

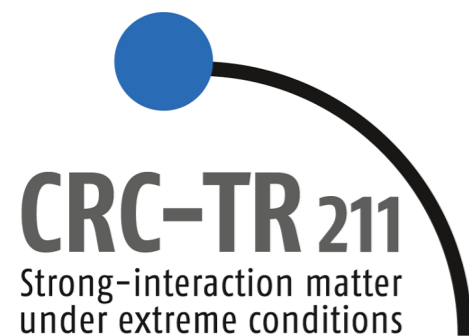
# Continuous temperature sampling in a single Monte-carlo Simulation

Parikshit Junnarkar

Technische Universität Darmstadt

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TECHNISCHE  
UNIVERSITÄT  
DARMSTADT

Work done in collaboration with

P. Thomas Jahn, Guy Moore and Daniel Robaina

***Multicanonical reweighting for the QCD topological susceptibility***

arXiv:2103.01069

*Phys Rev. D* 104 (2021) 1, 014502

- ▶ Motivation and general idea of the method.
- ▶ Application to QCD  $\chi_{\text{top}}(\beta)$  at high temperatures.
- ▶ Simulation algorithm.
- ▶ Results and comparison with the standard method.

- ▶ Consider a lattice calculation of  $\langle \mathcal{O} \rangle_\beta$

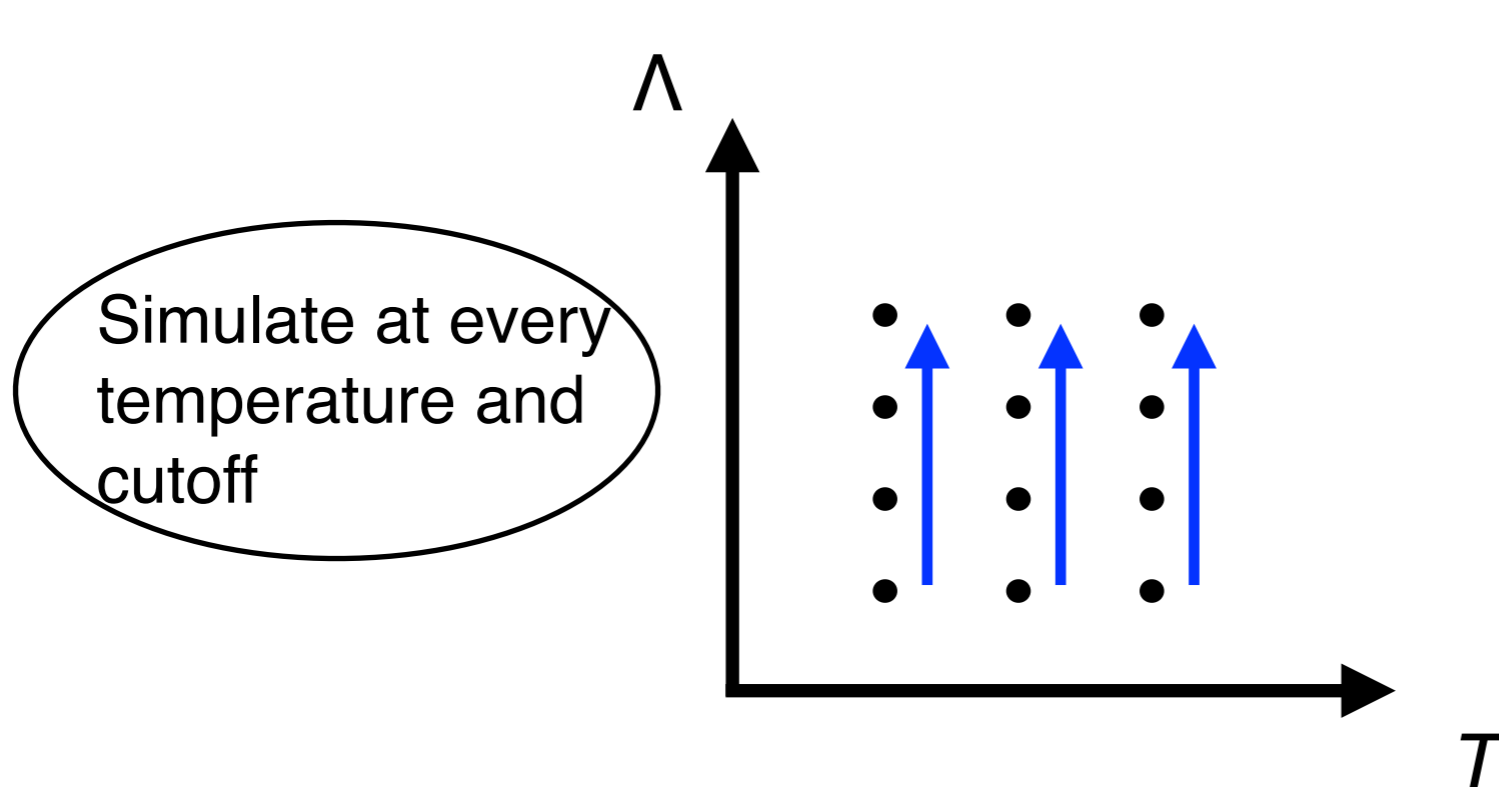
$$\langle \mathcal{O} \rangle_\beta = \frac{1}{\mathcal{Z}} \int \mathcal{D}U e^{-\beta S[U]} \mathcal{O}[U]$$

