

Universal Properties of the Vortex in the (2+1)-d $O(2)$ through dualization

Joao C. Pinto Barros

AEC, Institute for Theoretical Physics, University of Bern

July 27, 2021

Lattice 2021

- The (2+1)-d $O(2)$ Scalar Field Theory;
 - broken phase;
 - vortex excitation;

- The (2+1)-d $O(2)$ Scalar Field Theory;
 - broken phase;
 - vortex excitation;
- The quantum vortex;
 - Dualization: vortex as a charged particle;
 - Non-perturbative quantization;

- The (2+1)-d $O(2)$ Scalar Field Theory;
 - broken phase;
 - vortex excitation;

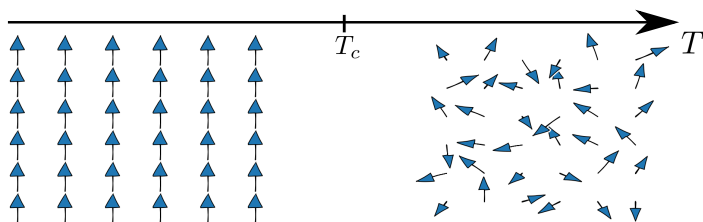
- The quantum vortex;
 - Dualization: vortex as a charged particle;
 - Non-perturbative quantization;

- Results and perspectives.

The (2+1)-d $O(2)$ model

Broken phase

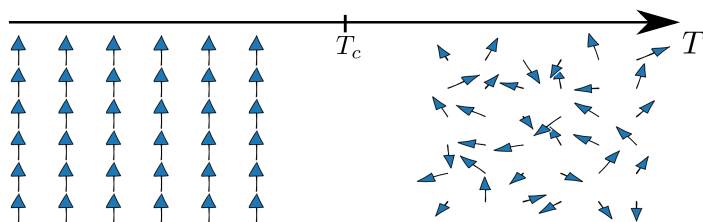
Symmetric Phase



The (2+1)-d $O(2)$ model

Broken phase

Symmetric Phase



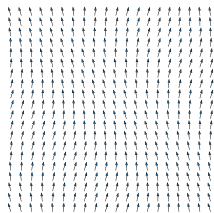
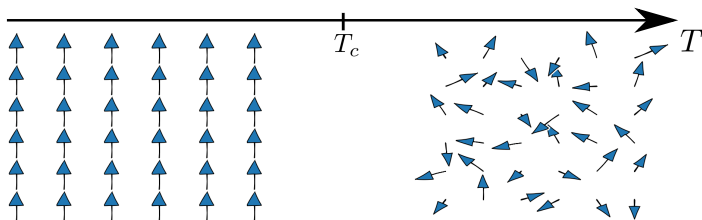
In the broken phase:

- Spontaneous breakdown of the $O(2)$ symmetry;
- Massless modes: Goldstone bosons.

The (2+1)-d $O(2)$ model

Broken phase

Symmetric Phase



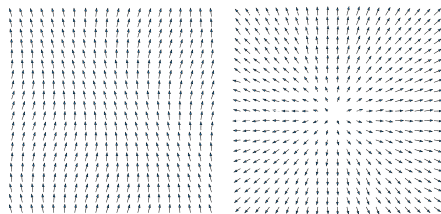
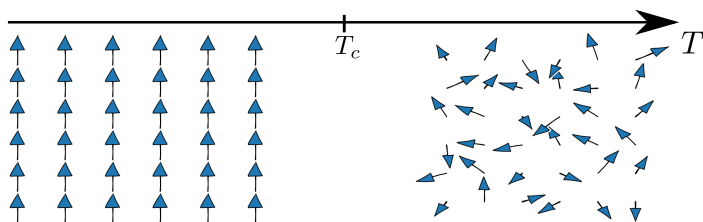
In the broken phase:

- Spontaneous breakdown of the $O(2)$ symmetry;
- Massless modes: Goldstone bosons.

