

Investigating a Renormalization Group Multigrid Approach for Domain Wall Fermions

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The ensembles used in this work were produced by Jiqun Tu.

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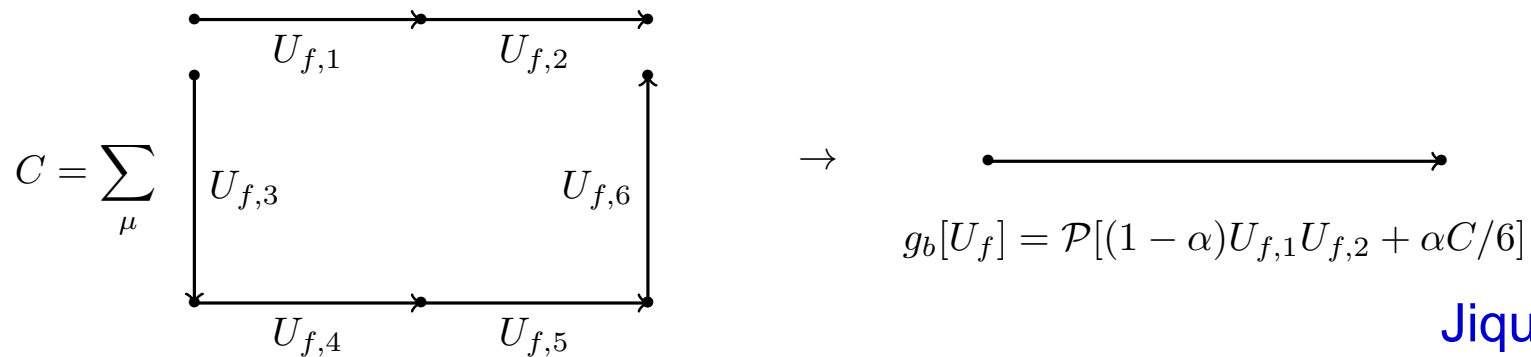
Motivation

- RBC-UKQCD ensemble generation using CG in production
 - * Shamir Domain wall and Mobius Domain Wall fermions used.
 - * Many optimizations employed: Hasenbusch masses, force gradient integrator, optimized code, ...
- In HMC/RHMC, no effective multigrid solver for (M)DWF to date
 - * Setup time for Hierarchically Deflated CG (Boyle) can be amortized during a trajectory, but no net speed-up (Boyle, McGlynn)
 - * In the next talk, Peter Boyle will talk about the general status and report on his efforts to combine multigrid ideas and domain decomposition within HMC.
- Given a fine ensemble, can a coarse ensemble with $2\times$ the lattice spacing and the same long distance physics, be used as a low-setup time preconditioner for fermion solves on the fine ensemble?
 - * Good a^2 scaling of MDWF + Iwasaki + DSDR ensembles intriguing
 - * Jiqun Tu generated such matched ensembles for other purposes [Lattice 2017, 10.1051/epjconf/201817502006, EPJ Web Conf. 175 (2018) 02006]

Matched Fine and Coarse Ensembles

	fine	coarse	blocked coarse
	$\langle O \rangle_f$	$\langle O \rangle_c$	$\langle O \rangle_c^b$
size	$24^3 \times 64 \times 12$	$12^3 \times 64 \times 12$	$12^3 \times 64 \times 12$
β	1.943	1.633	-
am_l	0.000787	0.008521	0.007494
am_h	0.019896	0.065073	0.064150
$a^{-1}(\text{GeV})$	2.001(18)	1.015(16)	1.010(16)
am_{res}	0.004522(12)	0.007439(86)	0.00847(21)
$m_\pi(\text{MeV})$	300(3)	307(5)	308(8)
$m_K(\text{MeV})$	491(5)	506(8)	507(11)
$m_\Omega(\text{MeV})$	1557(71)	1652(27)	1685(52)
$f_\pi(\text{MeV})$	138(2)	147(2)	151(3)
$f_K(\text{MeV})$	155(2)	166(3)	169(4)

