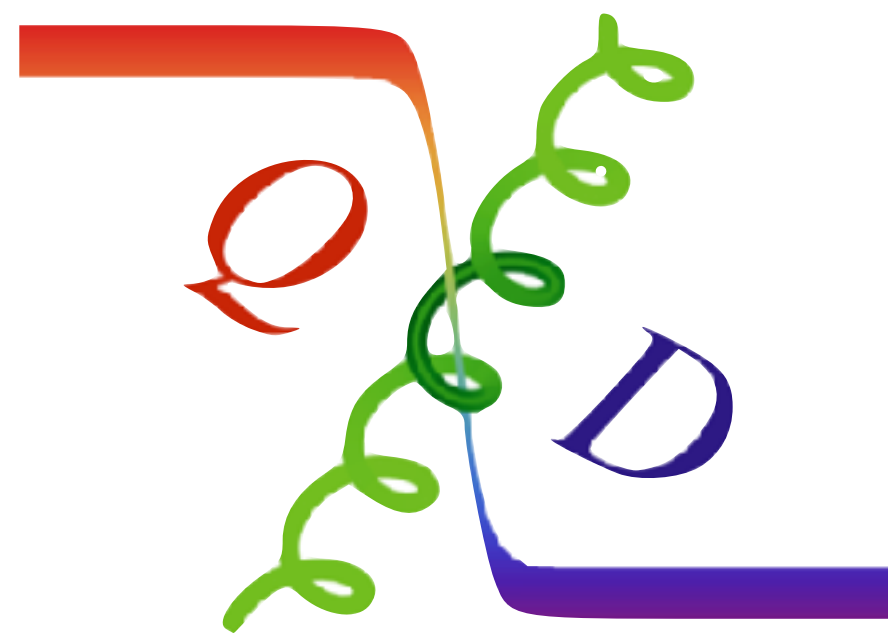


MOM and SMOM renormalization using the chiral fermion



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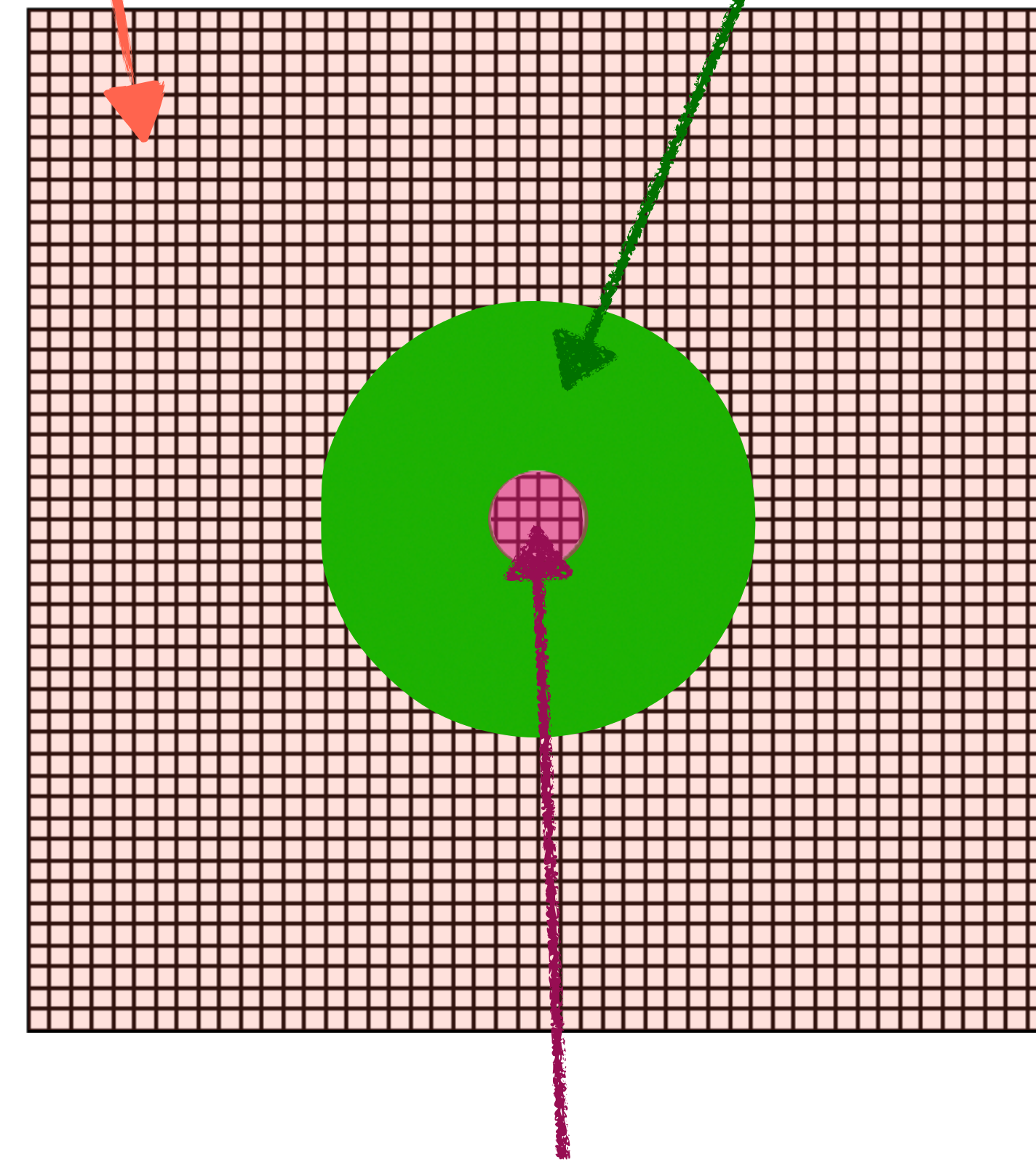


Renormalization

- The phenomenological studies are carried out under the modified minimum subtraction scheme ($\overline{\text{MS}}$).
- But it requires the modification of the dimension and then can not be done on the lattice!
- Thus an intermediate regularization-independent scheme is essential to convert the lattice result to $\overline{\text{MS}}$ scheme.

IR region can only be calculated using the Lattice regularization

Intermediate region which can be calculated using kinds of regularizations



UV region which suffers from obvious discretizing effects

Effective renormalization constant

- Obtain the regularization independent renormalization constant non-perturbatively:

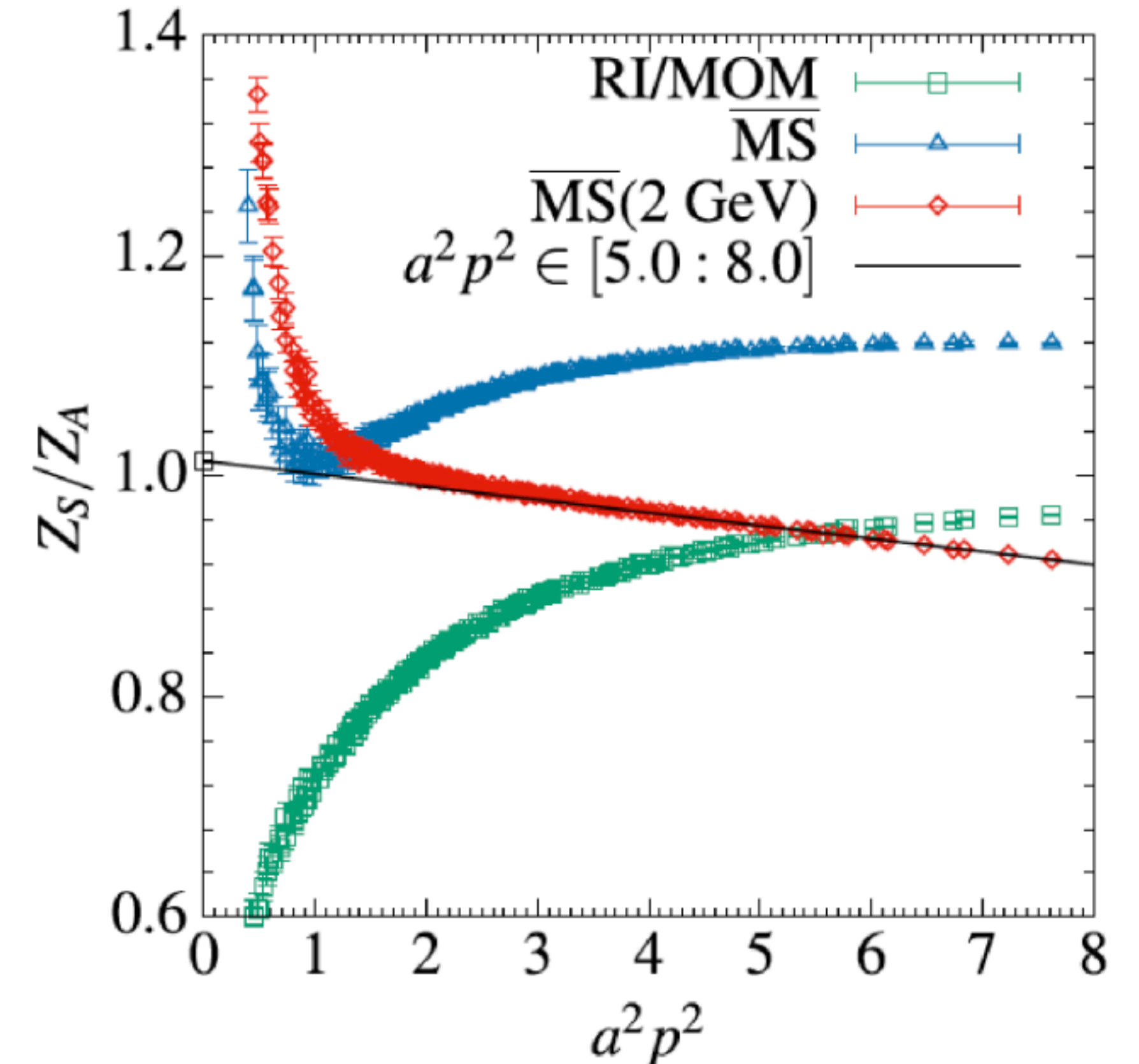
$$Z_S^{\text{MOM}}(Q, a) = 1 - \frac{\alpha_s C_F}{4\pi} [-3\log(a^2 Q^2) - \xi + b_S] + \mathcal{O}(\alpha_s^2, a^2 Q^2)$$

