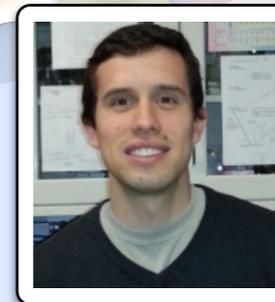


HYBRID ANALOG-DIGITAL QUANTUM SIMULATIONS FOR QUANTUM FIELD THEORIES

ZOHREH DAVOUDI

UNIVERSITY OF MARYLAND, COLLEGE PARK

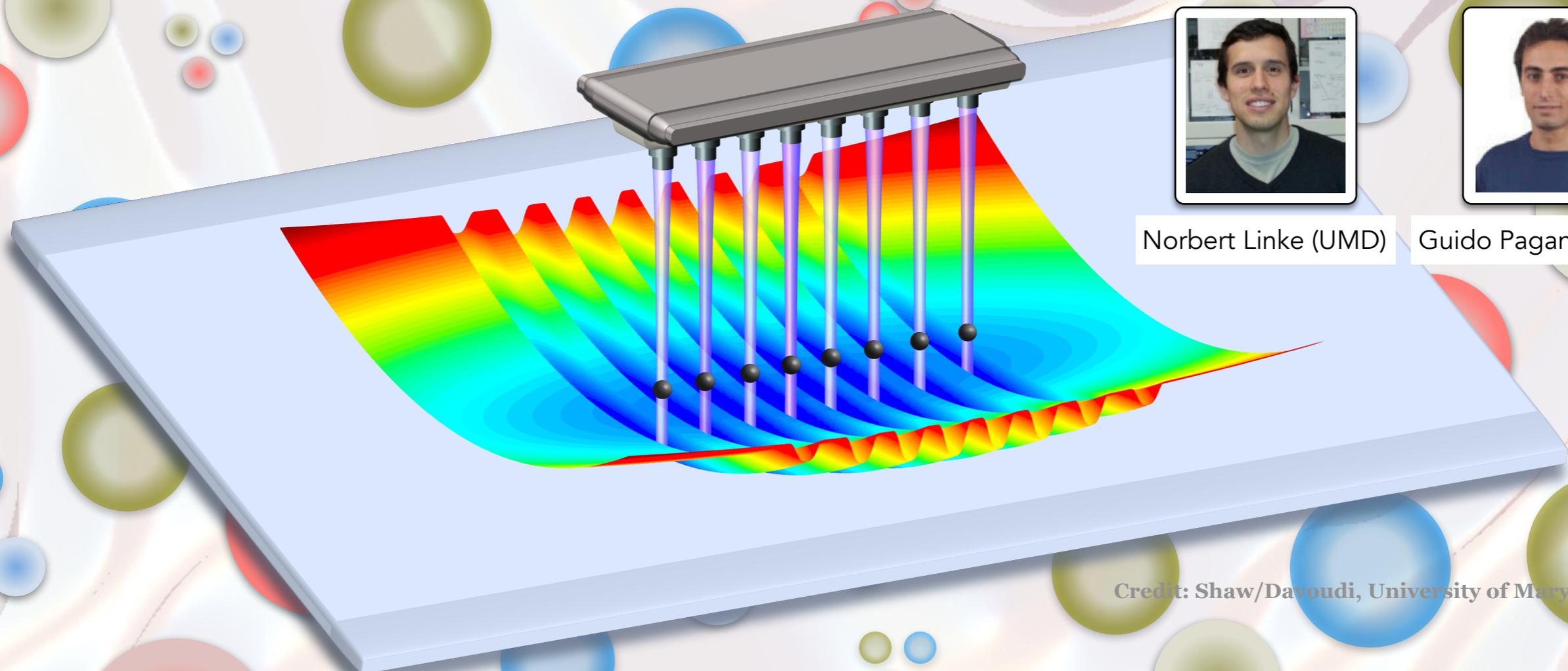
Work in collaboration with:



Norbert Linke (UMD)



Guido Pagano (Rice U)



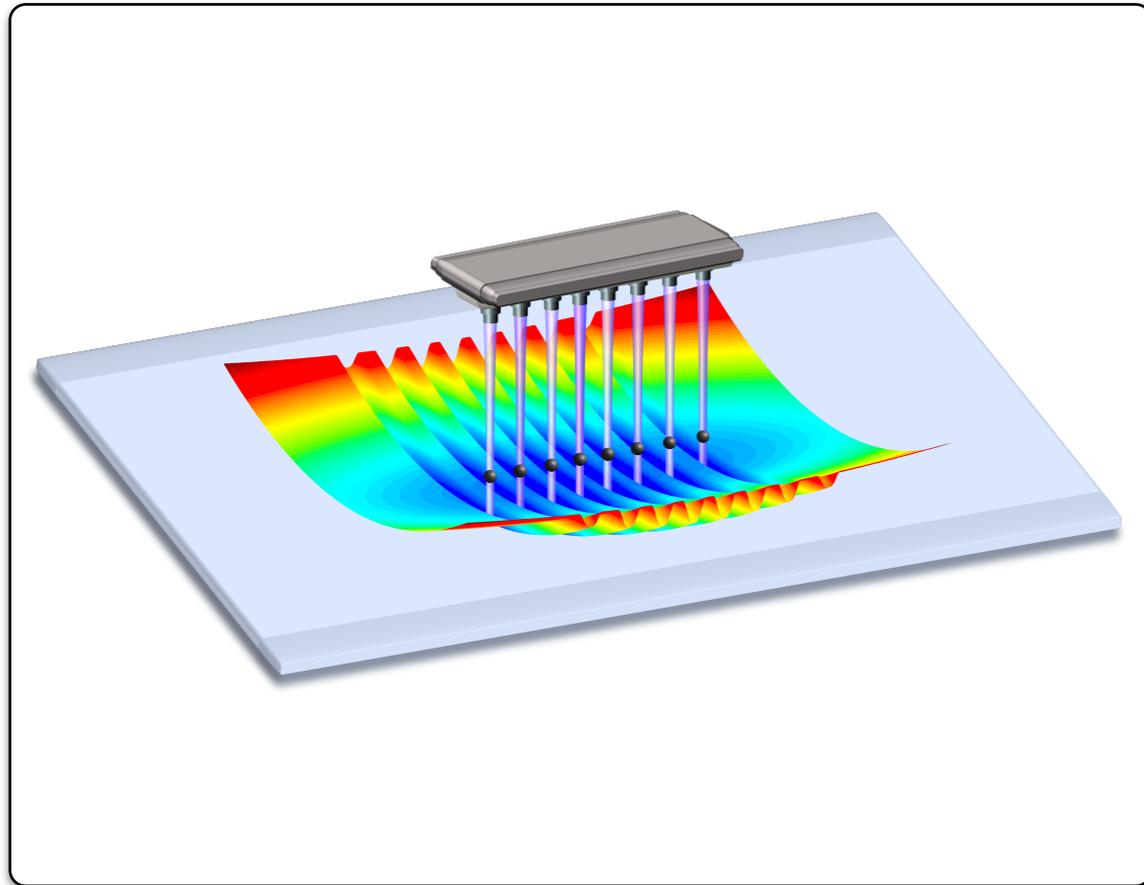
Credit: Shaw/Davoudi, University of Maryland

The 38th International Symposium on Lattice Field Theory
Zoom/Gather@MIT
July 26-30, 2021

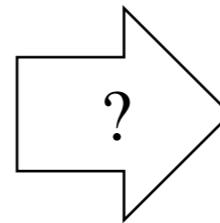
TO PUT THINGS IN
THE CONTEXT...

THE END GOAL FOR US: QUANTUM SIMULATION OF QUANTUM CHROMODYNAMICS!

