

# Determination of the light, strange and charm quark masses using twisted mass fermions



*Constantia Alexandrou*



ETMC: C. A., S. Bacchio, G. Bergner, M. Constantinou, M. Di Carlo, P. Dimopoulos, J. Finkenrath, E. Fiorenza, R. Frezzotti, M. Garofalo, K. Hadjiyiannakou, B. Kostrzewa, G. Koutsou, K. Jansen, V. Lubicz, M. Mangin-Brinet, F. Manigrasso, G. Martinelli, E. Papadiofantous, F. Pittler, G.C. Rossi, F. Sanfilippo, S. Simula, C. Tarantino, A. Todaro, C. Urbach, U. Wenger, arXiv:[2104.13408](https://arxiv.org/abs/2104.13408),



**STIMULATE**  
European Joint Doctorates

# Approach

- \* Perform analysis of  $N_f=2+1+1$  twisted mass fermions gauge ensembles simulated
  - ❖ at three values of the lattice spacing
  - ❖ with a range of light quark masses in the range of 135 MeV to 350 MeV
  - ❖ strange and charm quark mass tuned to approximately their physical value
- \* Employ  $N_f=4$  twisted mass fermion gauge ensembles for the computation of the renormalisation constants - **see talk by M. Di Carlo, Friday 6:15 am**
- \* Extract the quark masses using two setups
  - ❖ the meson sector
  - ❖ the baryon sector
- \* Main improvements as compared to our previous work
  - ❖ two ensembles at the physical point
  - ❖ more accurate determination of renormalisation constants
  - ❖ extraction using baryon sector

# Lattice ensembles

$N_f=2+1+1$  twisted mass fermions with a clover term

Generated by the European Twisted Mass Collaboration (ETMC)

Abdel-Rehim *et al.* (ETMC) PRD 95, 094515 (2017), arXiv: 1507.05068

C. A. *et al.* (ETMC) PRD 98, 054518 (2018), arXiv:1807.00495

C. A. *et al.* (ETMC), arXiv: 2104.06747

Ensemble	$L^3 \times T$	MDUs	$a\mu_\ell$	$af_\pi$	$m_\pi L$	$m_N/m_\pi$	$m_\pi$ [MeV]
$\beta = 1.726, c_{SW} = 1.74, a\mu_\sigma = 0.1408, a\mu_\delta = 0.1521, w_0/a = 1.8352$ (35)							
cA211.53.24	$24^3 \times 48$	5026	0.00530	0.07106 (36)	3.99	–	346.4 (1.6)
cA211.40.24	$24^3 \times 48$	5298	0.00400	0.06809 (30)	3.47	–	301.6 (2.1)
cA211.30.32	$32^3 \times 64$	10234	0.00300	0.06674 (15)	4.01	4.049 (14)	261.1 (1.1)
cA211.12.48	$48^3 \times 96$	2936	0.00120	0.06133 (33)	3.85	5.685 (28)	167.1 (0.8)
$\beta = 1.778, c_{SW} = 1.69, a\mu_\sigma = 0.1246864, a\mu_\delta = 0.1315052, w_0/a = 2.1299$ (16)							
cB211.25.32	$32^3 \times 64$	3959	0.00250	0.05652 (38)	3.35	4.104 (36)	253.3 (1.4)
cB211.25.48	$48^3 \times 96$	5246	0.00250	0.05726 (12)	5.02	4.124 (17)	253.0 (1.0)
cB211.14.64	$64^3 \times 128$	6187	0.00140	0.05477 (12)	5.02	5.119 (36)	189.8 (0.7)
<b>cB211.072.64</b>	$64^3 \times 128$	3161	0.00072	0.05267 (14)	3.62	6.760 (30)	<b>136.8 (0.6)</b>
$\beta = 1.836, c_{SW} = 1.6452, a\mu_\sigma = 0.106586, a\mu_\delta = 0.107146, w_0/a = 2.5045$ (17)							
cC211.20.48	$48^3 \times 96$	2000	0.00200	0.04892 (13)	4.13	4.244 (25)	245.73 (98)
<b>cC211.06.80</b>	$80^3 \times 160$	3207	0.00060	0.04504 (10)	3.78	6.916 (19)	<b>134.3 (0.5)</b>

A-ensembles  
B-ensembles  
C-ensembles

Not used in the baryon sector

Lattice spacing from pion sector

- Use Osterwalder-Seiler strange and charm valence quarks  $\rightarrow$  tune using i) meson and ii) baryon masses
- Lattice spacing determine from i)  $f_\pi$  and ii) the nucleon mass  $m_N$  in the so symmetric limit

# Pseudoscalar renormalisation constant $Z_P$

For details see talk by M. Di Carlo, Friday 6:15 am

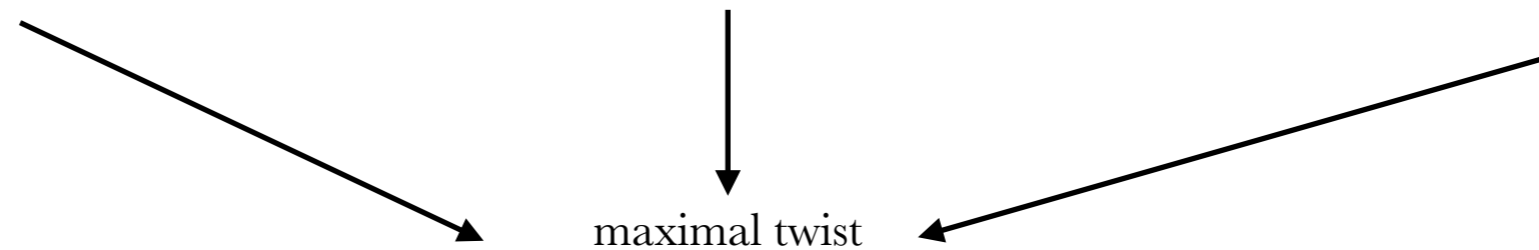
In the twisted mass fermion (TMF) formulation at maximal twist the renormalised mass is given by  $m_f = \frac{\mu_f}{Z_P}$

bare twisted mass parameter of quark f

- We employ the non-perturbative, mass independent RI'-MOM renormalisation scheme defined at the chiral limit  $\rightarrow$  need to generate  $N_f=4$  gauge ensembles close to the chiral limit and at the same lattice spacing as the three sets of  $N_f=2+1+1$  ensembles

G. Martinelli, C. Pittori, C. T. Sachrajda, M. Testa, and A. Vladikas, NPB445 (1995) arXiv:hep-lat/9411010

$\beta = 1.726$ $24^3 \times 48$			$\beta = 1.778$ $24^3 \times 48$			$\beta = 1.836$ $32^3 \times 64$		
$a\mu_{\text{sea}}$	$am_{PS}$	$am_{PCAC}$	$a\mu_{\text{sea}}$	$am_{PS}$	$am_{PCAC}$	$a\mu_{\text{sea}}$	$am_{PS}$	$am_{PCAC}$
0.0060	0.1689(15)	$-4.1(1.4) \times 10^{-4}$	0.0075	0.1748(15)	$-2.3(0.8) \times 10^{-5}$	0.0050	0.1276(14)	$-4.3(3.1) \times 10^{-5}$
0.0080	0.1905(11)	$-4.3(1.1) \times 10^{-5}$	0.0088	0.1871(18)	$-8.6(8.0) \times 10^{-5}$	0.0065	0.1447(14)	$+5.9(2.1) \times 10^{-5}$
0.0100	0.2155(12)	$+1.5(1.3) \times 10^{-4}$	0.0100	0.2006(18)	$-1.6(0.8) \times 10^{-4}$	0.0080	0.1585(14)	$+1.6(0.3) \times 10^{-4}$
0.0115	0.2289(12)	$+1.7(1.1) \times 10^{-4}$	0.0115	0.2158(11)	$+0.2(9.5) \times 10^{-5}$	0.0095	0.1744(12)	$+2.0(0.3) \times 10^{-4}$



