

Thermal phase structure of dimensionally reduced super-Yang–Mills

David Schaich (U. Liverpool)



Lattice 2021, July 29

[arXiv:2003.01298](https://arxiv.org/abs/2003.01298) and more to come with Raghav G. Jha and Anosh Joseph

Overview

Dimensional reduction simplifies lattice supersymmetry, retaining interesting dynamics and holographic dualities

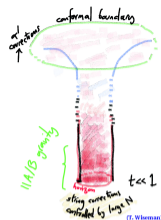
Motivations for lattice susy & dimensional reduction

Berenstein–Maldacena–Nastase (BMN) [hep-th/0202021]
deformation of super-Yang–Mills quantum mechanics

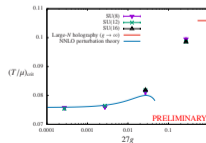
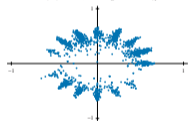
BMN phase diagram results & future plans

These slides: davidshaich.net/talks/2107Lattice.pdf

Feel free to follow up on Gather / Slack or directly!



BMN SU(16) $N_c = 16$ $g = 0.01$ $T/\mu = 0.108$



General motivations: Lattice supersymmetry

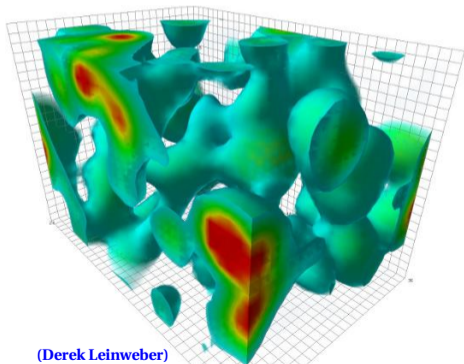
Lattice field theory promises first-principles predictions
for strongly coupled supersymmetric QFTs

BSM



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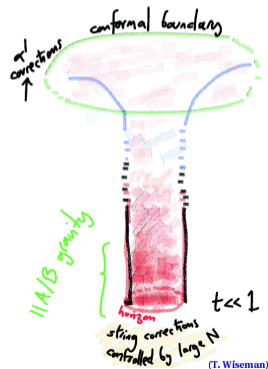
QFT



(Derek Leinweber)

BMN phase diagram

Holography



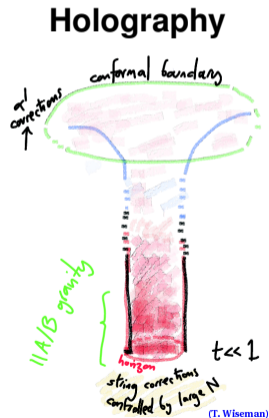
(T. Wiseman)

General motivations: Lattice supersymmetry

Lattice field theory promises first-principles predictions
for strongly coupled supersymmetric QFTs

Today

Focus on 0+1 dimension as testbed for holography



Further motivations: Super-Yang–Mills QM

Supersymmetry is a space-time symmetry, ($I = 1, \dots, \mathcal{N}$)
adding spinor generators Q_α^I and $\bar{Q}_{\dot{\alpha}}^I$ to translations, rotations, boosts

$\{Q_\alpha^I, \bar{Q}_{\dot{\alpha}}^J\} = 2\delta^{IJ}\sigma_{\alpha\dot{\alpha}}^\mu P_\mu$ broken in discrete space-time
→ relevant susy-violating operators



Further motivations: Super-Yang–Mills QM

Supersymmetry broken in discrete space-time \longrightarrow fine-tuning challenges

Simplify via dimensional reduction \longrightarrow significant recent progress

Compactifying all spatial dimensions,

4d $\mathcal{N} = 4$ SU(N) super-Yang–Mills \longrightarrow quantum mechanics of $N \times N$ matrices

[Goldstone, Hoppe, de Wit, Nicolai, \dots , 1980s]

[Holographic duality conjecture by Banks–Fischler–Shenker–Susskind, [hep-th/9610043](https://arxiv.org/abs/hep-th/9610043)]

$$\mathbf{S}_{\text{BFSS}} = \frac{N}{4\lambda} \int d\tau \text{Tr} \left[- (D_\tau \mathbf{X}_i)^2 + \Psi_\alpha^T \gamma_{\alpha\beta}^\tau D_\tau \Psi_\beta - \frac{1}{2} \sum_{i < j} [\mathbf{X}_i, \mathbf{X}_j]^2 + \frac{1}{\sqrt{2}} \Psi_\alpha^T \gamma_{\alpha\beta}^i [\mathbf{X}_i, \Psi_\beta] \right]$$

