

# Quantum Counter Terms for Lattice Field Theory on Curved Manifolds (Lattice 2021)

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7/27/2021

# Background

We seek to probe the non-perturbative regime of  $\phi^4$  theory on a discretized  $S^2$  lattice close to the strong coupling Wilson-Fisher IR fixed point.

Expanding on the results of Brower et al. 1803.08512 / PRD 98, 014502 (2018) which studies the same model with perturbative counterterms [1].

# Background

Start with continuum  $\phi^4$  action defined on a smooth Riemann manifold  $\mathcal{M}$ .

$$S = \int_{\mathcal{M}} d^d x \sqrt{g} \left( \frac{1}{2} g^{\mu\nu} \partial_\mu \phi(x) \partial_\nu \phi(x) + \frac{1}{2} m^2 \phi^2(x) + \lambda \phi^4(x) \right)$$

Introduce a simplicial complex to discretize the manifold. The lattice action is

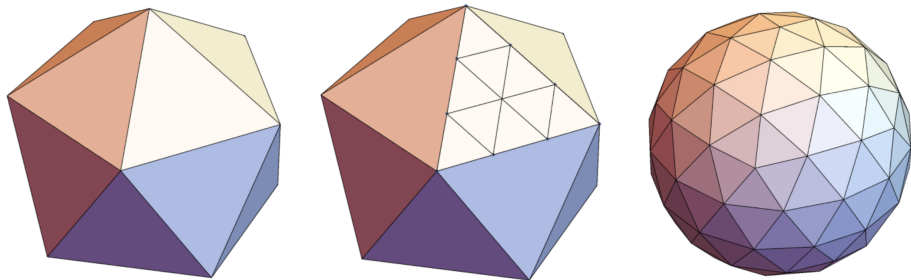
$$S = \frac{1}{2} \sum_{\langle xy \rangle} \frac{V_{xy}}{\ell_{xy}^2} (\phi_x - \phi_y)^2 + \sum_x \sqrt{g_x} \left( \frac{1}{2} m^2 \phi_x^2 + \lambda \phi_x^4 \right)$$

Site/link prefactors are determined via finite element method [1].

$\sqrt{g_x}$  is the normalized area associated with site  $x$ .

# Background

Discretize  $\mathbb{S}^2$  using a refined icosahedron.

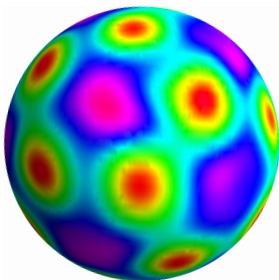


Approach continuum limit by increasing refinement (refinement of 3 is shown).

# Spherical Symmetry Breaking

Quantum theory has UV divergences due to quantum fluctuations.

Spherical symmetry of UV and IR operators is lost in the continuum limit.

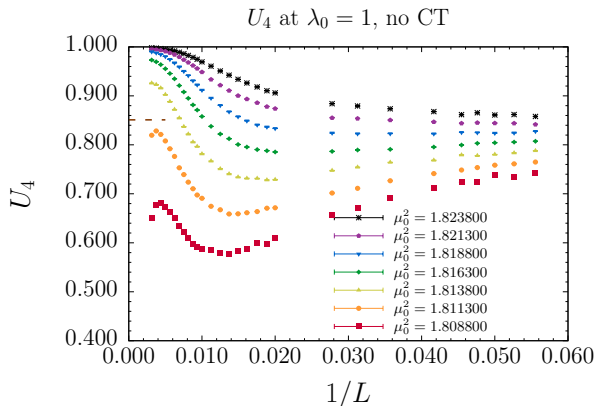


$\langle \phi_x^2 \rangle$  from Monte Carlo simulation showing breaking of spherical symmetry near the critical surface.

# Unstable Critical Surface Without Counterterms

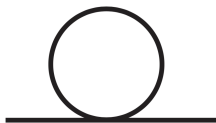
Critical surface is not well-defined as the continuum limit is approached.