- For the subtraction of the pion pole, a partially quenched approach is employed using additional values of the valence quarks at fixed sea quark mass
- We use two fitting Ansatzes for removing order  $(ap)^2$  discretisation effects and two  $p^2$ -ranges  $\rightarrow$  resulting  $Z_P$  all consistent - average over all of them
- Convert to  $\overline{MS}$  and evolve to 2 GeV or 3 GeV

# Scale setting

**Meson sector:** see also talk by P. Dimopoulos Thursday, July 29th, 05:30

- Use iso-symmetric values of pion mass and decay constant

C.A. *et al.* (2021), arXiv:2104.06747

$$m_{\pi}^{isoQCD} = 135.0(2) \text{ MeV} \quad \text{and} \quad f_{\pi}^{isoQCD} = 130.4(2) \text{ MeV}$$

FLAG report, S. Aoki *et al.*, Eur Phys. J. 77 (2017), arXiv:1607.00299

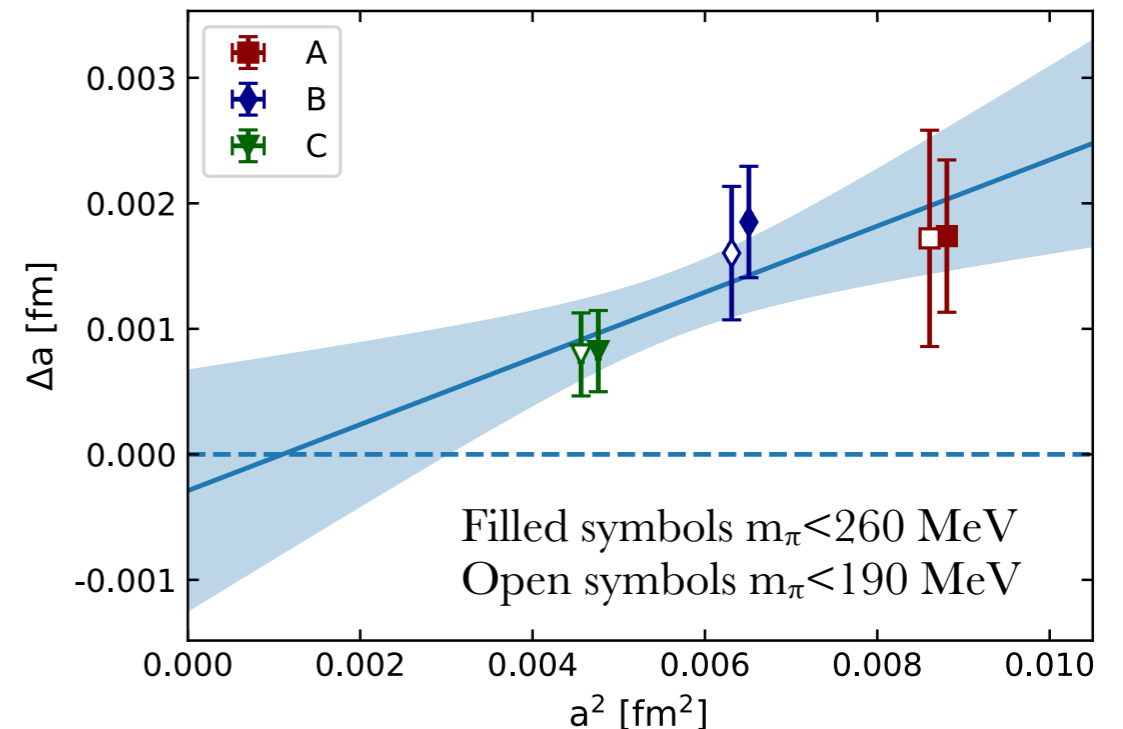
- NLO SU(2) chiral perturbation theory, correct for volume effects and take continuum limit of  $m_{\pi}$  and  $f_{\pi}$  in units of  $w_0$
- Determine  $w_0$  in the continuum limit and use  $w_0/a$  to extract  $a$

## Baryon sector:

- Use iso-symmetric values of pion and nucleon mass  $m_N^{isoQCD} = 0.9389 \text{ GeV}$
- SU(2) chiral perturbation theory to one-loop

$$(a_i m_N) = a_i m_N^0 - 4c_1 \frac{(a_i m_{\pi})^2}{a_i} - \frac{3g_A^2}{16\pi f_{\pi}^2} \frac{(a_i m_{\pi})^3}{a_i^2}$$

