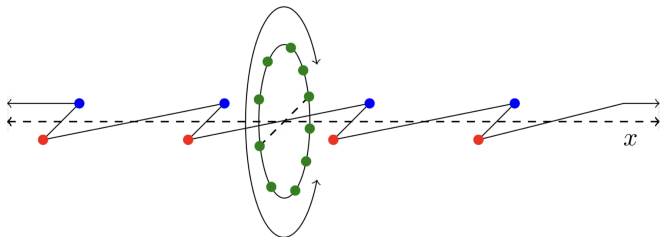


Lattice July 2021

Quantum Algorithms for Simulating the Lattice Schwinger Model



Alexander F. Shaw

University of Maryland - College Park

Based on work of same title in Quantum 4, 306 (2020)
by AF Shaw, P Lougovski, JR Stryker and N Wiebe

Why Quantum Simulation?

Elements we can simulate from
First Principles

Periodic Table of the Elements

1	2											10	11	12	13	14	15	16	17	18	
H	He											Ne	Ar	Kr	Xe	Rn					
Li	Be											B	C	N	O	F	Ne				
Na	Mg											Al	Si	P	S	Cl	Ar				
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr				
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe				
Cs	Ba	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn					
Fr	Ra	Rf	Db	Sg	Bh	Hs	Mt	Ds	Rg	Cn	Nh	Fl	Mc	Lv	Ts	Og					
Lanthanide Series		57	58	59	60	61	62	63	64	65	66	67	68	69	70	71	72				
Actinide Series		89	90	91	92	93	94	95	96	97	98	99	100	101	102	103	104				

What we'd like to simulate

*[Yamazaki et. al (2010) PhysRevD.81.111504]

- ▶ Using fundamental models to calculate physics at longer length scales is hard.

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Cs		Ba										Hf				Ta									
Fr		Ra										Rf				Db									
La		Ce										Pr				Nd									
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- ▶ Access different physics (real-time dynamics)

Why Quantum Simulation?

Elements we can simulate from
First Principles

Periodic Table of the Elements

The image shows a standard periodic table of elements. The elements Hydrogen (H) and Helium (He) are circled in green. A red line outlines the rest of the periodic table, from Lithium (Li) to Oganesson (Og). A small asterisk is placed next to Helium (He).

What we'd like to simulate

*[Yamazaki et. al (2010) PhysRevD.81.111504]

- ▶ Using fundamental models to calculate physics at longer length scales is hard.
- ▶ Access different physics (real-time dynamics)
- ▶ Classical methods have been very successful, but hit roadblocks. (exponential resource scalings, sign problems)

Goal: Estimate Observables from QCD

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How to, using real-time dynamics:

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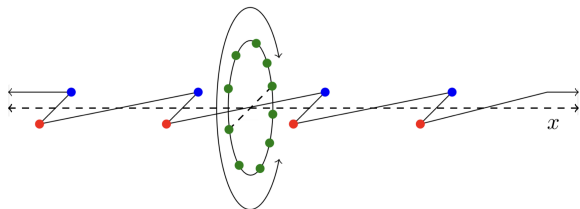
How hard is this to do on a quantum computer? Depends on the computational model.

	Cheap	Expensive
Near-term (NISQ)	Single-qubit rotations	Entangling gates (CNOT)
Far-term (Fault-Tolerant)	Clifford + CNOT gates	T-gates

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The Lattice Schwinger Model*



odd sites - electrons

even sites - positrons

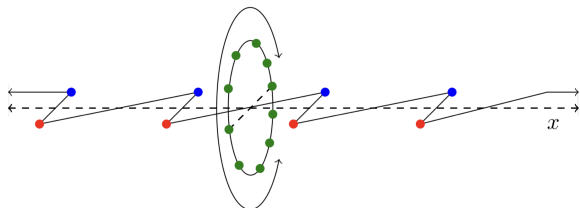
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\implies
continuum limit