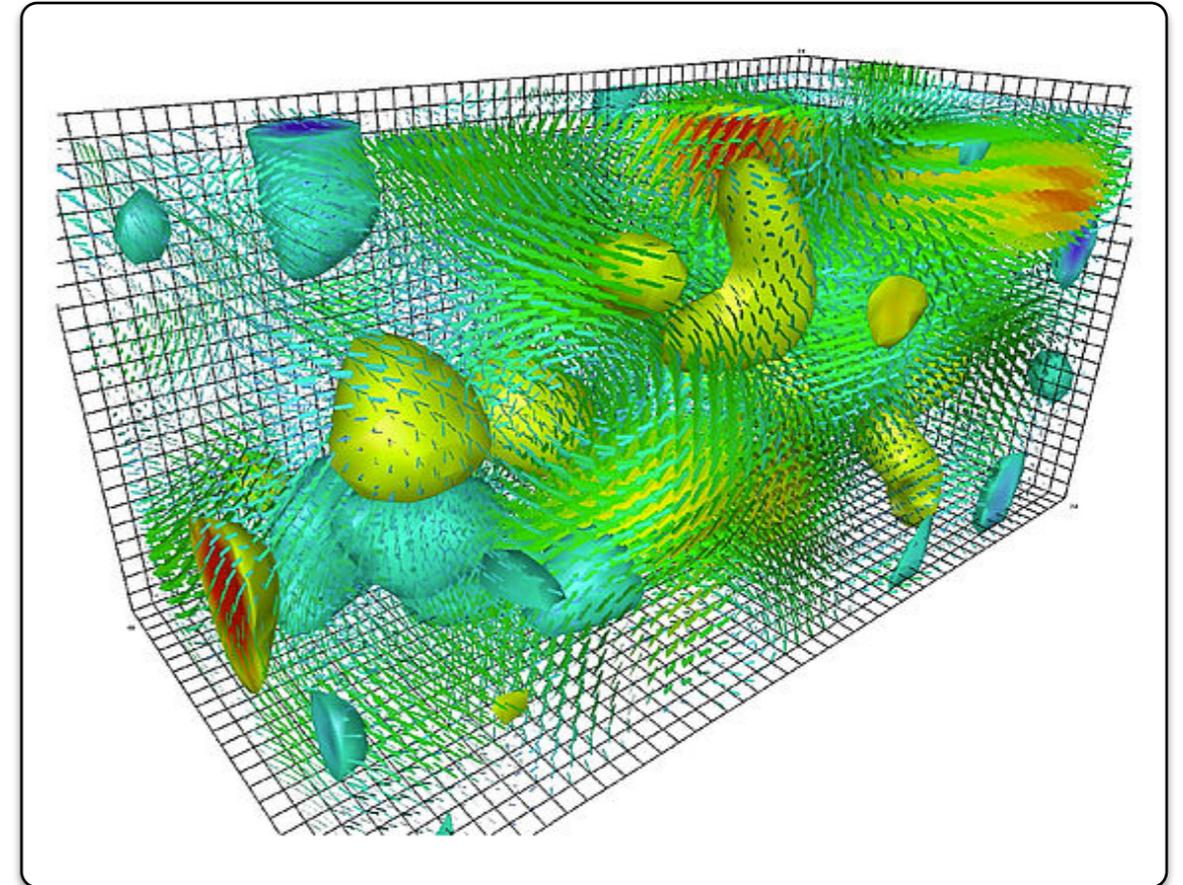
A controlled quantum system



CREDIT: ANDREW SHAW, UNIVERSITY OF MARYLAND



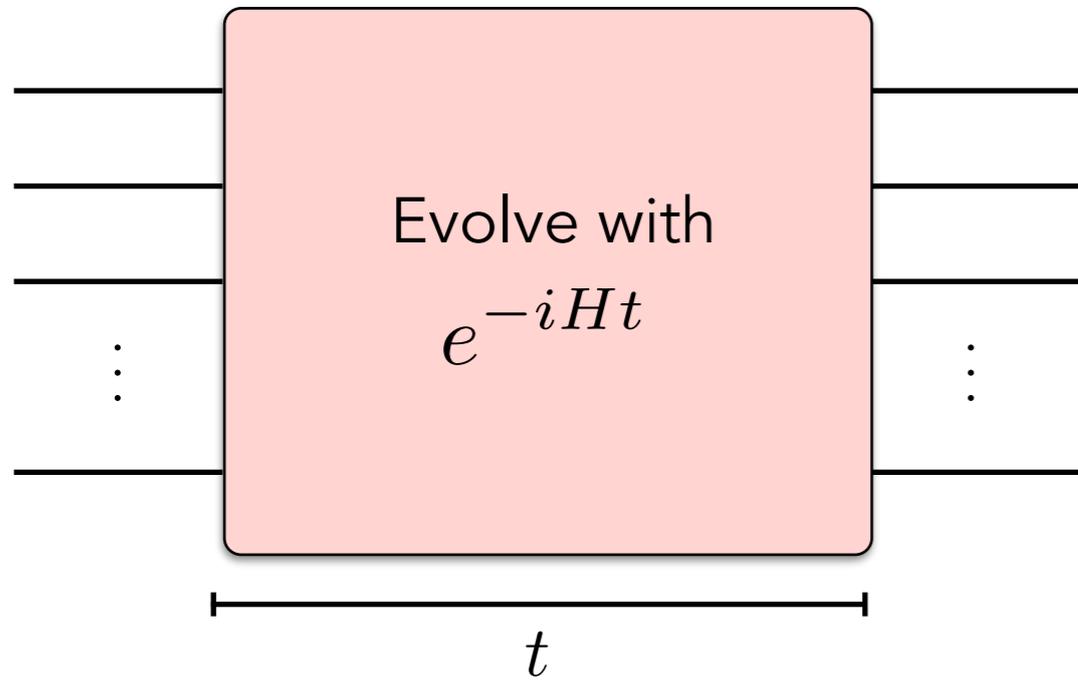
Strong-interaction physics



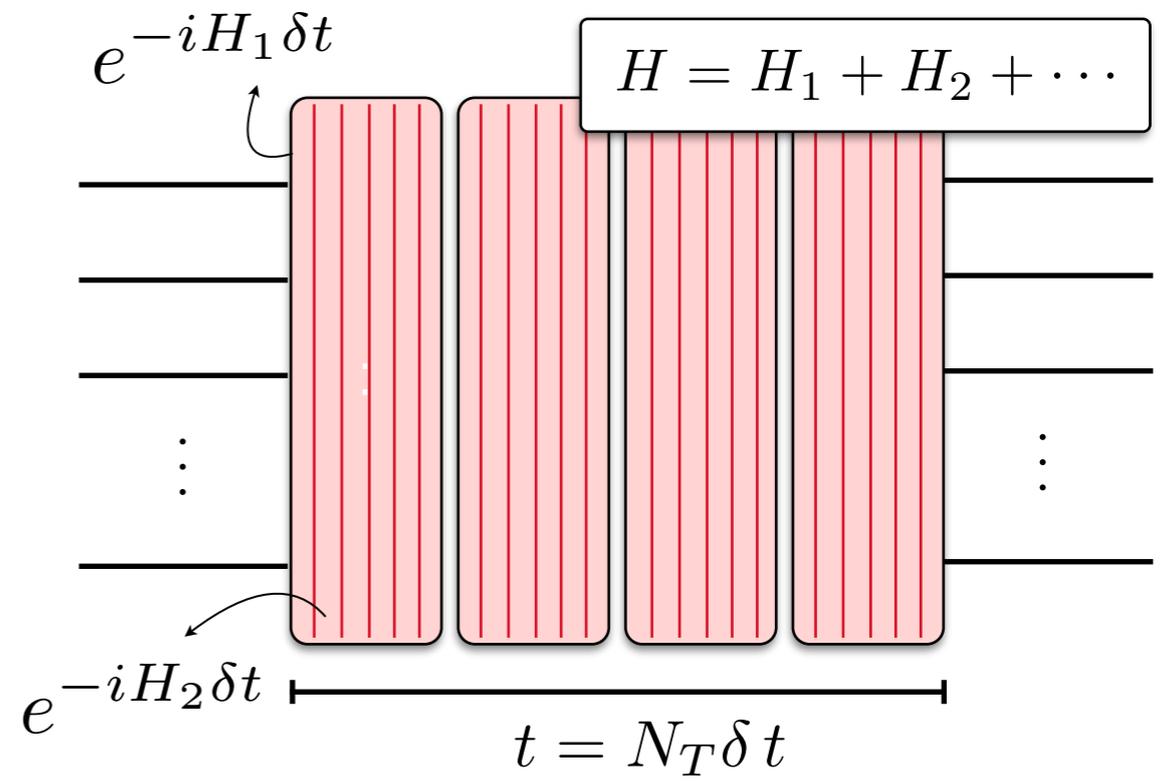
COPY RIGHT: UNIVERSITY OF ADELAIDE

DIFFERENT APPROACHES TO QUANTUM SIMULATION

Analog

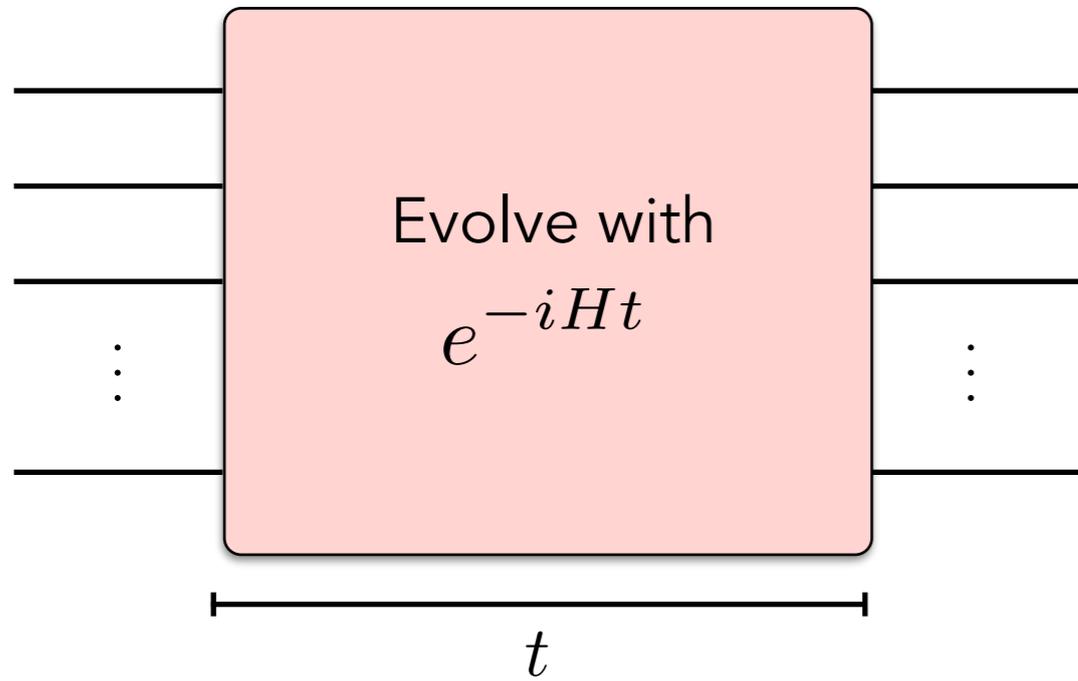


Digital

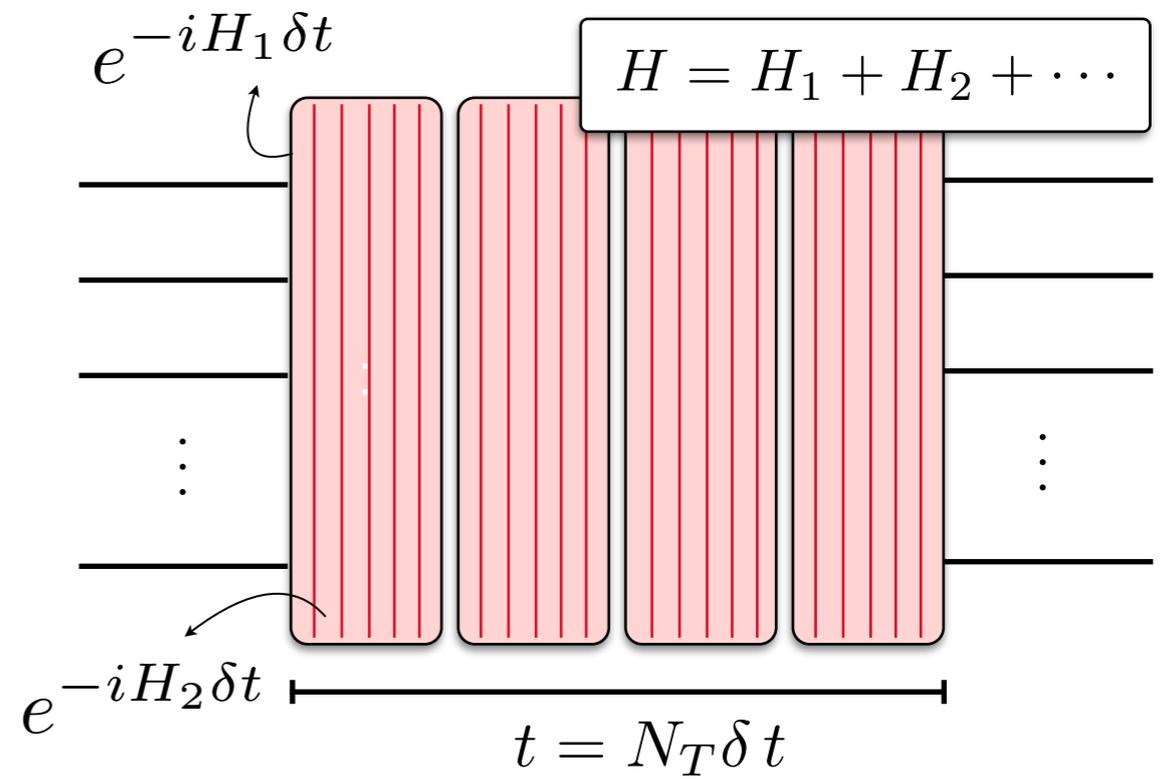


DIFFERENT APPROACHES TO QUANTUM SIMULATION

Analog



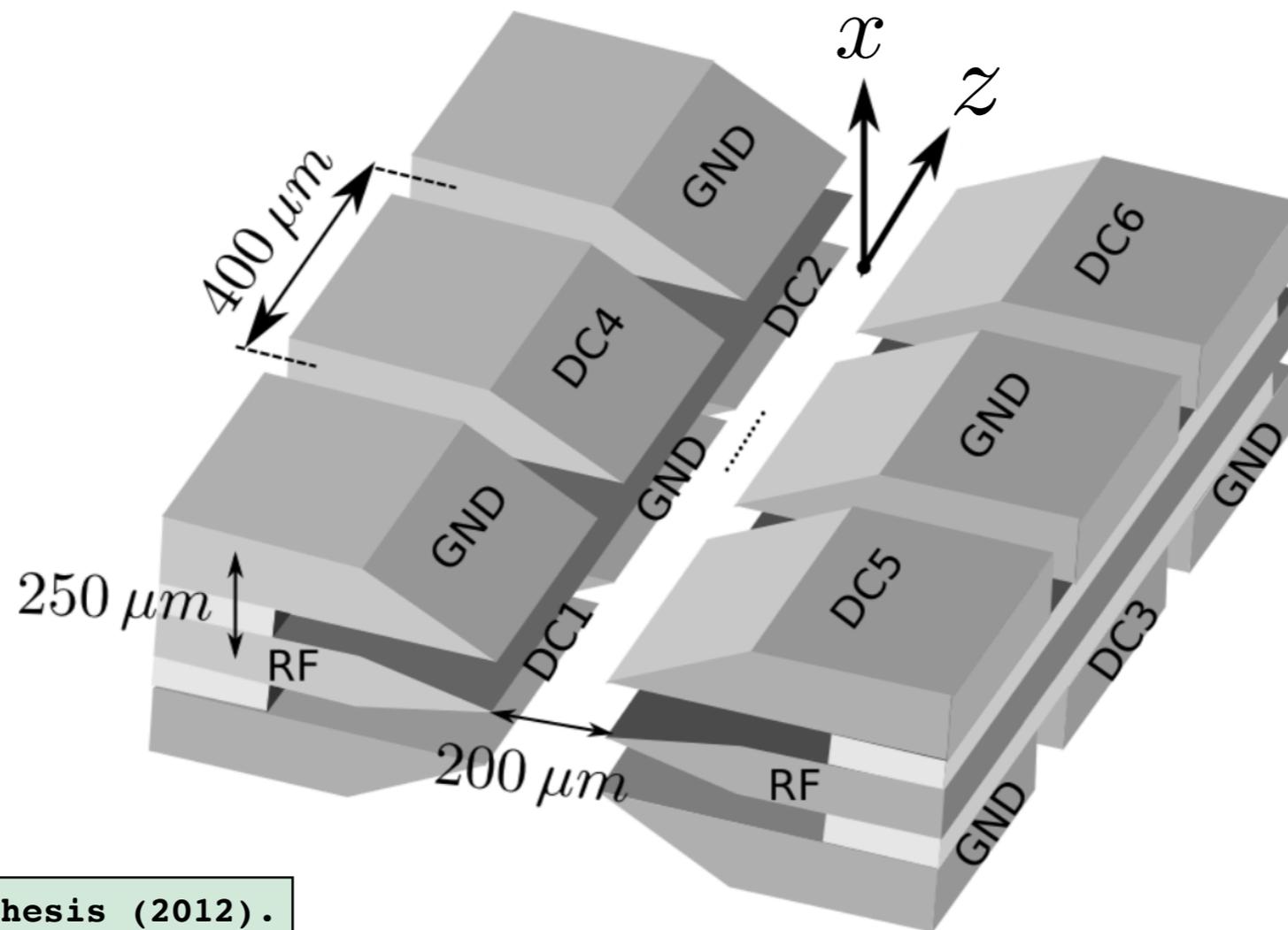
Digital



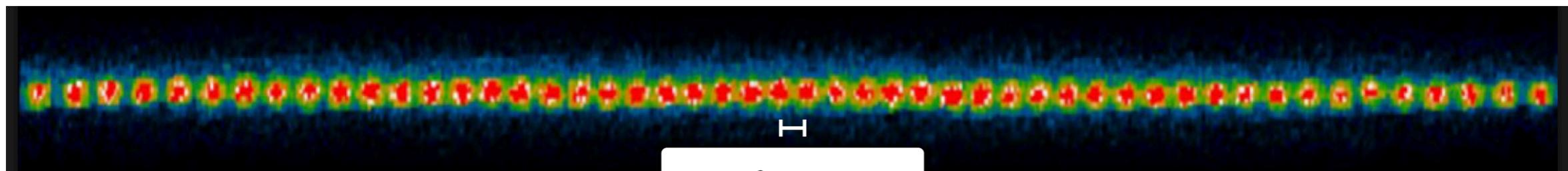
Analog-Digital

A QUICK TOUR TO THE
UNDERLYING PHYSICS OF
TRAPPED-ION SIMULATORS

A RADIO-FREQUENCY PAUL TRAP:



Islam, UMD PhD Thesis (2012).



\sim few μm

UMD, Monroe Lab.

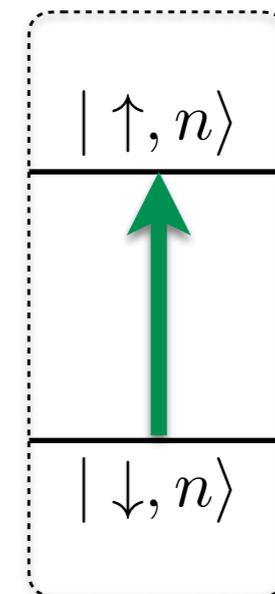
ION-LASER INTERACTIONS

Wineland et al, J.Res.Natl.Inst.Stand.Tech. 103 (1998) 259, Schneider et al, Rep. Prog. Phys. 75 024401 (2012)

$\mathcal{O}(\eta^0)$:

$$H_{\text{carr}} = -\frac{\Omega}{2} (\sigma^+ e^{-i\phi} + \sigma^- e^{i\phi})$$

One-qubit operations



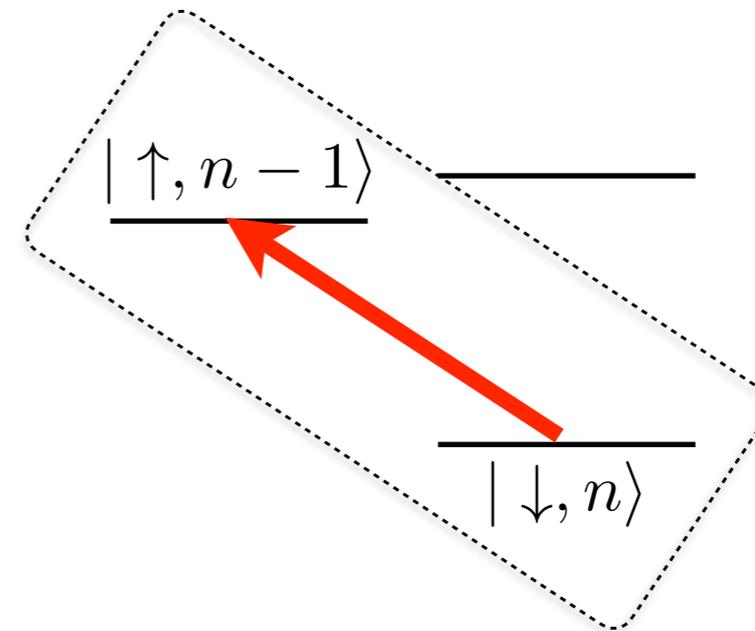
ION-LASER INTERACTIONS

Wineland et al, J.Res.Natl.Inst.Stand.Tech. 103 (1998) 259, Schneider et al, Rep. Prog. Phys. 75 024401 (2012)

$\mathcal{O}(\eta^1)$:

$$H_{rsb} \approx \frac{i}{2} \eta \Omega [a^\dagger \sigma^- e^{i\phi} - a \sigma^+ e^{-i\phi}]$$

Spin-phonon transitions



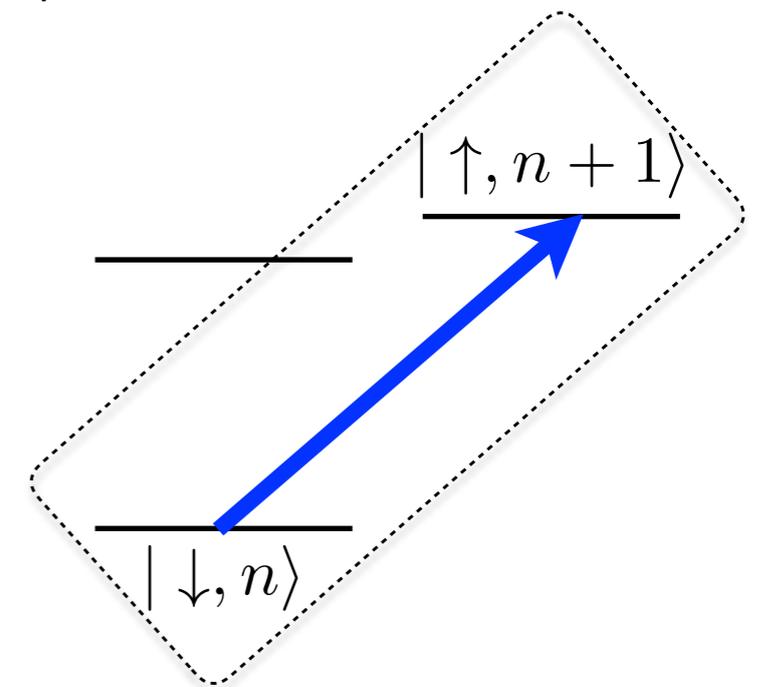
ION-LASER INTERACTIONS

Wineland et al, J.Res.Natl.Inst.Stand.Tech. 103 (1998) 259, Schneider et al, Rep. Prog. Phys. 75 024401 (2012)

$\mathcal{O}(\eta^1)$:

$$H_{bsb} \approx \frac{i}{2} \eta \Omega [a \sigma^- e^{i\phi} - a^\dagger \sigma^+ e^{-i\phi}]$$

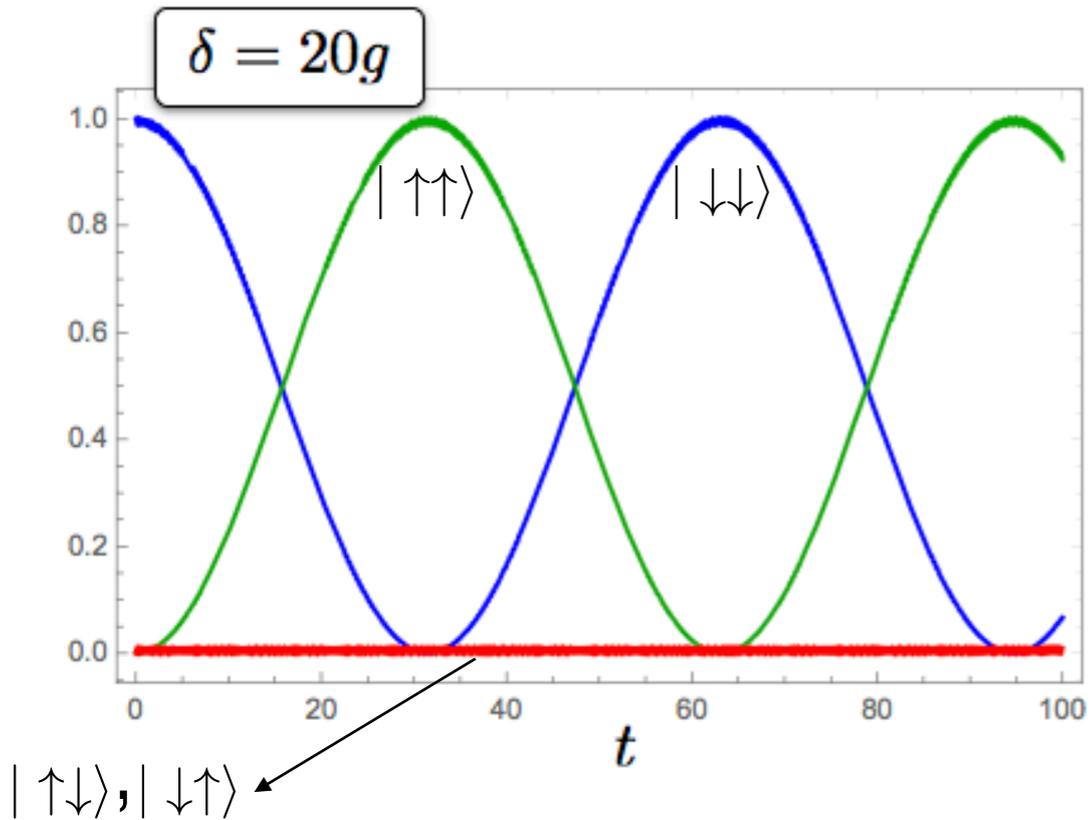
Spin-phonon transitions



TWO-QUBIT ENTANGLING OPERATION

Adiabatic elimination technique and the use of sideband transitions effectively couples two spins

