Precision Calculation of the *x*dependence of PDFs from Lattice QCD

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YONG ZHAO JUL. 28, 2021



In collaboration with Xiang Gao, Andrew Hanlon, Nikhil Karthik, Swagato Mukherjee, Peter Petreczky, Philipp Scior, Sergey Syritsyn, in preparation.

Large-momentum effective theory (LaMET)

• The quasi-PDF: X. Ji, PRL 110 (2013); SCPMA57 (2014).

$$\tilde{f}(y, P^{z}) = \int_{-\infty}^{\infty} \frac{dz}{4\pi} e^{iz(yP^{z})} \langle P | \bar{\psi}(z) \Gamma W[z, 0] \psi(0) | P \rangle$$

 Direct calculation of x-dependence through largemomentum expansion:

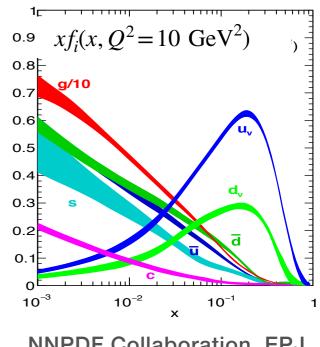
$$f(x,\mu) = \int_{-\infty}^{\infty} \frac{dy}{|y|} C\left(\frac{x}{y}, \frac{\mu}{yP^{z}}\right) \tilde{f}(y, P^{z}, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^{2}}{(xP^{z})^{2}}, \frac{\Lambda_{\text{QCD}}^{2}}{((1-x)P^{z})^{2}}\right)$$

- X. Xiong, X. Ji, J.-H. Zhang and YZ, PRD 90 (2014);
- Y.-Q. Ma and J. Qiu, PRD98 (2018), PRL 120 (2018);
- T. Izubuchi, X. Ji, L. Jin, I. Stewart, and YZ, PRD98 (2018).

Precision calculation = controlled systematics.

- Lattice: excited states, $a \rightarrow 0$, physical m_{π} , $L \rightarrow \infty$, etc.;
- Perturbative matching (currently at NNLO) and resummation;
- Power corrections, controllable within [x_{min}, x_{max}].