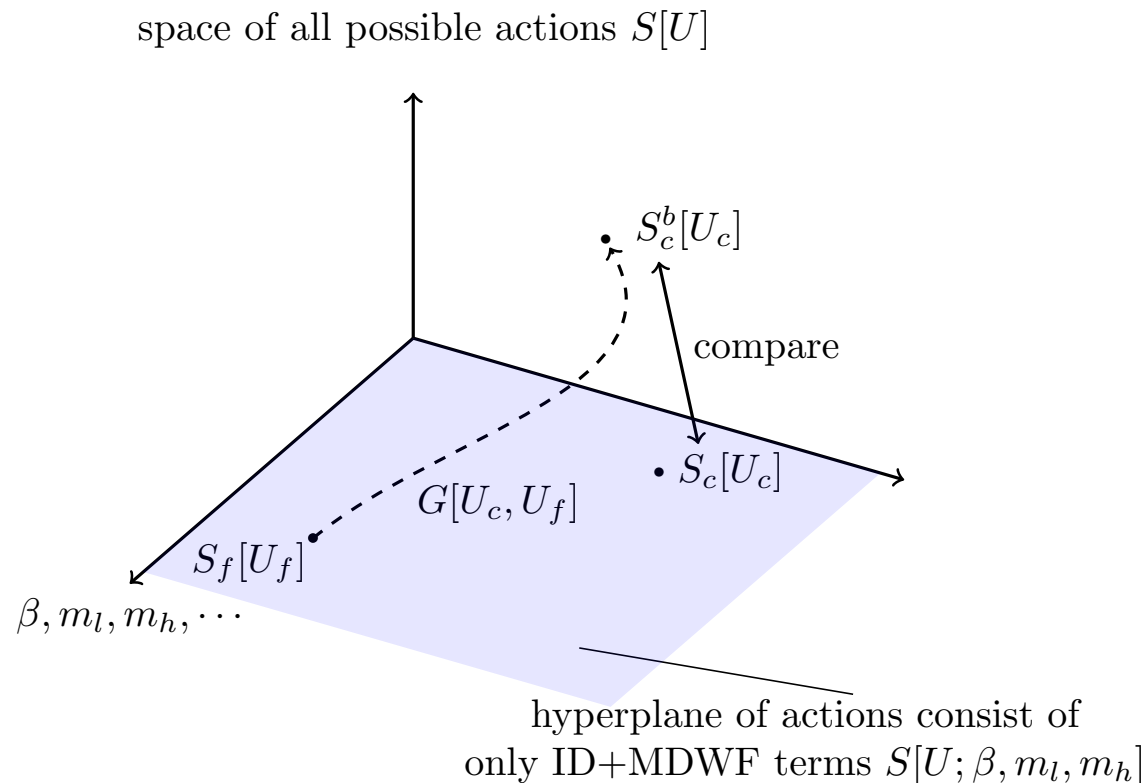
Blocked coarse ensemble generated from the fine ensemble with an APE style RG-blocking



Jiqun Tu

Rationale

- Small $O(a^2)$ scaling violations for coarse (1 GeV) and fine (2 GeV) ensembles
 - * This implies they lie on essentially the same RG trajectory.
- Simple RG blocking creates a coarse lattice from a given fine lattice
 - * Blocking fast to do numerically
- Since fine and blocked-coarse lattice have approximately identical physics, can the blocked lattice be used as a preconditioner for fine lattice DWF solves?



Preconditioned CG

Preconditioned CG algorithm

Result: solution for $Ax = b$

$$r_0 = b - Ax_0$$

$$z_0 = M^{-1}r_0$$

$$p_0 = z_0$$

$$k = 0$$

while $|r_{k+1}| > \epsilon$ **do**

$$\alpha_k = \frac{\langle r_k | z_k \rangle}{\langle p_k | A | p_k \rangle}$$

$$x_{k+1} = x_k + \alpha_k p_k$$

$$r_{k+1} = r_k - \alpha_k A p_k$$

$$z_{k+1} = M^{-1} r_{k+1}$$

$$\beta_k = \frac{\langle z_{k+1} | r_{k+1} \rangle}{\langle z_k | r_k \rangle}$$

