

Hadronic vacuum polarization of the muon on 2+1+1-flavor HISQ ensembles: an update.

Lattice 2021,

MIT – JULY 26-30, 2021

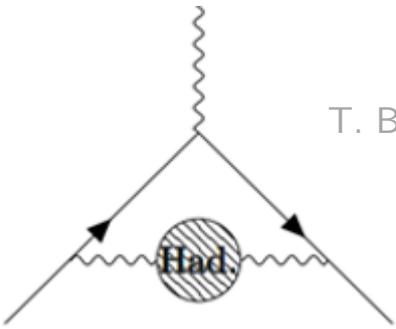
Shaun Lahert – University of Illinois Urbana-Champaign
on behalf of the Fermilab Lattice, HPQCD, and MILC collaborations.

7/26/2021

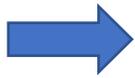
Leading Order Hadronic Vacuum Polarization

❖ Energy-momentum representation.

FT: $\text{---} \circlearrowleft \text{---} \sim \hat{\Pi}_{\text{Had.}}(q^2),$
 $J_i(0) \quad \text{Had.} \quad J_i(x)$



T. Blum, Phys. Rev. Lett. 91, 052001



$$a_{\mu}^{\text{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} dq^2 K(q^2) \hat{\Pi}_{\text{Had.}}(q^2)$$

❖ Time-momentum representation.

Bernecker & Meyer, 2011

$$a_{\mu}^{\text{HVP,LO}} = 4\alpha^2 \int_0^{\infty} dt \tilde{K}(t) C(t) \quad , \quad C(t) = \sum_i^3 \int d^3x \langle J_i(x) J_i(0) \rangle$$

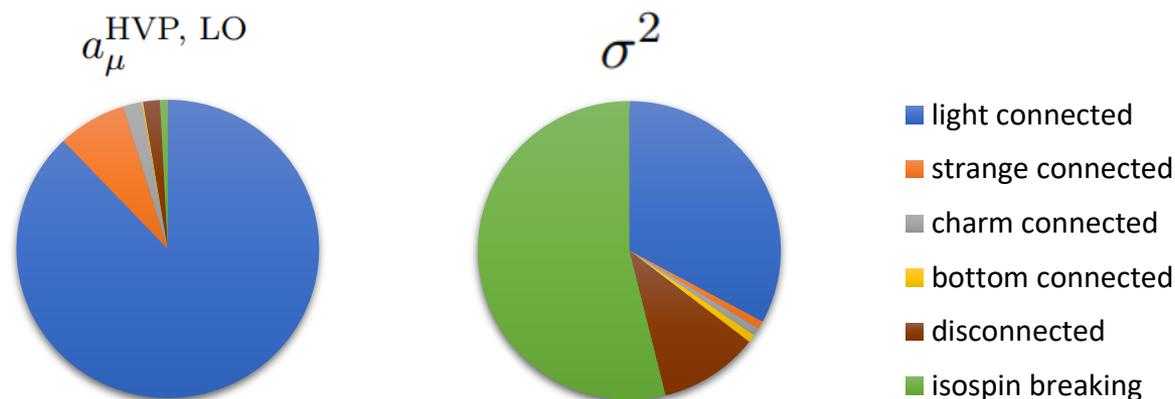
$$J_i = \sum_f Q_f \bar{q}_f \gamma_i q_f$$

- ❖ Contributions separated by
 - ❖ connected & disconnected wick contractions
 - ❖ quark flavor
 - ❖ Isospin symmetric & breaking contributions