- Calculate the matching coefficient perturbatively and obtain the result at $\overline{\text{MS}}$ scale Q :

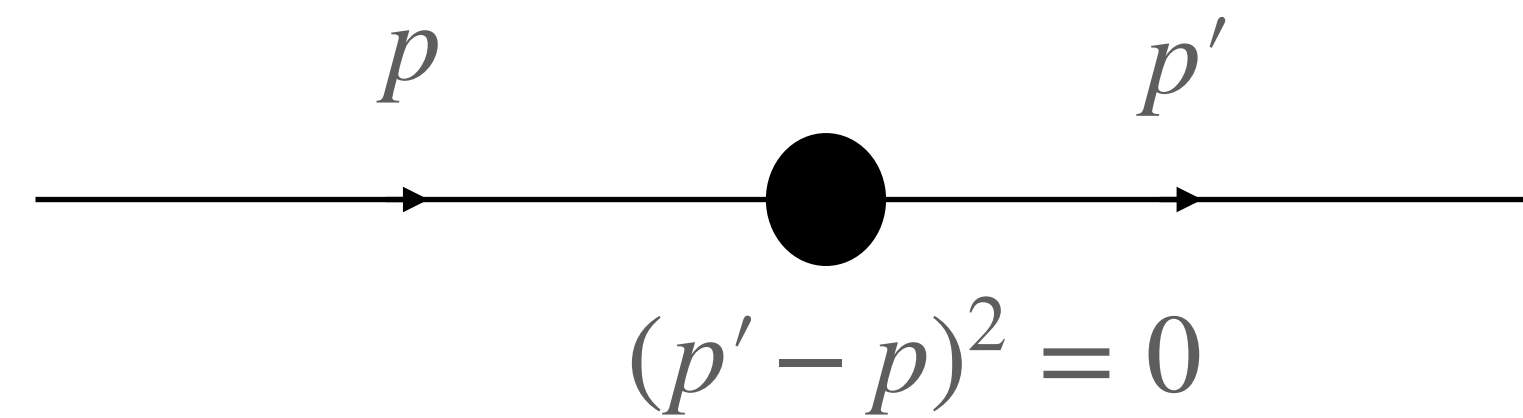
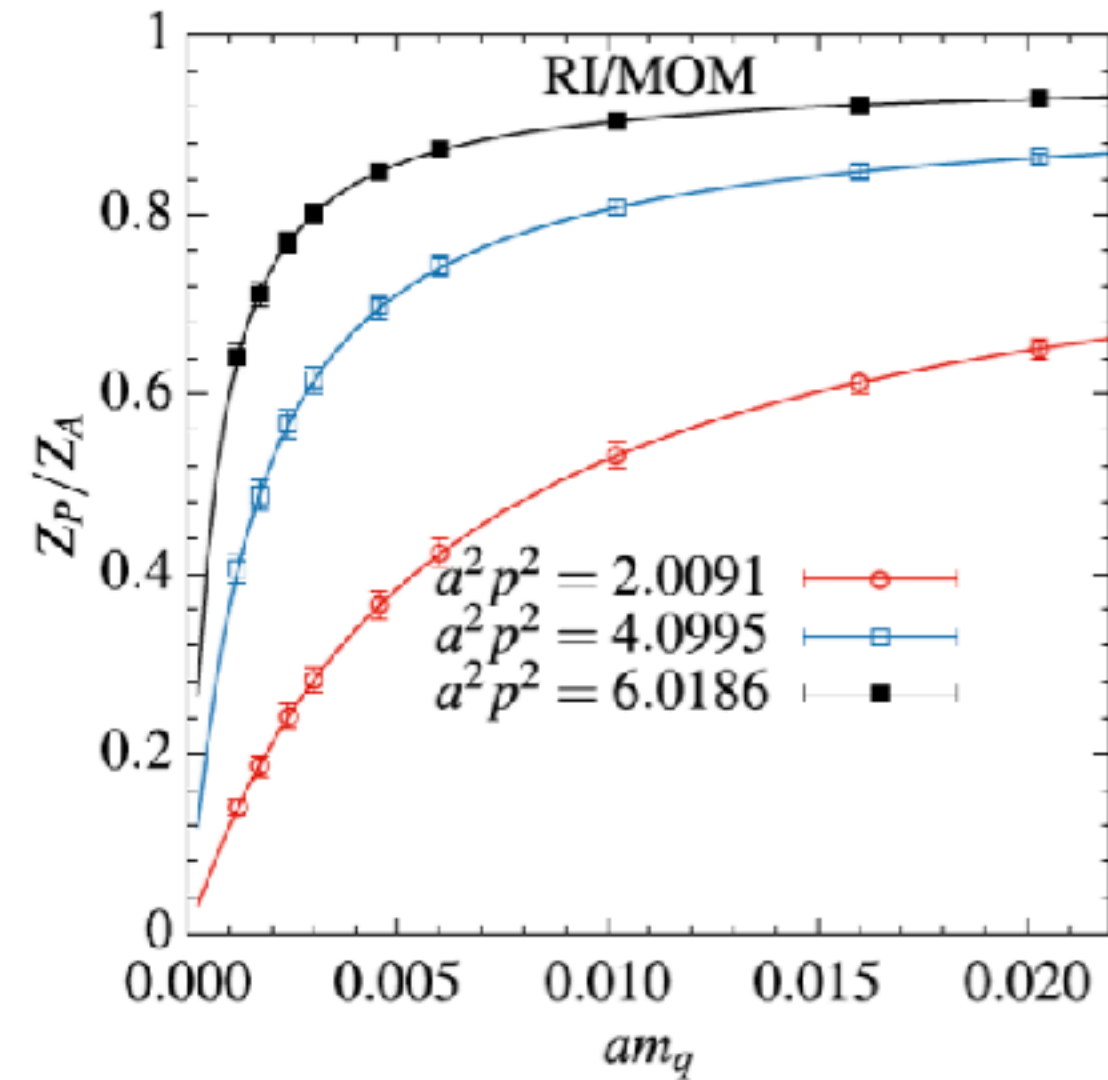
$$Z_S^{\overline{\text{MS}}}(Q, a) = 1 - \frac{\alpha_s C_F}{4\pi} [-3\log(a^2 Q^2) - 5 + b_S] + \mathcal{O}(\alpha_s^2, a^2 Q^2)$$

- Obtain the result at $\overline{\text{MS}}$ scale μ with the scale evolution:

$$Z_S^{\overline{\text{MS}}}(\mu, a) = 1 - \frac{\alpha_s C_F}{4\pi} [-3\log(a^2 \mu^2) - 5 + b_S] + \mathcal{O}(\alpha_s^2, a^2 Q^2)$$