Single coupling  
corresponds to  
single temperature

- ▶ We are interested in the temperature and cutoff dependence

- ▶ Consider a lattice calculation of  $\langle \mathcal{O} \rangle_\beta$

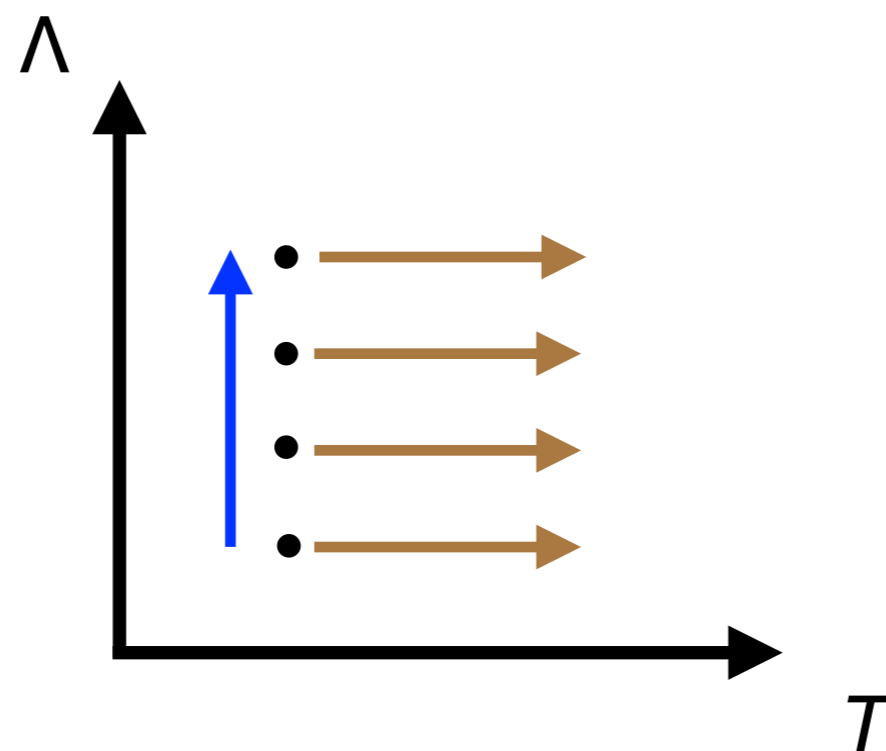
$$\langle \mathcal{O} \rangle_\beta = \frac{1}{\mathcal{Z}} \int \mathcal{D}U e^{-\beta S[U]} \mathcal{O}[U]$$



Could be expensive if the temperature range is really large

Perhaps there is a better way ?

- ▶ A novel way .... Sample temperatures continuously !
- ▶ Employ a reweighting in temperature at a fixed cutoff



Could be very cost effective

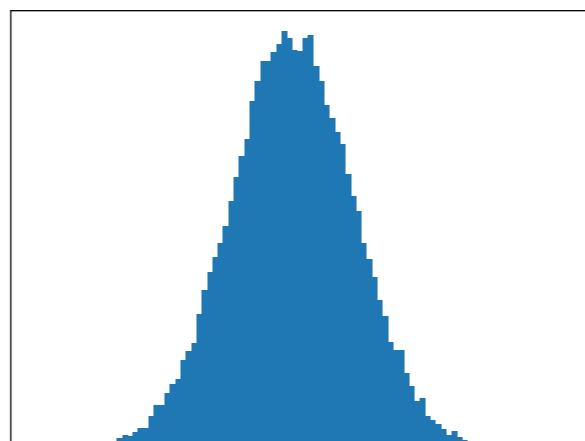
- ▶ Consider a lattice calculation of  $\langle \mathcal{O} \rangle_\beta$

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Single  $\beta$  simulates single temperature



Action histogram  
approx gaussian

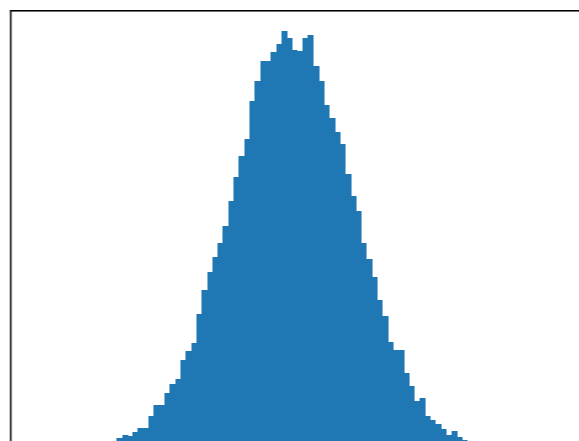


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Single  $\beta$  simulates single temperature

To simulate continuous temperatures, we need to simulate continuous  $\beta$



Action histogram approx gaussian

- ▶ Consider a lattice calculation of  $\langle \mathcal{O} \rangle_\beta$

$$\langle \mathcal{O} \rangle_\beta = \frac{1}{\mathcal{Z}} \int \mathcal{D}U e^{-\beta S[U]} \mathcal{O}[U] \quad \mathcal{Z} = \int \mathcal{D}U e^{-\beta S[U]}$$