16 fermions Ψ_α

9 scalars $\mathbf{X}_i \longleftrightarrow$ SO(9) global symmetry

Supersymmetric mass deformation

Berenstein–Maldacena–Nastase, [hep-th/0202021](#)

Split the 9 scalars $X_i \rightarrow 3X_I$ and $6X_A$ with different masses
 \implies break $SO(9) \rightarrow SO(6) \times SO(3)$ symmetry

Dimensionful mass parameter $\mu \rightarrow$ dimensionless coupling $g \equiv \lambda/\mu^3$

$$\mathcal{S}_{\text{BMN}} = \mathcal{S}_{\text{BFSS}} - \frac{N}{4\lambda} \int d\tau \text{Tr} \left[\left(\frac{\mu}{3} X_I \right)^2 + \left(\frac{\mu}{6} X_A \right)^2 + \frac{\mu}{4} \psi_\alpha^T \gamma_{\alpha\beta}^{123} \psi_\beta - \frac{\sqrt{2}\mu}{3} \epsilon_{\text{IJK}} X_I X_J X_K \right]$$

Lift flat directions while preserving all 16 supersymmetries

Simple discretization provided in [public parallel code](#)

Finite-temperature phase diagram and holography

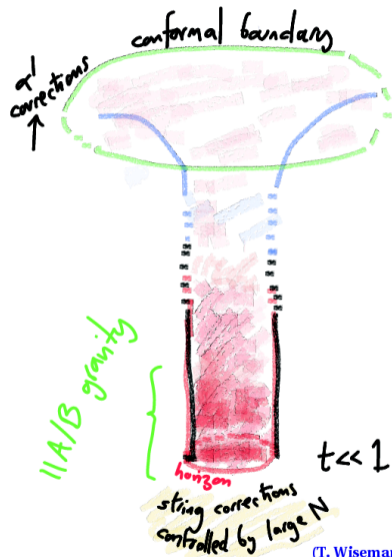
Holographic duality conjecture

Thermodynamics of maximal SYM
at large N and strong coupling



Black holes in string theory

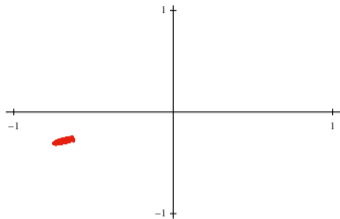
Numerical dual-supergravity [arXiv:1411.5541]
→ first-order BMN deconfinement transition



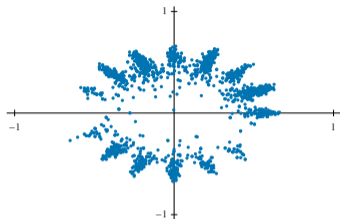
(T. Wiseman)

Deconfinement signalled by Polyakov loop

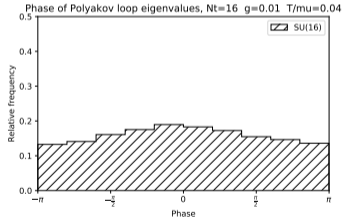
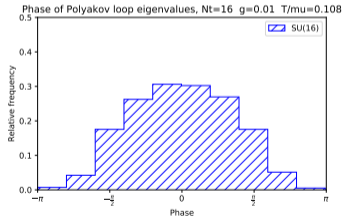
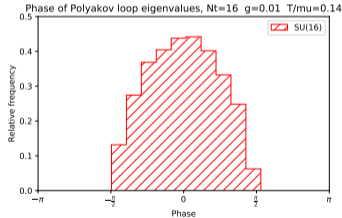
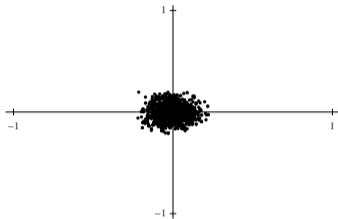
BMN SU(16) $N_\tau = 16$ $g = 0.01$ $T/\mu = 0.14$



BMN SU(16) $N_\tau = 16$ $g = 0.01$ $T/\mu = 0.108$



BMN SU(16) $N_\tau = 16$ $g = 0.01$ $T/\mu = 0.04$



Weak-coupling perturbation theory \longrightarrow first-order BMN deconfinement transition

The lattice phase diagram game

Gauge groups $SU(8)$, $SU(12)$, $SU(16)$

×

Lattice sizes $N_T = 8, 16, 24$

×

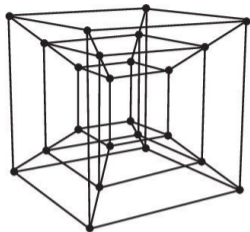
Dimensionless couplings $g = \lambda/\mu^3 = 0.00001, 0.0001, 0.001, 0.01$

