

D→P Semileptonic Decays with Highly Improved Staggered Quarks

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Lattice 2021 — 29 July 2021



Outline

- Scope of the talk
- The gauge ensembles
- Setup of calculation
- Preliminary results
- Summary and next steps

Special thanks to my colleagues in the all-HISQ subgroup:

Carleton DeTar, Aida El-Khadra, Elvira Gámiz, Zech Gelzer,
Steve Gottlieb, Andreas Kronfeld, Andrew Lytle, Jim Simone

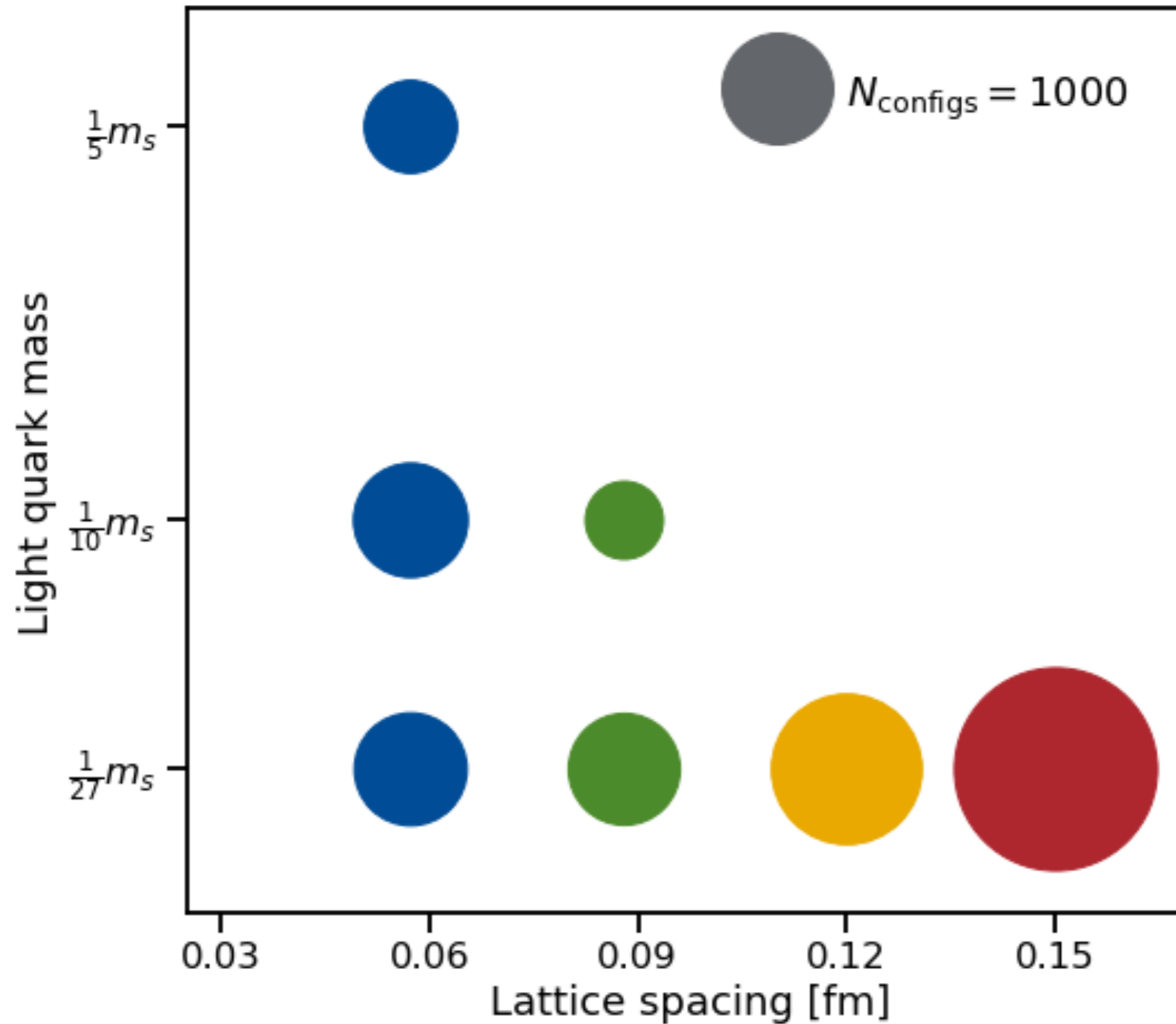


Scope: the all-HISQ campaign

- HISQ ensembles with $N_f=(2+1+1)$ dynamical sea quarks generated by the MILC collaboration
- Valence quarks:
 - Light and strange quarks match the sea
 - Heavy quarks: range from $0.9 m_c$ up to cutoff ($ma \sim 1$)
- **Campaign Goal:** form factors for decays of B, D mesons to pseudoscalars
 - ▶ D mesons: $D_{(s)} \rightarrow \pi, K$
 - ▶ B mesons: $B_{(s)} \rightarrow D_{(s)}, \pi, K$
- Eventual target: lattice spacings from 0.15 fm—0.03 fm
- **This talk:** preliminary results for
 - ▶ 4 lattice spacings: 0.15, 0.12, 0.09, 0.06 fm
 - ▶ Decays of D mesons only
- For B decays: see Andrew Lytle's talk in this session
- All 3pt functions are fully blinded

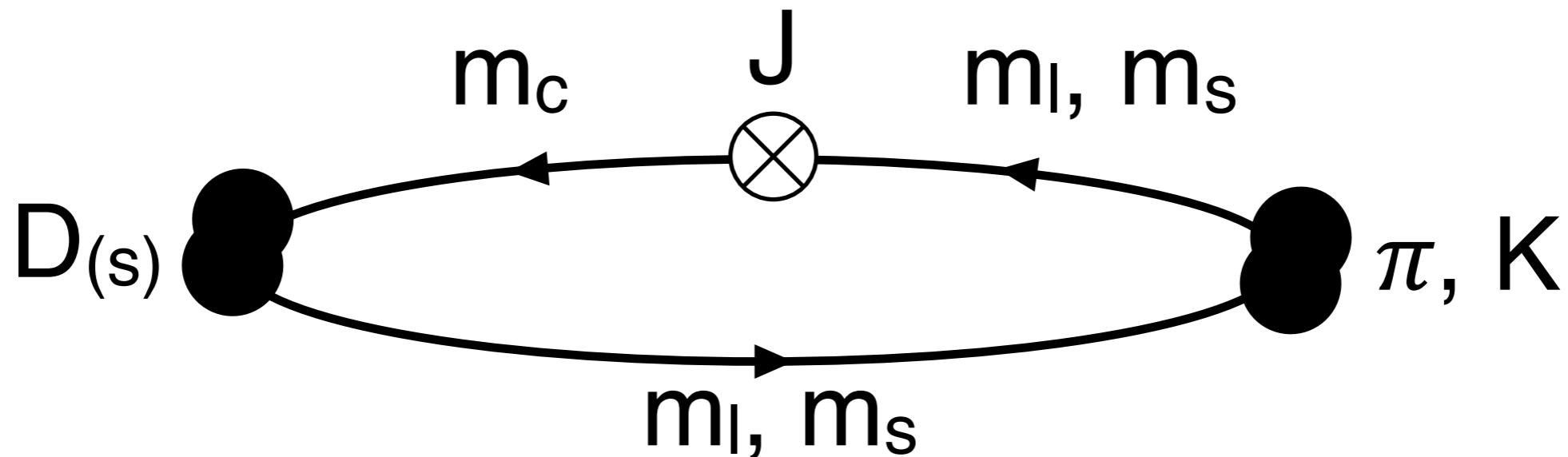


The gauge ensembles





3pt correlators



- We compute scalar, vector, and tensor currents
- Kinematic setup:
 - Work in rest frame of mother hadron $D_{(s)}$
 - Vary the energy of the recoiling daughter hadron π, K

