

Renormalisation constants of quark bilinear operators in QCD with dynamical u, d, s and c quarks



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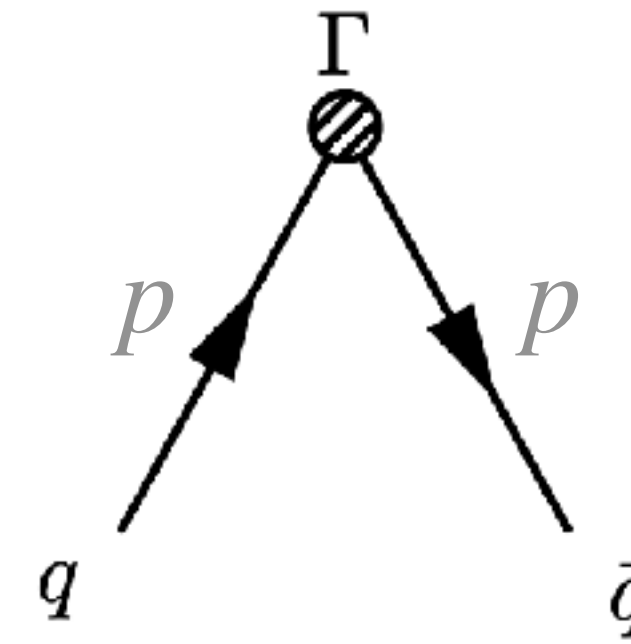
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Renormalisation in the RI'-MOM scheme

Numerical procedure:

- Building irreducible vertices from propagators
- Reduction of cut-off effects
 - ▶ selection of “democratic momenta”
 - ▶ subtraction of $O(g^2 a^\infty)$ effects [new]
- ⊙ Chiral extrapolation of vertices
- Perturbative evolution to a common scale
- ⊙ Extraction of the RCs through appropriate fit procedure

Two independent RI'-MOM analyses currently ongoing with complementary lattice setups and procedures



$$O_\Gamma = \bar{q}_1 \Gamma q_2$$

$$\Gamma = \begin{matrix} \mathbf{1}, & \gamma_5, & \gamma_\mu, & \gamma_\mu \gamma_5, & \sigma_{\mu\nu} \\ S & P & V & A & T \end{matrix}$$

$$\langle p | O_\Gamma^R | p \rangle \Big|_{p^2=\mu^2} = \langle p | O_\Gamma | p \rangle_{\text{tree}}$$

$$Z_q^{-1} Z_O \text{Tr}[\Lambda_O P_O] \Big|_{p^2=\mu^2} = 1$$

with

$$Z_q = \begin{cases} -\frac{i}{4} \text{Tr} \left[\frac{\partial S^{-1}(p)}{\partial \psi} \right]_{p^2=\mu^2} & \text{RI} \\ -\frac{i}{4} \text{Tr} \left[\frac{\not{p} S^{-1}(p)}{p^2} \right]_{p^2=\mu^2} & \text{RI}' \end{cases}$$

G. Martinelli et al. - NPB 445 (1995) 81

Hadronic contaminations in RI-MOM RC estimators

Let us study the **large** p^2 behaviour of the 2-fermion Green function ($p^2 \gg \Lambda_{\text{QCD}}^2$)

$$G_O(p^2) = \int d^4x d^4y e^{-ip(x-y)} \langle q_1(x) O_{12}^\Gamma(0) \bar{q}_2(y) \rangle$$

$$O_{12}^\Gamma(0) = \bar{q}_1(0) \Gamma q_2(0)$$

$$\sim S_q(p) \Lambda_O(p^2) S_q(p) + \sum_H w_H^\Gamma \frac{\Lambda_{\text{QCD}}^4}{(p^2)^2 m_H^2}$$

$$w_H^\Gamma \sim \frac{\langle 0 | O_{21}^\Gamma | H \rangle \langle H | O_{12}^\Gamma | 0 \rangle}{\Lambda_{\text{QCD}}^4} \quad (\text{dimensionless})$$

contr. from **quark** states $|x| \sim 1/|p|$ **OR** $|y| \sim 1/|p|$

 contr. from **hadronic** states $|x-y| \sim 1/|p|$
 $|x+y| \sim \text{large}$

 $S_q(p) \sim 1/\not{p}$ quark propagator

Hadronic contaminations in Λ_O are suppressed as $\mathcal{O}(\Lambda_{\text{QCD}}^2/p^2)$ but have **different chiral behaviour**:

$$\Gamma = \gamma_5$$

$$\frac{1}{p^2} \frac{\Lambda_{\text{QCD}}^4}{m_\pi^2}$$

Goldstone boson pole

$$\Gamma = \gamma_\mu \gamma_5$$

$$\frac{1}{p^2} \frac{m_\pi^2}{\Lambda_{\text{QCD}}^2} \frac{\Lambda_{\text{QCD}}^4}{m_\pi^2} \rightarrow \frac{\Lambda_{\text{QCD}}^2}{p^2}$$

