

Proton momentum and angular momentum decomposition with overlap fermions

Gen Wang

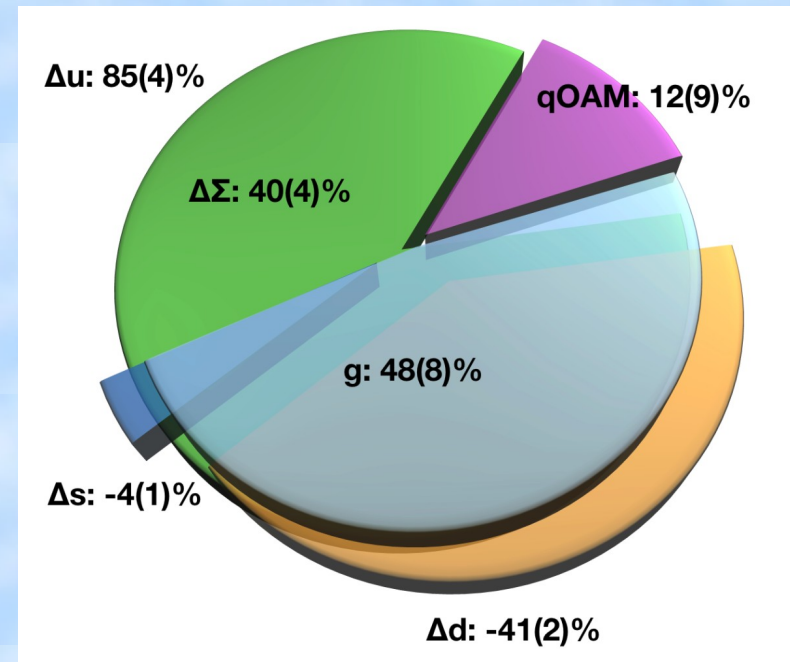
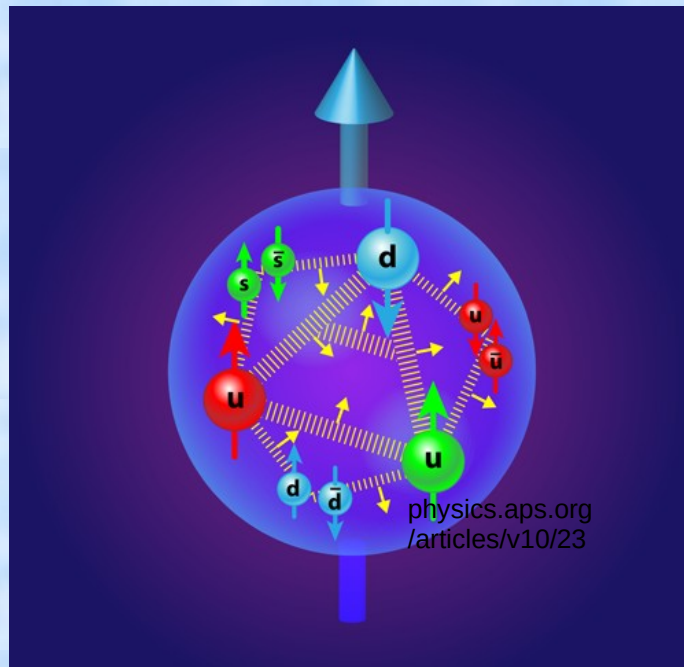
Yi-Bo Yang, Jian Liang, Terrence Draper, and Keh-Fei Liu

χ QCD Collaboration



Motivation

Only 30% of the proton's spin is carried by the spin of quark constituents [1]



[1] C. A. Aidala, S. D. Bass, D. Hasch, and G. K. Mallot, "The Spin Structure of the Nucleon," Rev. Mod. Phys. 85, 655 (2013)

[2] J. Liang, Y.-B. Yang, T. Draper, M. Gong, and K.-F. Liu, Phys. Rev. D 98, 074505 (2018), arXiv:1806.08366

