Confinement and Chiral Symmetry in the $SU(3)$ Instanton-dyon Ensemble

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Pure $SU(3)$: 2102.11321 [hep-ph]; Two-flavor QCD: in progress
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Outline

Instantons

- Instantons are 4D topological solitons related to tunneling between vacua of different N_{CS}
- Instantons generate chiral symmetry breaking, but not confinement; monopoles needed
- Instantons do not interact directly with the holonomy and are charge neutral

Instanton-dyons

- At $\langle P \rangle \neq 0$, the instanton is seen to be made of N_c constituents $\rightarrow$ the dyons
- Dyons have non-zero magnetic charge and non-integer topological charge
- Their action $S_i = \frac{8\pi^2}{g^2}\nu_i$ is ≈ 4 near T_c, making them suitable for using semiclassical methods
- Dyon action sets temperature scale $S_0 = \frac{8\pi^2}{g^2} = \frac{11}{3} N_c \ln(T/\Lambda)$

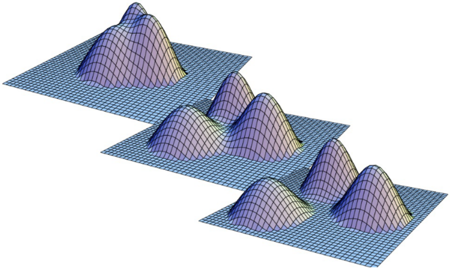
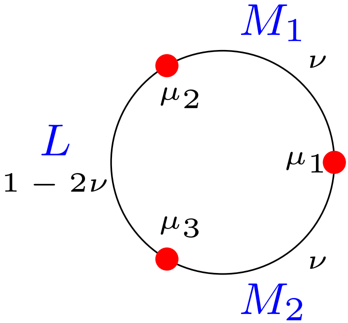


Figure: [Kraan, van Baal (1998)]

Polyakov Loop and Holonomy

- $\langle P \rangle = \frac{1}{3} + \frac{2}{3}\cos(2\pi\nu)$
- $\langle P \rangle = 0$, $\nu = \frac{1}{3} \rightarrow$ confined phase
- Dyons have individual actions $S_{M1} = S_{M2} = S_0\nu$, $S_L = S_0(1 - 2\nu)$
- Antidyons have the same properties with opposite magnetic and topological charges



Free Energy Density

$$
\begin{aligned}
f = {} & \frac{4\pi^2}{3}(2(\nu(1-\nu))^2 + (2\nu(1-2\nu))^2) \\
& - 4n_M \ln\left[\frac{d_\nu e}{n_M}\right] - 2n_L \ln\left[\frac{d_{1-2\nu} e}{n_L}\right] + \frac{\ln(8\pi^3 N_M^2 N_L)}{\tilde{V}_3} + \Delta f
\end{aligned} \tag{1}
$$

Compute Δf by Monte-Carlo integration

Minimize $f(T, \nu, n_M, n_L)$ to get $f(T)$, $\nu(T)$, $n_M(T)$, $n_L(T)$

Dyon Interactions

- $\Delta S^{d\bar{d}}_{class} = -\frac{S_0 C_{d\bar{d}}}{2\pi}(\frac{1}{rT} - 2.75\pi\sqrt{\nu_i\nu_j}e^{-1.408\pi\sqrt{\nu_i\nu_j}rT})$
- $C_{d\bar{d}} = 2$ for same kind, -1 for different kinds of dyons
- $\Delta S^{core}_{class} = \frac{\nu V_0}{1+e^{2\pi\nu T(r-r_0)}}$
- Core is used at distances smaller than $x_0 = 2\pi\nu rT$, core radius goes as $1/\nu$
- Diakonov Determinant:

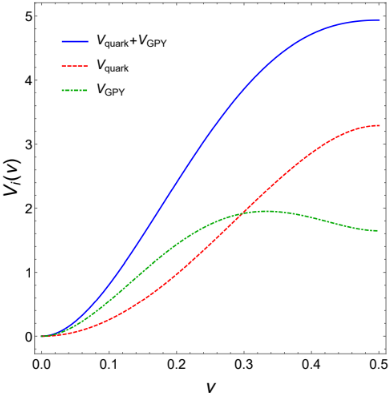
$$G_{im,jn} = \delta_{ij}\delta_{mn}(4\pi\nu_m - \sum_{k\neq i}\frac{2}{T|r_{i,m} - r_{k,m}|} + \sum_{k}\frac{1}{T|r_{i,m} - r_{k,p\neq m}|}) + \frac{2\delta_{mn}}{T|r_{i,m} - r_{j,n}|} - \frac{\delta_{m\neq n}}{T|r_{i,m} - r_{j,n}|}, \quad (2)$$