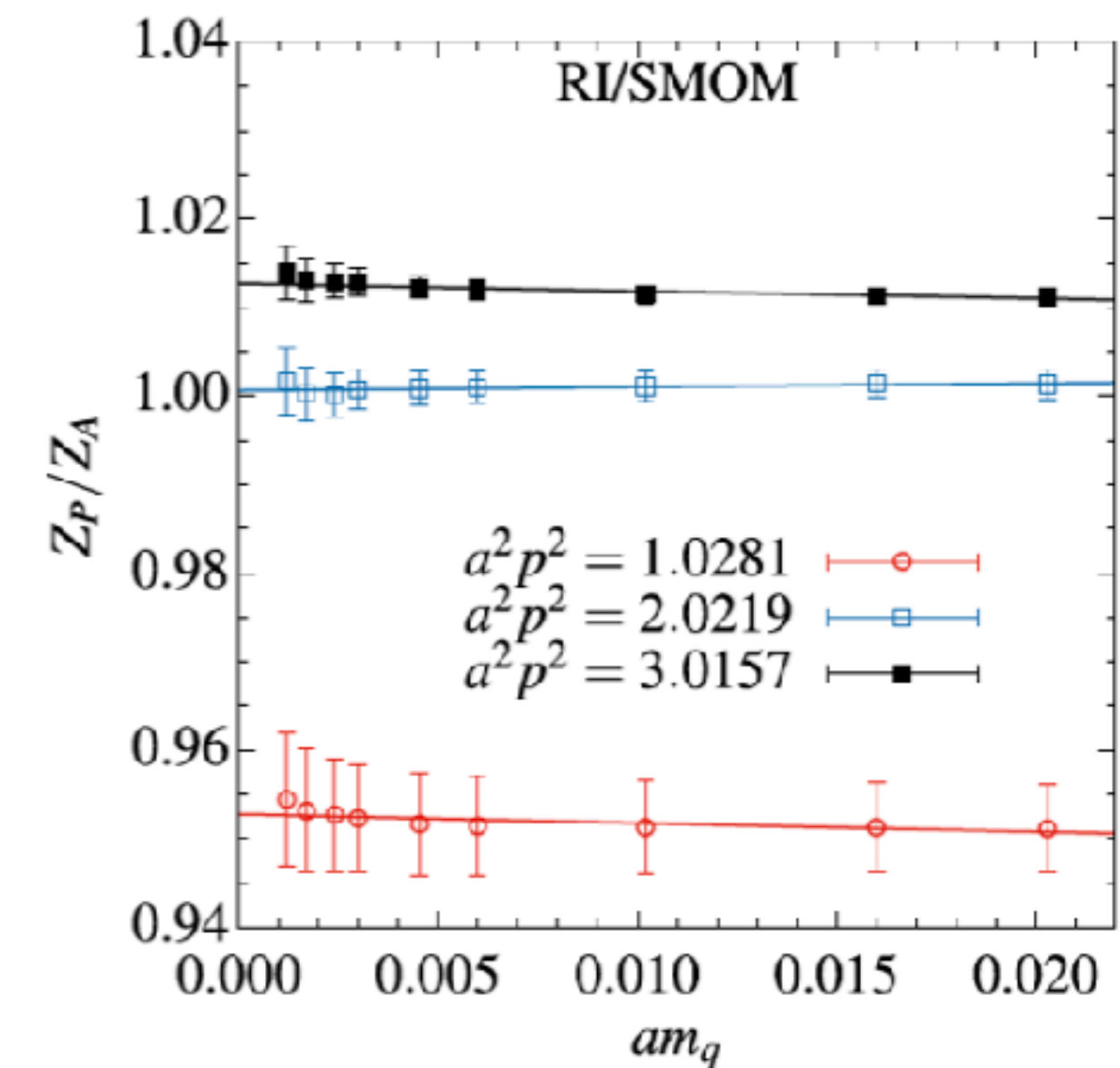
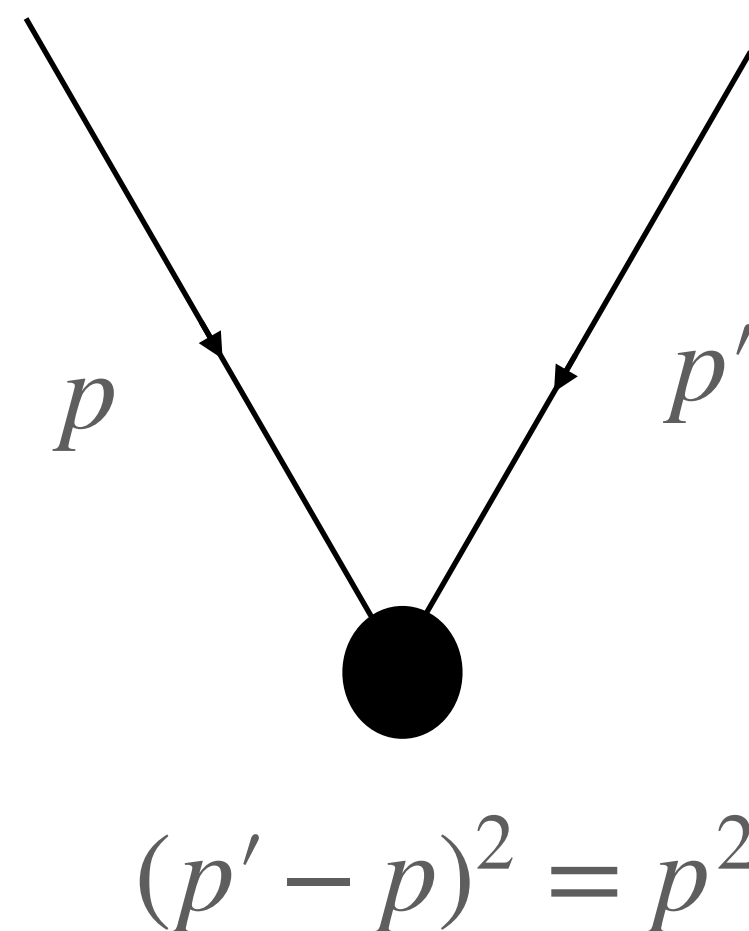


MOM v.s. SMOM



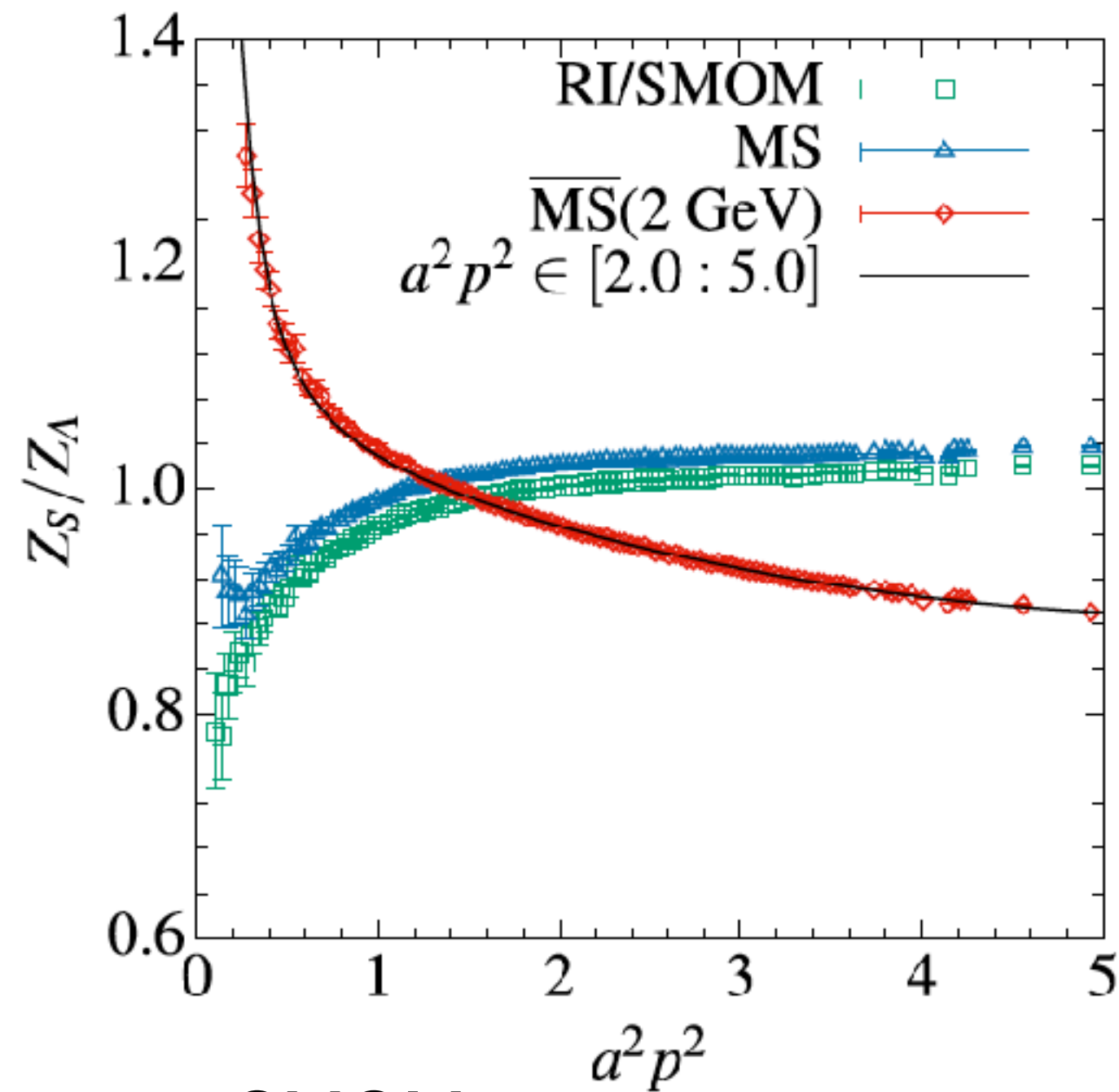
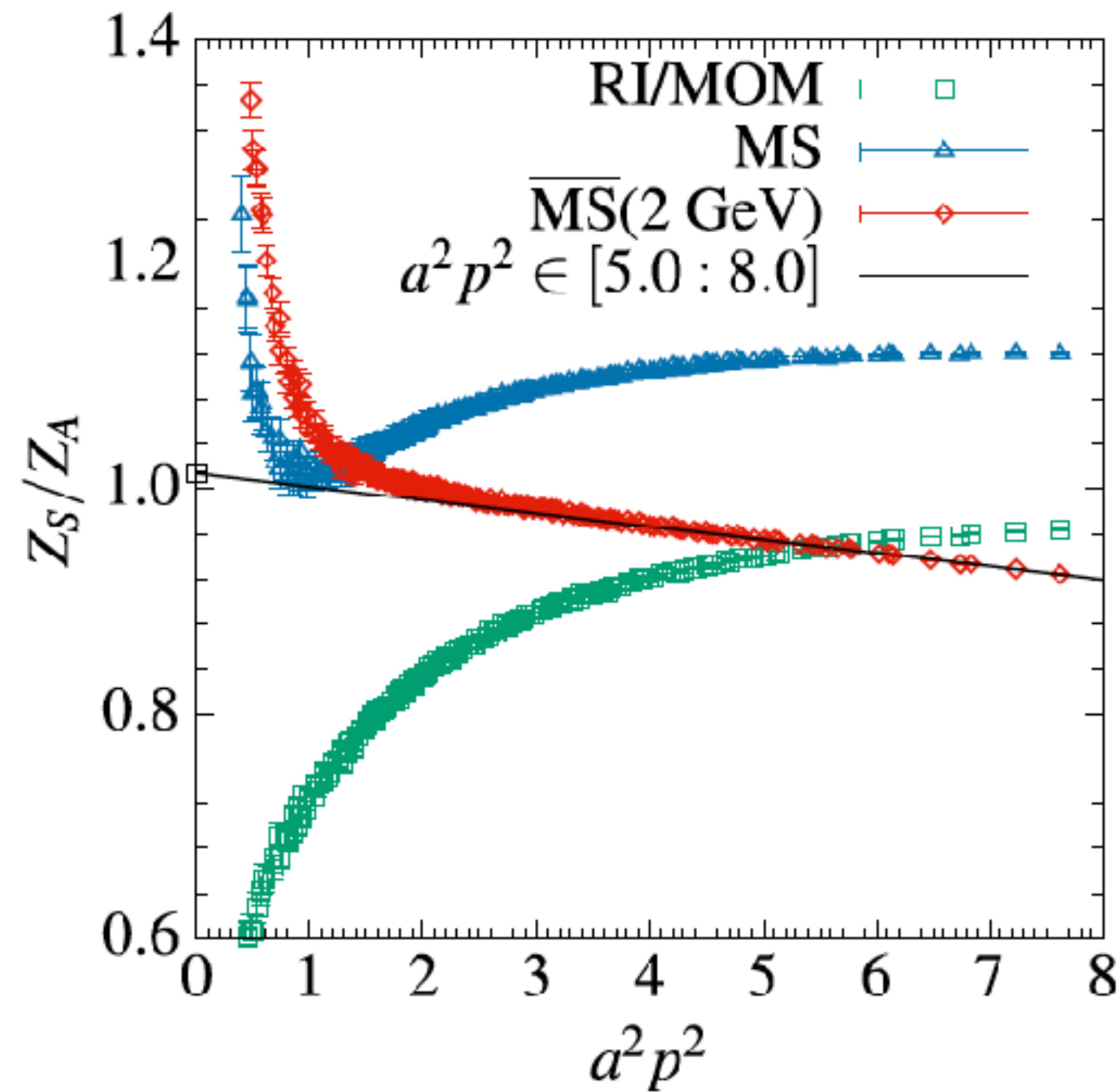
- **MOM scheme:** The zero momentum transfer at the current can introduce additional mixing/non-perturbative effect;

