

Symmetric mass generation in lattice gauge theory

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Lattice 2021

Outline

Can fermions acquire mass without breaking flavor symmetries ?

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$SO(4)$ invariant reduced staggered fermion model in $D = 4$ has massless and massive phases without any intervening symmetry broken phase

Action and Symmetries

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$$\mathcal{S} = \sum_{x,\mu} \psi^a [\eta_\mu \Delta_\mu \delta_{ab} + \mathbf{G} \phi_{ab}^+] \psi^b + \sum_x \frac{1}{4} \phi_+^2 - \frac{\kappa}{2} \sum_{x,\mu} [\phi_x^+ \phi_{x+\mu}^+ + \phi_x^+ \phi_{x-\mu}^+]$$

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- We use the isomorphism $SO(4) = SU_+(2) \times SU_-(2)$ for introducing Yukawa field ϕ^{ab} where $\phi_+^{ab} = P_+ \phi^{ab} = \frac{1}{2}(\phi^{ab} + \frac{1}{2} \epsilon_{abcd} \phi^{cd})$

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(Reduced) staggered fermion model

$SO(4)$ flavor symmetry doesn't allow a naive fermion bilinear and Z_2 symmetry ($\phi^+ \rightarrow \pm\phi^+, \psi \rightarrow \pm i\epsilon(x)\psi$) forbids any possible spontaneous symmetry breaking scenario.

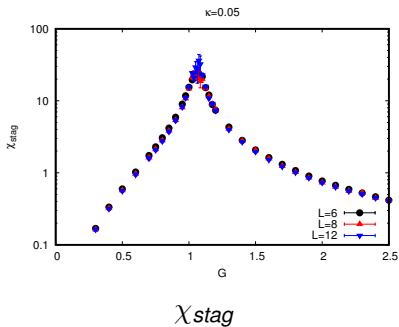
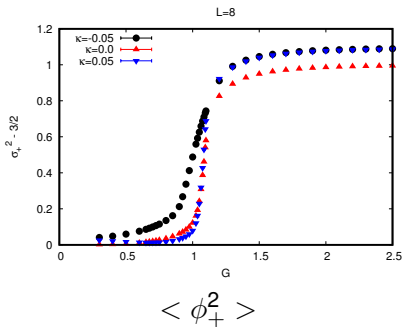
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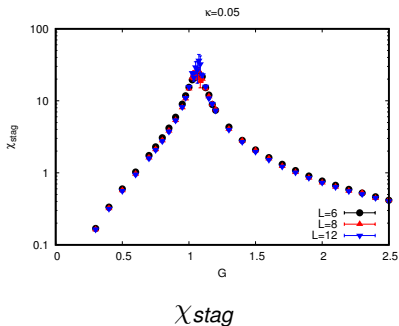
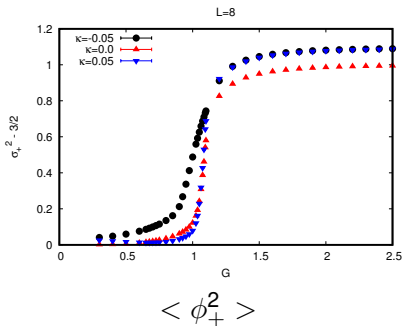
With non-zero κ we can achieve a direct transition between the massless and massive phase without breaking the global $SO(4)$ symmetry and the discrete Z_2 symmetry.

Phase structure from numerical simulation

Phase structure from numerical simulation



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Connected susceptibility

$\chi_{stag} = \frac{1}{V} \sum_{x,y} \langle \epsilon(x) \psi^a(x) \psi^b(x) \epsilon(y) \psi^a(y) \psi^b(y) \rangle$ shows no volume scaling while we have evidence for four-fermion condensate for non-zero κ

Extending this to a gauge model

Is it possible to retain this phase structure if we make one of the $SU(2)$ subgroups a gauge symmetry.