Dyons and Quarks

- We now include $N_f = 2$ flavors of dynamical massless quarks
- Dynamical $\rightarrow$ we include the fermionic determinant in the partition function
- In physical QCD the deconfinement transition is a smooth crossover, unlike the pure $SU(3)$ theory
- We can study the Dirac eigenvalue spectrum and thus the chiral condensate and chiral symmetry
- $\Sigma = |\langle \bar{q} q \rangle| = \lim_{V \to \infty} \lim_{\lambda \to 0} \pi \rho(\lambda)$

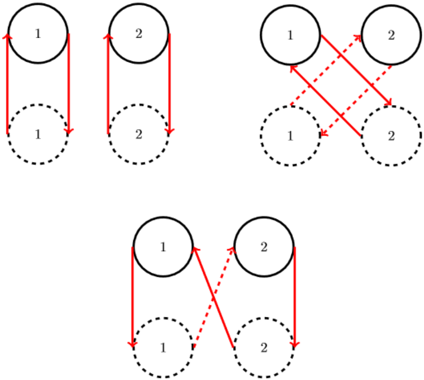
New Partition Function

- Two new terms in our partition function
- Quark perturbative potential $V_{quark} = -N_f \frac{4\pi^2}{3}(2\nu^4 - \nu^2)$
- Fermionic determinant $\det(i\not{D})^{N_f} \simeq \det(\hat{T})^{N_f}$
- Consider only the space spanned by the dyons' zero modes



Hopping Matrix

- Quarks can 'hop' from dyon to antidyon with amplitude T_{ij}
- $\hat{T} = \begin{pmatrix} 0 & T_{ij} \\ -T_{ji} & 0 \end{pmatrix}$
- $T_{ij} =$ $\bar{\nu} c' \exp(-\sqrt{11.2 + (\pi \bar{\nu} r T)^2})$
- $\det(\hat{T}) = T_{11}^2 T_{22}^2 + T_{12}^2 T_{21}^2 - 2 T_{11} T_{21} T_{22} T_{12}$



The Simulation Settings

- Compute free energy by standard integration over a dummy parameter in 10 steps $\lambda = 0.1, ...1$
- Perform $\mathcal{O}(50000)$ simulations with different inputs
- For each value of $S_0(T)$, fit to find minimizing input parameters
- Simulations run with $N_D = 120$ dyons in a 3D periodic box setup

	min.	max.	step size	no. of steps
S_0	8	21	1	14
ν	0.1	0.35	$0.01\bar{6}$	16
n_M	0.15	0.6	0.03	16
N_M/N_L	1.0	30.0	varied	16

The First-Order Transition

- Phase transition at $S_0 = 13.18$
- Jumps from $\langle P \rangle = 0$ to $\langle P \rangle \simeq 0.4$ ($\nu = 1/3$ to $\nu \simeq 0.23$)

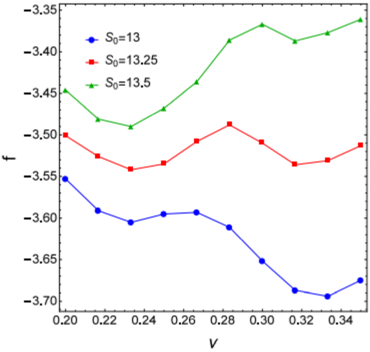


Figure: $N_f = 0$

Polyakov Loop and T_c

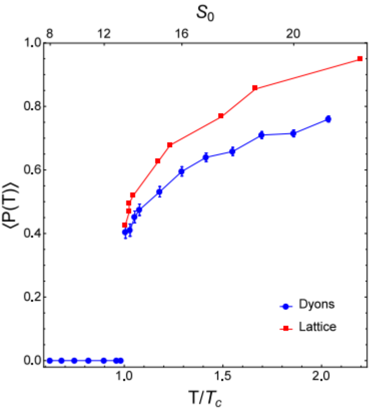


Figure: $N_f = 0$, Lattice data from Kaczmarek et. al. (2002)