- **SMOM scheme:** Requiring the momentum transfer to be non-zero can avoid such a problem.



Our previous SMOM result

Yujiang Bi, et.al, χ QCD, PRD97(2018)094501, 1710.08678



- $a=0.114 \text{ fm}$
- $48^3 \times 96$
- Overlap on DWF
- $m_\pi^{\text{sea}} = 139 \text{ MeV}$

• MOM:

• Linear $a^2 p^2$ dependence

• Larger matching correction:

$$\frac{Z_S^{\overline{\text{MS}}}}{Z_S^{\text{MOM}}}(4\text{GeV}) = 1 + 0.092 + 0.047 + 0.028 + \mathcal{O}(\alpha_s^4)$$

• SMOM:

• Non-linear $a^2 p^2$ dependence,

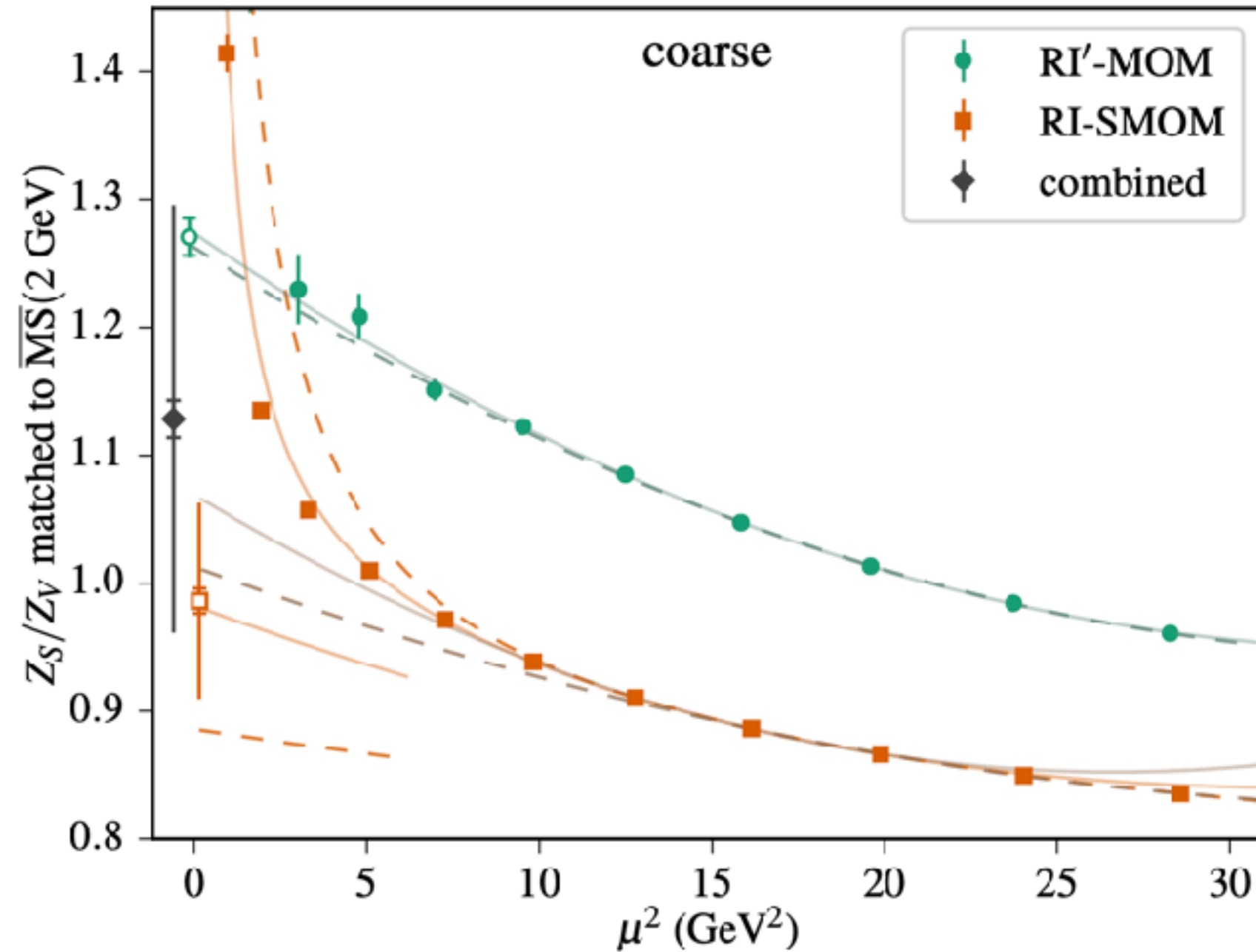
• Smaller matching correction:

$$\frac{Z_S^{\overline{\text{MS}}}}{Z_S^{\text{SMOM}}}(4\text{GeV}) = 1 + 0.011 + 0.003 + \mathcal{O}(\alpha_s^3)$$

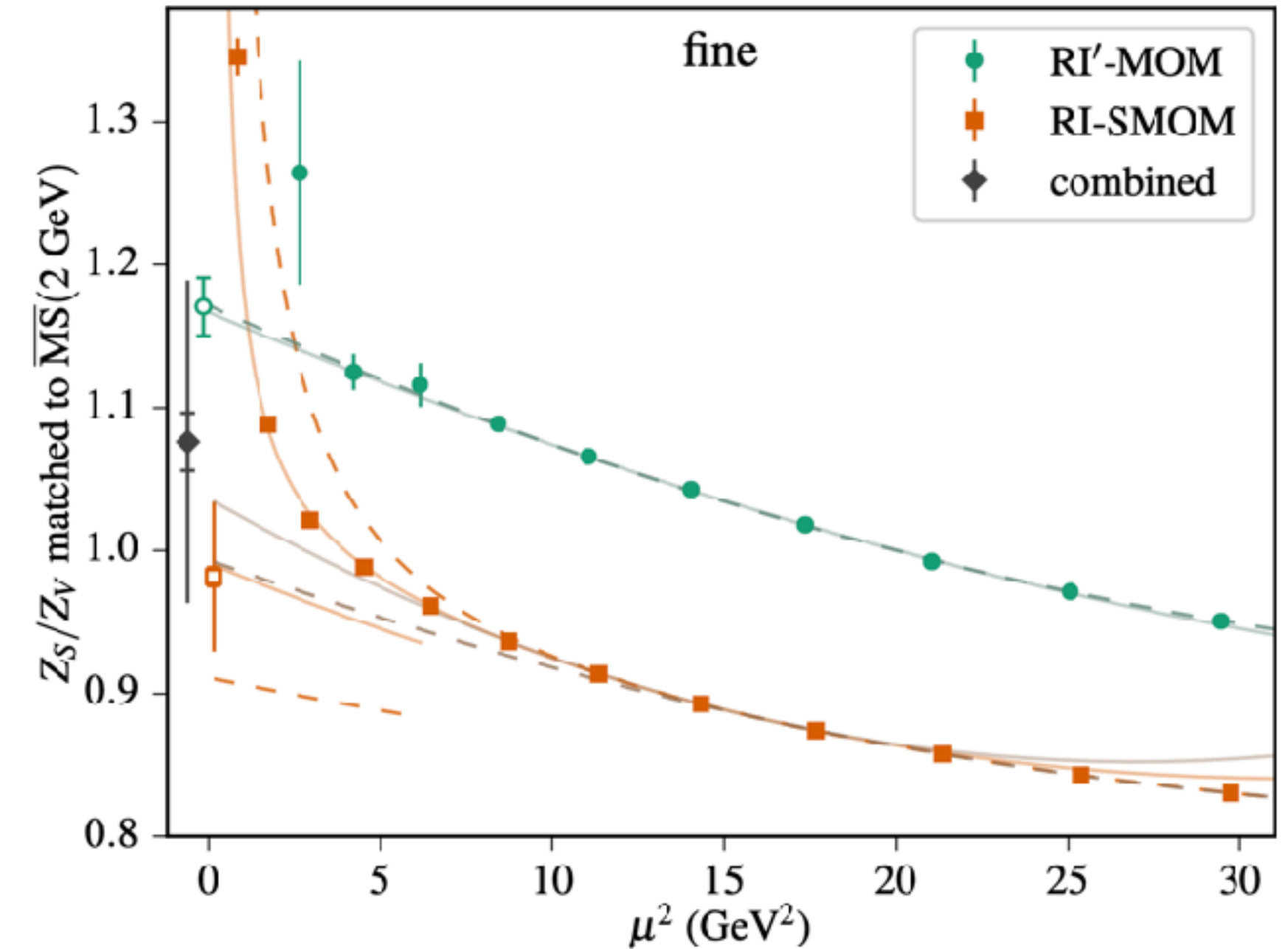
Observation from the other group

- Similar issue is also observed using the clover fermion;
- While the difference between the MOM and SMOM scheme seems to be even larger.