The (2+1)-d $O(2)$ model

Broken phase

Symmetric Phase



In the broken phase:

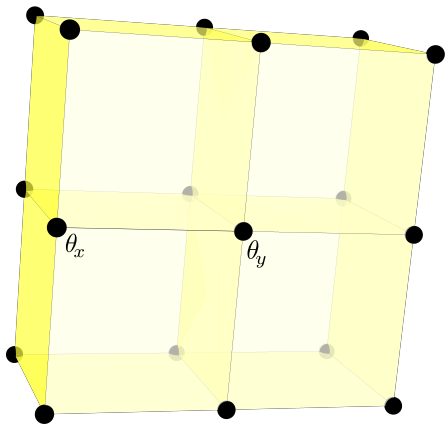
- Spontaneous breakdown of the $O(2)$ symmetry;
- Massless modes: Goldstone bosons.

- Classically, the mass is logarithmically divergent $E(R) \sim \pi v^2 \log \frac{R}{R_0}$;

- Classically, the mass is logarithmically divergent $E(R) \sim \pi v^2 \log \frac{R}{R_0}$;
- Semi-classical approximation quantizes only collective degrees of freedom;
 - vortex has an infinite mass;
J.-M. Duan, Phys. Rev. B49 (1994) 12381
D. J. Thouless, J. R. Anglin, Phys. Rev. Lett. 99 (2007) 105301.
 - but a finite mass associated to the core;
J.-M. Duan, Phys. Rev. B49 (1994) 12381.
G. Baym, E. Chandler, J. Low Temp. Phys. 50 (1983) 57
 - and recently it was argued that the mass does not diverge at all;
 - breaks translation invariance;
G. Delfino, W. Selke, A. Squarcini, Phys. Rev. Lett. 122 (2019) 050602

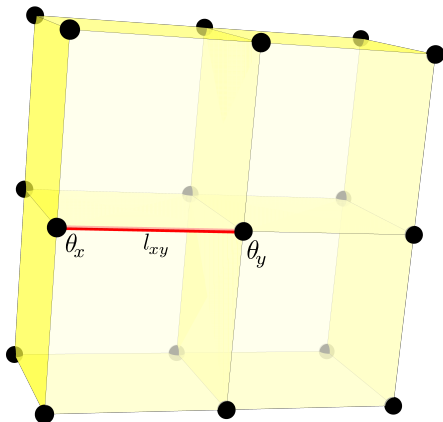
- Classically, the mass is logarithmically divergent $E(R) \sim \pi v^2 \log \frac{R}{R_0}$;
- Semi-classical approximation quantizes only collective degrees of freedom;
 - vortex has an infinite mass;
J.-M. Duan, Phys. Rev. B49 (1994) 12381
D. J. Thouless, J. R. Anglin, Phys. Rev. Lett. 99 (2007) 105301.
 - but a finite mass associated to the core;
J.-M. Duan, Phys. Rev. B49 (1994) 12381.
G. Baym, E. Chandler, J. Low Temp. Phys. 50 (1983) 57
 - and recently it was argued that the mass does not diverge at all;
 - breaks translation invariance;
G. Delfino, W. Selke, A. Squarcini, Phys. Rev. Lett. 122 (2019) 050602
- Fully non-perturbative approach is non-trivial:
 - vortex correlation not readily amenable to numerical simulations;
 - single vortex never occurs at finite periodic volume.

