

Stout-Smearing/Gradient Flow and c_{SW} at One Loop Order

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Goal:

Calculate the clover coefficient c_{5W} to one loop order with stout-smearing.

Previous work:

- *Sheikholeslami, Wohlert 1985* [Nuc1.Phys.B 259 (1985) 572]
→ introduce clover term, note $c_{5W} = 1$ at tree level.
- *Wohlert 1987* [DESY-87-069]
→ calculation of c_{5W} to one loop using twisted boundary conditions.
- *Aoki, Kuramashi 2003* [arXiv:hep-lat/0306015]
→ calculation of c_{5W} to one loop for different improved gauge actions using a small gluon mass as regulator.
- *Horsley, Perlt, Rakow, Schierholz, Schiller 2008* [arXiv:0807.0345]
→ calculation of c_{5W} to one loop for SLINC-fermions (i.e. stout-smearred links in the Wilson part of the action but not in the clover term).

$\mathcal{O}(a)$ -improved action:

$$\mathcal{S}_I = \mathcal{S}_{\text{Wilson}} + c_{\text{SW}} \cdot \sum_x \sum_{\mu < \nu} \bar{\psi}(x) \frac{1}{2} \sigma_{\mu\nu} F_{\mu\nu}(x) \psi(x)$$

$F_{\mu\nu}$: lattice field strength tensor i.e. clover shaped sum of plaquettes.

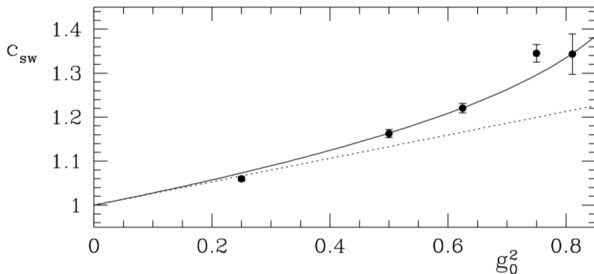
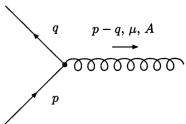


Figure: Non-perturbative determination of c_{SW} (data points), one-loop value (dashed line) and rational fit (solid line). Graphic from *Lüscher et al. 1996* [[hep-lat/9609035](#)]

Perturbative determination of $c_{\text{SW}} = c_{\text{SW}}^{(0)} + g_0^2 c_{\text{SW}}^{(1)} + \mathcal{O}(g_0^4)$ via qqq-vertex function



$$\Lambda_\mu^a(p, q) = \sum_{L=0}^{\infty} g_0^{2L+1} \Lambda_\mu^{a(L)}(p, q),$$

At tree level:

$$\begin{aligned} \Lambda_\mu^{a(0)}(p, q) &= (V_1^a)_\mu(p, q) \\ &= -g_0 T^a \left(i\gamma_\mu + a \left(\frac{1}{2}(p_\mu + q_\mu) - \frac{i}{2} c_{\text{SW}}^{(0)} \sum_\nu \sigma_{\mu\nu} (p_\nu - q_\nu) \right) \right) + \mathcal{O}(a^2) \end{aligned}$$

Sandwich with on-shell spinors:

$$\begin{aligned} \bar{u}(q) \Lambda_\mu^{a(0)}(p, q) u(p) &= -g_0 T^a \bar{u}(q) \left(i\gamma_\mu + \frac{a}{2} (1 - c_{\text{SW}}^{(0)}) (p_\mu + q_\mu) \right) u(p) + \mathcal{O}(a^2). \\ &\Rightarrow c_{\text{SW}}^{(0)} = 1 (= r) \end{aligned}$$

At one-loop level:

$$g_0^3 \Lambda_\mu^{a(1)} = -g_0^3 T^a \left(\gamma_\mu F_1 + a \not{q} \gamma_\mu F_2 + a \gamma_\mu \not{p} F_3 + a(p_\mu + q_\mu) G_1 + a(p_\mu - q_\mu) H_1 \right)$$

F_2 and F_3 do not contribute on-shell, H_1 vanishes due to symmetry arguments.