- ▶ Replace the weight  $\beta S$

$$\begin{aligned} \mathcal{Z}(\beta) &= \int \mathcal{D}U e^{-\beta S} \\ &\quad \downarrow \\ &= \int \mathcal{D}U e^{-W(S)} \end{aligned} \quad \begin{aligned} W(S) &= \beta S \\ W'(S) &= \frac{dW(S)}{dS} = \beta \end{aligned}$$

Continuous  
Temperatures

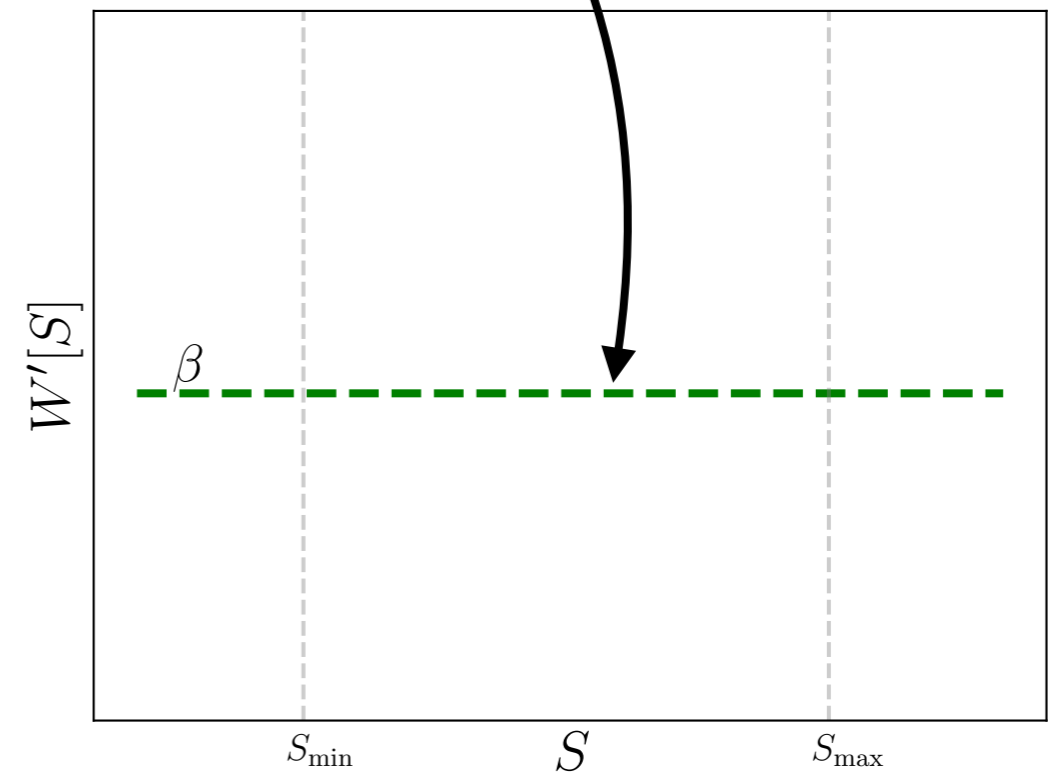


Continuous  
gauge couplings  $\beta'$ 's

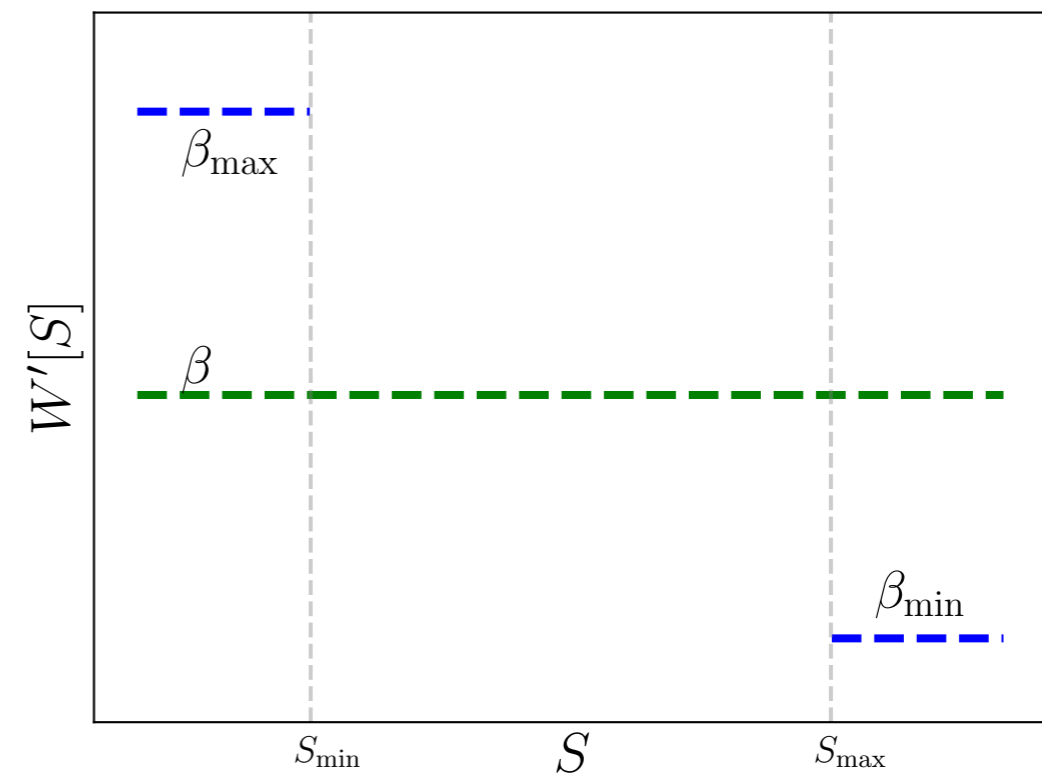