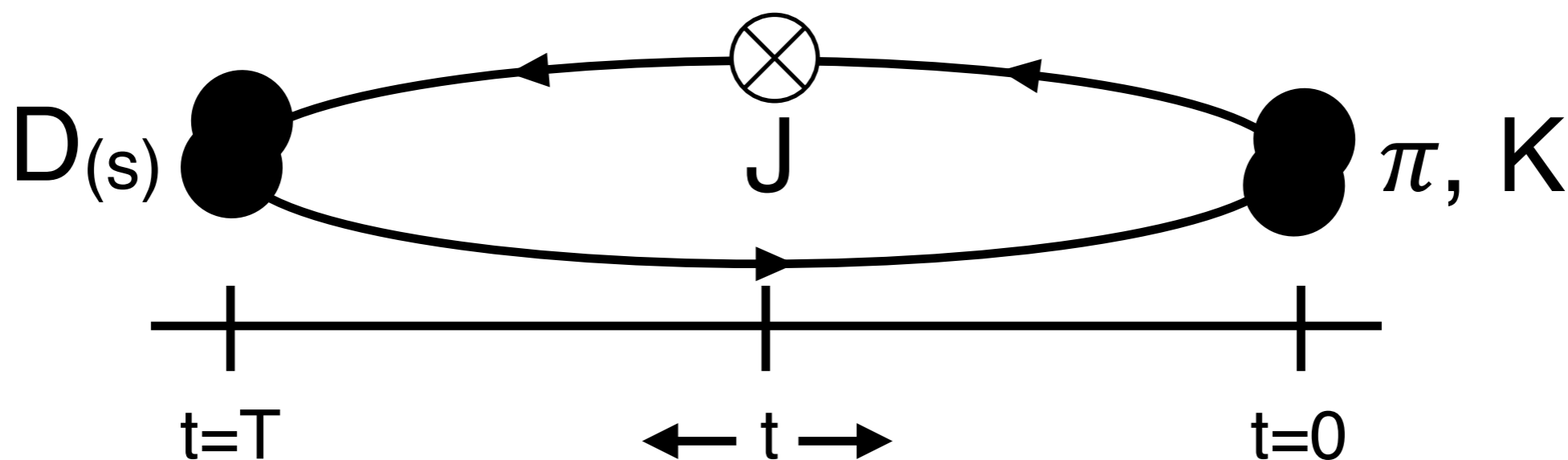


3pt correlators

- As usual, the spectral decomposition reads:

$$\begin{aligned}
 C_3(t; T) &= \langle \mathcal{O}_D(T) J(t) \mathcal{O}_\pi(0) \rangle \\
 &\sim \langle 0 | \mathcal{O}_D | D \rangle \underbrace{\langle D | J | \pi \rangle}_{\text{desired form factor}} \langle \pi | \mathcal{O}_\pi | 0 \rangle e^{-m_D(T-t)} e^{-m_\pi t}
 \end{aligned}$$

(bare) transition matrix element ~ desired form factor





Form factors from 2pt + 3pt

- Recall the ratio $R(T, t)$

$$R(T, t) = (\text{factors}) \times \frac{C_3(T, t)}{\sqrt{C_2^L(t)C_2^H(T-t)}} e^{+E_L t/2} e^{+E_H(T-t)/2}$$

$\xrightarrow{1 \ll t \ll T} f_J$

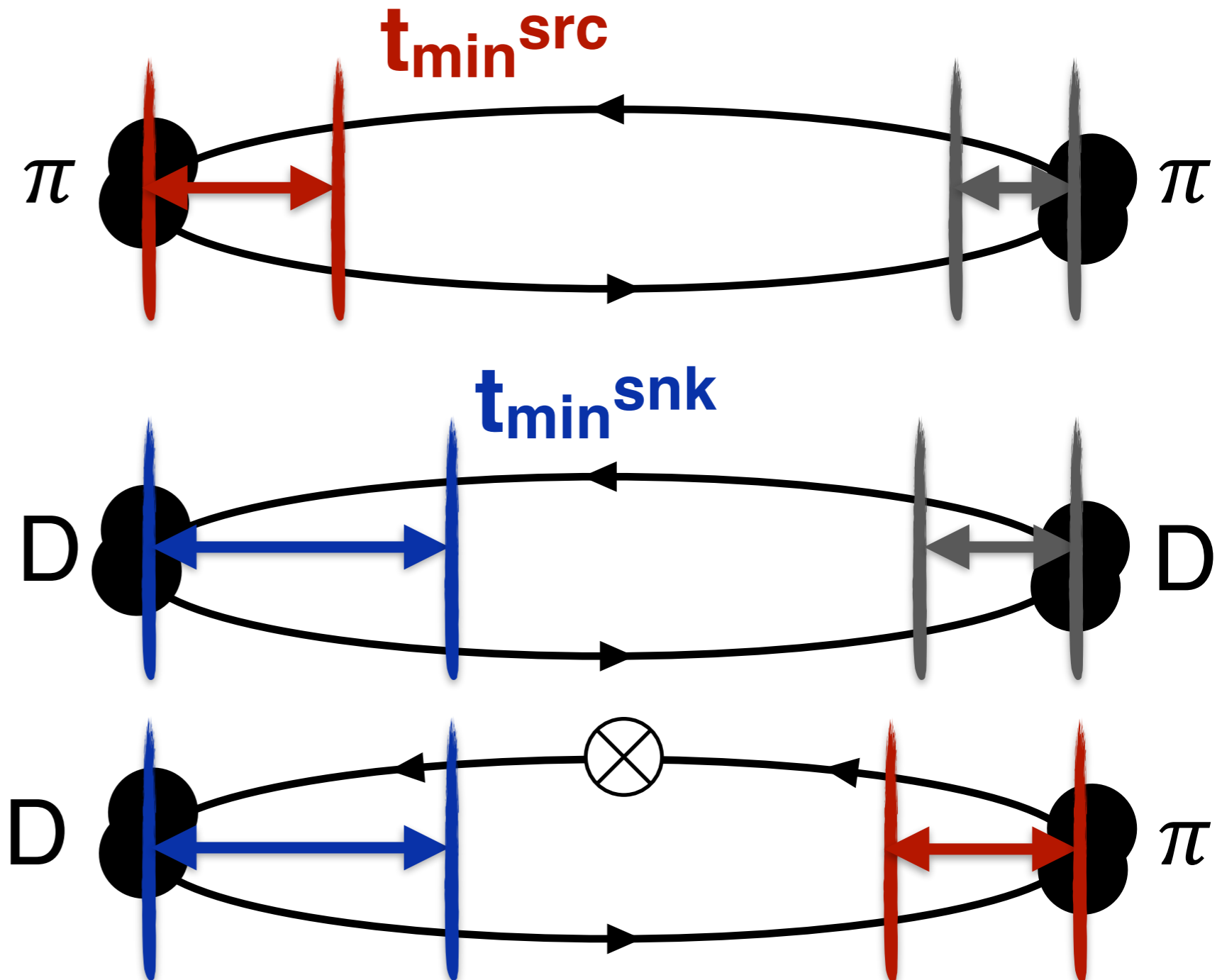
- The ratio asymptotically approaches the form factors
- Our analysis of 2pt and 3pt functions fits the full spectral decomposition
- The ratio gives a valuable visual check on fit results