We measure  $U_4 = \frac{3}{2} \left( 1 - \frac{\langle m^4 \rangle}{3\langle m^2 \rangle^2} \right)$  where  $m = \sum_x \sqrt{g_x} \phi_x$



## Perturbative Counter Terms

In 2d, 1-loop diagram is logarithmically UV divergent. Higher order diagrams are finite.



Performing the perturbative expansion in  $\lambda$ , renormalization requires the addition of *local* mass counterterms. To first order,

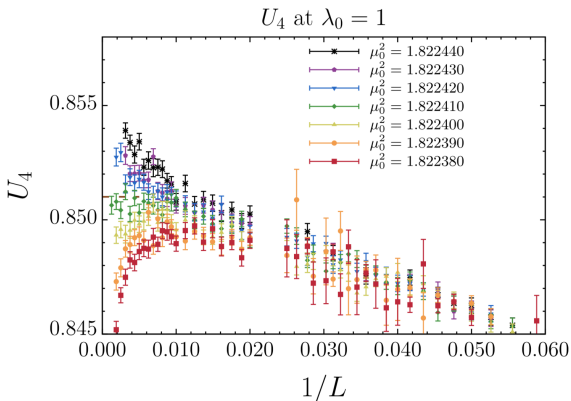
$$\delta m_x^2 = 6Q\lambda \log(\sqrt{g_x})$$

where  $Q = \sqrt{3}/8\pi$  and  $\sqrt{g_x}$  is the normalized area associated with site  $x$ .

$$S \rightarrow S + \frac{1}{2} \sum_x \sqrt{g_x} \delta m_x^2 \phi_x^2$$

# Restoration of Critical Surface

Spherical symmetry and critical surface are restored at  $\lambda \leq 1$  after adding local perturbative counterterms.

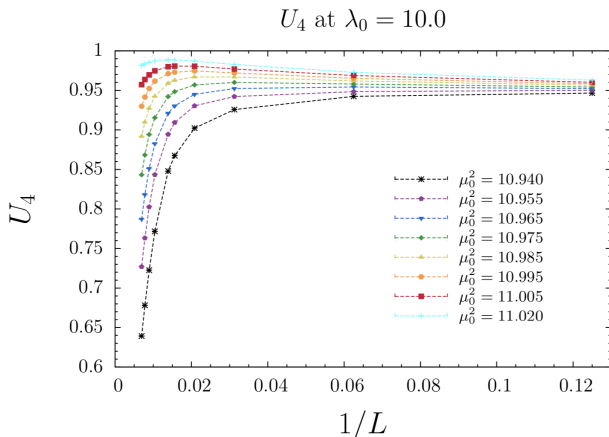


Binder cumulant approaches the correct continuum value of  $U_4^* \simeq 0.851$



# Critical Surface Near Wilson-Fisher Fixed Point

Spherical symmetry remains broken at  $\lambda > 1$ . (Gasbarro, 2019 [2])

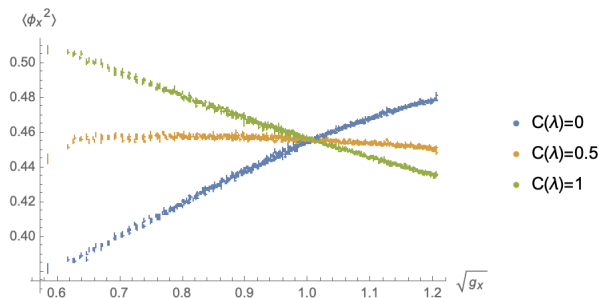


Unable to approach Wilson-Fisher fixed point ( $10 \lesssim \lambda_{WF} \lesssim 20$ )

# Non-Perturbative Counter Terms

Conjecture that symmetry can be restored at strong coupling by adjusting strength of counterterms as a function of  $\lambda$ .