The quantum vortex as a quantum particle



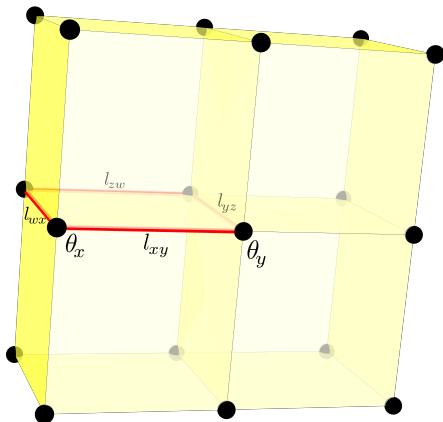
$$s(\theta_x - \theta_y)$$

The quantum vortex as a quantum particle



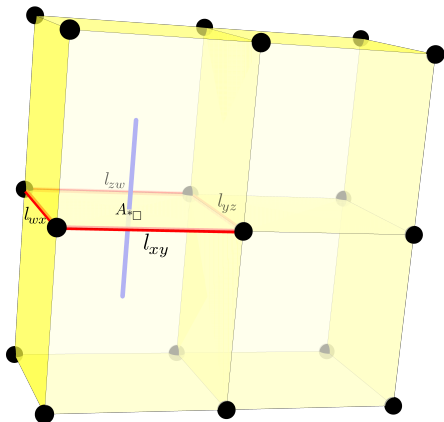
$$s(\theta_x - \theta_y) \rightarrow s(l_{xy})$$

The quantum vortex as a quantum particle



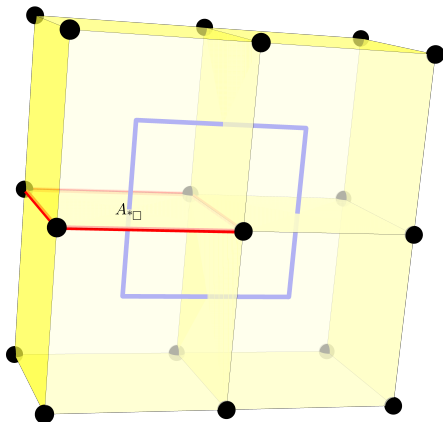
$s(l_{xy})$ with constraints $l_{xy} + l_{yz} + l_{zw} + l_{wx} = 0$

The quantum vortex as a quantum particle



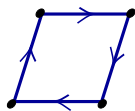
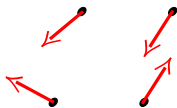
$$s(l_{xy}, A_{*\square}) \text{ with } \delta(l_{xy} + l_{yz} + l_{zw} + l_{wx}) = \sum_{A_{*\square} \in \mathbb{Z}} e^{iA_{*\square}(l_{xy} + l_{yz} + l_{zw} + l_{wx})}$$

The quantum vortex as a quantum particle



Integrate out $l_{xy} \Rightarrow$ Gauge Theory : $\tilde{\mathfrak{s}}(A_{*\square})$

The dualization picture



Continuous spin model	DUAL	Scalar QED
$O(2)$ global symmetry	\Leftrightarrow	$U(1)$ local symmetry
Weak/Strong coupling	\Leftrightarrow	Strong/Weak coupling
Symmetric phase	\Leftrightarrow	Higgs phase
Broken phase	\Leftrightarrow	Coulomb phase
<u>Goldstone bosons</u>	\Leftrightarrow	<u>Photons</u>
<u>Vortex</u>	\Leftrightarrow	<u>Charged scalar</u>

J. Fröhlich, G. Morchio, F. Strocchi, Phys. Lett. B89 (1979) 61

D. Buchholz, K. Fredenhagen, Commun. Math. Phys. 84 (1982) 1

D. Buchholz, Commun. Math. Phys. 85 (1982) 49

D. Buchholz, J. E. Roberts, Commun. Math. Phys. 330 (2014) 935

The charged particle

- Vortex correlation \Leftrightarrow charged particle (χ_x) correlation;

The charged particle

- Vortex correlation \Leftrightarrow charged particle (χ_x) correlation;
- The proper gauge invariant operator creates the particle along with a cloud of soft photons

$$\chi_x^C = e^{i\Delta^{-1}\delta A_x} \chi_x$$

The charged particle

- Vortex correlation \Leftrightarrow charged particle (χ_x) correlation;
- The proper gauge invariant operator creates the particle along with a cloud of soft photons

$$\chi_x^C = e^{i\Delta^{-1}\delta A_x} \chi_x$$

$$\chi_x^C = e^{i\Delta^{-1}\delta A_x} \chi_x$$

χ_x

 Charged particle

Charged particle

+

Cloud of photons

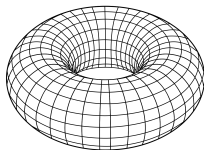
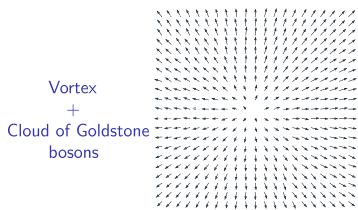