$$\Gamma \neq \{\gamma_5, \gamma_\mu \gamma_5\}$$

$$\frac{\Lambda_{\text{QCD}}^2}{p^2}$$

G. Martinelli *et al.* - NPB 445 (1995) 81
+ ETMC (*in preparation*)

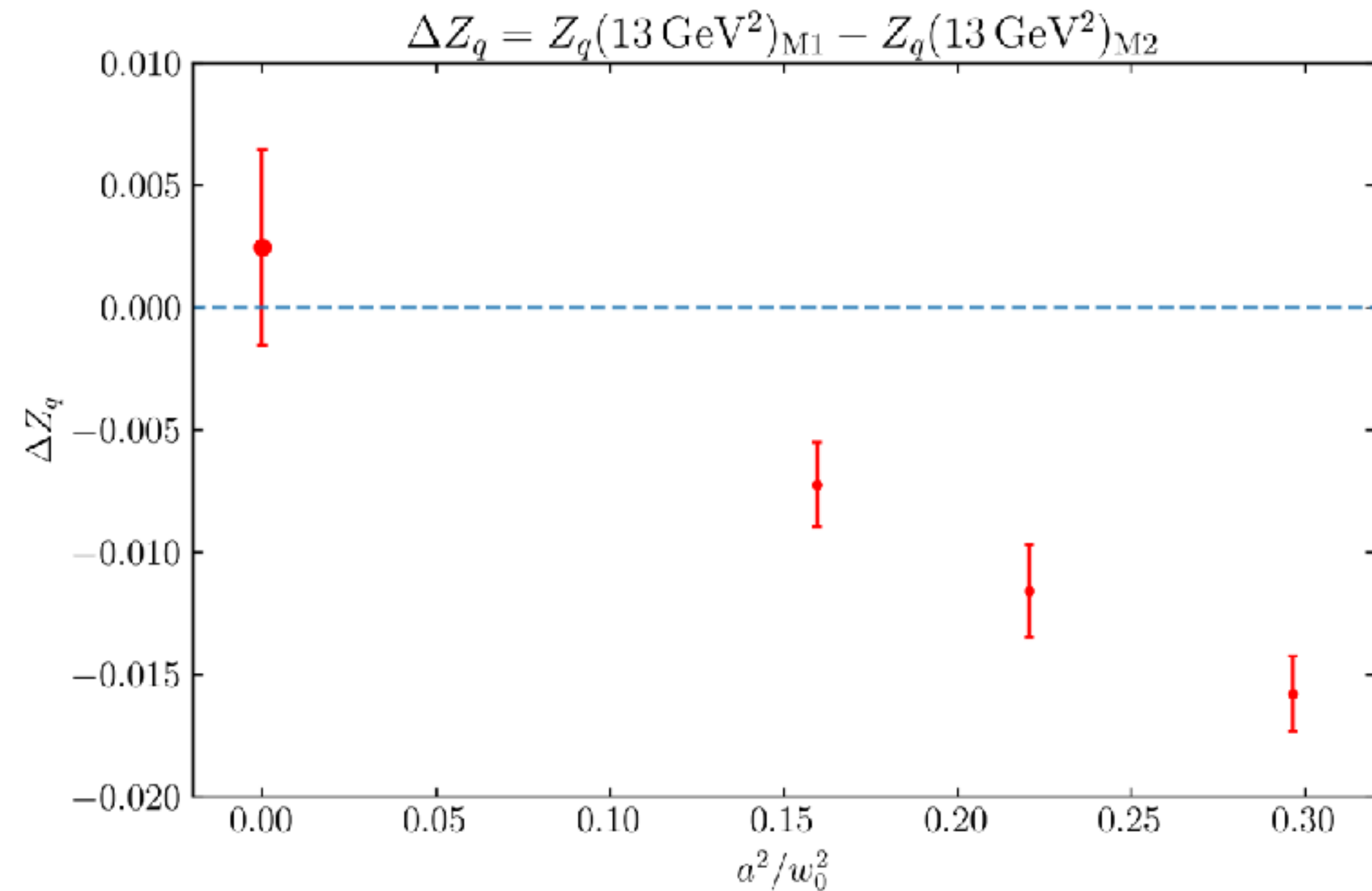
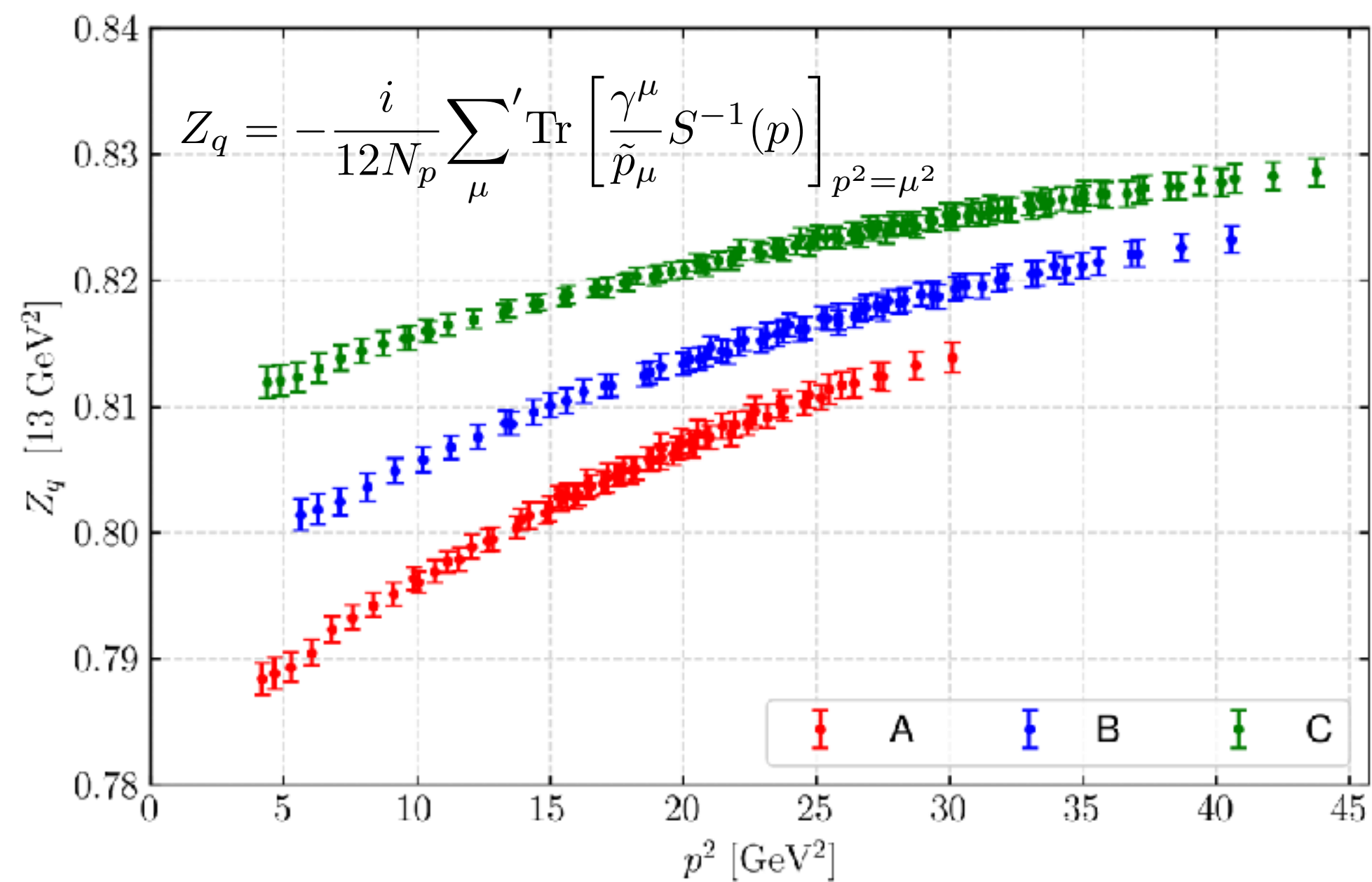
Determination of Z_q

- **no hadronic contaminations & mass dependence:** extremely mild (\Rightarrow constant fit)
- clean extraction using the following **fit methods:**

$$\zeta_q(p_{\text{ref}}^2; a^2 p^2) = \boxed{Z_q(p_{\text{ref}}^2)} + d_2^{(q)}(a^2 p^2)$$