Fixed by physical value of the nucleon mass



The lattice spacings from meson and baryon sector agree in the continuum limit

# Determination of quark masses in the meson sector (I)

We use the lattice spacings determined in the meson sector

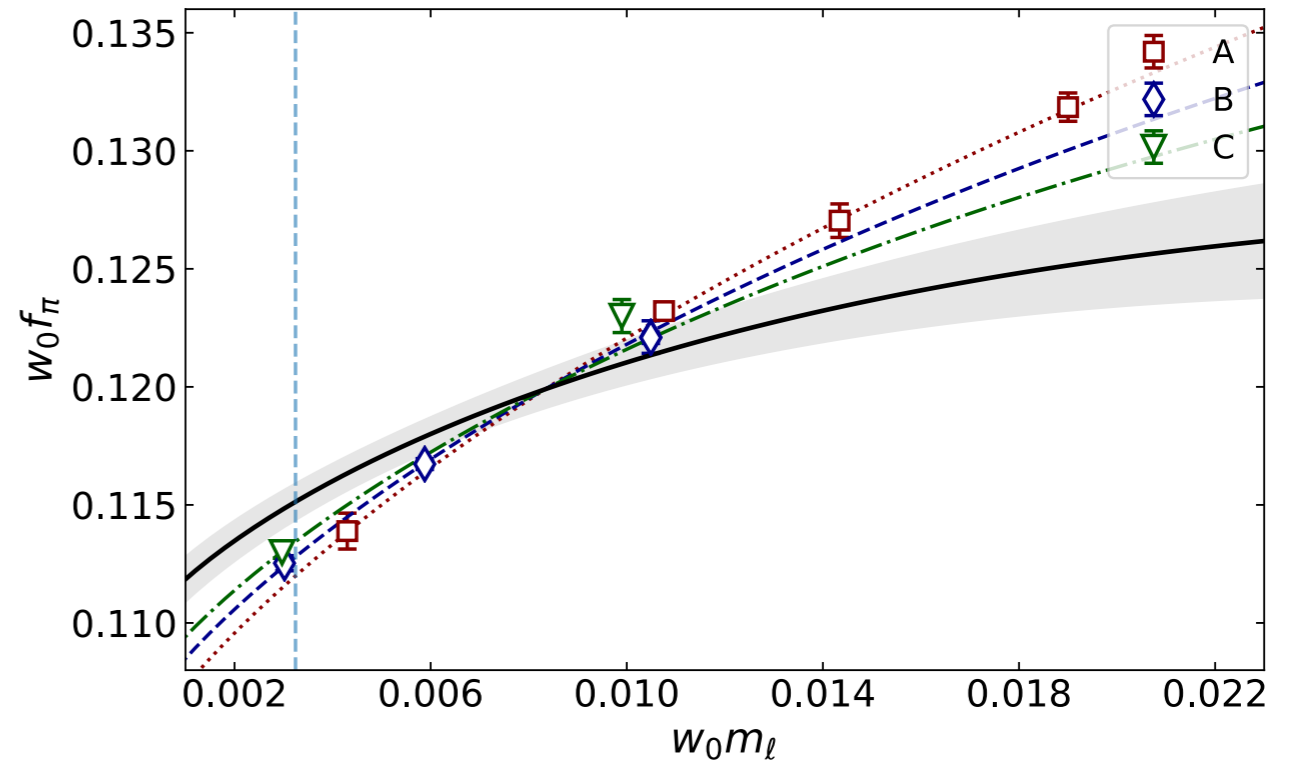
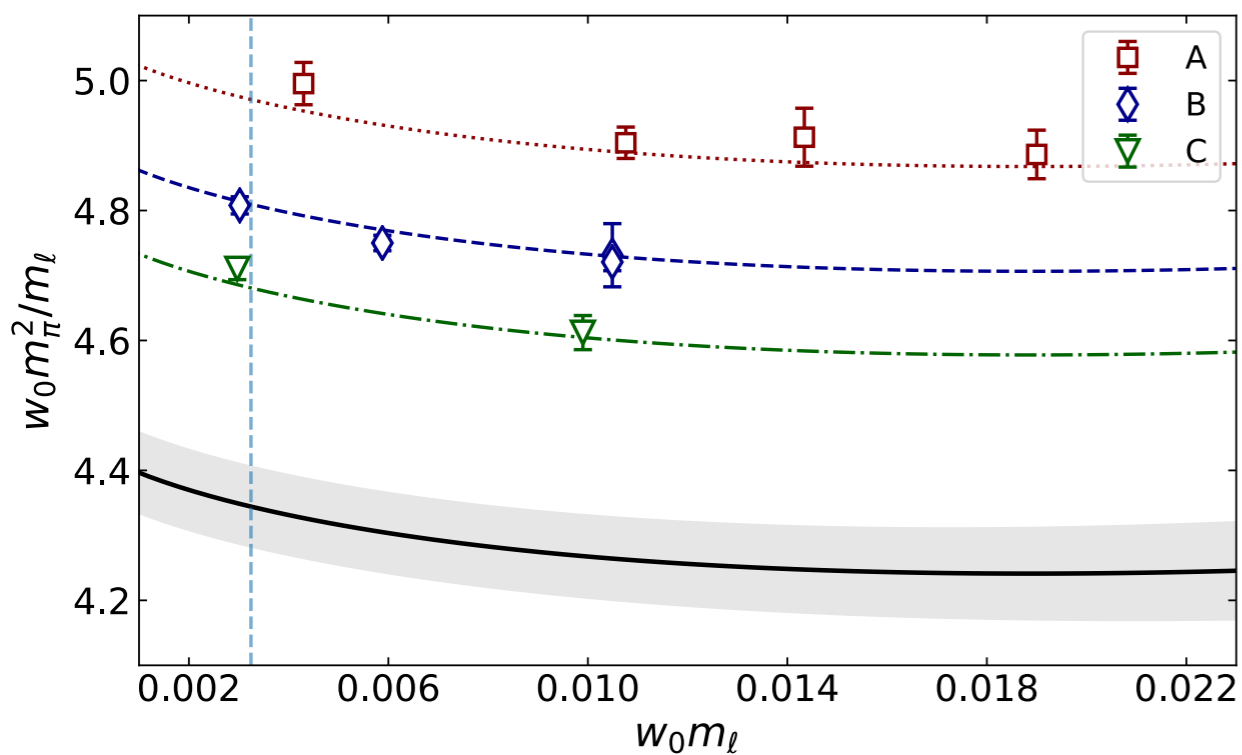
## Light (u/d) quark mass

$$(m_\pi w_0)^2 = 2(Bw_0)(m_\ell w_0) \left[ 1 + \xi_\ell \log \xi_\ell + P_1 \xi_\ell + P_2 a^2/w_0^2 \right] K_{M^2}^{FSE}$$

$$(f_\pi w_0) = (fw_0) \left[ 1 - 2\xi_\ell \log \xi_\ell + P_3 \xi_\ell + P_4 a^2/w_0^2 + a^2 m_\ell P_5 \right] K_f^{FSE}$$

Finite volume corrections

$$\xi_\ell = \frac{2Bm_\ell}{(4\pi f)^2}, \quad P_1 = -\bar{\ell}_3 - 2 \log(m_\pi^{\text{isoQCD}}/(4\pi f)), \quad P_3 = 2\bar{\ell}_4 + 4 \log(m_\pi^{\text{isoQCD}}/(4\pi f))$$