- Vortex correlation function: $\langle \chi_x^C \chi_y^{C\dagger} \rangle$

P. A. M. Dirac, *Canad. J. Phys.* 33 (1955) 650

J. Fröhlich, P. A. Marchetti, *Euro. Phys. Lett.* 2 (1986) 933

The charged particle at finite volume

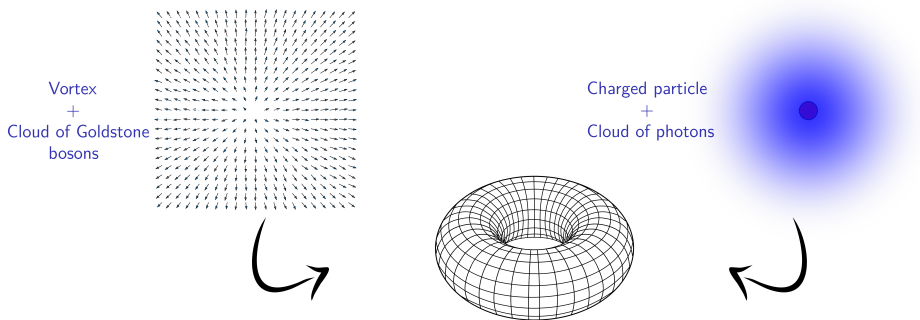


Charged particle
+
Cloud of photons



- No net charge on the torus;

The charged particle at finite volume



- No net charge on the torus;
- C-periodic boundary conditions (C* boundary conditions):

$$\begin{aligned}A_\mu(x + L\hat{i}) &= -A_\mu(x) - \partial_\mu\varphi_i(x) \\ \chi(x + L\hat{i}) &= \chi(x)^* e^{i\varphi_i(x)}\end{aligned}$$

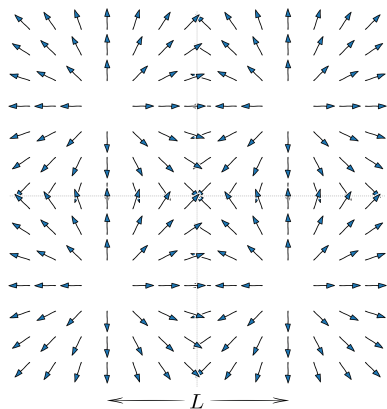
U.-J. Wiese, Nucl. Phys. B375 (1992) 45

B. Lucini, A. Patella, A. Ramos, N. Tantalò, JHEP 1602 (2016) 076C

I. Campos, P. Fritzsche, M. Hansen, M. K. Marinkovic, A. Patella, A. Ramos, N. Tantalò, Eur. Phys. J. C (2020) 80:195

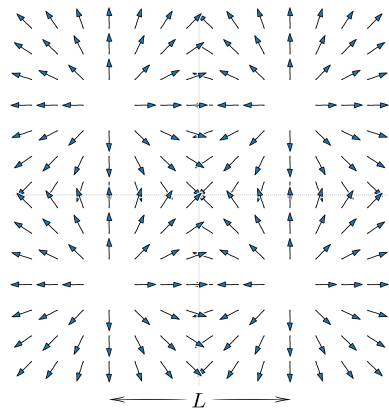
The C-periodic vortex: mass and charge computation