M1: linear fit of ζ_O in $p^2 \in \mathcal{J}_{M1}(p_{\text{ref}}^2)$

M2: constant fit of ζ_O in $p^2 \in \mathcal{J}_{M2}(p_{\text{ref}}^2)$



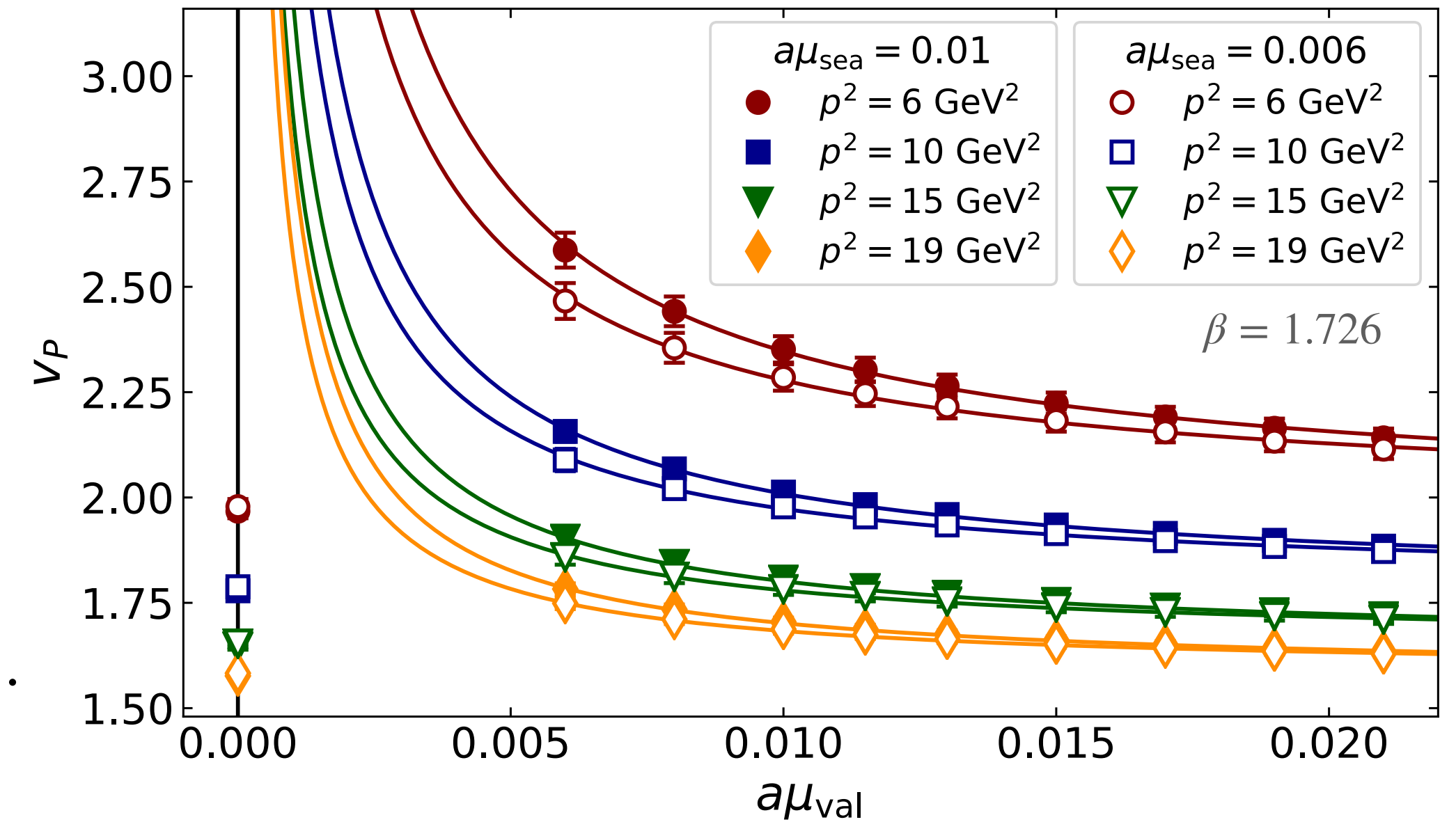
Determination of Z_P

Chiral extrapolation & the Goldstone pole

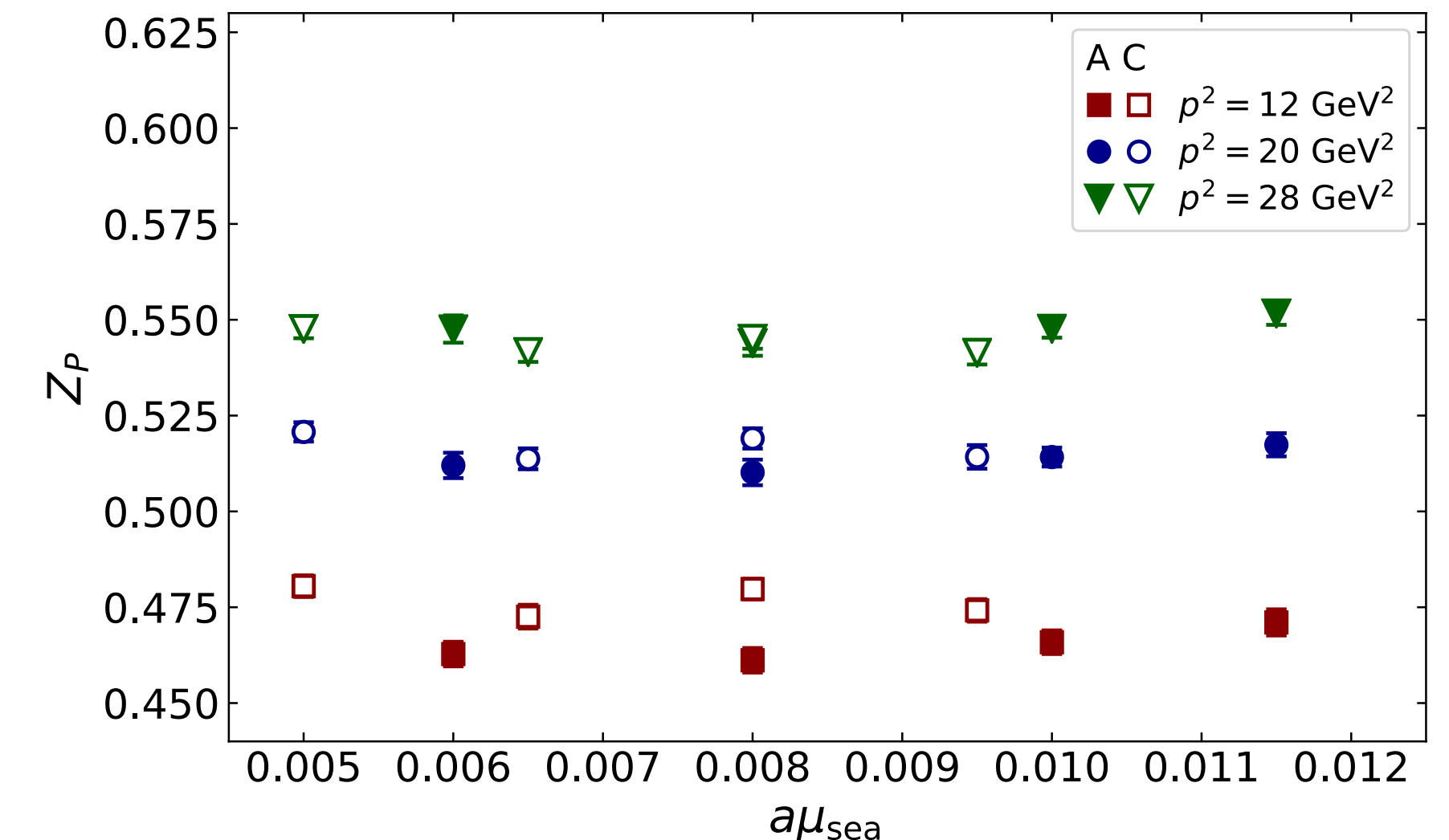
In our partially quenched setup $[m_\pi^2]_{\text{val}} \sim \mu_{\text{val}}$ (at LO in PQ- χ PT)

$$v_P(p^2, \mu_{\text{val}}, \mu_{\text{sea}}) = \underbrace{\mathcal{V}_P(p^2, \mu_{\text{sea}})}_{\text{cannot be disentangled}} + \underbrace{\frac{K'}{p^2}}_{\text{pion pole}} + \underbrace{\frac{K}{p^2} \frac{1}{\mu_{\text{val}}}}_{\text{pion pole}} + \underbrace{\frac{K''}{p^2} \mu_{\text{val}}}_{\text{mostly } \mathcal{O}(a^2)} + \dots$$

- ▶ **extrapolate** to $\mu_{\text{val}} = 0$ **without** including K'' in the fit:
 \Rightarrow absorb the $\mathcal{O}(a^2)$ [+ tiny $\mathcal{O}(a^0)$] into $Z_P \leftrightarrow$ better statistical precision
- ▶ residual dependence on $\mu_{\text{sea}} \leftrightarrow$ statistical fluctuations
 \Rightarrow **average** of $Z_P(\mu_{\text{sea}}, p^2)$ over μ_{sea}
- ▶ get estimate of $Z_P(p^2)$ in the (full) chiral limit



C. Alexandrou et al. (ETM) - arXiv: 2104.13408 [hep-lat]



Determination of Z_P

After chiral extrapolation & evolution

Residual dependence on p^2 due to lattice artefacts & possibly also to $O(a^0)$ residual hadronic contaminations $\sim 1/p^2$

Safety checks against hadronic contaminations:

- **Fitting** the pole

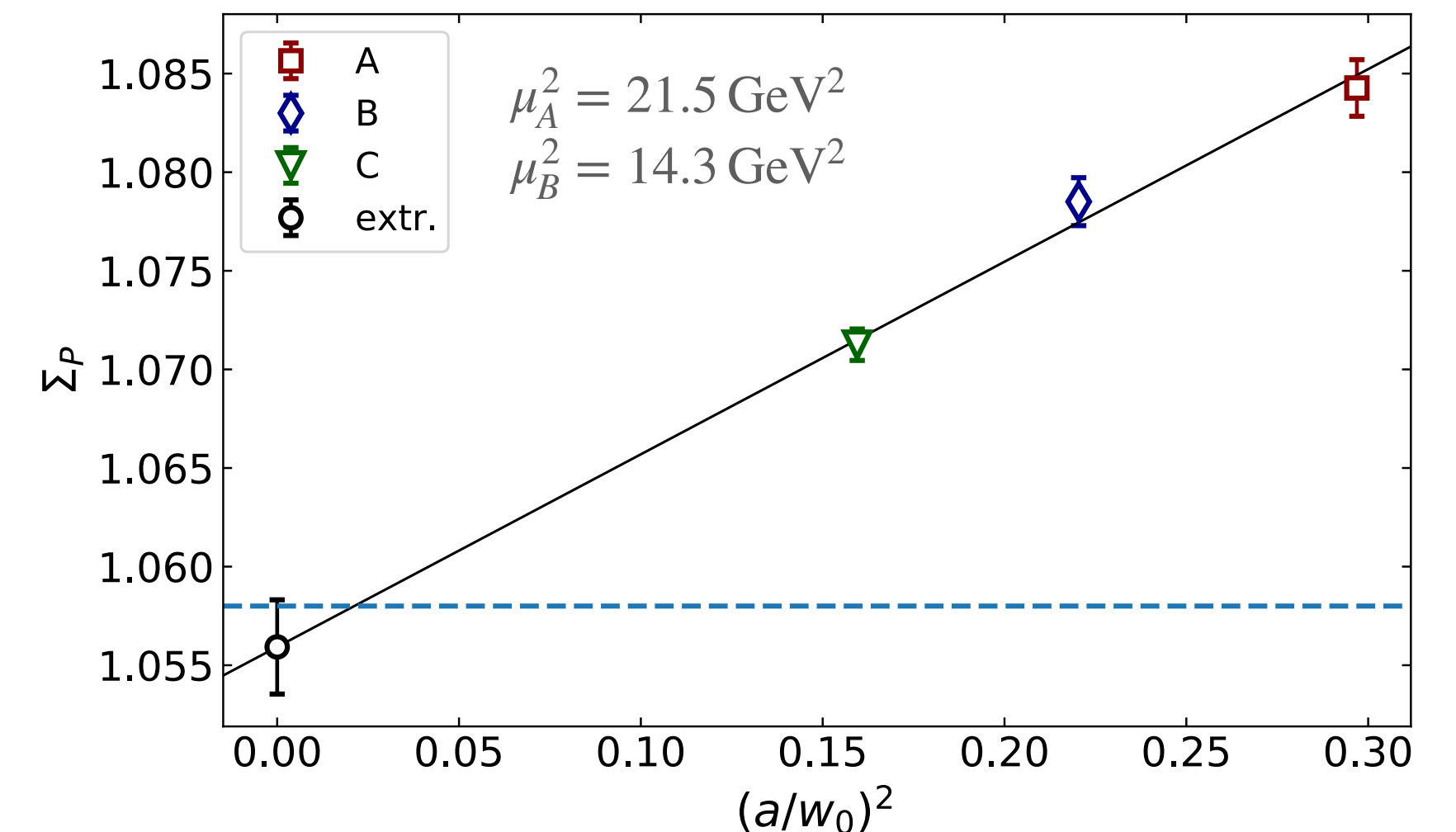
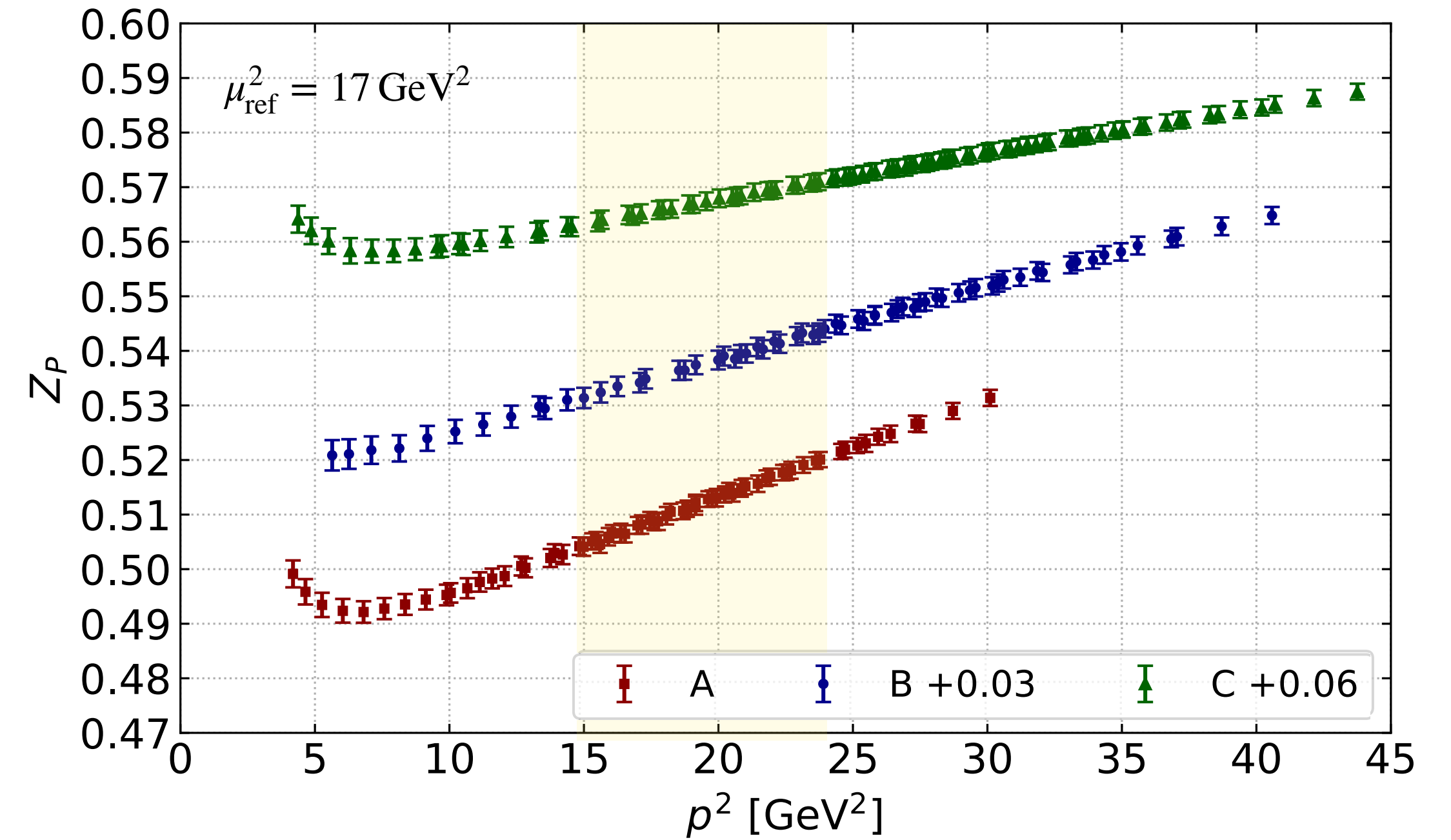
$$Z_P(\mu_{\text{ref}}^2) = z_0 + z_1 p^2 + \frac{z_{-1}}{p^2}$$

β	z_{-1}
1.726	-0.011 (43)
1.778	0.033 (42)
1.836	-0.021 (30)

\Rightarrow corresponding parameter z_{-1} compatible with **zero**

- **Step scaling function** $\Sigma_P(\mu_A^2, \mu_B^2) = Z_P(\mu_A^2)/Z_P(\mu_B^2)$

