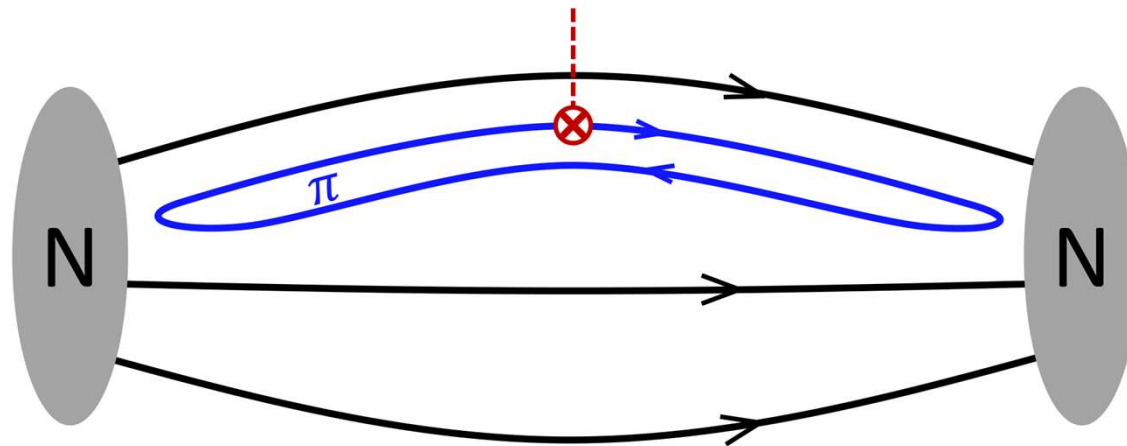


# Flavor Diagonal Nucleon Charges

Sungwoo Park,

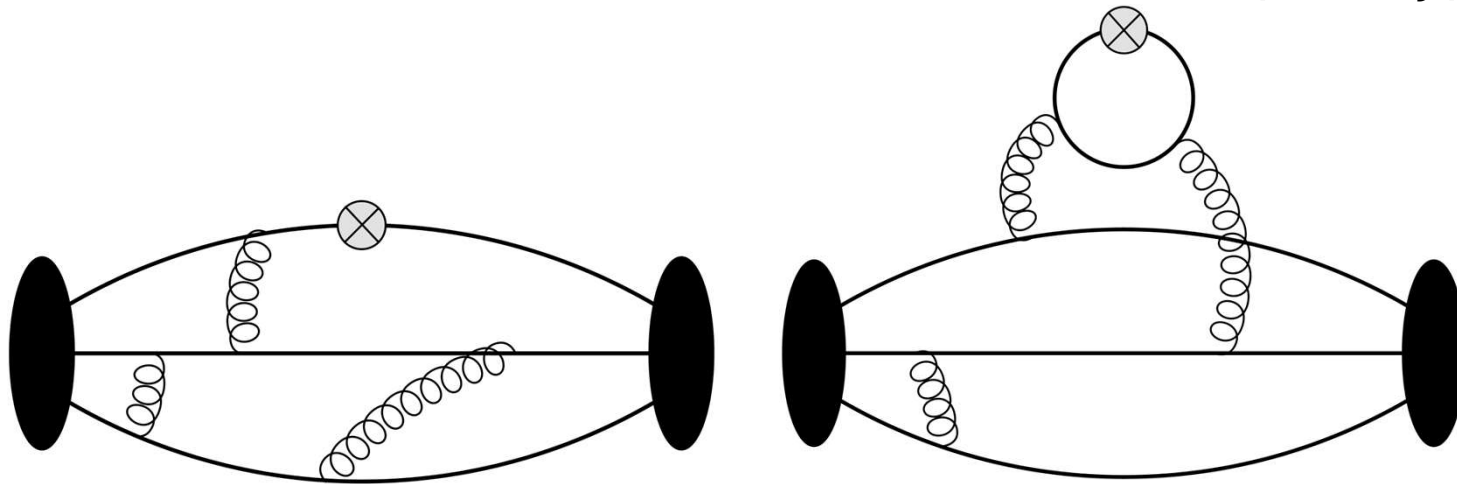
T. Bhattacharya, R. Gupta, S. Mondal, B. Yoon, H-W Lin, R. Zhang

LANL, MSU



Flavor diagonal matrix elements require high precision measurements of quark bilinear operators within the nucleon state for both “connected” and “disconnected” 3-point correlation functions,

Nucleon charges  $g_A$ ,  $g_T$ , and  $g_S$  obtained from ME  $\langle N | \bar{q}_i \Gamma q_j | N \rangle$



**Connected**

**Disconnected**

# Outline of talk

- Methodology for calculation of 3-point functions very well established
  - Signal in connected versus disconnected
- Removing excited state contributions (ESC)
  - Fits using the spectral decomposition to connected plus disconnected contributions (for a scalar charge  $g_S^{u+d}$ )
  - What states contribute? Do  $N\pi$  /  $N\pi\pi$  states contribute?
- Renormalization
  - Constructing the full mixing matrix
- Chiral-Continuum-Finite-Volume extrapolations
  - Chiral fits to  $\sigma_{\pi N}$ , the nucleon sigma term

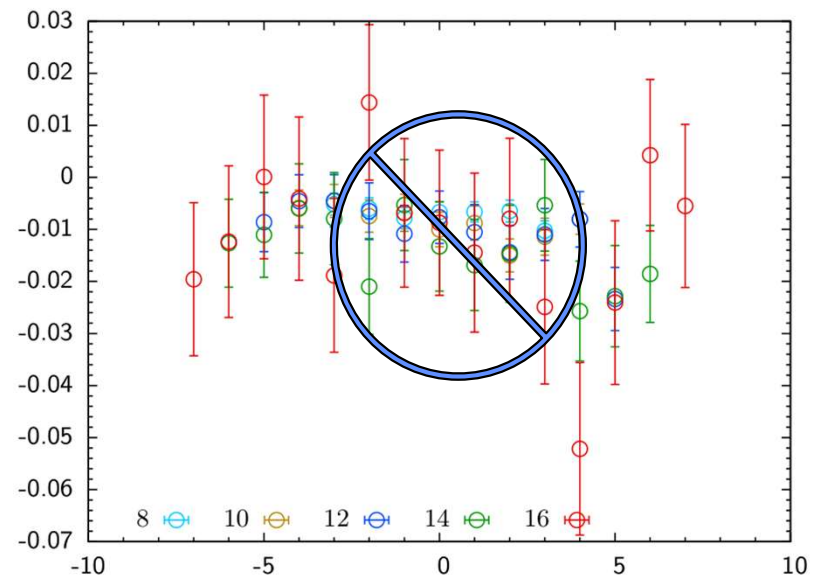
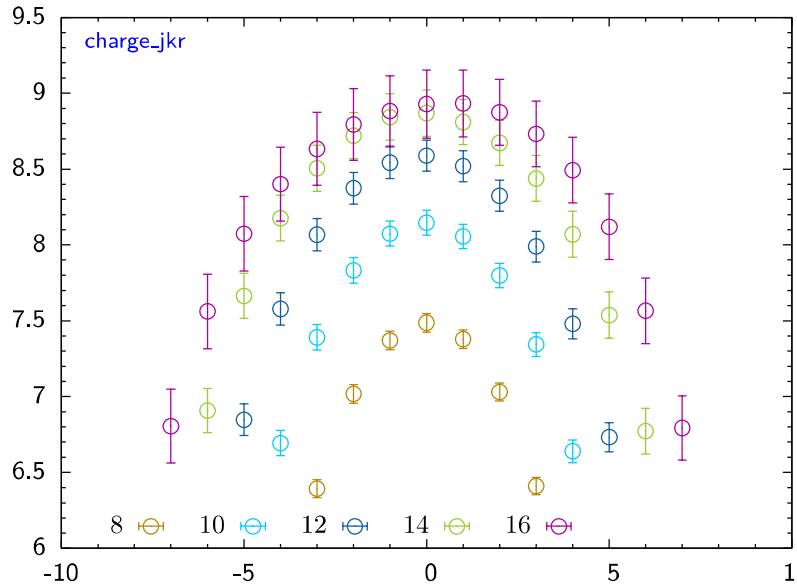
# All discussion based on Clover-on-HISQ data