$$p_{k+1} = z_{k+1} + \beta_k p_k$$

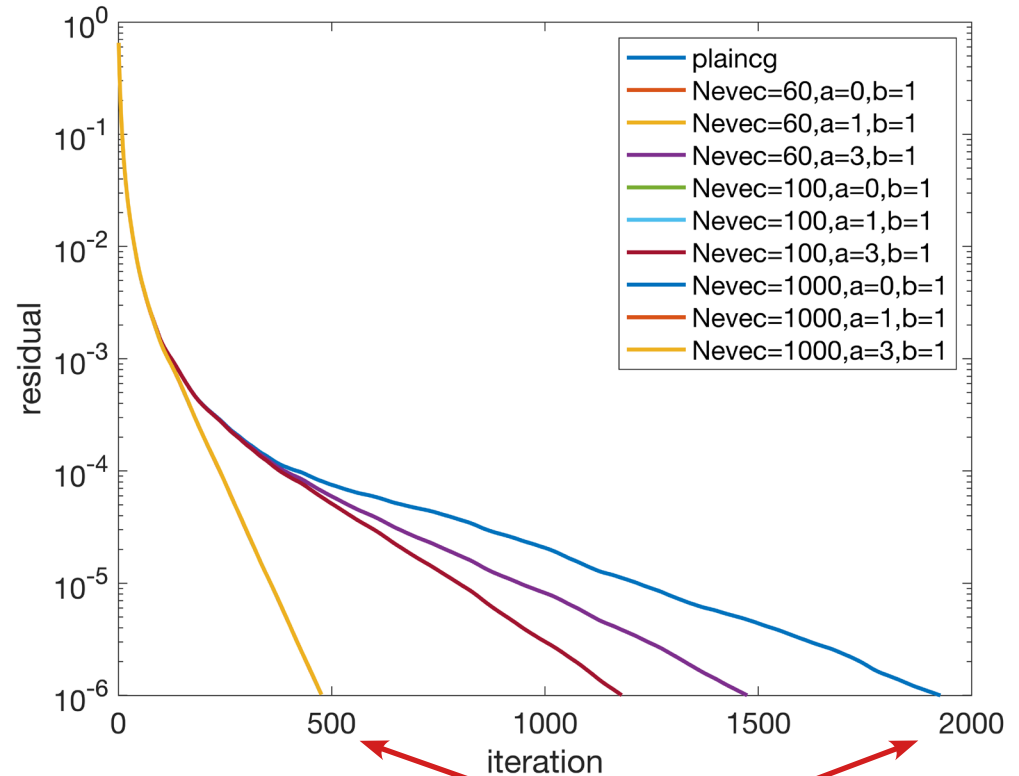
$$k = k + 1$$

end

- Working with $A = D^\dagger D$
- Precondition with fine lattice eigenvectors

$$M^{-1} = a + b \left(1 + \sum_i^N |v_{h,i}\rangle \langle v_{h,i}| \left(\frac{1}{\lambda_{h,i}} - 1 \right) \right)$$

Results for fine lattice CG solve with mass = m_s



4x fewer iterations even with quark mass of m_s

Coarse Eigenvector Preconditioning

Want to change from fine eigenvector $v_{h,i}$ preconditioner

$$M^{-1} = a + b\left(1 + \sum_i^N |v_{h,i}\rangle\langle v_{h,i}| \left(\frac{1}{\lambda_{h,i}} - 1\right)\right)$$

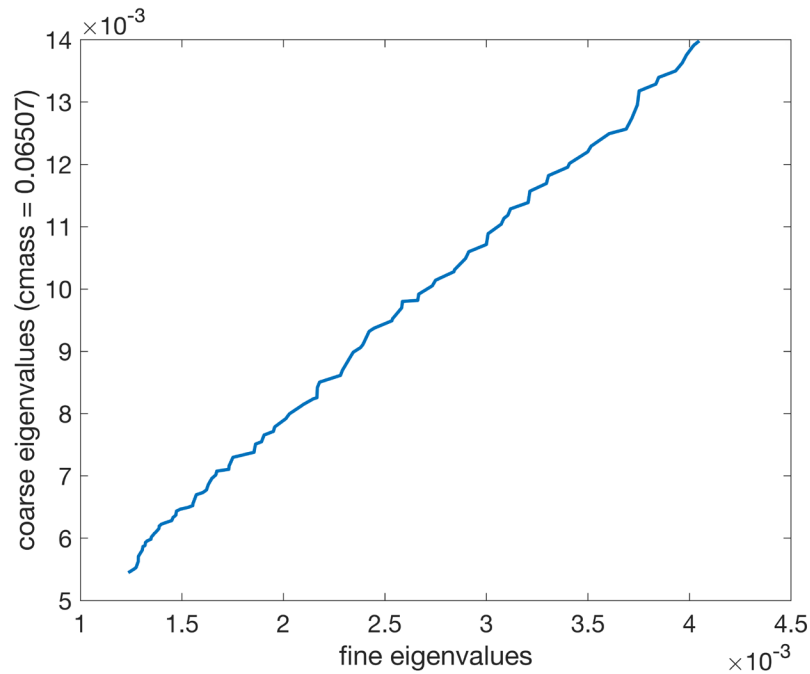
to one based on coarse eigenvectors $v_{2h,i}$ from blocked coarse lattice

$$M^{-1} = 1 + bPI\left(\sum_i^N |\psi_{2h,i}\rangle\langle\psi_{2h,i}| \frac{1}{\lambda_{2h,i}}\right)RP$$