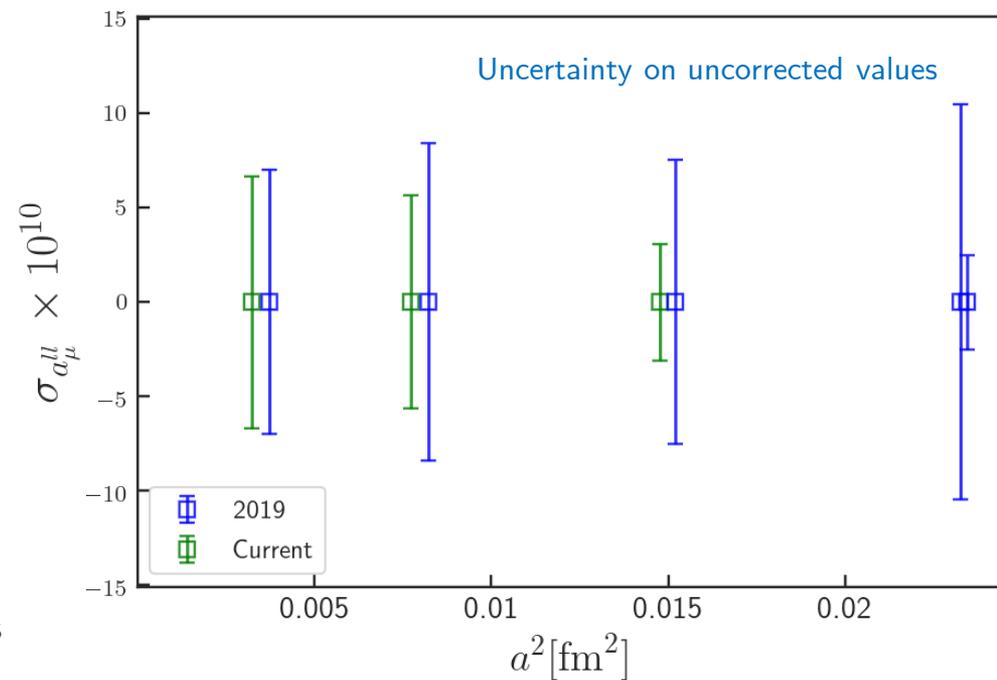
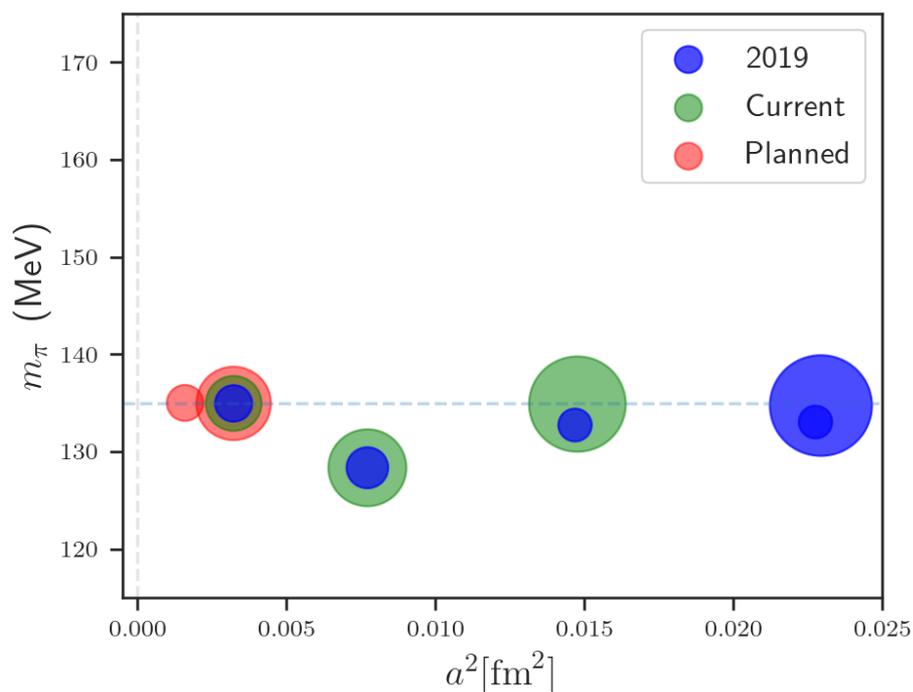
$$10^{10} a_{\mu}^{\text{HVP,LO}} = 699(15)_{u,d}(1)_{s,c,b}$$

PhysRevD (2019)
101.034512



- ❖ Light-quark, connected contribution (90% of total). Uncertainties dominated by statistics, continuum extrapolation, scale-setting & finite volume effects.
- ❖ This talk will focus on the calculation of this contribution including an exclusive long-distance tail calculation.
- ❖ Isospin breaking and disconnected contributions: [Poster by Craig McNeile](#)