Physical  $M_\pi$  Ensemble:  $a \approx 0.09 \text{ fm}$ ,  $M_\pi = 135 \text{ MeV}$ ,  $M_\pi L = 3.9$

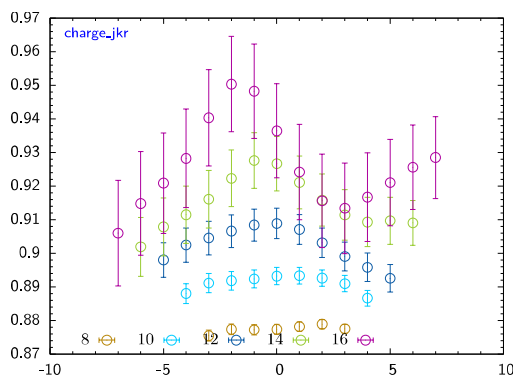
- Source-sink time separation  $\tau = 8, 10, 12, 14, 16$  (0.7~1.4fm)
- 1290/1270 configs
- 128 measurements of connected per config
- 10,000 sources for disconnected loop per config

# When are n-state ESC fits reliable?

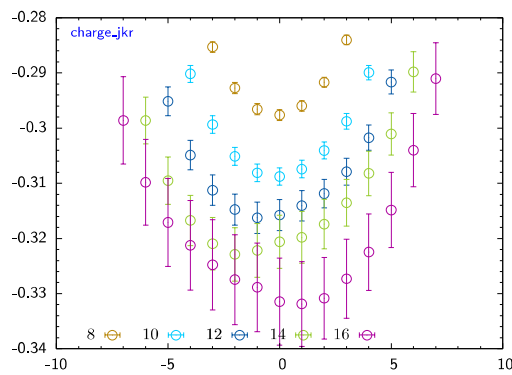
- Spectral decomposition of 3-point function tells us
  - Data for a given  $\tau$  should be symmetric about  $\tau/2$
  - Convergence should be monotonic for large enough  $\tau$  especially when only “one” excited state contribution is left
  - Only positive parity intermediate states contribute in the fits