$$\delta m_x^2 \rightarrow C(\lambda)\delta m_x^2$$

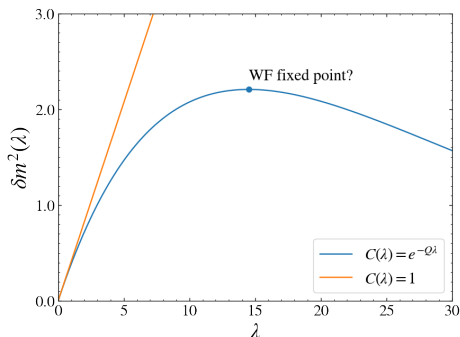


Plotted at  $\lambda = 10$ , refinement of 96.

# Non-Perturbative Counter Terms

We obtain good results with  $C(\lambda) = e^{-Q\lambda}$ . Counterterms become

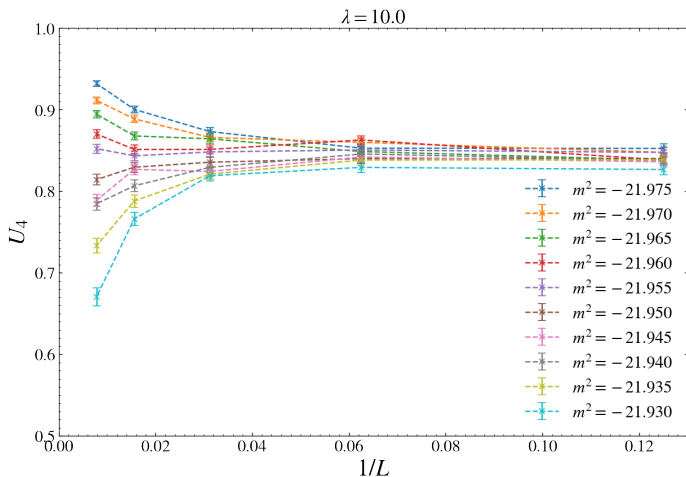
$$\delta m_x^2(\lambda) = 6Q\lambda e^{-Q\lambda} \log(\sqrt{g_x})$$



Maximum at  $\lambda_{WF} \simeq 14.5$ , possibly corresponding to the Wilson-Fisher fixed point.

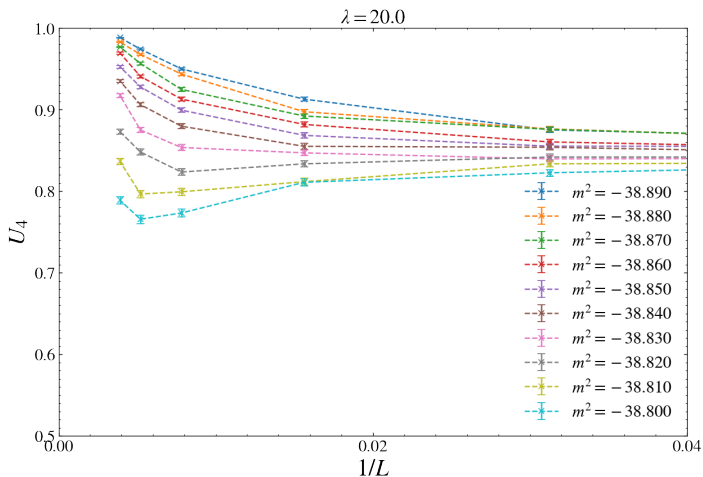
# Critical Surface with Non-Perturbative Counter Terms

Restoration of critical surface is greatly improved at  $\lambda = 10$ .



# Critical Surface with Non-Perturbative Counter Terms

However, obstruction still exists at larger  $\lambda$



# Conclusion

- Stability of critical surface near the strong coupling Wilson-Fisher fixed point is improved by tuning the strength of the perturbative counterterms.
- Further tuning is required to fully restore symmetry in the continuum limit.
- Exact form of non-perturbative counterterms is not well understood, but may provide direct access to the Wilson-Fisher fixed point.