Simulation Details



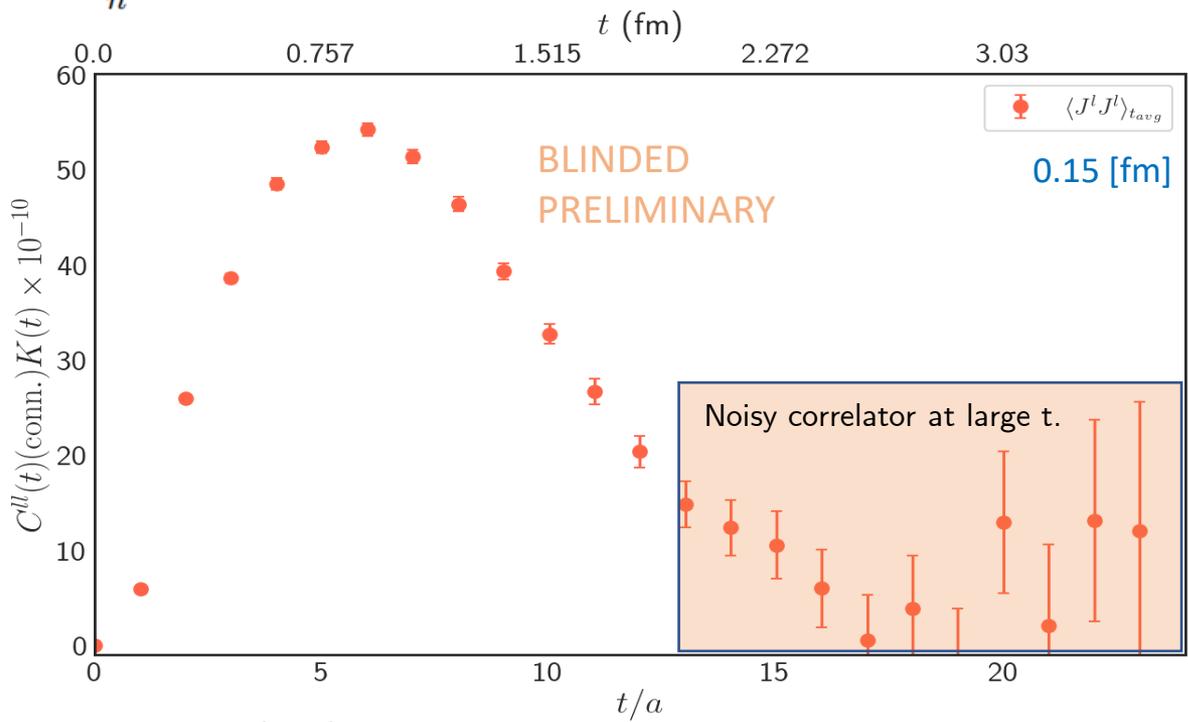
- ❖ Nf=2+1+1 HISQ ensembles with physical light-quark masses.
- ❖ 4 lattice spacings [0.06-0.15] fm, planning to add a 5th lattice spacing at 0.042 fm.
- ❖ All ensembles (except largest spacing) with $m_\pi L > 3.7$.

Exclusive two-pion study

$$a_{\mu}^{ll}(\text{conn.}) = 4\alpha^2 \int_0^{\infty} dt \tilde{K}(t) C(t)_{\text{conn.}}^{ll}, \quad C(t)_{\text{conn.}}^{ll} = \sum_i^3 \int d^3x \langle J_i^l(x) J_i^l(0) \rangle_{\text{conn.}}$$

$$C(t)_{\text{conn.}}^{ll} = \sum_n \left| \langle 0 | J_i^l | n \rangle \right|^2 e^{-E_n t}$$

Lowest energy states are $\pi\pi$.



$\langle 0 | J_i^l | \pi\pi \rangle$ small -> signal-to-noise problem.



Perform dedicated spectroscopy study with $\pi\pi$ operators that overlap strongly with these states to precisely determine the energies and amplitudes.

Staggered two-pion operators

Staggered pion operator

$$\pi^-(\vec{p})_\xi = \sum_x e^{i\vec{p}\cdot\vec{x}} \bar{u}(x) \gamma_5 \otimes \gamma_\xi d(x)$$

Zero momentum taste irreps:

Scalar & Pseudo-scalar (1D)

Vector & Pseudo-vector (3D)

Tensor (3D+3D)

Non-zero momentum taste irreps:

1D irreps are Wilson like.

3D irreps can split (1D+2D) based on momentum & taste direction alignment.