Continuous  $W'(S)$

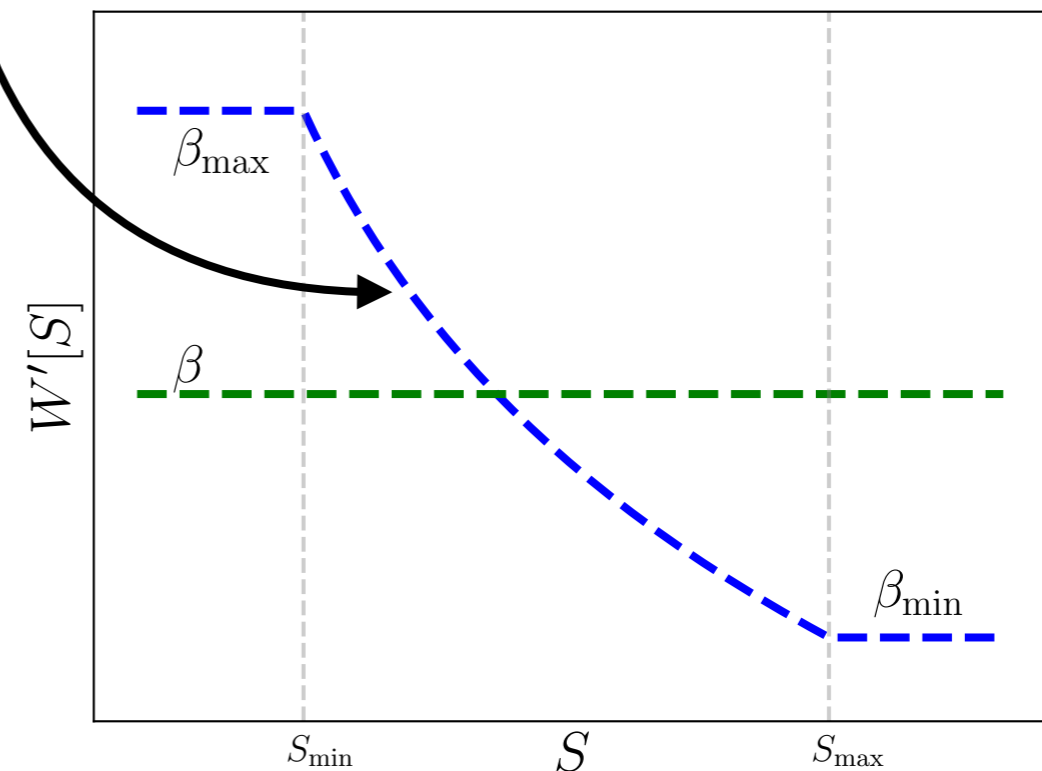
- ▶ Standard single temperature/coupling MC



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- ▶ Define the range for continuous temperatures



- ▶ Standard single temperature /  $\beta$  MC
- ▶ A continuous sampling in  $S$  within the range samples continuous  $\beta$



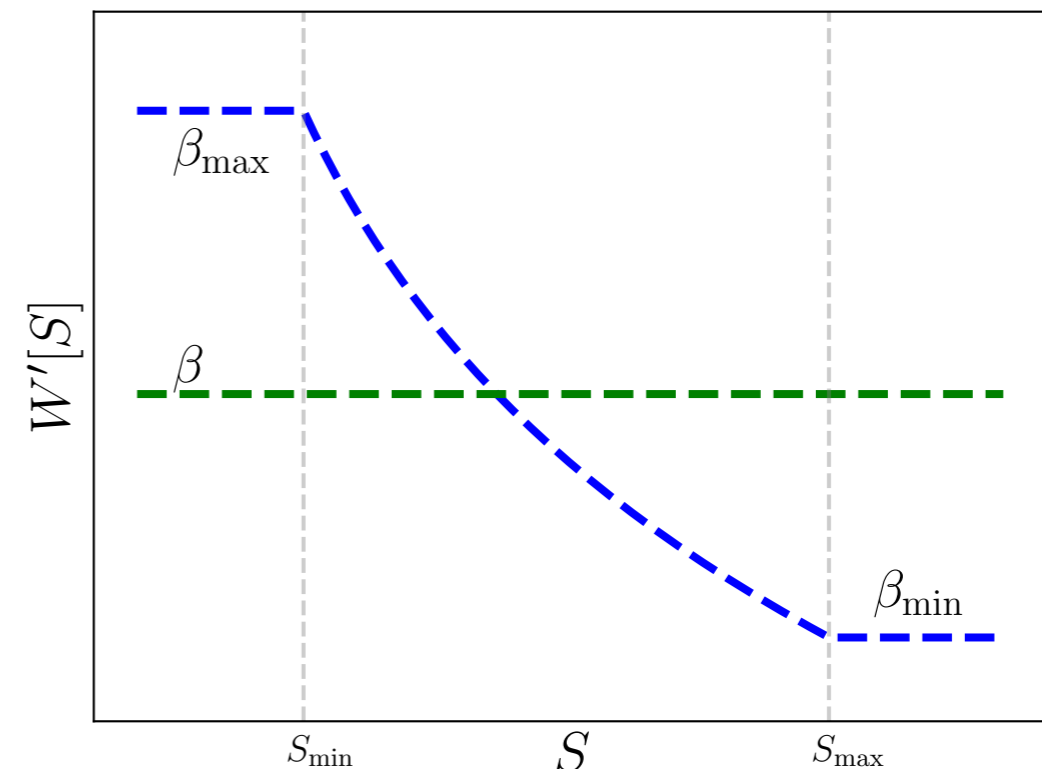
- ▶ Standard single temperature /  $\beta$  MC