# Future Directions

- Tune non-perturbative counterterms to improve restoration of spherical symmetry at larger  $\lambda$ .
- Improve statistics, larger lattices (USQCD allocation on Fermilab lattice cluster).
- Measure projection of operators onto spherical harmonics to improve quantitative measurement of spherical symmetry breaking.
- Explore other manifolds (boundary AdS theory w/ Cameron Coghburn et al.)
- Explore other theories, add fermions (previous talk by Rich Brower).

# References

- [1] Richard C. Brower, Michael Cheng, Evan S. Weinberg, George T. Fleming, Andrew D. Gasbarro, Timothy G. Raben, and Chung-I Tan. Lattice  $\phi^4$  field theory on Riemann manifolds: Numerical tests for the 2-d Ising CFT on  $\mathbb{S}^2$ . *Phys. Rev. D*, 98(1):014502, 2018.
- [2] Andrew David Gasbarro. Studies of conformal behavior in strongly interacting quantum field theories, 2019.

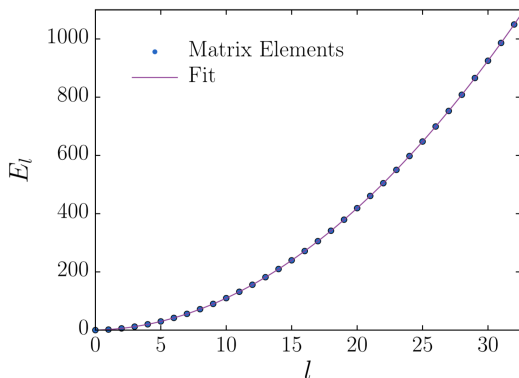


# Classical Theory

The free action ( $\lambda = 0$ ) in terms of a quadratic interaction matrix  $M_{xy}$  is

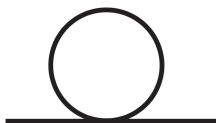
$$S_{\text{free}} = \sum_{x,y} \phi_x M_{xy} \phi_y$$

Eigenvalues of  $M_{xy}$  projected onto spherical harmonics approach the appropriate values  $\ell(\ell + 1)$  and degeneracy  $2\ell + 1$  as the continuum limit is approached (shown for refinement of 128).



## Perturbative Counter Terms

In 2d, 1-loop diagram is logarithmically UV divergent. Higher order diagrams are finite.



Performing the perturbative expansion in  $\lambda$ , renormalization requires the addition of *local* mass counterterms  $\delta m_x^2 = 6\lambda\delta G_{xx}$  where  $\delta G_{xx}$  is the diagonal of  $[M_{xy}]^{-1}$  after subtracting out the position-independent contribution to remove the logarithmic divergence.

$$\delta G_{xx} = [M_{xx}]^{-1} - \sum_y \sqrt{g_y} [M_{yy}]^{-1}$$

# Perturbative Counter Terms

Continuum form of  $\delta G_{xx}$  is

$$\delta G_{xx} = Q \log \left( \frac{a^2}{a_x^2} \right) = -Q \log(\sqrt{g_x})$$

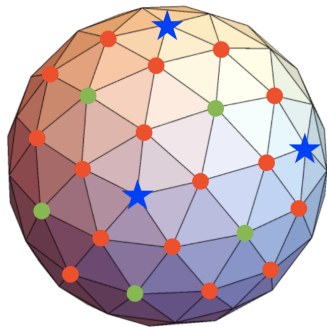
where  $Q = \sqrt{3}/8\pi$  and  $\sqrt{g_x}$  is the normalized area associated with site  $x$ . Continuum counterterm is then

$$\delta m_x^2 = 6Q\lambda \log(\sqrt{g_x})$$

In most cases, calculating  $[M_{xy}]^{-1}$  numerically results in a more uniform handling of the 12 exceptional points than this continuum form.

# Icosahedrally Distinct Points

We take advantage of icosahedral symmetry by binning in groups of “distinct” icosahedral points.



Shown for a refinement of 3. The stars are the 12 “exceptional” points.