle 0 | J_i^l | n \rangle \right|^2 e^{-E_n t} \quad \text{Lowest energy states are } \pi\pi.$$



$\langle 0 | J_i^l | \pi\pi \rangle$  small -> signal-to-noise problem.



*Perform dedicated spectroscopy study with  $\pi\pi$  operators that overlap strongly with these states to precisely determine the energies and amplitudes.*

# Staggered two-pion operators

Staggered pion operator

$$\pi^-(\vec{p})_\xi = \sum_x e^{i\vec{p}\cdot\vec{x}} \bar{u}(x) \gamma_5 \otimes \gamma_\xi d(x)$$

*Zero momentum taste irreps:*

*Scalar & Pseudo-scalar (1D)*

*Vector & Pseudo-vector (3D)*

*Tensor (3D+3D)*

*Non-zero momentum taste irreps:*

*1D irreps are Wilson like.*

*3D irreps can split (1D+2D) based on momentum & taste direction alignment.*

Two-pion operators

$$\mathcal{O}_{\pi\pi}(0) = \sum_{\substack{\xi_1, \xi_2 \\ \vec{p} \in \{p\}^*}} CG_{\text{stag, iso.}} \pi(\vec{p})_{\xi_1} \pi(-\vec{p})_{\xi_2}$$

Clebsch-Gordon's from irreducible representations of staggered lattice symmetry group.

$$SU_I(2) \times \left( T_N^3 \rtimes \{ \Xi_\mu, C_0 \} \rtimes \{ \tilde{R}_{ij}, I_S \} \right) = SU(2) \times \left( Z_N^3 \rtimes \Gamma_{4,1} \rtimes W_3 \right)$$

*Wigner's method twice ~ little groups of little groups.*

$$CG = \text{unitarity} : \sum_{g \in G} \left[ D^{(\pi_1)}(g) \times D^{(\pi_2)}(g) \right] AD^{(J)}(g)^\dagger$$

*-Sakata, J.Math.Phys. 15 (1974) 1702-1709*

*Taste scalar vector current "one-link" -> taste diagonal two pion states.*

# Staggered two-pion states on the lattice

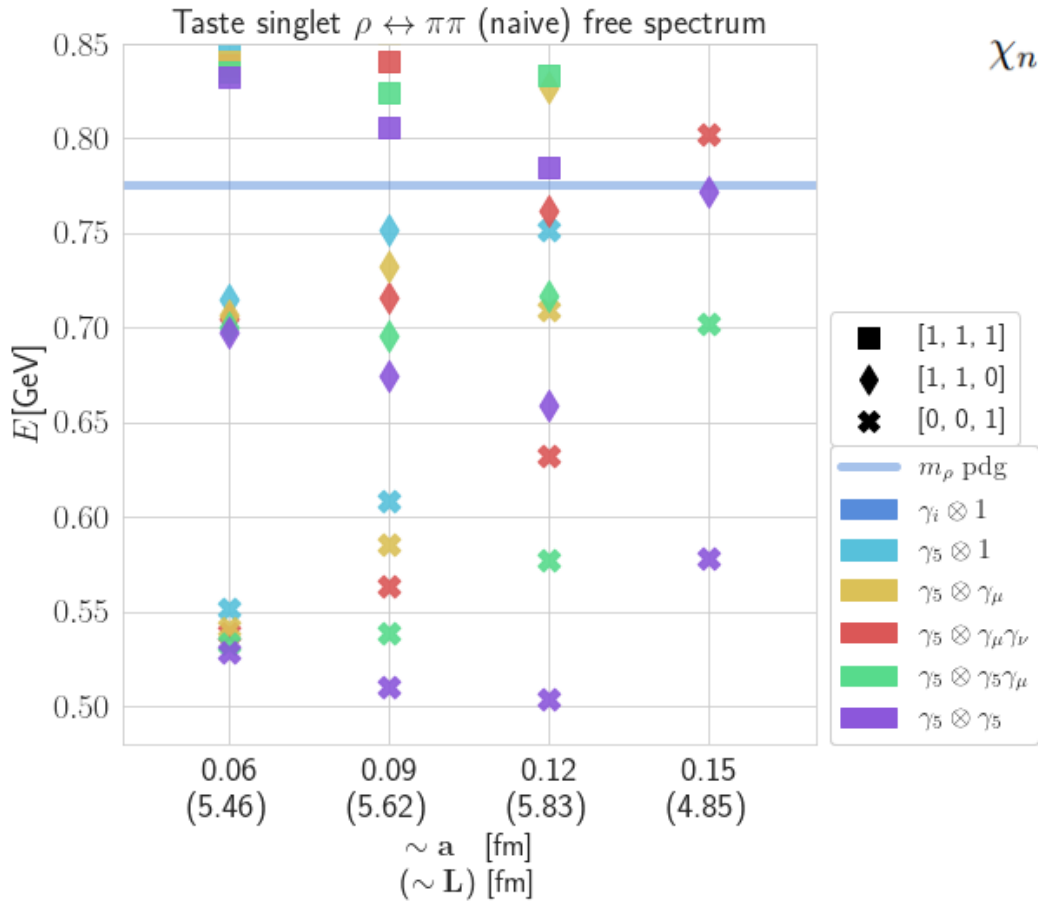
$$\mathbf{C}(t) = \begin{pmatrix} C(t)_{J, \vec{J} \rightarrow J, \vec{J}} & C(t)_{J, \vec{J} \rightarrow \pi\pi} \\ C(t)_{\pi\pi \rightarrow J, \vec{J}} & C(t)_{\pi\pi \rightarrow \pi\pi} \end{pmatrix} \quad \longrightarrow \quad \mathbf{C}(t)v = \lambda \mathbf{C}(t_0)v$$