i = u, d, c, s, t, b $x \xrightarrow{P} \xrightarrow{P} P$ i = g

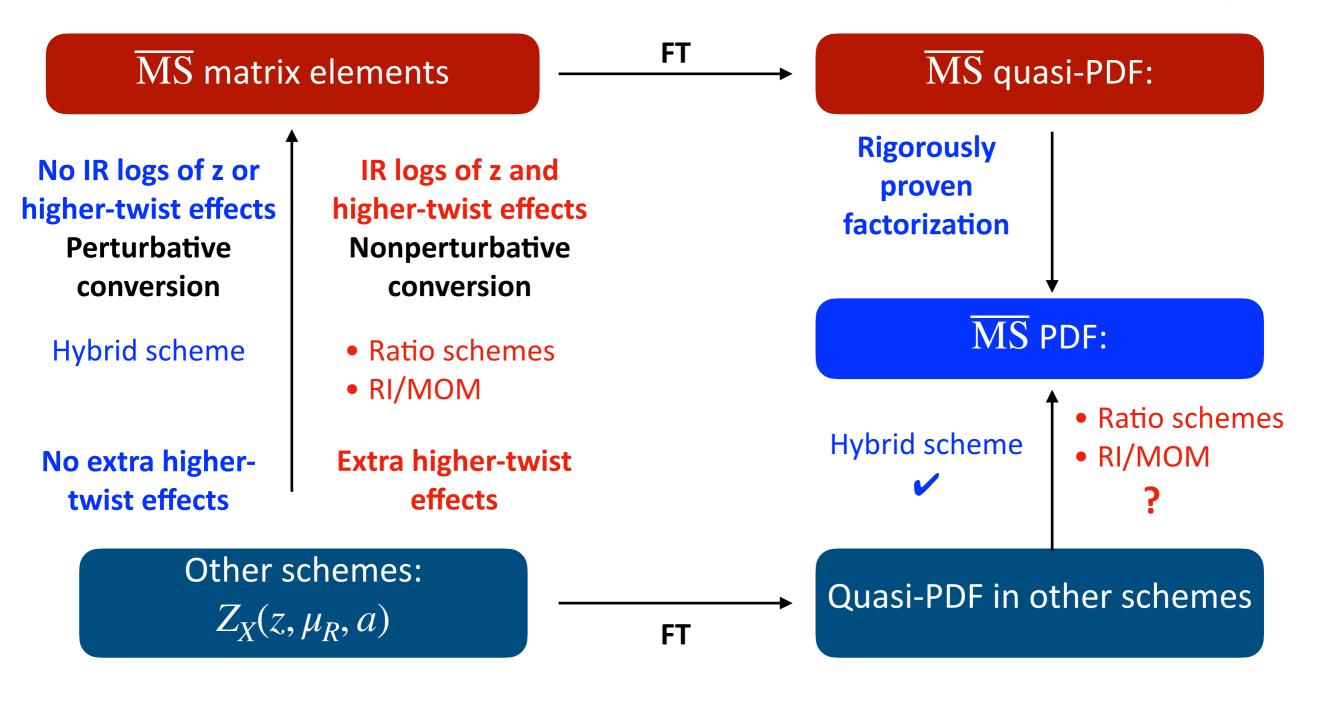


NNPDF Collaboration, EPJ C77 (2017)

Renormalization and factorization

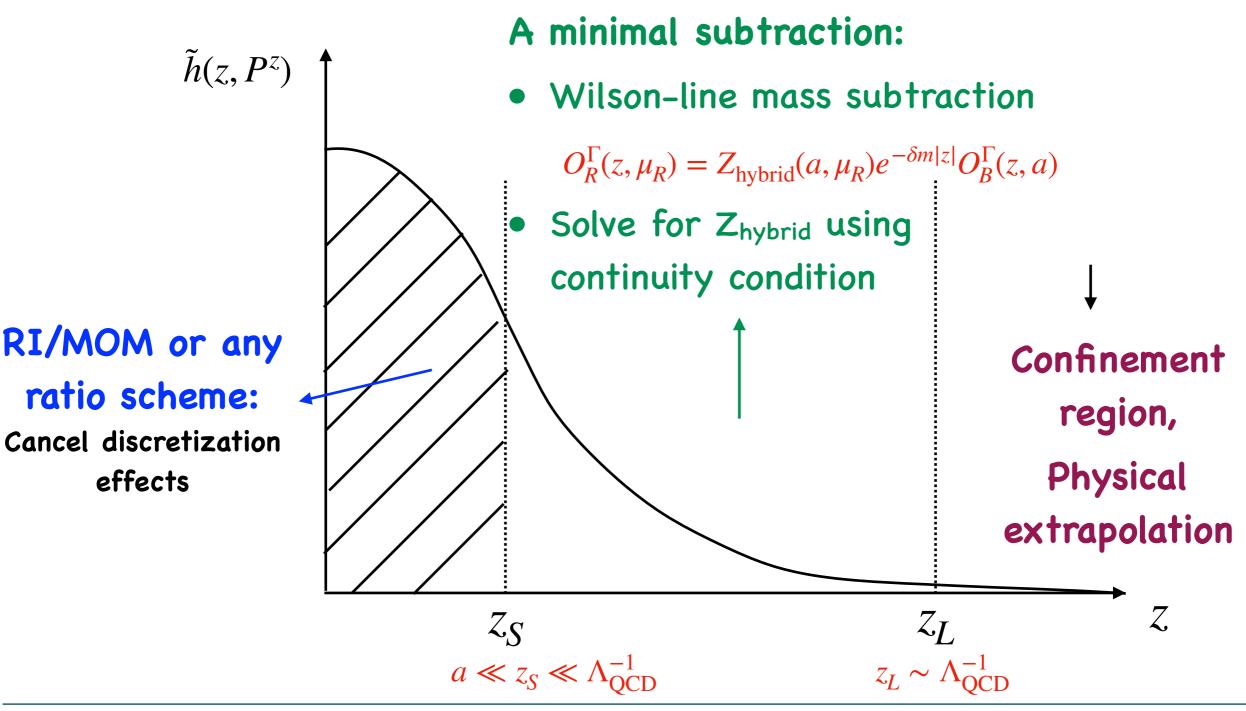
• Coordinate space:

• Momentum space/x-space:



Hybrid renormalization scheme

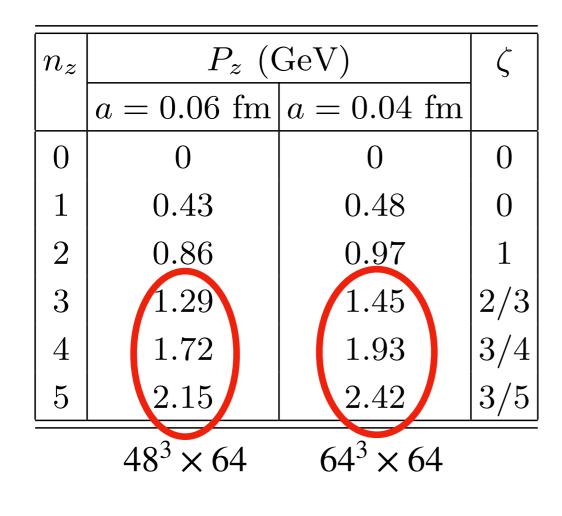
See X. Ji, YZ, et al., NPB 964 (2021) and references therein.



 $O_{R}^{\Gamma}(z,a) = e^{-\delta m(a)|z|} Z_{O}(a,\mu) O_{R}^{\Gamma}(z,\mu)$

Lattice calculation

• Wilson-clover fermion on 2+1 flavor HISQ configurations.



X. Gao, YZ, et al., PRD102 (2020).
X. Gao, YZ, et al., 2102.01101.

$$m_{\pi} = 300 \text{ MeV}$$

Wilson-line mass renormalization

Polyakov loop

 $\begin{array}{c|c} \langle \Omega | & & \uparrow \\ \leftarrow & T \rightarrow \infty & - \end{array} \end{array} | \Omega \rangle & \propto \exp[-V(R)T] \end{array}$

Renormalization condition:

$$V^{\text{lat}}(r,a) \bigg|_{r=r_0} + 2\delta m(a) = 0.95/r_0$$
$$\delta m(a) = \frac{1}{a} \sum_n c_n \alpha_s^n (1/a) + \delta m_0^{\text{lat}}$$
$$\delta m_0^{\text{lat}} \sim \Lambda_{\text{QCD}}$$

C. Bauer, G. Bali and A. Pineda, PRL108 (2012).

 $a\delta m(a = 0.04 \text{ fm}) = 0.1508(12)$ $a\delta m(a = 0.06 \text{ fm}) = 0.1586(8)$ $a\delta m(a = 0.076 \text{ fm}) = 0.1597(16)$ A. Bazavov et al., TUMQCD, PRD98 (2018).