- ▶ Computing  $Z(\beta_0)$  from  $\int \mathcal{D}U e^{-W(S)}$  Assuming  $W(S)$  is known

$$Z(\beta_0) = \int \mathcal{D}U e^{-W(S)} e^{W(S) - \beta_0 S}$$

$$\propto \sum_i e^{W(S_i) - \beta_0 S_i}$$

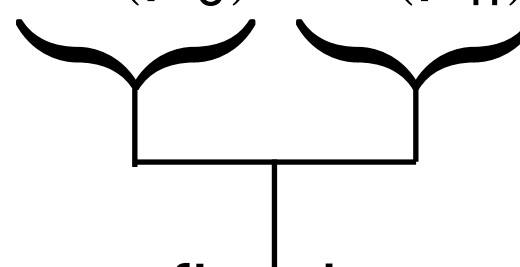
- ▶ Overall normalisation unknown
- ▶ Regularisation issues of  $Z(\beta_0)$



- ▶ Compute QCD  $\chi_{\text{top}}(\beta)$  at high temperatures :

$$\chi_{\text{top}}(\beta) a^4(\beta) = \frac{1}{V_L} \frac{Z_1(\beta)}{Z_0(\beta)}$$

- ▶ Actually compute ratios of  $\chi_{\text{top}}(\beta)$

$$\frac{\chi(\beta_h) a^4(\beta_h)}{\chi(\beta_c) a^4(\beta_c)} = \frac{Z_1(\beta_h)}{Z_1(\beta_c)} \frac{Z_0(\beta_c)}{Z_0(\beta_h)}$$


Allows two fixed topology simulations

- ▶ Compute QCD  $\chi_{\text{top}}(\beta)$  at high temperatures :

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- ▶ Ratio in temperature reweighting :

$$\frac{\chi_{\text{top}}(\beta_h) a(\beta_h)^4}{\chi_{\text{top}}(\beta_c) a(\beta_c)^4} = \frac{Z_1(\beta_h)}{Z_1(\beta_c)} \frac{Z_0(\beta_c)}{Z_0(\beta_h)} = \left( \frac{(\sum_{iQ} e^{W_Q[S_{iQ}]} e^{-\beta_h S_{iQ}}) (\sum_i e^{W[S_i]} e^{-\beta_c S_i})}{(\sum_{iQ} e^{W_Q[S_{iQ}]} e^{-\beta_c S_{iQ}}) (\sum_i e^{W[S_i]} e^{-\beta_h S_i})} \right)$$

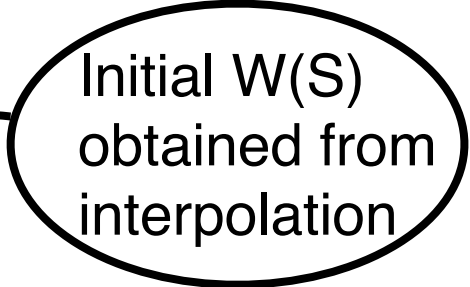


- ▶ We would like to generate  $dP[U] = \frac{\mathcal{D}U e^{-W(S[U])}}{\int \mathcal{D}U e^{-W(S[U])}}$

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▶ Construct the MD Hamiltonian :

$$\mathcal{H}(\pi, U) = \sum_{\mu, x} \frac{1}{2} (\pi_{\mu}(x))^2 + W(S[U])$$



Initial  $W(S)$   
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interpolation

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- ▶ Solve the Hamilton's EOM :

$$\frac{dU}{dt} = -i\pi U \quad \frac{d\pi}{dt} = iU^\dagger \frac{\partial W[S[U]]}{\partial U} = iU^\dagger \frac{\partial W[S[U]]}{\partial S[U]} \frac{\partial S[U]}{\partial U}$$

$\beta$  changes during  
trajectory evolution

- ▶ Metropolis accept/reject step :

$$\Delta \mathcal{H} = \mathcal{H}(\pi_f, U_f) - \mathcal{H}(\pi_i, U_i)$$

- ▶  $W(S)$  needs to be computed in a separate simulation :

- ▶ Initial  $W(S)$  obtained from interpolation :

$$(\beta_{\min}, \beta_{\text{mid}}, \beta_{\max}) \longrightarrow (S_{\max}, S_{\text{mid}}, S_{\min})$$

- ▶ Action  $S$  is split in  $N$  intervals  $(S_1, S_2 \dots S_N)$

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$$W(S) = \begin{cases} W_i + s_r x (W_{i+1} - W_i), & S_{\max} < S_i < S_{\min} \\ \beta_{\max} S, & S < S_{\min} \\ \beta_{\min} S, & S > S_{\max} \end{cases}$$