2pt + 3pt

- π : $(n+n_0)$ states
- D : $(m+m_0)$ states
- Fix distances

t_{\min}^{src} and t_{\min}^{snk}
for all correlators
in physical units



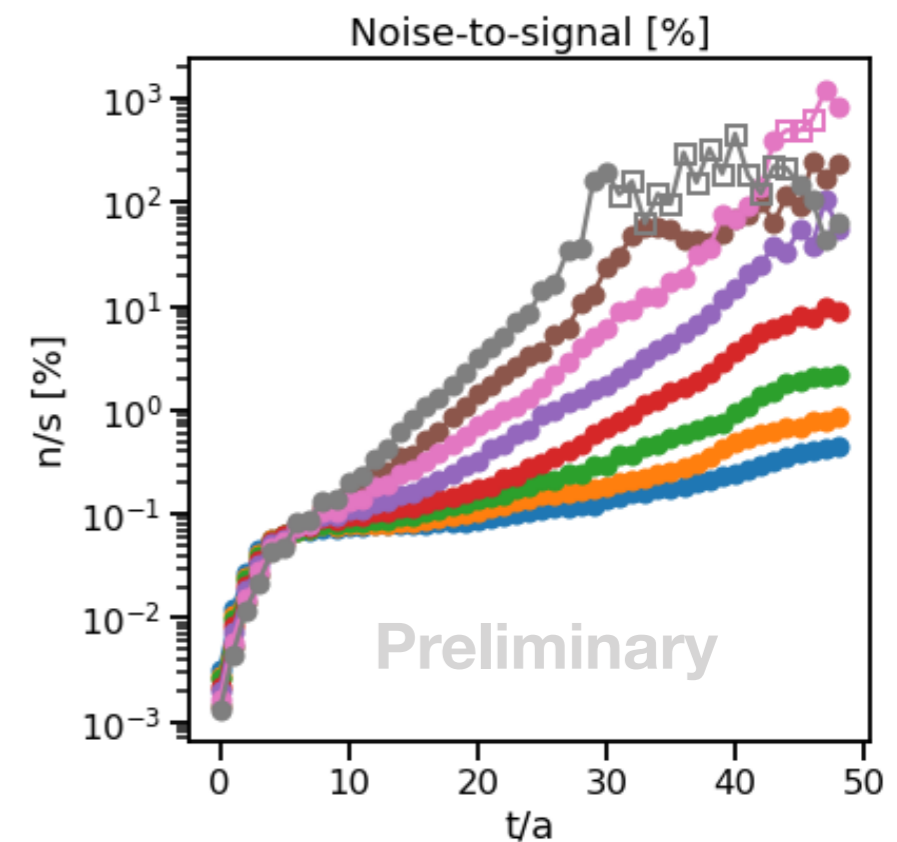
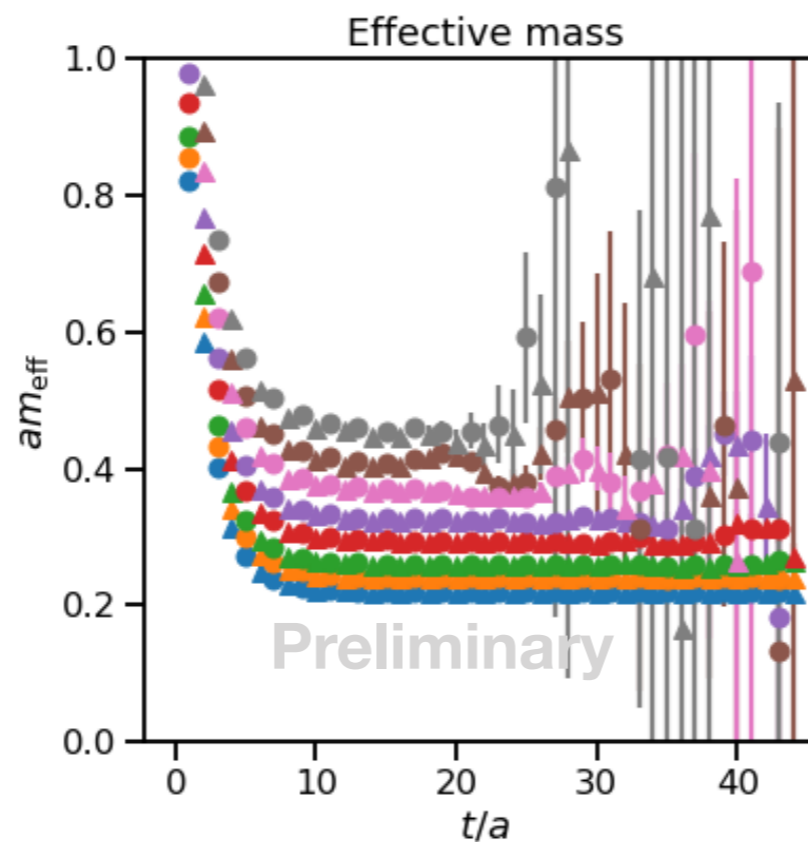
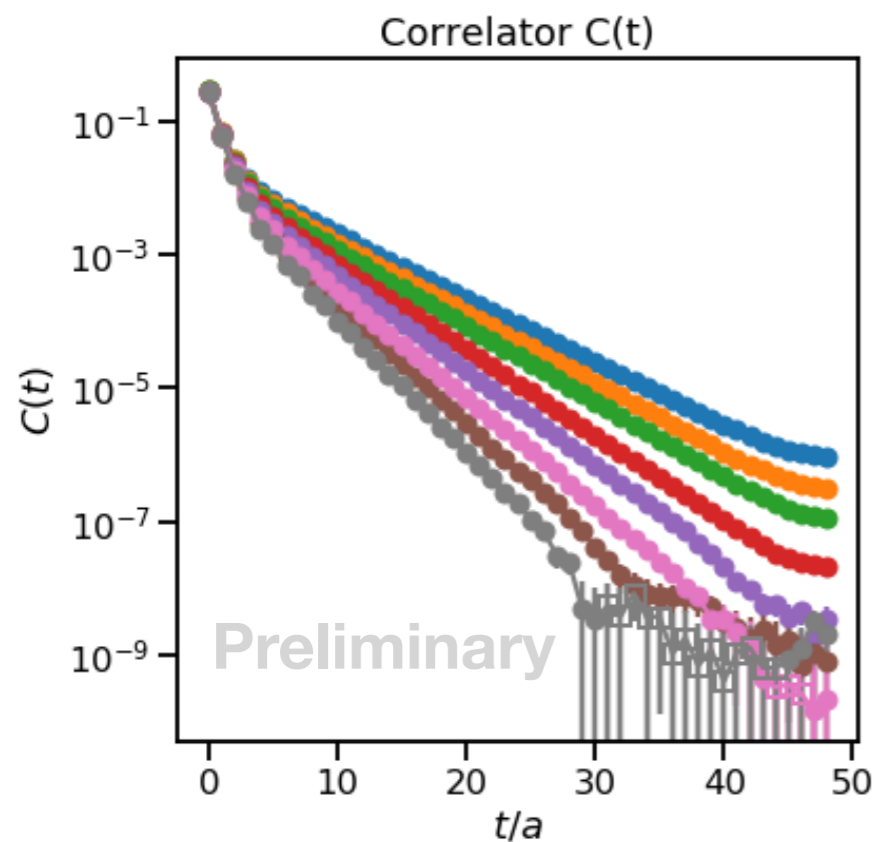


Preliminary results



2pt functions

- Consider kaon 2pt functions at $a=0.09$ fm
- Fit each 2pt function
- Find appropriate values for t_{\min} and number of states





2pt functions

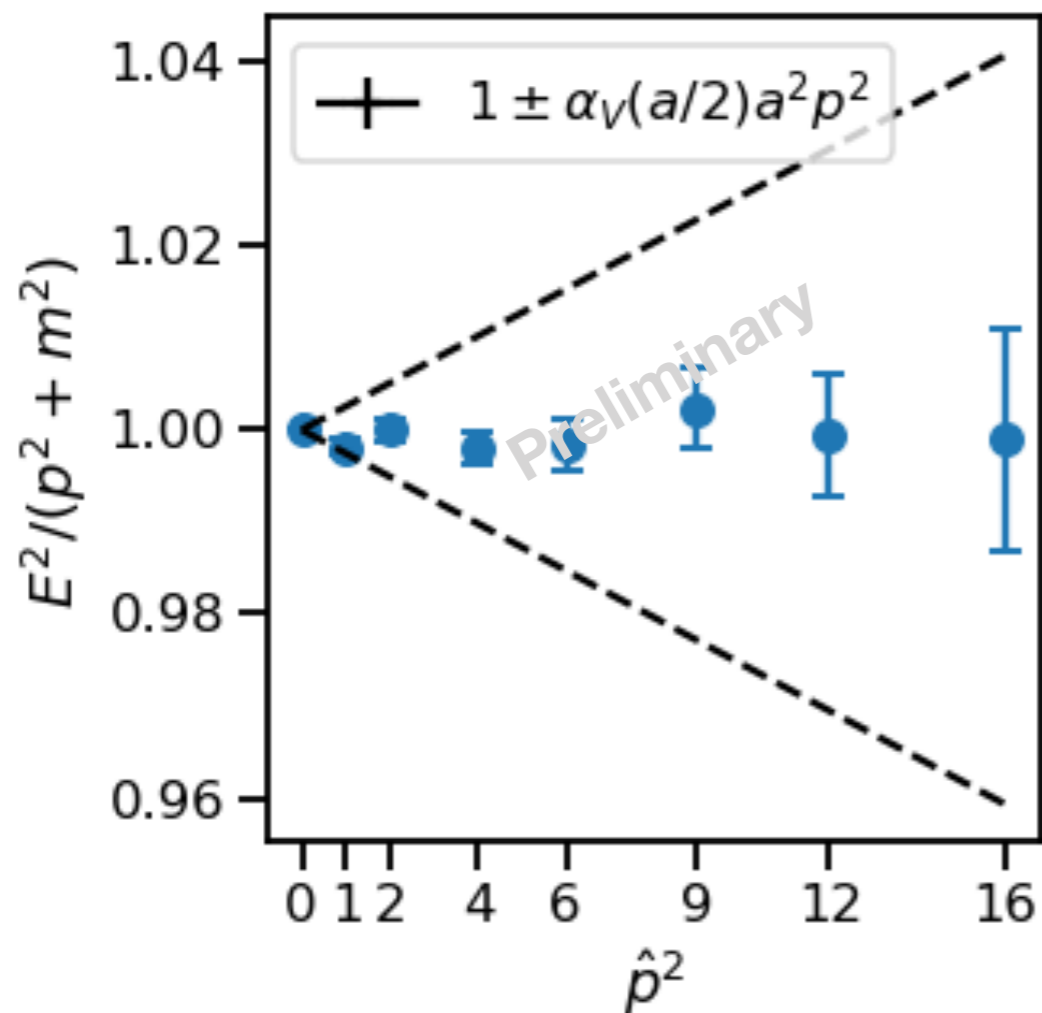
- Consider kaon 2pt functions at $a=0.09$ fm
- Fit each 2pt function
- Find appropriate values for t_{\min} and number of states
- Check relativistic dispersion relation