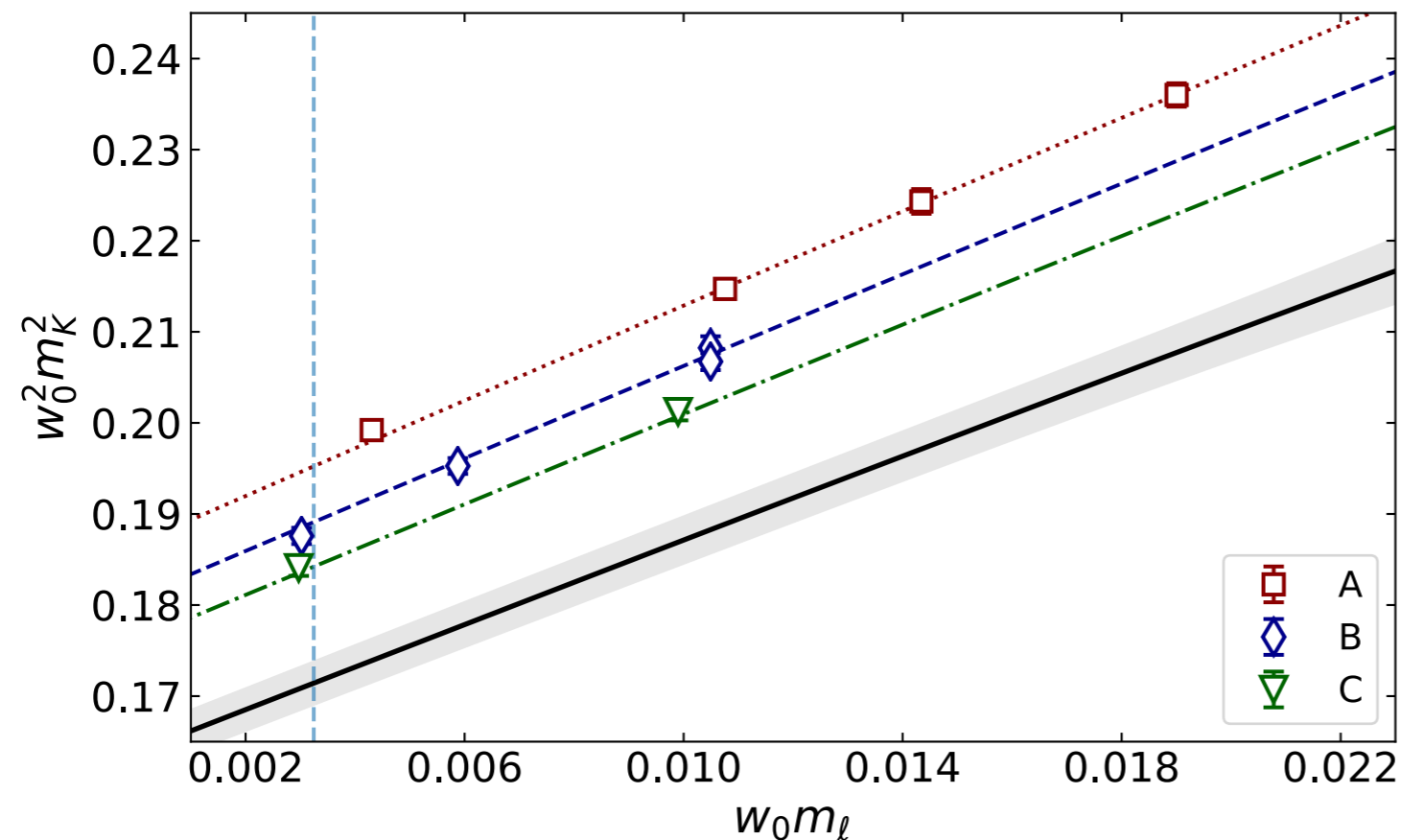


# Determination of quark masses in the meson sector (II)

## Strange quark mass

- Use as input  $m_K^{\text{isoQCD}} = 494.2(3)$  MeV
- Use three reference values of  $m_s$  and NLO ChPT

$$m_K^2 = a + b m_s w_0 \quad (m_K w_0)^2 = P_0 (m_\ell w_0 + m_s w_0) [1 + P_1 m_\ell w_0 + P_2 m_\ell^2 w_0^2 + P_3 a^2 / w_0^2]$$



# Determination of quark masses in the meson sector (III)

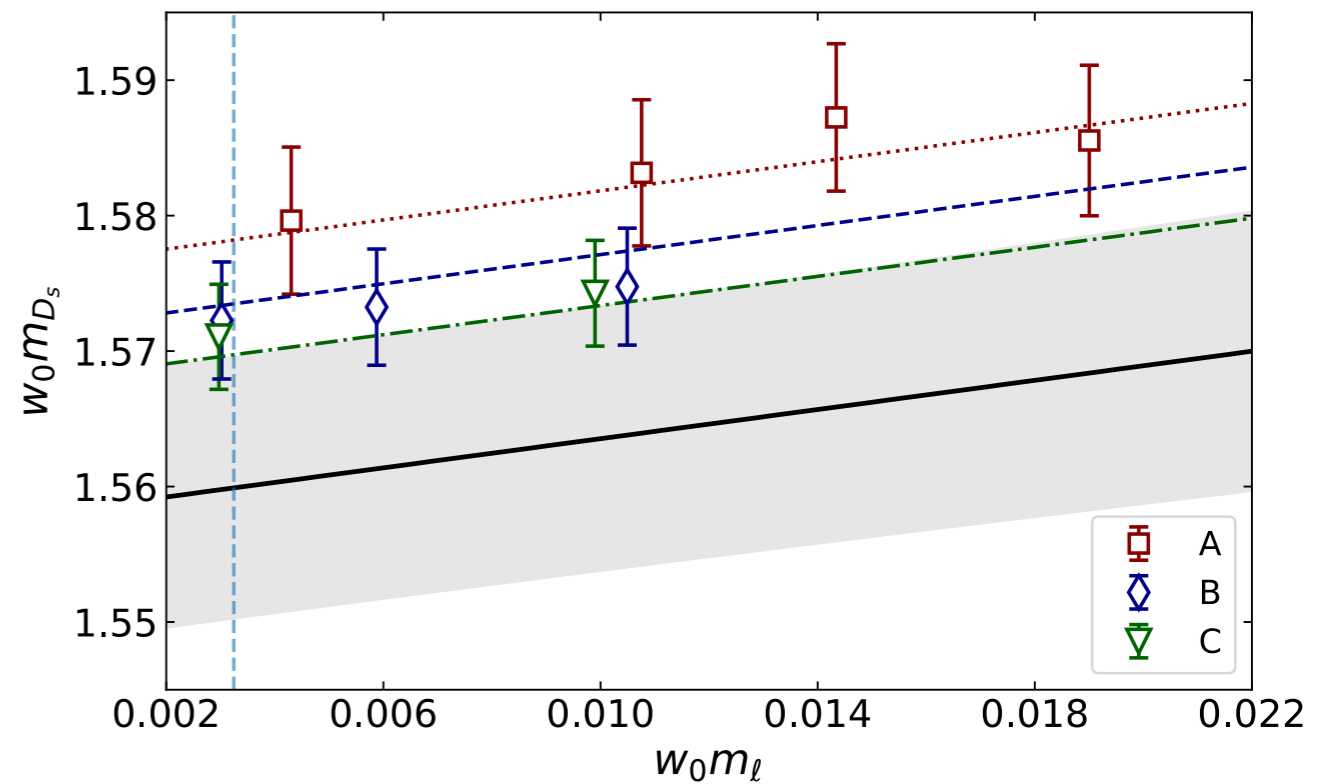
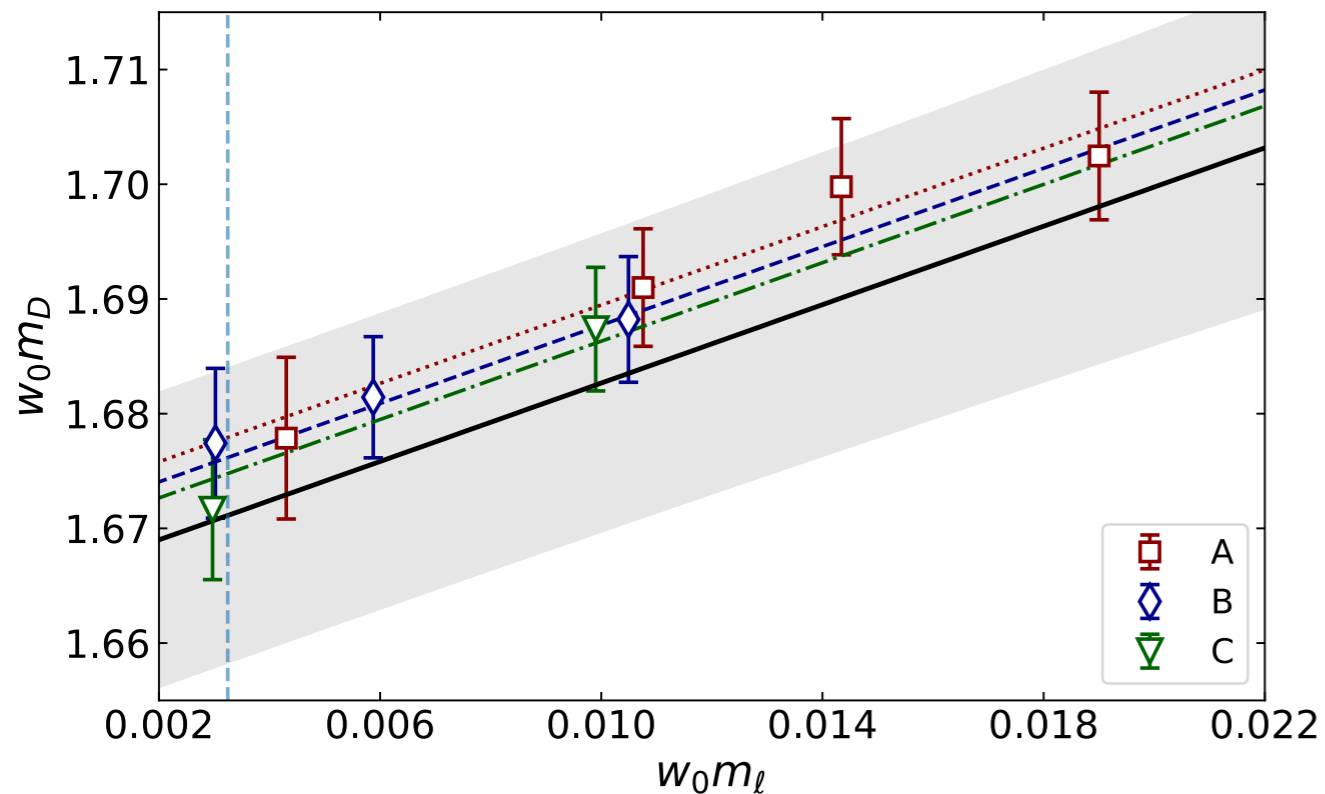
## Charm quark mass