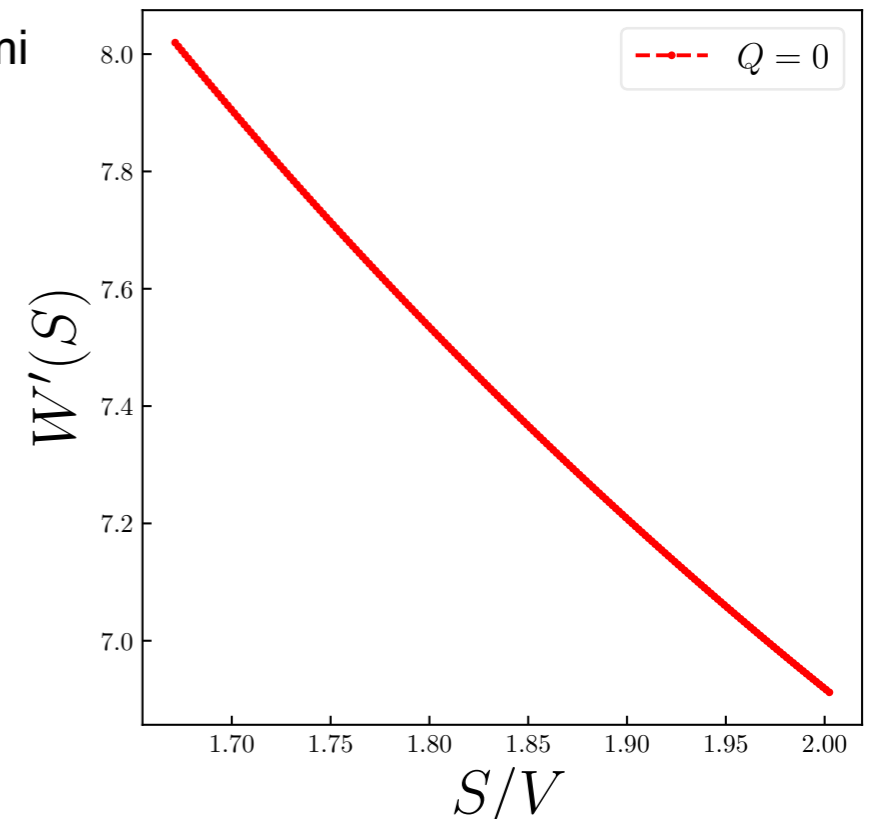
- ▶ After  $S$  traverses all intervals and back,  $s_r$  is reduced
- ▶ Separate simulation needed for  $Q=0,1$