×

Scan dimensionless temperatures T/μ across transition

= 333 ensembles

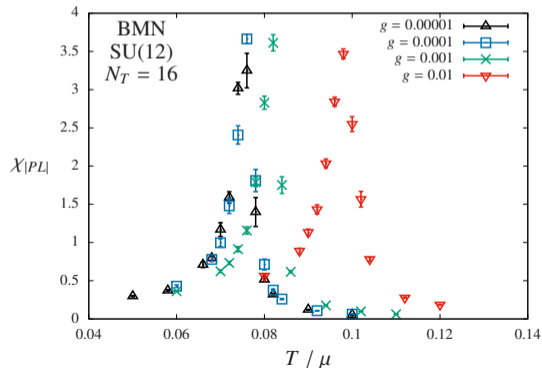
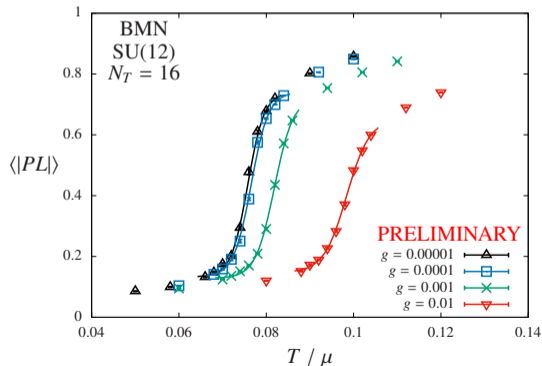
[up to 50,000 MD time units per ensemble]



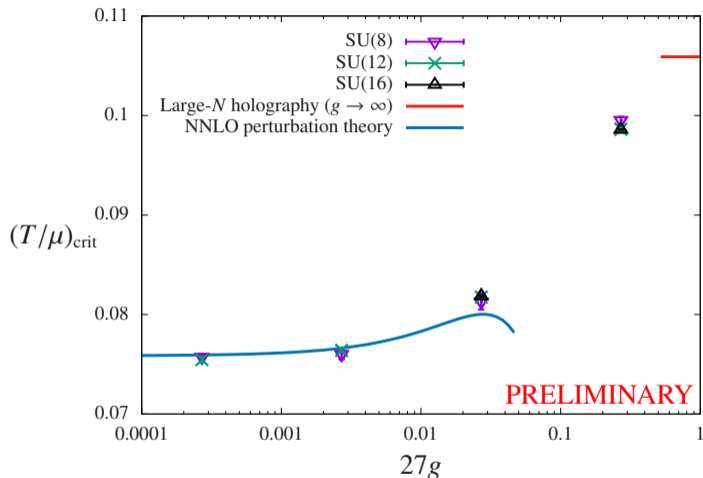
Polyakov loop and susceptibility

Extract critical T/μ by fitting Polyakov loop to simple sigmoid function

Results match peaks in Polyakov loop susceptibility



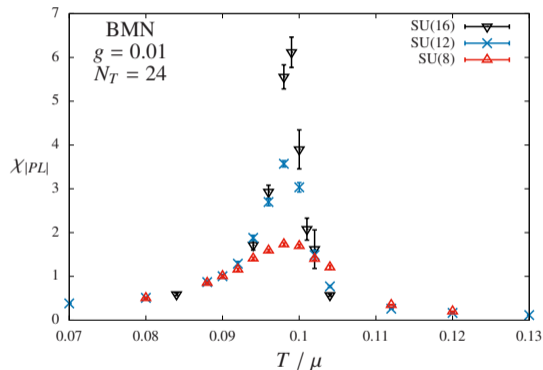
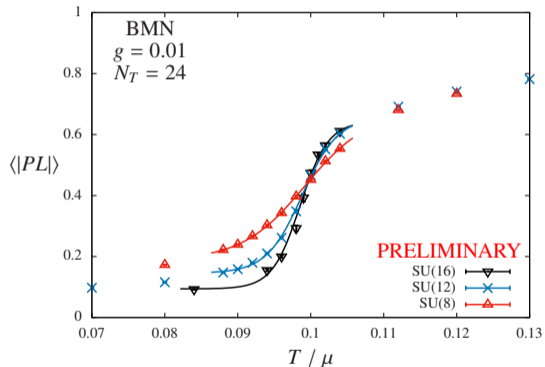
Today's main result: Phase diagram