[3] Y.-B. Yang, J. Liang, Y.-J. Bi, Y. Chen, T. Draper, K.-F. Liu, and Z. Liu, Phys. Rev. Lett. 121, 212001 (2018), arXiv:1808.08677

Nucleon Energy-momentum tensor

Traceless, symmetric energy-momentum tensor (EMT) between two nucleon state to T1, T2 and T3 form factors

$$\begin{aligned} \langle p', s' | \mathcal{T}^{\{\mu\nu\}q,g} | p, s \rangle &= \frac{1}{2} \bar{u}(p', s') \left[T_1(q^2) (\gamma^\mu \bar{p}^\nu + \gamma^\nu \bar{p}^\mu) \right. \\ &\quad \left. + \frac{1}{2m} T_2(q^2) (i(\bar{p}^\mu \sigma^{\nu\alpha} + \bar{p}^\nu (\sigma^{\mu\alpha}))) + \frac{1}{m} T_3(q^2) q^\mu q^\nu \right]^{q,g} u(p, s) \end{aligned}$$

T1 and T1+T2 form factors at zero momentum transfer give Momentum and angular momentum fractions

$$\mathcal{T}^{\{4i\}q,g} \quad \Longrightarrow \quad \langle x \rangle^{q,g} = T_1(0)^{q,g} \quad \langle J \rangle^{q,g} = \frac{1}{2} [T_1(0) + T_2(0)]^{q,g}$$

Mixing and Renormalization

T1, T2 and T3 form factors are renormalized by

$$\begin{aligned}T^{u,d}(\text{CI})^R &= Z_{QQ}^{\overline{\text{MS}}}(\mu)T^{u,d}(\text{CI}), \\T^{u,d,s}(\text{DI})^R &= Z_{QQ}^{\overline{\text{MS}}}(\mu)T^{u,d,s}(\text{DI}) + \delta Z_{QQ}^{\overline{\text{MS}}}(\mu) \sum_{q=u,d,s} [T^q(\text{CI}) + T^q(\text{DI})] \\&\quad + Z_{QG}^{\overline{\text{MS}}}(\mu)T^g(\text{DI}), \\T^g(\text{DI})^R &= Z_{GQ}^{\overline{\text{MS}}}(\mu) \sum_{q=u,d,s} [T^q(\text{CI}) + T^q(\text{DI})] + Z_{GG}^{\overline{\text{MS}}}T^g(\text{DI}),\end{aligned}$$

Non-perturbative renormalization constants include mixing calculated in [1]

| Lattice | Z_{QQ} | δZ_{QQ} | Z_{QG} | Z_{GQ} | Z_{GG} |
|---------|------------|-----------------|-----------|------------|------------|
| 32ID | 1.25(0)(2) | 0.018(2)(2) | 0.017(17) | 0.57(3)(6) | 1.29(5)(9) |

Normalization

With momentum and angular momentum conservation, the momentum and angular momentum sum rules are

$$\langle x \rangle^q + \langle x \rangle^g = T_1(0)^q + T_1(0)^g = 1$$

$$J^q + J^g = \frac{1}{2} \{ [T_1(0)^q + T_2(0)^q] + [T_1(0) + T_2(0)]^g \} = \frac{1}{2}$$

We will use these two sum rules as the normalization conditions for the local current