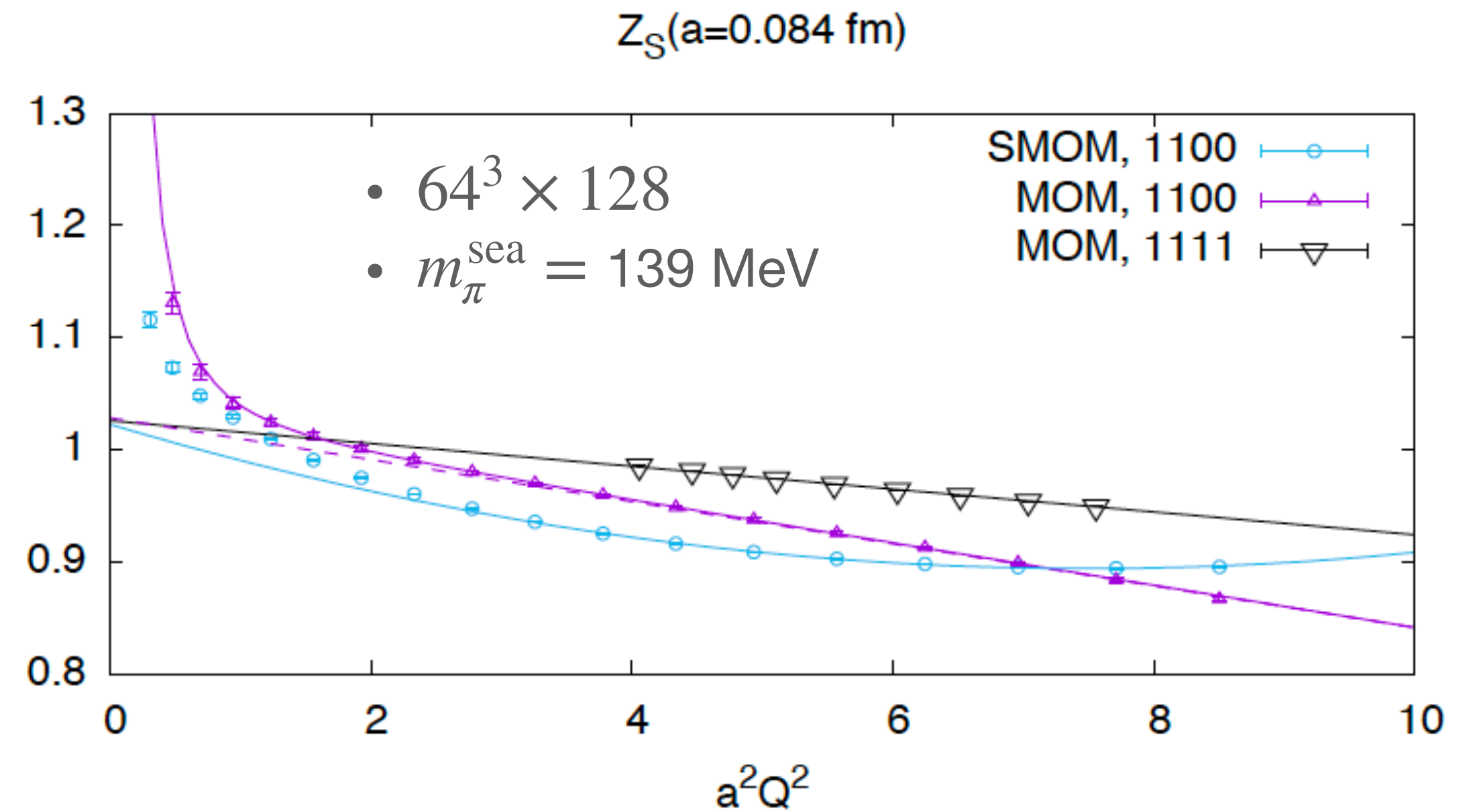
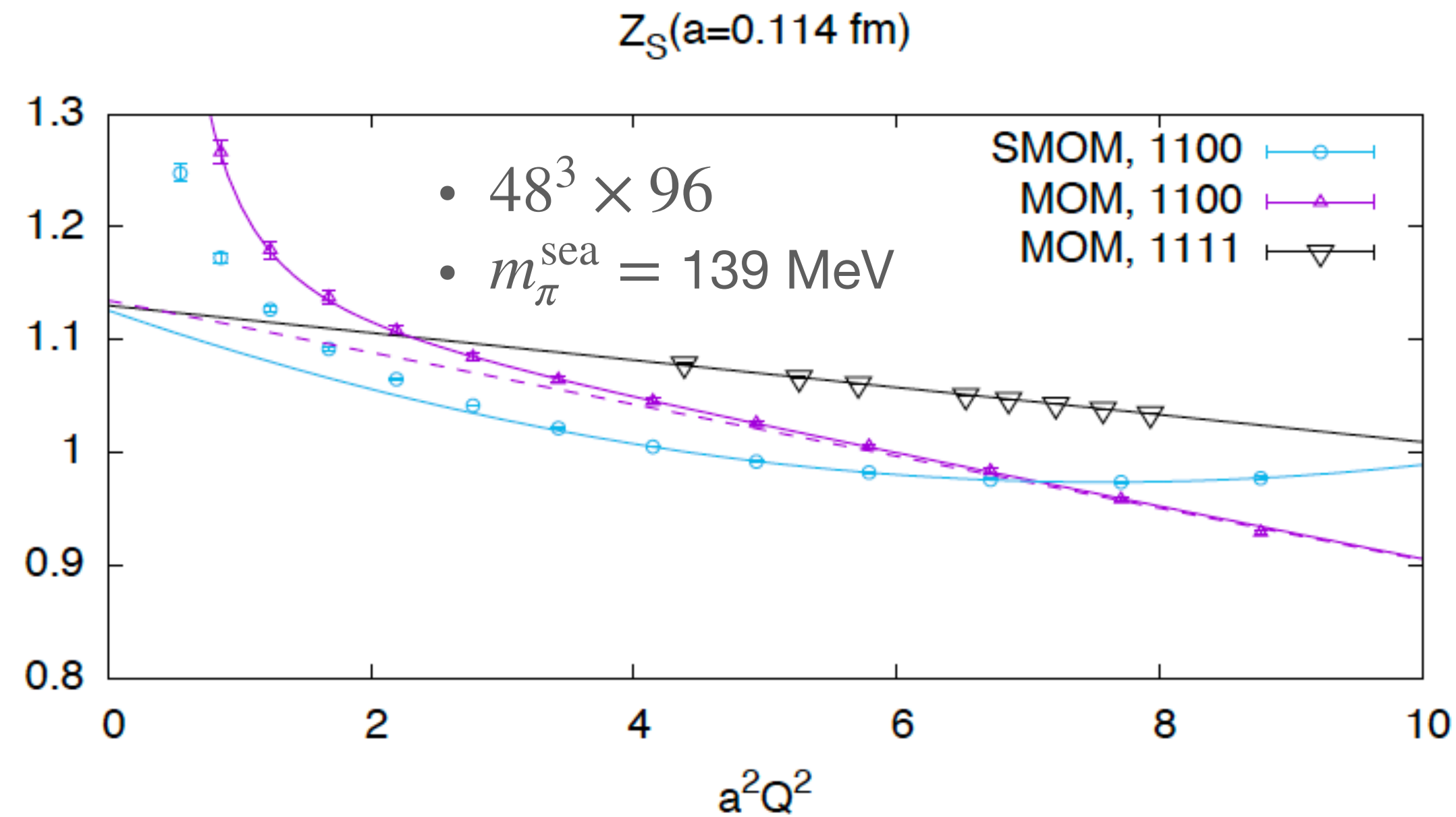
- $a=0.116$ fm
- 48^4
- Clover
- $m_\pi^{\text{sea}} = 137$ MeV