- Use as input  $m_D^{\text{isoQCD}} = 1867.0(4)$  MeV ,  $m_{D_s}^{\text{isoQCD}} = 1969.0(4)$  MeV .

- Use three reference values of  $m_c$  and the polynomials

$$m_{D_s} = a + bm_c w_0 \quad m_D = P_0 + P_1 m_\ell w_0 + P_2 a^2 / w_0^2$$

$$m_{D_s} = P_0^s + P_1^s m_\ell w_0 + P_2^s a^2 / w_0^2$$





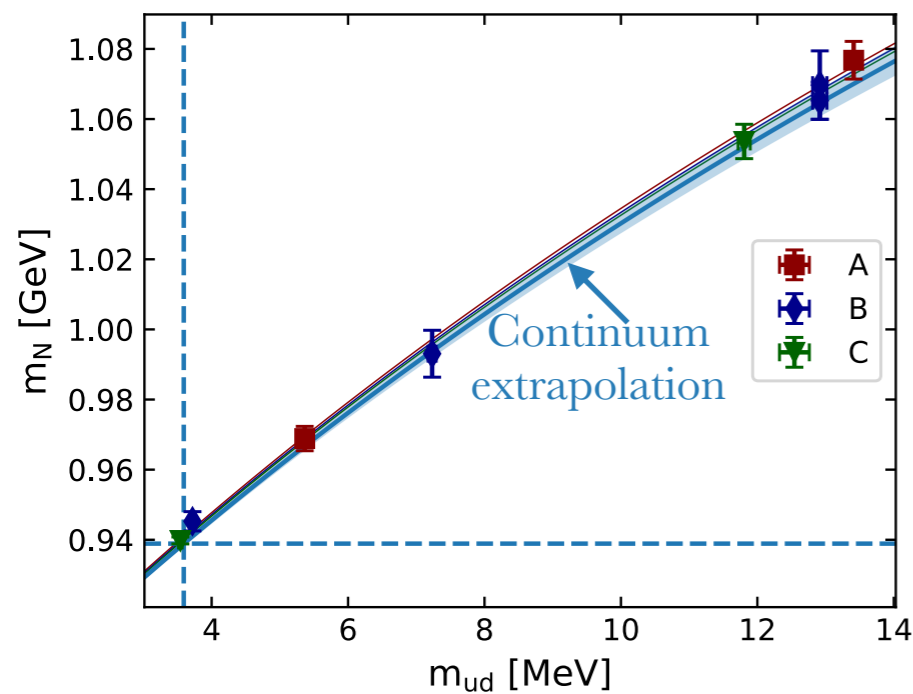
# Determination of quark masses in the baryon sector (I)

We use the lattice spacings determined from the nucleon mass

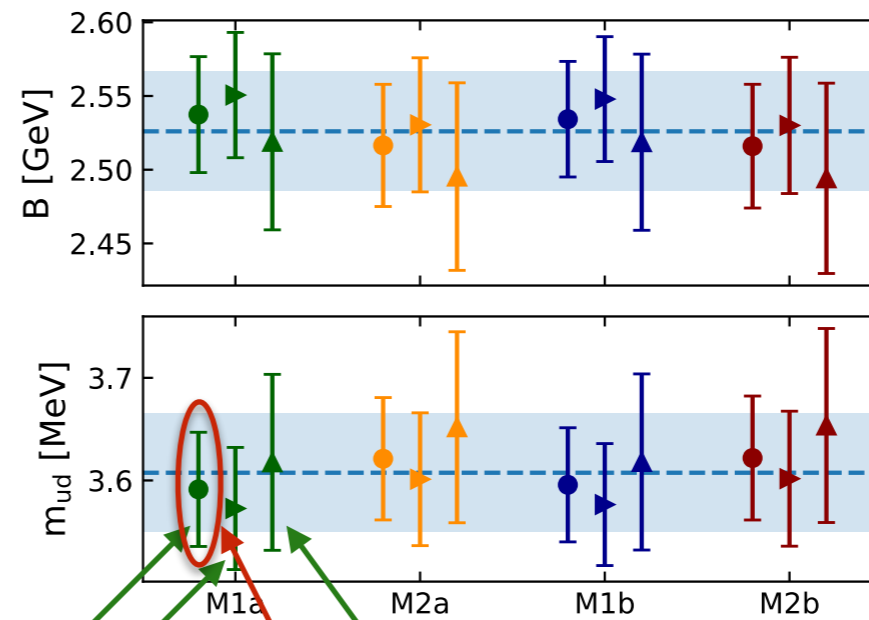
## Light (u/d) quark mass

To one-loop order  $m_\pi^2 = 2Bm_{ud}(1 + c_2a^2)$  i.e. two fit parameters B and  $c_2$

$$\longrightarrow m_N(m_{ud}) = m_N^0 - 4c_1 (2Bm_{ud}(1 + c_2a^2)) - \frac{3g_A^2}{16\pi f_\pi^2} (2Bm_{ud}(1 + c_2a^2))^{3/2}$$