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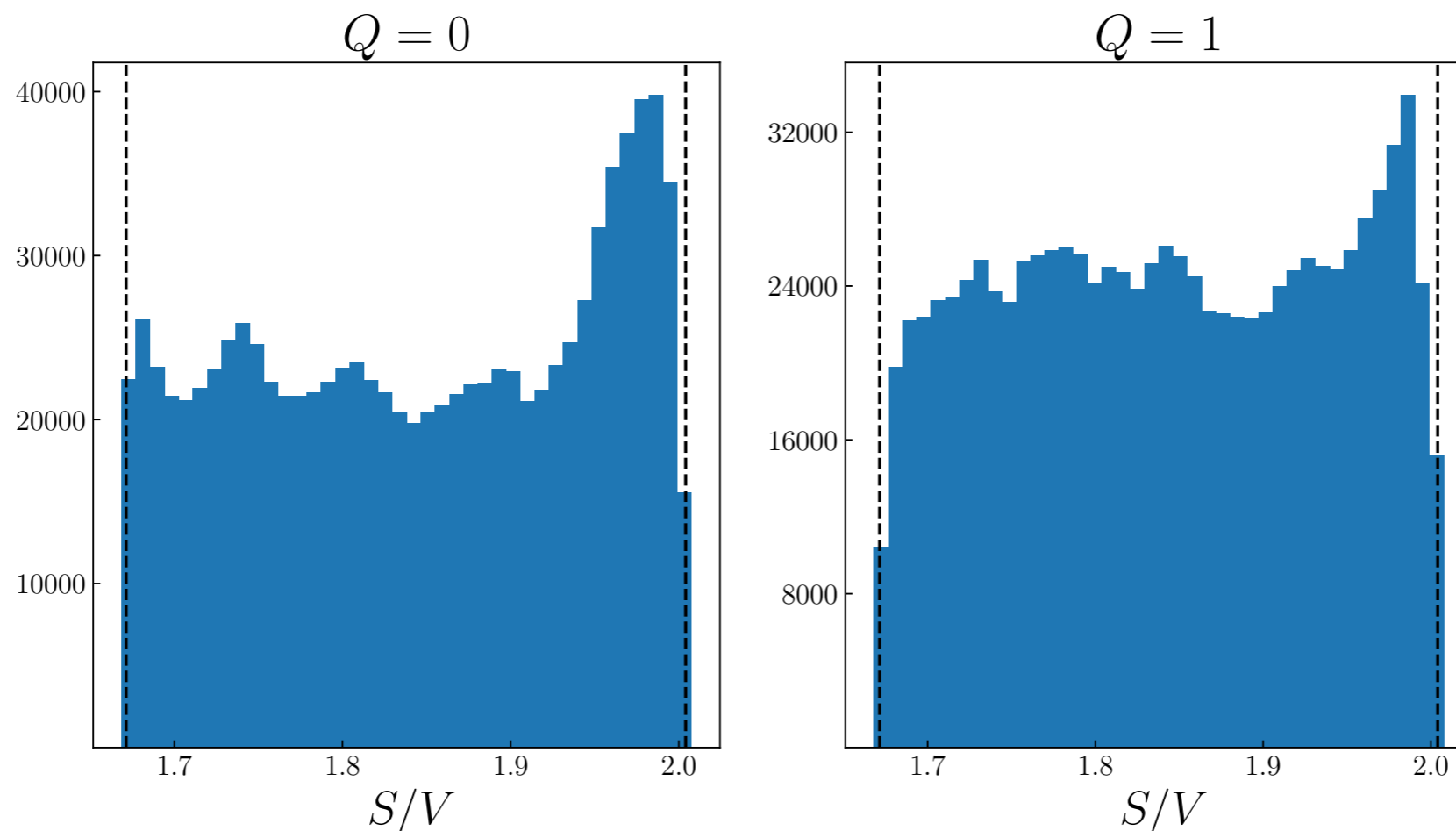
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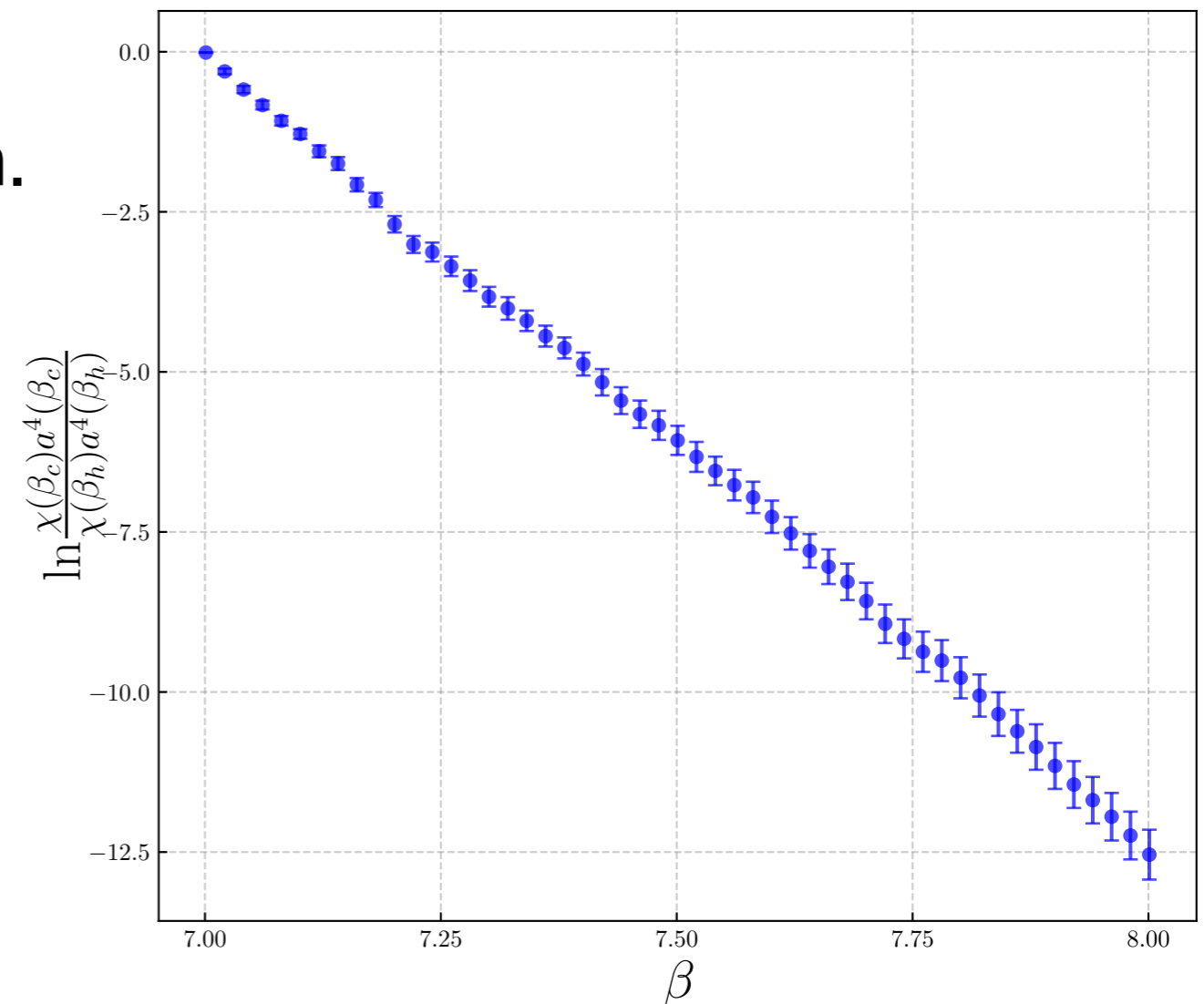