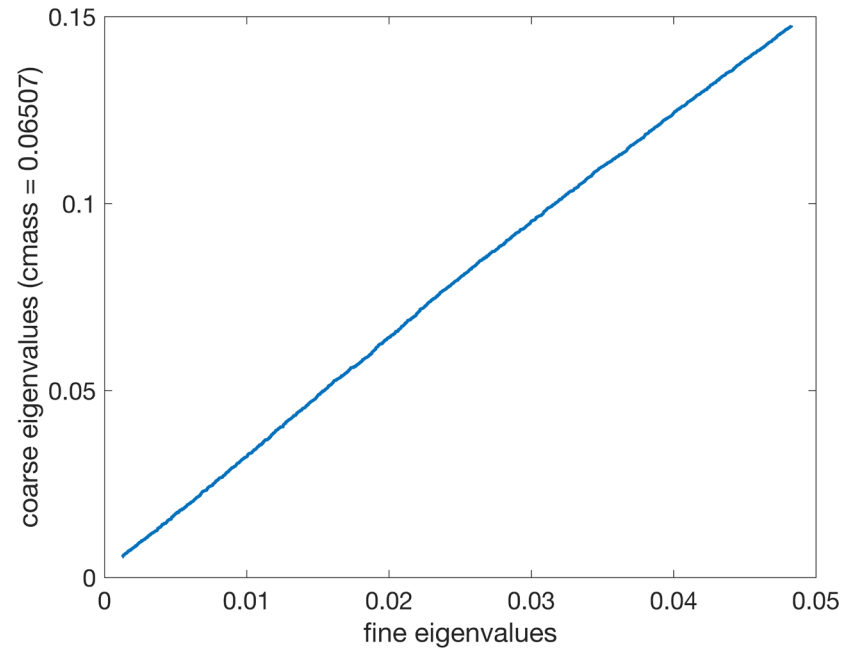
Work in Landau gauge

- * **R** is a restriction operator, **I** is an interpolation or prolongation operator
- * **P** is a smoother or filter

Comparing Eigenvalue Spectrum



First 100 eigenvalues

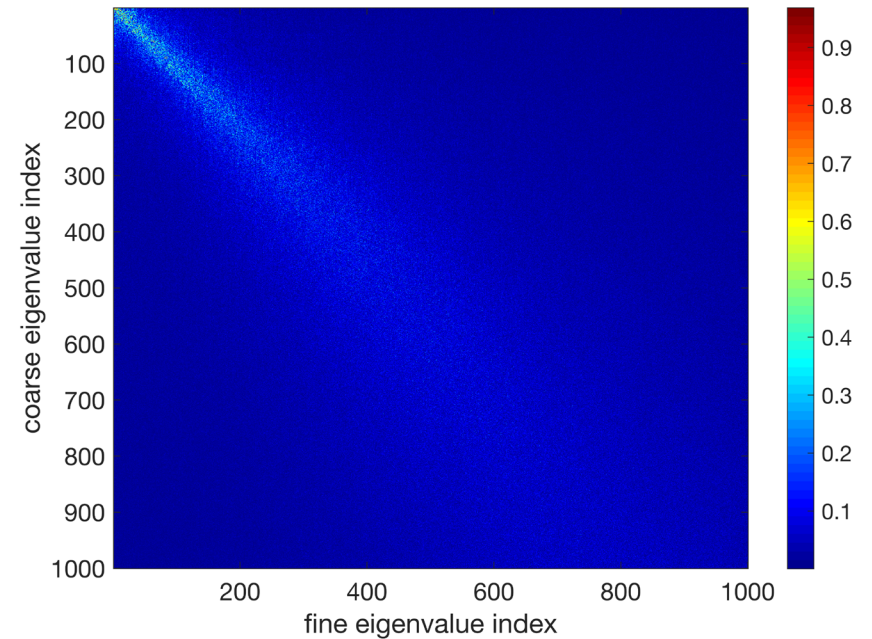
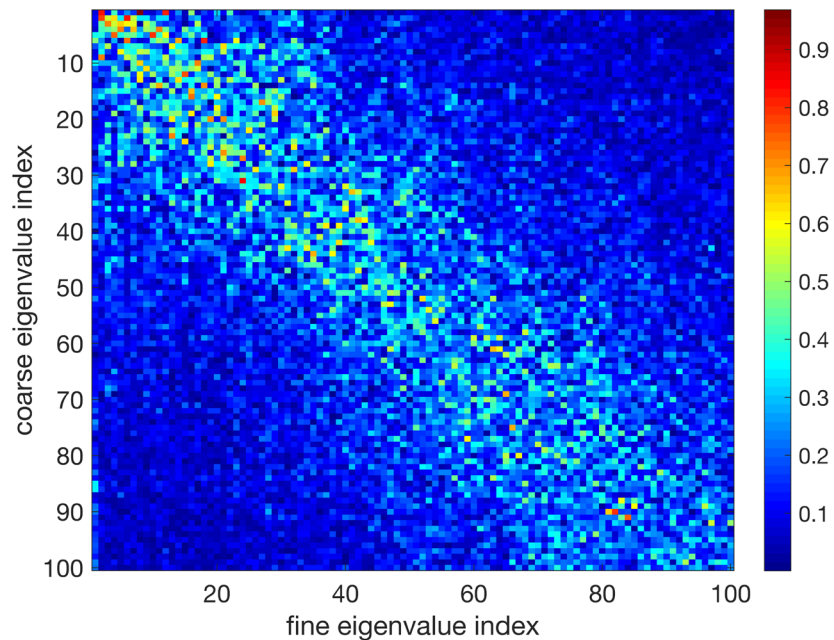