- $a=0.093$ fm
- 64^4
- Clover
- $m_\pi^{\text{sea}} = 133$ MeV



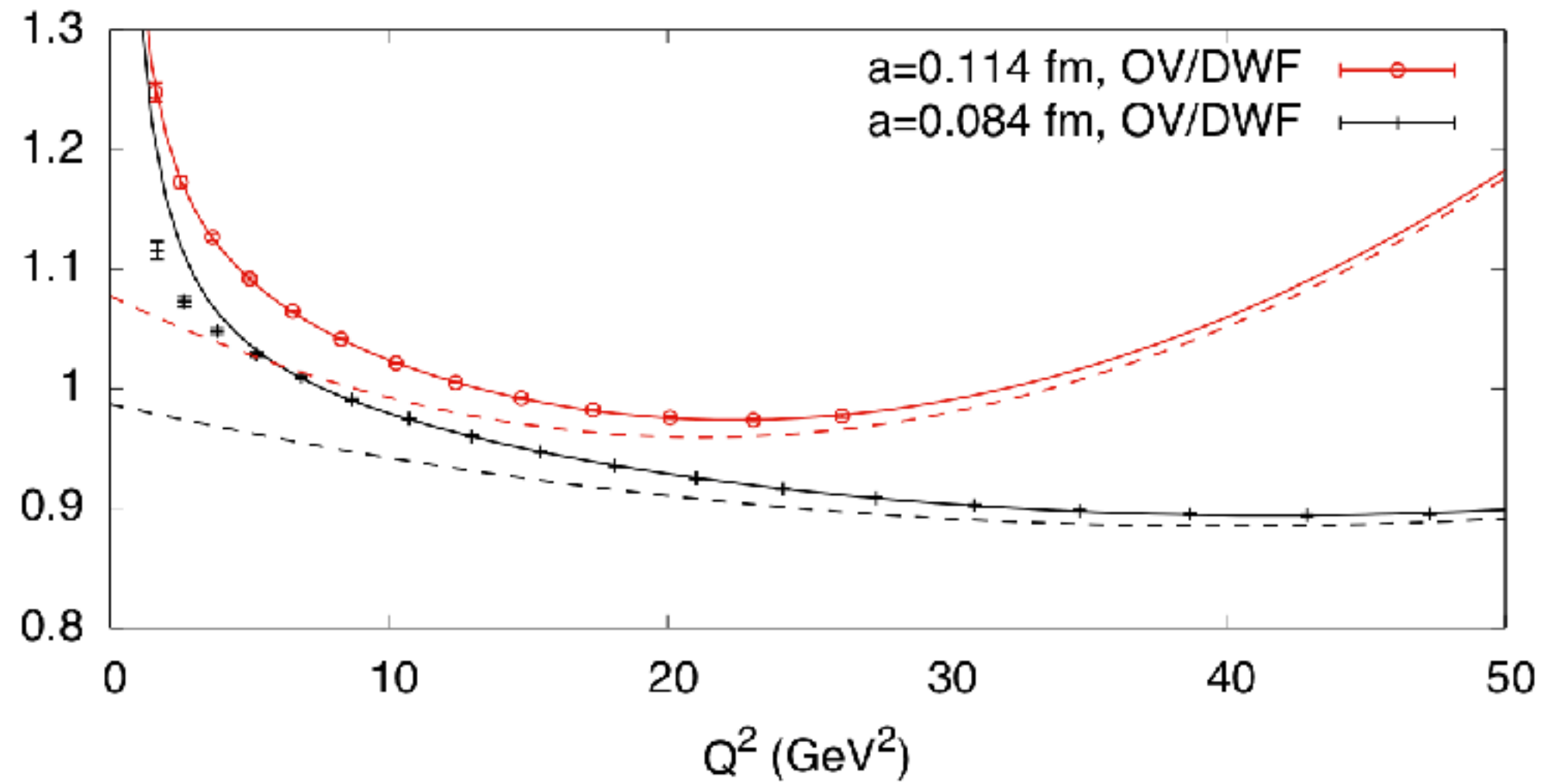
OV@DWF at different lattice spacing



- Use the polynomial form in a range of $Q^2 \in [Q_{min}^2, 9/a^2]$ and tune the Q_{min} to make $\chi^2/\text{d.o.f} \sim 1$ for the SMOM result;
- The results from different scheme and momentum setups seem to be consistent with each other;
- **But.....**

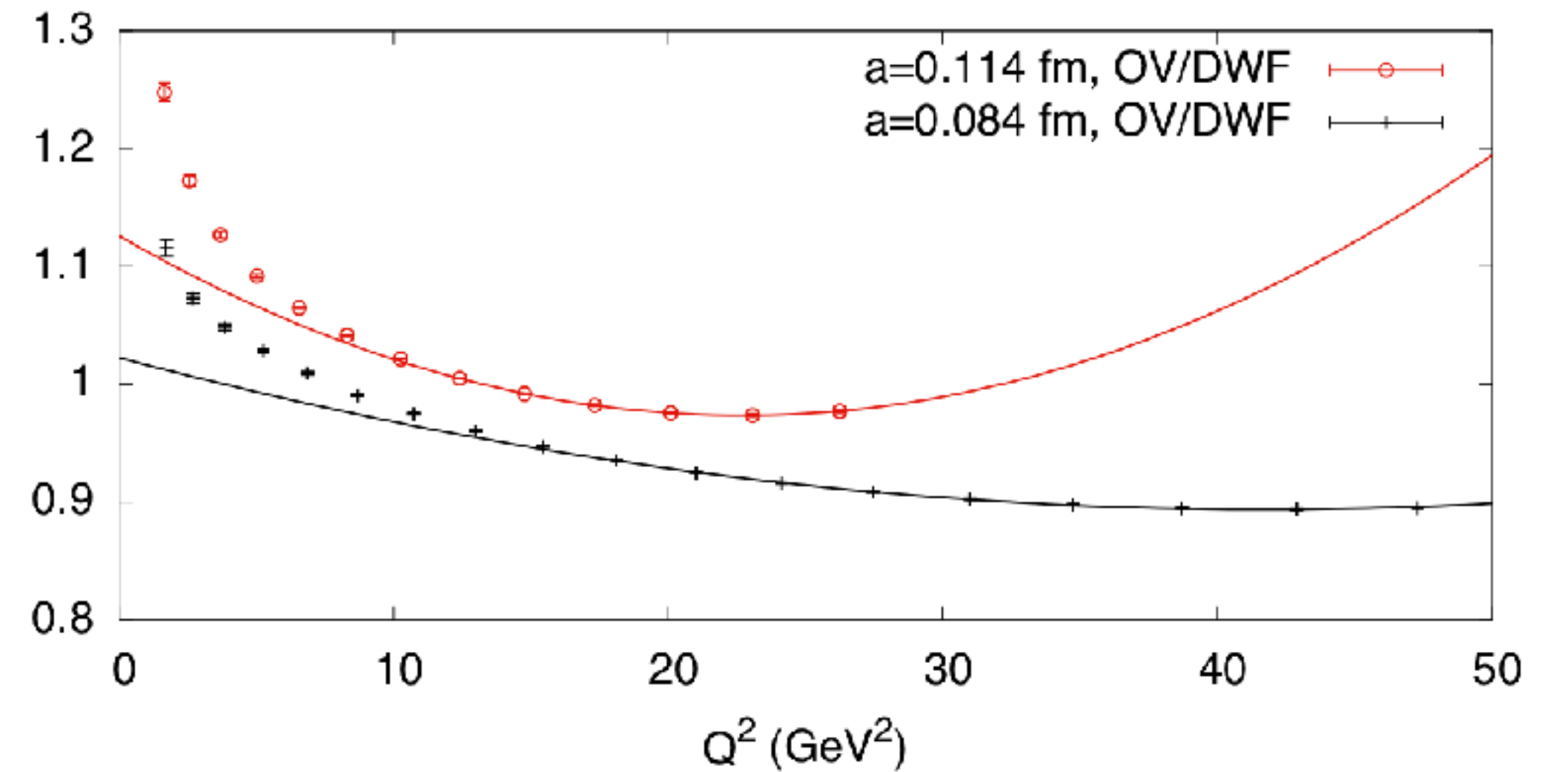
Fitting Ansatz Dependence

$Z_S(Q,a)$, with $1/Q^2$ term



$$Z_S^{\text{SMOM},b}(Q^2) = \frac{c_{-1}^{\text{SMOM}}}{Q^2} + Z_S + c_1^{\text{SMOM}} a^2 Q^2 + c_2^{\text{SMOM}} a^4 Q^4$$

$Z_S(Q,a)$, without $1/Q^2$ term



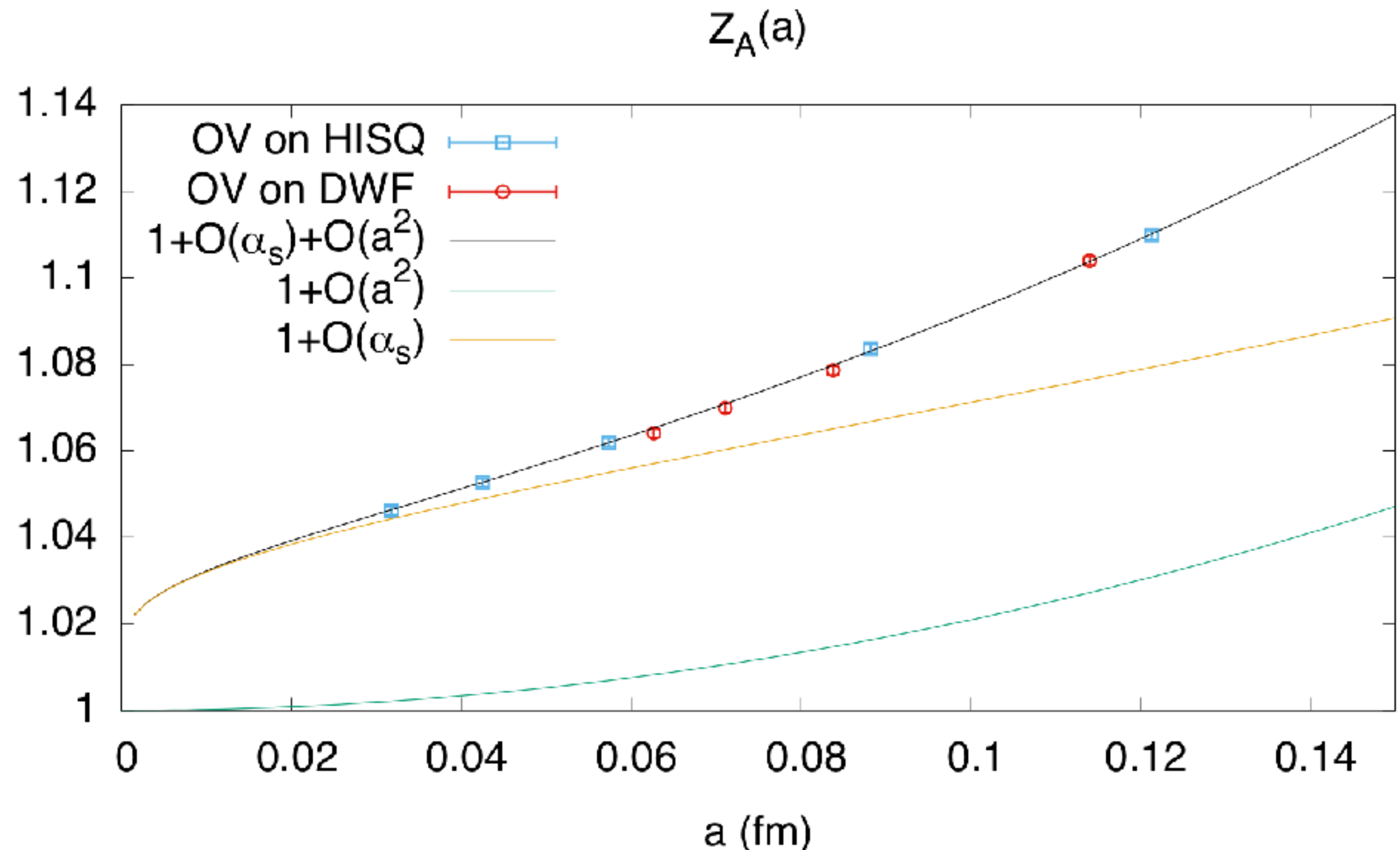
$$Z_S^{\text{SMOM},a}(Q^2) = Z_S + c_1^{\text{SMOM}} a^2 Q^2 + c_2^{\text{SMOM}} a^4 Q^4$$

