

Nucleon Axial Form Factor from Domain Wall on HISQ

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On behalf of the CalLat collaboration

UC Berkeley/LBNL

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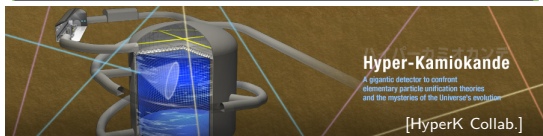
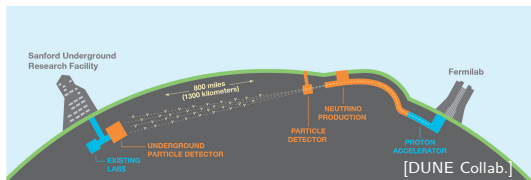
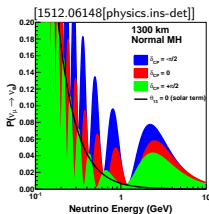
38th International Symposium on Lattice Field Theory

CalLat Team Members

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- ▶ Jason Chang
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- ▶ Jinchen He
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- ▶ Colin Morningstar
- ▶ Amy Nicholson
- ▶ Jiqun Tu
- ▶ Pavlos Vranas
- ▶ André Walker-Loud



Motivation from Neutrino Oscillation Expt

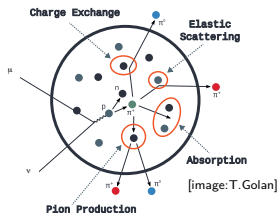
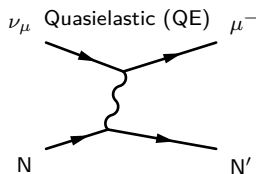
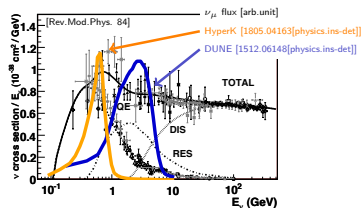


Flagship neutrino oscillation experiments
(DUNE, HyperK)

$V - A$ interaction \implies weak current

Secondary beam $\implies \nu$ flux probes range of E_ν

Axial Form Factor



Weak current \implies large nuclear targets

Multiple interaction regimes in play

Ideal factorization: Nucleon / Nuclear

Quasielastic scattering:

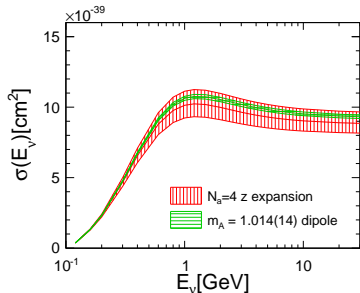
theoretically simple process

dominates low energy interactions

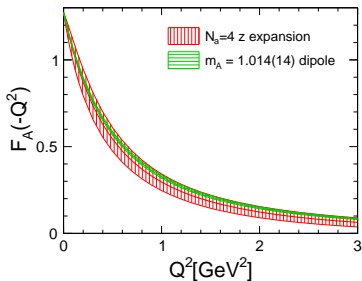
primary signal measurement

QE F_A from z Expansion

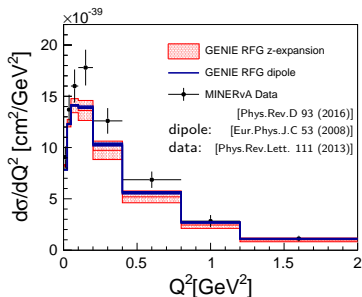
- ▶ Dipole $F_A(Q^2) = g_A/(1 + Q^2/m_A^2)^2$
strict Q^2 shape,
inconsistent w/ QCD
- ▶ Dipole FF ansatz significantly
underestimates FF uncertainty



- ▶ Nucl. xsec uncertainty from FF
same size as data-MC tensions
- ▶ Source of tensions unclear
btw. nucleon/nuclear



- ▶ Model-independent parameterization
Order of mag. larger $\delta\sigma$



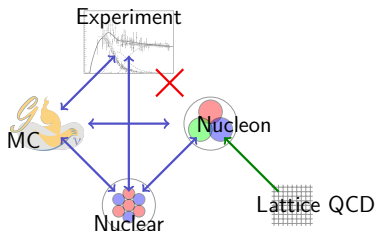
LQCD as Disruptive Technology

Ideal: Modern high stats ν -D₂ scattering bubble chamber expt

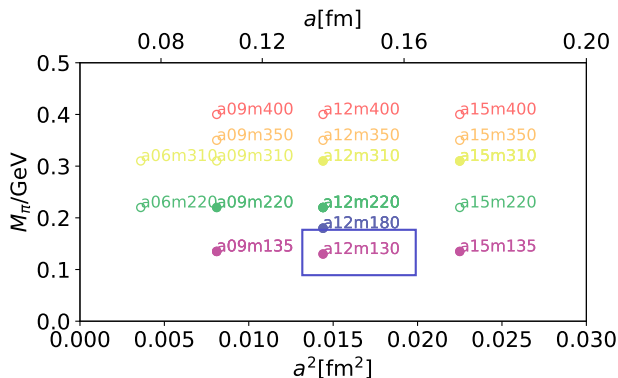
Some community push, safety concerns

⇒ LQCD as a alternative/complement to expt

- ✓ No nuclear effects
- ✓ Realistic uncertainty estimates
- ✓ Systematically improvable
- ✓ Computers are (relatively) inexpensive