- ▶ Separate simulation needed for  $Q=0, 1$

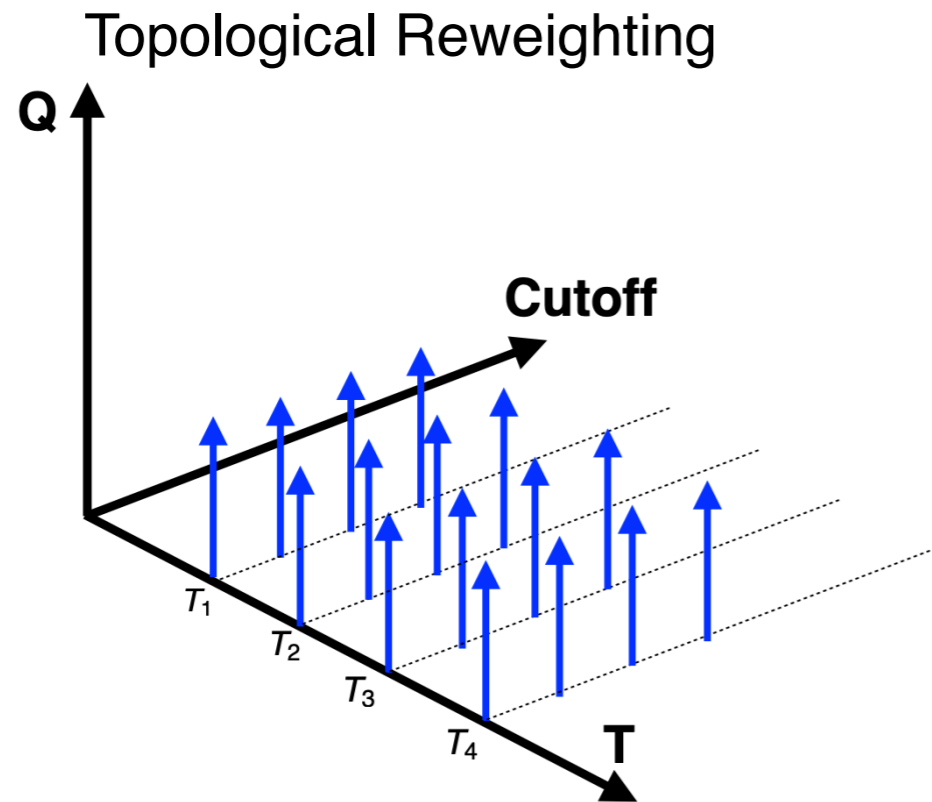
- ▶ Simulation with Wilson gauge action
- ▶ Temperature range  $2.5T_c - 9.4T_c$
- ▶ Lattice geometry  $N_t \times N_{x,y} \times N_z = 10 \times 32^2 \times 36$



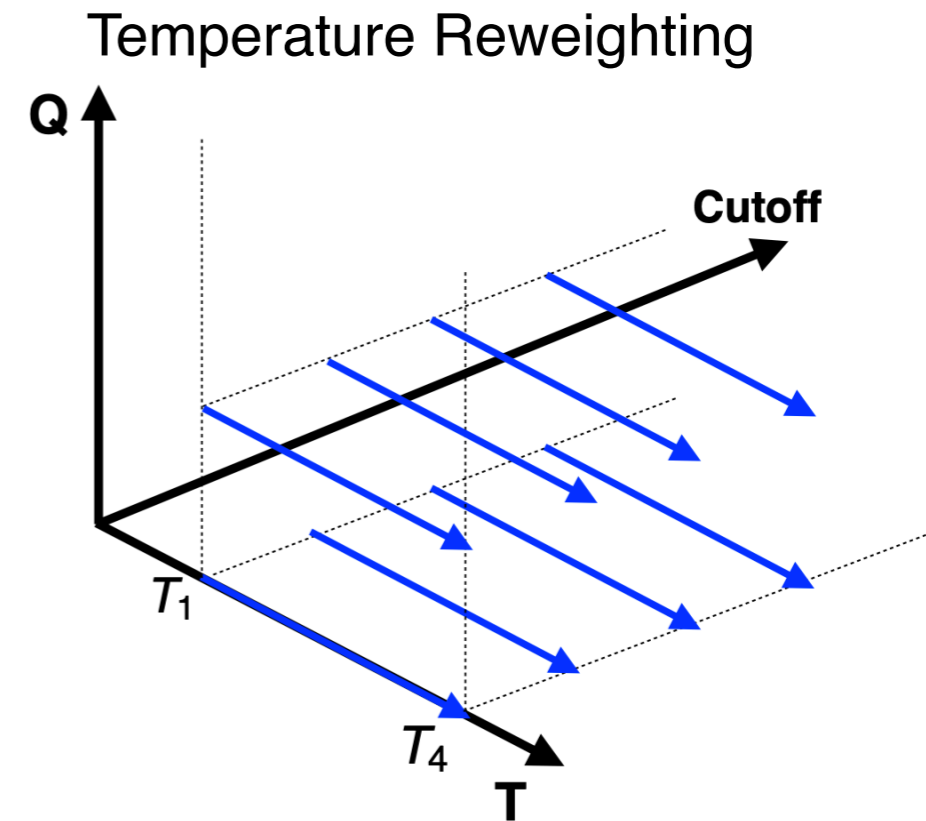
- ▶ Method accomplishes continuous sampling of temperatures.
- ▶ Results in agreement with topological reweighting.
- ▶ Scale agnostic determination.
- ▶ Method comparable to current methods for high statistics determination.
- ▶ Inefficient for less precise determination







- ▶ Needs to be simulated at every cutoff and temperature.
- ▶ No of simulations =  $x$
- ▶ Simulation for reweighting function cheap
- ▶ Cheaper for lower statistics



- ▶ Needs to be simulated at two topologies and every cutoff.
- ▶ No of simulations =  $x / 2$
- ▶ Simulation for reweighting function very expensive
- ▶ Expensive for lower statistics

Thank you for your attention !