Solve GEVP:

$\chi_n = (v_n)_i \mathcal{O}_i$  with best overlap with state  $n$ .

$$\langle \chi_n \chi_n^\dagger \rangle = \sum_n Z_n^2 e^{-E_n t}$$

$$\langle \chi_n J_l^\dagger \rangle = \sum_n Z_n \langle 0 | J_l | n \rangle e^{-E_n t}$$



❖ Operators for all states up to rho energy.

❖ Need all taste irrep rows to account for irrep splitting at non-zero mom.

$$E_{\pi_\xi \pi_\xi} = 2\sqrt{m_\xi^2 + p^2}, \quad \vec{p} = \frac{2\pi}{L} [n_x, n_y, n_z]$$



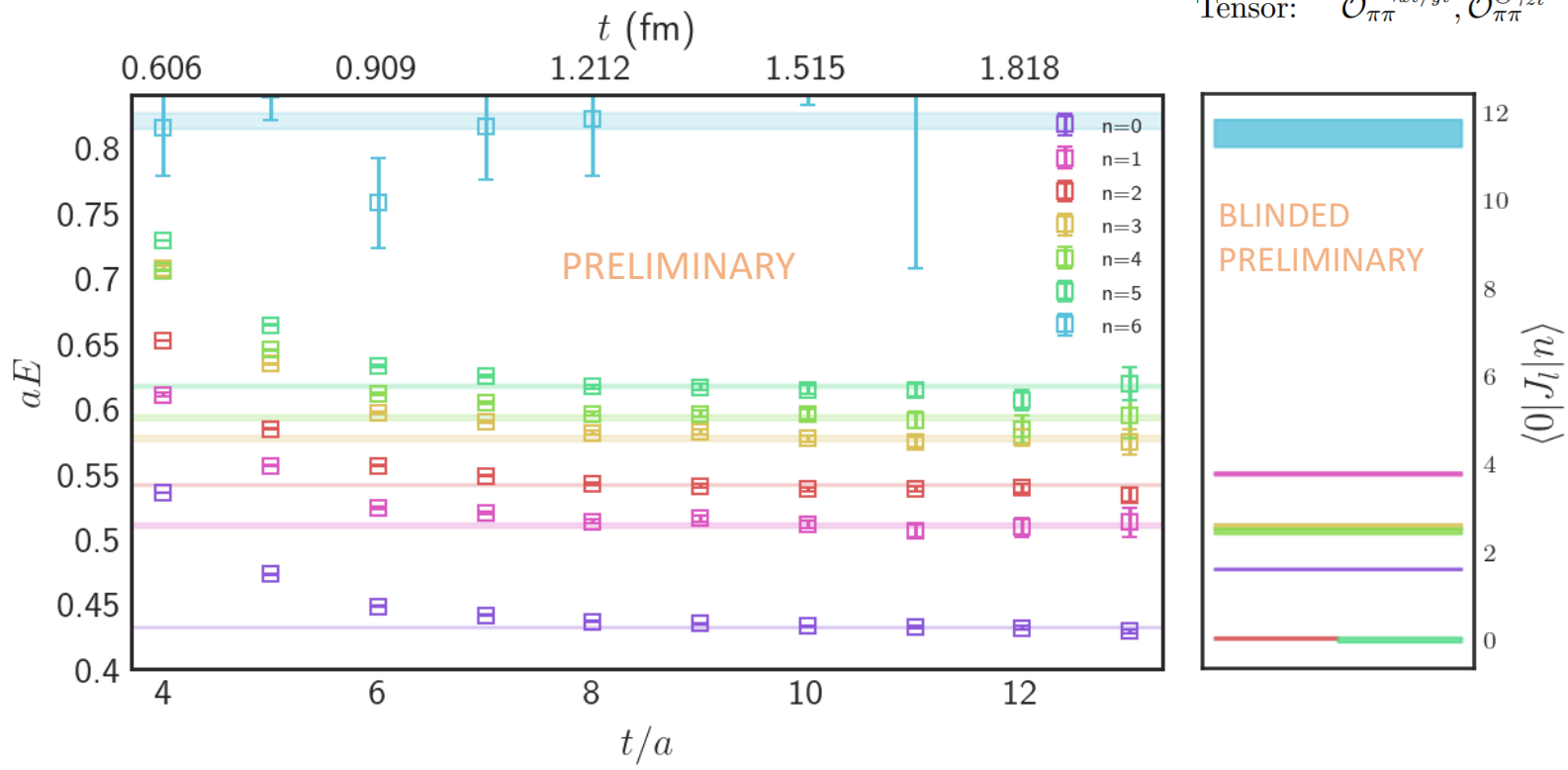
# Staggered two-pion states on the lattice: GEVP Results