- Fit in a range of $Q^2 \in [Q_{min}^2, 9/a^2]$ and tune the Q_{min} to make $\chi^2/\text{d.o.f} \sim 1$;
- The discretization errors are close to each other in different cases;
- The difference between Z_S with different fitting ansatz become closer at smaller lattice spacing.

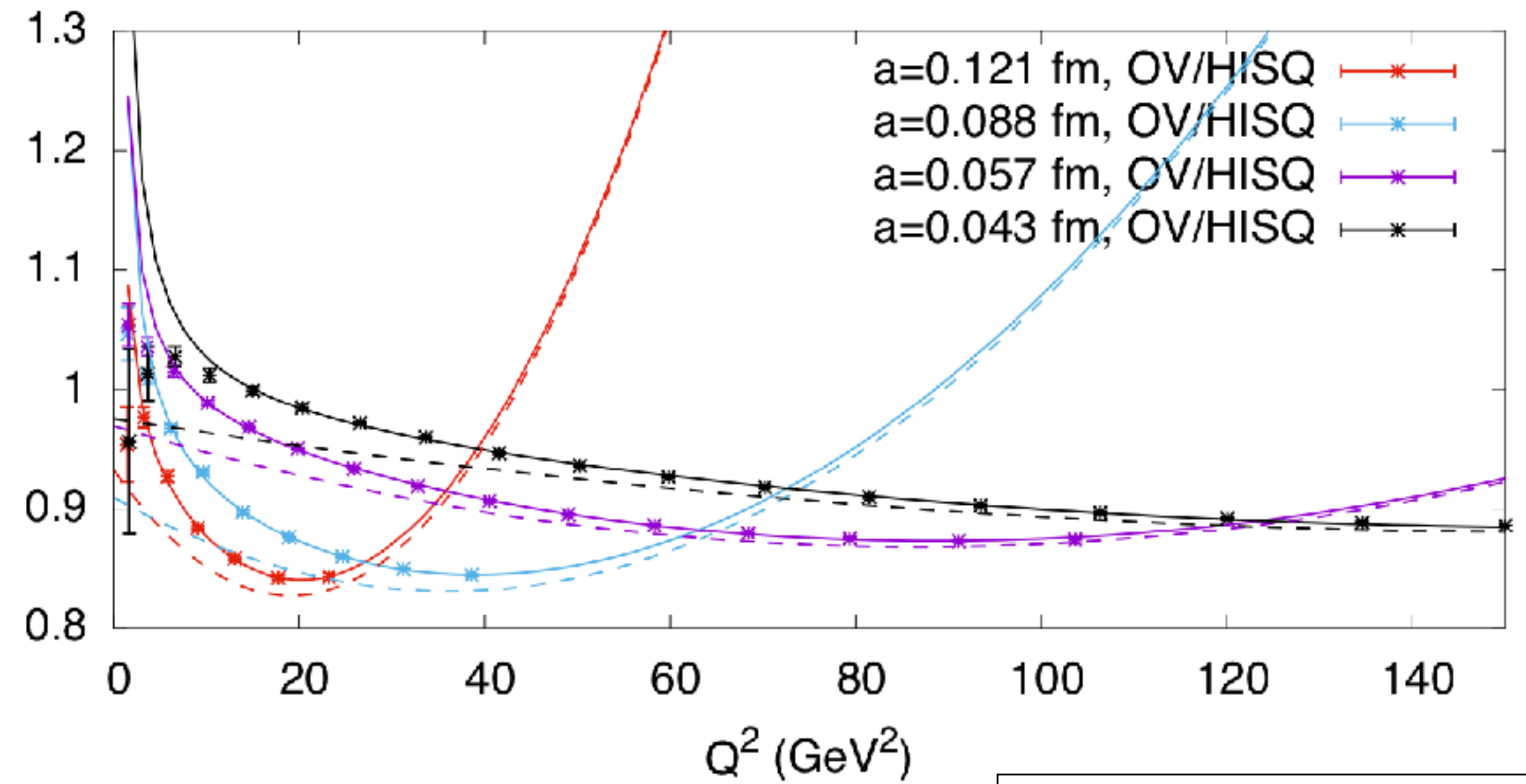
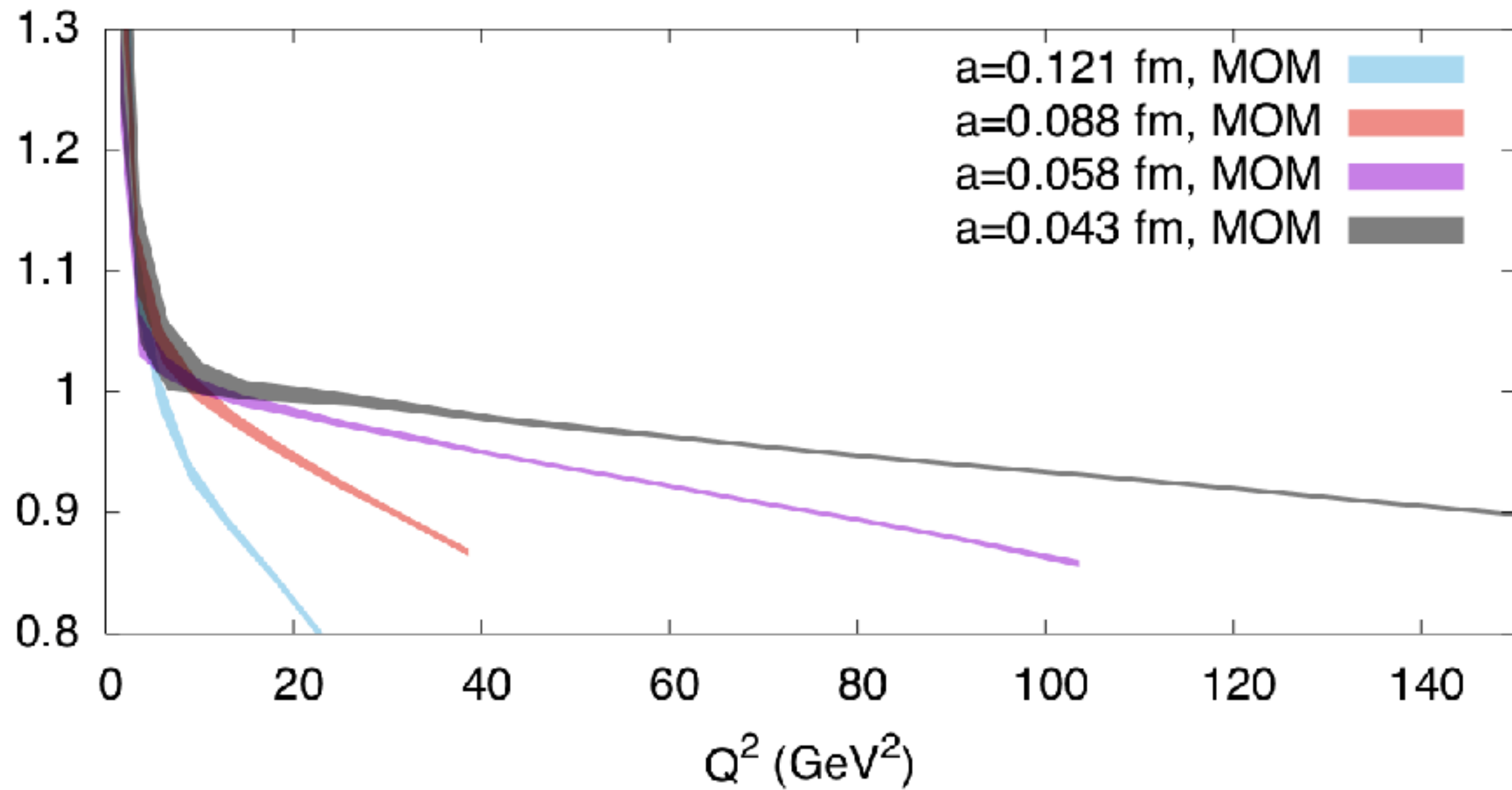
a (fm)	$ Q_{min} $ (GeV)	c_1^{SMOM} (GeV ²)	Z_S	c_1^{SMOM}	c_2^{SMOM}
0.114 fm	3.2		1.126(2)	-0.040(1)	0.0026(1)
	1.2	0.32(1)	1.078(3)	-0.033(1)	0.0024(1)
0.084 fm	4.1	–	1.023(2)	-0.034(1)	0.0023(1)
	2.0	0.37(2)	0.987(3)	-0.028(1)	0.0020(1)

(Axial-)Vector current normalization

- The axial-vector current normalization constants using the overlap fermion seems to be insensitive to the fermion and gauge action in sea;
- We can consider the OV@HISQ cases to further investigate the lattice spacing dependence.