$$N_q \langle x \rangle_R^q + N_g \langle x \rangle_R^g = 1$$

$$N_q J_R^q + N_g J_R^g = \frac{1}{2}$$

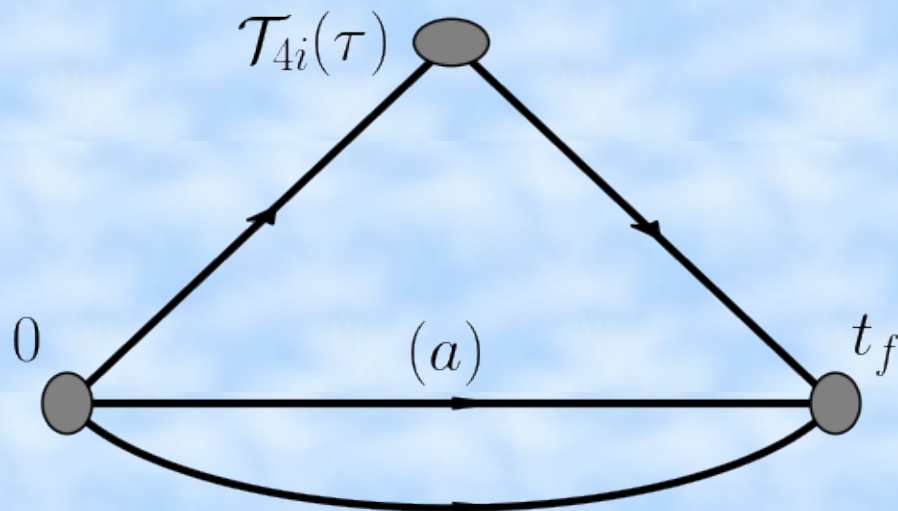
As T2 is too small and noisy

$$N_q = N_g = N$$

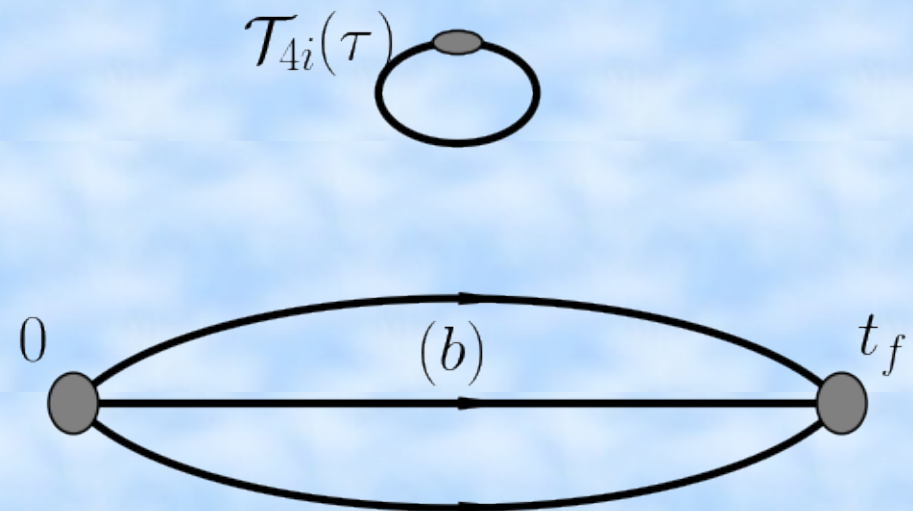
Lattice Simulations

- Lattice

- 32ID--Domain Wall 2+1 Lattice, $32^3 * 64$, $a=0.143$ fm, Pion 172 MeV
- Overlap Fermions with six different valence quark masses



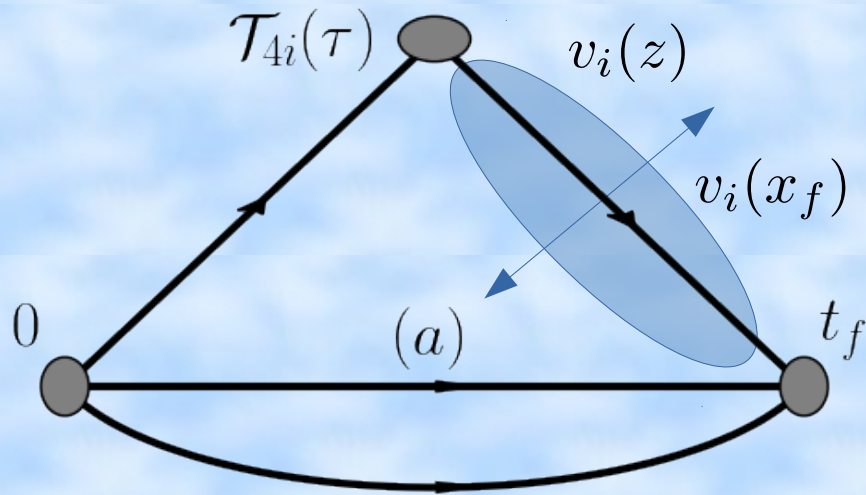
Connected insertions (CI)
for light quarks