\Rightarrow small $O(a^2)$ and **agreement** with continuum PT (N³LO)



Determination of Z_P

Final results

RI'-MOM results:

RI'-MOM				
β	M1a	M2a	M1b	M2b
1.726	0.4774(24)	0.5079(24)	0.4917(26)	0.5301(24)
1.778	0.4812(32)	0.5042(26)	0.4944(27)	0.5255(23)
1.836	0.4899(26)	0.5053(23)	0.5046(27)	0.5240(24)

@17 GeV²
@21 GeV²
 $p^2 \in (15, 19) \text{ GeV}^2$
 $p^2 \in (18, 24) \text{ GeV}^2$

Perturbative evolution intermediate scale:

RI'-MOM, $\mu_{\text{ref}}^2 = 19 \text{ GeV}^2$				
β	M1a	M2a	M1b	M2b
1.726	0.4849(24)(35)	0.5159(24)(37)	0.4851(26)(32)	0.5229(24)(34)
1.778	0.4888(33)(35)	0.5121(26)(37)	0.4877(27)(32)	0.5184(23)(34)
1.836	0.4976(26)(36)	0.5133(23)(37)	0.4978(27)(33)	0.5169(24)(34)

Perturbative conversion to $\overline{\text{MS}}$ scheme:

$\overline{\text{MS}}, \mu_{\text{ref}}^2 = 19 \text{ GeV}^2$				
β	M1a	M2a	M1b	M2b
1.726	0.569(3)(5)	0.605(3)(5)	0.569(3)(5)	0.614(3)(5)
1.778	0.574(4)(5)	0.601(3)(5)	0.572(3)(5)	0.608(3)(5)
1.836	0.584(3)(5)	0.602(3)(5)	0.584(3)(5)	0.607(3)(5)

Detailed analysis & results reported in recent ETMC paper on quark masses:

C. Alexandrou et al. (ETM) - arXiv: 2104.13408 [hep-lat]
 (see also **C. Alexandrou's talk** - Wed 28/7 - 13:00 local time)

Final error budget:

$$\boxed{\text{RI'-MOM error} \lesssim 0.5\%} + \boxed{\text{evolution \& conversion to } \overline{\text{MS}} \text{ } O(0.5\% - 0.9\%)}$$

- ⇒ Systematic uncertainty on **perturbative formula** for evolution and matching to $\overline{\text{MS}}$ dominant, but can be improved
- ⇒ Reduction in RI'-MOM error possible thanks to the **analysis discussed** and to the **subtraction of $O(g^2 a^\infty)$ effects**

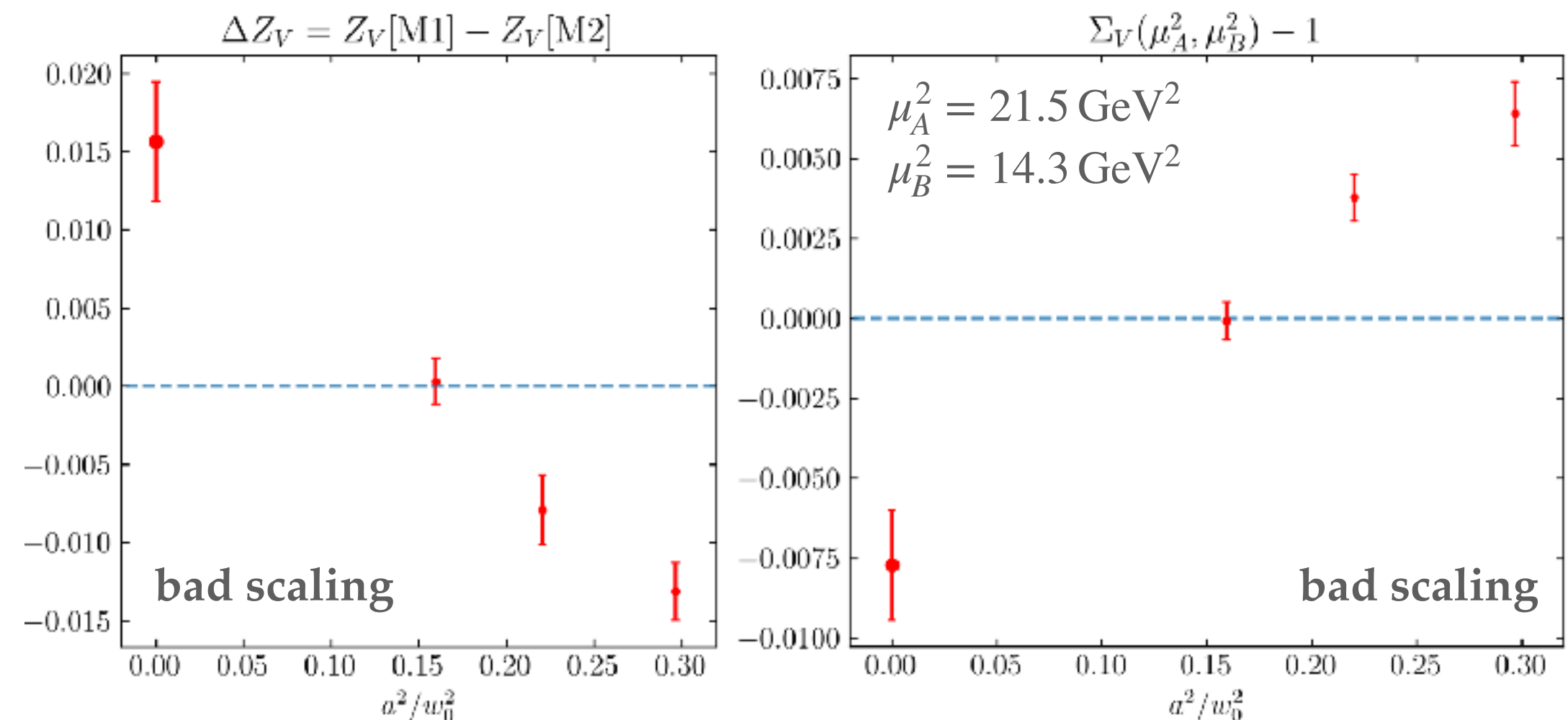
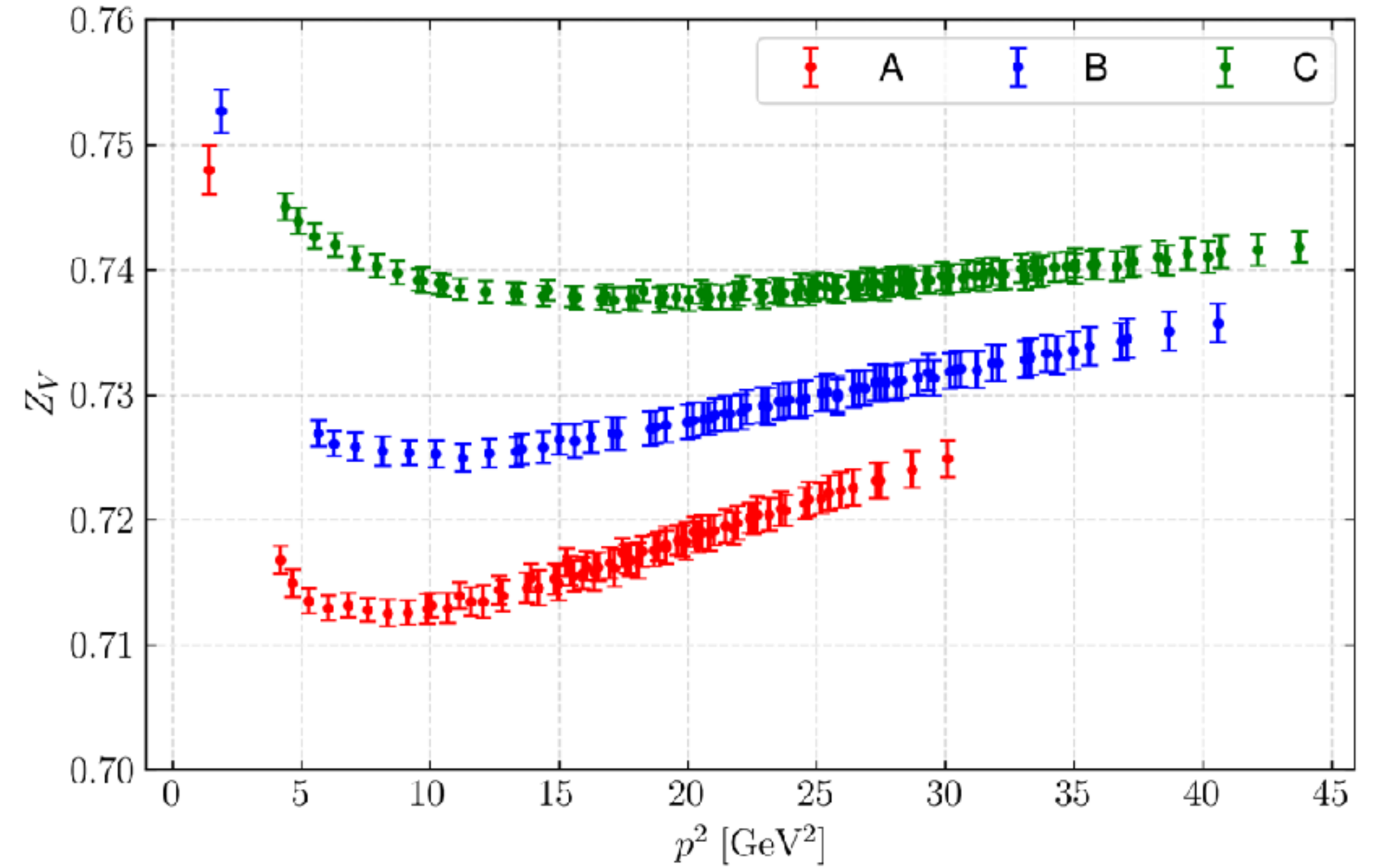
C. Alexandrou et al. (ETM) - PRD 95 (2017) 3