Reproduce weak-coupling pert. theory, approach strong-coupling holography

Dependence on $SU(N)$ gauge group

Larger $N \rightarrow$ sharper transition at same critical T/μ



Susceptibility peaks consistent with expected first-order $\sim N^2$ scaling

Recap and outlook

Dimensional reduction simplifies lattice supersymmetry, retaining interesting dynamics and holographic dualities

BMN quantum mechanics is interesting target

with [public code](#) available

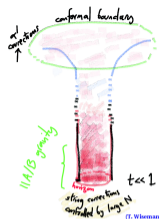
Phase diagram results agree with weak-coupling pert. theory, approach strong-coupling holography

More investigations underway:

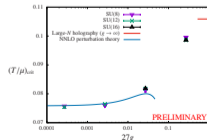
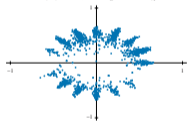
Gauge invariance \rightarrow global symmetry [[1802.02985](#)]

Remove fermions \rightarrow bosonic BMN

More observables & predictions for holography



BMN SU(16) $N_c = 16$ $g = 0.01$ $T/\mu = 0.108$



Thanks for your attention!

Collaborators

Raghav G. Jha, Anosh Joseph

Funding and computing resources



UK Research
and Innovation



Backup: Breakdown of Leibniz rule on the lattice

$$\{Q_\alpha, \bar{Q}_{\dot{\alpha}}\} = 2\sigma_{\alpha\dot{\alpha}}^\mu P_\mu = 2i\sigma_{\alpha\dot{\alpha}}^\mu \partial_\mu \text{ is problematic}$$

$$\implies \text{try finite difference } \partial\phi(x) \longrightarrow \Delta\phi(x) = \frac{1}{a} [\phi(x+a) - \phi(x)]$$

Crucial difference between ∂ and Δ

$$\begin{aligned}\Delta[\phi\eta] &= a^{-1} [\phi(x+a)\eta(x+a) - \phi(x)\eta(x)] \\ &= [\Delta\phi]\eta + \phi\Delta\eta + a[\Delta\phi]\Delta\eta\end{aligned}$$

Full supersymmetry requires Leibniz rule $\partial[\phi\eta] = [\partial\phi]\eta + \phi\partial\eta$

only recovered in $a \rightarrow 0$ continuum limit for any local finite difference

Backup: Breakdown of Leibniz rule on the lattice

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Supersymmetry vs. locality 'no-go' theorems

by Kato–Sakamoto–So [[arXiv:0803.3121](https://arxiv.org/abs/0803.3121)] and Bergner [[arXiv:0909.4791](https://arxiv.org/abs/0909.4791)]

Complicated constructions to balance locality vs. supersymmetry

Non-ultralocal product operator \rightarrow lattice Leibniz rule but not gauge invariance

D'Adda–Kawamoto–Saito, [arXiv:1706.02615](https://arxiv.org/abs/1706.02615)

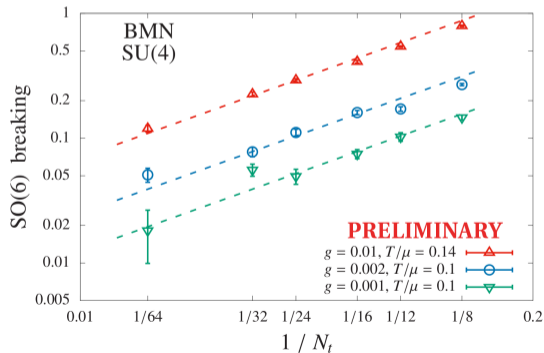
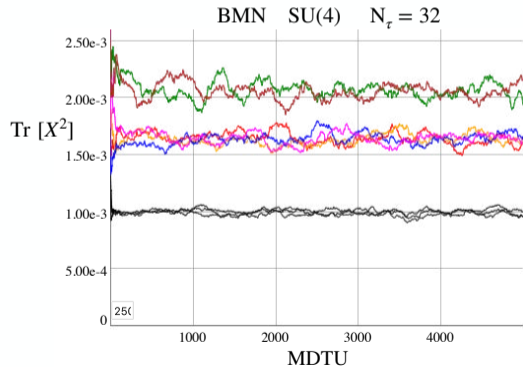
Cyclic Leibniz rule \rightarrow partial lattice supersymmetry but only $(0+1)d$ QM so far

Kadoh–Kamei–So, [arXiv:1904.09275](https://arxiv.org/abs/1904.09275)

Backup: Discretization artifacts

Discretization artifacts break $SO(6)$ part of BMN $SO(6) \times SO(3)$ symmetry

Breaking increases with g but vanishes in $N_T \rightarrow \infty$ continuum limit



Backup: Large- N limit