Disconnected insertions (DI)
quarks and glue

CI 3pt with FFT

The stochastic-sandwich method with low mode substitution used to estimate the propagator from current to sink can be improved further with FFT



$$S(z|x_f) = S^L(z|x_f) + S^H(z|x_f),$$

$$S^L(z|x_f) = \sum_{\lambda_i \leq \lambda_c} \frac{1}{\lambda_i + m} v_i(z) v_i^\dagger(x_f)$$

The corresponding “low mode” part of the 3pt can be written as

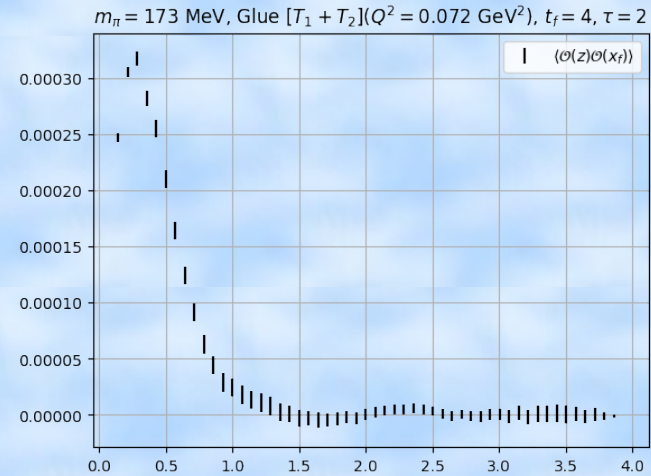
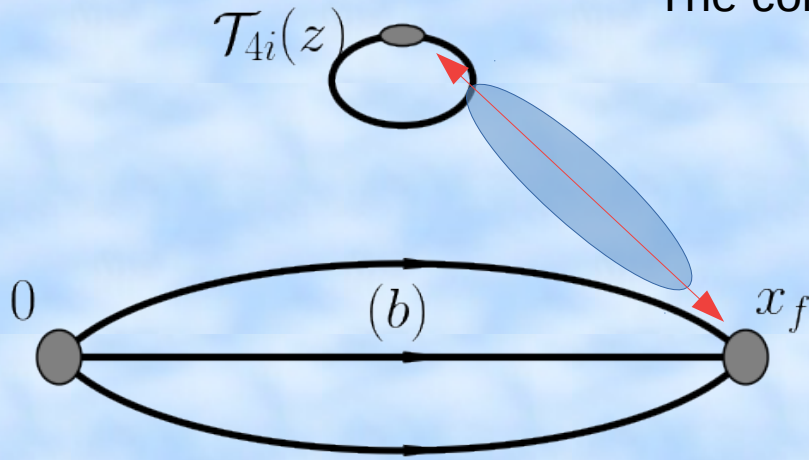
$$C_{\text{CI,3pt}}^{u/d} = \left\langle \sum_{\lambda_i \leq \lambda_c} \text{Tr} \left[\frac{1}{\lambda_i + m} G_i^L(\vec{q}, \tau) F_i^{L,u/d}(\vec{p}_f, t_f) \right] \right\rangle$$

Similar separations could also be done with the stochastic “high mode” part

DI parts

Cluster-decomposition error (CDER) [1] technique are used for DI parts

The correlations satisfy [2] $\langle \mathcal{O}(z)\mathcal{O}(x_f) \rangle_s \leq Ar^{-\frac{2}{3}}e^{-Mr}$