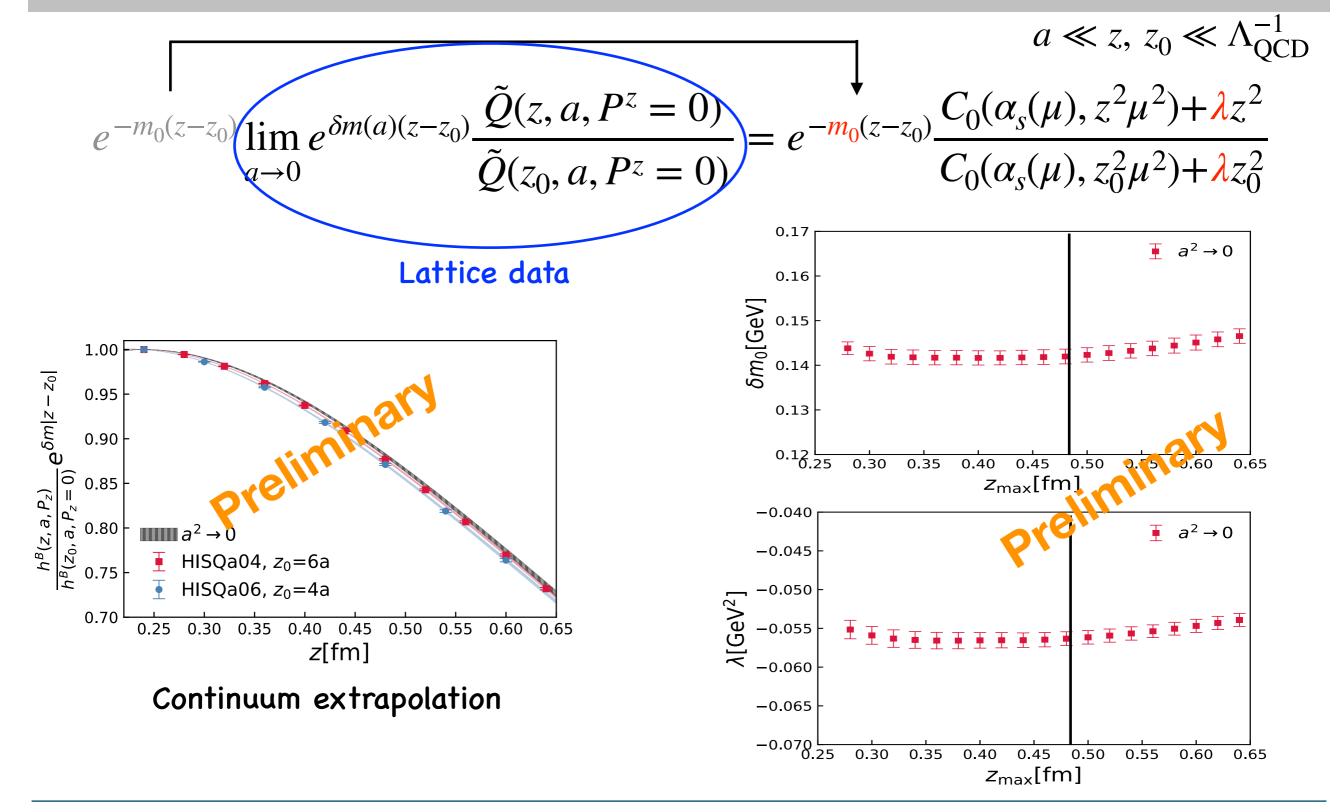
Wilson-line mass renormalization

• Check of continuum limit:

 $O_B^{\Gamma}(z, a) = e^{-\delta m|z|} Z_O(a) O_R^{\Gamma}(z)$

$$\lim_{a \to 0} e^{-\delta m(z-z_0)} \frac{\tilde{h}(z, a, P^z = 0)}{\tilde{h}(z_0, a, P^z = 0)} = \frac{\tilde{h}(z, P^z = 0, \mu)}{\tilde{h}(z_0, P^z = 0, \mu)} \qquad z, z_0 \gg a$$

Matching to the MSbar scheme

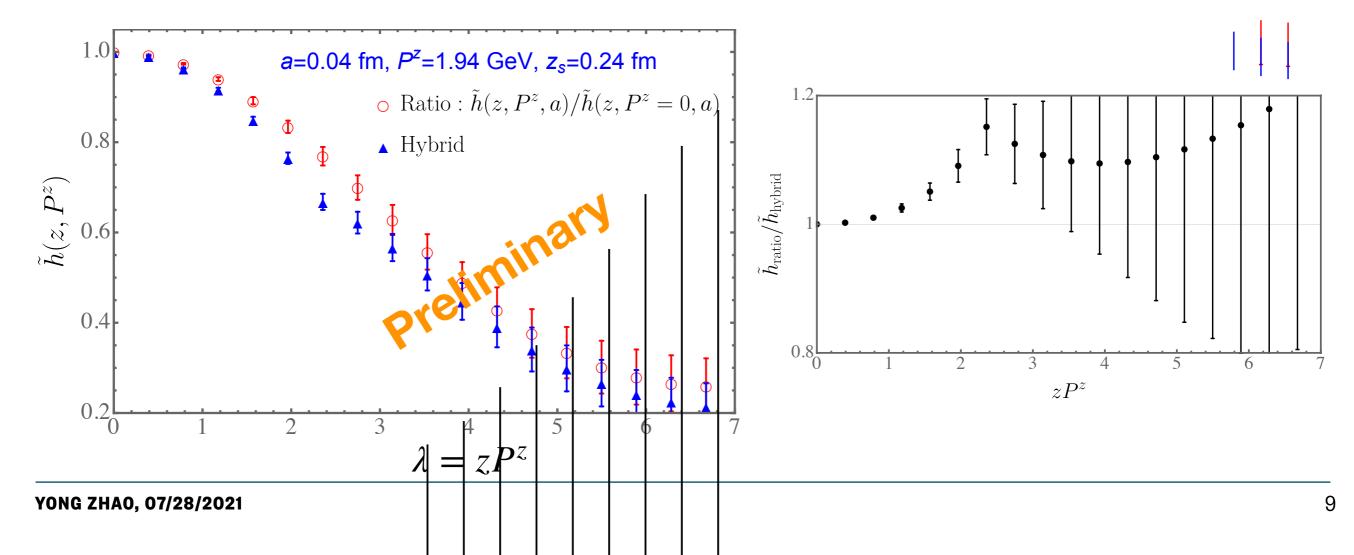


Renormalized matrix element

Power-correction corrected ratio at short distance:

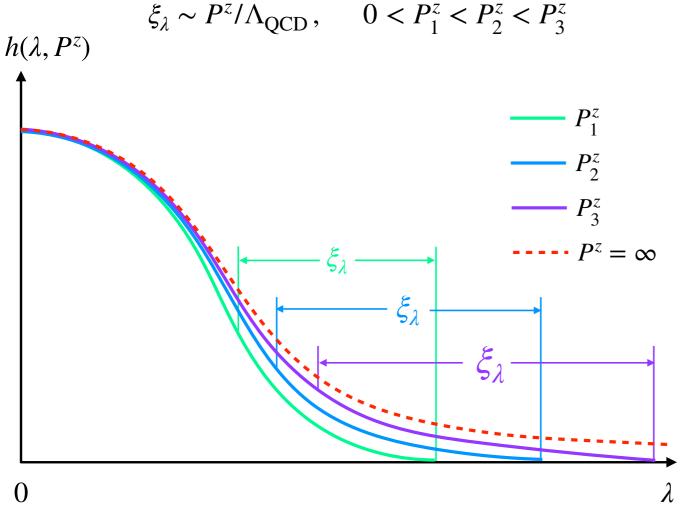
 $\lim_{a \to 0} \frac{\tilde{h}(z, a, P^z)}{\tilde{h}(z, a, P^z = 0)} \frac{C_0(\alpha_s(\mu), z^2 \mu^2) + \lambda z^2}{C_0(\alpha_s(\mu), z^2 \mu^2)} = \frac{\tilde{h}(z, P^z, \mu)}{C_0(\alpha_s(\mu), z^2 \mu^2)} \qquad 0 \le z \le z_S$

• Hybrid scheme matrix element:



Physical extrapolation beyond *z*_L

- When z is larger than the hadron size, $\tilde{h}(z, P^z) \propto e^{-z/\xi_z}$ with $\xi_z \sim 1/\Lambda_{\rm QCD}$; in $\lambda = zP^z$ space, the correlation length $\xi_\lambda = P^z \xi_z \sim P^z / \Lambda_{\rm QCD}$;
- As $P^z \to \infty$, $\xi_{\lambda} \to \infty$, only the twist-2 contribution survives, $\tilde{h}(\lambda, P^z) \sim 1/\lambda^d$, which is determined by the small-x Regge behavior of the PDF.



Therefore, if

- P^z is not very large, e.g. 2–5 GeV, use an exponential form;
- P^z is very large, use an algebraic form.