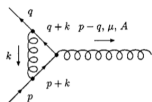
$$\begin{aligned} & g_0^3 \bar{u}(q) \left(\frac{a}{2} c_{\text{SW}}^{(1)} (p_\mu + q_\mu) T^a + \Lambda_\mu^{a(1)}(p, q) \right) u(p) \\ &= g_0^3 \bar{u}(q) \left(i \gamma_\mu F_1 + \frac{a}{2} (p_\mu + q_\mu) (c_{\text{SW}}^{(1)} - 2G_1) T^a \right) u(p) + \mathcal{O}(p^2, q^2) + \mathcal{O}(a^2) \\ &\Rightarrow c_{\text{SW}}^{(1)} = 2G_1 \end{aligned}$$

Extract G_1 :

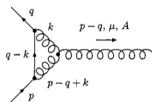
$$G_1 = -\frac{1}{8} \text{Tr} \left(\left(\frac{\partial}{\partial p_\mu} + \frac{\partial}{\partial q_\mu} \right) \Lambda_\mu^{a(1)} - \left(\frac{\partial}{\partial p_\nu} - \frac{\partial}{\partial q_\nu} \right) \Lambda_\mu^{a(1)} \gamma_\nu \gamma_\mu \right)_{p, q \rightarrow 0}^{\mu \neq \nu},$$

Diagrams:

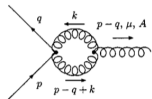
$$\Lambda_{\mu}^{a(1)} = \sum_{i=a, \dots, f} \Lambda_{\mu}^{a(1)(i)} = \sum_{i=a, \dots, f} \int_{-\pi}^{\pi} \frac{d^4 k}{(2\pi)^4} I_{\mu}^{a(i)}(p, q, k).$$



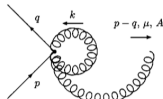
(a)



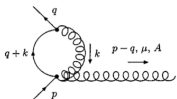
(b)



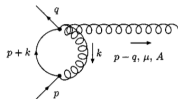
(c)



(d)



(e)



(f)

Example: Diagram (c)

$$I_{\mu}^{a(c)} = \sum_{b,c} \sum_{\nu,\rho} V_{2\nu\rho}^{bc}(p, q, q-p-k) G(k) G(p-q+k) V_{g3\nu\rho\mu}^{bca}(p-q+k, -k, q-p)$$

Extract divergent part:

$$G_1^{(i)} = \int_{-\pi}^{\pi} \frac{d^4 k}{(2\pi)^4} [I_{\mu}^{(i)}(k) - \theta(\pi^2 - k^2) J_{\mu}^{(i)}(k, 0)] + \frac{\Omega_3}{(2\pi)^4} \int_0^{\pi} J_{\mu}^{(i)}(k, \mu^2) k^3 dk$$

$$J^{(c)} = \frac{3}{4} N_{C\text{CSW}}^{(0)} \frac{2\pi^2}{(2\pi)^4} \int_0^{\pi} \frac{k^3}{(k^2 + \mu^2)^2} dk = \frac{3}{4} N_{C\text{CSW}}^{(0)} \frac{1}{16\pi^2} \left(\ln \left(\frac{\pi^2}{\mu^2} \right) - 1 \right) + \mathcal{O}(\mu^2)$$

$$\Rightarrow 2G_1^{(c)} = \frac{9}{2} L - 0.0813095$$

Diagram	Divergent part	Constant part	Aoki, Kuramashi
(a)	$-L/3$	0.00457196	0.004572(2)
(b)	$-9L/2$	0.0830768	0.08311(3)
(c)	$9L/2$	-0.0813095	$-0.08133(3)$
(d)	0	0.297394537	0.29739454(1)
(e)	$L/6$	-0.0175746	$-0.017574(1)$
(f)	$L/6$	-0.0175746	$-0.017574(1)$
Sum	0	0.268588292	0.26858825(1)

- Stout-smearing

$$U_\mu^{(1)}(x) = e^{i\rho Q_\mu(x)} U_\mu(x)$$

with

$$Q_\mu(x) = \frac{1}{2i} \left(W_\mu(x) - \frac{1}{3} \text{Tr}[W_\mu(x)] \right)$$

$$W_\mu(x) = \sum_{\nu \neq \mu} (P_{\mu\nu}^\dagger - P_{\mu\nu})$$

- Iterated smearing $\rightarrow U_\mu^{(n)}(x)$
- Wilson flow is generated by infinitesimal stout-smearings

$$\begin{aligned} \partial_t U_\mu(x, t) &= -g_0^2 \{ \partial_{x,\mu} \mathcal{S}_W[U(x, t)] \} U_\mu(x, t) = iQ_\mu(x, t) U_\mu(x, t) \\ U_\mu(x, 0) &= U_\mu(x) \end{aligned}$$

- Limit $n \rightarrow \infty$, $\rho \rightarrow 0$ with $n\rho = t = \text{const.}$ (t in lattice units, i.e. t/a^2).