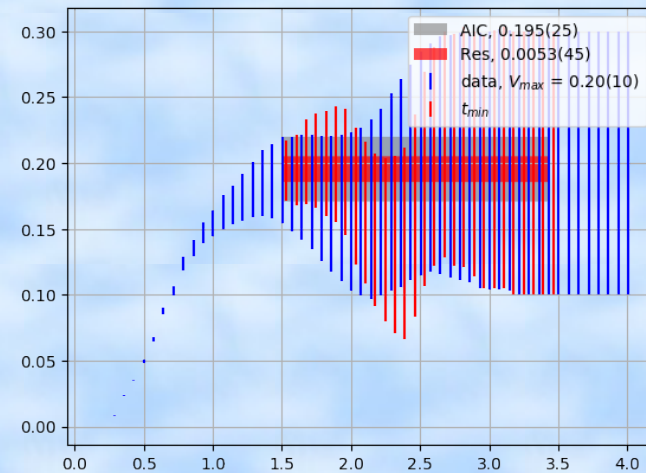


Use Akaike information criterion (AIC) weight factors to average different fit results

$$AIC = \exp \left[-\frac{1}{2}(\chi^2 - 2n_{\text{dof}}) \right]$$

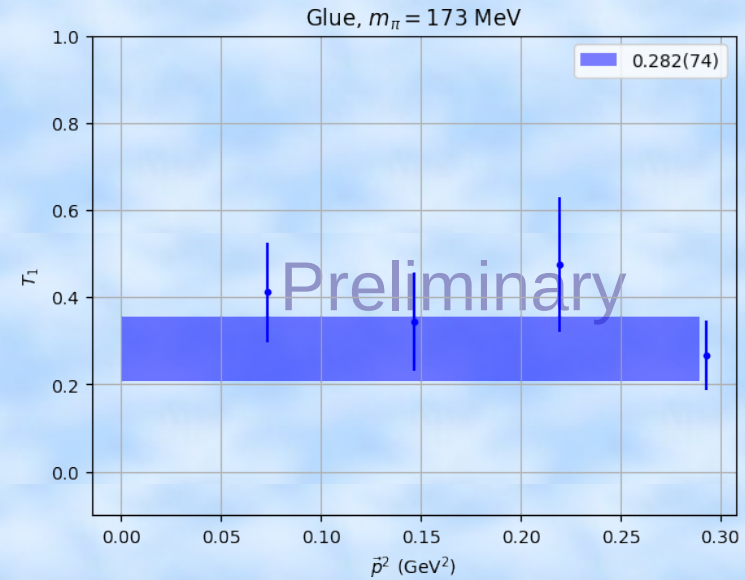
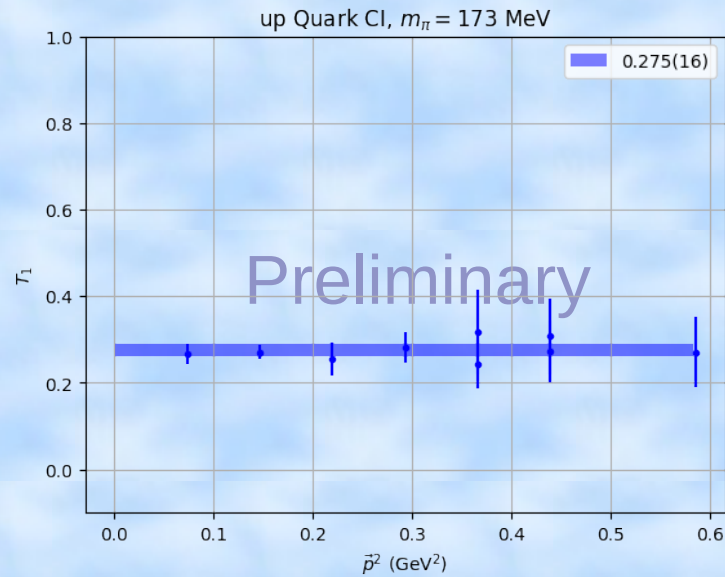
The residue of the correlators from 1.5 fm

$$Res = \sum_{\substack{r < r_{\text{max}} \\ r > r_{\text{cut}}}} Ar^{-\frac{3}{2}}e^{-Mr}$$