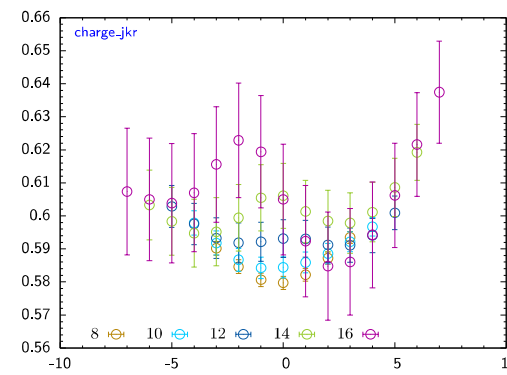
$$g_A^{u,conn}$$



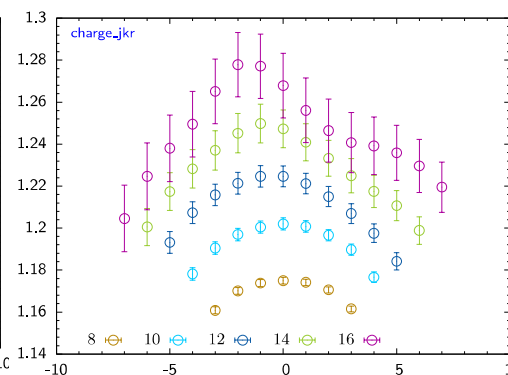
$$g_A^{d,conn}$$



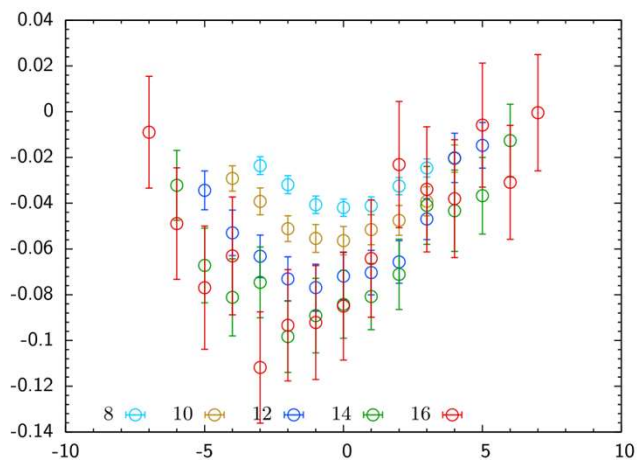
$$g_A^{u+d,conn}$$



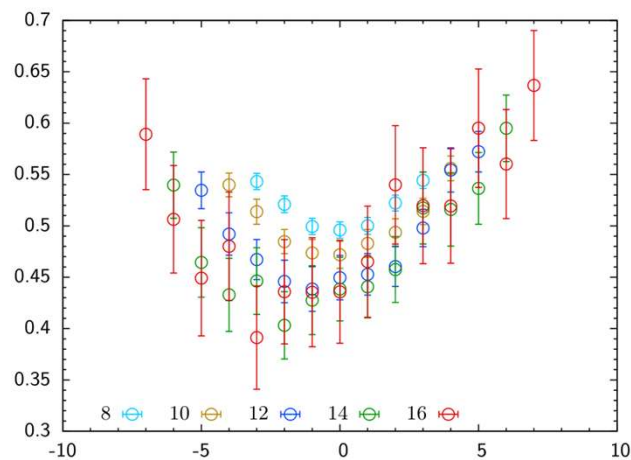
$$g_A^{u-d}$$



$$g_A^{l,disc}$$

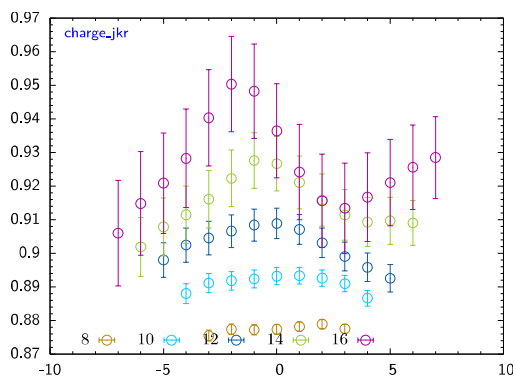


$$g_A^{u+d} = g_A^{u+d,conn} + 2g_A^{l,disc}$$

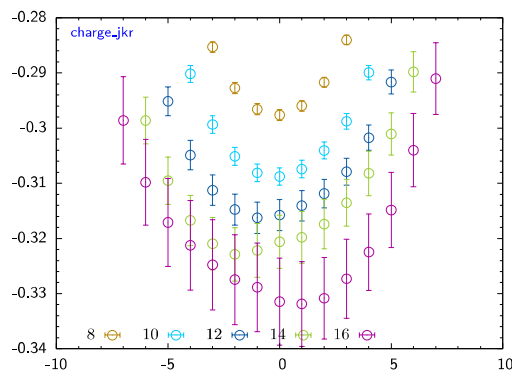


Also see “anatomy of ESC” in appendix D in arXiv:2103.05599

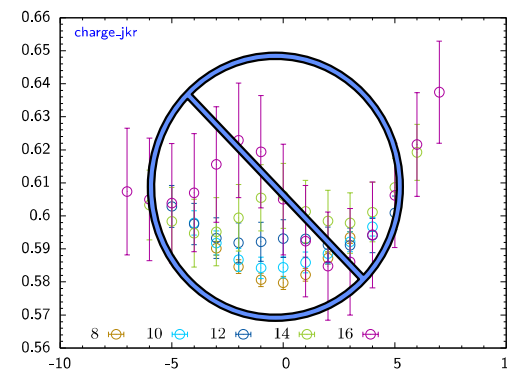
$$g_A^{u,conn}$$



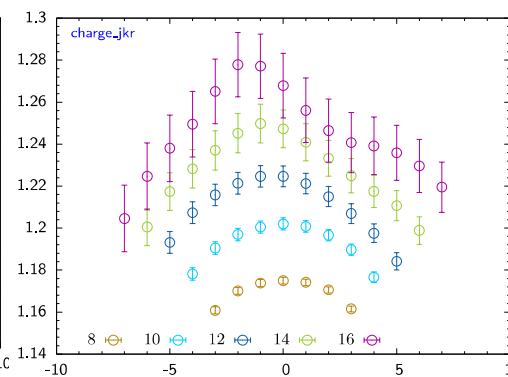
$$g_A^{d,conn}$$



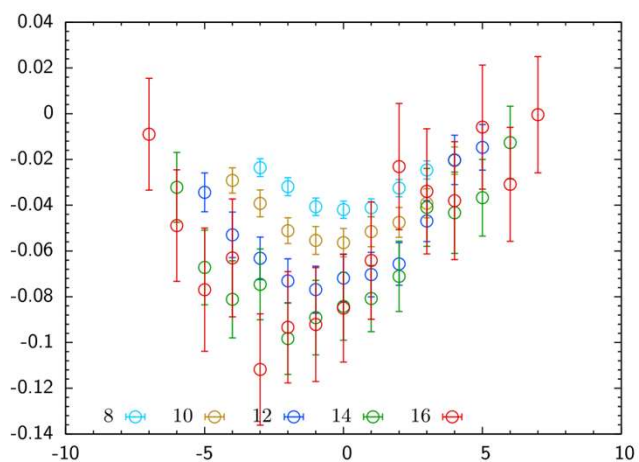
$$g_A^{u+d,conn}$$



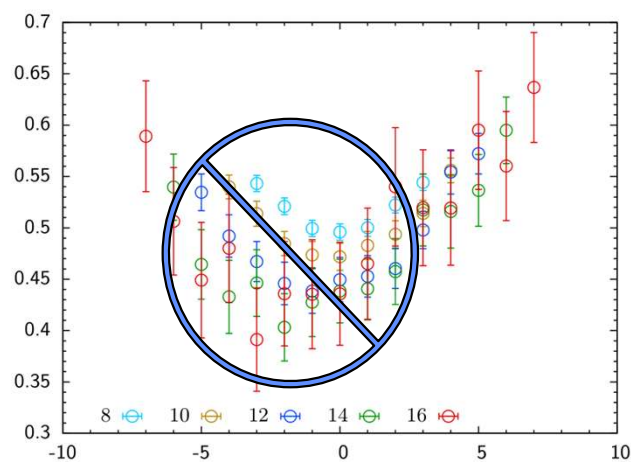
$$g_A^{u-d}$$



$$g_A^{l,disc}$$

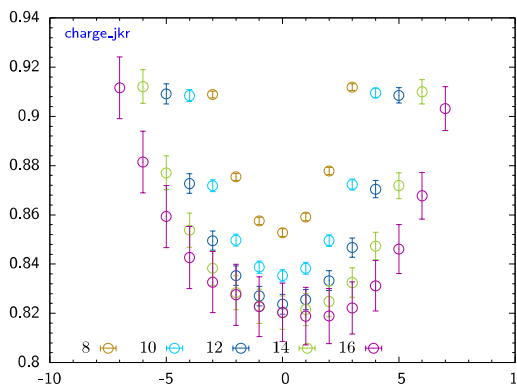
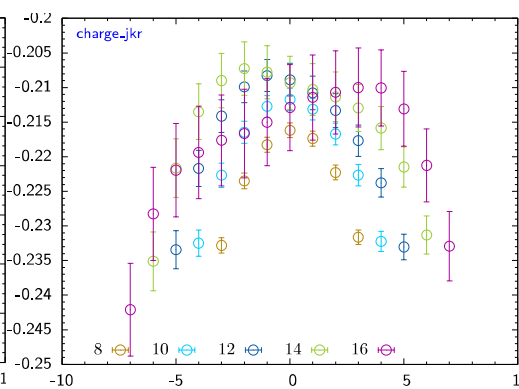
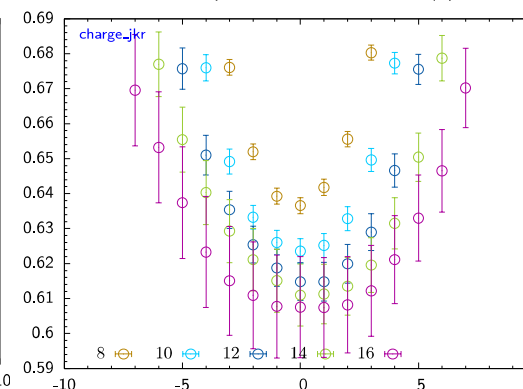
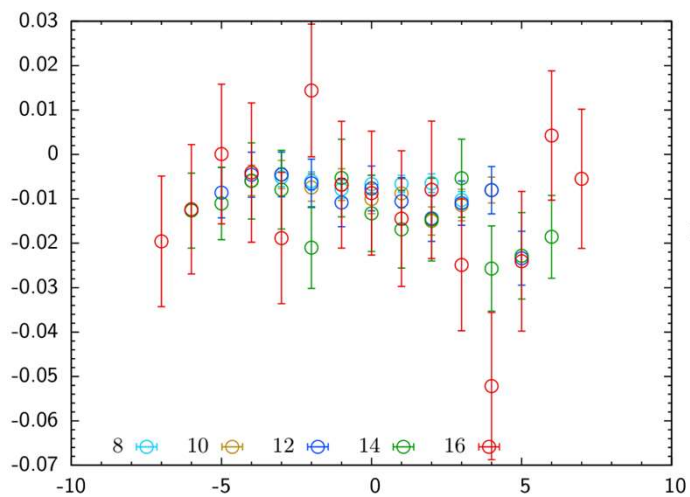
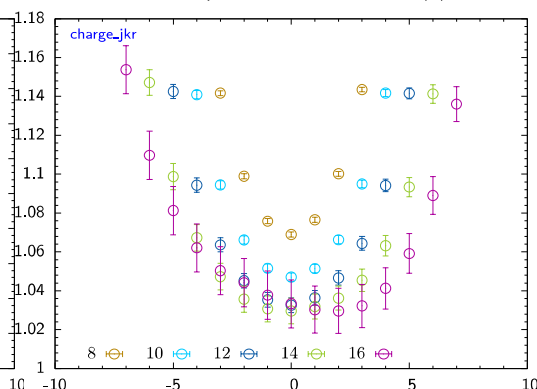
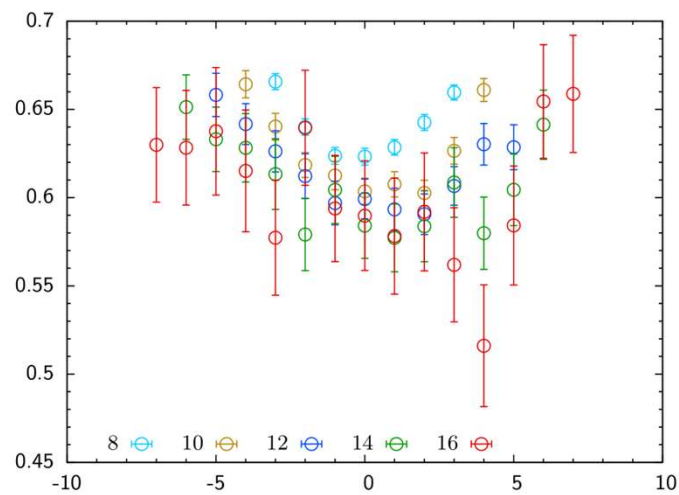