First 1000 eigenvalues

Eigenvalue densities are very similar and eigenvalues differ by $3\times$

Comparing Eigenvectors

Calculate magnitude of inner product of fine eigenvector $v_{h,i}$ with interpolated blocked coarse eigenvector $Iv_{2h,j}$



Individual eigenvectors are not in one-to-one correspondence.

Compare Low Mode Subspace

Defining the coarse and fine lattice inverses as

$$S_c = I \sum_i^N |\psi_{2h,i}\rangle \langle \psi_{2h,i}| \frac{1}{\lambda_{2h,i}} R$$

$$S_f = \sum_i^N |\psi_{h,i}\rangle \langle \psi_{h,i}| \frac{1}{\lambda_{h,i}}$$

we calculate

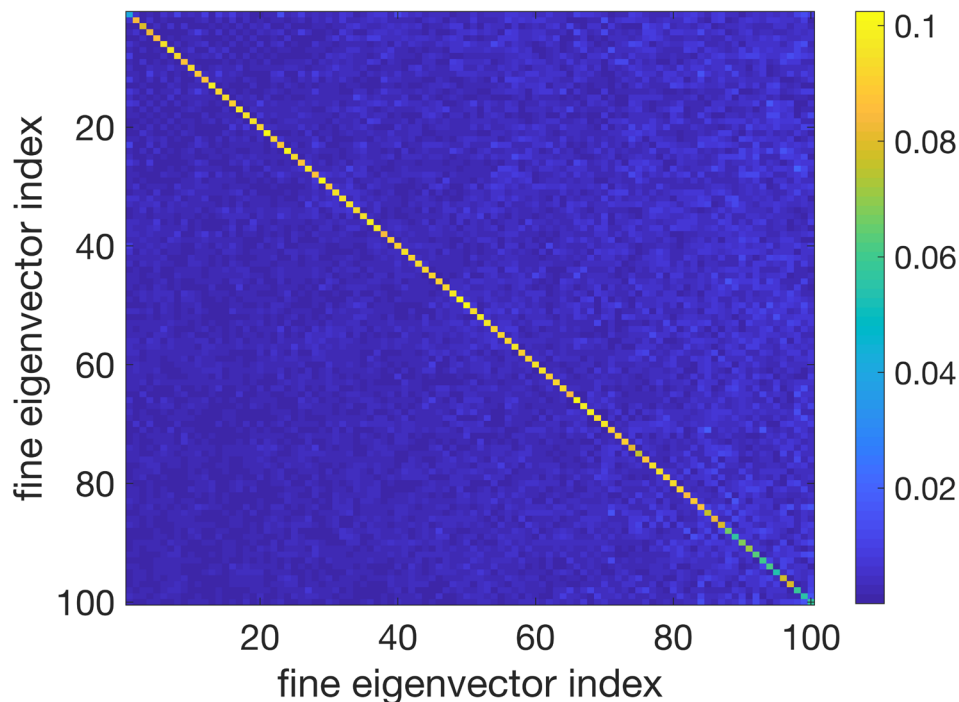
$$X = I \sum_i^N |\psi_{2h,i}\rangle \langle \psi_{2h,i}| \frac{1}{\lambda_{2h,i}} R \sum_j^N |\psi_{h,j}\rangle \langle \psi_{h,j}| \lambda_{h,j}$$