## Dependence on the extraction of $Z_P$



Fit using  $m_\pi < 260$  MeV

Increase lower fit range by one for extraction of  $m_N$

Fit using  $m_\pi < 190$  MeV

take deviation from the mean of the other two as systematic error

Average over the values using the different  $Z_P$

# Determination of quark masses in the baryon sector (II)

## Strange quark mass

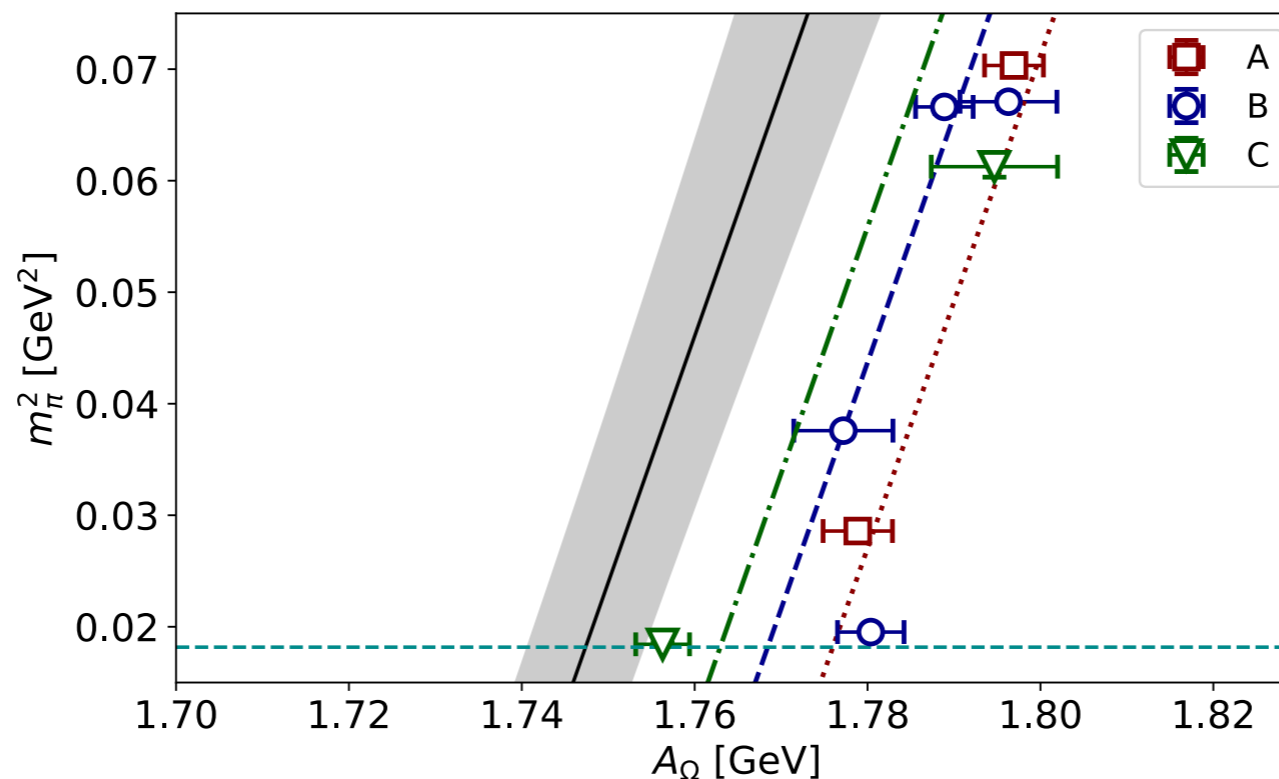
- Use as input the physical value of the  $\Omega$ -mass,  $m_{\Omega}^{(phys.)} = 1672.5(3)$  MeV and parameterise linearly its dependence on  $m_s$  around a reference mass  $\tilde{m}_s$

$$m_{\Omega} = A_{\Omega} + B_{\Omega} (m_s - \tilde{m}_s)$$

- Employ two fully consistent approaches:
  - Expand directly the parameters  $A_{\Omega}$  and  $B_{\Omega}$  and perform extrapolation in  $m_{\pi}$  and continuum limit

$$A_{\Omega}(a, m_{\pi}^2) = c_1 + c_2 m_{\pi}^2 + c_3 a^2, \quad B_{\Omega}(a, m_{\pi}^2) = c'_1 + c'_2 m_{\pi}^2$$

$$\longrightarrow m_s = \tilde{m}_s + \frac{m_{\Omega}^{(phys.)} - A_{\Omega}(0, m_{\pi}^{(phys.)})}{Z_P B_{\Omega}(0, m_{\pi}^{(phys.)})}$$



# Determination of quark masses in the baryon sector (II)

## Strange quark mass

- Use as input the physical value of the  $\Omega$ -mass,  $m_{\Omega}^{(\text{phys.})} = 1672.5(3)$  MeV and parameterise linearly its dependence on  $m_s$  around a reference mass  $\tilde{m}_s$

$$m_{\Omega} = A_{\Omega} + B_{\Omega} (m_s - \tilde{m}_s)$$

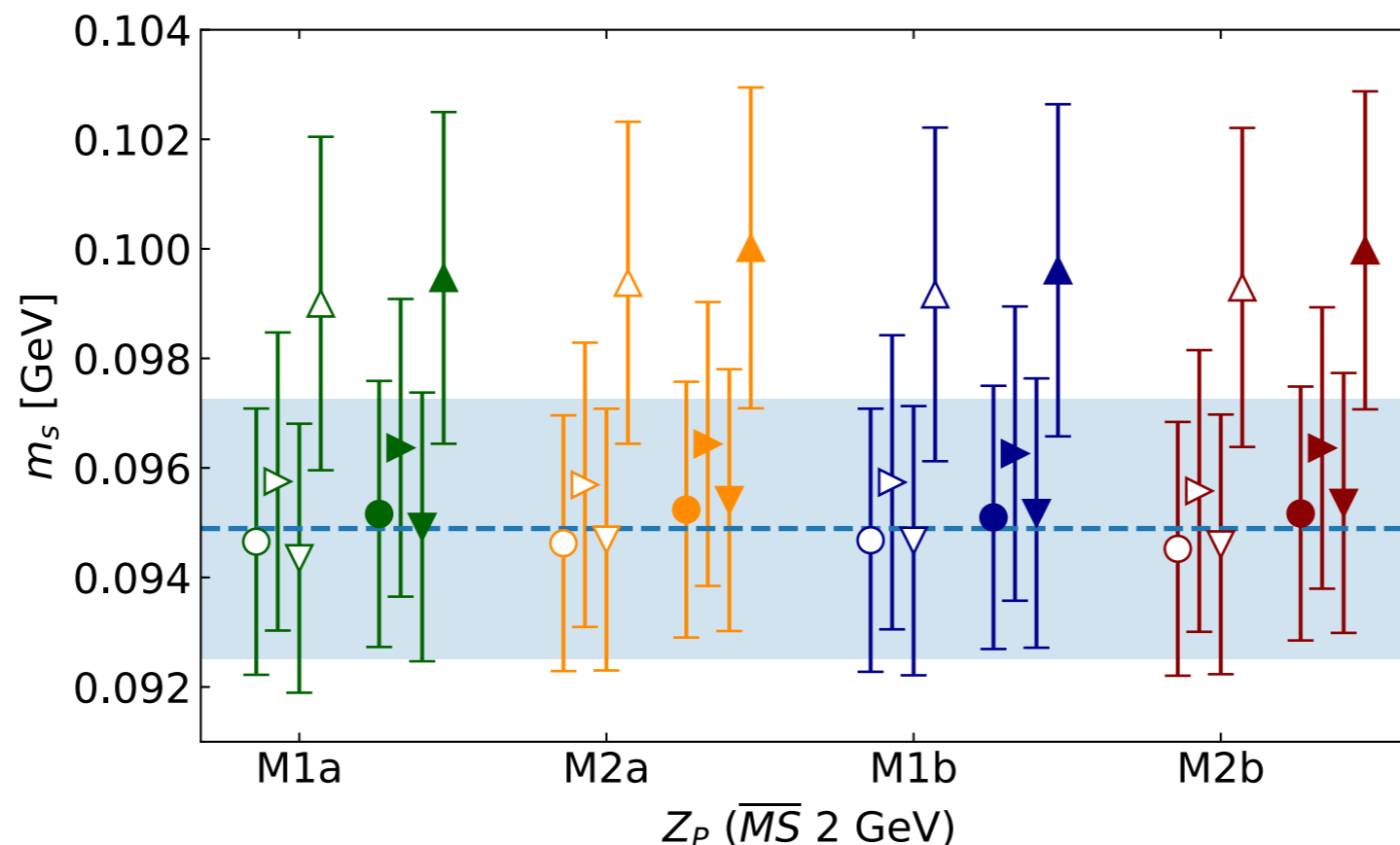
- Employ two fully consistent approaches:

II. adopt an iterative strategy:

- start by fixing a value of the renormalized strange quark mass  $m_s$  in physical units for all ensembles and extrapolate to the continuum limit and physical point using

$$m_{\Omega} = m_{\Omega}^{(0)} - 4c_{\Omega}^{(1)} m_{\pi}^2 + d_{\Omega}^{(2)} a^2$$

- iterate changing the value of  $m_s$  until the resulting value of the mass of  $\Omega$  at the physical point and continuum limit matches the physical value



# Determination of quark masses in the baryon sector (III)

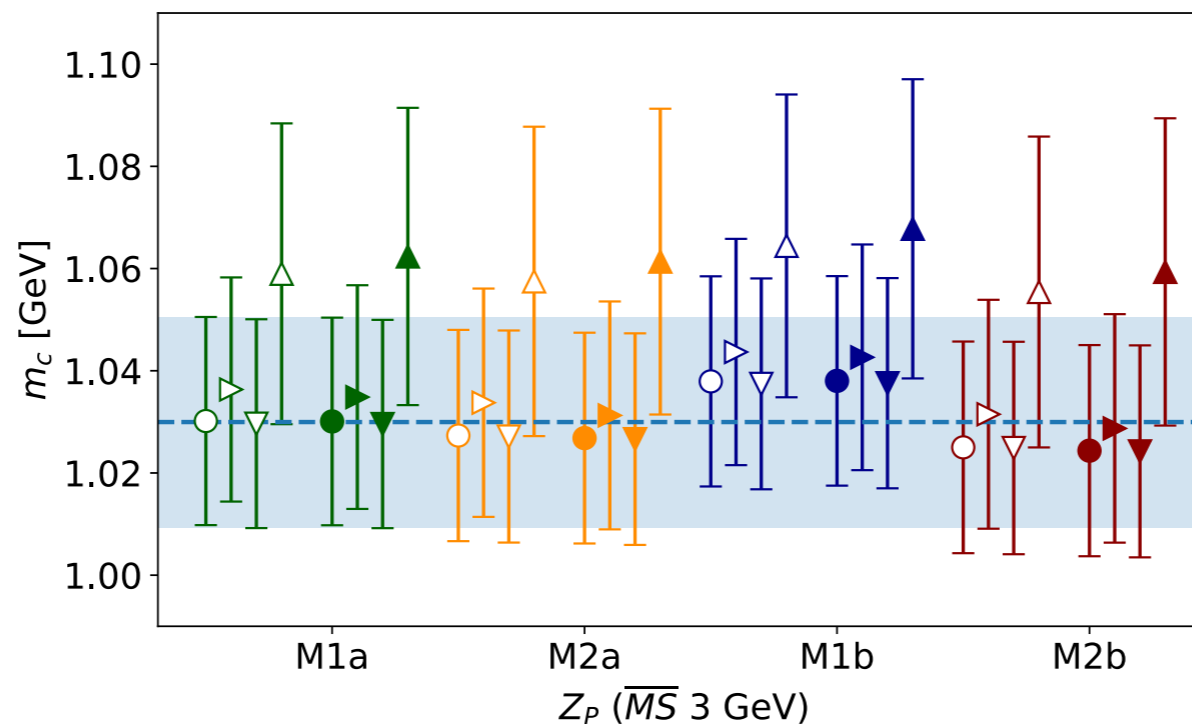
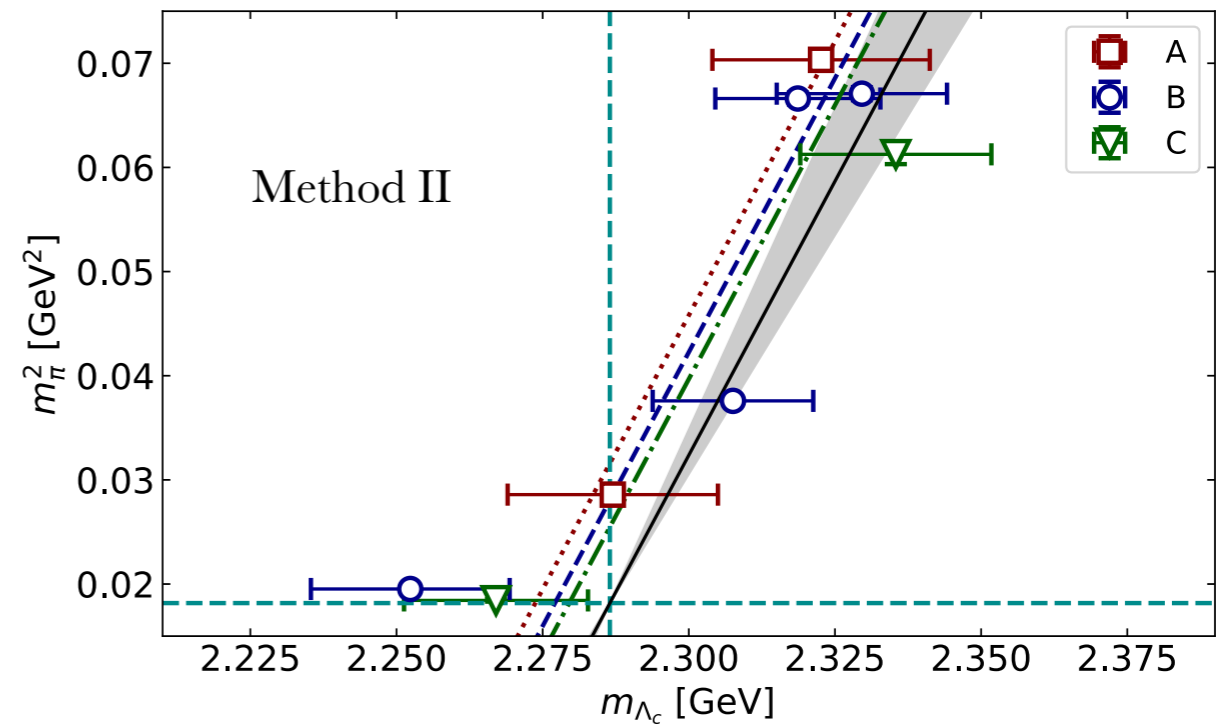
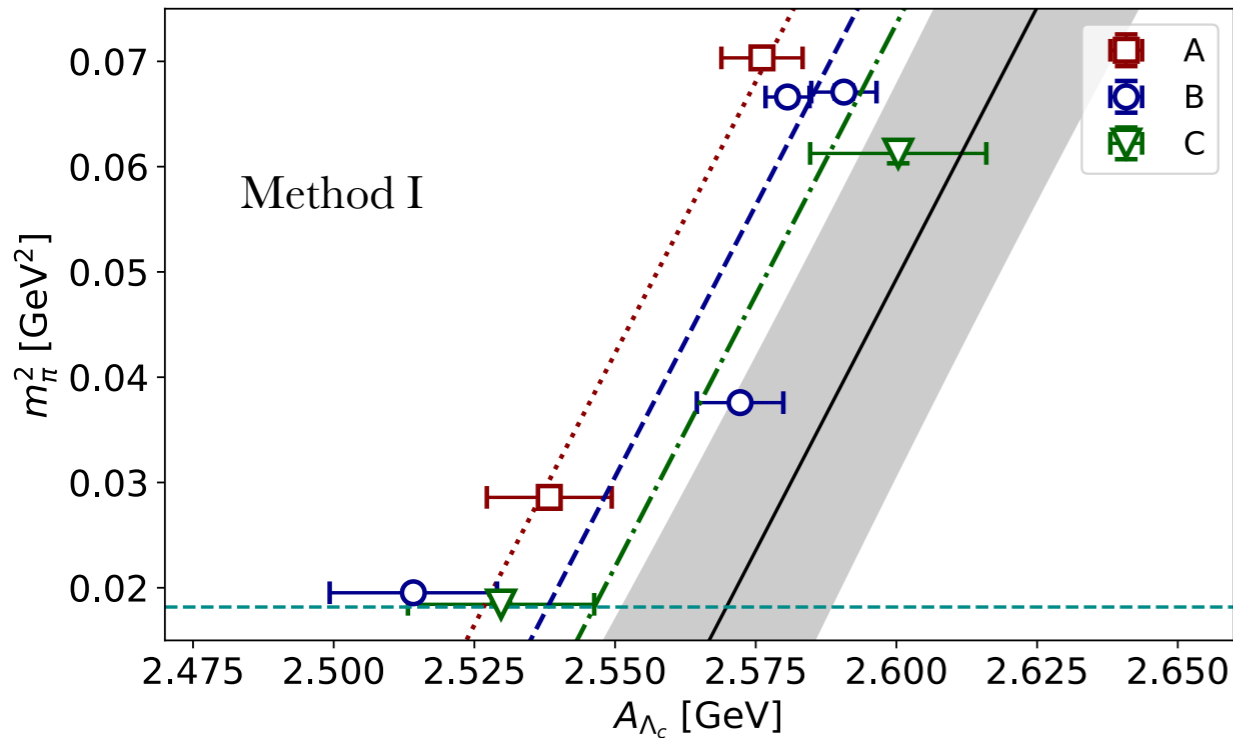
## Charm quark mass