Action

$$S_F = \sum_{x,\mu} \frac{1}{2} \eta_\mu(x) \text{Tr}[\psi^\dagger(x) U_\mu(x) \psi(x + \mu) V_\mu^\dagger(x) - \psi^\dagger(x) U_\mu^\dagger(x - \mu) \psi(x - \mu) V_\mu(x - \mu)]$$

Symmetries

$$\begin{aligned}\psi &\rightarrow G(x) \psi(x) H^\dagger \\ U_\mu(x) &\rightarrow G(x) U_\mu(x) G^\dagger(x + \mu) \\ V_\mu(x) &\rightarrow H(x + \mu) V_\mu H^\dagger(x)\end{aligned}$$

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We gauge the symmetry under which the scalar is a singlet.

Yukawa term and auxiliary field action

$$S_Y = \sum_x \text{Tr}(\phi(x)\psi(x)\psi^\dagger(x)) + \frac{1}{2\lambda^2} \sum_x \text{Tr}(\phi(x))^2$$

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Integrating out the scalar field generates four-fermion term

$$\sum_x \text{Tr}(\psi(x)\psi^\dagger(x)\psi(x)\psi^\dagger(x))$$

Continued

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 $V_\mu(x) \rightarrow \pm V_\mu(x), \psi(x) \rightarrow \epsilon(x)\psi(x)$. Polyakov line is a good order parameter.

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Reduction to $SO(4)$ model

Using the pseudo-reality condition for $SU(2)$ elements $\psi^\dagger = \sigma_2 \psi^T \sigma_2$ and $\psi = \sigma_\mu \chi_\mu$ it is easy to show that this model reduces to the aforementioned $SO(4)$ model in $\beta_G, \beta_H \rightarrow \infty$ limit.

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As β_H increases the lattice spacing a decreases and the physical temperature $T = \frac{1}{La}$ increases with de-confinement once temperature exceeds the strong scale.

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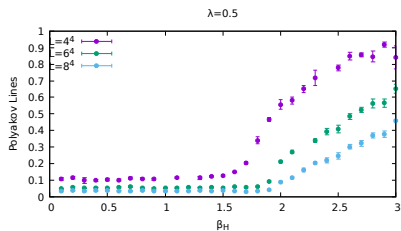
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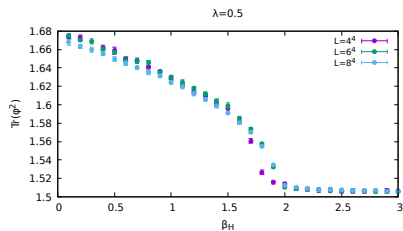
As β_H increases the lattice spacing a decreases and the physical temperature $T = \frac{1}{La}$ increases with de-confinement once temperature exceeds the strong scale. We keep Yukawa coupling small and scan for four-fermion condensate as β_H is varied.

Numerical results

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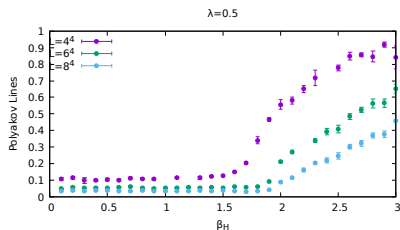


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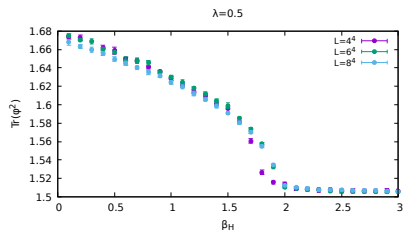


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$$\langle \text{Tr}(\phi^2) \rangle$$

$\text{Tr}(\phi^2)$ is used as a proxy for four-fermion condensate. Here the four-fermion condensate is a result of gauge interactions since the Yukawa coupling is below the critical λ_c to drive the transition.

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- Spectrum of composite states in the confining regime requires further work.
- Can this mechanism be used to decouple (have infinitely massive) right-handed states and obtain a chiral theory in the continuum limit ?

Thank you!

Collaborators

Simon Catterall , Goksu Can Toga

Funding and computing resources



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