$|\uparrow\uparrow\rangle$



ω_m

$|\downarrow\uparrow\rangle$

$I \delta$

g

g

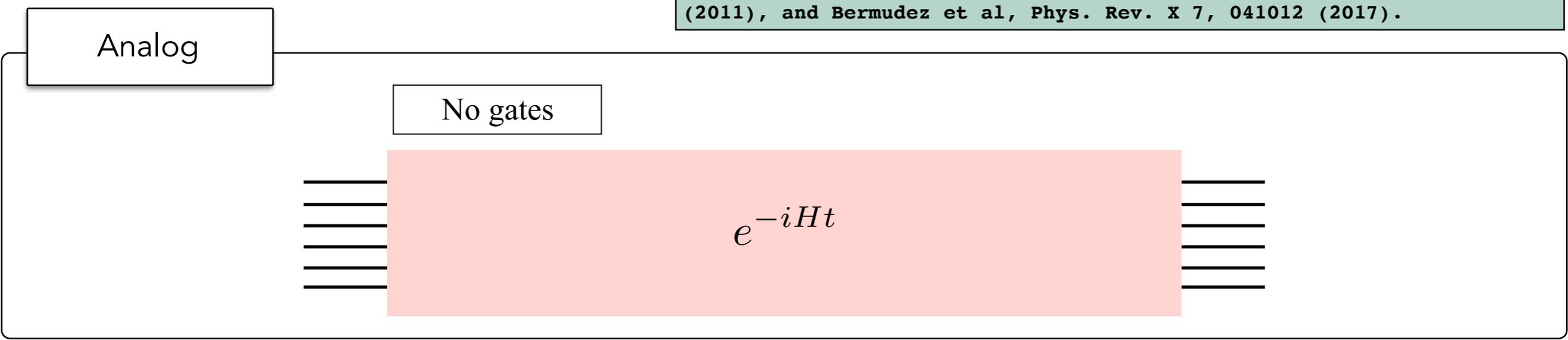
$|\downarrow\downarrow\rangle$

$$H_{\text{eff}} \propto \sigma_x^{(i)} \otimes \sigma_x^{(j)}$$

DIGITAL, ANALOG, AND HYBRID
MODES OF TRAPPED-ION SIMULATORS

Analog

No gates



The diagram shows a quantum circuit with a single gate. On the left, there are five horizontal lines representing qubits. These lines enter a large pink rectangular gate. The gate is labeled with the mathematical expression e^{-iHt} . On the right side of the gate, the five lines exit. The entire circuit is enclosed in a rounded rectangular frame. A box labeled 'Analog' is connected to the left side of the frame, and a box labeled 'No gates' is positioned above the gate.

$$e^{-iHt}$$

ZD, Hafezi, Monroe, Pagano, Seif, Shaw, Phys. Rev. Research, 2, 023015 (2020), arXiv: 1908.03210 [quant-ph].

Andrade, ZD, Grass, Hafezi, Pagano, Seif, arXiv: 2107.xxxx [quant-ph].

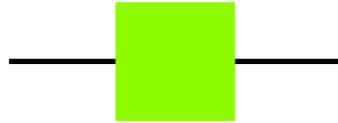
See also: Bermudez et al, Pays.Rev.A79, 060303 R (2009).

See also interesting work by:

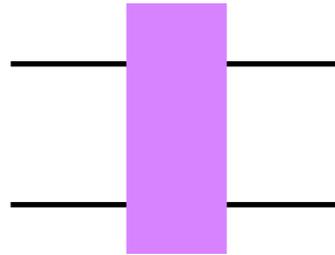
Lamata et al, Phys. Rev. Lett. 98 253005 (2007), Gerritsma et al, Nature 463, 68 (2010), Casanova et al, Phys. Rev. Lett. 107, 260501 (2011), and Bermudez et al, Phys. Rev. X 7, 041012 (2017).

Digital

Single-spin gates

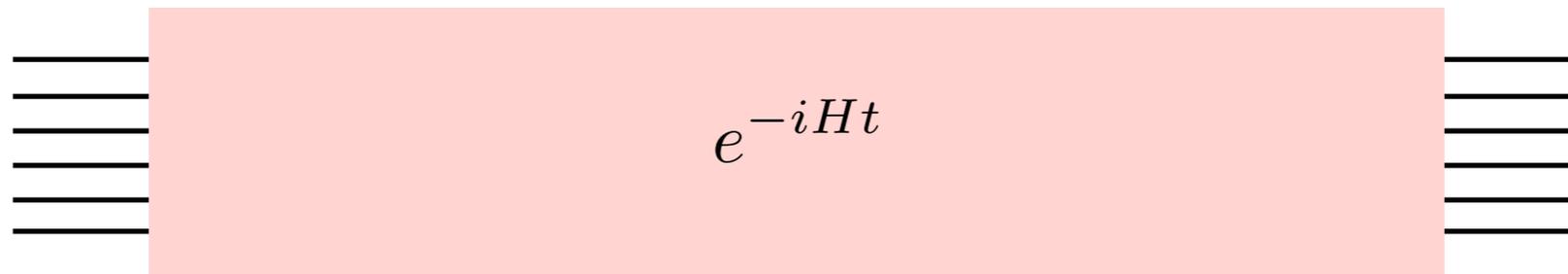


Two-spin gate (MS)



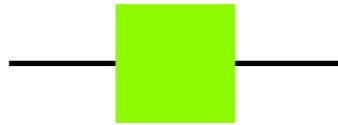
Analog

No gates

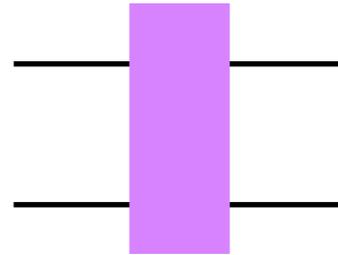


Digital

Single-spin gates



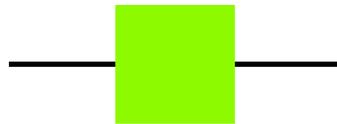
Two-spin gate (MS)



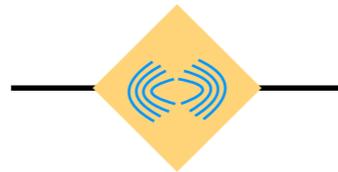
Analog-Digital

ZD, Linke, Pagano, arXiv:2104.09346 [quant-ph].

Single-spin gates



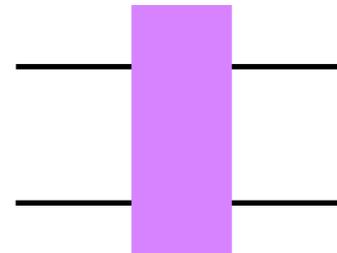
Spin-(normal) phonon gate



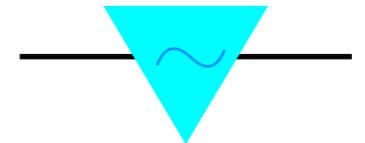
Spin-(local) phonon gate



Two-spin gate (MS)



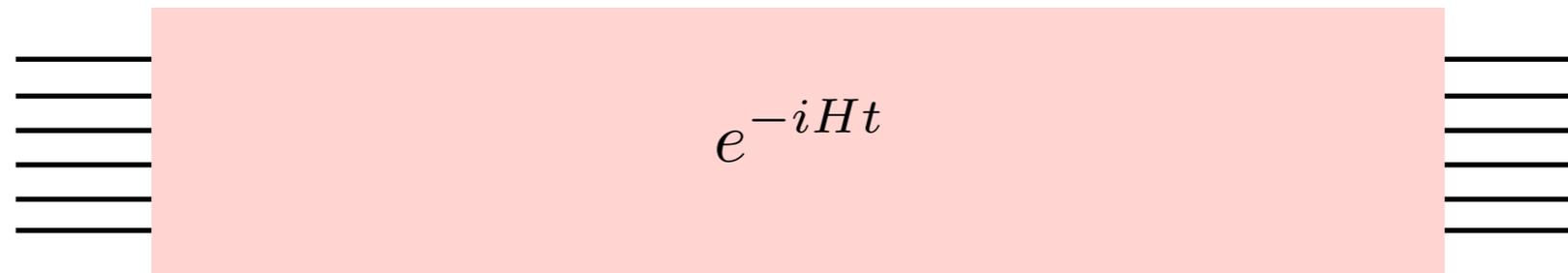
Standing-wave gate



Analog

For standing-wave tool, see: Porras and Cirac, Phys. Rev. Lett. 93, 263602 (2004).