Determination of Z_V

$$v_V(p^2, \mu_{\text{val}}, \mu_{\text{sea}}) = \mathcal{V}_V(p^2, \mu_{\text{sea}}) + \frac{C'}{p^2} + \frac{C''}{p^2} \mu_{\text{val}} + \dots$$

- **Chiral extrapolation:** linear mass dependence tiny
 \Rightarrow constant fit in μ_{val} and μ_{sea} improve stat. uncertainty
 + tiny $O(a^0)$ mass dependence quoted as systematics ($\lesssim 1\%$)
- **Hadronic contaminations** $\sim 1/p^2$ are *finite* in the chiral limit and are **relevant** even at high p^2
 \Rightarrow methods **M1** and **M2** not efficient \leftrightarrow bad scaling!
 \Rightarrow need to fit directly the $\sim 1/p^2$ pole on a large range of p^2

$$\mathbf{M3} : \zeta_V(p^2) = Z_V + d_2^{(V)}(a^2 p^2) + \frac{\epsilon_V}{p^2}$$



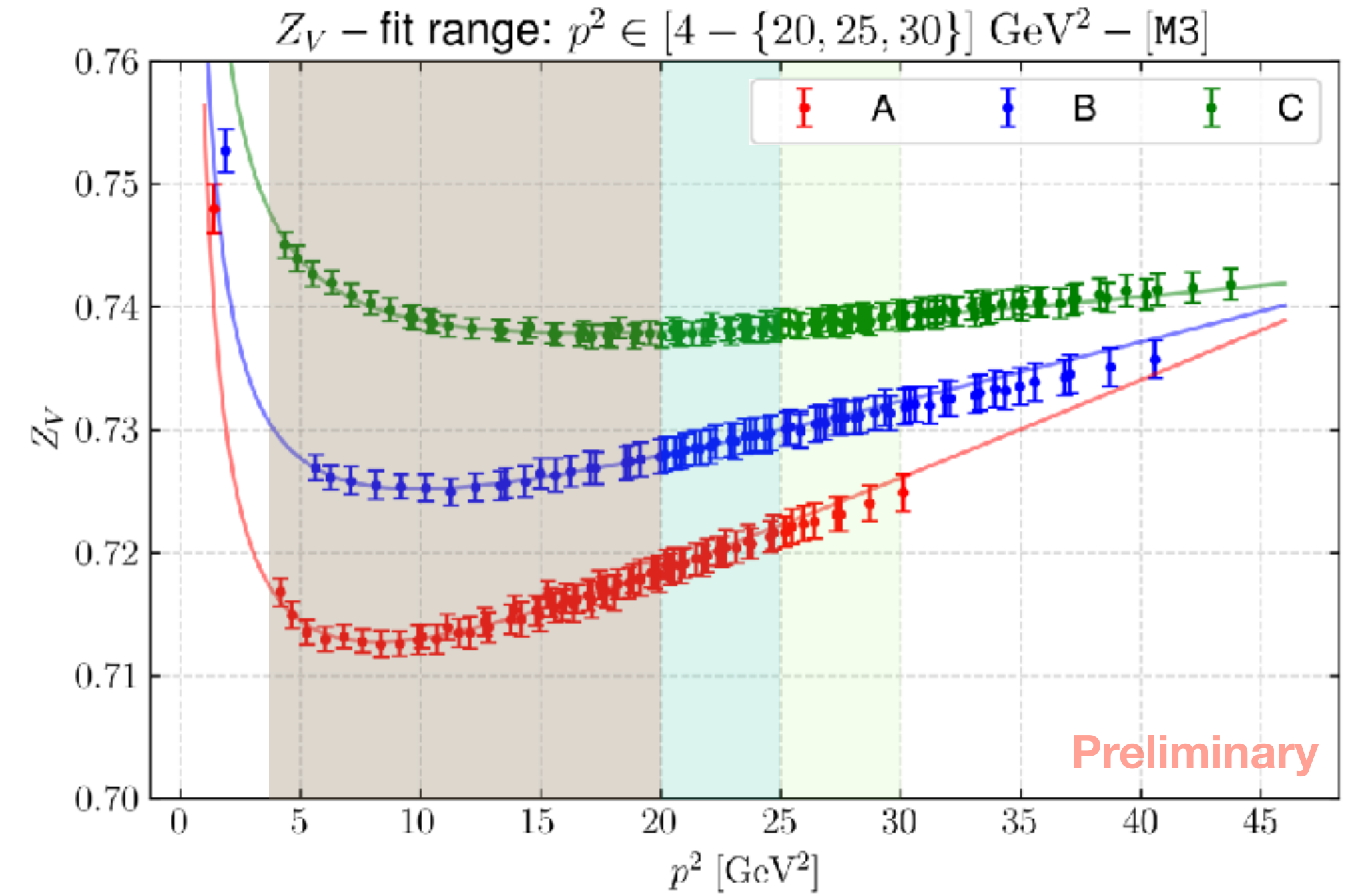
Determination of Z_V

Results & comparison with WTI

M3: $\zeta_V(p^2) = Z_V + d_2^{(V)}(a^2 p^2) + \frac{\epsilon_V}{p^2}$

Preliminary

β	Z_V	$d_2^{(V)}$	ϵ_V [GeV ²]
1.726	0.6989 (15) (10)	0.0038 (3)	0.057 (6)
1.778	0.7148 (22) (10)	0.0032 (4)	0.051 (11)
1.836	0.7306 (13) (10)	0.0018 (4)	0.061 (6)