Perturbative expansion of the smeared link variable:

At leading order:

$$U_\mu^{(1)}(x) = 1 + ig_0 T^a A_\mu^{a(1)}(x) + \mathcal{O}(g_0^2)$$
$$A_\mu^{a(1)}(x) = A_\mu^a(x) + \rho(Q_\mu^a(x))^{\text{LO}}$$

Fourier transform: ($\hat{k}_\mu = 2 \sin(\frac{1}{2} k_\mu)$)

$$A_\mu^{a(1)}(k) = \sum_\nu \left(f(k) \delta_{\mu\nu} - (f(k) - 1) \frac{\hat{k}_\mu \hat{k}_\nu}{\hat{k}^2} \right) A_\nu^a(k)$$

$$A_\mu^{a(n)}(k) = \sum_\nu \left(f(k)^n \delta_{\mu\nu} - (f(k)^n - 1) \frac{\hat{k}_\mu \hat{k}_\nu}{\hat{k}^2} \right) A_\nu^a(k)$$

$$A_\mu^a(k, t) = \sum_\nu \left(e^{-t\hat{k}^2} \delta_{\mu\nu} - (e^{-t\hat{k}^2} - 1) \frac{\hat{k}_\mu \hat{k}_\nu}{\hat{k}^2} \right) A_\nu^a(k, 0)$$

$$f(k) = 1 - \rho \hat{k}^2$$

Perturbative expansion of the smeared link variable:

At next-to-leading order:

$$U_{\mu}^{(1)}(x) = 1 + ig_0 T^a A_{\mu}^{a(1)}(x) - \frac{g_0^2}{2} \left(\frac{1}{2} \{T^a, T^b\} A_{\mu}^{a(1)}(x) A_{\mu}^{b(1)}(x) + \frac{1}{2} [T^a, T^b] A_{\mu}^{ab(1)}(x) \right) + \mathcal{O}(g_0^3)$$

$$A_{\mu}^{ab(1)}(x) = 4\rho \cdot \left[\frac{1}{2} (Q_{\mu}^a(x))^{\text{LO}} A_{\mu}^b(x) + (Q_{\mu}^{ab}(x))^{\text{NLO}} \right]$$

Fourier transform:

$$A_{\mu}^{ab(1)}(k_1, k_2) = 4\rho \sum_{\nu, \rho} g_{\mu\nu\rho}(k_1, k_2) A_{\nu}^a(k_1) A_{\rho}^b(k_2),$$

with

$$\begin{aligned} g_{\mu\nu\rho}(k_1, k_2) = & \delta_{\mu\nu} \sin\left(\frac{1}{2}(2k_{1\rho} + k_{2\rho})\right) \cos\left(\frac{1}{2}k_{2\mu}\right) \\ & - \delta_{\mu\rho} \sin\left(\frac{1}{2}(2k_{2\nu} + k_{1\nu})\right) \cos\left(\frac{1}{2}k_{1\mu}\right) \\ & - \delta_{\nu\rho} \sin\left(\frac{1}{2}(k_{1\mu} - k_{2\mu})\right) \cos\left(\frac{1}{2}(k_{1\nu} + k_{2\nu})\right). \end{aligned}$$

Perturbative expansion of the smeared link variable:

At next-to-leading order:

$$A_\mu^{ab(1)}(k_1, k_2) = 4\rho \sum_{\nu, \rho} g_{\mu\nu\rho}(k_1, k_2) A_\nu^a(k_1) A_\rho^b(k_2)$$
$$A_\mu^{ab(n)}(k_1, k_2) = 4\rho \sum_{m=0}^{n-1} \left[\sum_{\nu, \rho} g_{\mu\nu\rho}(k_1, k_2) A_\nu^{a(m)}(k_1) A_\rho^{b(m)}(k_2) \right].$$

$$A_\mu^{ab}(k_1, k_2, t) = 4 \sum_{\nu, \rho} g_{\mu\nu\rho}(k_1, k_2)$$
$$\times \sum_{\alpha\beta} \left[(e^{-t\hat{k}_1^2} - 1) \delta_{\nu\alpha} - (e^{-t\hat{k}_1^2} - 1) \frac{\hat{k}_{1\nu} \hat{k}_{1\alpha}}{\hat{k}_1^2} \right]$$
$$\times \left[(e^{-t\hat{k}_2^2} - 1) \delta_{\rho\beta} - (e^{-t\hat{k}_2^2} - 1) \frac{\hat{k}_{2\rho} \hat{k}_{2\beta}}{\hat{k}_2^2} \right] A_\alpha^a(k_1, 0) A_\beta^b(k_2, 0)$$
$$= 4 \sum_{\nu, \rho} g_{\mu\nu\rho}(k_1, k_2) [A_\nu^a(k_1, t) - A_\nu^a(k_1, 0)] [A_\rho^b(k_2, t) - A_\rho^b(k_2, 0)]$$

Next steps:

- Smearing/GF at order g_0^3 .
- Derive Feynman-rules with smearing (especially from the clover term).
- Calculate $c_{\text{SW}}^{(1)}$ depending on ρ , n or t .