Callat Computation



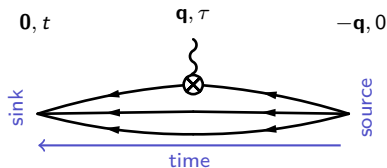
► Open symbol: g_A computation

► Filled symbol: existing $g_A(Q^2)$ data

Multiple volumes on many ensembles

⇒ Only a12m130 in this talk

Setup



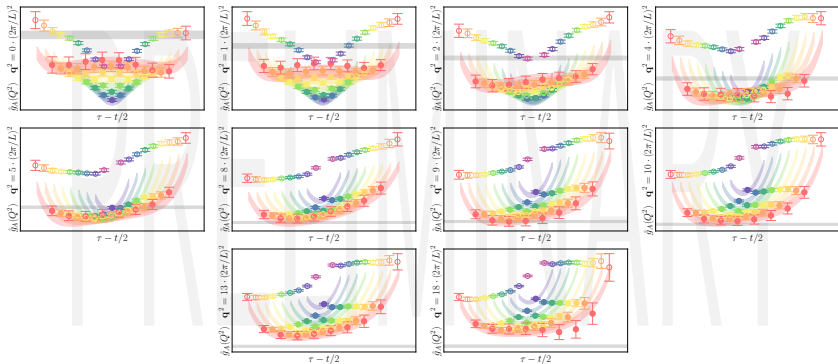
$$\mathcal{R}_{\mathcal{A}_z}(t, \tau, \mathbf{q}) = \frac{C_{\mathcal{A}_z}^{3\text{pt}}(t, \tau, \mathbf{q})}{\sqrt{C^{2\text{pt}}(t - \tau, \mathbf{0}) C^{2\text{pt}}(\tau, \mathbf{q})}} \sqrt{\frac{C^{2\text{pt}}(\tau, \mathbf{0})}{C^{2\text{pt}}(t, \mathbf{0})} \frac{C^{2\text{pt}}(t - \tau, \mathbf{q})}{C^{2\text{pt}}(t, \mathbf{q})}}$$
$$\xrightarrow{t - \tau, \tau \rightarrow \infty} \frac{1}{\sqrt{2E_{\mathbf{q}}(E_{\mathbf{q}} + M)}} \left[-\frac{q_z^2}{2M} \hat{g}_P(Q^2) + (E_{\mathbf{q}} + M) \hat{g}_A(Q^2) \right]$$

$$\mathcal{A}_z \text{ w/ } q_z = 0 \implies \mathcal{R}_{\mathcal{A}_z}(t, \tau, \mathbf{q}) \rightarrow \sqrt{\frac{E_{\mathbf{q}} + M}{2E_{\mathbf{q}}}} \hat{g}_A(Q^2)$$

\implies No induced pseudoscalar

\implies Simplified analysis of $\hat{g}_A(Q^2)$

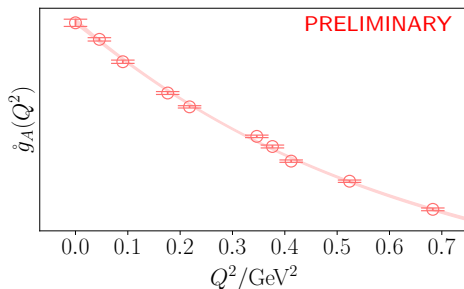
Three-Point Function Ratios



Global 3-state exponential fit \times 10 momenta w/ $|q_i| \leq 3 \cdot (2\pi/L)$

- ▶ colors: $t/a \in \{3, \dots, 12\}$ ▶ gray band: fit $\hat{g}_A(Q^2)$ ▶ All $q_z = 0$
- ▶ horizontal: τ , centered about midpoint
- ▶ vertical: $\sqrt{2E_{\mathbf{q}}/(E_{\mathbf{q}} + M)} \mathcal{R}_{\mathcal{A}_z}(t_{\text{sep}}, \tau, \mathbf{q}) \rightarrow \hat{g}_A(Q^2)$
- ▶ Minimum time separation $\tau/a, (t - \tau)/a \geq 2$

Axial Form Factor Fit



Plotted: 4-parameter z expansion + 4 sum rules

$\delta g_A/g_A$: 0.6%; $\delta r_A^2/r_A^2$: 13%

Cf. $\delta r_A^2/r_A^2 \sim 48\%$ from D_2 [Phys.Rev.D 93 (2016)]

Results yield $M_{A,\text{dipole}} \gtrsim 1.0 \text{ GeV}$

Prelim dipole fit underestimates uncertainty \times few

$$\implies r_A^2 = -(6/g_A)dF_A/dQ^2 \Big|_{Q^2=0}; \quad r_{A,\text{dipole}}^2 = (1 \text{ GeV}^2/M_{A,\text{dipole}}^2) \times 0.466\text{fm}^2$$

Outlook

- ▶ Lattice QCD is the most practical approach to determine weak matrix elements for neutrino oscillation
- ▶ CalLat data for $g_A(Q^2)$ on several ensembles
- ▶ Large-scale global fits to several momenta @ M_π physical
- ▶ Good agreement across the board
- ▶ Future: more momenta, $p_z \neq 0$, other currents & tests of PPD...

Thanks for your attention!