Two-pion operators

$$\mathcal{O}_{\pi\pi}(0) = \sum_{\substack{\xi_1, \xi_2 \\ \vec{p} \in \{p\}^*}} CG_{\text{stag, iso.}} \pi(\vec{p})_{\xi_1} \pi(-\vec{p})_{\xi_2}$$

Clebsch-Gordon's from irreducible representations of staggered lattice symmetry group.

$$SU_I(2) \times \left(T_N^3 \rtimes \{ \Xi_\mu, C_0 \} \rtimes \{ \tilde{R}_{ij}, I_S \} \right) = SU(2) \times \left(Z_N^3 \rtimes \Gamma_{4,1} \rtimes W_3 \right)$$

Wigner's method twice ~ little groups of little groups.

$$CG = \text{unitarity} : \sum_{g \in G} \left[D^{(\pi_1)}(g) \times D^{(\pi_2)}(g) \right] AD^{(J)}(g)^\dagger$$

-Sakata, J.Math.Phys. 15 (1974) 1702-1709

Taste scalar vector current "one-link" -> taste diagonal two pion states.

Staggered two-pion states on the lattice

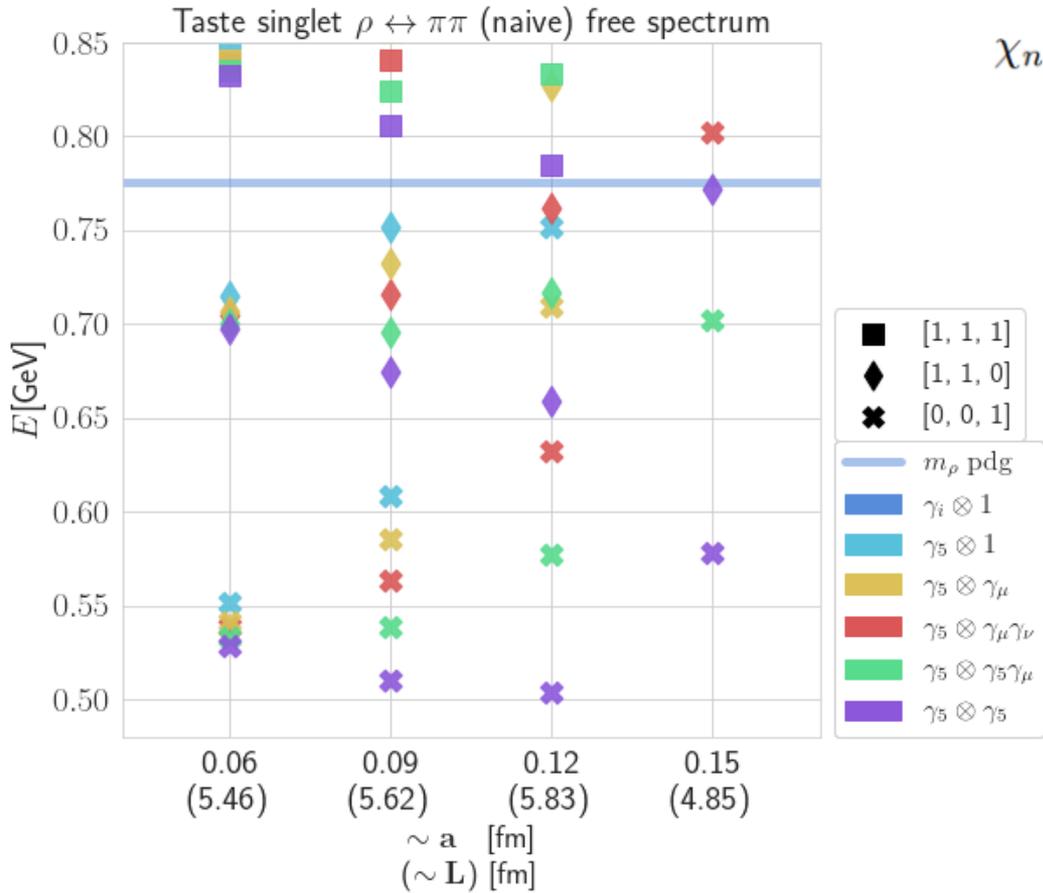
$$\mathbf{C}(t) = \begin{pmatrix} C(t)_{J,\vec{J} \rightarrow J,\vec{J}} & C(t)_{J,\vec{J} \rightarrow \pi\pi} \\ C(t)_{\pi\pi \rightarrow J,\vec{J}} & C(t)_{\pi\pi \rightarrow \pi\pi} \end{pmatrix} \quad \longrightarrow \quad \mathbf{C}(t)v = \lambda \mathbf{C}(t_0)v$$