$$g_A^{u+d} = g_A^{u+d,conn} + 2g_A^{l,disc}$$

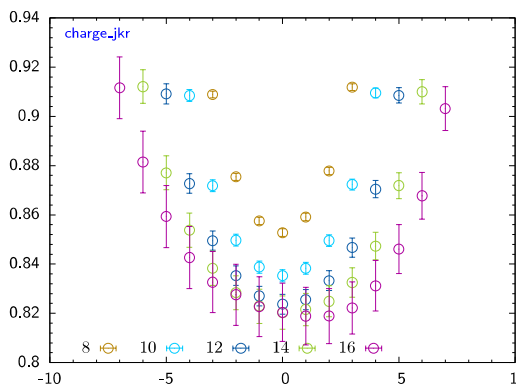
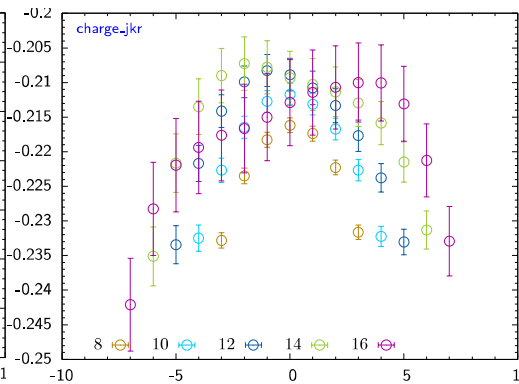
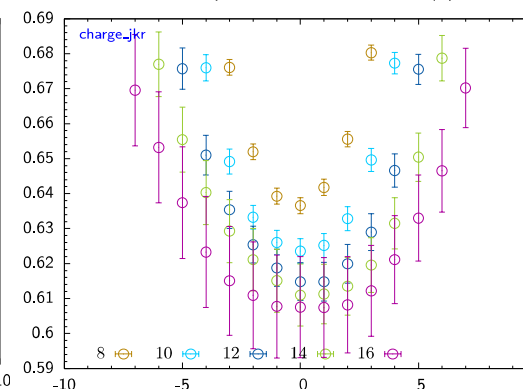
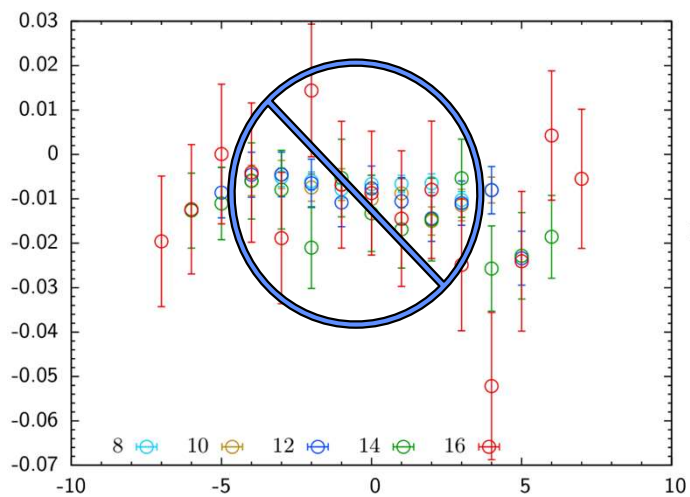
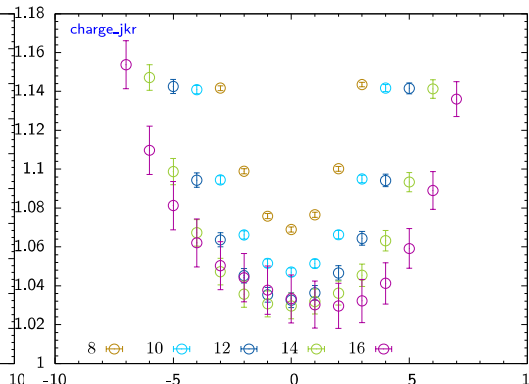
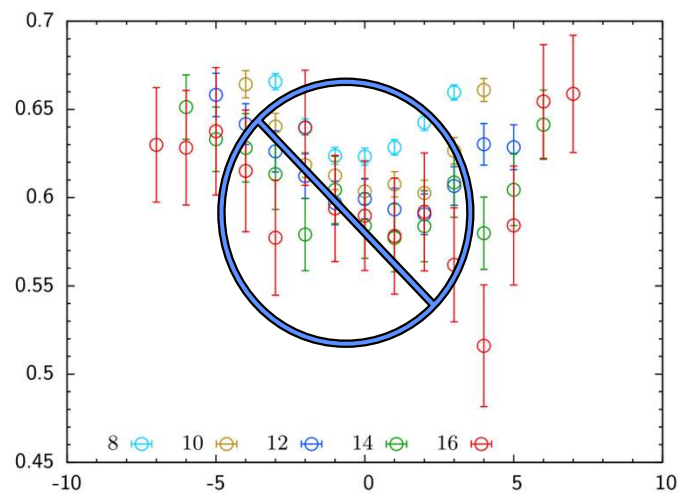


## $g_A$

- Signal in  $g_A^{u,conn}$ ,  $g_A^{d,conn}$ ,  $g_A^{l,disc}$  versus  $\tau$  improves together
- ESC in  $g_A^{u-d}$  adds
  - ⇒ Excited-state fits are reasonable
  - ⇒ Open issue of contribution of  $N\pi$  state. (arXiv:2103.05599)
- ESC in  $g_A^{u+d,conn}$  subtracts
  - ⇒  $\tau$  dependence hard to resolve
- $g_A^{l,disc}$  converges to a more negative value  $\approx -0.1$ 
  - ⇒ Makes the quark contribution to the proton spin smaller