$$\langle \chi_n \chi_n^\dagger \rangle = \sum_n Z_n^2 e^{-E_n t}, \quad \langle \chi_n J_l^\dagger \rangle = \sum_n Z_n \langle 0 | J_l | n \rangle e^{-E_n t}$$

Energies and Amplitudes from combined fit.

## Operator Basis at 0.15 [fm]

Vector current:	$\mathcal{O}_J, \tilde{\mathcal{O}}_J$	
Pseudo-scalar:	$\mathcal{O}_{\pi\pi}^{\otimes 5}$	[0, 0, 1], [1, 1, 0]
Pseudo-vector:	$\mathcal{O}_{\pi\pi}^{\otimes 5x/y}, \mathcal{O}_{\pi\pi}^{\otimes 5z}$	[0, 0, 1]
Tensor:	$\mathcal{O}_{\pi\pi}^{\otimes \gamma_{xt/yt}}, \mathcal{O}_{\pi\pi}^{\otimes \gamma_{zt}}$	[0, 0, 1]

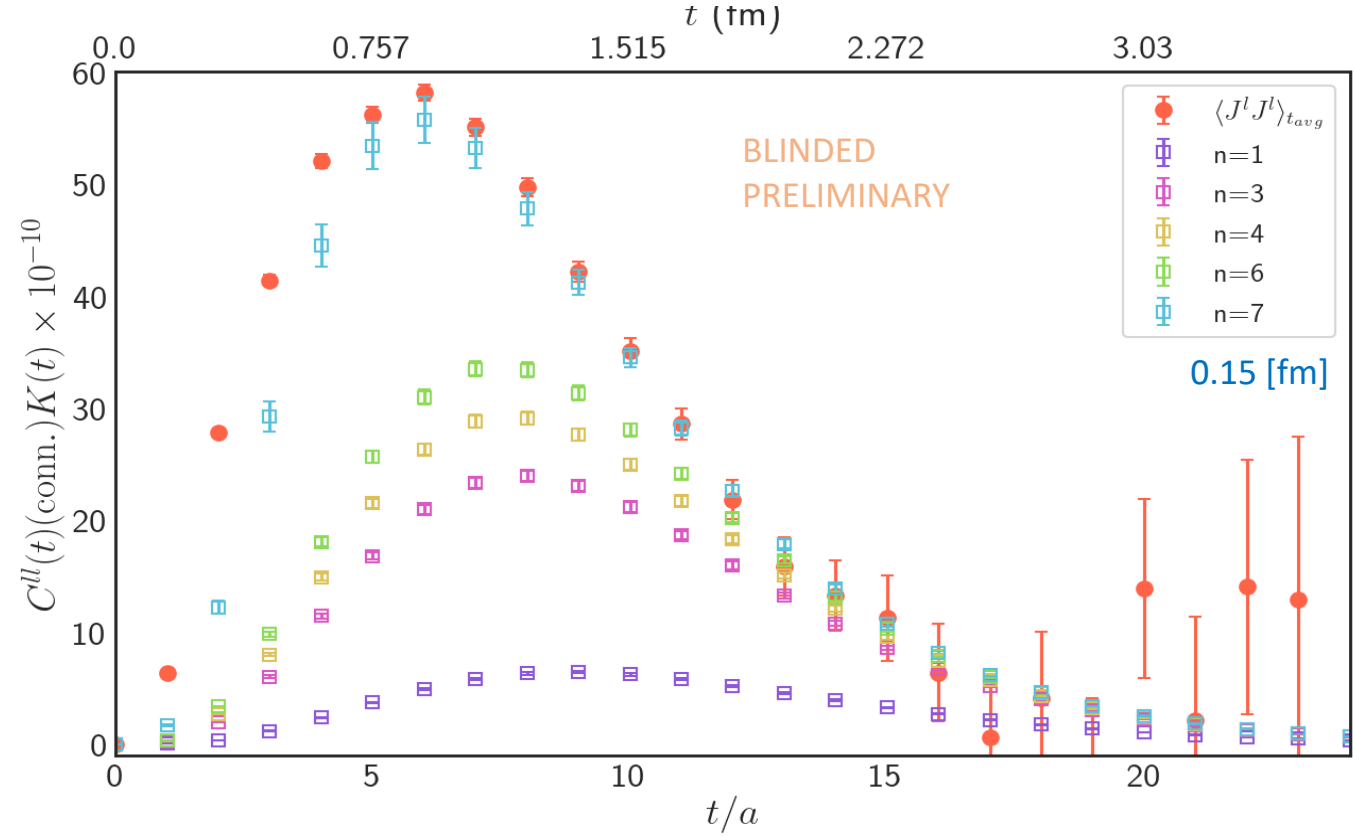


Effective mass plot for diagonal correlation functions, bands are fit results.

- ❖ Pseudo-vector & tensor taste irrep splitting's (1&2, 4&5) are non-degenerate.
- ❖ Overlap amplitudes heavily suppressed for certain states (2&5).

# Correlator Reconstruction

$$C(t)_{\text{conn.}}^{ll} = \sum_n \left| \langle 0 | J_i^l | n \rangle \right|^2 e^{-E_n t}$$



Vector current is time-averaged to suppress oscillating contributions.

- ❖ Raw:  $\frac{\sigma_{a_{\mu}^{ll}}}{a_{\mu}^{ll}} \times 100$  7.8%
- ❖ Bounding Method  $t_c = 2.65$  fm: 4.9%
- ❖ Reconstruction,  $n=6$ ,  $t^* = 1.82$  fm, two-pion states beneath rho : 1.5%
- ❖ Reconstruction,  $n=7$ ,  $t^* = 1.67$  fm, Including the rho: 1.4%

# Window Analysis

*Restrict integral to some Euclidean time region to emphasize different physics / lattice effects.*

$$a_{\mu, \text{win}} = 4\alpha^2 \int_0^\infty dt C(t) \tilde{K}(t) \Theta(t, t_1, t_2, \Delta)$$

$$\Theta(t, t_1, t_2, \Delta) = \frac{1}{2} \left[ \tanh\left(\frac{t-t_1}{\Delta}\right) - \tanh\left(\frac{t-t_2}{\Delta}\right) \right]$$