Periodic

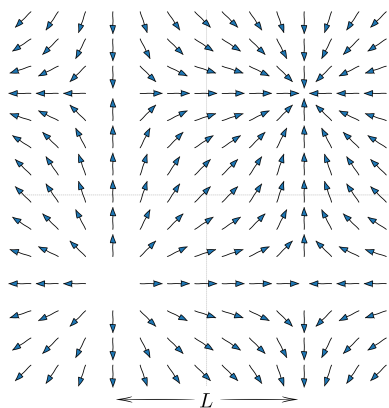


The C-periodic vortex: mass and charge computation

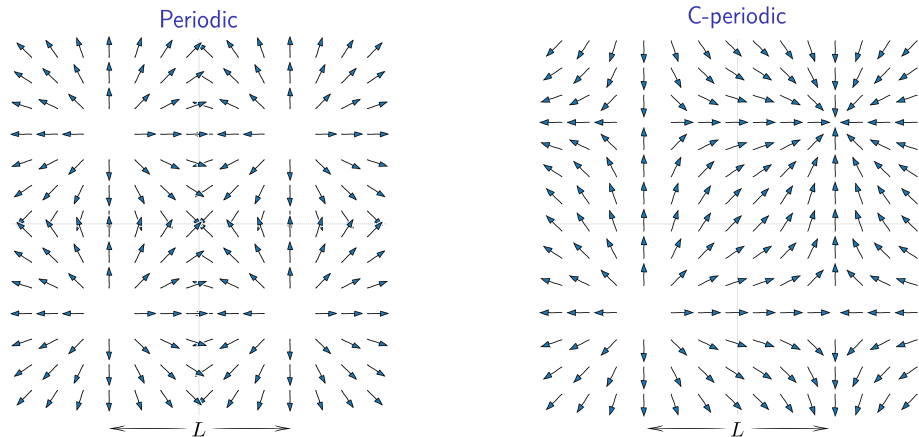
Periodic



C-periodic

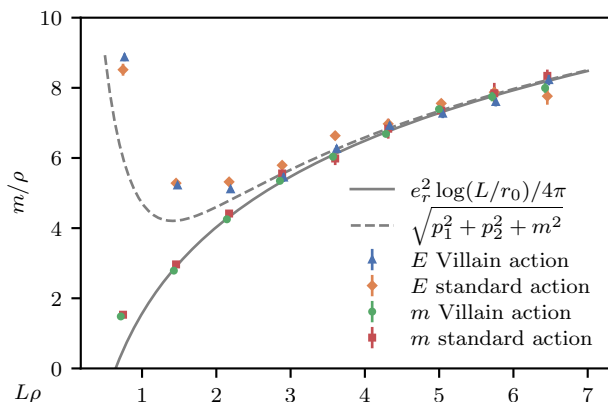


The C-periodic vortex: mass and charge computation



Vortices interact with their C-periodic copy

$$m = \frac{e^2}{4\pi} \log \left(\frac{L}{r_0} \right) \rightarrow \text{Determine the charge}$$



- Log divergent mass;
- Universal vortex charge: $e_r^2 = 3.58(8) \times (4\pi\rho)$;
- Breaking of Lorentz invariance: $E = m + \frac{\rho^2}{2m_k} \Rightarrow \frac{m_k}{m} = 0.71(3)$ (for $L\rho = 1.43(2)$).

- Computation of the vortex mass and vortex charge
 - The logarithmic divergent mass survives quantization;
 - The vortex charge gets renormalized.

M. Hornung, JCPB, and U-J. Wiese arXiv:2106.16191 (2021).



- Computation of the vortex mass and vortex charge
 - The logarithmic divergent mass survives quantization;
 - The vortex charge gets renormalized.

M. Hornung, JCPB, and U-J. Wiese arXiv:2106.16191 (2021).



- Future directions
 - Study the other side of the phase transition;

- Computation of the vortex mass and vortex charge
 - The logarithmic divergent mass survives quantization;
 - The vortex charge gets renormalized.

M. Hornung, JCPB, and U.-J. Wiese arXiv:2106.16191 (2021).



- Future directions
 - Study the other side of the phase transition;
 - Gauge the original $O(2)$ symmetry;

- Computation of the vortex mass and vortex charge
 - The logarithmic divergent mass survives quantization;
 - The vortex charge gets renormalized.

M. Hornung, JCPB, and U-J. Wiese arXiv:2106.16191 (2021).



- Future directions
 - Study the other side of the phase transition;
 - Gauge the original $O(2)$ symmetry;
 - Study the $SU(2)$ quantum vortex.

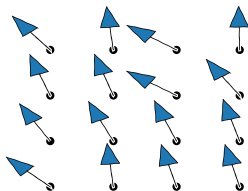
- Computation of the vortex mass and vortex charge
 - The logarithmic divergent mass survives quantization;
 - The vortex charge gets renormalized.