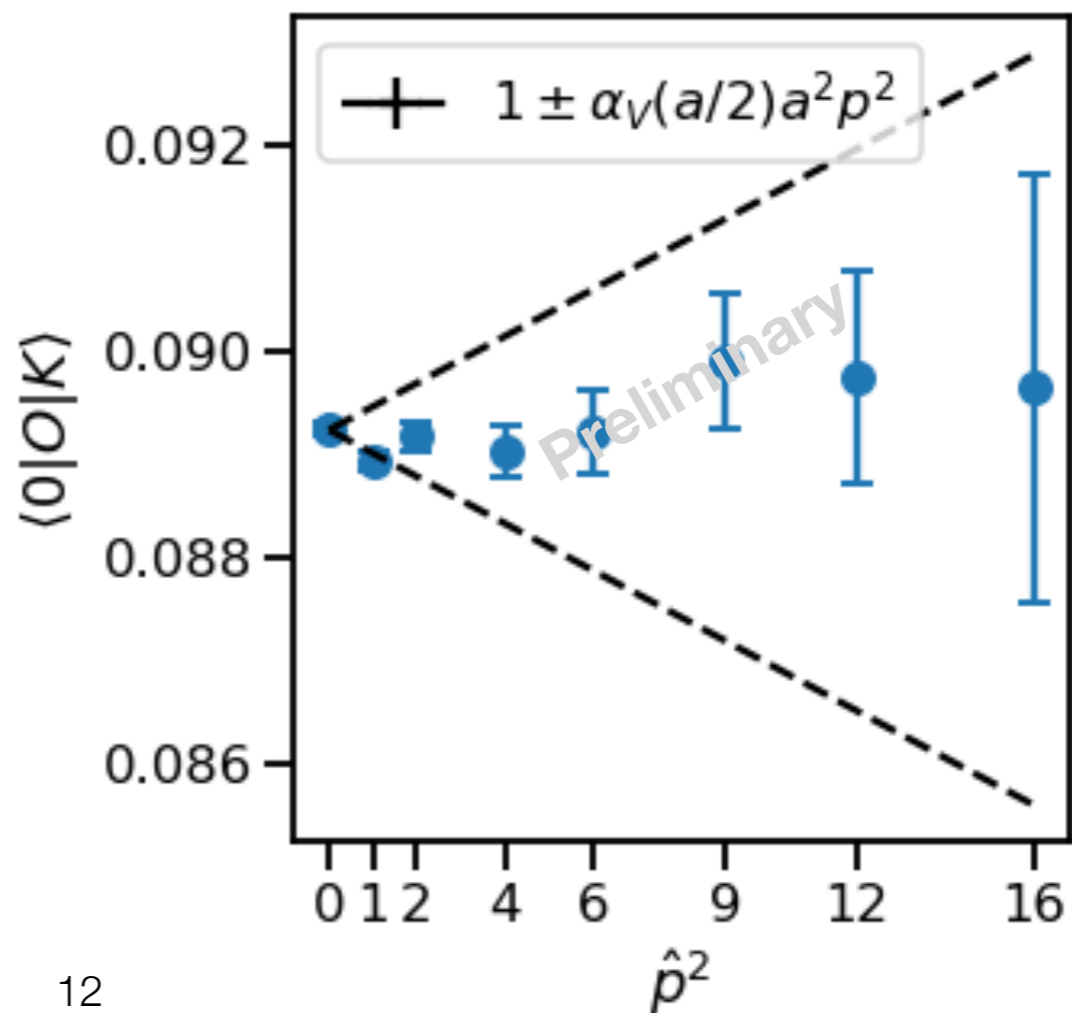
2pt functions

- Consider kaon 2pt functions at $a=0.09$ fm
- Check relativistic dispersion relation