$g_T^{u,conn}$  $g_T^{d,conn}$  $g_T^{u+d,conn}$  $g_T^{u-d}$  $g_T^{l,disc}$ 

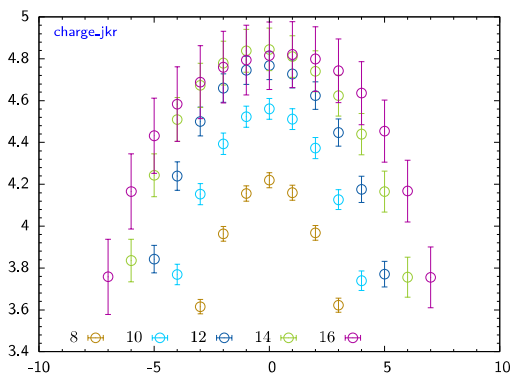
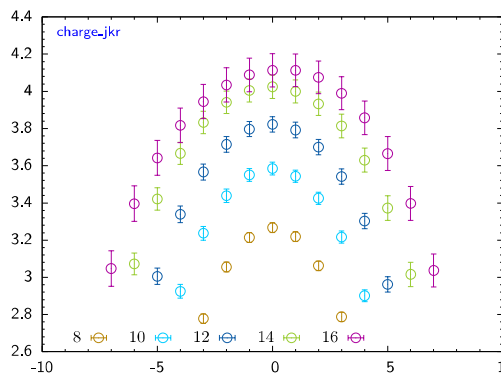
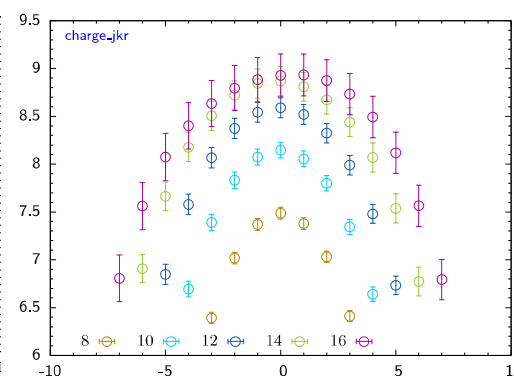
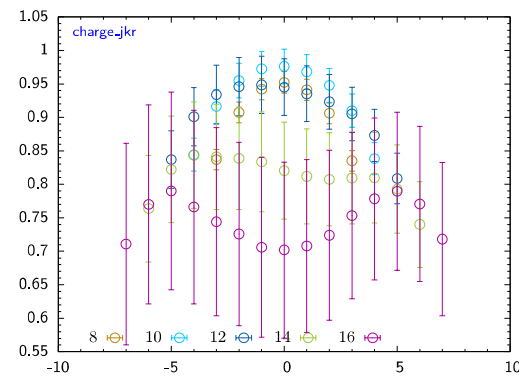
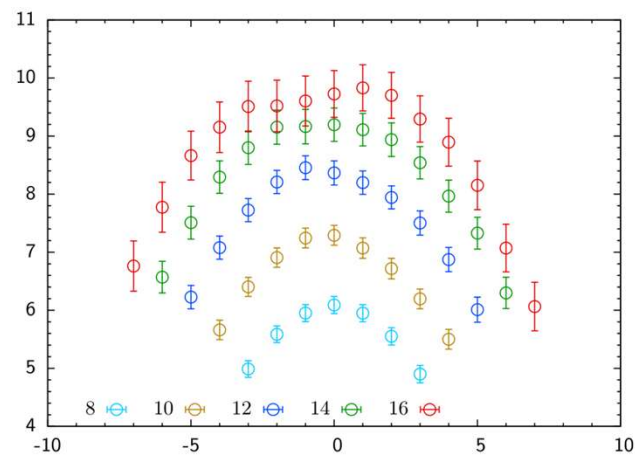
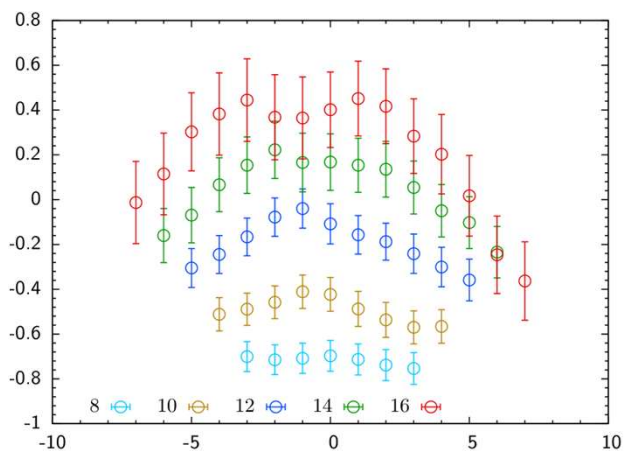
$$g_T^{u+d} = g_T^{u+d,conn} + 2g_T^{l,disc}$$

$g_T^{u,conn}$  $g_T^{d,conn}$  $g_T^{u+d,conn}$  $g_T^{u-d}$  $g_T^{l,disc}$ 

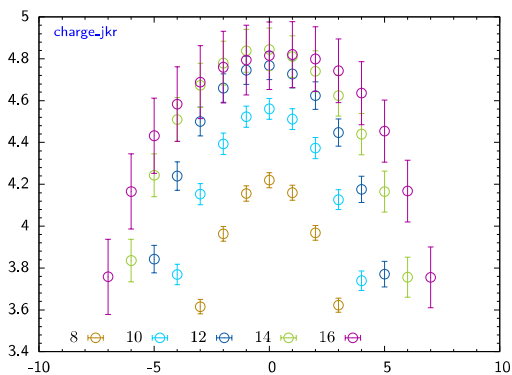
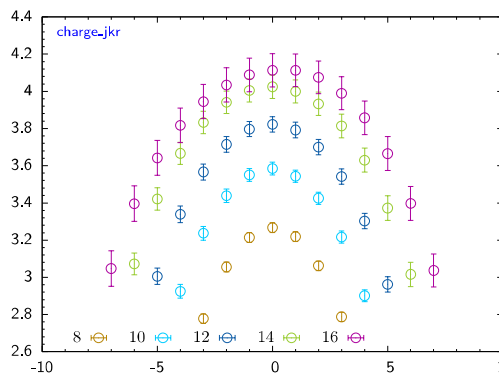
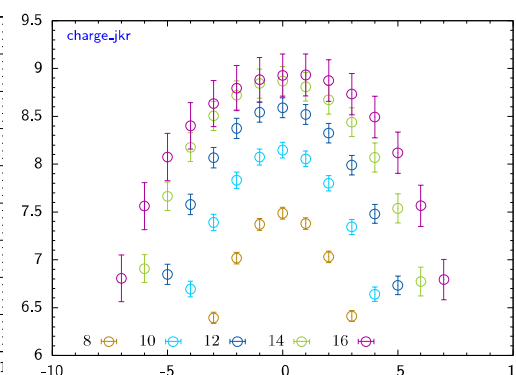
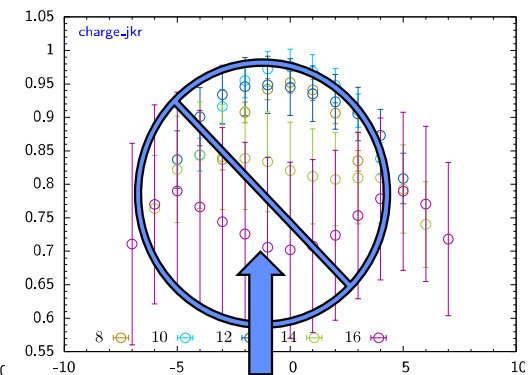
$$g_T^{u+d} = g_T^{u+d,conn} + 2g_T^{l,disc}$$

## $g_T$

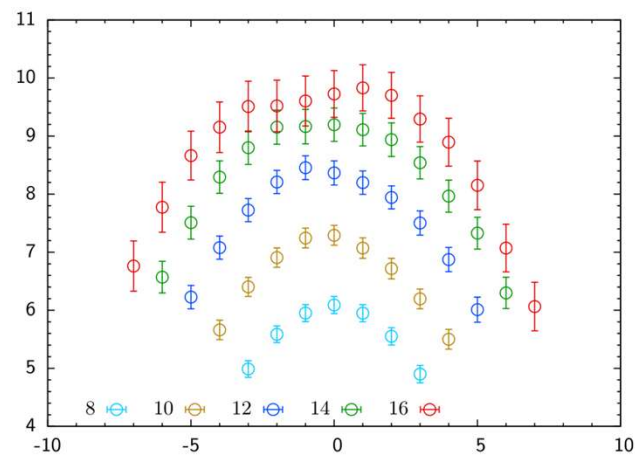
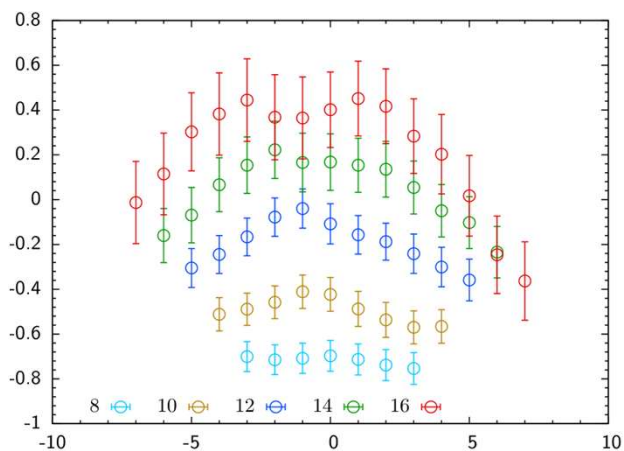
- ESC in  $g_T^{u,conn}$  larger than in  $g_T^{d,conn}$   
⇒ ESC  $g_T^{u,conn}$  dominates all combinations
- ESC in  $g_T^{u-d}$  and  $g_T^{u+d,conn}$  mainly driven by  $g_T^{u,conn}$   
⇒ Excited-state fits are reasonable
- $g_T^{l,disc} \sim -0.01$  and noisy  
⇒ A small contribution but makes  $g_T^{u+d}$  noisy