M. Hornung, JCPB, and U-J. Wiese arXiv:2106.16191 (2021).



- Future directions
 - Study the other side of the phase transition;
 - Gauge the original $O(2)$ symmetry;
 - Study the $SU(2)$ quantum vortex.





$$\text{Standard: } e^{-S(\{\theta\})} = \prod_{\langle x,y \rangle} e^{\frac{1}{g^2} \cos(\theta_x - \theta_y)}$$

$$\text{Villain: } e^{-S(\{\theta\})} = \prod_{\langle x,y \rangle} \sum_{n_{xy} \in \mathbb{Z}} e^{-\frac{1}{2g^2} (\theta_x - \theta_y + 2\pi n_{xy})^2}$$

The dualized action

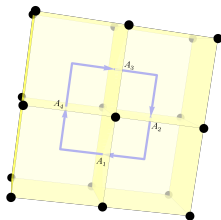
- The final gauge theory will depend on the initial action;
- We probe two actions:
 - Villain action

$$e^{-S(\{\theta\})} = \prod_{\langle x,y \rangle} \sum_{n_{xy} \in \mathbb{Z}} e^{-\frac{1}{2g^2} (\theta_x - \theta_y + 2\pi n_{xy})^2} \stackrel{\text{Dual}}{\Leftrightarrow} e^{-\tilde{S}(\{A\})} = \prod_{\square} e^{-\frac{g^2}{2} F_{\square}^2}$$

- Standard action

$$e^{-S(\{\theta\})} = \prod_{\langle x,y \rangle} e^{\frac{1}{g^2} \cos(\theta_x - \theta_y)} \stackrel{\text{Dual}}{\Leftrightarrow} e^{-\tilde{S}(\{A\})} = \prod_{\square} I_{F_{\square}} \left(\frac{2}{g^2} \right)$$

$$F_{\square} = A_1 + A_2 - A_3 - A_4$$



- Integer gauge theory

Dual action : $A_{x\mu} \in \mathbb{Z}$

$$\tilde{S}(\{A\}) = \sum_{\square} \tilde{s}(F_{\square})$$

- Integer gauge theory

Dual action : $A_{x\mu} \in \mathbb{Z}$

$$\tilde{S}(\{A\}) = \sum_{\square} \tilde{s}(F_{\square})$$

- It can be regarded as $(2 + 1)$ QED coupled to a charged scalar field

J. Fröhlich, P. A. Marchetti, Euro. Phys. Lett. 2 (1986) 933

Scalar QED : $\bar{A}_{x\mu} \in \mathbb{R}$

$$S_{\text{QED}}(\{\bar{A}\}, \{\chi\}) = \sum_{\square} s(\bar{F}_{\square}) - \frac{\kappa}{2} \sum_{x,\mu} \left(\chi_{x+\hat{\mu}}^* e^{i\bar{A}_{x\mu}} \chi_x + \chi_x^* e^{-i\bar{A}_{x\mu}} \chi_{x+\hat{\mu}} \right)$$

- Integer gauge theory

Dual action : $A_{x\mu} \in \mathbb{Z}$

$$\tilde{S}(\{A\}) = \sum_{\square} \tilde{s}(F_{\square})$$

- It can be regarded as (2 + 1) QED coupled to a charged scalar field
J. Fröhlich, P. A. Marchetti, Euro. Phys. Lett. 2 (1986) 933