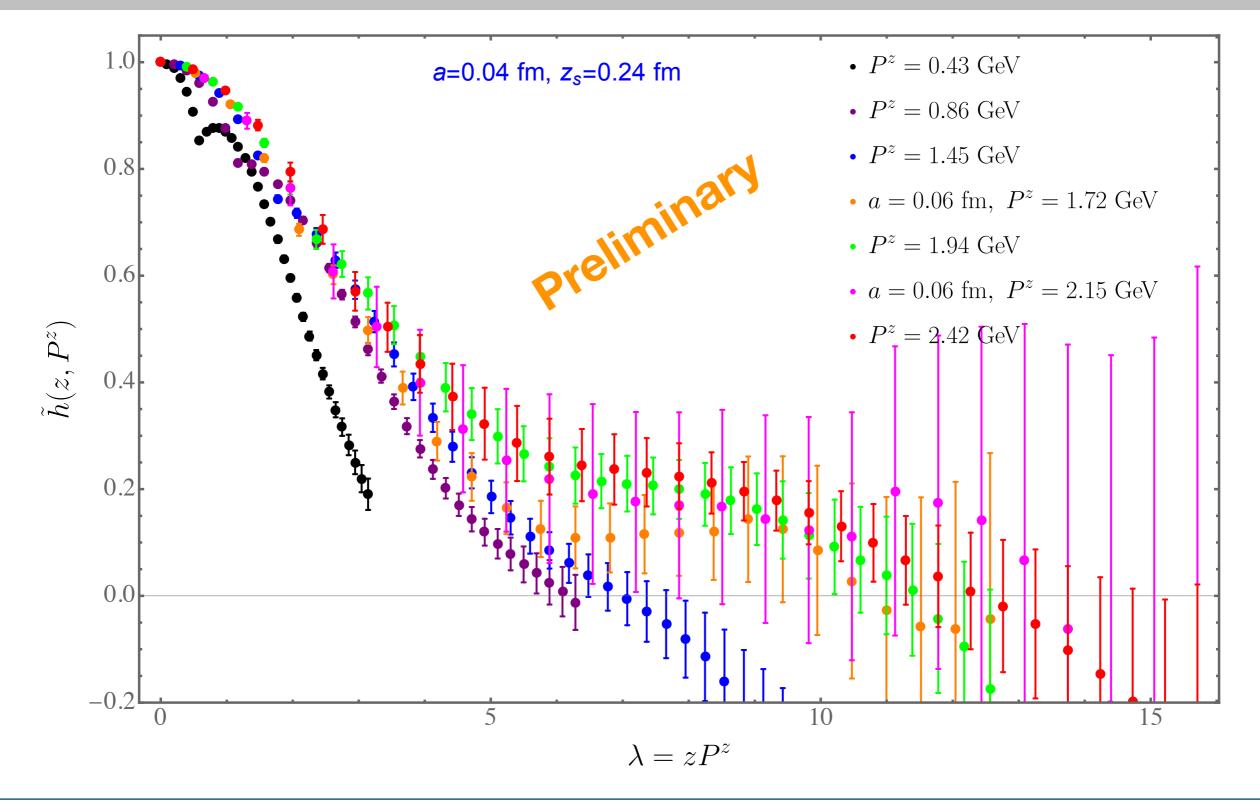
Physical extrapolation beyond *z*_L

Impact of exponential extrapolation:

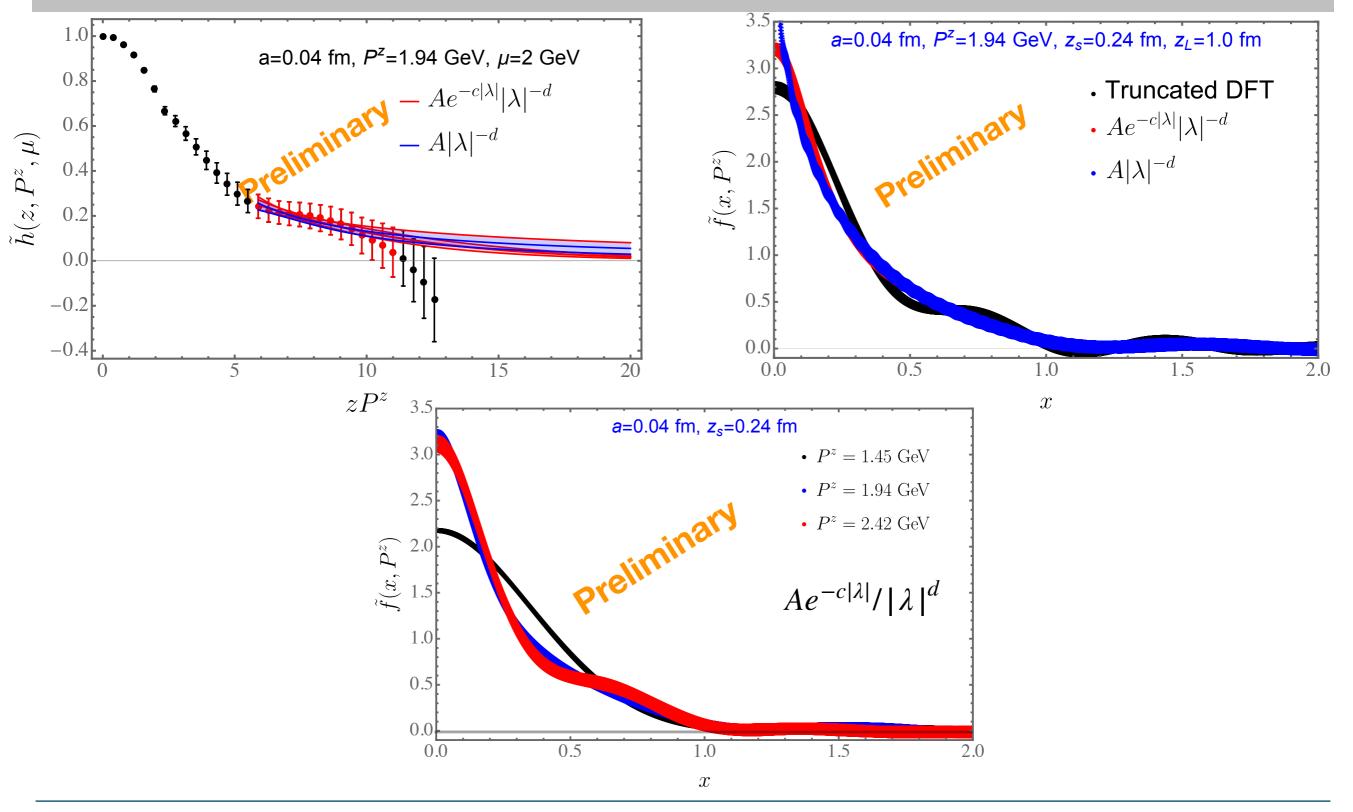
- Remove unphysical oscillation from Fourier transform;
- Affects the $x \rightarrow 0$ and $x \rightarrow 1$ regions most, but not $[x_{min}, x_{max}]$ where LaMET can have controlled prediction.