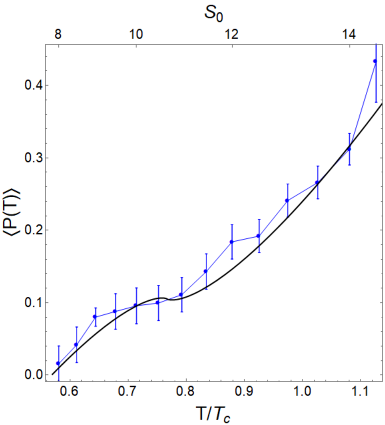


Figure: $N_f = 2$

Dyon Densities

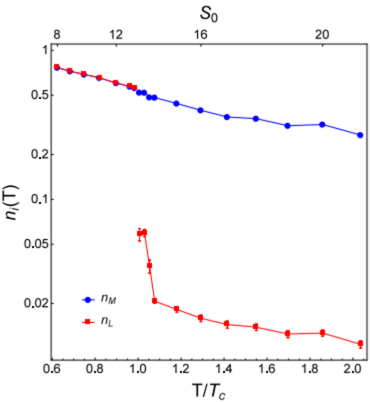


Figure: $N_f = 0$

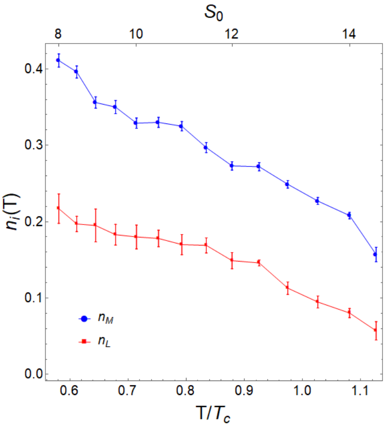


Figure: $N_f = 2$

Zero-Mode Zone

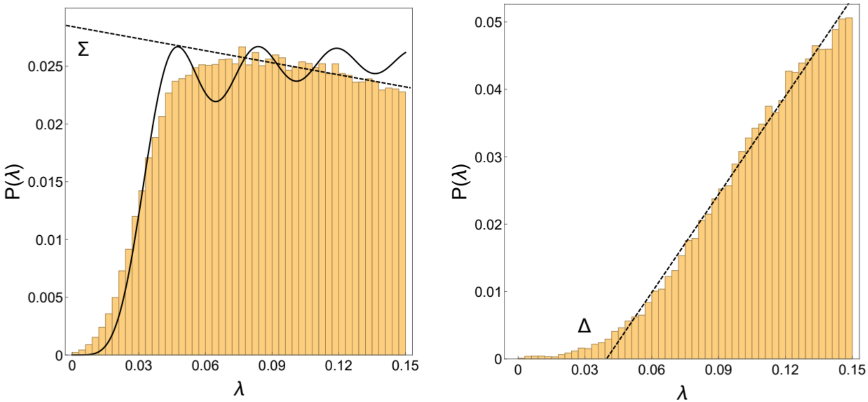
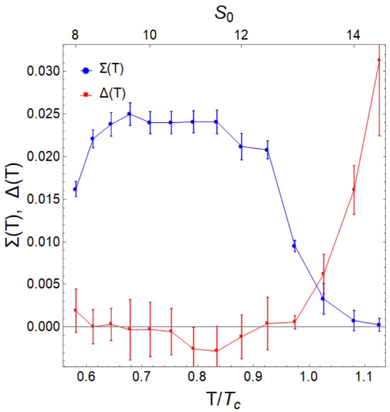


Figure: Left: $S_0 = 8$, Right: $S_0 = 14$

Low T: there is a finite density $\rho(0) \to$ broken chiral symmetry
High T: an eigenvalue gap Δ appears, restored chiral symmetry

Chiral Phase Transition



- The condensate quickly goes to zero at the same time a non-zero eigenvalue gap forms
- From this, we estimate $S_0(T_c) = 13.25$
- Should also belong to $O(4)$ universality class

Conclusions

- The dyon ensemble does possess a first-order deconfinement phase transition to the same value of $\langle P \rangle$ as the lattice for the pure $SU(3)$ theory and is compatible with a second-order transition when $N_f = 2$
- High dyon density leads to collectivization of zero modes and generates a nonzero chiral condensate at below T_c and finite eigenvalue gap above T_c
- Not discussed: Topological susceptibility, dyon correlations, trace deformed theory, nonzero quark mass