$g_S^{u,conn}$  $g_S^{d,conn}$  $g_S^{u+d,conn}$  $g_S^{u-d}$  $g_S^{l,disc}$ 

$$g_S^{u+d} = g_S^{u+d,conn} + 2g_S^{l,disc}$$

$g_S^{u,conn}$  $g_S^{d,conn}$  $g_S^{u+d,conn}$  $g_S^{u-d}$ 

Needs very  
high statistics

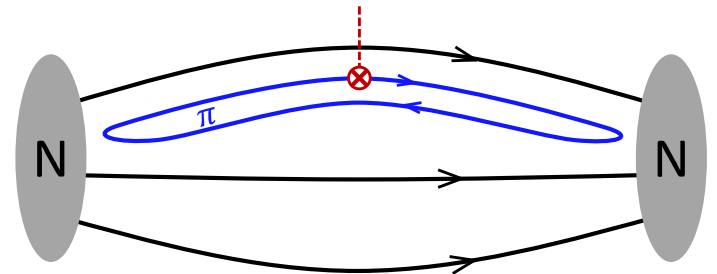
 $g_S^{l,disc}$ 

$$g_S^{u+d} = g_S^{u+d,conn} + 2g_S^{l,disc}$$

## $g_S$

- ESC in  $g_S^{u,conn}$  and  $g_S^{d,conn}$  are similar
- ESC in  $g_S^{u-d}$  cancels
  - ⇒ Excited-state fits need high statistics
- ESC in  $g_A^{u+d,conn}$  add
  - ⇒  $\tau$  dependence well-resolved
- $g_S^{l,disc}$  has a good signal
  - ⇒ Makes a big contribution

$\chi$ PT analysis  
indicates  $N\pi/N\pi\pi$   
states have  
significant coupling



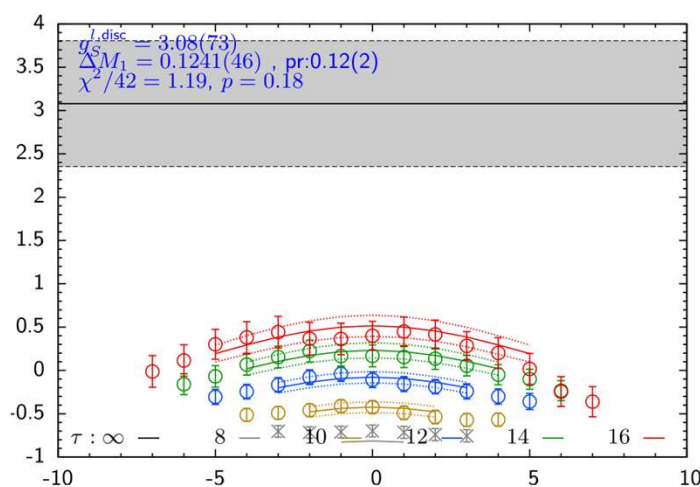
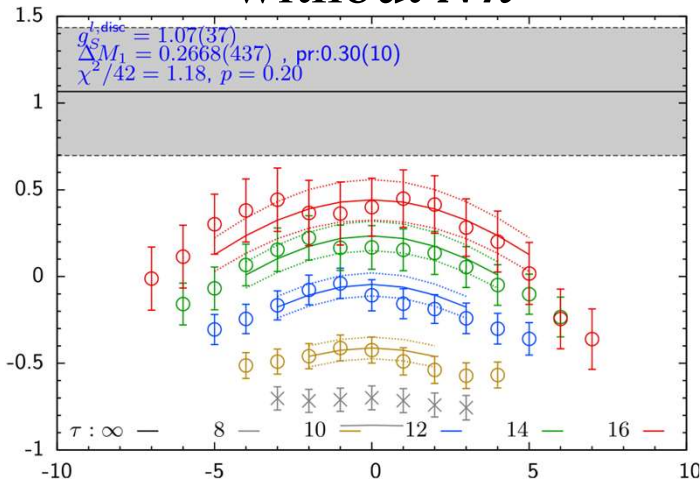
without  $N\pi$

with  $N\pi$

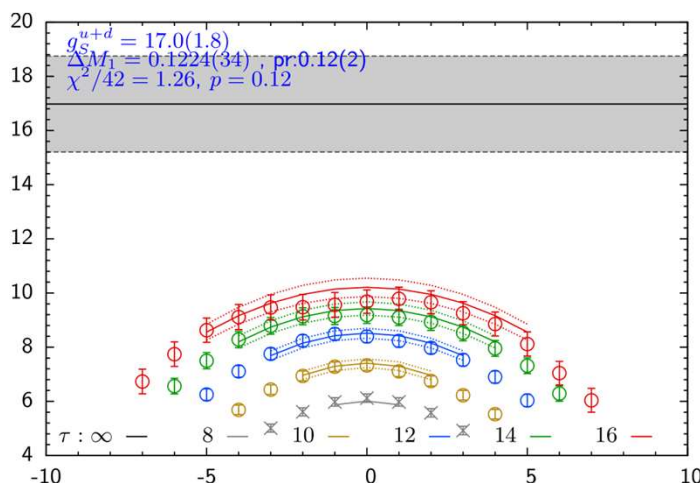
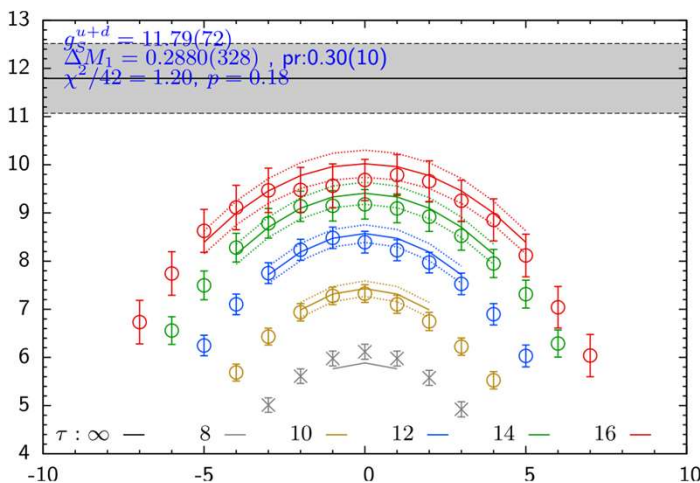
$a \approx 0.09 \text{ fm}$   
 $M_\pi \approx 135 \text{ MeV}$