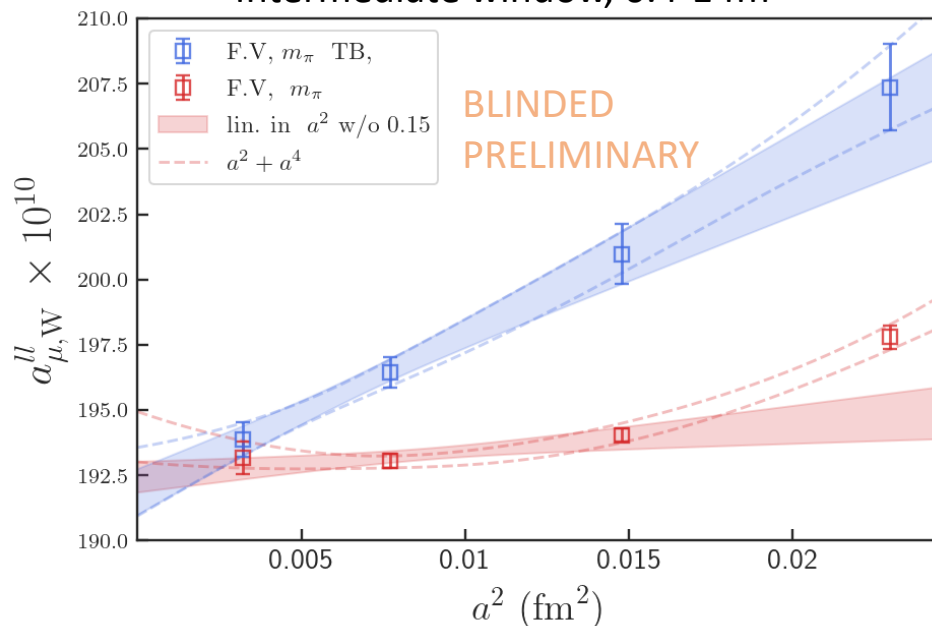
T. Blum et al,  
arXiv:1801.07224

Window	$[t_1, t_2]$ fm	$\Delta$ [fm]
SD	$[-\infty, 0.4]$	0.15
W	$[0.4, 1]$	0.15
LD	$[1, \infty]$	0.15

- SD: short-distance, discretization effects prominent.
- LD: long-distance (two-pion states), finite volume effects prominent.
- W: less sensitive to these extremum effects, easier to compute with precision to compare different lattice calculations.

**Goal of lattice g-2 community: agreement on these quantities.**

Intermediate window, 0.4-1 fm



- ❖ Corrections from  $\rho - \gamma - \pi\pi$  model (leading order).
- ❖ Good consistency between extrapolations of data with(out) discretization effect corrections.

# Conclusions

Goal of the Fermilab Lattice, HPQCD & MILC collaboration:

Compute the HVP contribution to  $g-2$  using lattice QCD with an uncertainty of  $< 0.5\%$ .

To achieve this, the light-quark, connected contribution must be known with commensurate precision as it accounts for  $\sim 90\%$  of the total value. To achieve this, we aim to

- ❖ Increase statistics at finest lattice spacings.
- ❖ Compute two-pion tail exclusively.
- ❖ Add a 5<sup>th</sup> lattice spacing (0.042 fm) to analysis.
- ❖ Reducing scale setting uncertainty through absolute & relative scale calculations. 1912.00028.
- ❖ Refining estimates on FV corrections & uncertainties.

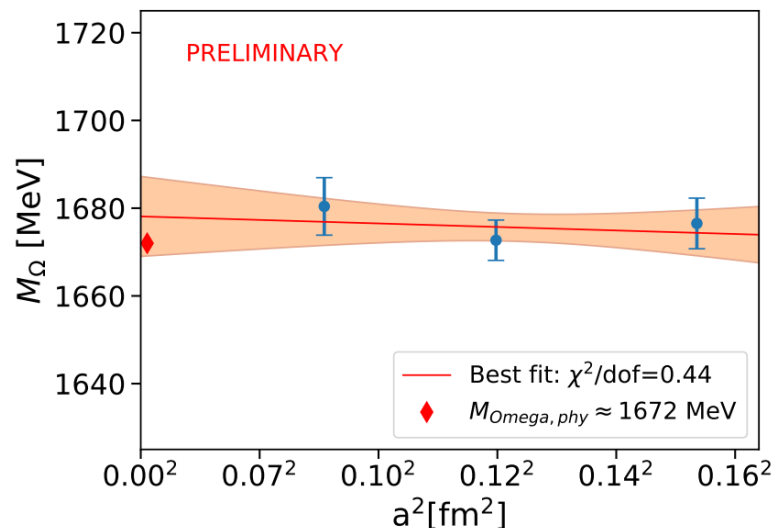
In parallel to the work presented here, calculations for

- ❖ Disconnected contribution: goal  $< 10\%$  total.
- ❖ IB effects: goal is explicit calculation of all leading-order effects.
  - ❖ SIB effects on connected & disconnected contributions.
  - ❖ Leading QED effects. [Poster from Craig McNeile for disconnected & IB results.](#)

**Thank you**

## Absolute scale parameter $M_\Omega$

- ❖ Insensitive to isospin breaking effects compared with  $f_\pi$ .
- ❖  $M_\Omega$  at 0.15, 0.12 and 0.09 fm, C Hughes et al., 1912.00028.
- ❖ Increasing statistics & extending to 0.06 fm.



## Relative scale parameter $w_0/a$

- ❖ Ongoing high statistics gradient flow study on all HISQ ensembles.
- ❖ 6 combinations of ensemble/flow/observable action discretization to study & quantify cut-off effects.

❖ Self-consistency check:

$$a_{\mu}^{ll}(\text{conn.}) = a_{\mu,SD}^{ll} + a_{\mu,W}^{ll} + a_{\mu,LD}^{ll}$$