 $\tilde{f}(x, P^z)$ $\lambda_{L1} < \lambda_{L2} < \lambda_{L3} < \infty$ λ_{L1} λ_{L3} $\lambda_I = \infty$ ${\mathcal X}$

Renormalized matrix elements

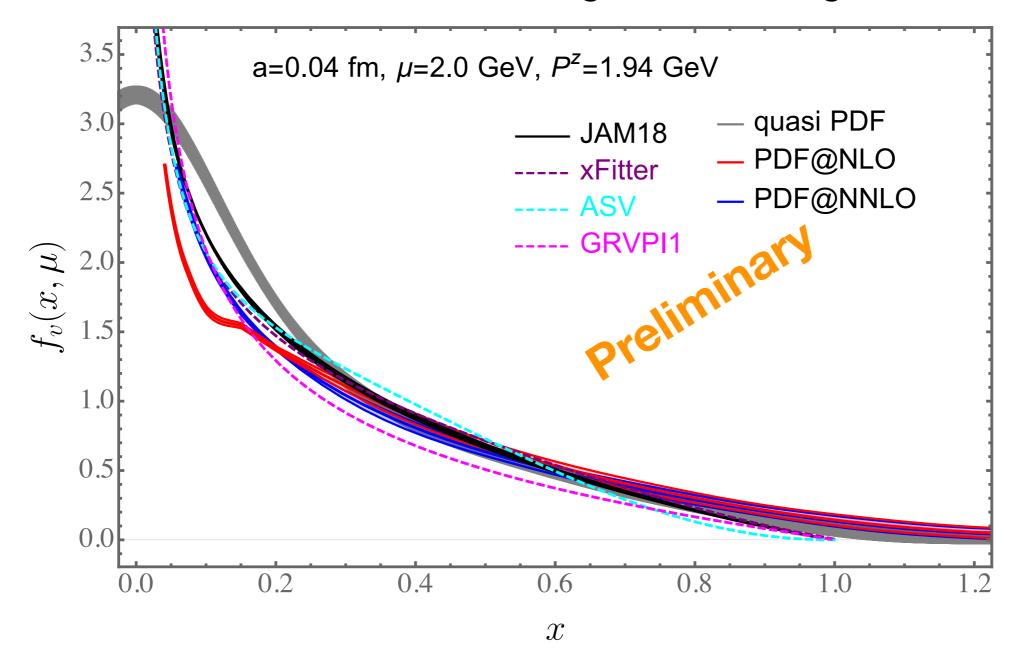


Physical extrapolation and Fourier transform



Perturbative matching at NNLO

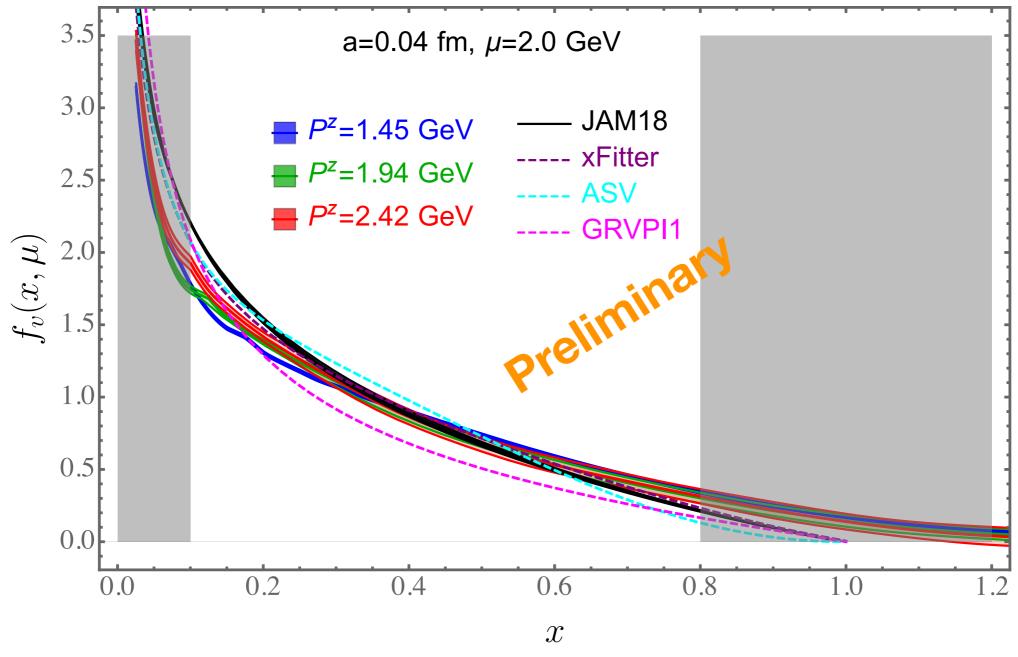
• Perturbative correction shows good convergence.



Error band assumes 100% correlation of points at all x during the matching.

YONG ZHAO, 07/28/2021

Comparison with phenomenology



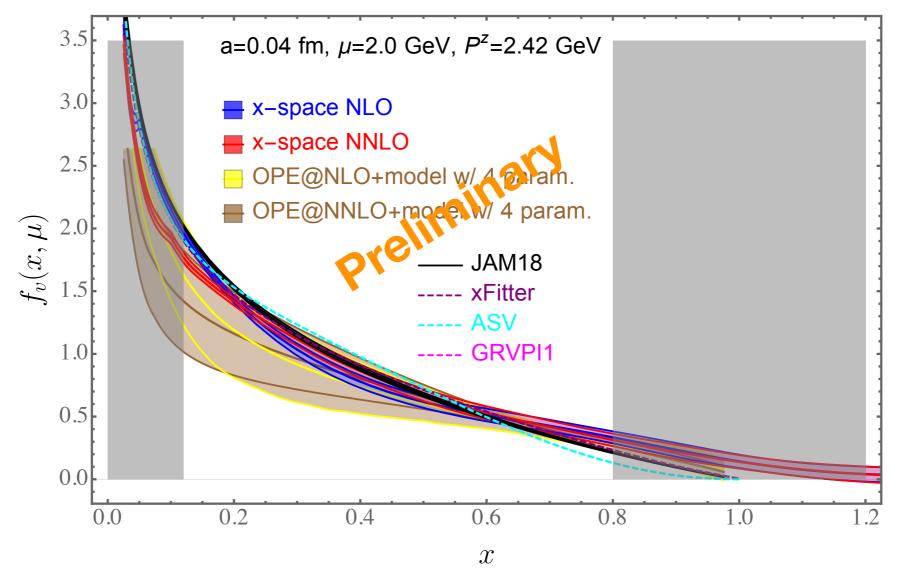
Gray band: x<0.12 or x>0.8. (A very preliminary estimate)

Error at $P^z=2.42$ GeV > 15% or variation in P^z > 1% between $P^z=1.94$ and 2.42 GeV.

Comparison with ratio scheme analysis in coordinate space

Ratio scheme with nonzero-momentum matrix elements in the dominator.

X. Gao, **YZ**, et al., PRD102 (2020).



Systematics to be included in x-space analysis: physical pion mass, lattice spacing dependence, variation of factorization scale (α_s), variation of z_L , etc.

Conclusion

- We have carried out lattice calculation of the x-dependence of pion valence PDF with the hybrid renormalization scheme;
- The Wilson-line mass correction can be well determined from lattice and matched to the MSbar scheme by using the NNLO OPE formula with quadratic power correction;
- NNLO matching shows good perturbative convergence;
- We demonstrate that we can predict the x-dependence with controlled systematic uncertainties within a subregion of x;
- Systematics to be analyzed: physical pion mass, lattice spacing dependence, variation of factorization scale (α_s), variation of z_L , etc.