$N(1)\pi(-1)$   
 or,  $N(0)\pi(0)\pi(0)$

$g_S^{l,\text{disc}}$



$g_S^{u+d} =$   
 $g_S^{u+d,\text{conn}}$   
 $+ 2g_S^{l,\text{disc}}$



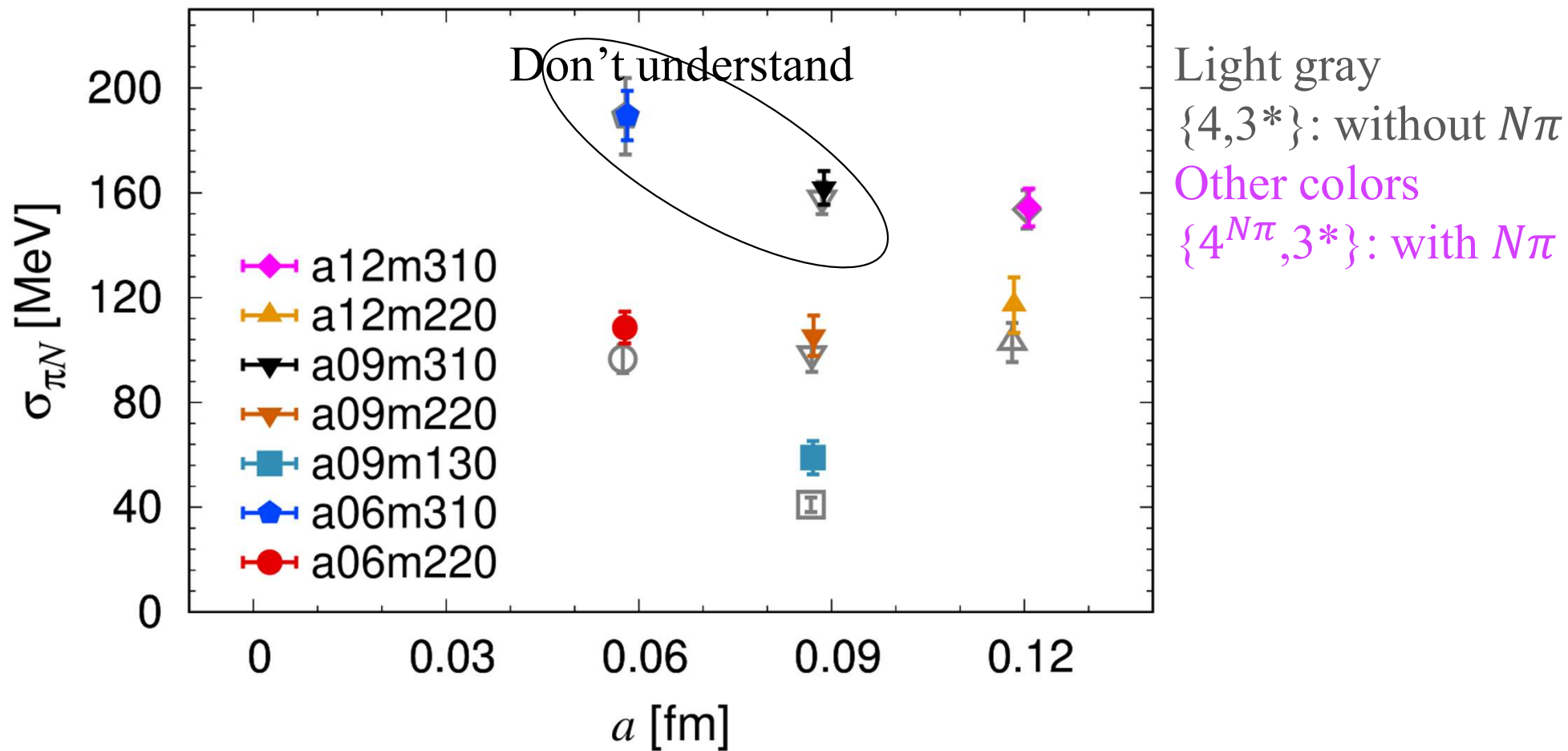
$\sigma_{\pi N} = m_l g_S^{u+d} \sim 40 \text{ MeV}$

$\sigma_{\pi N} = m_l g_S^{u+d} \sim 60 \text{ MeV}$

## $g_S$ and the nucleon sigma-term

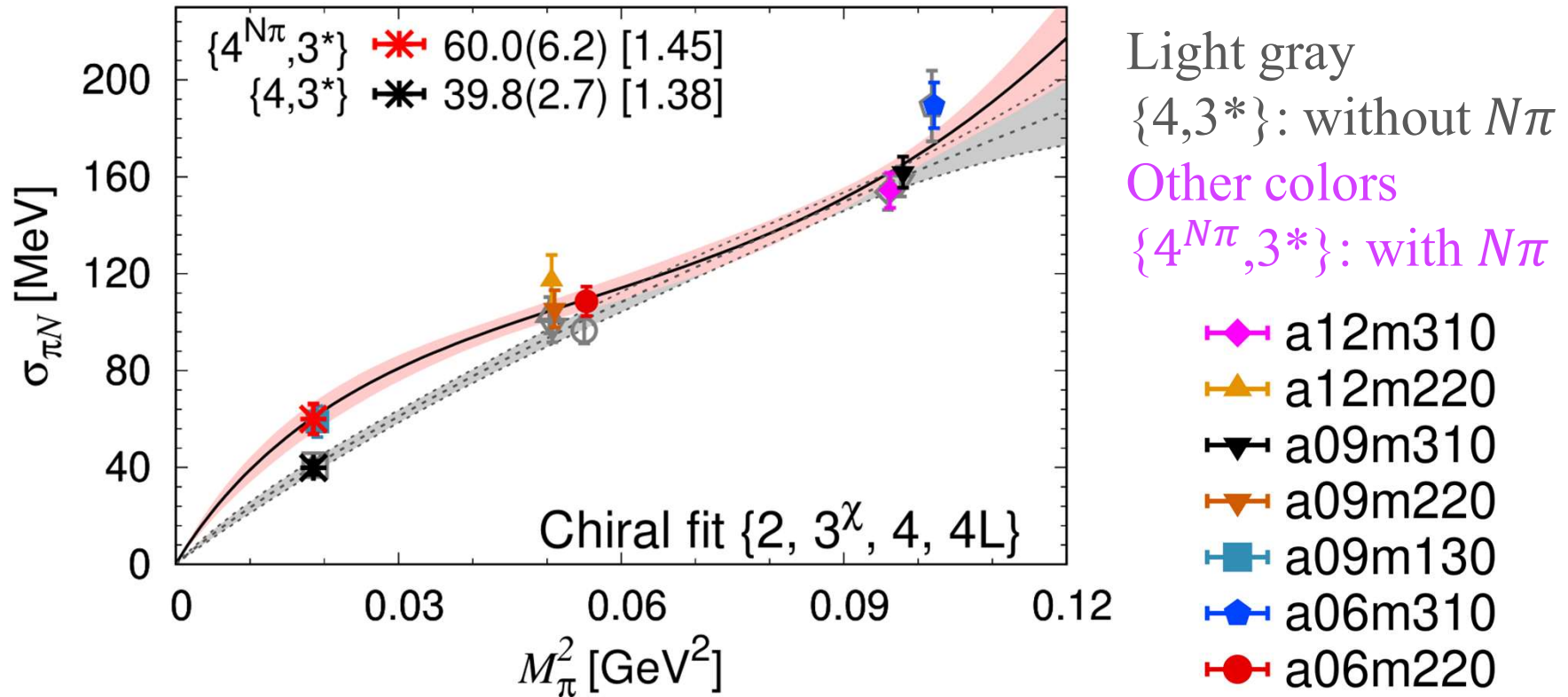
- Signal in  $g_S^{u+d,conn}$  and  $g_S^{l,disc}$  is good  
⇒  $\tau$  dependence well-resolved
- ESC in  $g_S^{u+d,total}$  from  $u$ ,  $d$ ,  $disc$  terms **adds**
- n-state fits to get the ground state ME are good
- ME (Nucleon  $\sigma$ -term) depends strongly on whether  $N\pi$  is included in excited-state fits
  - $\chi^2$  are equally good
- $\chi$ PT analysis implies significant contribution of  $N\pi$  and  $N\pi\pi$  state

*a*-dependence is not clear



# Chiral fits

arXiv:2105.12095



$$\sigma_{\pi N} = d_2 M_\pi^2 + d_3^\chi M_\pi^3 + d_4 M_\pi^4 + d_{4L} M_\pi^4 \log M_\pi^2$$