✓ $E^2 = p^2 + m^2$



✓ $\langle 0|O|K\rangle = \text{const}$





2pt functions

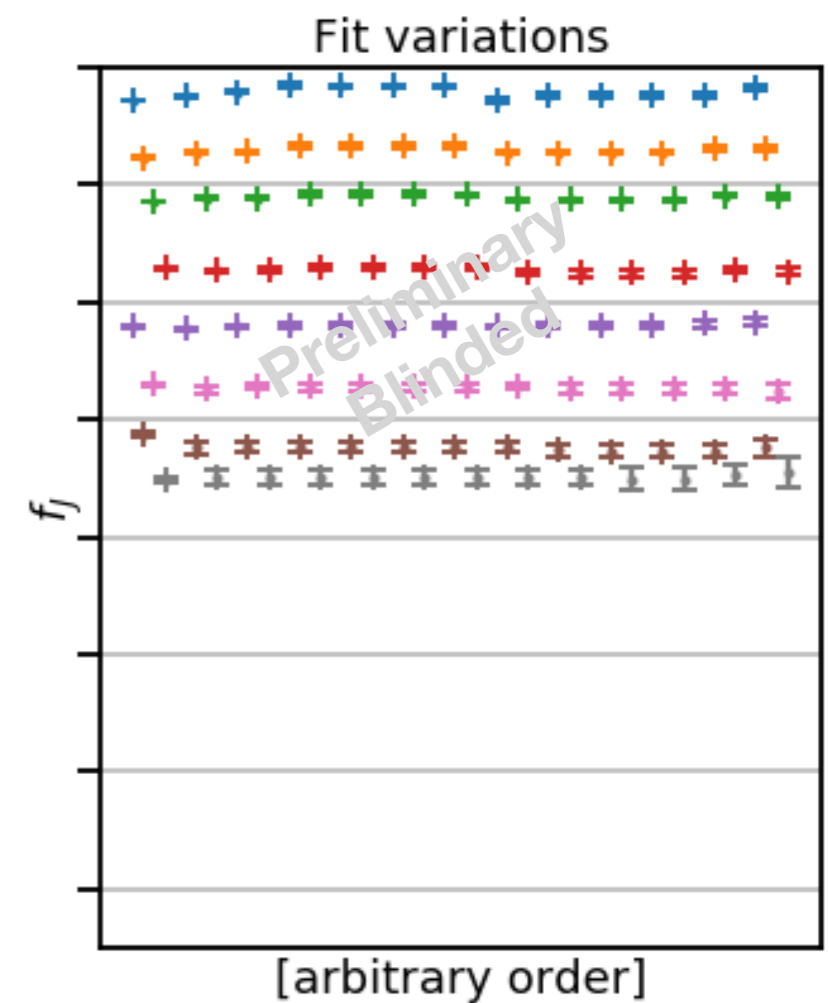
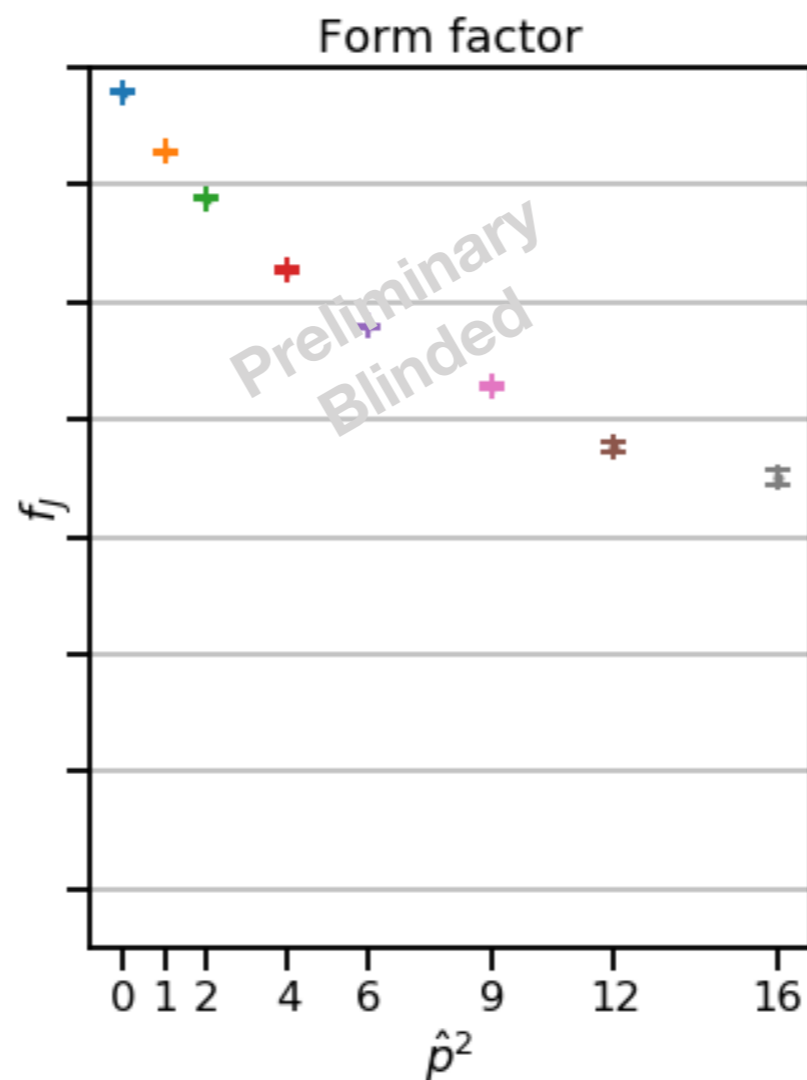
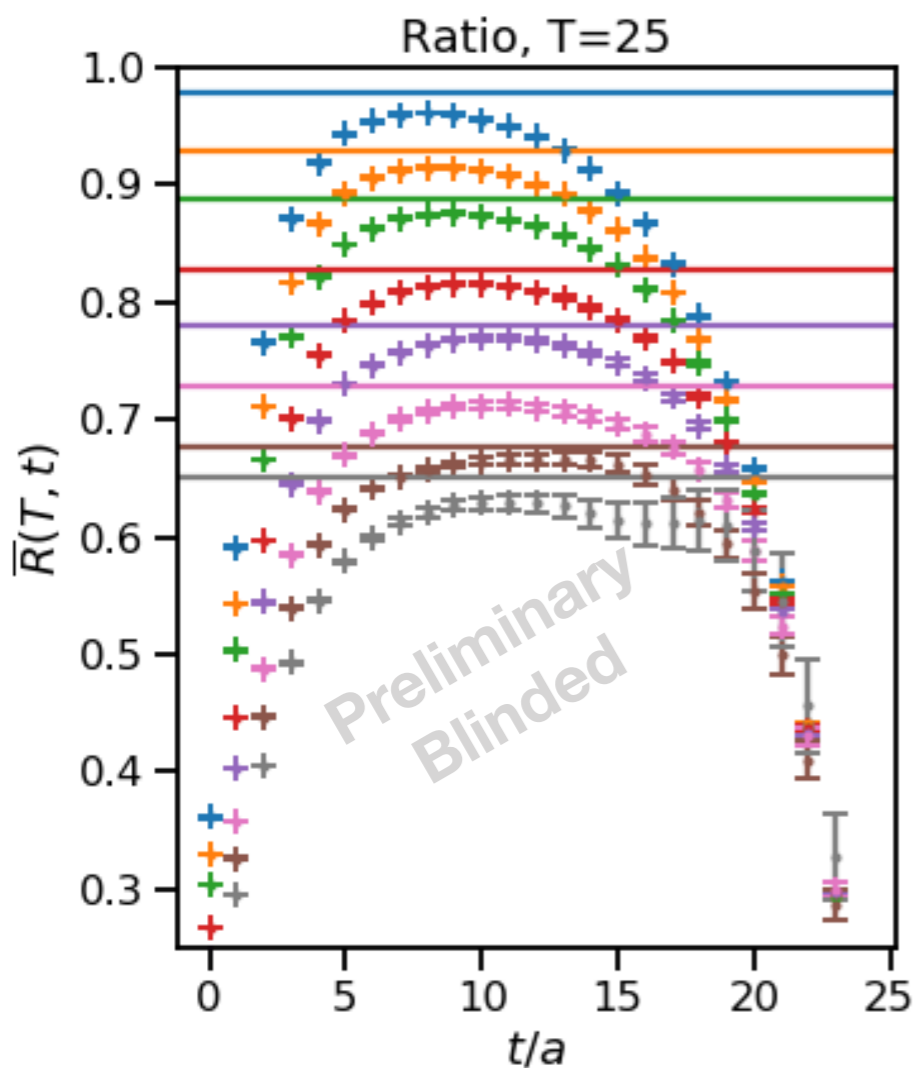
- Consider kaon 2pt functions at $a=0.09$ fm
- Check relativistic dispersion relation
- Use knowledge from 2pt functions as inputs for 3pt fits
 - ▶ Fix t_{\min} in physical units
 - ▶ Fix number of states $(n+n_0)$ and $(m+m_0)$
 - ▶ Impose knowledge $E^2/(\mathbf{p}^2+m^2) = 1 + \mathcal{O}(\alpha_s(ap)^2)$
 - ▶ Verify results visually for sanity



Form factor fits

- $D_s \rightarrow K$, physical masses, $a=0.09$ fm
- Joint fit with two 2pt functions and five 3pt functions

Scalar form factor $f_0, D_s \rightarrow K$





Form factor fits

- $D_s \rightarrow K$, physical masses, $a=0.09$ fm
- Joint fit with two 2pt functions and five 3pt functions
- Results are qualitatively very similar for
 - ▶ Different currents (V^i , V^4 , and T^{ij})
 - ▶ Different masses (i.e., $D \rightarrow K$ and $D \rightarrow \pi$)
 - ▶ Different lattice spacings
- Note: 3pt results are fully blinded



Renormalization

- PCVC: $\nabla_{\mu} V^{\mu} = (m_1 - m_2)S$
- This Ward Identity guarantees absolute normalization for the scalar form factor: $Z_S \equiv 1$.
- Z-factors are still needed for the local, one-link vector currents.
- PCVC defines the Z-factors non-perturbatively

$$\begin{aligned}
 (M_D - M_{\pi}) \langle D | V_{\text{lat}}^4 | \pi \rangle Z_{V^4} + \mathbf{q} \cdot \langle D | \mathbf{V}_{\text{lat}} | \pi \rangle Z_{V^i} & \\
 = (m_1 - m_2) \langle D | S_{\text{lat}} | \pi \rangle &
 \end{aligned}$$