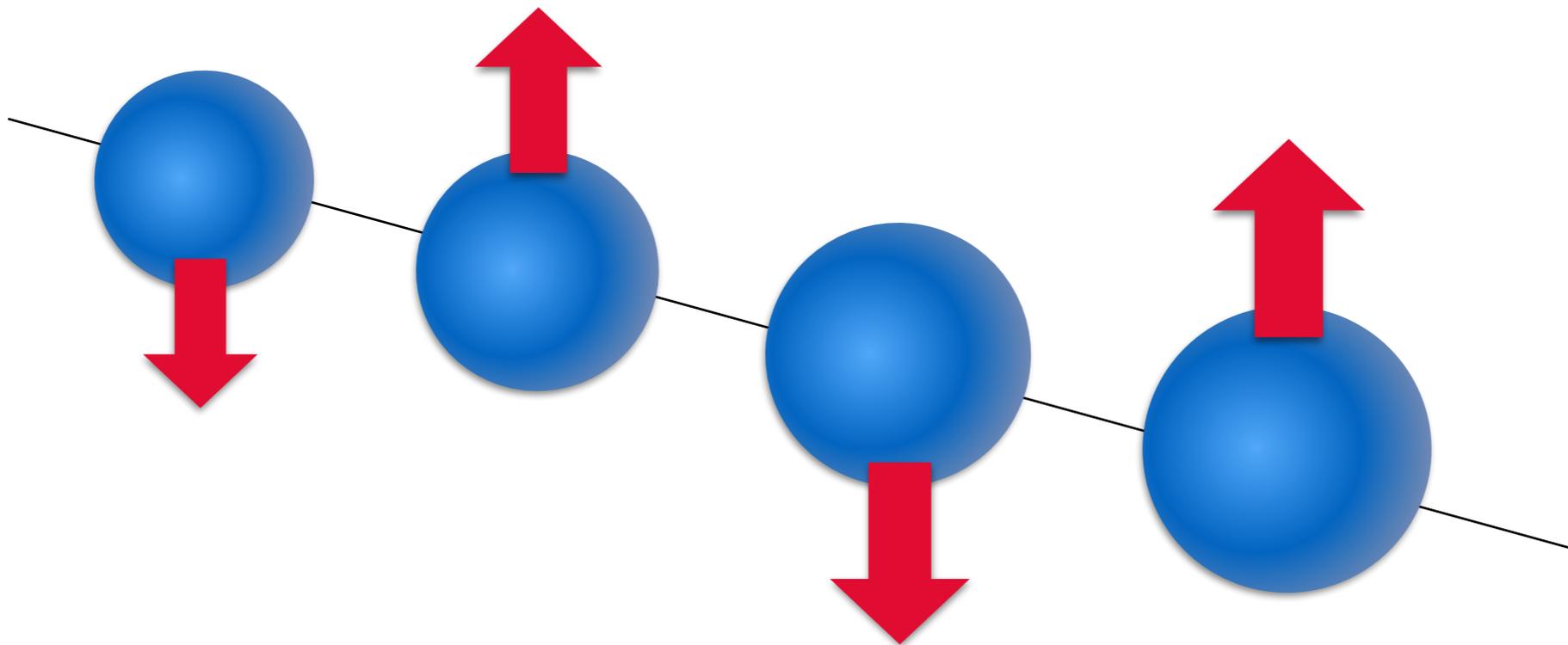
No gates

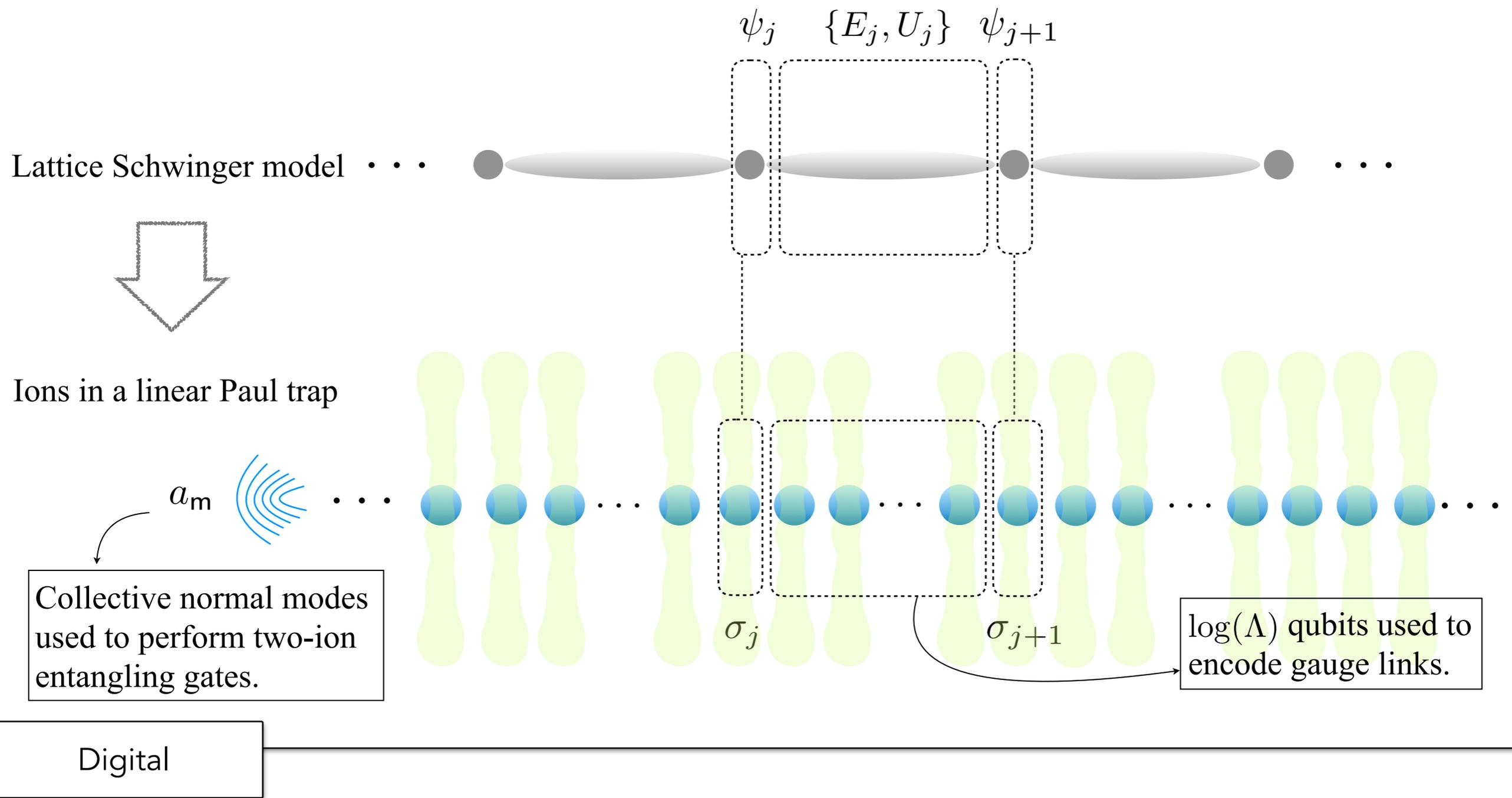


A LATTICE GAUGE THEORY EXAMPLE
STUDIED WITHIN DIGITAL AND HYBRID
MODES OF THE TRAPPED-ION SIMULATOR

LATTICE SCHWINGER MODEL: A TESTBED FOR QUANTUM
SIMULATION OF LATTICE GAUGE THEORIES

$$H = -ix \sum_{n=1}^{N-1} [\psi_n^\dagger U_n \psi_{n+1} - \text{h.c.}] + \sum_{n=1}^{N-1} E_n^2 + \mu \sum_{n=1}^N (-1)^n \psi_n^\dagger \psi_n$$





$$H = -ix \sum_{n=1}^{N-1} [\psi_n^\dagger U_n \psi_{n+1} - \text{h.c.}] + \sum_{n=1}^{N-1} E_n^2 + \mu \sum_{n=1}^N (-1)^n \psi_n^\dagger \psi_n$$

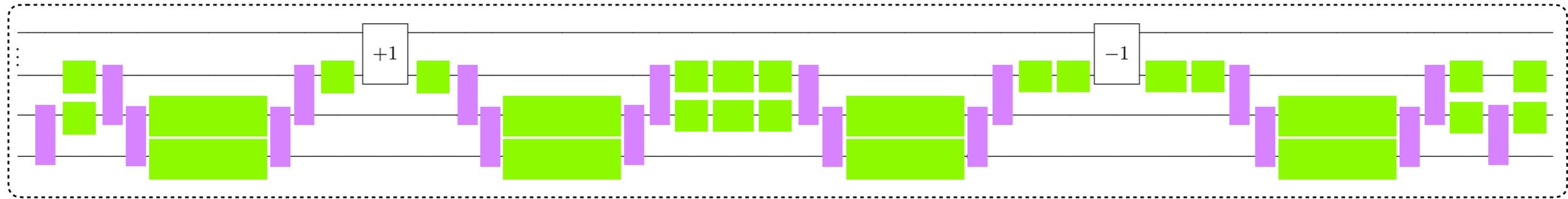
Lattice Schwinger model ...

$$\psi_j \quad \{E_j, U_j\} \quad \psi_{j+1}$$

Circuit and recourse analysis

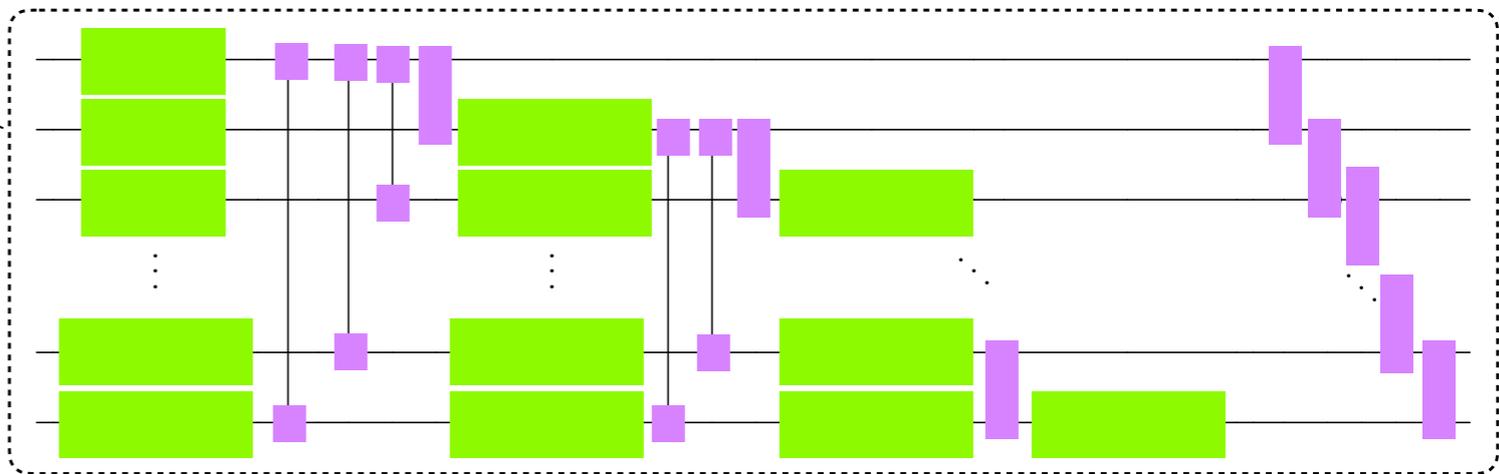
Alexander Shaw's talk in this conference.

Shaw, Lougovski, Stryker, Wiebe, Quantum 4, 306 (2020)



Sample gauge-fermion interaction block

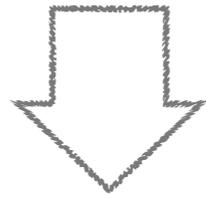
Part of electric field interactions acting on gauge DOF registers