- Use a similar approach to the one used for  $m_s$  with input the physical value of the mass of  $\Lambda_c$

$$m_{\Lambda_c}^{(phys.)} = 2286.5(1) \text{ MeV}$$

$$m_{\Lambda_c} = A_{\Lambda_c} + B_{\Lambda_c} (m_c - \tilde{m}_c)$$

$$m_{\Lambda_c} = m_{\Lambda_c}^{(0)} + c_{\Lambda_c}^{(1)} m_\pi^2 + d_{\Lambda_c}^{(2)} a^2$$



Fully consistent results from methods I (open symbols) & II (filled symbols)

# Results

- ✱ We take a weighted average over values from different  $Z_P$
- ✱ In the meson sector we average also over results from using different scale settings ( $w_0, t_0/w_0, \sqrt{t_0}$ ), ensembles for which  $m_\pi < 190$  MeV (setting  $P_3=0$ ), and using only two of the ensembles providing the  $\chi^2/\text{dof} < 2.5$ . For the charm we also average the results using D and D<sub>s</sub>

To compute the mean and error we use

$$\bar{x} = \sum_{i=1}^N w_i \bar{x}_i, \quad \sigma_{\text{stat}}^2 = \sum_{i=1}^N w_i \sigma_i^2, \quad \sigma_{\text{syst}}^2 = \sum_{i=1}^N w_i (\bar{x}_i - \bar{x})^2, \quad w_i = \frac{1/\sigma_i^2}{\sum_{j=1}^N 1/\sigma_j^2}$$

- ✱ In the baryon sector for the strange and charm quark masses we average over the results obtained from methods I and II. The systematic error is due to the chiral extrapolation and the fit range in extracting the input mass

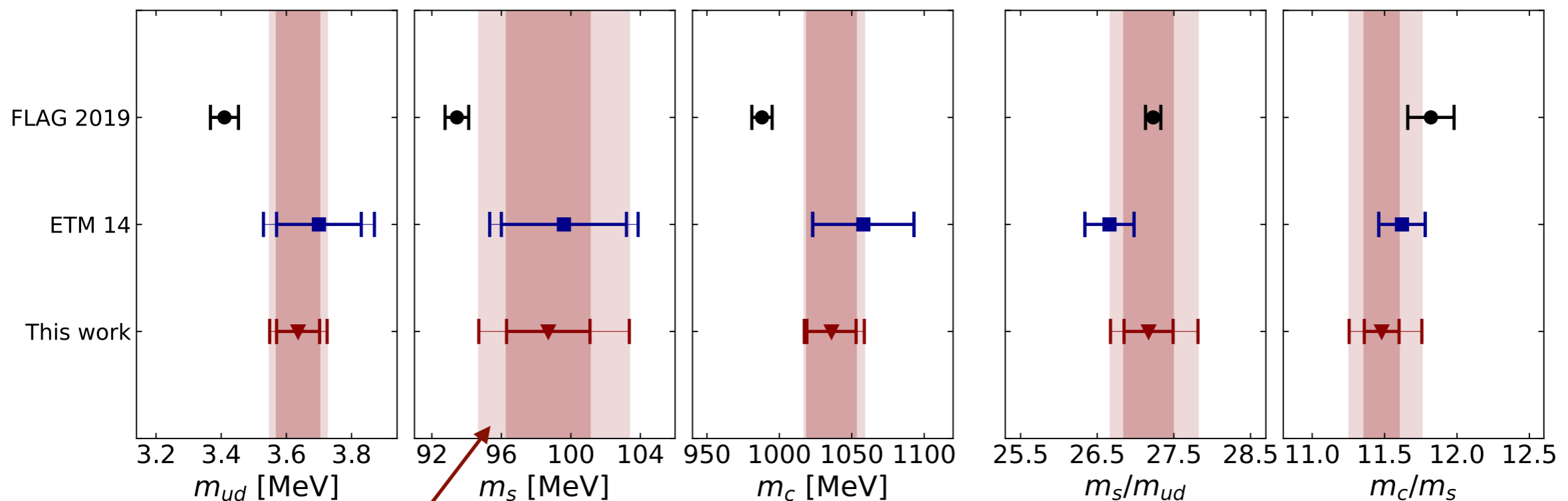
	$m_{ud}$ [MeV]	$m_s$ [MeV]	$m_c$ [MeV]	$m_s/m_{ud}$	$m_c/m_s$
Meson sector	3.689(80)(66)	101.0(1.9)(1.4)	1039(15)(8)	27.30(24)(14)	11.43(9)(10)
Baryon sector	3.608(58)( $^{+32}_{-19}$ )	94.9(2.4)( $^{+4.1}_{-1.0}$ )	1030(21)( $^{+22}_{-5}$ )	26.30(61)( $^{+1.17}_{-0.33}$ )	12.04(31)( $^{+58}_{-15}$ )

- ✱ Biggest tension between the meson and baryon determinations is seen for  $m_s$  for which we have a large systematic error in the baryon sector many due to the chiral extrapolation

# Conclusions

## \* Final values

Final results	$3.636(66)^{(+60)}_{(-57)}$	$98.7(2.4)^{(+4.0)}_{(-3.2)}$	$1036(17)^{(+15)}_{(-8)}$	$27.17(32)^{(+56)}_{(-38)}$	$11.48(12)^{(+25)}_{(-19)}$
FLAG 2019	3.410(43)	93.44(68)	988(7)	27.23(10)	11.82(16)



Large error due to the tension between the determinations in the meson and baryon sectors

## \* Next step: Add a third lattice spacing at the physical point



Would allow a continuum extrapolation directly at the physical point  
Avoiding chiral extrapolation is important for the baryon sector