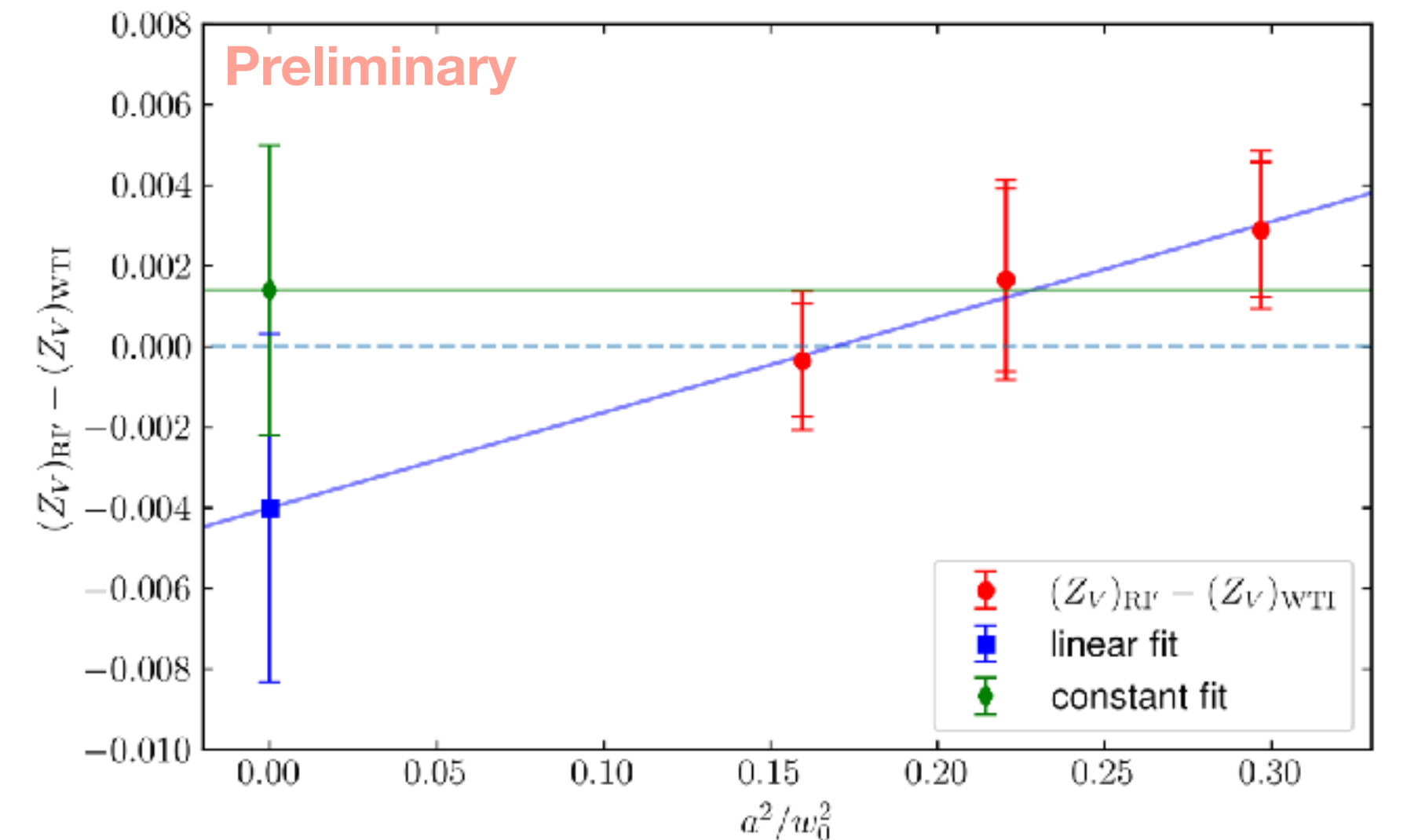
Z_V can also be evaluated via the **PCAC Ward-Takahashi identity**:

$$Z_V \sum_{\mathbf{x}} \langle \tilde{\partial}_0 V_0(t, \mathbf{x}) P_5^\dagger(0) \rangle^{(x)} = (\mu_1 + \mu_2) \sum_{\mathbf{x}} \langle P_5(t, \mathbf{x}) P_5^\dagger(0) \rangle^{(x)}$$

$$(A_R)_\mu = Z_V A_\mu^{(q)} = -i Z_V V_\mu^{(x)}$$

Preliminary

β	Z_V [WTI]
1.726	0.6960 (7)
1.778	0.7131 (6)
1.836	0.7310 (5)



Determination of Z_A/Z_V & comparison with alternative methods

M3: $\zeta_{A/V}(p^2) \simeq \frac{Z_A}{Z_V} + d_2^{(A/V)}(a^2 p^2) + \frac{\epsilon_{A/V}}{p^2}$

Z_A/Z_V can also be evaluated using different regularisations of the fermionic action:

$$\begin{cases} q_f = \exp(i\gamma_5 r_f \pi/4) \chi_f \\ \bar{q}_f = \bar{\chi}_f \exp(i\gamma_5 r_f \pi/4) \end{cases}$$

$$\langle \alpha | (O_\Gamma)_R | \beta \rangle = Z_{O_{\hat{\Gamma}}} \langle \alpha | O_{\hat{\Gamma}} | \beta \rangle^{\text{tm}} + \mathcal{O}(a^2) = Z_{O_{\hat{\Gamma}}} \langle \alpha | O_{\hat{\Gamma}} | \beta \rangle^{\text{OS}} + \mathcal{O}(a^2)$$