$d_3^\chi$ : fixed to  $\chi$ PT prediction

## Chiral fits

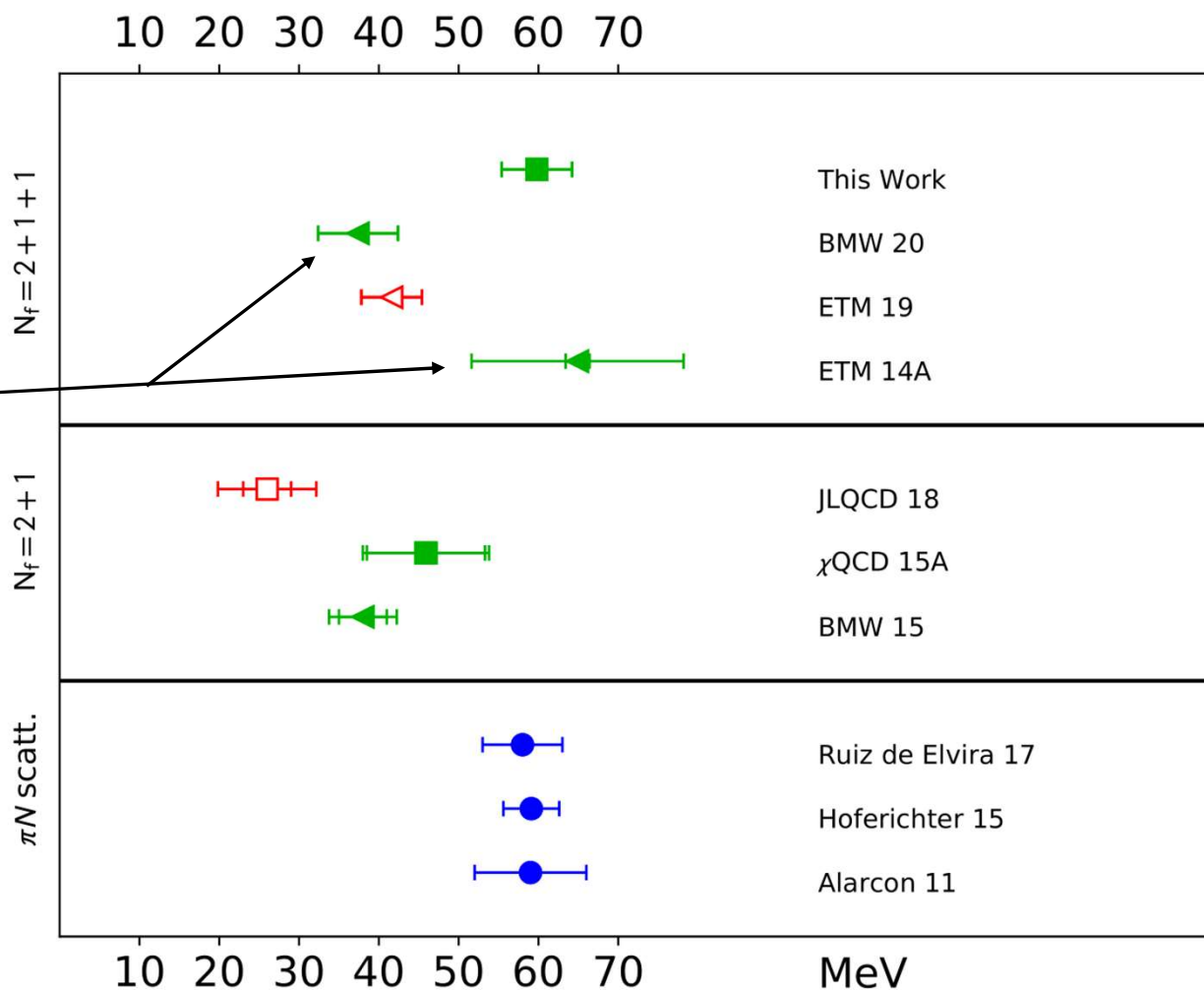
without  $N\pi$ with  $N\pi$ 

chiral fit	$\{4, 3^*\}_{\text{sim}}$						$\{4^{N\pi}, 3^*\}_{\text{sim}}$					
	$d_2$ (GeV $^{-1}$ )	$d_3$ (GeV $^{-2}$ )	$d_4$ (GeV $^{-3}$ )	$d_{4L}$ (GeV $^{-3}$ )	$\frac{\chi^2}{\text{dof}}$	$\sigma_{\pi N}$ (MeV)	$d_2$ (GeV $^{-1}$ )	$d_3$ (GeV $^{-2}$ )	$d_4$ (GeV $^{-3}$ )	$d_{4L}$ (GeV $^{-3}$ )	$\frac{\chi^2}{\text{dof}}$	$\sigma_{\pi N}$ (MeV)
$\chi\text{PT}$	4.44	-8.55	-	11.35	-	-	4.44	-8.55	-	11.35	-	-
$\{2, 3\}$	2.53(20)	-2.88(71)	-	-	1.12	39.0(2.0)	3.29(29)	-5.15(99)	-	-	2.23	47.3(2.9)
$\{2, 3, 4\}$	2.80(76)	-5.3(6.5)	5(13)	-	1.37	39.7(2.8)	6.6(1.5)	-32(12)	50(23)	-	1.56	59.6(6.3)
$\{2^x, 3^x, 4\}$	4.44	-8.55	-2.69(44)	-	18.8	58.96(15)	4.44	-8.55	-1.30(43)	-	4.91	59.42(14)
$\{2, 3^x, 4, 4L\}$	2.99(43)	-8.55	6(12)	-3.4(7.4)	1.38	39.8(2.7)	5.50(91)	-8.55	49(21)	27(13)	1.45	60.0(6.2)
$\{2, 3^x, 4, 4L^x\}$	3.83(12)	-8.55	30.1(1.4)	11.35	1.91	44.0(1.8)	4.44(17)	-8.55	24.1(1.9)	11.35	1.45	53.3(2.5)
$\{2^x, 3^x, 4, 4L\}$	4.44	-8.55	43.3(4.7)	20.2(2.1)	3.33	48.2(1.1)	4.44	-8.55	25.4(5.7)	11.9(2.5)	1.44	52.9(1.4)
$\{2^x, 3^x, 4, 4L^x\}$	4.44	-8.55	23.10(44)	11.35	5.87	52.91(15)	4.44	-8.55	24.13(43)	11.35	1.20	53.25(14)

$$\sigma_{\pi N} = d_2 M_\pi^2 + d_3 M_\pi^3 + d_4 M_\pi^4 + d_{4L} M_\pi^4 \log M_\pi^2$$

Coefficients with  $N\pi$  “consistent” with  $\chi\text{PT}$  predicted  $d_i$

$$\sigma_{\pi N}$$



Tension in  
F-H method.

Both with  
 $M_\pi \gtrsim 200$   
MeV data

## Summary

- Higher statistics needed to resolve  $\tau$  dependence in
  - $g_A^{l,disc}$  ,  $g_A^{u+d,conn}$
  - $g_T^{l,disc}$  ,  $g_T^{u+d,conn}$
  - $g_S^{u-d,conn}$
- Evidence of large contribution of  $N\pi$  &  $N\pi\pi$  in ES fits to  $g_S^{u+d}$   
Contribution becomes more significant as  $M_\pi \rightarrow 135$  MeV.
- $\sigma_{\pi N}$  changes from  $\sim 40$  MeV to  $\sim 60$  MeV on including the  $N\pi$  and  $N\pi\pi$  excited states [arXiv:2105.12095]
- Need more data to improve the continuum & finite volume extrapolation of  $\sigma_{\pi N}$