If the subspaces were identical and complete, we would find $X =$ the identity

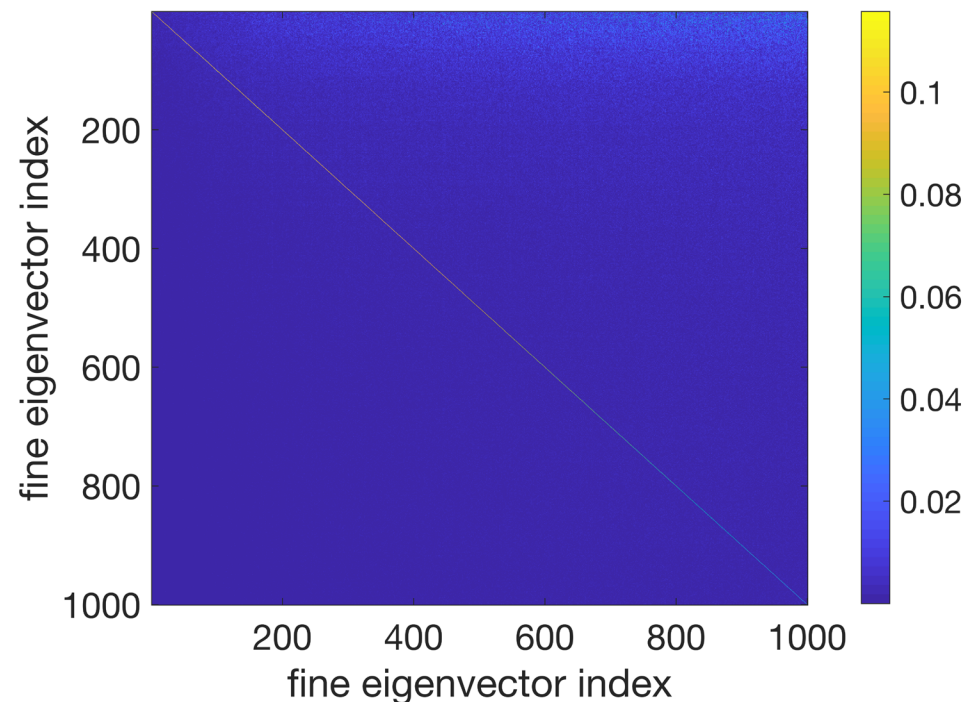
Comparison of Low Mode Subspaces

$$X = I \sum_i^N |\psi_{2h,i}\rangle \langle \psi_{2h,i}| \frac{1}{\lambda_{2h,i}} R \sum_j^N |\psi_{h,j}\rangle \langle \psi_{h,j}| \lambda_{h,j}$$