Solve GEVP:

$\chi_n = (v_n)_i \mathcal{O}_i$ with best overlap with state n .

$$\langle \chi_n \chi_n^\dagger \rangle = \sum_n Z_n^2 e^{-E_n t}$$

$$\langle \chi_n J_l^\dagger \rangle = \sum_n Z_n \langle 0 | J_l | n \rangle e^{-E_n t}$$



❖ Operators for all states up to rho energy.

❖ Need all taste irrep rows to account for irrep splitting at non-zero mom.

$$E_{\pi_\xi \pi_\xi} = 2\sqrt{m_\xi^2 + p^2}, \quad \vec{p} = \frac{2\pi}{L} [n_x, n_y, n_z]$$

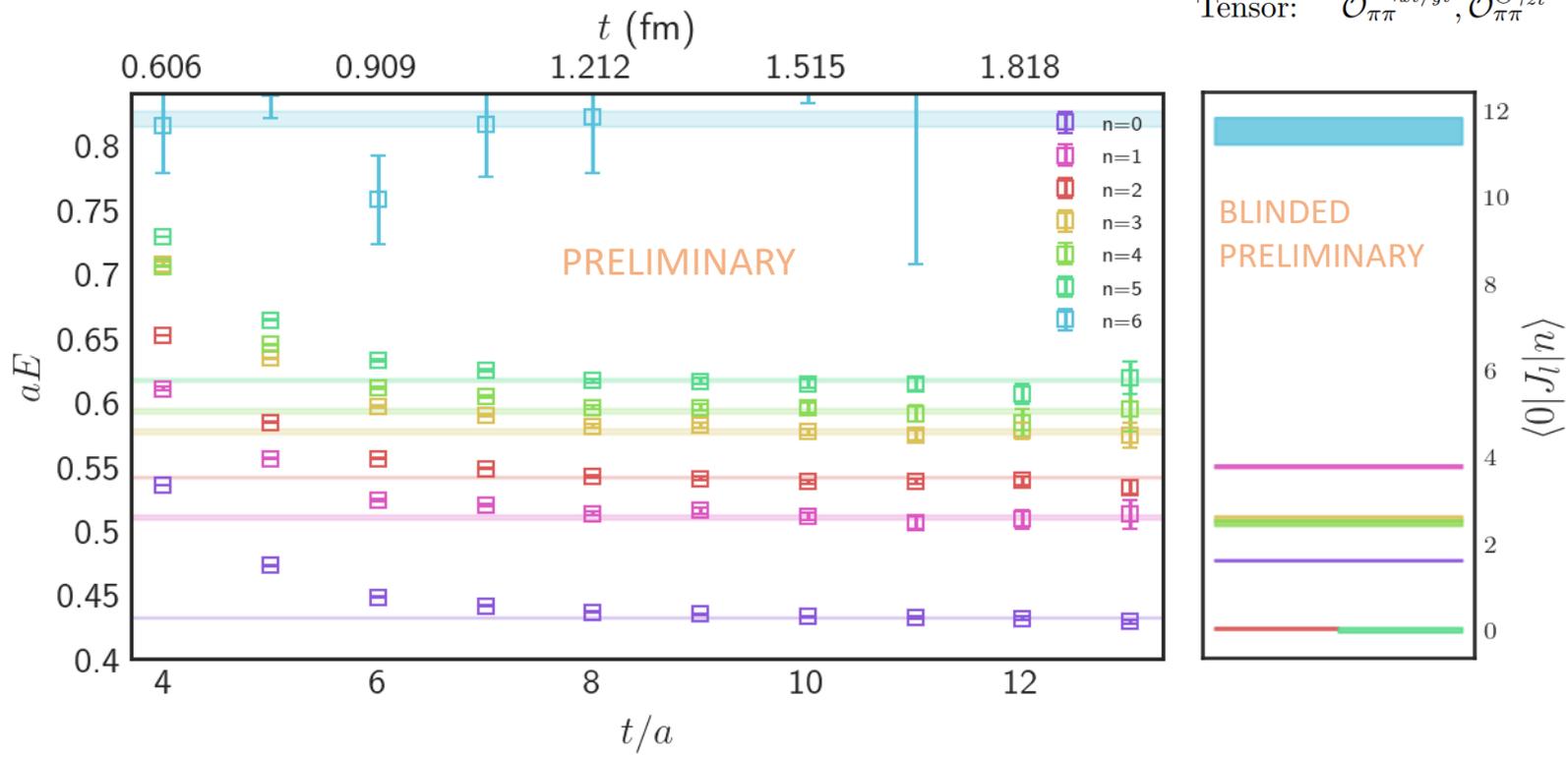
Staggered two-pion states on the lattice: GEVP Results