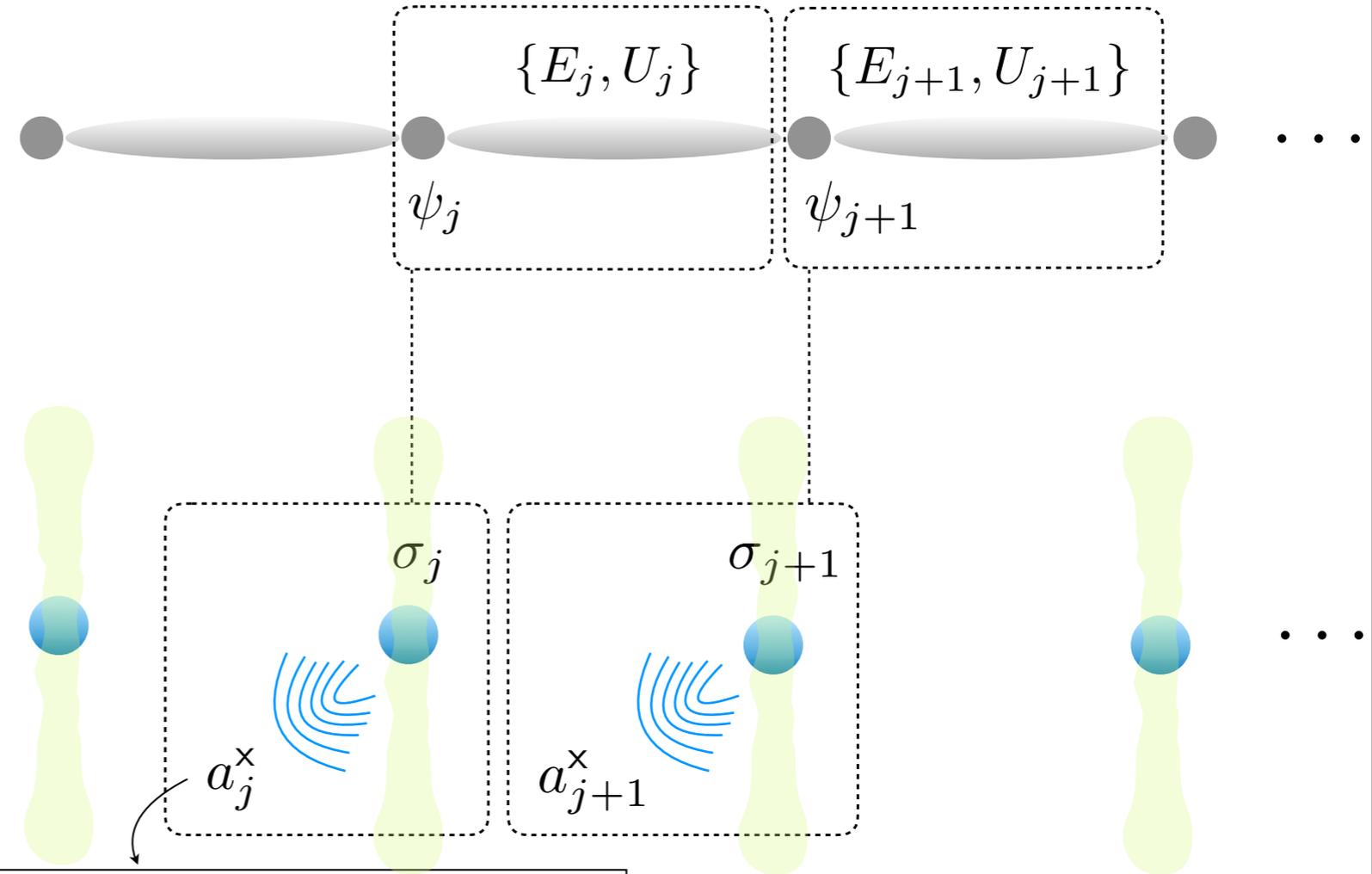
Near term cost

	$\delta_g = 10^{-3}$		$\delta_g = 10^{-4}$		$\delta_g = 10^{-5}$		$\delta_g = 10^{-6}$		$\delta_g = 10^{-7}$	
	$\tilde{\epsilon}^2$	CNOT								
$x = 10^{-2}$	—	7.3e4	—	1.6e5	—	3.4e5	—	7.3e5	5.6e-2	1.6e6
$x = 10^{-1}$	—	1.6e4	—	3.5e4	—	7.5e4	5.9e-2	1.6e5	2.7e-3	3.5e5
$x = 1$	—	4.6e3	—	9.9e3	1.0e-1	2.1e4	4.7e-3	4.6e4	2.2e-4	9.9e4
$x = 10^2$	—	2.8e3	8.3e-1	6.1e3	3.8e-2	1.3e4	1.8e-3	2.8e4	8.2e-5	6.0e4

Lattice Schwinger model ...



Ions in a linear Paul trap



Collective normal modes used to perform two-ion entangling gates.

Local transverse modes used to encode the dynamic of the gauge fields.

Analog-Digital

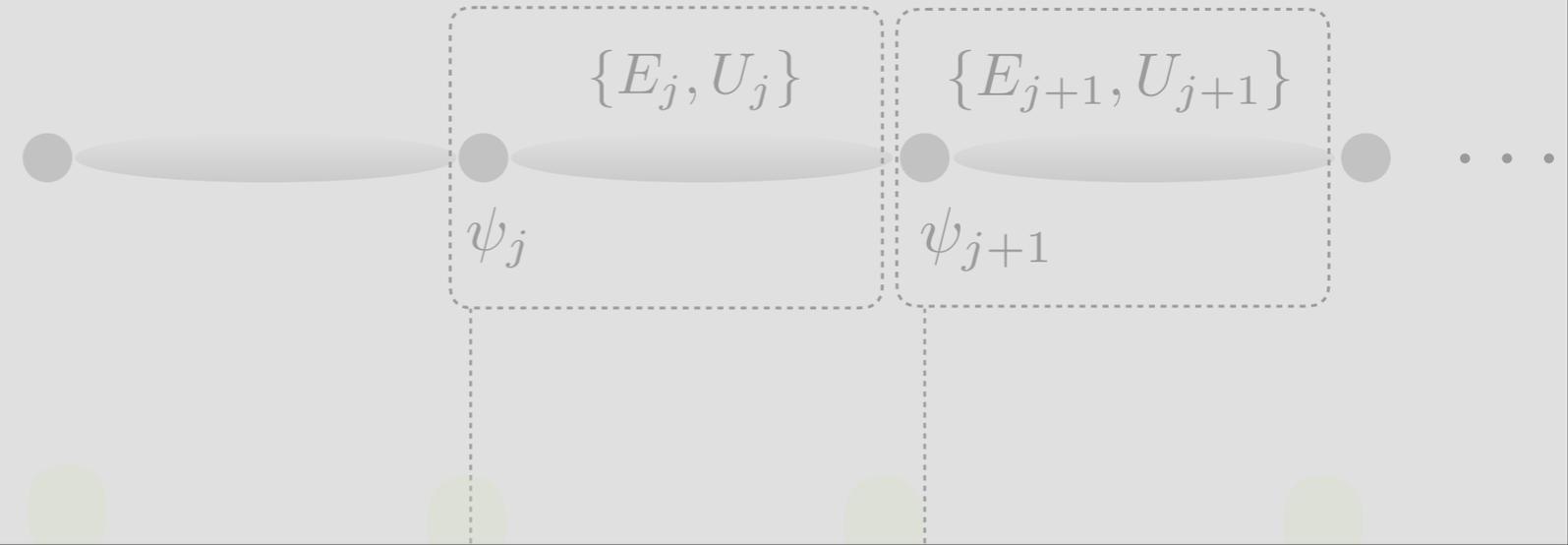
ZD, Linke, Pagano, arXiv:2104.09346 [quant-ph].

See Yang et al, Phys. Rev. A 94, 052321 (2016) for the highly-occupied bosonic model of the Schwinger model.

See also Casanova et al, Phys. Rev. Lett. 108, 190502 (2012), Lamata et al, EPJ Quant. Technol. 1, 9 (2014), and Mezzacapo et al, Phys. Rev. Lett. 109, 200501 (2012) for analog-digital approaches to other interacting fermion-boson theories.

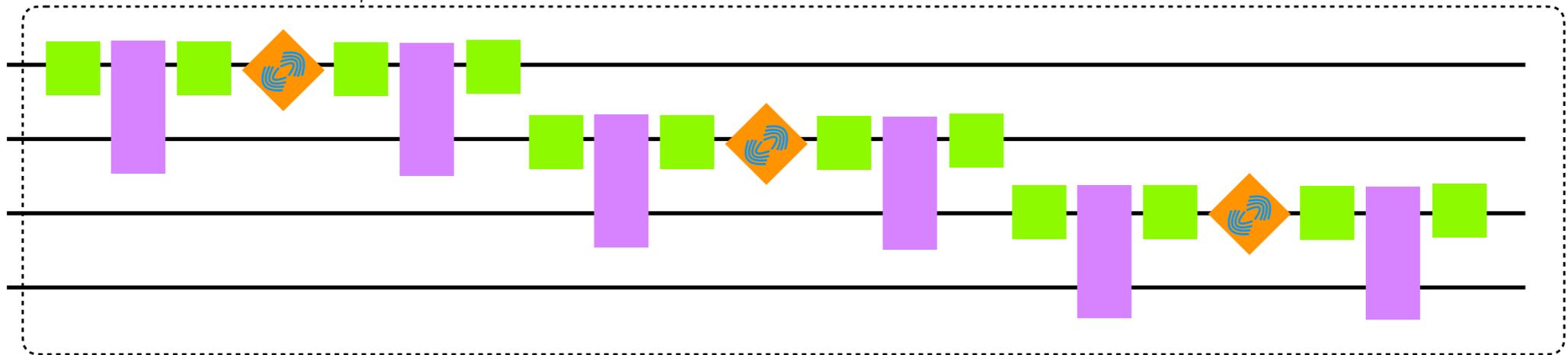
$$H = -ix \sum_{n=1}^{N-1} [\psi_n^\dagger U_n \psi_{n+1} - \text{h.c.}] + \sum_{n=1}^{N-1} E_n^2 + \mu \sum_{n=1}^N (-1)^n \psi_n^\dagger \psi_n$$

Lattice Schwinger model ...