QED in (1+1)D

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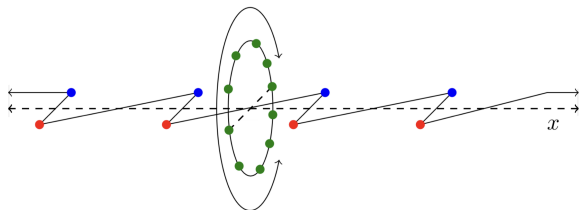
QED in (1+1)D

Simplest model that replicates features of QCD (confinement, spontaneous breaking of chiral symmetry).

*[Kogut, Susskind (1975) 10.1103/PhysRevD.11.395]

Computing time evolution

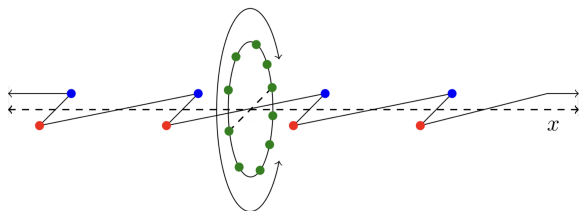
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Computational basis = occupation basis.

Computing time evolution

The Lattice Schwinger Model

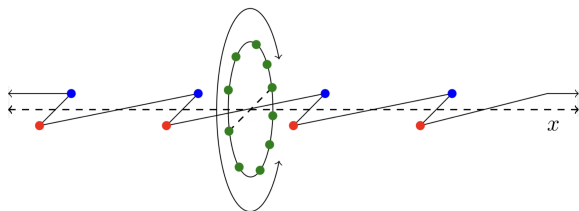


Computational basis = occupation basis. EX: for a fermion on an odd site (after Jordan-Wigner):

computational basis	occupation basis
$ 0\rangle$	presence of electron
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Computing time evolution

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The value of the electric field is represented as a binary integer in a qubit register, one for each link.

Computing time evolution

Decompose the Schwinger Model Hamiltonian

$$\begin{aligned} H_{\text{Schwinger}} &= \sum_r \hat{E}_r^2 + \mu \sum_r (-)^r \hat{\psi}_r^\dagger \hat{\psi}_r + \times \sum_r \left[\hat{U}_r \hat{\psi}_r^\dagger \hat{\psi}_{r+1} - \hat{U}_r^\dagger \hat{\psi}_r \hat{\psi}_{r+1}^\dagger \right] \\ &= \underbrace{H_{\text{E-field}} + H_{\text{mass}}}_{\text{diagonal in computational basis}} + \underbrace{H_{\text{interaction}}}_{\text{linear combination of tensor product of pauli matrices}} \\ &= \sum_{j=1}^m H_j \quad \leftarrow \quad \text{Simulatable!} \end{aligned}$$

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$$V_{\text{trotter step}}(t) := \prod_{j=1}^m e^{-iH_j t/2} \prod_{j=m}^1 e^{-iH_j t/2}$$

$$\begin{aligned} \left\| e^{-iH_{\text{Schwinger}} t} - V_{\text{trotter step}}(t) \right\| &\leq \frac{1}{12} \sum_{x,y>x,z>x} \left\| [[H_x, H_y], H_z] \right\| t^3 \\ &+ \frac{1}{24} \sum_{x,y>x} \left\| [[H_x, H_y], H_x] \right\| t^3 \doteq \delta(\text{lattice parameters, evolution time}). \end{aligned}$$

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- ▶ Decomposition of $H_{\text{Schwinger}} = \sum_j H_j$

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- ▶ **Upper bound** on sufficient quantum-computational resources required for an arbitrary simulation:

COST(total evolution time, error in final state, lattice parameters)

$$\in \tilde{O}\left(\frac{N^{3/2} T^{3/2} \Lambda x^{1/2}}{\delta^{1/2}}\right)$$

where N = lattice size, T = evolution time, Λ = E-field cutoff, and $x = 1/(ag)^2$ (a lattice spacing, g coupling constant)