$$\langle \chi_n \chi_n^\dagger \rangle = \sum_n Z_n^2 e^{-E_n t}, \quad \langle \chi_n J_l^\dagger \rangle = \sum_n Z_n \langle 0 | J_l | n \rangle e^{-E_n t}$$

Energies and Amplitudes from combined fit.

Operator Basis at 0.15 [fm]

Vector current:	$\mathcal{O}_J, \tilde{\mathcal{O}}_J$	
Pseudo-scalar:	$\mathcal{O}_{\pi\pi}^{\otimes 5}$	[0, 0, 1], [1, 1, 0]
Pseudo-vector:	$\mathcal{O}_{\pi\pi}^{\otimes 5x/y}, \mathcal{O}_{\pi\pi}^{\otimes 5z}$	[0, 0, 1]
Tensor:	$\mathcal{O}_{\pi\pi}^{\otimes \gamma_{xt/yt}}, \mathcal{O}_{\pi\pi}^{\otimes \gamma_{zt}}$	[0, 0, 1]

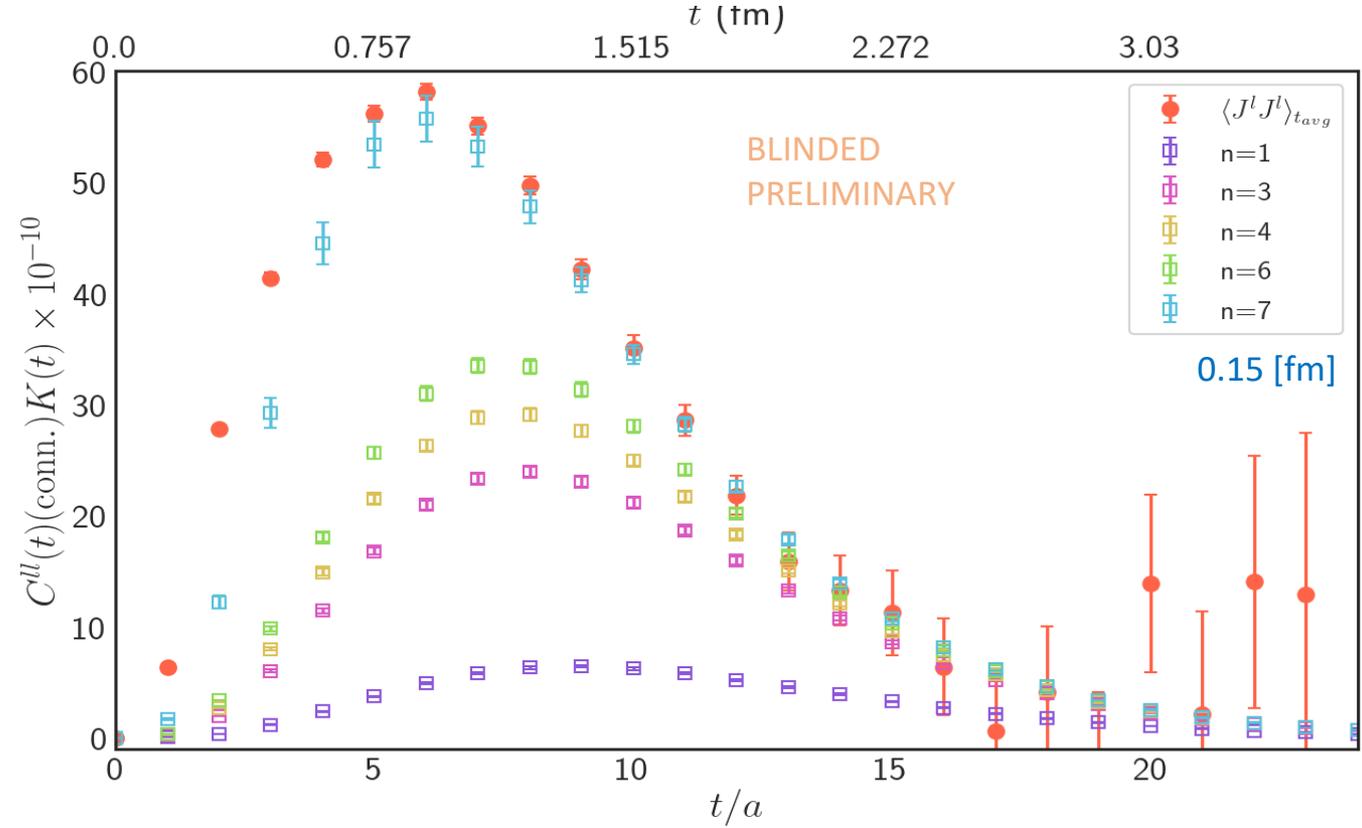


Effective mass plot for diagonal correlation functions, bands are fit results.

- ❖ Pseudo-vector & tensor taste irrep splitting's (1&2, 4&5) are non-degenerate.
- ❖ Overlap amplitudes heavily suppressed for certain states (2&5).

Correlator Reconstruction

$$C(t)_{\text{conn.}}^{ll} = \sum_n \left| \langle 0 | J_i^l | n \rangle \right|^2 e^{-E_n t}$$



Vector current is time-averaged to suppress oscillating contributions.

- ❖ Raw: $\frac{\sigma_{a_{\mu}^{ll}}}{a_{\mu}^{ll}} \times 100$ 7.8%
- ❖ Bounding Method $t_c = 2.65$ fm: 4.9%
- ❖ Reconstruction, $n=6$, $t^* = 1.82$ fm, two-pion states beneath rho : 1.5%
- ❖ Reconstruction, $n=7$, $t^* = 1.67$ fm, Including the rho: 1.4%

Window Analysis

Restrict integral to some Euclidean time region to emphasize different physics / lattice effects.

$$a_{\mu, \text{win}} = 4\alpha^2 \int_0^\infty dt C(t) \tilde{K}(t) \Theta(t, t_1, t_2, \Delta)$$

$$\Theta(t, t_1, t_2, \Delta) = \frac{1}{2} \left[\tanh\left(\frac{t-t_1}{\Delta}\right) - \tanh\left(\frac{t-t_2}{\Delta}\right) \right]$$