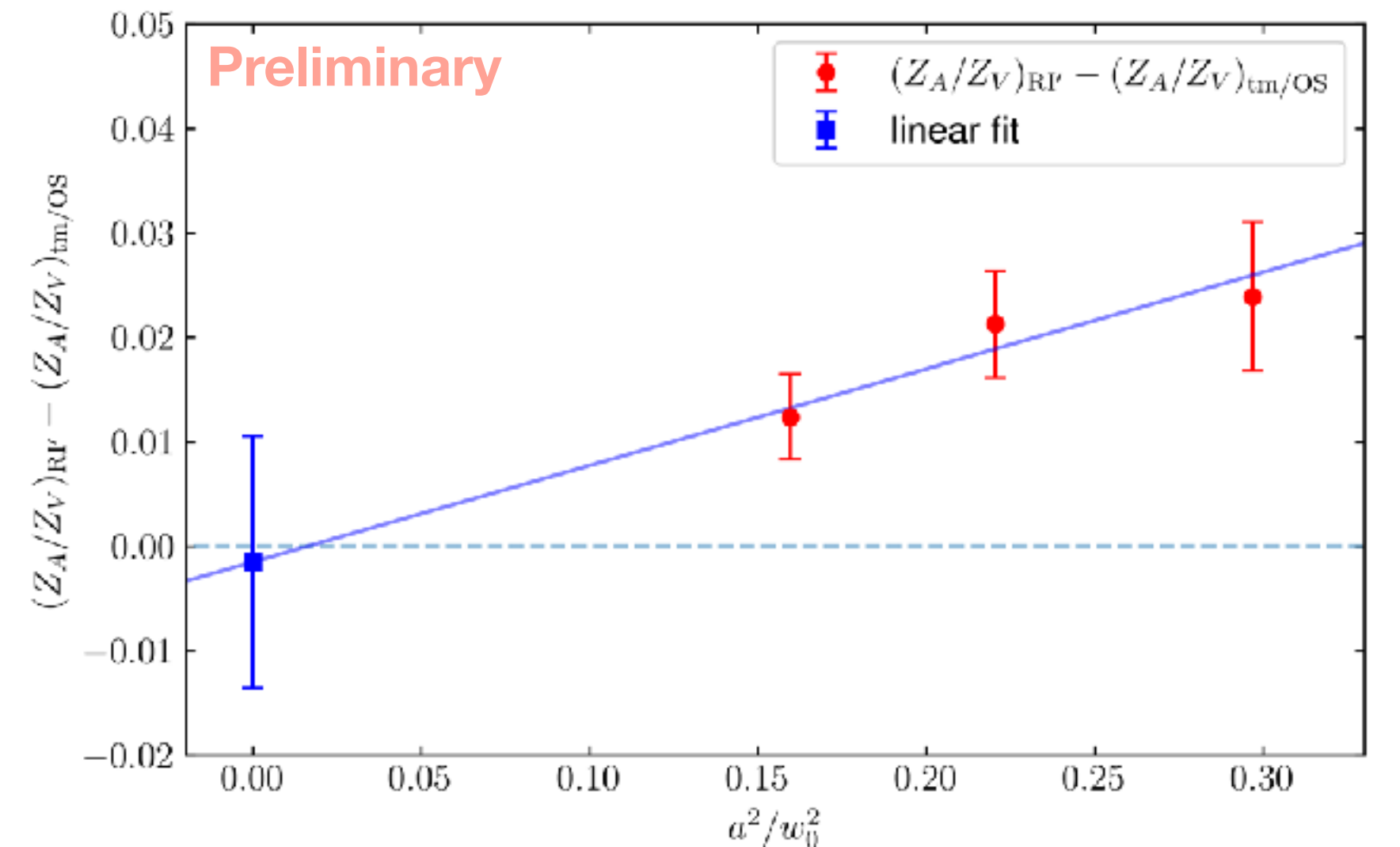
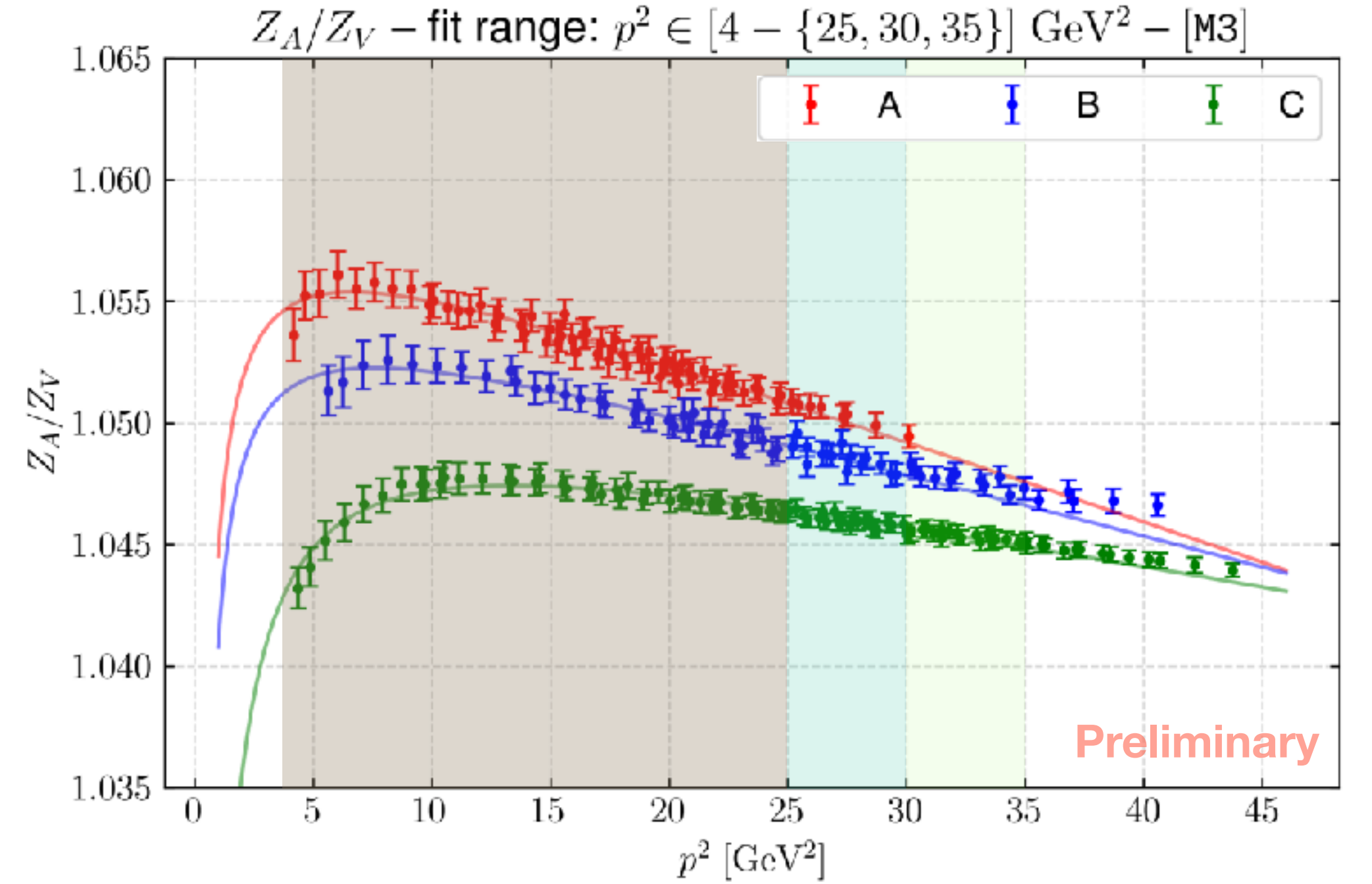
$r_u = -r_d = 1$ $r_u = r_d = 1$

$$\langle 0 | (A_0)_R | \pi \rangle \Rightarrow Z_V \langle \alpha | A_0 | \beta \rangle^{\text{tm}} = Z_A \langle \alpha | A_0 | \beta \rangle^{\text{OS}} + \mathcal{O}(a^2)$$

$A_\mu = \bar{q}_u \gamma_\mu \gamma_5 q_d$

Preliminary

β	Z_A/Z_V [RI']	Z_A/Z_V [tm/OS]
1.726	1.0599 (13)	1.036 (7)
1.778	1.0563 (11)	1.035 (5)
1.836	1.0524 (8)	1.040 (4)



Indirect evaluation of Z_A and Z_S

Combining RI-MOM and alternative methods

We can extract Z_A and Z_S with higher precision by **combining** RI'-MOM results and the ones from alternative methods:

$$Z_A = \left(\frac{Z_A}{Z_V} \right)_{\text{RI}'} \cdot (Z_V)_{\text{WTI}}$$

$O(0.1\%)$ dominated by $(Z_A/Z_V)_{\text{RI}'}$

can be reduced using tm/OS method
& increasing statistics

$$Z_S = \left(\frac{Z_S}{Z_P} \right)_{\text{tm/OS}} \cdot (Z_P)_{\text{RI}'}$$

$O(1\%)$ dominated by $(Z_P)_{\text{RI}'}$

can be reduced improving
perturbative results

... and Z_T ?

In progress: accuracy to be established, depending on hadronic contaminations and PT-evolution known only at 2-loops

expected to be **$O(1\%)$**

Conclusions

- Reaching **few per mille precision** on RCs in the RI'-MOM scheme is possible!
- **Subtraction of $O(g^2 a^\infty)$ artefacts & careful control of hadronic contaminations** (less suppressed wrt RI-SMOM: *but please check there too!*) is crucial
- Combination with **alternative methods** for RGI RCs ($Z_V, Z_A/Z_V, Z_S/Z_P$) allows to reach even higher precision (using WTI & pushing correlator statistics - no hadron contam.)
- **Two independent RI'-MOM analyses** in the ETMC to better control systematics + cross-checks of raw data & intermediate steps [*publication in preparation*]

Thanks for the attention!

& thanks to all the ETM collaborators

in particular, for the material presented here:

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