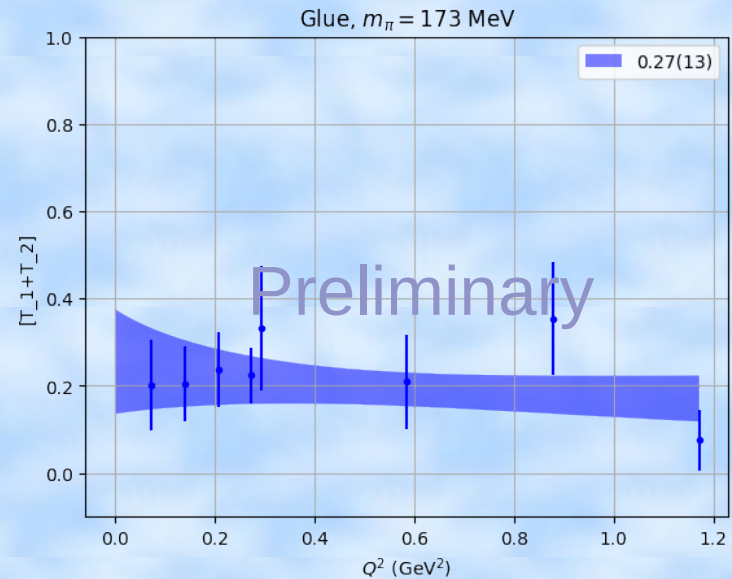
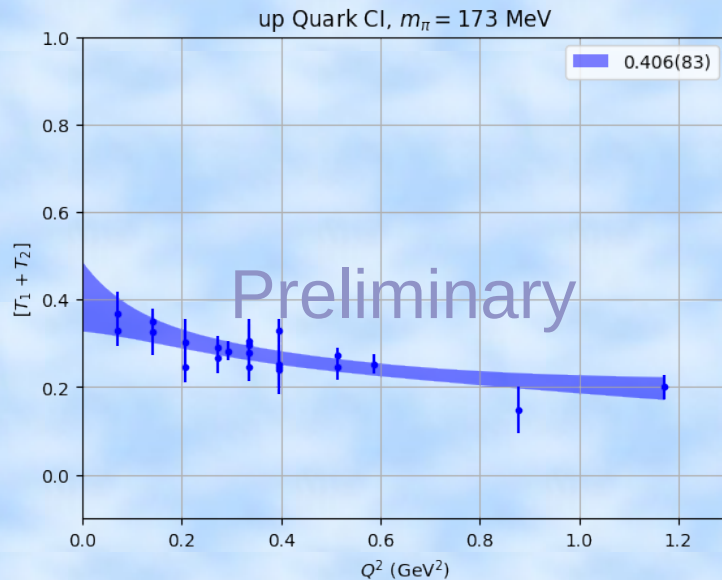


[1] K.-F. Liu, J. Liang, and Y.-B. Yang, Phys. Rev. D 97, 034507 (2018), arXiv:1705.06358
 [2] H. Araki, K. Hepp, and D. Ruelle, Helv. Phys. Acta 35, 164 (1962)