OV@HISQ at 310 MeV



Fangcheng He, et.al, χ QCD, in preparation

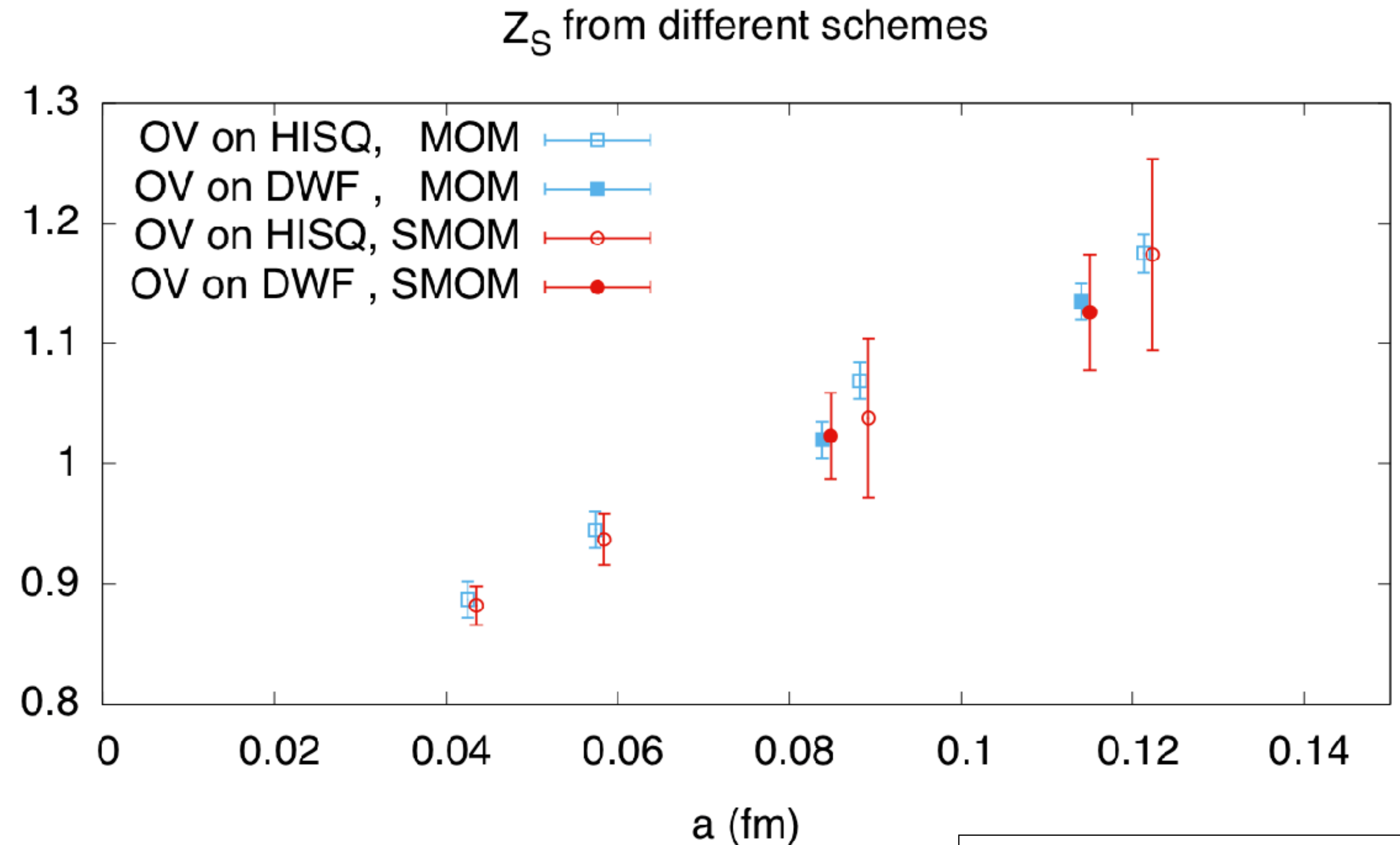
$Z_S^{\overline{\text{MS}}}(2\text{GeV})$ is normalized by its $a^2 Q^2$ extrapolated value to show the residual Q^2 dependence:

- The residual Q^2 dependence becomes weaker at smaller lattice spacing.

$Z_S^{\overline{\text{MS}}}(2\text{GeV})$ results

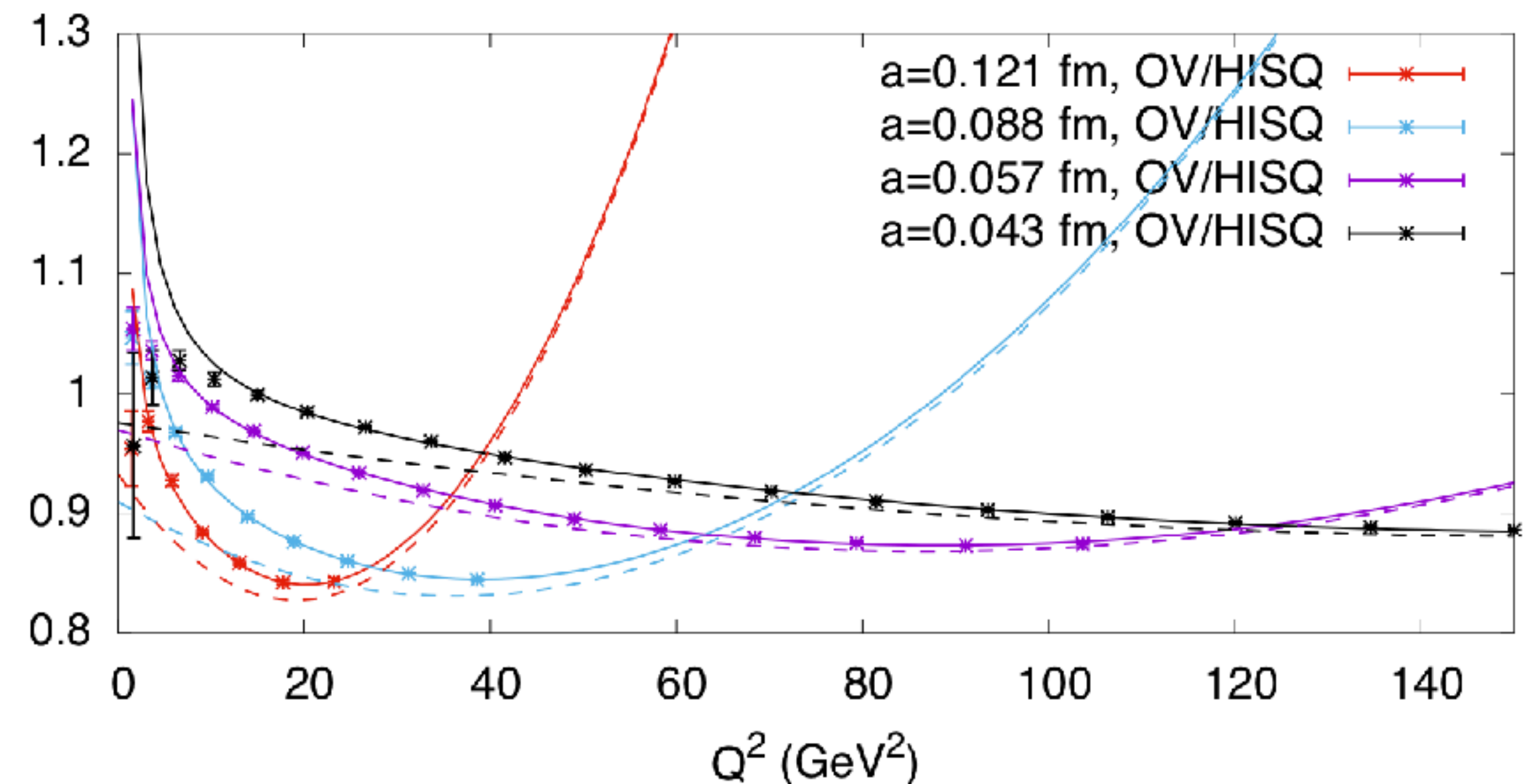
The uncertainty is majorly the systematic one from:

- Missing 4-loop matching ($\sim 1.5\%$ at 4 GeV) in the MOM case;
- Difference from two kinds of the fitting ansatz in the SMOM case.
- Such an uncertainty would vanishes for both the schemes in the continuum and provide consistent $Z_S^{\overline{\text{MS}}}(2\text{GeV})$.



Summary

- The SMOM renormalization for the scalar current can have additional systematic uncertainty from the fitting ansatz;
- Such an uncertainty would vanishes in the continuum, if we drop the small Q and do the extrapolation from large Q .
- It would be essential to if anyone is using the SMOM scheme for the scalar current/quark mass renormalization.



Chiral condensate in the $N_f = 2$ chiral limit

- Blue for pion mass spectrum;
- Orange for the Dirac operator spectrum;
- Our result seems to be much smaller the previous results using the SMOM scheme;
- The recent ETM result using the MOM scheme is in the middle.