Data and Conclusion (Room for Improvement)

We compile these results further with **estimation of mean positron/electron density** :

Near-term (NISQ) Simulation (no ancilla)

	$\delta_g = 10^{-3}$		$\delta_g = 10^{-4}$		$\delta_g = 10^{-5}$		$\delta_g = 10^{-6}$		$\delta_g = 10^{-7}$	
	$\tilde{\epsilon}^2$	CNOT	$\tilde{\epsilon}^2$	CNOT	$\tilde{\epsilon}^2$	CNOT	$\tilde{\epsilon}^2$	CNOT	$\tilde{\epsilon}^2$	CNOT
$x = 10^{-2}$	—	7.3e4	—	1.6e5	—	3.4e5	—	7.3e5	5.6e-2	1.6e6
$x = 10^{-1}$	—	1.6e4	—	3.5e4	—	7.5e4	5.9e-2	1.6e5	2.7e-3	3.5e5
$x = 1$	—	4.6e3	—	9.9e3	1.0e-1	2.1e4	4.7e-3	4.6e4	2.2e-4	9.9e4
$x = 10^2$	—	2.8e3	8.3e-1	6.1e3	3.8e-2	1.3e4	1.8e-3	2.8e4	8.2e-5	6.0e4

$x = (ag)^{-2}$ with a lattice spacing and g coupling constant

δ_g = error in CNOT channel

$\tilde{\epsilon}^2$ = worst case mean square error in mean positron density

Data and Conclusion (Room for Improvement)

Far-term Simulation (N is size of lattice, Λ is electric cutoff)

Upper Bounds on T-gate Cost of Specific Simulations ($\mu = 1, \epsilon^2 = 0.1$)

	Short Time ($T = 10/x$)		Long Time ($T = 1000/x$)	
	Sampling	Estimating	Sampling	Estimating
$N = 4, \Lambda = 2$				
Strong Coupling ($x = 0.1$)	$6.5 \cdot 10^7$	$2.4 \cdot 10^{11}$	$8.8 \cdot 10^{10}$	$3.3 \cdot 10^{14}$
Weak Coupling ($x = 10$)	$5.0 \cdot 10^6$	$1.8 \cdot 10^{10}$	$7.0 \cdot 10^9$	$2.6 \cdot 10^{13}$
$N = 16, \Lambda = 2$				
Strong Coupling ($x = 0.1$)	$7.2 \cdot 10^8$	$2.5 \cdot 10^{12}$	$9.4 \cdot 10^{11}$	$3.3 \cdot 10^{15}$
Weak Coupling ($x = 10$)	$5.6 \cdot 10^7$	$1.9 \cdot 10^{11}$	$7.6 \cdot 10^{10}$	$2.7 \cdot 10^{14}$
$N = 16, \Lambda = 4$				
Strong Coupling ($x = 0.1$)	$1.9 \cdot 10^9$	$6.3 \cdot 10^{12}$	$2.3 \cdot 10^{12}$	$8.1 \cdot 10^{15}$
Weak Coupling ($x = 10$)	$9.6 \cdot 10^7$	$3.2 \cdot 10^{11}$	$1.2 \cdot 10^{11}$	$4.2 \cdot 10^{14}$
$N = 64, \Lambda = 2$				
Strong Coupling ($x = 0.1$)	$6.6 \cdot 10^9$	$2.1 \cdot 10^{13}$	$8.5 \cdot 10^{12}$	$2.9 \cdot 10^{16}$
Weak Coupling ($x = 10$)	$5.2 \cdot 10^8$	$1.6 \cdot 10^{12}$	$6.9 \cdot 10^{11}$	$2.3 \cdot 10^{15}$
$N = 64, \Lambda = 4$				
Strong Coupling ($x = 0.1$)	$1.7 \cdot 10^{10}$	$5.4 \cdot 10^{13}$	$2.0 \cdot 10^{13}$	$6.9 \cdot 10^{16}$
Weak Coupling ($x = 10$)	$8.7 \cdot 10^8$	$2.7 \cdot 10^{12}$	$1.1 \cdot 10^{12}$	$3.6 \cdot 10^{15}$
$N = 64, \Lambda = 8$				
Strong Coupling ($x = 0.1$)	$4.5 \cdot 10^{10}$	$1.5 \cdot 10^{14}$	$5.3 \cdot 10^{13}$	$1.8 \cdot 10^{17}$
Weak Coupling ($x = 10$)	$1.5 \cdot 10^9$	$4.6 \cdot 10^{12}$	$1.7 \cdot 10^{12}$	$5.8 \cdot 10^{15}$