T1(0) and T1+T2 form factors



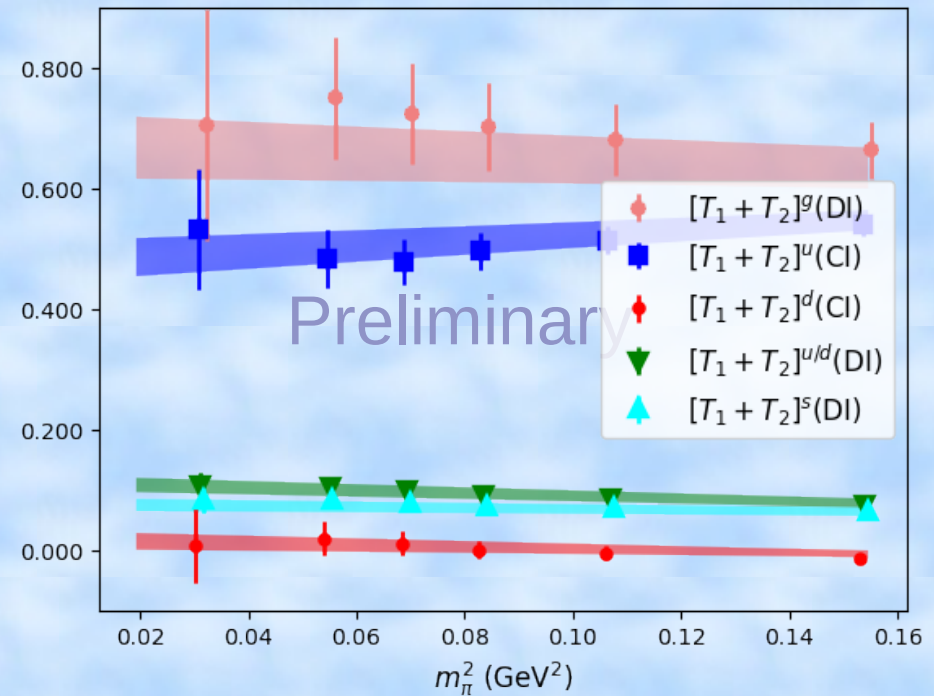
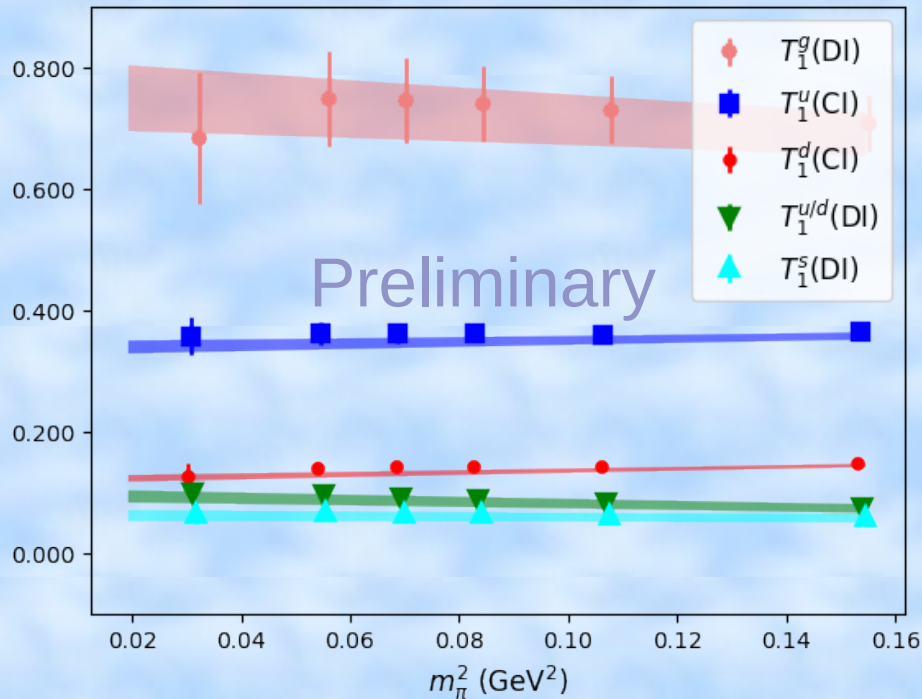
Averaged over results from different nucleon initial momenta for T1(0)