Scalar QED : $\bar{A}_{x\mu} \in \mathbb{R}$

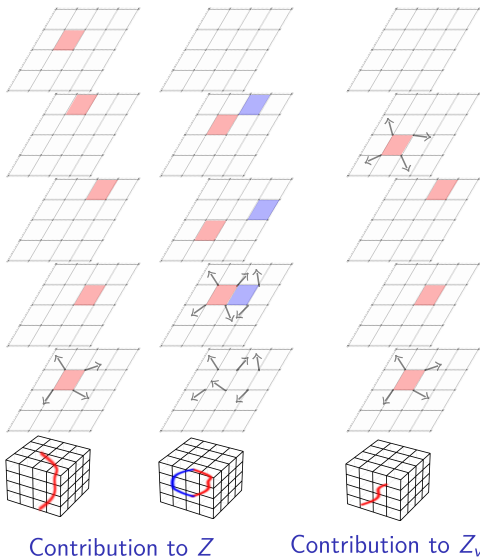
$$S_{\text{QED}}(\{\bar{A}\}, \{\chi\}) = \sum_{\square} s(\bar{F}_{\square}) - \frac{\kappa}{2} \sum_{x,\mu} \left(\chi_{x+\hat{\mu}}^* e^{i\bar{A}_{x\mu}} \chi_x + \chi_x^* e^{-i\bar{A}_{x\mu}} \chi_{x+\hat{\mu}} \right)$$

- Limit $\kappa \rightarrow +\infty$ and unitary gauge, $\chi_x = 1$:

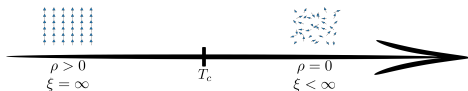
$$-\frac{\kappa}{2} \sum_{x,\mu} \left(\chi_{x+\hat{\mu}}^* e^{i\bar{A}_{x\mu}} \chi_x + \chi_x^* e^{-i\bar{A}_{x\mu}} \chi_{x+\hat{\mu}} \right) \rightarrow -\kappa \sum_{x,\mu} \cos \bar{A}_{x\mu} \Rightarrow \bar{A}_{x\mu} \in 2\pi\mathbb{Z}$$

The vortex as a charged particle

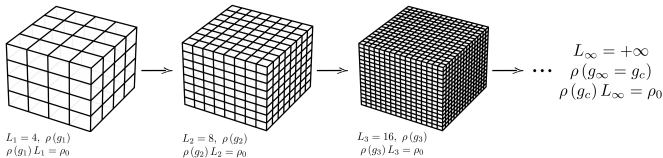
$$\langle \chi_x^C \chi_y^{C\dagger} \rangle \xrightarrow[\text{back}]{\text{Dualize}} \frac{Z_v}{Z}$$



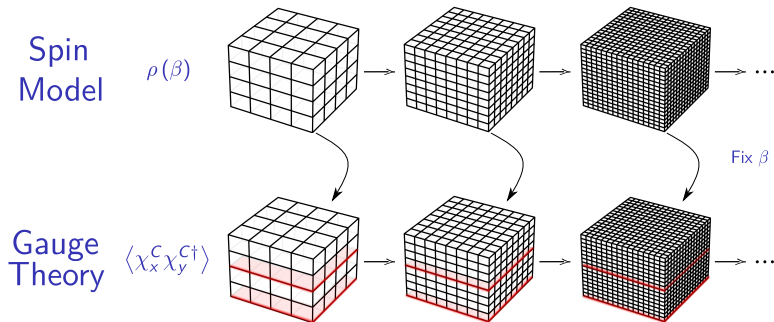
- The spin stiffness characterizes the distance to the critical point;

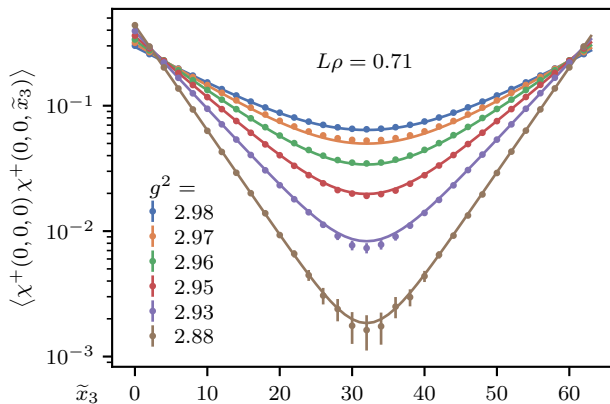


- Approaching the continuum limit



Approach the continuum limit at a finite volume characterized by $\rho_0 = \rho L$





Results: approaching the continuum limit