Renormalization

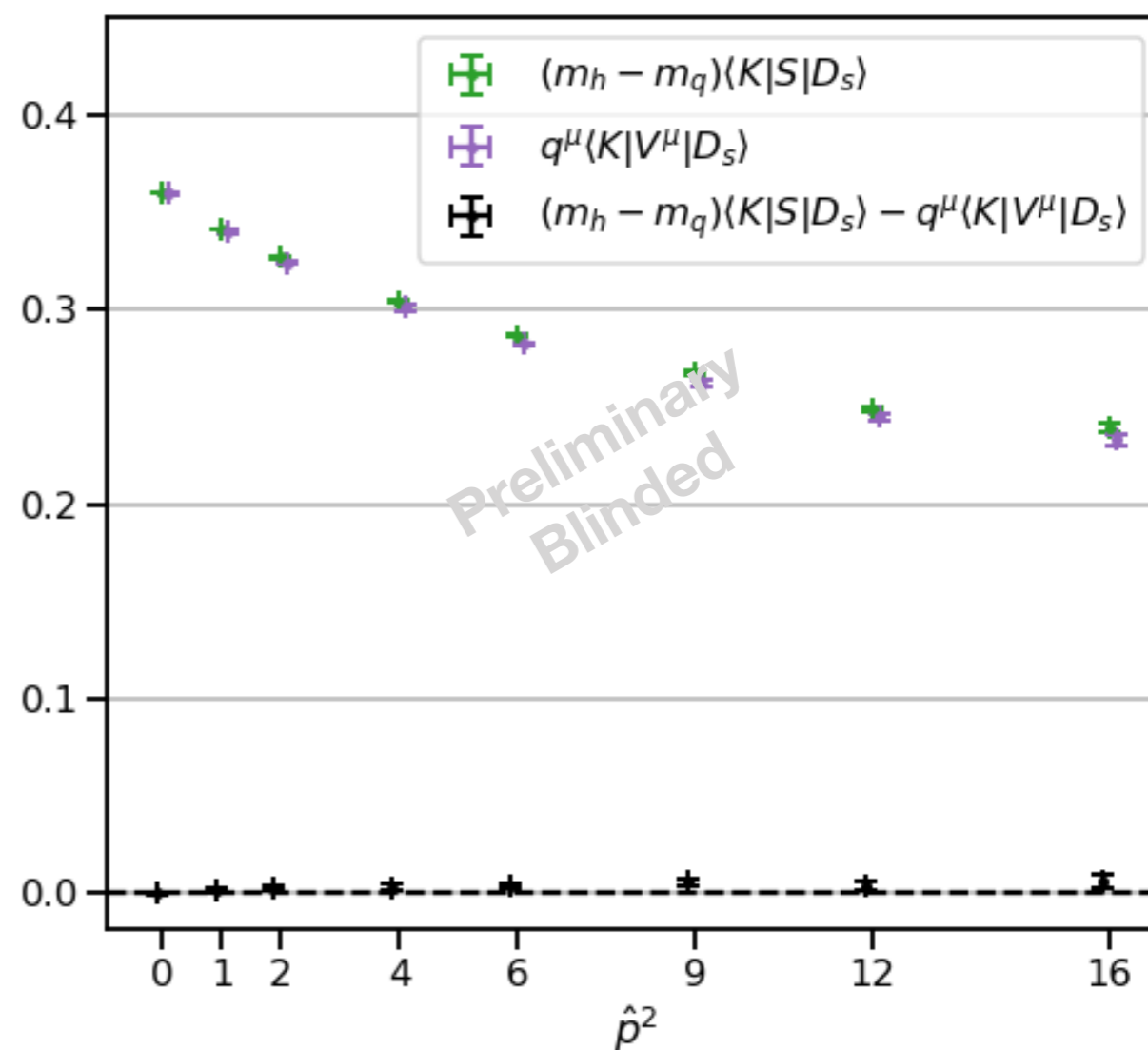
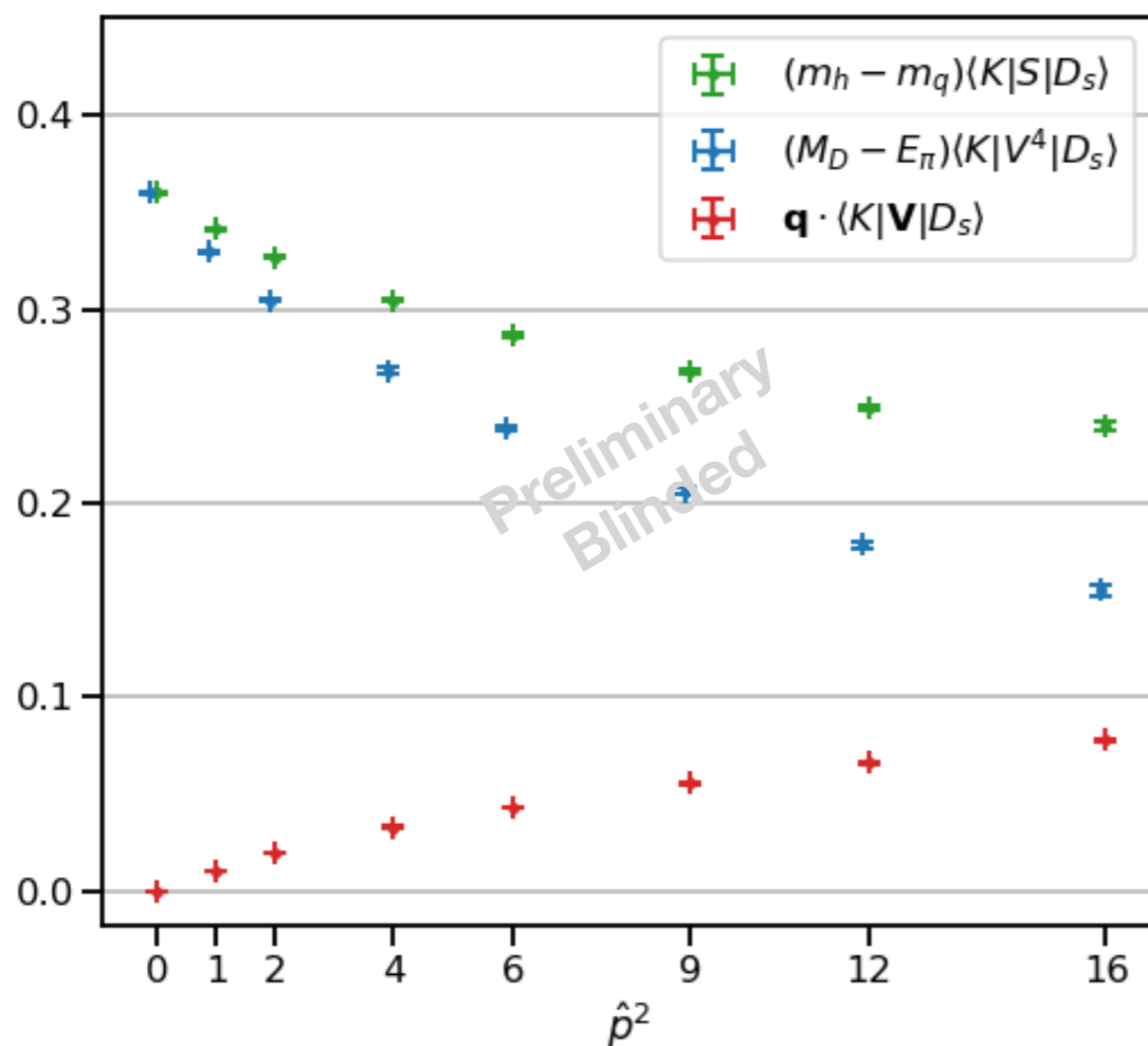
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- This Ward Identity guarantees absolute normalization for the scalar form factor: $Z_S \equiv 1$.
- Z-factors are still needed for the local, one-link vector currents.
- PCVC defines the Z-factors non-perturbatively
- Expect: $Z_{V_i}, Z_{V_4} \rightarrow 1$ as $a \rightarrow 0$



Renormalization

- $$(m_h - m_q) \langle K|S|D_s \rangle = (M_{D_s} - E_K) \langle K|V^4|D_s \rangle + \mathbf{q} \cdot \langle K|V|D_s \rangle$$