z-expansion fit to extrapolate to zero momentum transfer

T1 and T1+T2 form factors Summary

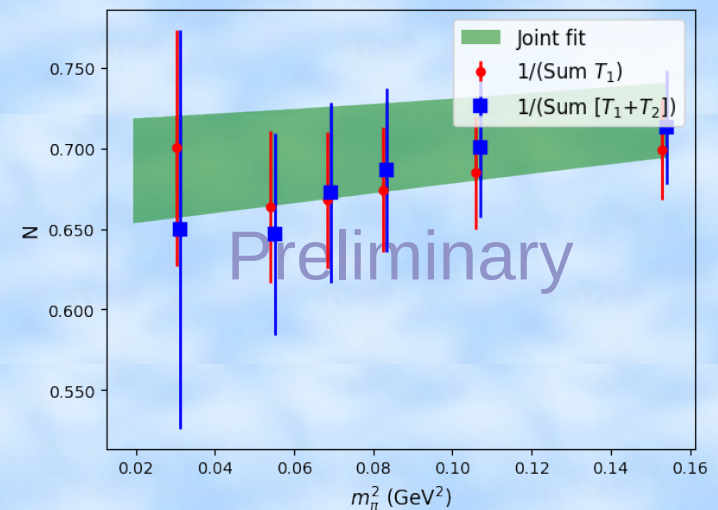
Simple linear extrapolations of each constituents under current statistics



The sum of renormalized T1 and T1+T2 form factors are consistent with each other at all pion masses

$$N \langle x \rangle_R^q + N \langle x \rangle_R^g = 1$$

$$N J_R^q + N J_R^g = \frac{1}{2}$$



Momentum and Angular momentum fractions

Summary table of the CI and DI parts for quark and gluon constituents

Preliminary

| | $u(\text{CI})$ | $d(\text{CI})$ | $[u + d](\text{CI})$ | $u/d(\text{DI})$ | $s(\text{DI})$ | glue |
|---------------------|----------------|----------------|----------------------|------------------|----------------|-----------|
| $\langle x \rangle$ | 0.233(12) | 0.0841(51) | 0.317(15) | 0.0637(58) | 0.0419(56) | 0.514(20) |
| $2J$ | 0.333(23) | 0.0097(96) | 0.343(24) | 0.0739(73) | 0.0512(62) | 0.458(22) |

Compare with previous calculation and phenomenological global fit results

Preliminary

| | u | d | $[u - d]$ | s | glue |
|-----------------------------------|-----------|------------|-----------|------------|-----------|
| $\langle x \rangle$ | 0.296(12) | 0.1477(69) | 0.149(10) | 0.0419(56) | 0.514(20) |
| $\langle x \rangle_{[1]}$ | 0.307(35) | 0.160(48) | 0.151(40) | 0.051(26) | 0.482(84) |
| $\langle x \rangle_{\text{CT14}}$ | 0.348(5) | 0.190(5) | 0.158(6) | 0.035(9) | 0.416(9) |

- [1] Y-B Yang, J. Liang, et al., χ QCD Collaboration, Phys. Rev. Lett. 121, 212001 (2018)
 [2] S. Dulat, et al., Phys. Rev. D, 93(3):033006, (2016)
 [3] M. Deka, T. Doi, Y-B Yang, et. al., χ QCD collaboration, PRD91, 014505 (2015)

Further Calculations

- FFT and low mode substitution has been successfully used with stochastic-sandwich method for CI to reach better statistics
- CDER technique greatly increase DI statistics with systematic errors under control
- A complete calculation of proton momentum and angular momentum fractions at several overlap valence pion masses has been done on one Lattice
- Extend the calculation to other lattice spacing and volumes to extrapolate to physical limit
- Extend to all kinematics to obtain T1, T2, T3 form factors at different Q^2

Thank You