Plots of X show it is primarily diagonal



First 100 eigenvalues

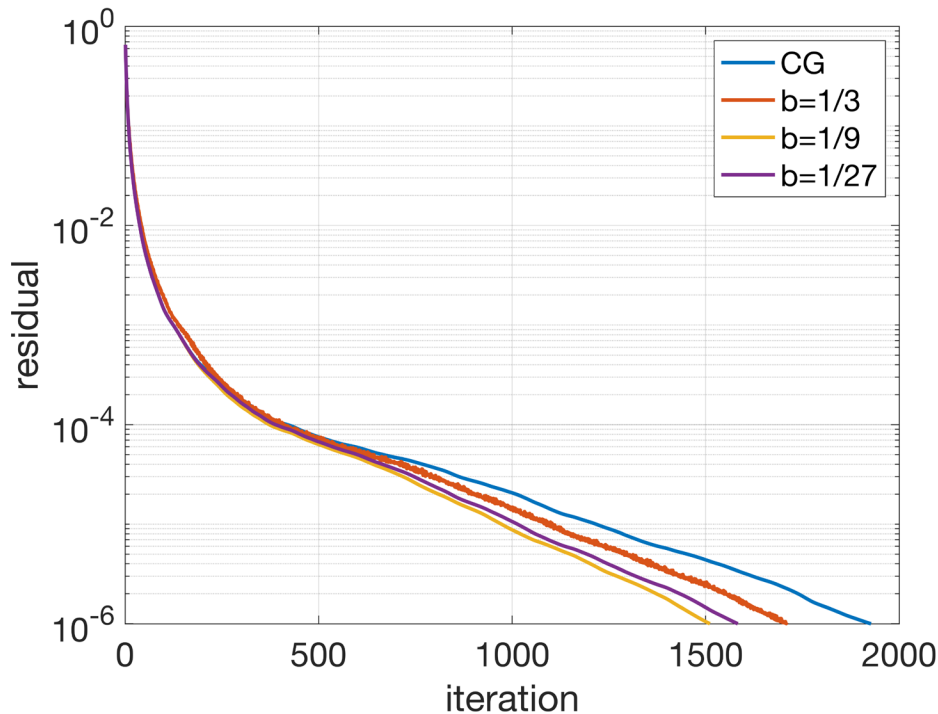


First 1000 eigenvalues

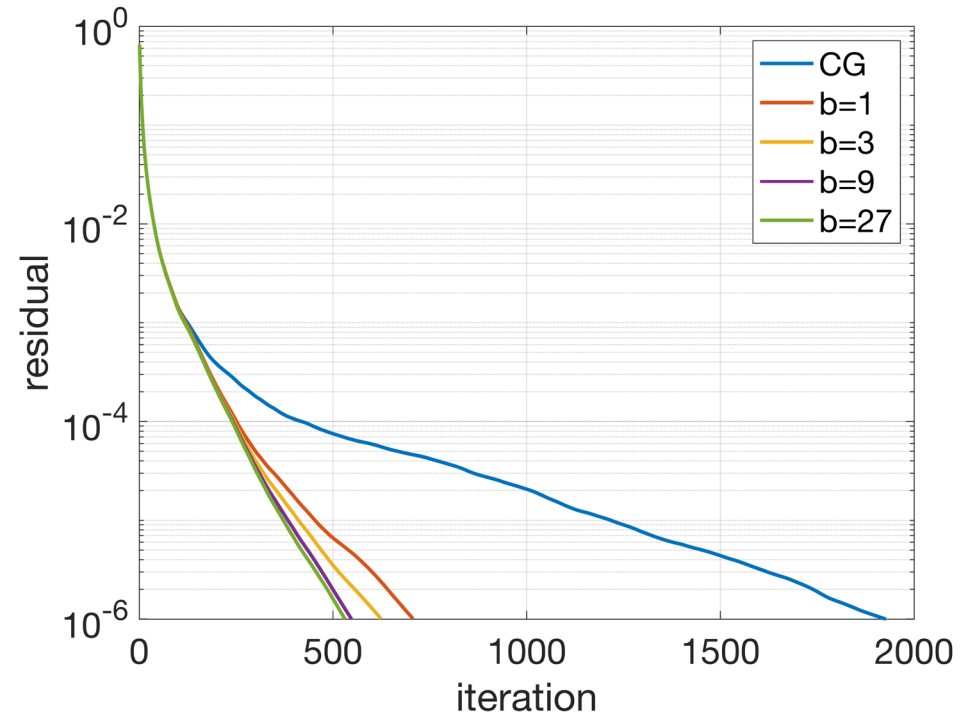
Individual eigenvectors are not in one-to-one correspondence, but the coarse low mode approximation to the inverse is quite similar to the high mode one.

Coarse Preconditioner

$$M^{-1} = 1 + bPI\left(\sum_i^N |\psi_{2h,i}\rangle\langle\psi_{2h,i}| \frac{1}{\lambda_{2h,i}}\right)RP$$



1000 coarse eigenvectors
No filter, $P = 1$
Some improvement



1000 coarse eigenvectors
Use fine eigenvectors for filter
4x improvement

$$P = \sum_i^N |\psi_{h,i}\rangle\langle\psi_{h,i}|$$

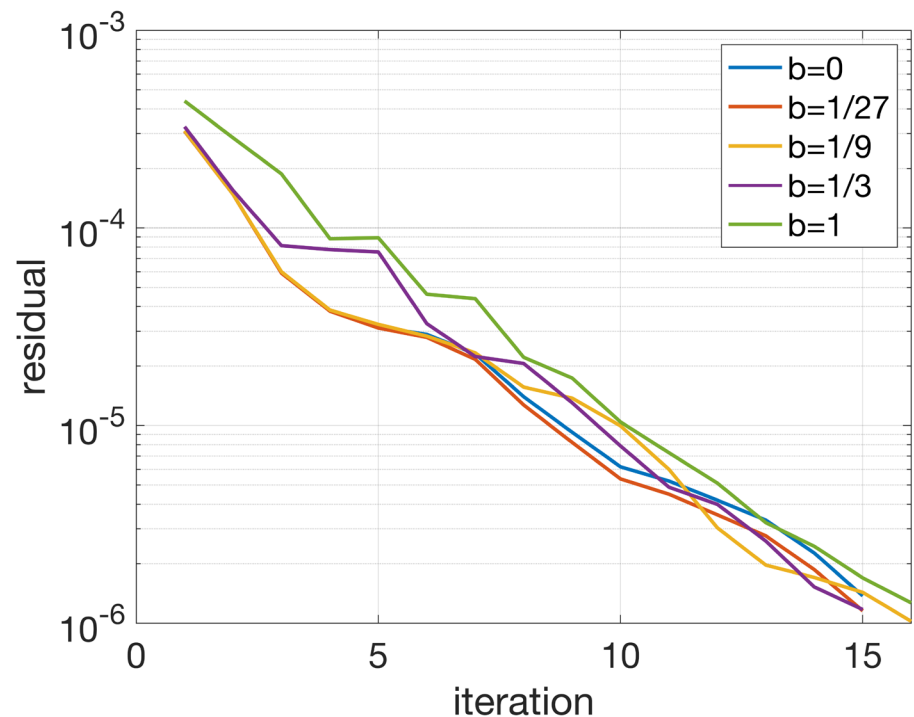
Exploring Various Filters/Smoothers

- Have tried Jacobi solver and Chebyshev polynomial filter as smoothers.
- Chebyshev polynomials work better in the smoother than Jacobi
- The total count of $D^\dagger D$ only modestly below the unpreconditioned case

Use these 5 steps as the preconditioner, setting $u = 0$ in step 1.