$$= q^\mu \langle K|V^\mu|D_s \rangle$$





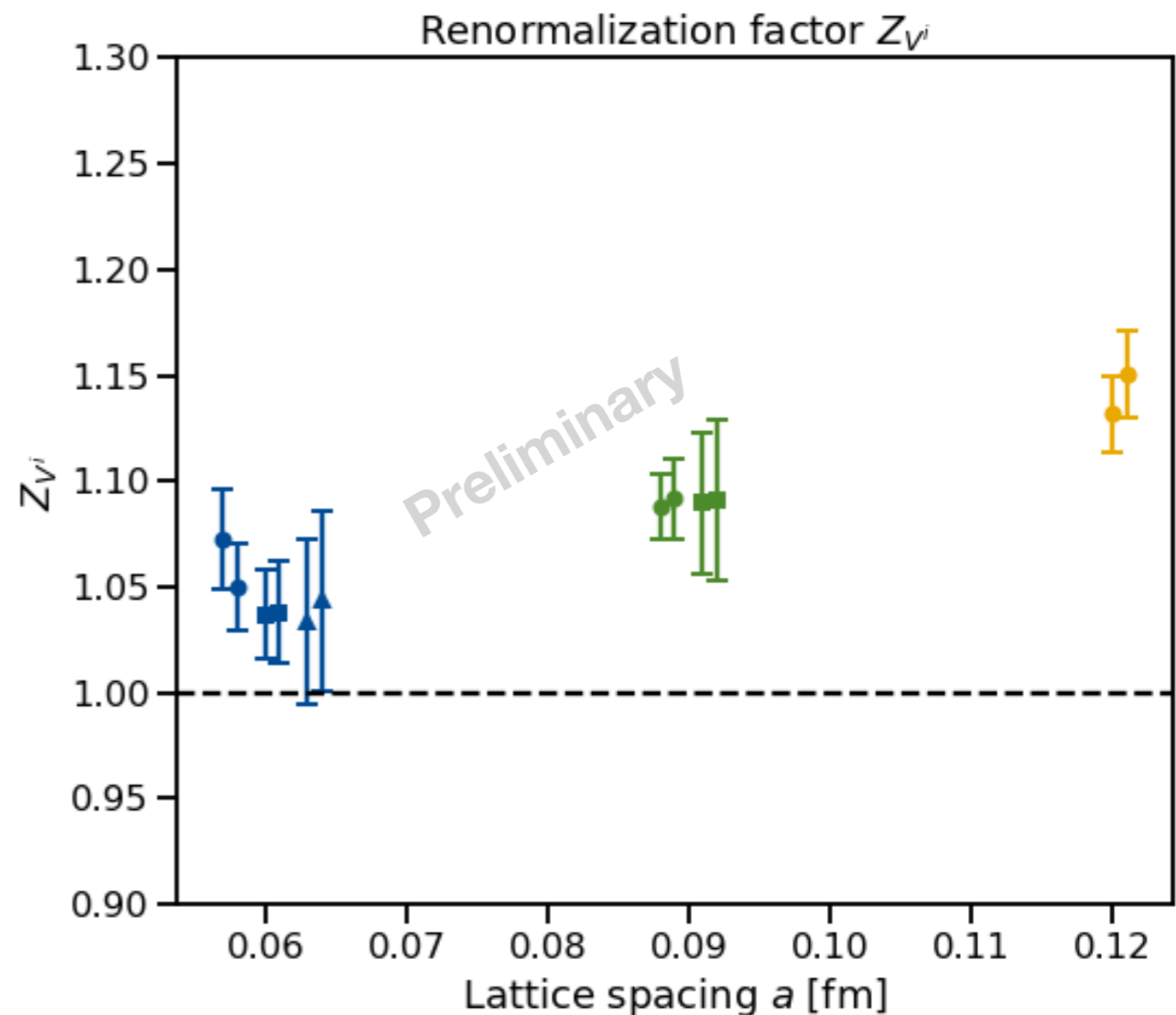
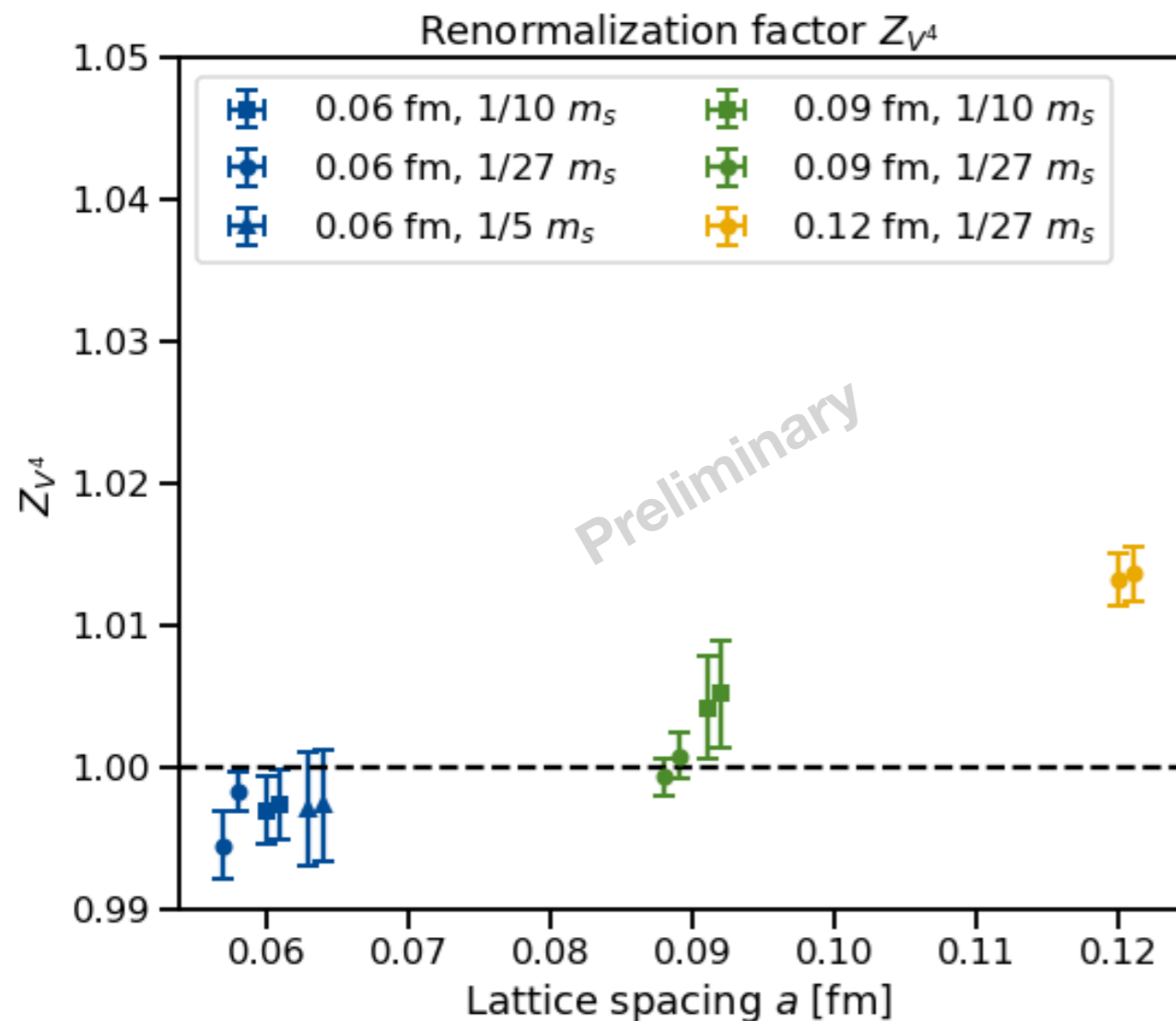
Renormalization

- $(m_h - m_q) \langle KISID_s \rangle = (M_{D_s} - E_K) \langle KIV^4ID_s \rangle + \mathbf{q} \cdot \langle KIVID_s \rangle$
 $= q^\mu \langle KIV^\mu ID_s \rangle$
- At least visually, the bare matrix elements already satisfy the Ward Identity well (even for rather large momenta)
- Many options exist to extract values for Z_{V4} and Z_{Vi}
- One possible solution: extract them from a fit



Renormalization

- Fit PCVC relation with Z_{V4} , Z_{Vi} as parameters:
- $(m_h - m_q) \langle KISID_s \rangle = Z_{V4} (M_{Ds} - E_K) \langle KIV^4ID_s \rangle + Z_{Vi} \mathbf{q} \cdot \langle KIVID_s \rangle$





Renormalized form factors

- Armed with Z-factors, we can construct form factors relevant for phenomenology. For example:

$$f_+(q^2) = [f_{\parallel}(q^2) + (M_H - E_L)f_{\perp}(q^2)]$$

$$f_0(q^2) = \frac{\sqrt{2M_H}}{M_H^2 - M_L^2} [(M_H - E_L)f_{\parallel}(q^2) - (E_L^2 - M_H^2)f_{\perp}(q^2)]$$

- Equivalent forms offer consistency checks

$$f_0(q^2) = \left(\frac{m_h - m_\ell}{M_H^2 - M_L^2} \right) \langle L|S|H \rangle \quad f_+(q^2) \propto "f_0 + f_{\perp}"$$