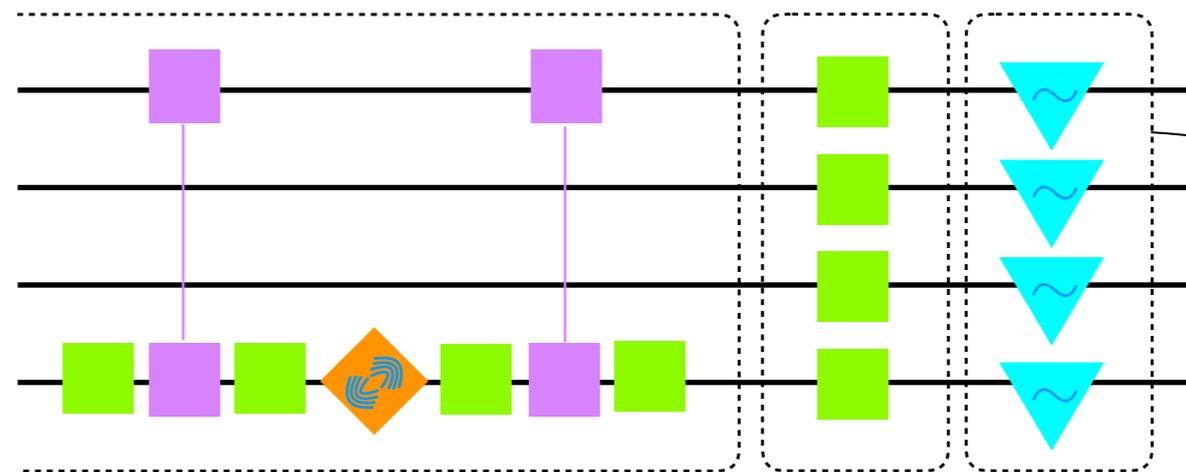
Analog-Digital

Fermion-gauge interactions



Collective used to pe entangling

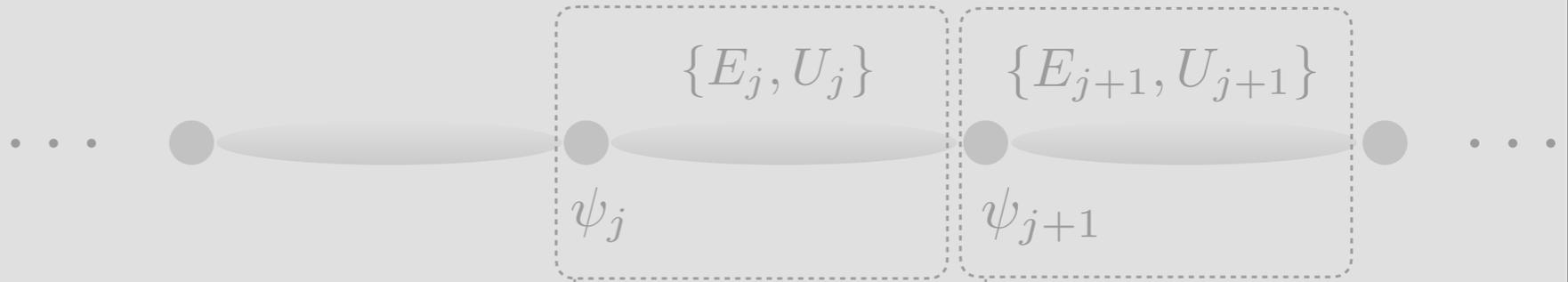
Gauge-field interactions



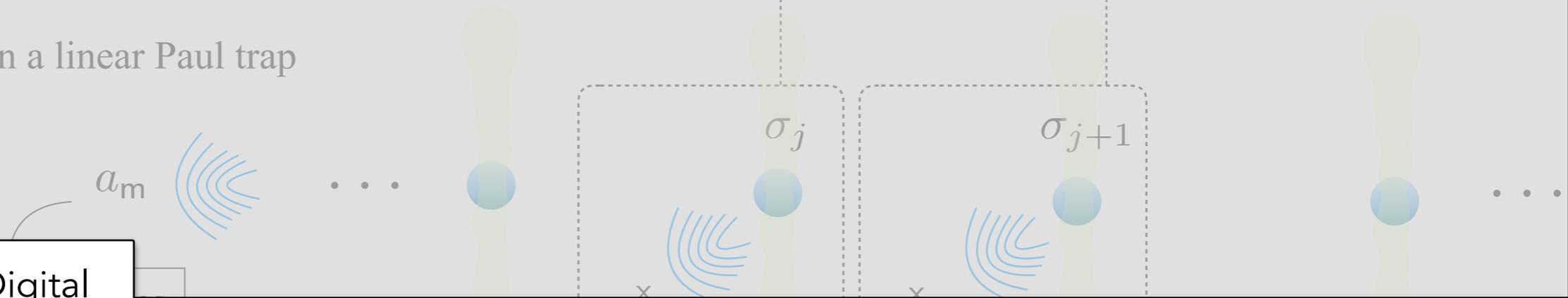
Fermion mass term

ZD, Linke, Pagano,
arXiv:2104.09346 [quant-ph].

Lattice Schwinger model



Ions in a linear Paul trap



Analog-Digital

ZD, Linke, Pagano, arXiv:2104.09346 [quant-ph].

Schwinger model			
	Fermion-gauge interaction	Fermion mass	Electric-field term
Analog-digital	$\mathcal{O}(N)$	$\mathcal{O}(1)$	$\mathcal{O}(N)$
Digital	$\mathcal{O}(N^2 (\log \Lambda)^2)$	$\mathcal{O}(1)$	$\mathcal{O}(N (\log \Lambda)^2)$

Let us compare the circuit structure of digital and analog-digital cases when gauge DOF are present.

A SCALAR FIELD THEORY COUPLED TO
FERMIONS (YUKAWA THEORY) WITH THE
HYBRID MODE OF THE SIMULATOR

$$H_{\text{Yukawa}}^{(I)} = \sum_{j=1}^N \left[\frac{i}{2b} (\psi_j^\dagger \psi_{j+1} - \psi_{j+1}^\dagger \psi_j) + m_\psi (-1)^j \psi_j^\dagger \psi_j \right]$$

$$H_{\text{Yukawa}}^{(II)} = b \sum_{j=1}^N \left[\frac{\Pi_j^2}{2} + \frac{(\nabla \varphi_j)^2}{2} + \frac{m_\varphi^2}{2} \varphi_j^2 \right]$$

$$H_{\text{Yukawa}}^{(III)} = gb \sum_{j=1}^N \psi_j^\dagger \varphi_j \psi_j,$$



Trapped-ion Hamiltonian

Model Hamiltonian

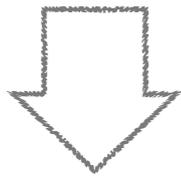
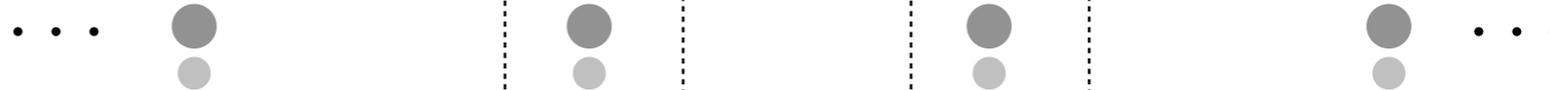
$$H_{\text{Yukawa}}^{(I)'} = \frac{1}{4b} \sum_{j=1}^N \sigma_j^x \sigma_{j+1}^x,$$

$$H_{\text{Yukawa}}^{(II)'} = \frac{1}{4b} \sum_{j=1}^N \sigma_j^y \sigma_{j+1}^y,$$

$$H_{\text{Yukawa}}^{(III)'} = \frac{m_\psi}{2} \sum_{j=1}^N (-1)^j \sigma_j^z + \text{const.},$$

$$H_{\text{Yukawa}}^{(IV)'} = \sqrt{\frac{g^2 b}{8N}} \sum_{j=1}^N (\mathbb{I}_j + \sigma_j^z) \sum_{m=1}^N \frac{1}{\sqrt{\varepsilon_m}} \times \\ (a_m^\dagger e^{-i \frac{2\pi j}{N} (m - \frac{N}{2} - 1)} + a_m e^{i \frac{2\pi j}{N} (m - \frac{N}{2} - 1)}) + \sum_{m=1}^N \varepsilon_m (a_m^\dagger a_m + \frac{1}{2}).$$

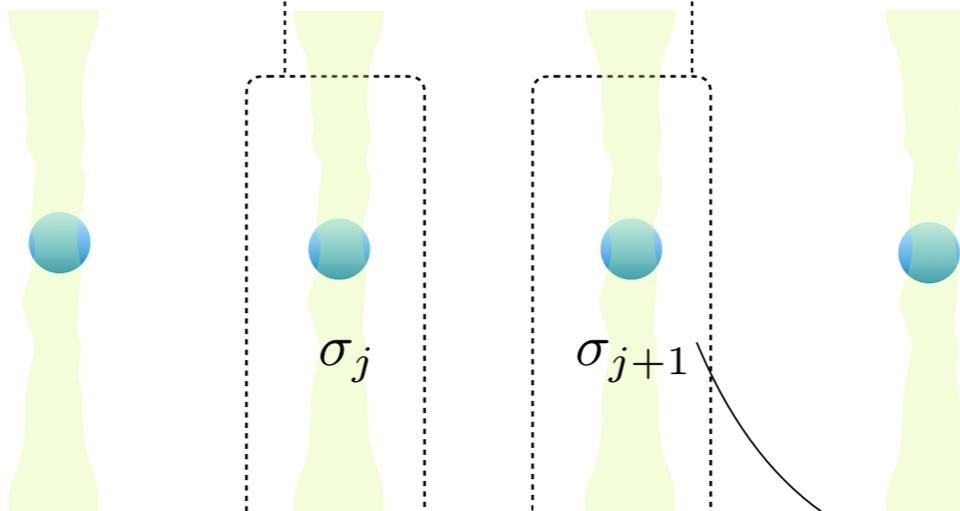
A Yukawa theory: scalar field coupled to fermions



Ions in a linear Paul trap



Collective normal modes are used to simulate the dynamic of scalar field and to perform spin-spin entangling gates.

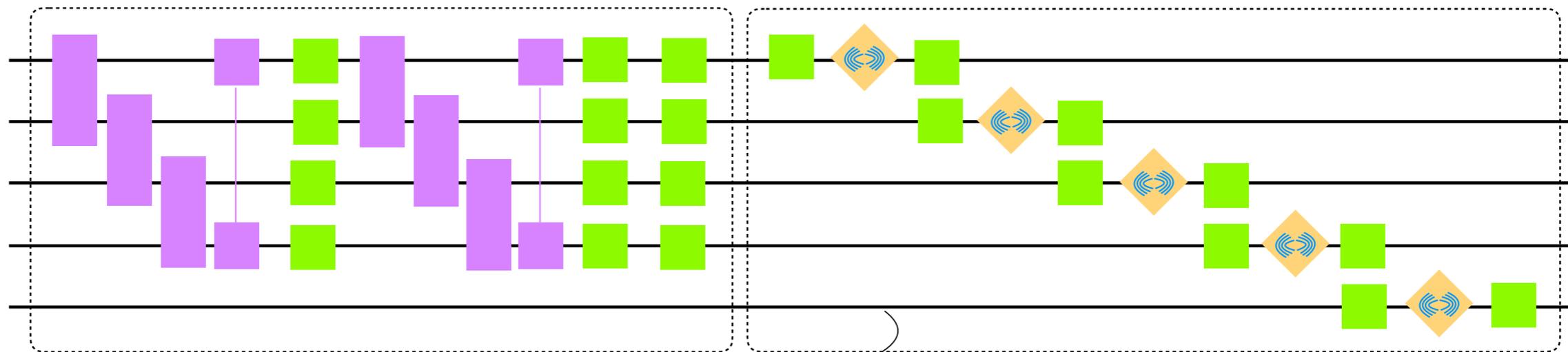
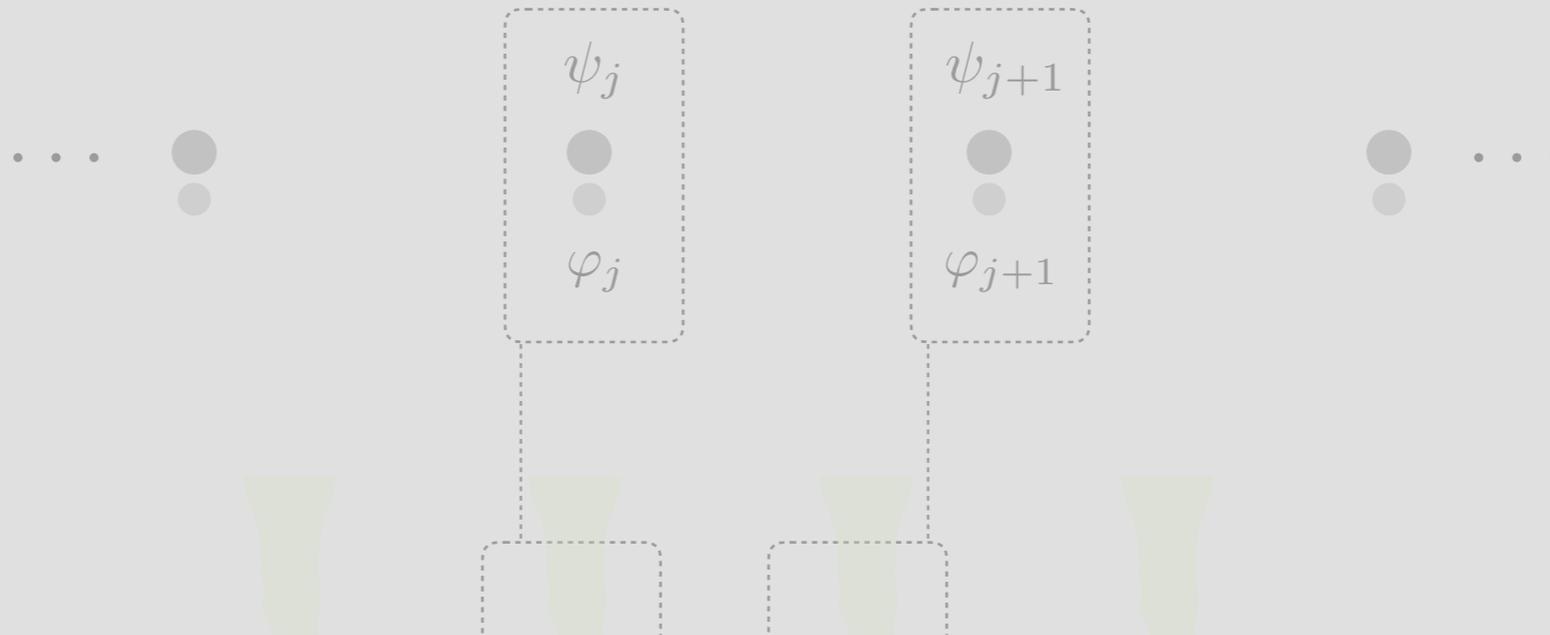


Internal states of the ion are used to encode the dynamic of fermions.