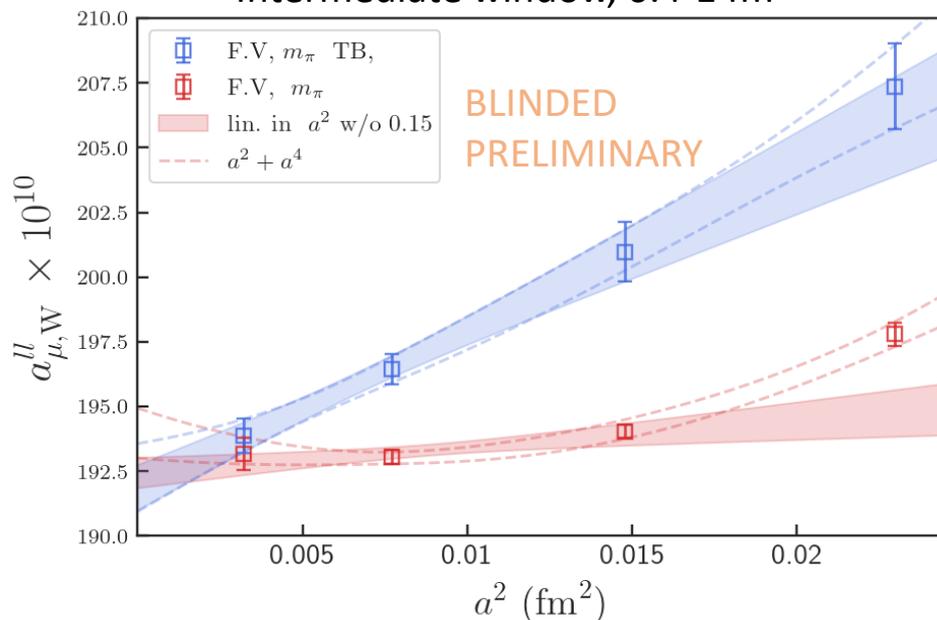
T. Blum et al,
arXiv:1801.07224

Window	$[t_1, t_2]$ fm	Δ [fm]
SD	$[-\infty, 0.4]$	0.15
W	$[0.4, 1]$	0.15
LD	$[1, \infty]$	0.15

- SD: short-distance, discretization effects prominent.
- LD: long-distance (two-pion states), finite volume effects prominent.
- W: less sensitive to these extremum effects, easier to compute with precision to compare different lattice calculations.

Goal of lattice g-2 community: agreement on these quantities.

Intermediate window, 0.4-1 fm



- ❖ Corrections from $\rho - \gamma - \pi\pi$ model (leading order).
- ❖ Good consistency between extrapolations of data with(out) discretization effect corrections.

Conclusions

Goal of the Fermilab Lattice, HPQCD & MILC collaboration:

Compute the HVP contribution to $g-2$ using lattice QCD with an uncertainty of $< 0.5\%$.

To achieve this, the light-quark, connected contribution must be known with commensurate precision as it accounts for $\sim 90\%$ of the total value. To achieve this, we aim to

- ❖ Increase statistics at finest lattice spacings.
- ❖ Compute two-pion tail exclusively.
- ❖ Add a 5th lattice spacing (0.042 fm) to analysis.
- ❖ Reducing scale setting uncertainty through absolute & relative scale calculations. 1912.00028.
- ❖ Refining estimates on FV corrections & uncertainties.

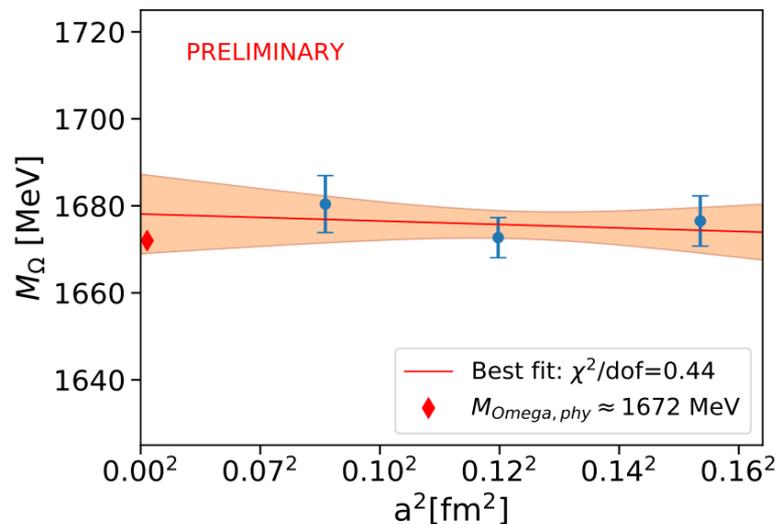
In parallel to the work presented here, calculations for

- ❖ Disconnected contribution: goal $< 10\%$ total.
- ❖ IB effects: goal is explicit calculation of all leading-order effects.
 - ❖ SIB effects on connected & disconnected contributions.
 - ❖ Leading QED effects. [Poster from Craig McNeile for disconnected & IB results.](#)

Thank you

Absolute scale parameter M_Ω

- ❖ Insensitive to isospin breaking effects compared with f_π .
- ❖ M_Ω at 0.15, 0.12 and 0.09 fm, C Hughes et al., 1912.00028.
- ❖ Increasing statistics & extending to 0.06 fm.



Relative scale parameter w_0/a

- ❖ Ongoing high statistics gradient flow study on all HISQ ensembles.
- ❖ 6 combinations of ensemble/flow/observable action discretization to study & quantify cut-off effects.

❖ Self-consistency check:

$$a_{\mu}^{ll}(\text{conn.}) = a_{\mu,SD}^{ll} + a_{\mu,W}^{ll} + a_{\mu,LD}^{ll}$$