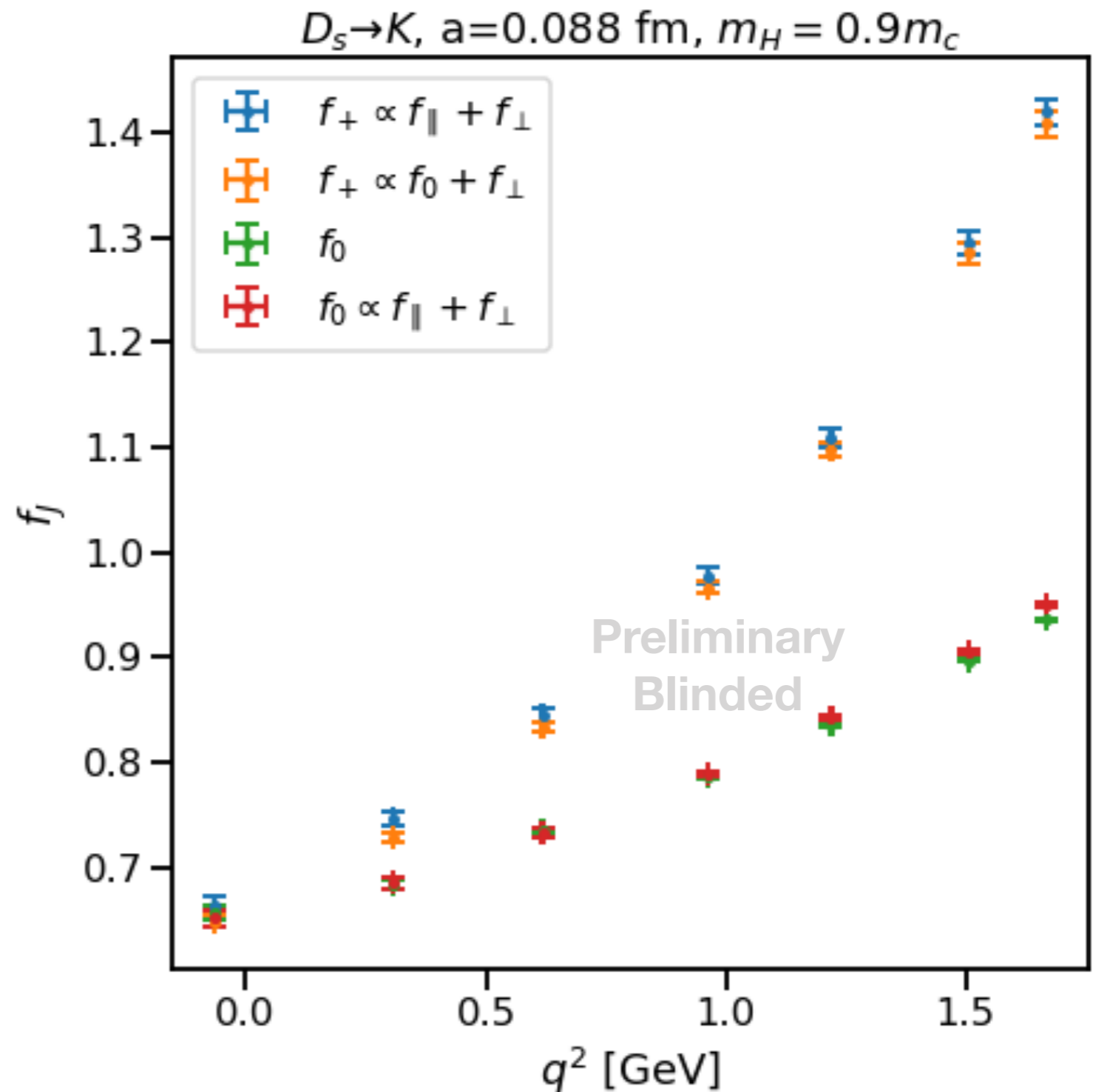
Renormalized form factors

Standard f_+ via “ $f_{\parallel} + f_{\perp}$ ”

Alternate f_+ via “ $f_0 + f_{\perp}$ ”

f_0 via scalar
matrix element

Standard f_0 via “ $f_{\parallel} + f_{\perp}$ ”





Renormalized form factors

Standard f_+ via “ $f_{\parallel} + f_{\perp}$ ”

Alternate f_+ via “ $f_0 + f_{\perp}$ ”

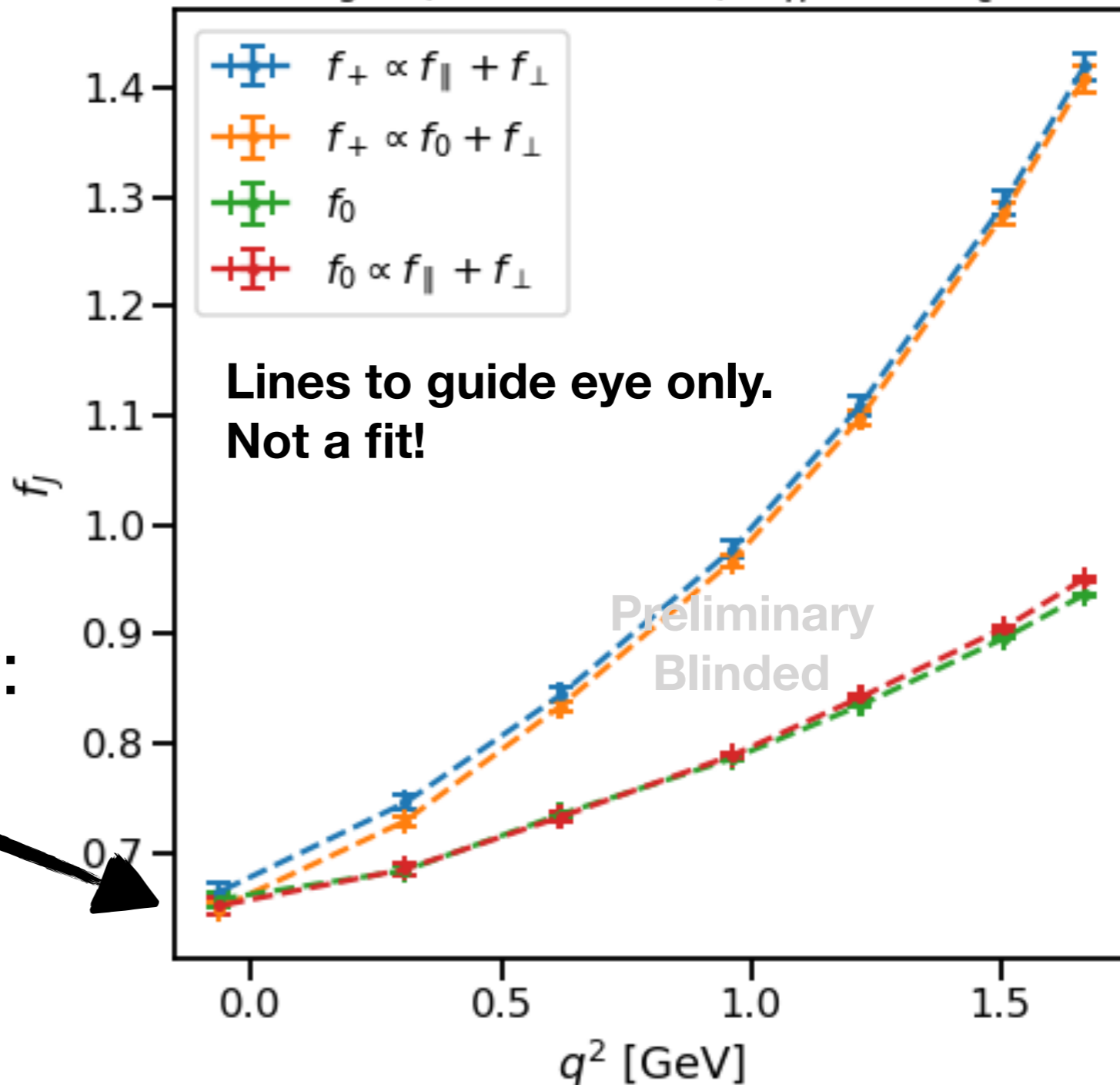
f_0 via scalar
matrix element

Standard f_0 via “ $f_{\parallel} + f_{\perp}$ ”

Check kinematic identity:
 $f_+(q^2) = f_0(q^2)$ at $q^2=0$ ✓



$D_s \rightarrow K$, $a=0.088$ fm, $m_H = 0.9m_c$





Summary and next steps

- We have presented preliminary results for $D_{(s)} \rightarrow \pi/K$
 - Matrix elements for S, V4, and V currents and range of momenta spanning full kinematic range
 - Non-perturbative renormalization via PCVC
 - Fully blinded results pass a variety of sanity checks
- In progress: chiral interpolation and continuum extrapolation



Backup slides



Three-point correlators

- We compute 40 current-momentum combinations for each mass triplet (m_{mother} , m_{daughter} , $m_{\text{spectator}}$)

momentum	\hat{p}^2	S	V^4	V^i			T^{ij}		
		f_0	f_{\parallel}	$f_{\perp}(1)$	$f_{\perp}(2)$	$f_{\perp}(3)$	$f_T(14)$	$f_T(24)$	$f_T(14)$
p000	0	✓	✓	✗	✗	✗	✗	✗	✗
p100	1	✓	✓	✓	✗	✗	✓	✗	✗
p110	2	✓	✓	✓	✓	✗	✓	✓	✗
p200	4	✓	✓	✓	✗	✗	✓	✗	✗
p211	6	✓	✓	✓	✓	✓	✓	✓	✓
p300	9	✓	✓	✓	✗	✗	✓	✗	✗
p222	12	✓	✓	✓	✓	✓	✓	✓	✓
p400	16	✓	✓	✓	✗	✗	✓	✗	✗

Tab. 4: Details about the correlation functions appearing the allHISQ campaign.