A Yukawa theory: scalar field coupled to fermions



Ions in a linear Paul trap



Free fermion terms

Ancilla ion

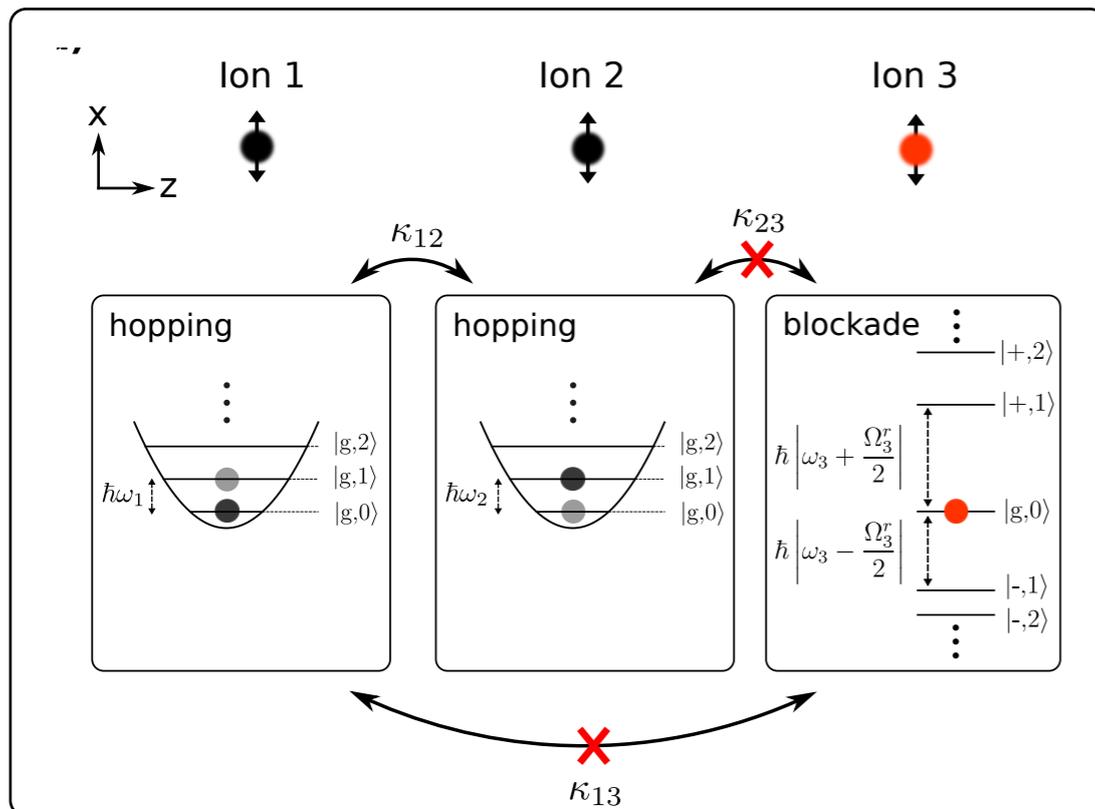
Free scalar field and fermion-scalar field interactions

Yukawa theory

	Fermion hopping	Fermion mass	Free scalar fields	Fermion scalar-field interaction
Analog-digital	$\mathcal{O}(N)$	$\mathcal{O}(1)$	$\mathcal{O}(1)$	$\mathcal{O}(1)$
Digital	$\mathcal{O}(N)$	$\mathcal{O}(1)$	$\mathcal{O}(1)$	$\mathcal{O}(N^2 (\log \Lambda)^2)$

IS PHONON CONTROL EXPERIMENTALLY
FEASIBLE? YES...AT LEAST FOR SMALL
SYSTEMS SO FAR!

Phonon hopping and blockade demonstrated in trapped-ion simulators:



Debnath et al, Phys. Rev. Lett. 120, 073001 (2018).

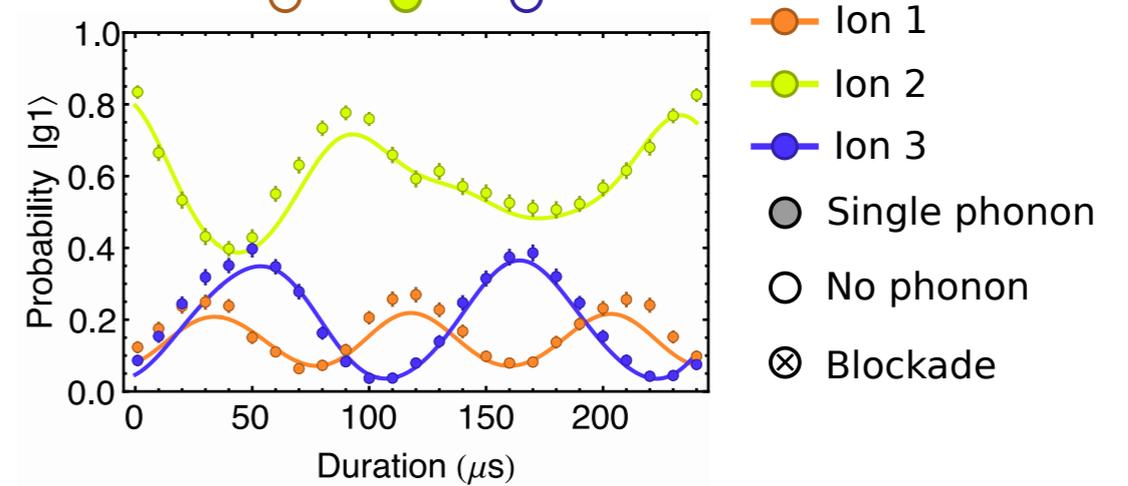


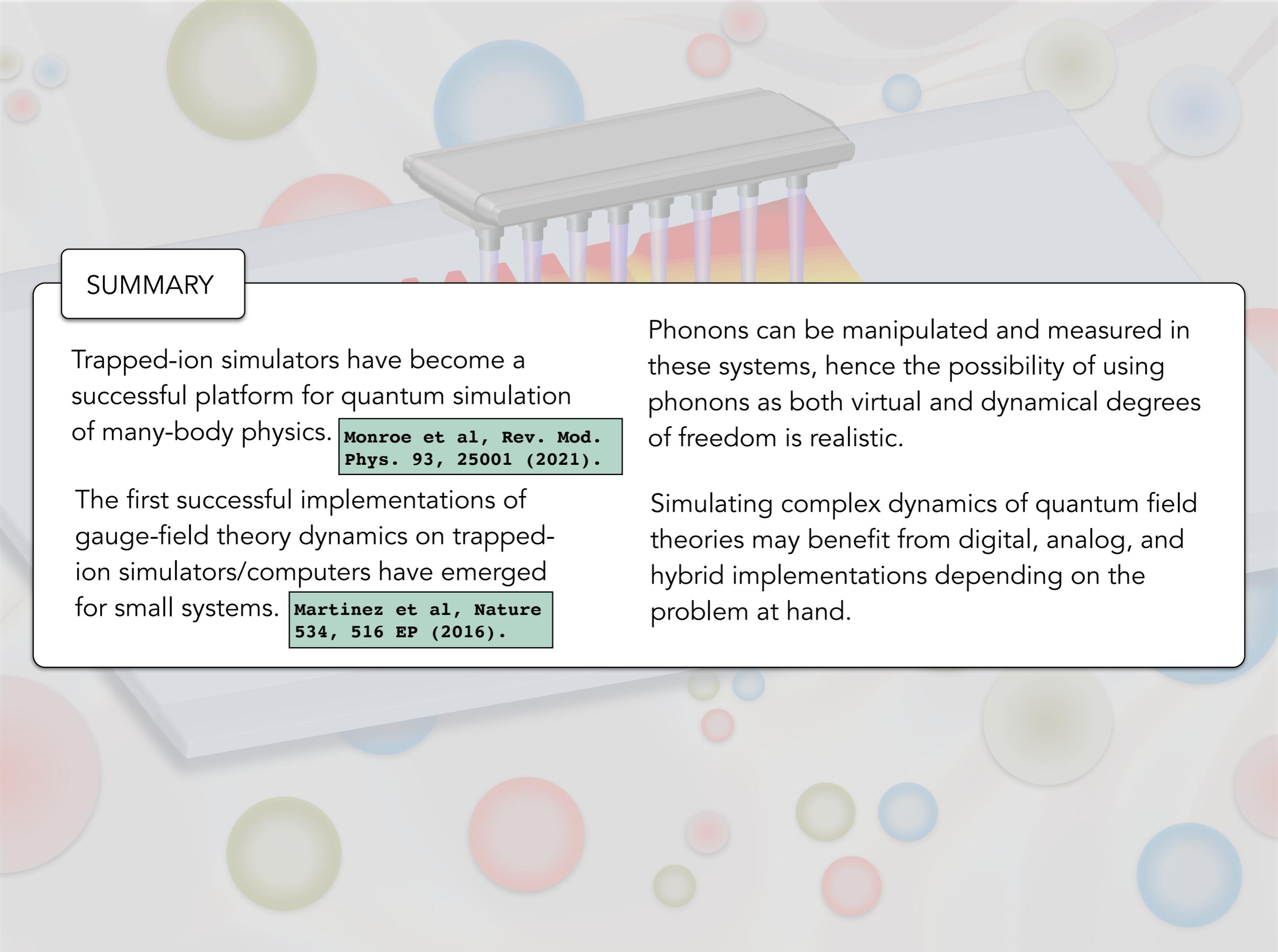
UNIVERSITY OF MARYLAND

Monroe-Linke Group

b)

ions 1 2 3





SUMMARY

Trapped-ion simulators have become a successful platform for quantum simulation of many-body physics.

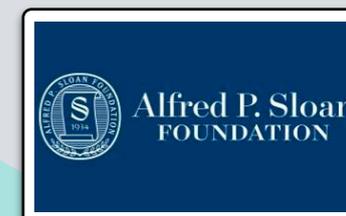
Monroe et al, Rev. Mod. Phys. 93, 25001 (2021).

The first successful implementations of gauge-field theory dynamics on trapped-ion simulators/computers have emerged for small systems.

Martinez et al, Nature 534, 516 EP (2016).

Phonons can be manipulated and measured in these systems, hence the possibility of using phonons as both virtual and dynamical degrees of freedom is realistic.

Simulating complex dynamics of quantum field theories may benefit from digital, analog, and hybrid implementations depending on the problem at hand.



THANK YOU