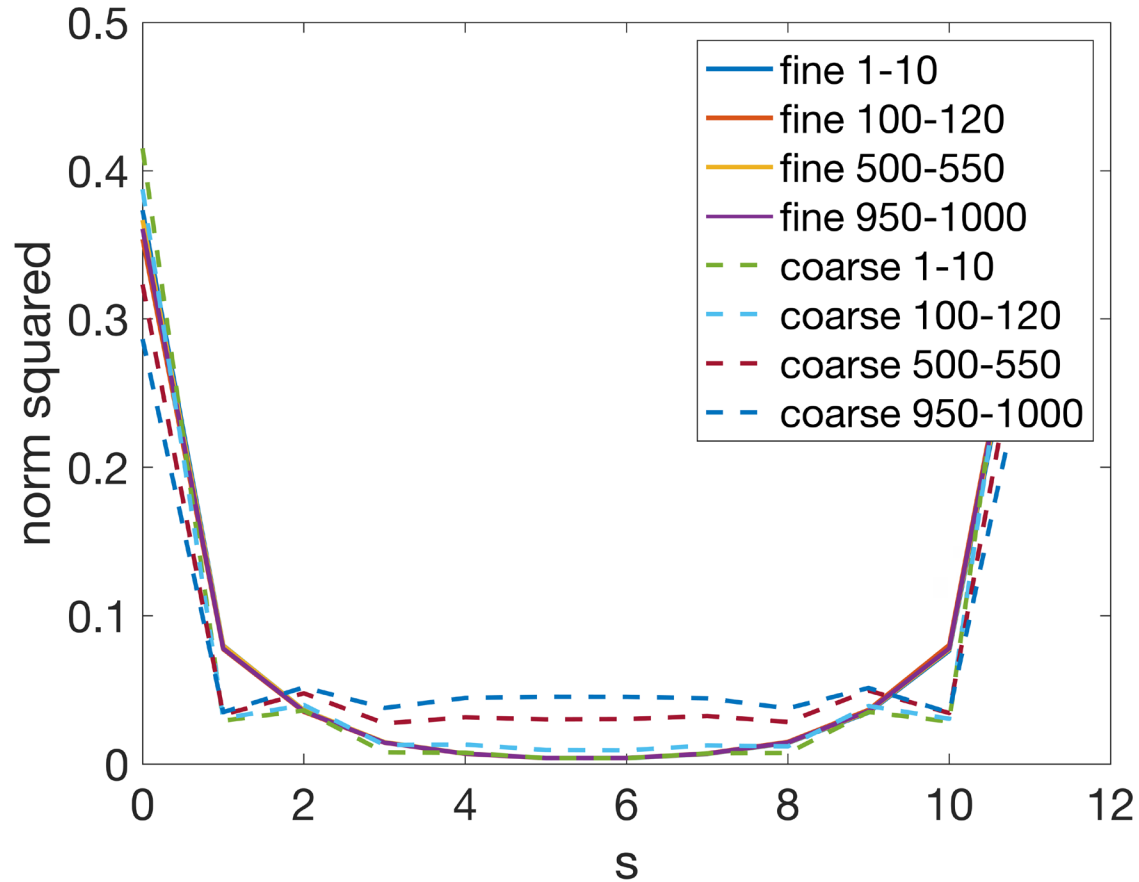
Jacobi and Chebyshev used for P

- 1) $u = P(r, u)$ (pre-smoothing)
- 2) $d = R(A_h u - r)$
- 3) $v = A_{2h}^{-1} d$
- 4) $u = u - b I v$
- 5) $u = P(r, u)$ (post-smoothing)



100 $D^\dagger D$ applications used in inner solver so total $D^\dagger D$ count is about 1500, similar to the 2000 iterations for CG.

Coarse and fine eigenmodes differ in 5th dimension



Summary

- APE-smearing style blocking produces a coarse lattice with $2\times$ the lattice spacing
- Low mode subspace on coarse lattice can be prolonged to a good approximation to low mode subspace on fine lattice.
- Using low mode subspace as a preconditioner for CG increases convergence rate - better results possible for m_{ud} than for m_s , as used here.
- Have tried various high-mode filters/smothers without much improvement.
- 5d structure of fine and coarse modes differs - perhaps a direction for improvement