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Questions on the Road to Quantum Simulating QCD:

- ▶ Can our algorithms generalize to higher dimensions and $SU(N)$?
- ▶ How do other algorithms (LCU, Qubitization) compare?
- ▶ Could other formulations of the model be beneficial? ($SU(2)$ - Quantum Link Model, Loop-String-Hadron)
- ▶ How can we best realize gauge invariance during simulation?

Our work is a **benchmark** to aid in answering these questions.

LCU - [Childs, Wiebe (2012) 10.26421/QIC12.11-12]

Qubitization - [Low, Chuang (2019) 10.22331/q-2019-07-12-163]

Quantum Link Model - [Chandrasekharan, Wiese (1996) 10.1016/S0550-3213(97)80041-7]

Loop-String-Hadron - [Raychowdhury, Stryker (2019) 10.1103/PhysRevD.101.114502]

Preliminary Results - SU(2):

Using abelian gauge invariant hopping decomposition* - Discussed in J. Stryker's talk **this evening, 10:45pm EST**:

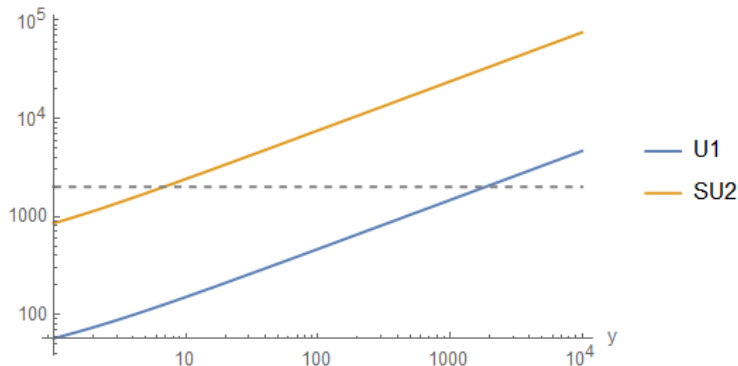
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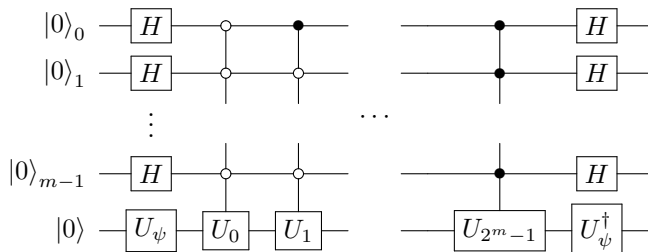
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$T(y, y, 1, y, 1/y)$ Trotter Steps



*[J. Stryker (2021) arXiv:2105.11548]

Using amplitude estimation to estimate $\langle \hat{O} \rangle$



If $\hat{O} = \sum_{a=0}^{2^m-1} \hat{U}_a$, \hat{U}_a unitary, then above circuit ends up in:

$\text{Prob}(\text{measure } |00\dots 0\rangle) = C \langle \hat{O} \rangle$, C known constant.